

# Quantum Geometry and Hawking Radiation

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[PMH-Kawai-Yokokura 2111.01967][PMH-Kawai 2207.07122]

[Chau-PMH-Kawai-Shao-Wang 2309.01638]

[PMH-Imamura-Kawai-Shao 2309.12936]

[PMH-Kawai-Shao 2411.01105]

[Chang-PMH-Lee-Shao 2412.02577]

[Ho-Shao-Yoda ...][Cheng-Ho-Kawai-Shao ...]

# spacetime uncertainty and UV-IR connection

Noncommutative Geometry:

$$[x, t] = i\ell_p^2 \quad \Rightarrow \quad \Delta x \Delta t \gtrsim \ell_p^2$$

$$\Rightarrow \quad \Delta x > \ell_p^2 \Delta E > \ell_p^2 \Delta p$$

$$\Delta x \gtrsim \ell_p^2 \Delta p \quad (\text{UV-IR connection})$$

Generalized Uncertainty Relation: [Amati-Ciafaloni-Veneziano 87, 89]  
[Konishi-Paffuti-Provero 90]

$$\Delta x \Delta p \gtrsim 1 + \ell_p^2 \Delta p^2$$

$$\text{Large } \Delta p \quad \rightarrow \quad \Delta x \gtrsim \ell_p^2 \Delta p$$

# spacetime uncertainty and UV-IR connection

$$\Delta x \gtrsim \ell_p^2 \Delta p \quad (\text{UV-IR connection})$$

- Large spatial uncertainty at trans-Planckian energies.  
Dramatic UV corrections with IR effects.  
Hawking radiation is insensitive to generic UV physics.  
But it is sensitive to UV-IR connection:
  - Hawking radiation stops around scrambling time.

# black hole formation and evaporation: conventional model

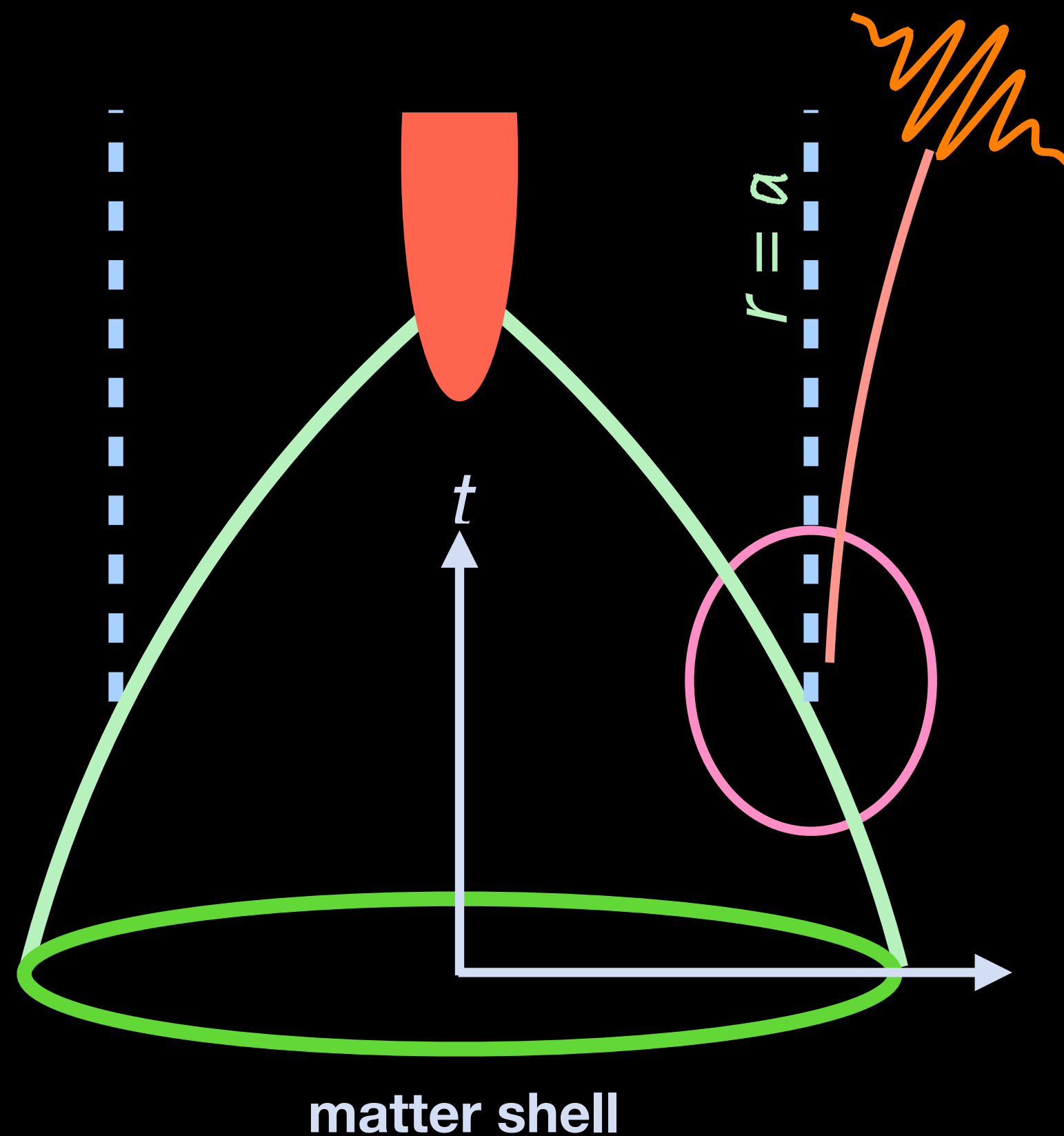
[Hawking 76]

$$a = 2\ell_p^2 M$$

Hawking radiation: mostly geometric information

info. in matter  $\rightarrow$  radiation?  
same as burning a piece of coal?

$$G_N = \ell_p^2, \quad \hbar = 1, \quad c = 1$$



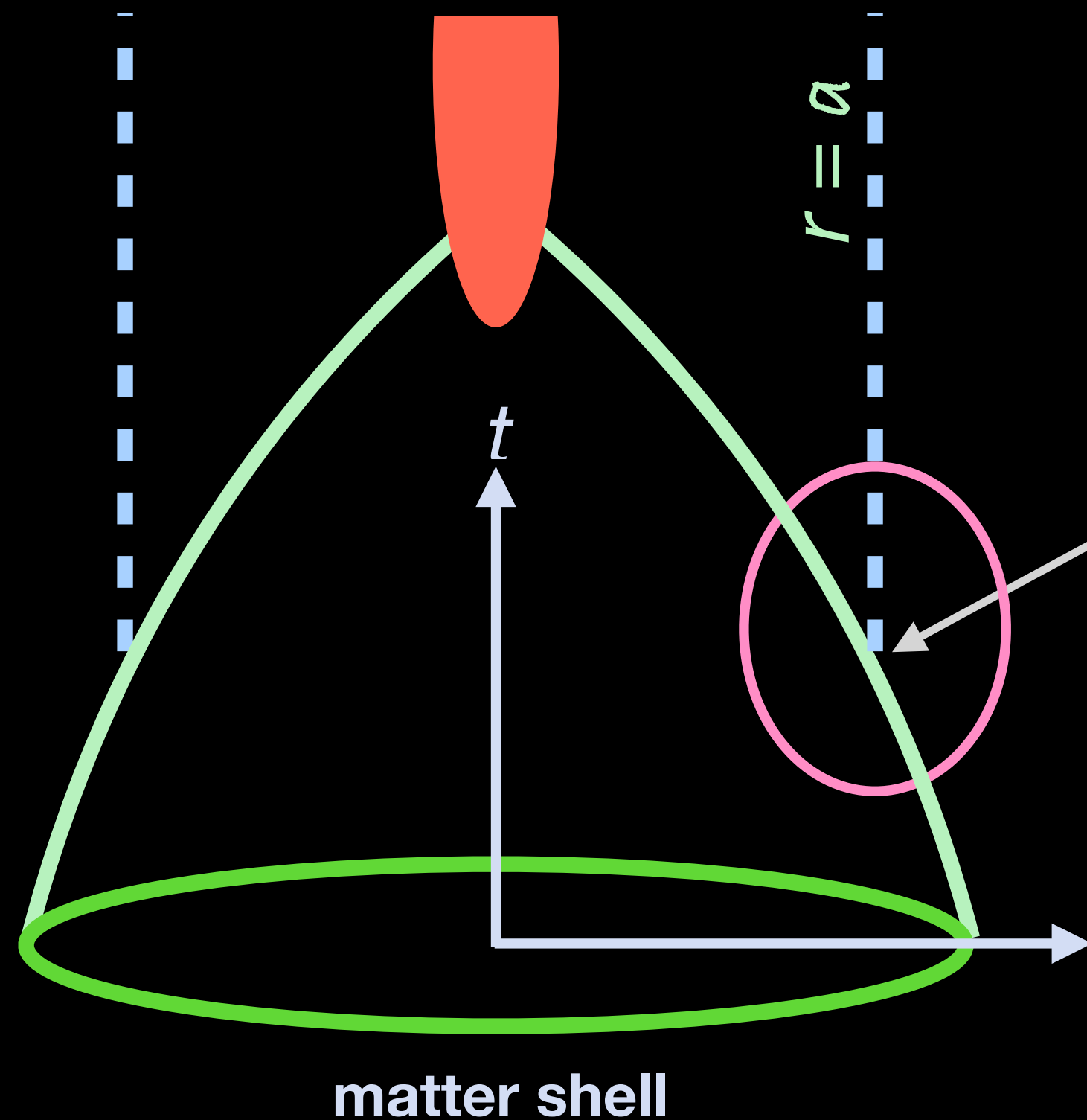
# black hole formation and evaporation: conventional model

[Hawking 76]

$$a = 2\ell_p^2 M$$

Equivalence principle:

We may have just crossed the horizon.



13directioner.deviantart.com

General Relativity  $\Rightarrow$  “uneventful” horizon

# black-hole information paradox

Hawking's calculation in low-energy effective theory (LEET)

- Hawking radiation carries (almost) no information and persists until (almost) complete evaporation
- information paradox

Q: Is Hawking's LEET calculation reliable?



# black-hole information paradox

[Quanta magazine "*The Fuzzball Fix for a Black Hole Paradox*" 2015]

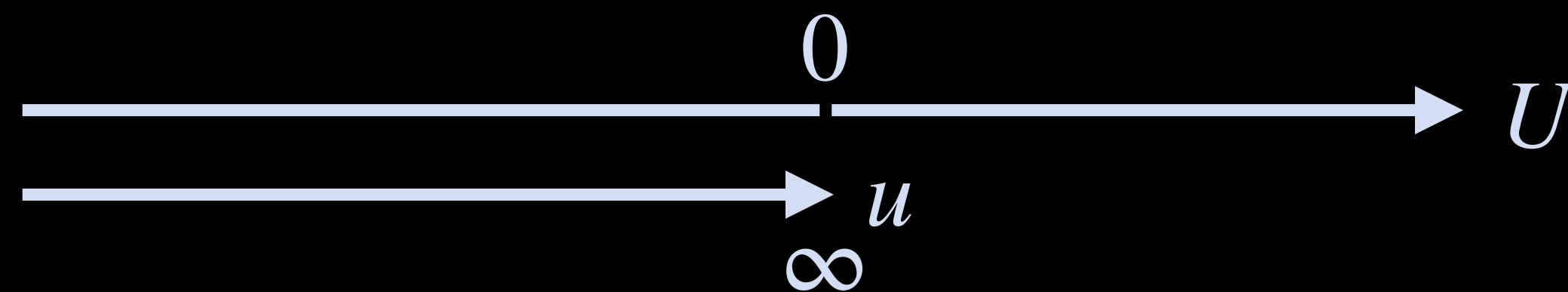
Polchinski said of Hawking's original premise:

*"... a situation that seems to violate the laws of quantum mechanics.  
... 'Quantum mechanics is modified. Find my mistake.' And nobody  
found his mistake."*

There have been countless papers on how Hawking radiation is robust,  
insensitive to UV modifications of LEET.

# Hawking radiation

[Hawking 74]



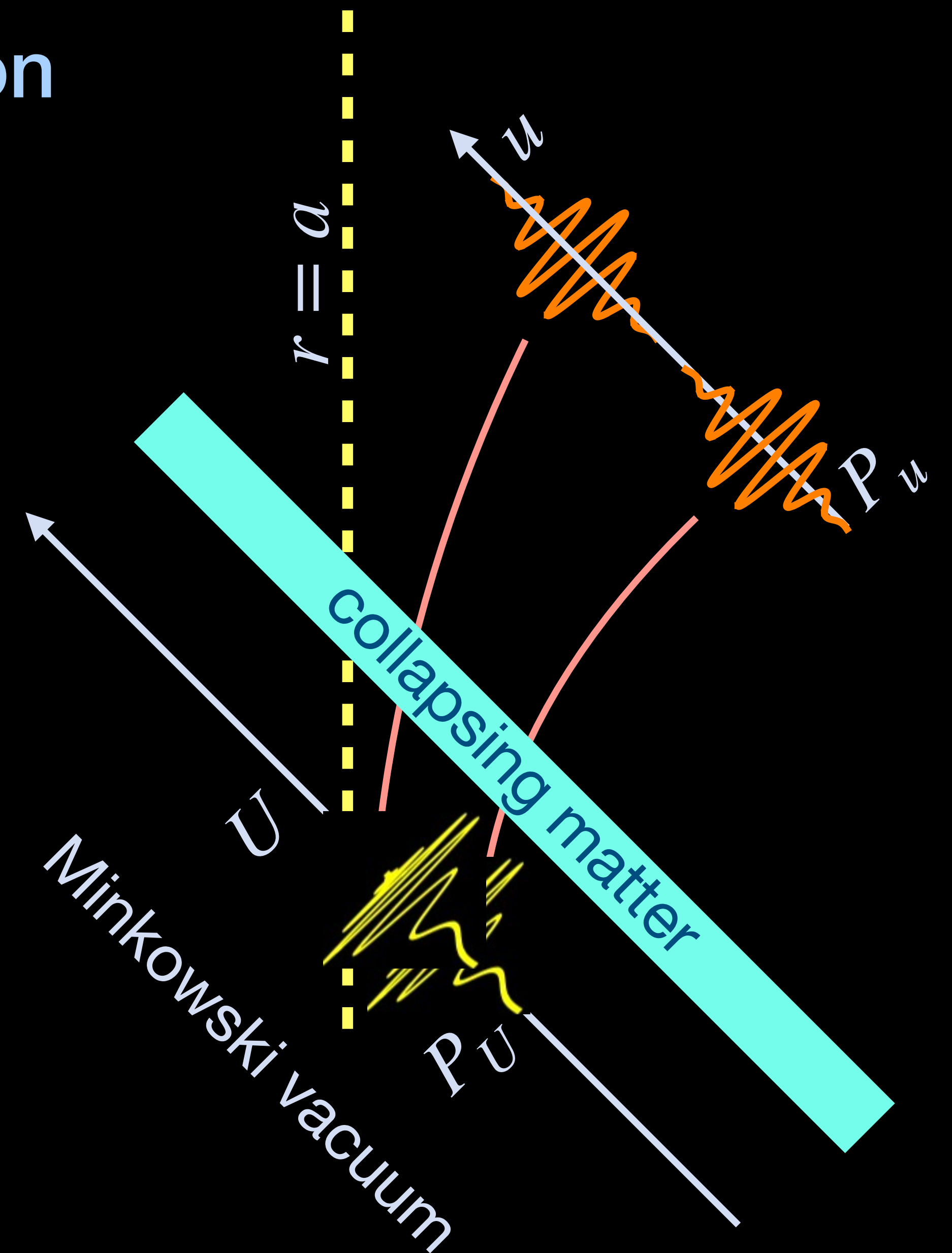
$$U = -2ae^{-u/2a}$$

$$\langle P_U \rangle \sim \langle P_u \rangle e^{u/2a}$$

Minkowski vacuum  $|0\rangle$  of the infinite past

$\simeq$  Hawking radiation at large distance

*Q: Is there a trans-Planckian Lorentz-invariant?*





# Hawking radiation

[Hawking 74]

outgoing massless scalar

$$\phi = \int_0^\infty \frac{d\Omega}{2\pi\sqrt{\Omega}r} \left( a_\Omega e^{-i\Omega U} + a_\Omega^\dagger e^{i\Omega U} \right) = \int_0^\infty \frac{d\omega}{2\pi\sqrt{\omega}r} \left( b_\omega e^{-i\omega u} + b_\omega^\dagger e^{i\omega u} \right)$$

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$$\Rightarrow \quad b_\omega = \int_0^\infty d\Omega \left( \alpha_{\omega\Omega} a_\Omega + \beta_{\omega\Omega} a_\Omega^\dagger \right), \quad b_\omega^\dagger = \int_0^\infty d\Omega \left( \alpha_{\omega\Omega}^* a_\Omega^\dagger + \beta_{\omega\Omega}^* a_\Omega \right).$$

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$$\Rightarrow b_\omega = \int_0^\infty d\Omega \left( \alpha_{\omega\Omega} a_\Omega + \beta_{\omega\Omega} a_\Omega^\dagger \right), \quad b_\omega^\dagger = \int_0^\infty d\Omega \left( \alpha_{\omega\Omega}^* a_\Omega^\dagger + \beta_{\omega\Omega}^* a_\Omega \right).$$

Unruh vacuum (for freely falling observers):  $a_\Omega |0\rangle = 0 \quad (\forall \Omega > 0)$

→ No Hawking radiation for freely falling observers.

Hawking radiation for distant observers:  $\langle 0 | b_\omega^\dagger b_{\omega'} | 0 \rangle > 0$ .

# higher-derivative interactions

[PMH-Yokokura 20][PMH-Kawai-Yokokura 21]

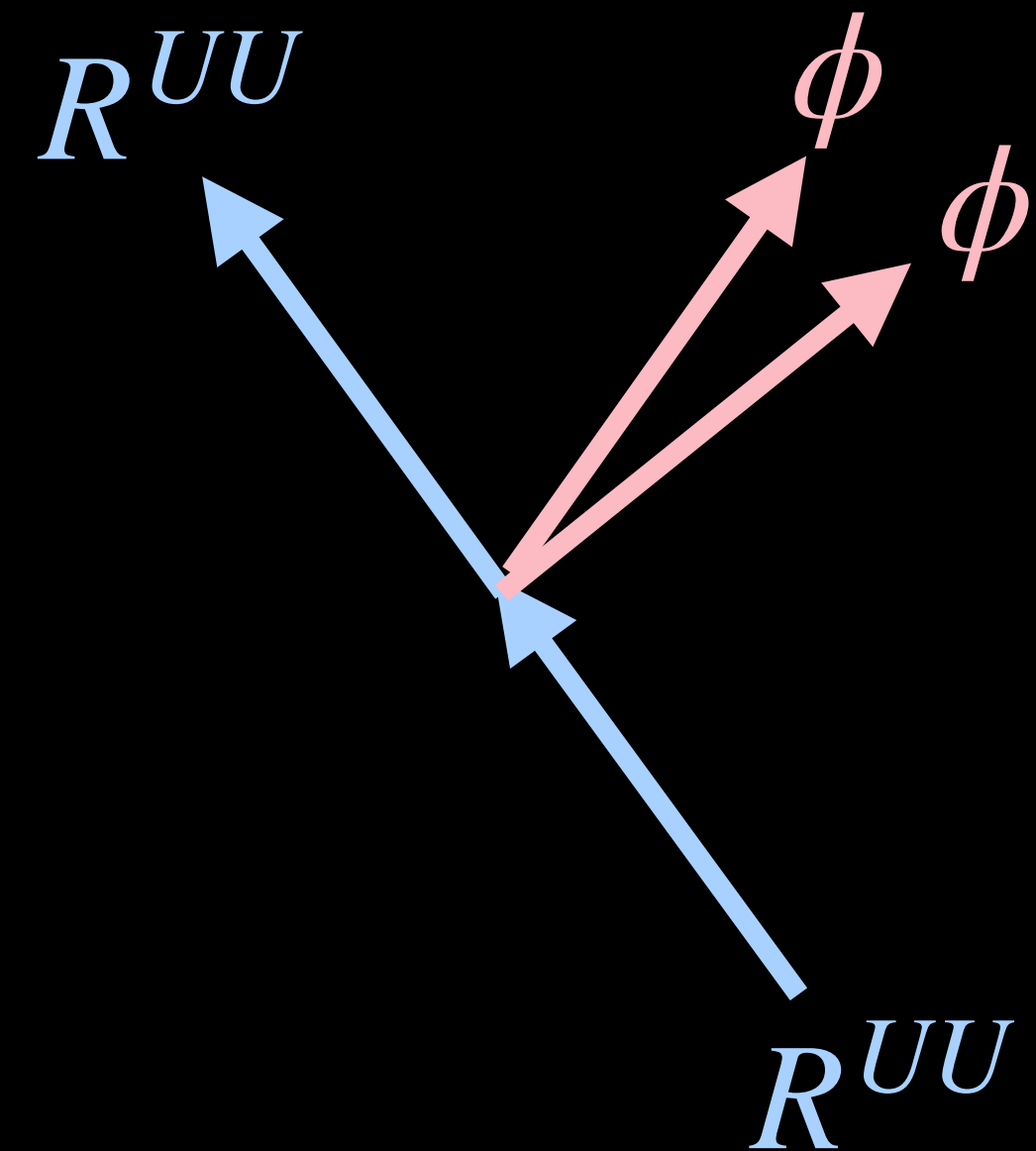
Ex:

$$S_n \equiv \frac{\lambda_n}{M_p^{4n-2}} \int d^4x \sqrt{-g} R^{\mu_1 \nu_1} \dots R^{\mu_n \nu_n} \left( \nabla_{\mu_1} \dots \nabla_{\mu_n} \phi \right) \left( \nabla_{\nu_1} \dots \nabla_{\nu_n} \phi \right)$$

$$\sim \lambda_n \left( \frac{s}{M_p^2} \right)^{2n-1}$$

$$s \equiv P_U P_V \sim \frac{1}{a^2} e^{u/2a}$$

$$u \sim a \log(a/\ell_p) \Rightarrow \text{Ampl.} \sim \mathcal{O}(1)$$

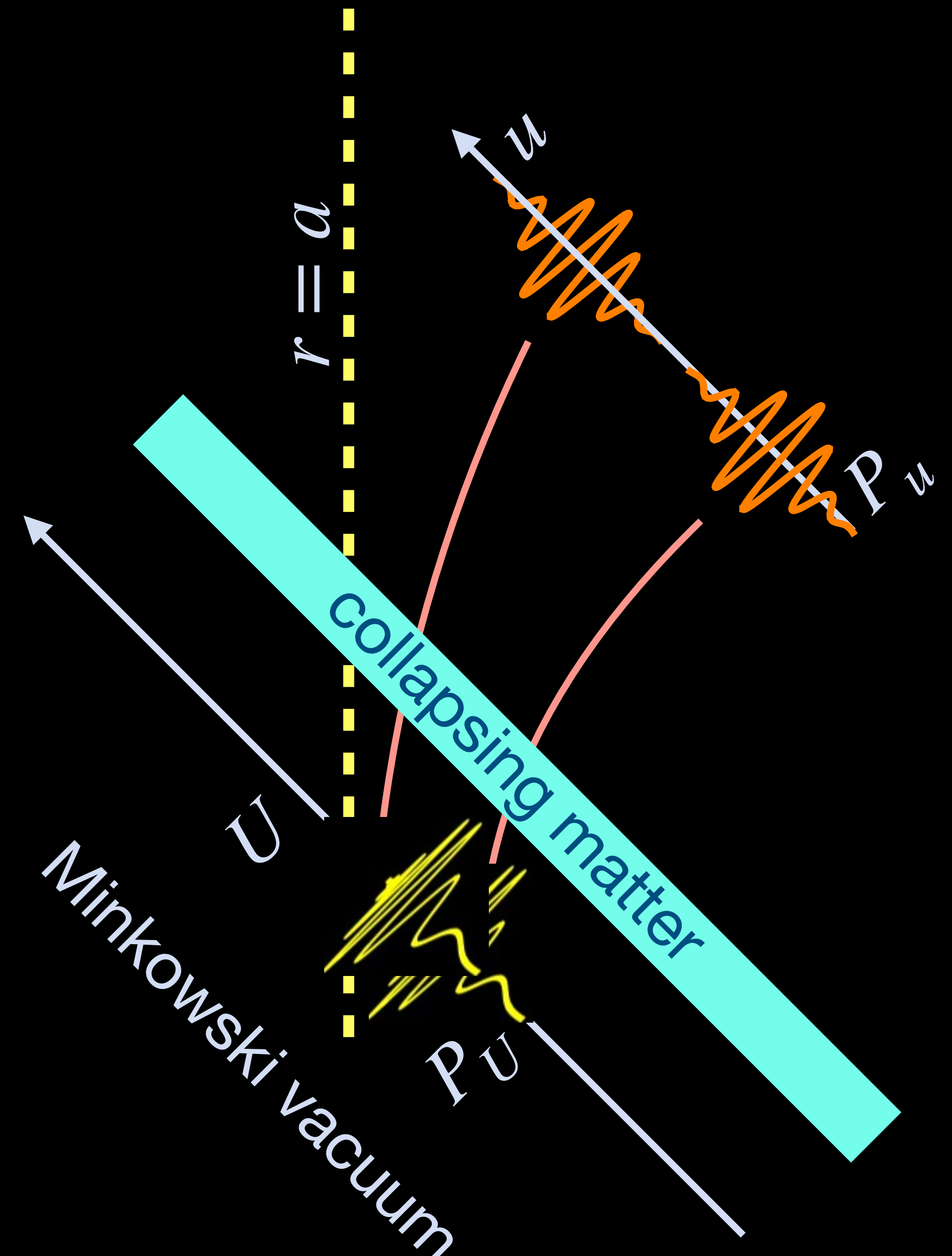




# black hole as a microscope

The detection of Hawking quanta explores trans-Planckian regions of the spacetime in Minkowski vacuum.

It needs UV theories for reliable predictions.



# black-hole information paradox

Q: Is Hawking's calculation reliable?

No. ← Trans-Planckian local Lorentz inv. after scrambling time.

# black-hole information paradox

Q: Is Hawking's calculation reliable?

No. ← Trans-Planckian local Lorentz inv. after scrambling time.

→ UV physics is relevant.

How does a given UV theory change HR?

1. HR as in LEET  $\Rightarrow$  still a paradox. ← literature
2. HR enhanced by informative radiation.

Q: Other logical possibilities?

## alternative resolution

[Chau-PMH-Kawai-Shao-Wang 23][PMH-Imamura-Kawai-Shao 23][PMH-Kawai-Shao 24]

3. Hawking radiation stops well before the Page time,  
say, around scrambling time:

$$t_{scr} \sim a \log(a/\ell_p) \ll t_{Page} \sim a^3/\ell_p^2$$

→ HR is negligible.



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$$t_{scr} \sim a \log(a/\ell_p) \ll t_{Page} \sim a^3/\ell_p^2$$

→ HR is negligible.

However, in the literature, “*HR is insensitive to UV.*”

Are there loopholes in the literature?

# Departure from the literature

In the literature, typically,

- the collapsing matter is ignored.

Time translation symmetry  $\Rightarrow$  Hawking radiation persists.

Hartle-Hawking or Boulware vacuum.

- time-dependent amplitude of Hawking radiation is not derived.

Hawking temperature is indeed robust.

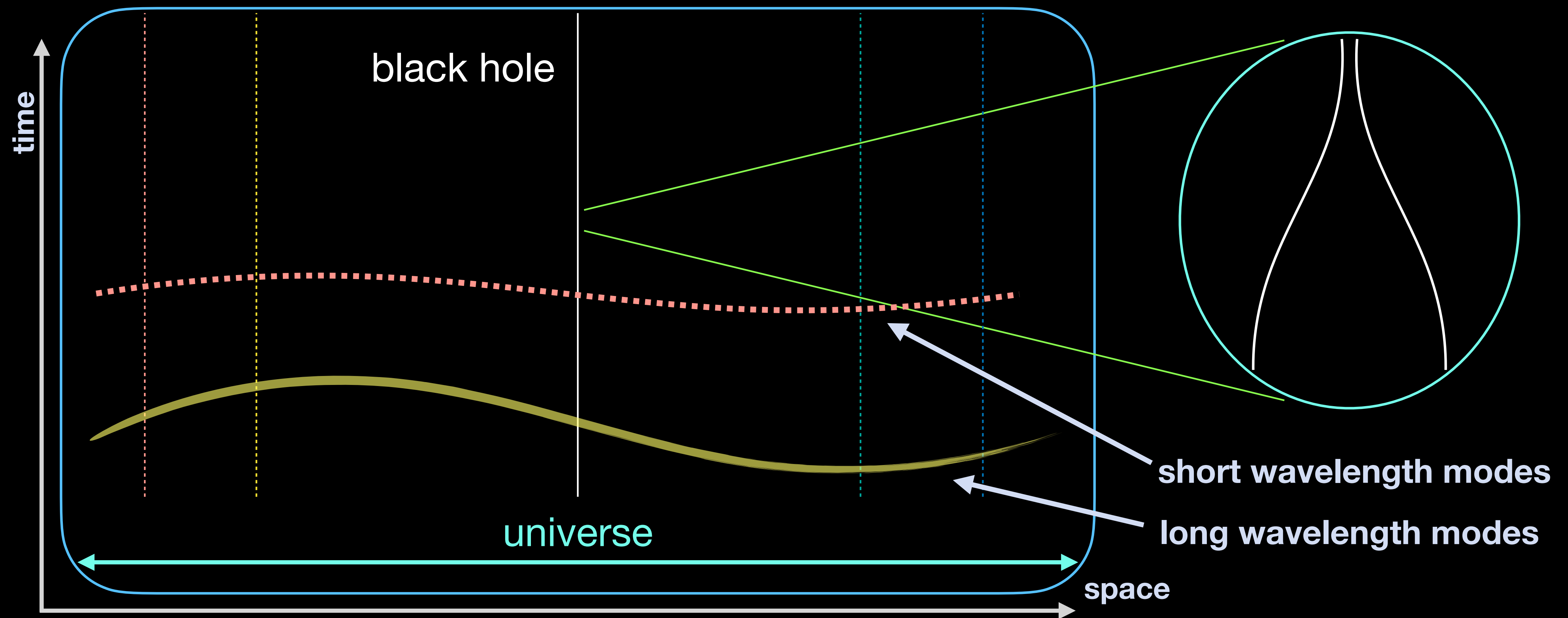
- there is no UV-IR relation in UV physics.

e.g. UV dispersion relations do not turn off Hawking radiation.

e.g. 2D or 3D theories.

Small-scale structures are irrelevant to low-energy modes.

UV-IR relation  $\rightarrow$  also to trans-Planckian modes.



# Generalized Uncertainty Principle (GUP)

[Amati-Ciafaloni-Veneziano 87, 89] [Konishi-Paffuti-Provero 90]

$$\Delta x \Delta p \gtrsim 1 + \ell_p^2 \Delta p^2$$

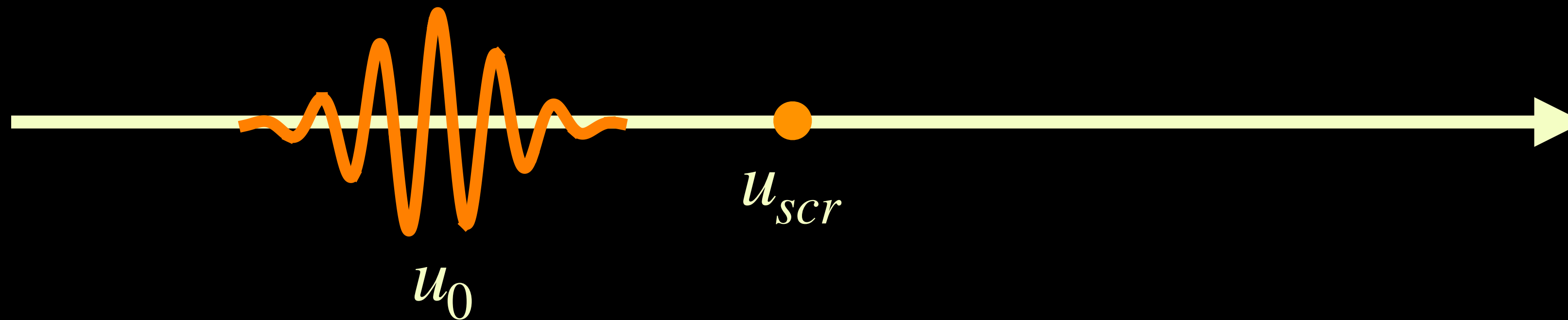
$$\Rightarrow \text{minimal length} \quad \Delta x \gtrsim 2\ell_p$$

$$\Delta p \gg \ell_p^{-1} \quad \Rightarrow \quad \Delta x \gtrsim \ell_p^2 \Delta p$$

Robustness of HR with GUP noted repeatedly in many works.

[0506110, 1212.6591, 1410.4115, 1410.5065, 1501.03256, 1501.06025, 1602.04304, 1704.03536, 1709.00637, 1903.01382, 2112.13573, 2306.03077].

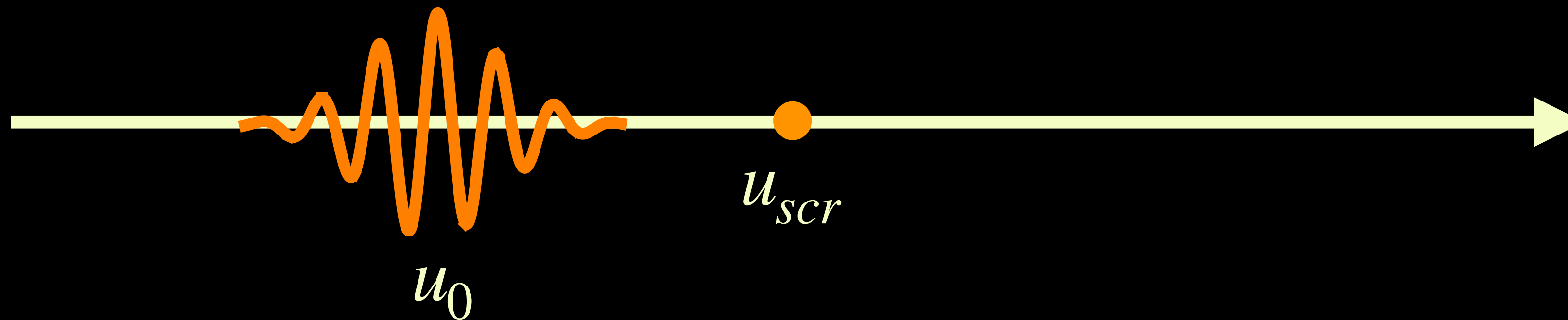




Hawking radiation = the VEV of the number operator for this wave packet:

$$\langle 0 | b_{\Psi}^{\dagger} b_{\Psi} | 0 \rangle \simeq \frac{1}{2} \frac{\omega_0}{e^{4\pi a \omega_0} - 1} \int_{-\infty}^{u_{scr}} du \left| \Psi_{(\omega_0, u_0)}(u) \right|^2$$

$$u_{scr} = 4a \log(a/\ell_p)$$

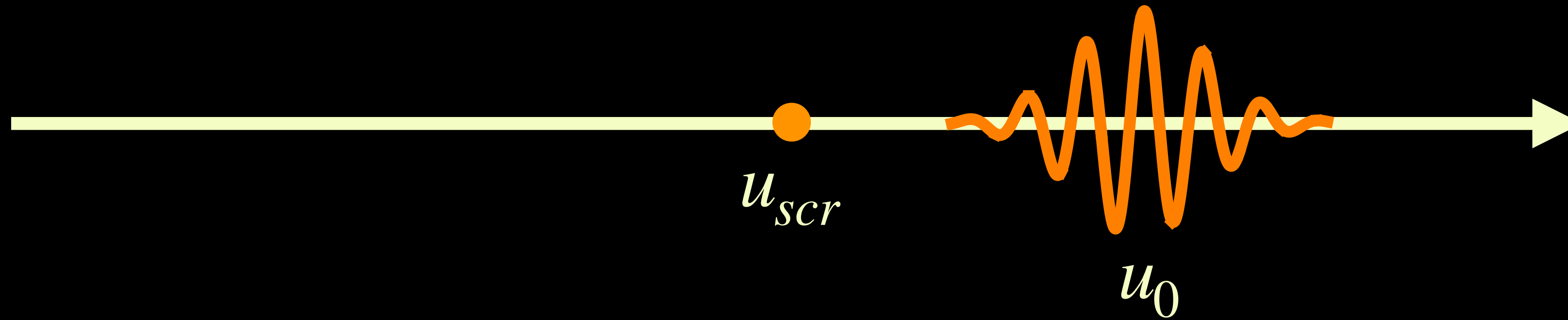


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$\uparrow$   $u_{scr} = 4a \log(a/\ell_p)$

Hawking temperature is the same.



Hawking radiation = the VEV of the number operator for this wave packet:

$$\langle 0 | b_{\Psi}^{\dagger} b_{\Psi} | 0 \rangle \simeq \frac{1}{2} \frac{\omega_0}{e^{4\pi a \omega_0} - 1} \int_{-\infty}^{u_{scr}} du \left| \Psi_{(\omega_0, u_0)}(u) \right|^2$$

→ 0 when  $u_0 \gg u_{scr} \equiv 2a \log(a/\ell_p)$  (scrambling time).

At late times, the wave packet has a large  $\Delta p$ , GUP implies a large  $\Delta x$ .

When  $\Delta x$  is much larger than  $a$ , it does not contribute to Hawking radiation.

[Chau-PMH-Kawai-Shao-Wang 23]

# comments

- Entanglement Island proposal [Almheiri-Engelhardt-Marolf-Maxfield 19, Penington 19]  
nonlocal but low-energy effective description (lower dimensions)  
UV description for  $D \geq 4$ ?
- AdS/CFT duality compatibility

## Applications:

- Primordial black holes as dark matter candidates. [Che-Yu's talk]
- Trans-Planckian Censorship Conjecture  
[Blamart-Laliberte-Brandenberger 23][Brandenberger-PMH-Kawai-Shao 24]



# conclusion

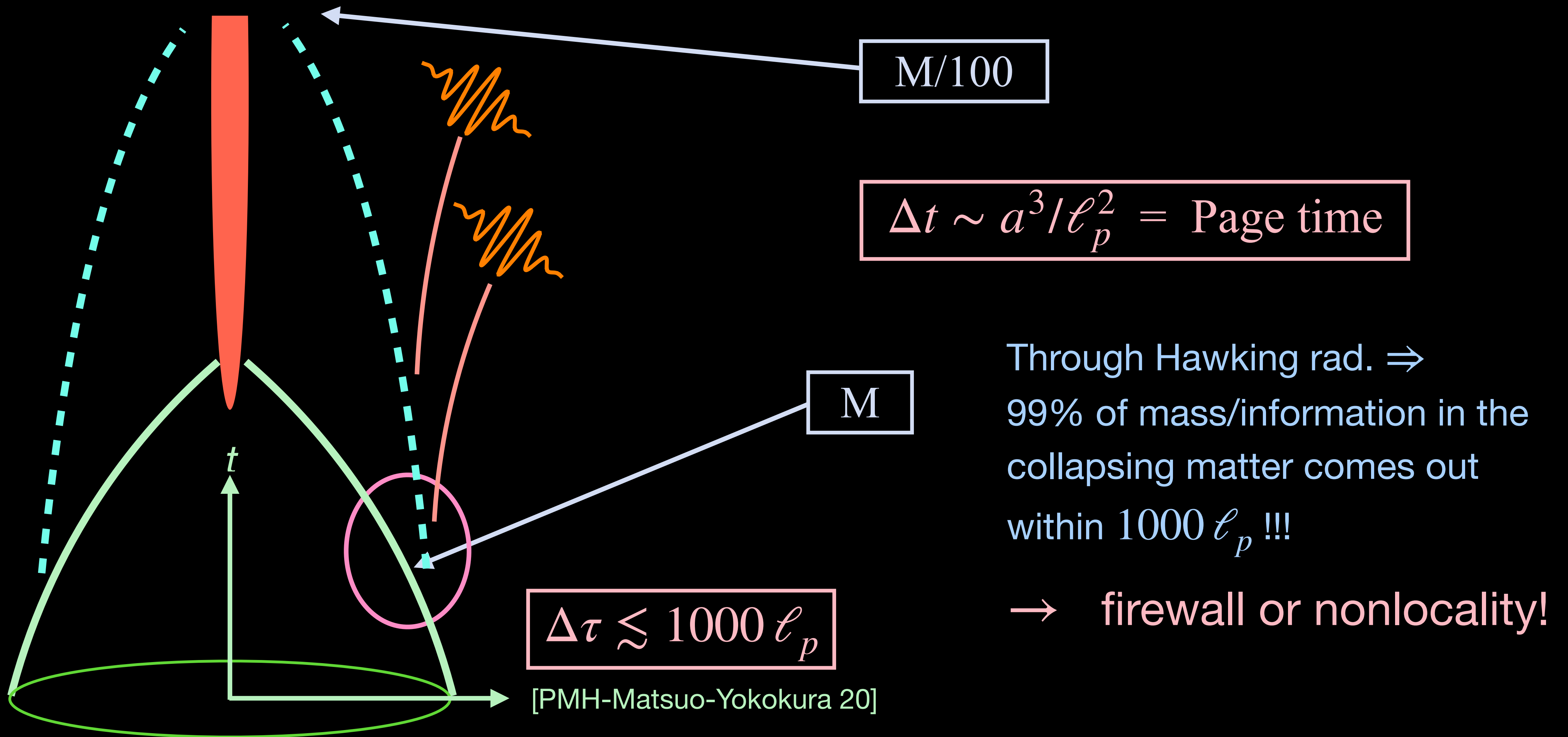
LEET breaks down at scrambling time for Hawking radiation.

Alternative proposal

- Spacetime uncertainty turns off Hawking radiation at scrambling time.
- Black holes are essentially classical.
- Hawking temperature not modified  $\rightarrow$  same BH entropy, etc.
- Spacetime uncertainty blurs the horizon.
  - $\rightarrow$  Other channels of informative radiation (cf. fuzzball)

*Thank you!*

# Information transfer in the conventional model





# Is UV theory relevant?

## Nice-slice argument\*

Near-horizon region remains approximately vacuum  
in low-energy effective theory.



# nice-slice argument

[Lowe-Polchinski-Sussking-Thorlacius-Uglum 95]

The curvature  $\sim \mathcal{O}(1/a)$ .

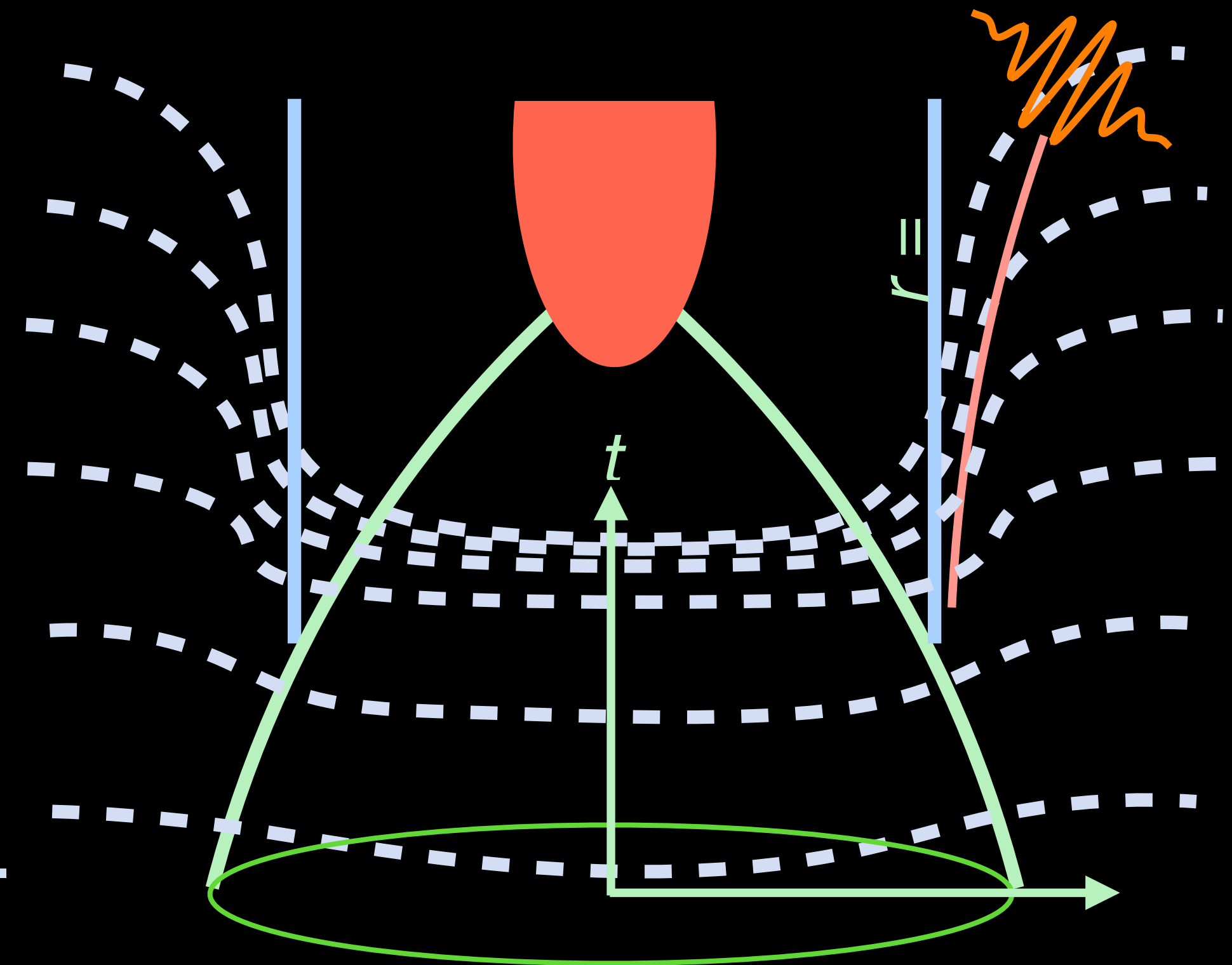
**Adiabatic theorem**  $\Rightarrow$

If the initial state is vacuum,  
excitations of energies  $\lesssim \mathcal{O}(1/a)$   
from time evolution.

$\Rightarrow$  Effective theory remains valid.

**Decoupling theorem**

Effective theory breaks down for trans-Planckian observations!





# Is UV theory relevant?

## Nice-slice argument

Near-horizon region remains approximately vacuum  
in low-energy effective theory. (OK)

→ UV physics is irrelevant?

LEET is not good for predicting UV events.

Q: Is Hawking radiation UV physics?

Q: Trans-Planckian local Lorentz invariant?



# another example: SFT

[PMH-Imamura-Kawai-Shao 23] [Chang-PMH-Lee-Shao 24]

$$S_{SFT} = \int d^D x \left[ \frac{1}{2} \phi_\alpha (\partial^2 - m_\alpha^2) \phi_\alpha + g_{\alpha\beta\gamma} \tilde{\phi}_\alpha \tilde{\phi}_\beta \tilde{\phi}_\gamma + \cdots \right],$$

$$\tilde{\phi}_\alpha \equiv e^{\frac{1}{2} \ell^2 \partial_\mu \partial^\mu} \phi_\alpha \quad \eta = ( - + + \cdots + )$$

→ exponential suppression of UV interactions in string theory.

$$S_{SFT} = \int d^D x \left[ \tilde{\phi}_\alpha (\partial^2 - m_\alpha^2) e^{-\ell^2 \partial^2} \tilde{\phi}_\alpha + g_{\alpha\beta\gamma} \tilde{\phi}_\alpha \tilde{\phi}_\beta \tilde{\phi}_\gamma + \cdots \right]$$

→ coupling to background suppressed for large- $k$  modes.

→  $\Delta U \Delta V \gtrsim \ell_E^2$  (spacetime uncertainty relation, cf. [Yoneya])

Hawking radiation turned off at scrambling time.

# compatibility with AdS/CFT duality

- Small AdS black hole

Before scrambling time: same classical black-hole geometry in all models.

↔ a thermal state with  $T_H$  at scrambling time [Danielsson et al 99])

After scrambling time: (before Hawking-Page transition)

- Conventional model: complete evaporation via HR in conventional model
- Our scenario: HR stops. If there are other decay channels ...

↔ Hawking temp increases towards Planck scale but details are unknown.

Entanglement entropy formula cannot predict Page time.

It is possible to have the “same” Page curve, but different Page time.

- Large AdS black hole [Maldacena 01]

No obvious paradox if the causal barrier (horizon) is removed.

# compatibility with AdS/CFT duality

