

# Asymptotically Safe Standard Model

Francesco Sannino



# Plan

- Meaning of fundamental
- Asymp. safe theories in 4D
- Large  $N_f$  asymptotic safety
- Asymptotically safe standard model
- Exact nonperturbative results for  $N=1$  supersymmetric safety
- Super GUTs with R-parity are trivial = unphysical
- SUSY-like radiative symmetry breaking in UV safe theories\*
- New paths in (astro)particle physics

# Standard Model

# Standard Model

**Fields:**

Gauge fields + fermions + scalars

# Standard Model

**Fields:**

Gauge fields + fermions + scalars

**Interactions:**

Gauge:  $SU(3) \times SU(2) \times U(1)$  at EW scale

# Standard Model

**Fields:**

Gauge fields + fermions + scalars

**Interactions:**

Gauge:  $SU(3) \times SU(2) \times U(1)$  at EW scale

Yukawa: Fermion masses/Flavour

Scalar self-interaction

Culprit: Higgs

# Gauge - Yukawa theories

$$L = -\frac{1}{2}F^2 + i\bar{Q}\gamma_\mu D^\mu Q + y(\bar{Q}_L H Q_R + \text{h.c.})$$

$$\text{Tr} [D H^\dagger D H] - \lambda_u \text{Tr} [(H^\dagger H)^2] - \lambda_v \text{Tr} [(H^\dagger H)]^2$$

# Gauge - Yukawa theories

$$L = \boxed{-\frac{1}{2}F^2 + i\bar{Q}\gamma_\mu D^\mu Q} + y(\bar{Q}_L H Q_R + \text{h.c.})$$

Gauge  $\boxed{\text{Tr} [D H^\dagger D H]} - \lambda_u \text{Tr} [(H^\dagger H)^2] - \lambda_v \text{Tr} [(H^\dagger H)]^2$

# Gauge - Yukawa theories

$$L = \boxed{-\frac{1}{2}F^2 + i\bar{Q}\gamma_\mu D^\mu Q} + \boxed{y(\bar{Q}_L H Q_R + \text{h.c.})} \quad \text{Yukawa}$$

Gauge  $\boxed{\text{Tr} [D H^\dagger D H]} - \lambda_u \text{Tr} [(H^\dagger H)^2] - \lambda_v \text{Tr} [(H^\dagger H)]^2$

# Gauge - Yukawa theories

$$L = \boxed{-\frac{1}{2}F^2 + i\bar{Q}\gamma_\mu D^\mu Q} + \boxed{y(\bar{Q}_L H Q_R + \text{h.c.})} \quad \text{Yukawa}$$

Gauge  $\text{Tr} [D H^\dagger D H]$  -  $\lambda_u \text{Tr} [(H^\dagger H)^2] - \lambda_v \text{Tr} [(H^\dagger H)]^2$

Scalar selfinteractions

# Gauge - Yukawa theories

$$L = \boxed{-\frac{1}{2}F^2 + i\bar{Q}\gamma_\mu D^\mu Q} + \boxed{y(\bar{Q}_L H Q_R + \text{h.c.})} \quad \text{Yukawa}$$

Gauge  $\text{Tr} [D H^\dagger D H]$  -  $\lambda_u \text{Tr} [(H^\dagger H)^2] - \lambda_v \text{Tr} [(H^\dagger H)]^2$

Scalar selfinteractions

4D: standard model, dark matter, ...

# Gauge - Yukawa theories

$$L = \boxed{-\frac{1}{2}F^2 + i\bar{Q}\gamma_\mu D^\mu Q} + \boxed{y(\bar{Q}_L H Q_R + \text{h.c.})} \quad \text{Yukawa}$$

Gauge  $\text{Tr} [D H^\dagger D H]$  -  $\lambda_u \text{Tr} [(H^\dagger H)^2] - \lambda_v \text{Tr} [(H^\dagger H)]^2$

Scalar selfinteractions

4D: standard model, dark matter, ...

Lower D: condensed matter, phase transitions, graphene

# Gauge - Yukawa theories

$$L = \boxed{-\frac{1}{2}F^2 + i\bar{Q}\gamma_\mu D^\mu Q} + \boxed{y(\bar{Q}_L H Q_R + \text{h.c.})} \quad \text{Yukawa}$$

Gauge  $\text{Tr} [D H^\dagger D H]$  -  $\lambda_u \text{Tr} [(H^\dagger H)^2] - \lambda_v \text{Tr} [(H^\dagger H)]^2$

Scalar selfinteractions

4D: standard model, dark matter, ...

Lower D: condensed matter, phase transitions, graphene

4D plus: extra dimensions, string theory, ...

# Gauge - Yukawa theories

$$L = \boxed{-\frac{1}{2}F^2 + i\bar{Q}\gamma_\mu D^\mu Q} + \boxed{y(\bar{Q}_L H Q_R + \text{h.c.})} \quad \text{Yukawa}$$

Gauge  $\boxed{\text{Tr } [D H^\dagger D H]} - \boxed{\lambda_u \text{Tr } [(H^\dagger H)^2] - \lambda_v \text{Tr } [(H^\dagger H)]^2}$

Scalar selfinteractions

4D: standard model, dark matter, ...

Lower D: condensed matter, phase transitions, graphene

4D plus: extra dimensions, string theory, ...

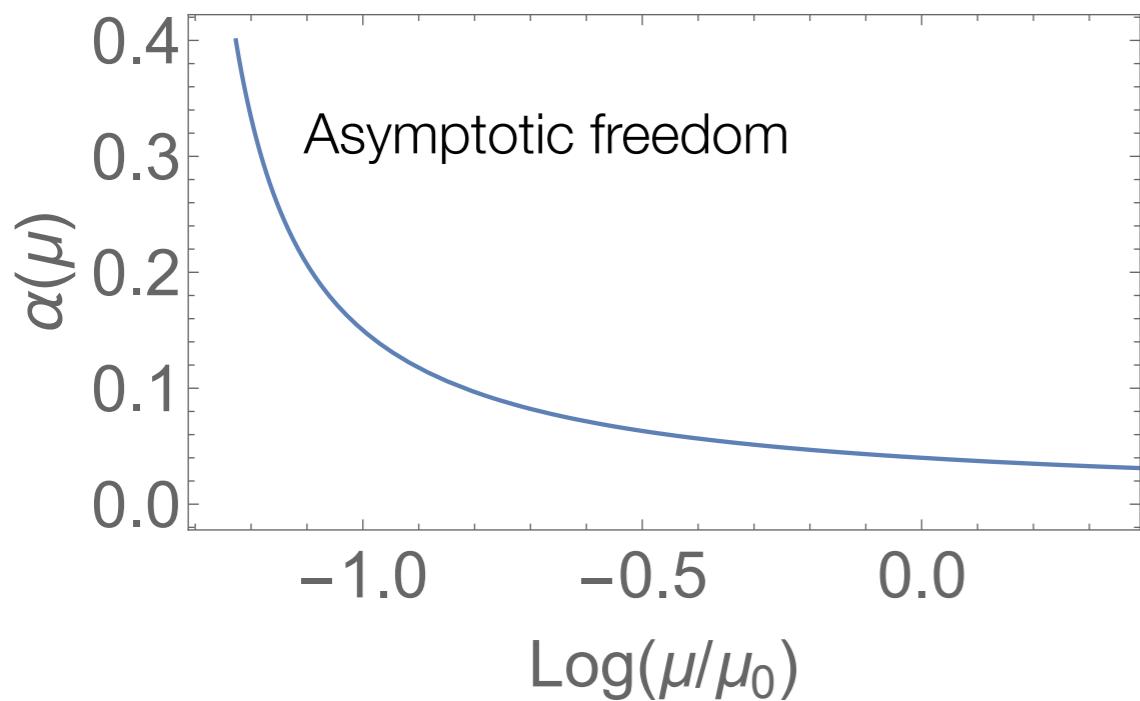
*Universal description of physical phenomena*

# Fundamental theory

*Wilson: A fundamental theory has an UV fixed point*

## Trivial fixed point

- ◆ Non-interacting in the UV
- ◆ Logarithmic scale depend.



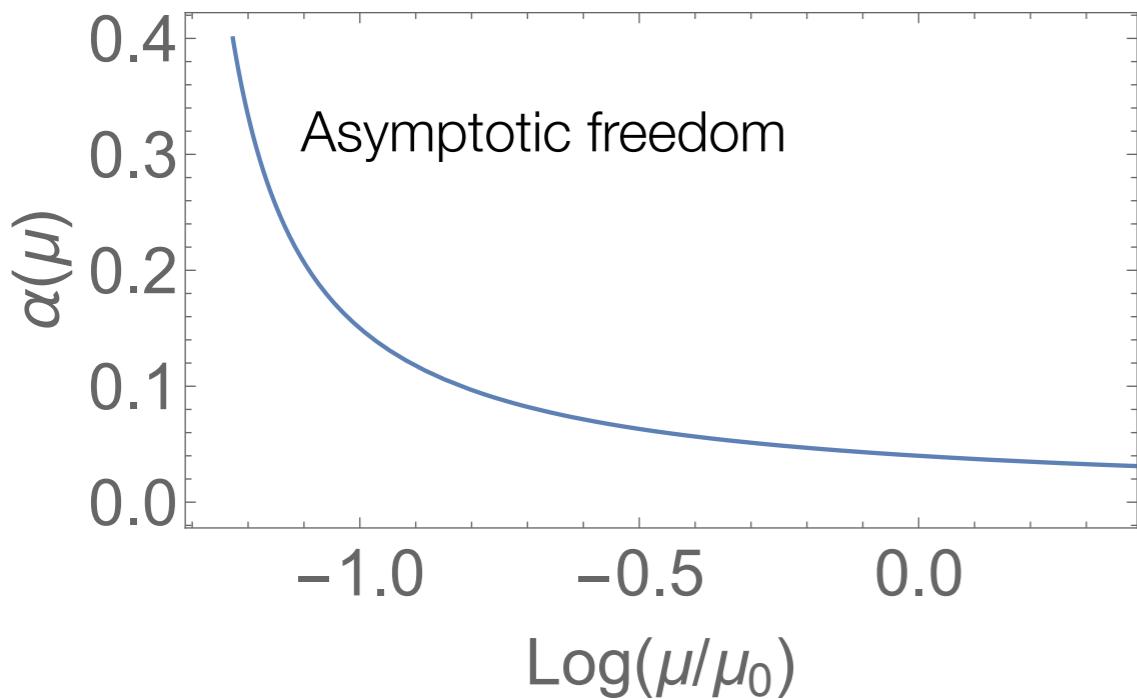
# Fundamental theory

*Wilson: A fundamental theory has an UV fixed point*

Trivial fixed point

Interacting fixed point

- ◆ Non-interacting in the UV
- ◆ Logarithmic scale depend.



# Fundamental theory

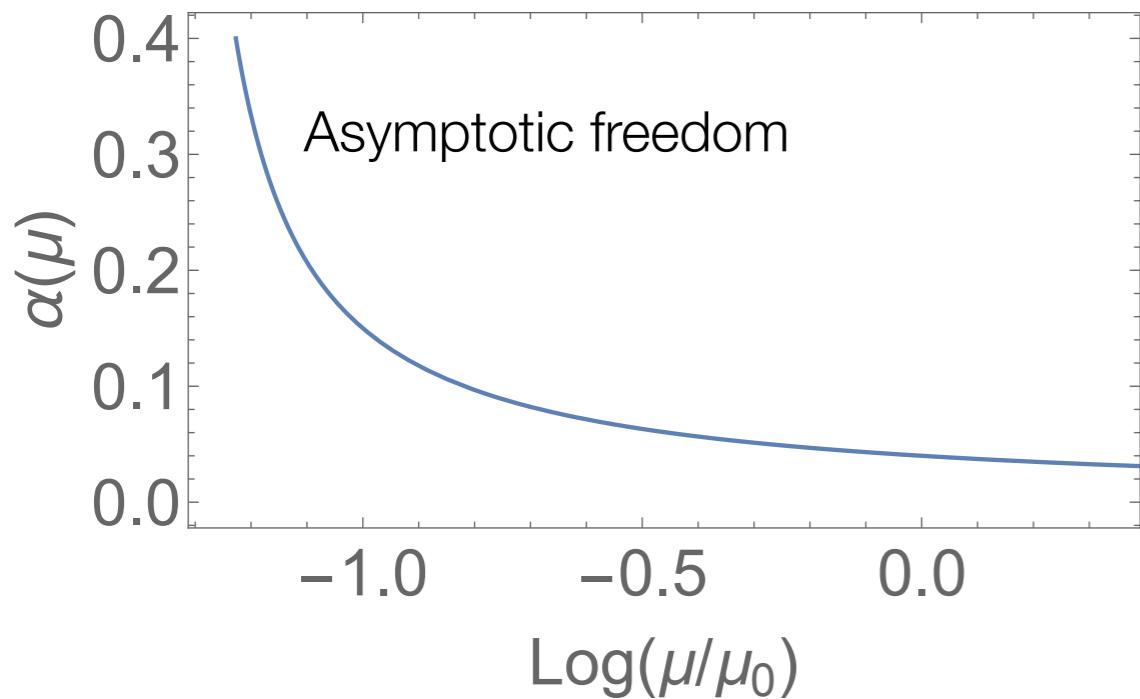
*Wilson: A fundamental theory has an UV fixed point*

## Trivial fixed point

- ◆ Non-interacting in the UV
- ◆ Logarithmic scale depend.

## Interacting fixed point

- ◆ Integrating in the UV
- ◆ Power law

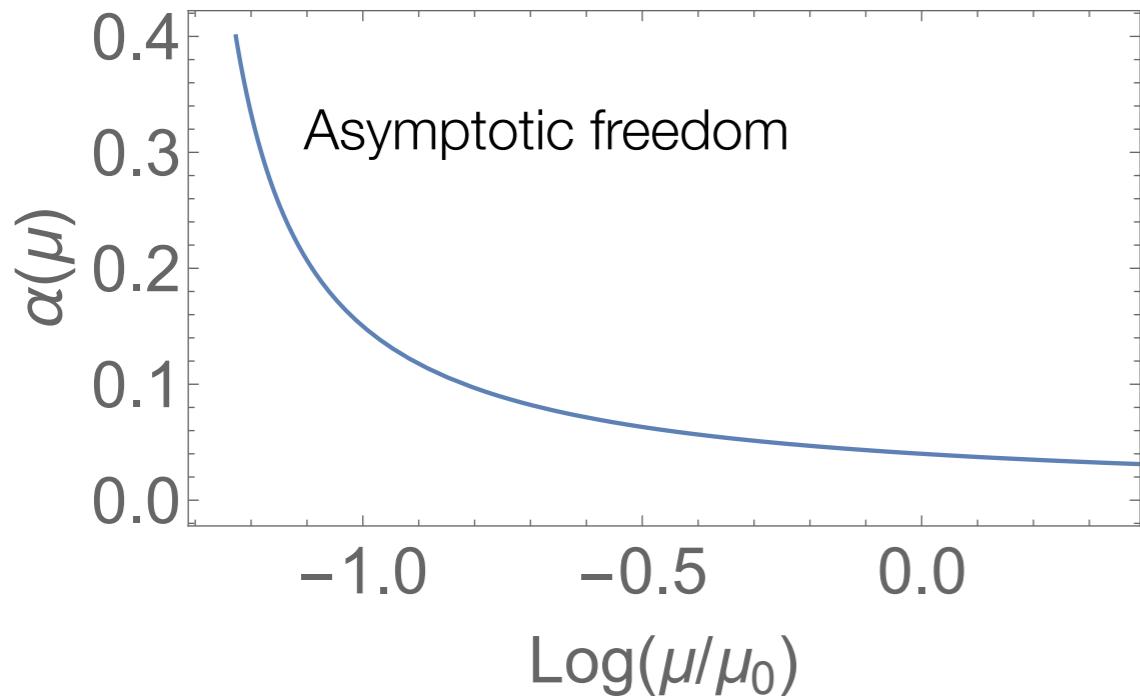


# Fundamental theory

*Wilson: A fundamental theory has an UV fixed point*

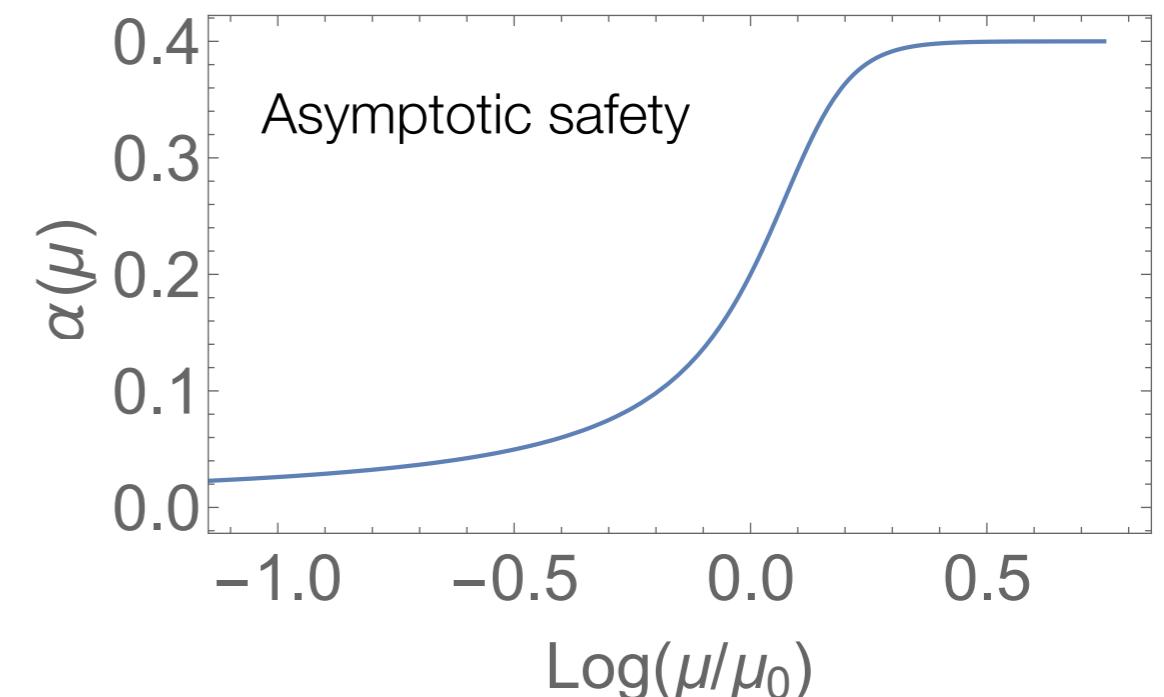
## Trivial fixed point

- ◆ Non-interacting in the UV
- ◆ Logarithmic scale depend.

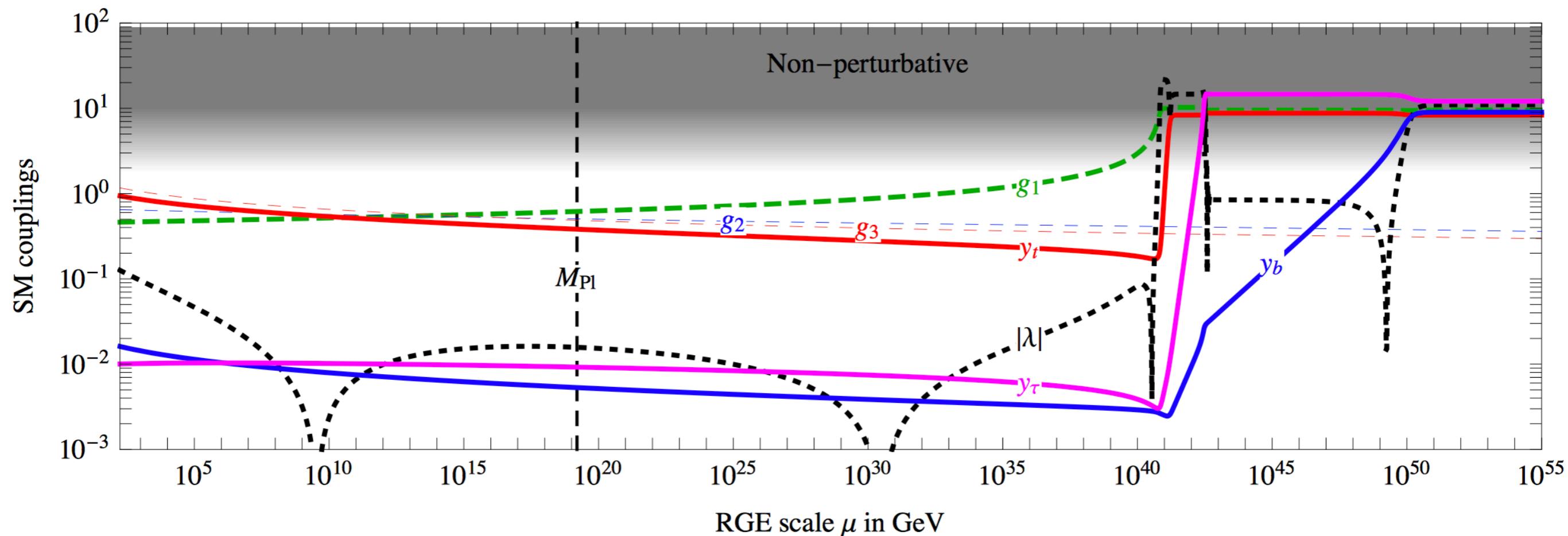


## Interacting fixed point

- ◆ Integrating in the UV
- ◆ Power law

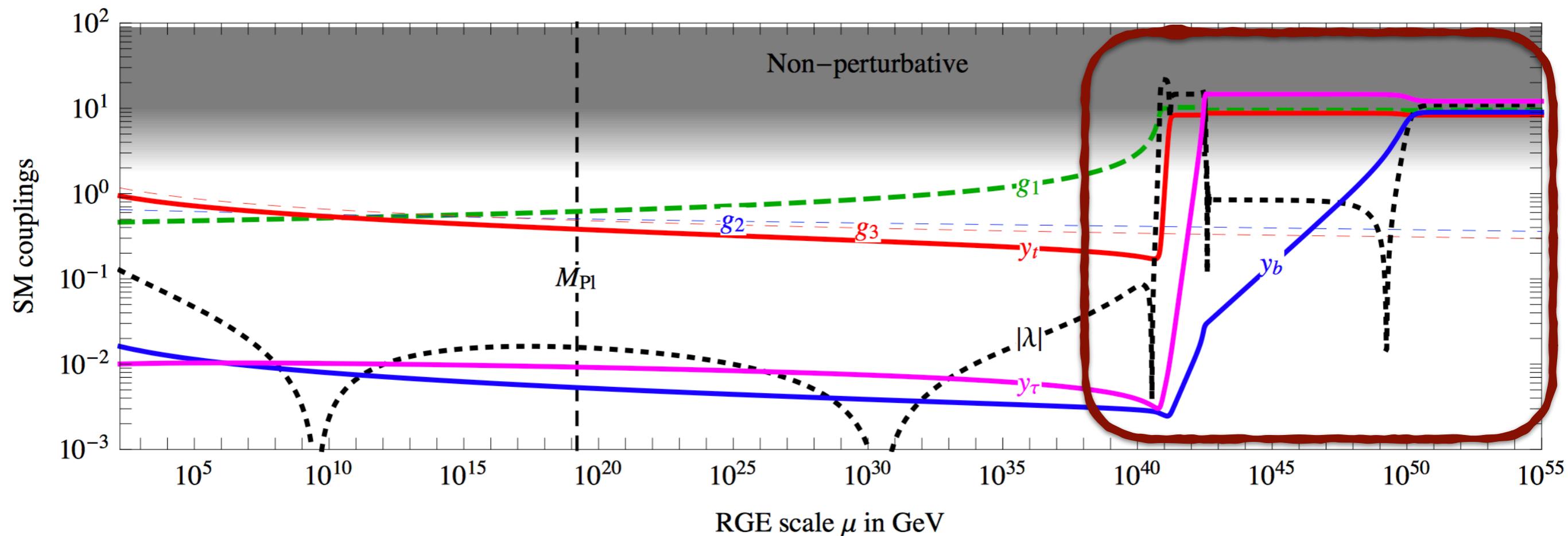


# Is the Standard Model safe?



SM RGE at 3 loops in  $g_{1,2,3}, y_t, \lambda$  and at 2 loops in  $y_{b,\tau}$

# Is the Standard Model safe?



SM RGE at 3 loops in  $g_{1,2,3}$ ,  $y_t$ ,  $\lambda$  and at 2 loops in  $y_{b,\tau}$

# Do theory like these exist?

*Precise and/or nonperturbative exact results for UV interacting fixed points*

# Exact 4D Interacting UV Fixed Point

Antipin, Gillioz, Mølgaard, Sannino 1303.1525 PRD

Litim and Sannino, 1406.2337, JHEP

Litim, Mojaza, Sannino, 1501.03061, JHEP

# Exact 4D Interacting UV Fixed Point

Antipin, Gillioz, Mølgaard, Sannino 1303.1525 PRD

Litim and Sannino, 1406.2337, JHEP

Litim, Mojaza, Sannino, 1501.03061, JHEP

$$L = -F^2 + i\bar{Q}\gamma \cdot DQ + y(\bar{Q}_L H Q_R + \text{h.c.}) +$$

$$\text{Tr} [\partial H^\dagger \partial H] - u \text{Tr} [(H^\dagger H)^2] - v \text{Tr} [(H^\dagger H)]^2$$

Fields	$SU(N_c)$	$SU_L(N_f)$	$SU_R(N_f)$	$U_V(1)$
$G_\mu$	Adj	1	1	0
$Q_L$	□	□	1	1
$Q_R^c$	□	1	□	-1
$H$	1	□	□	0

# Veneziano Limit

Litim and Sannino, 1406.2337, JHEP

Litim, Mojaza, Sannino, 1501.03061, JHEP

- ◆ Normalised couplings

$$\alpha_g = \frac{g^2 N_C}{(4\pi)^2}, \quad \alpha_y = \frac{y^2 N_C}{(4\pi)^2}, \quad \alpha_h = \frac{u N_F}{(4\pi)^2}, \quad \alpha_v = \frac{v N_F^2}{(4\pi)^2}$$

$$\frac{v}{u} = \frac{\alpha_v}{\alpha_h N_F}$$

At large  $N$

$$\frac{N_F}{N_C} \in \Re^+$$

# Veneziano Limit

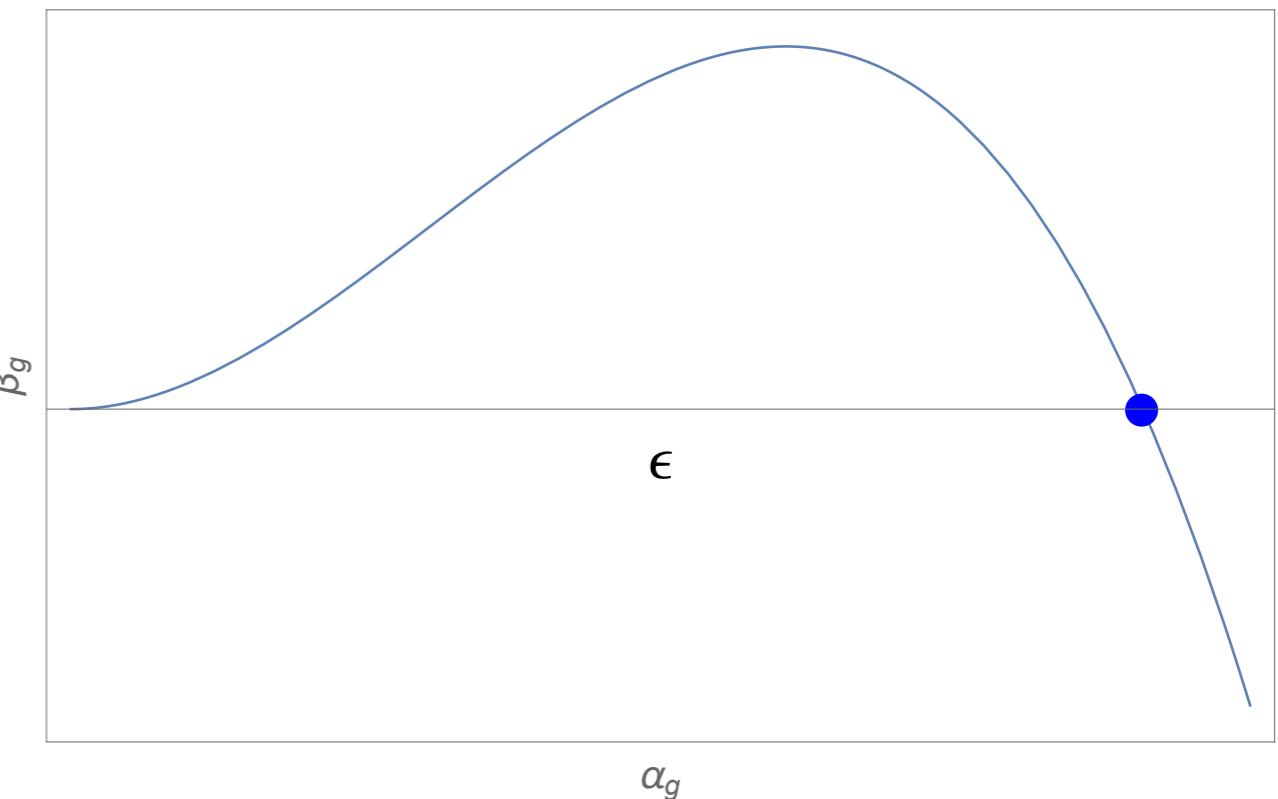
- ◆ Normalised couplings

$$\alpha_g = \frac{g^2 N_C}{(4\pi)^2}, \quad \alpha_y = \frac{y^2 N_C}{(4\pi)^2}, \quad \alpha_h = \frac{u N_F}{(4\pi)^2}, \quad \alpha_v = \frac{v N_F^2}{(4\pi)^2}$$

$$\frac{v}{u} = \frac{\alpha_v}{\alpha_h N_F}$$

At large  $N$

$$\frac{N_F}{N_C} \in \Re^+$$



# Veneziano Limit

- ◆ Normalised couplings

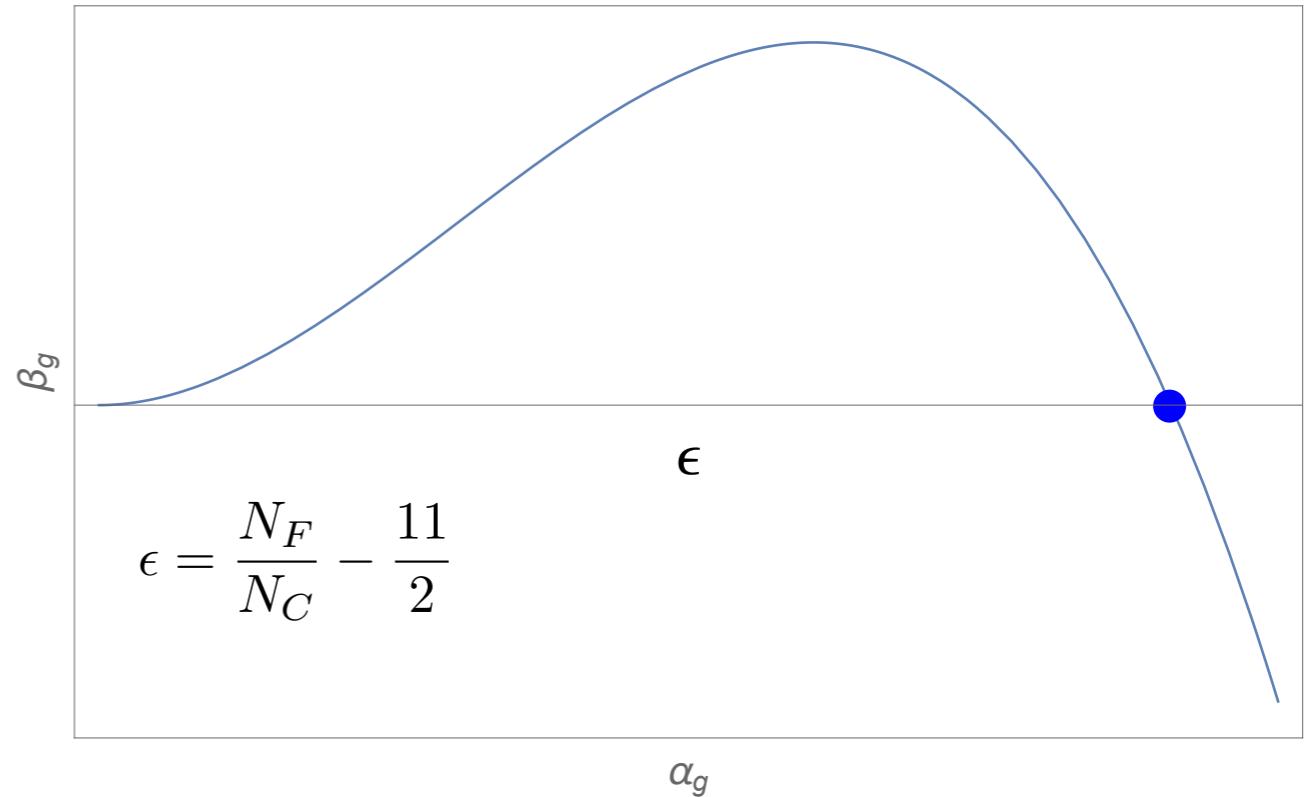
$$\alpha_g = \frac{g^2 N_C}{(4\pi)^2}, \quad \alpha_y = \frac{y^2 N_C}{(4\pi)^2}, \quad \alpha_h = \frac{u N_F}{(4\pi)^2}, \quad \alpha_v = \frac{v N_F^2}{(4\pi)^2}$$

$$\frac{v}{u} = \frac{\alpha_v}{\alpha_h N_F}$$

At large  $N$

$$\frac{N_F}{N_C} \in \Re^+$$

$$\epsilon = \frac{N_F}{N_C} - \frac{11}{2}$$



# Veneziano Limit

- ◆ Normalised couplings

$$\alpha_g = \frac{g^2 N_C}{(4\pi)^2}, \quad \alpha_y = \frac{y^2 N_C}{(4\pi)^2}, \quad \alpha_h = \frac{u N_F}{(4\pi)^2}, \quad \alpha_v = \frac{v N_F^2}{(4\pi)^2}$$

$$\frac{v}{u} = \frac{\alpha_v}{\alpha_h N_F}$$

At large N

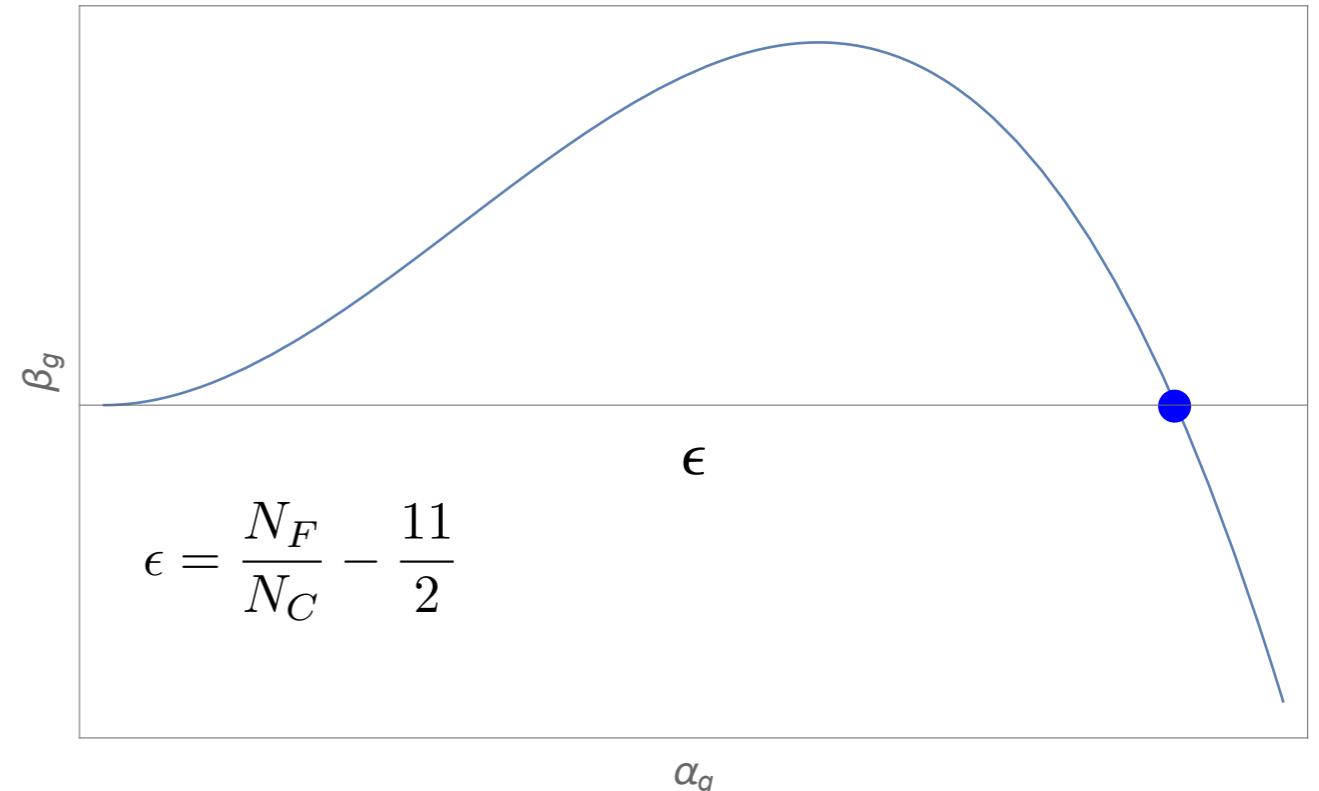
$$\frac{N_F}{N_C} \in \Re^+$$

Impossible in Gauge Theories with Fermions alone

Caswell, PRL 1974

Litim and Sannino, 1406.2337, JHEP

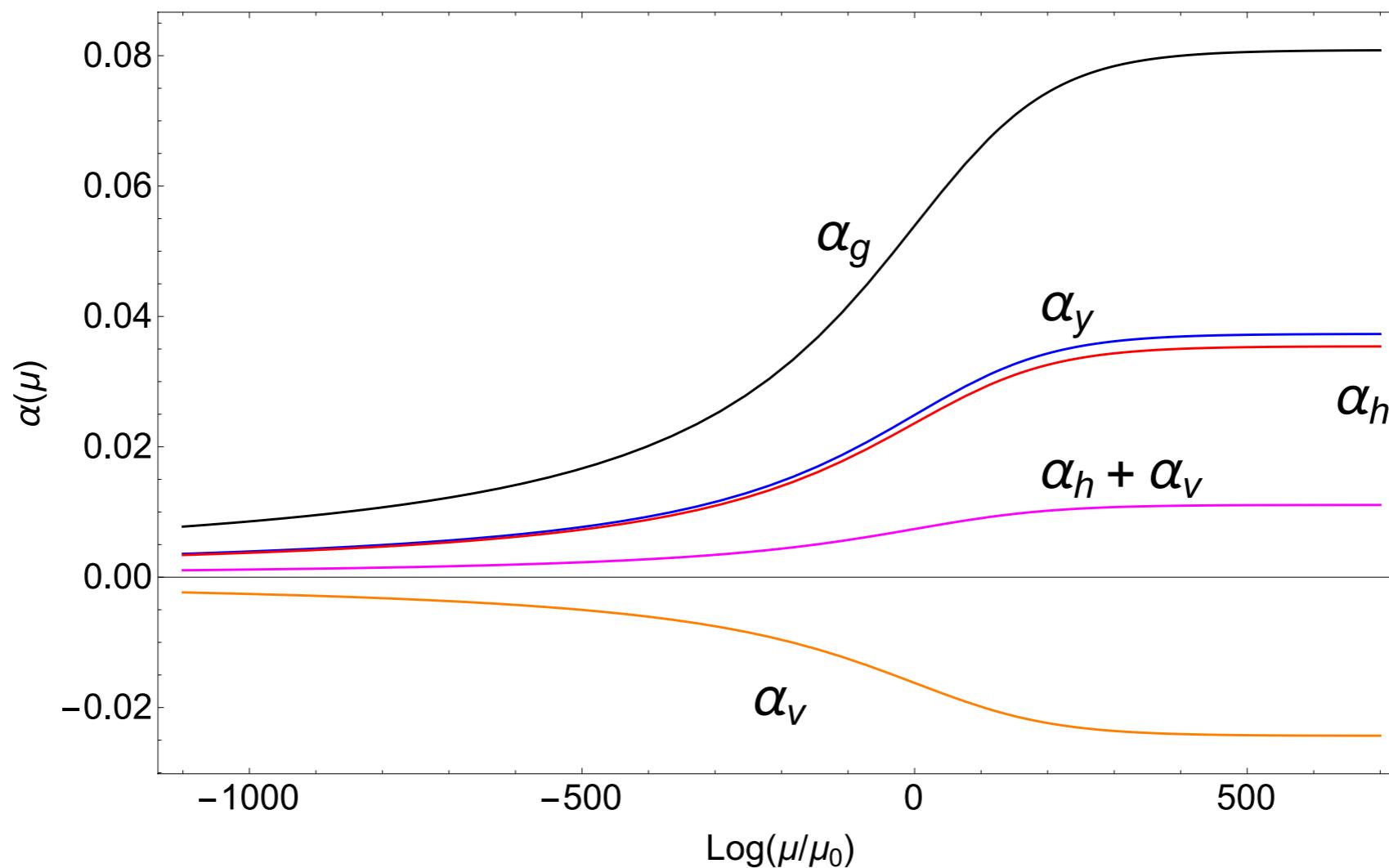
Litim, Mojaza, Sannino, 1501.03061, JHEP



# Complete asymptotic safety

Litim and Sannino, 1406.2337, JHEP

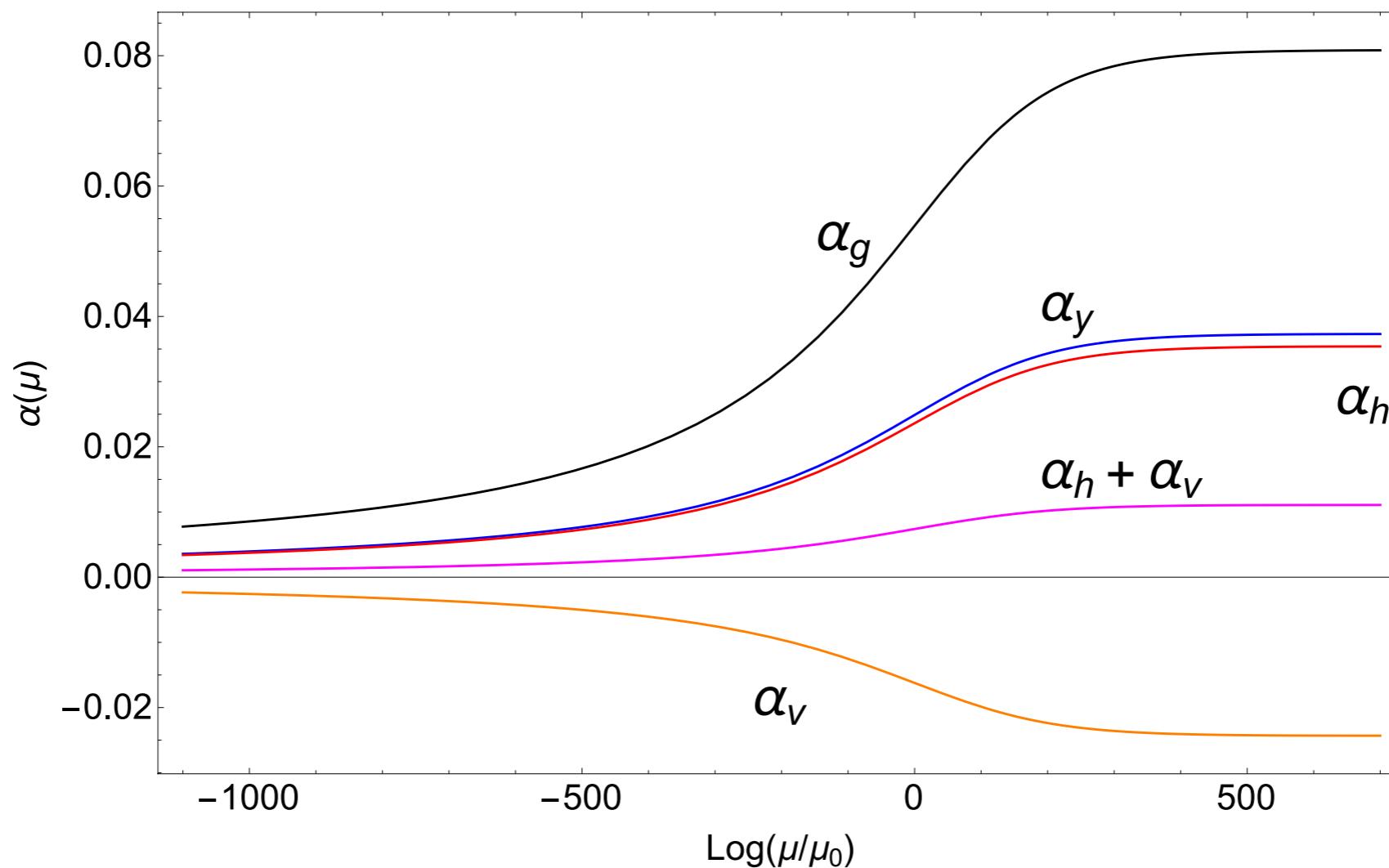
Gauge + fermion + scalars theories can be fund. at any energy scale



# Complete asymptotic safety

Litim and Sannino, 1406.2337, JHEP

Gauge + fermion + scalars theories can be fund. at any energy scale

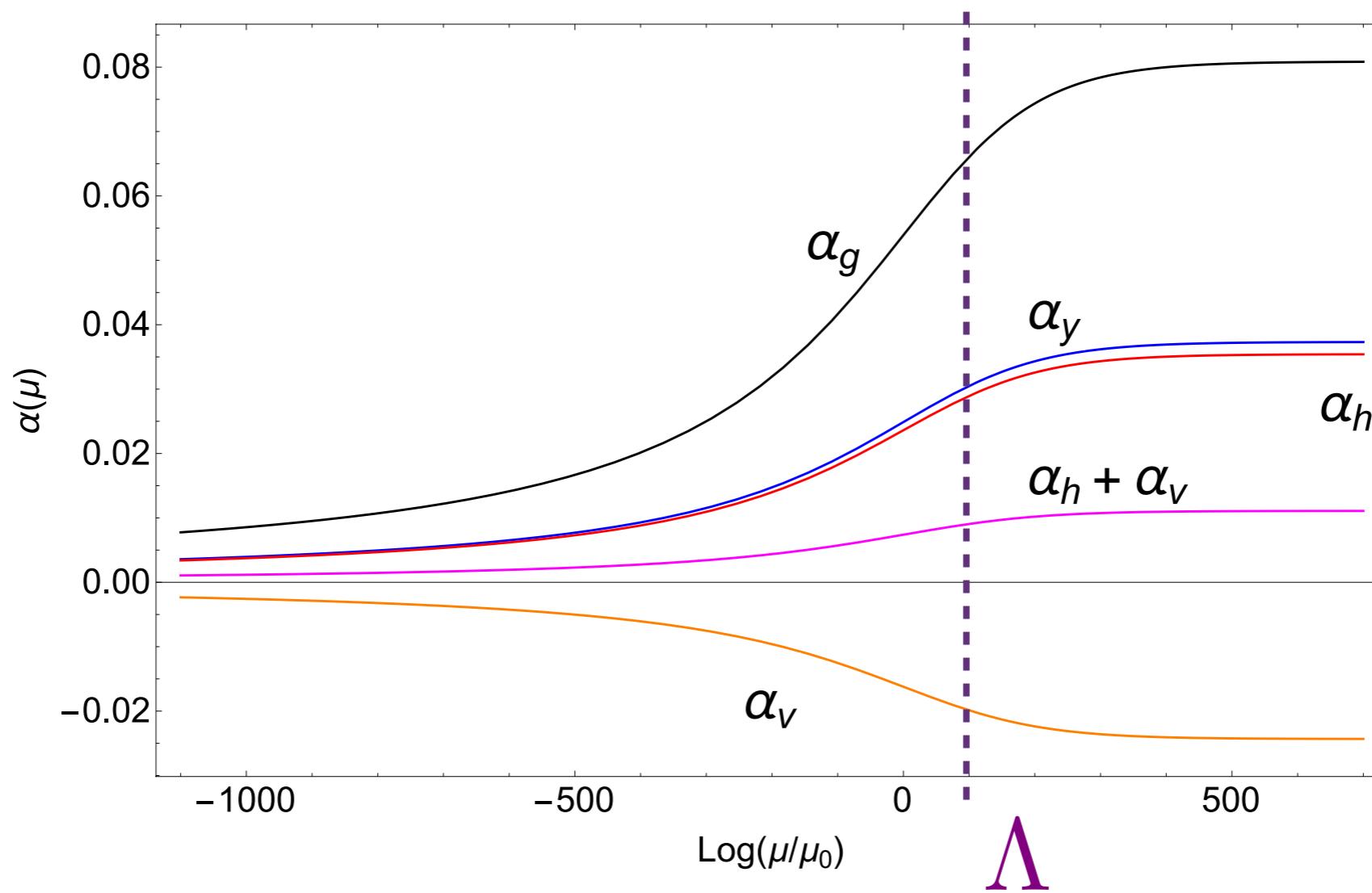


Scalars are needed perturbatively to make the theory fundamental

# Complete asymptotic safety

Litim and Sannino, 1406.2337, JHEP

Gauge + fermion + scalars theories can be fund. at any energy scale



Scalars are needed perturbatively to make the theory fundamental

# Gauged Higgs UV Fixed Point

Pelaggi, Sannino, Strumia, Vigiani, 1701.01453

Fields	Gauge symmetries		Global symmetries	
	Spin	$SU(N_c)$	$U(N_F)_L$	$U(N_F)_R$
$\psi$	1/2	□	□	1
$\bar{\psi}$	1/2	□	1	□
$S$	0	1	□	□
$H$	0	□	1	1
$N$	1/2	1	1	□
$N'$	1/2	1	□	1

# Gauged Higgs UV Fixed Point

Pelaggi, Sannino, Strumia, Vigiani, 1701.01453

Fields	Gauge symmetries		Global symmetries	
	Spin	$SU(N_c)$	$U(N_F)_L$	$U(N_F)_R$
$\psi$	1/2	□	□	1
$\bar{\psi}$	1/2	□	1	□
$S$	0	1	□	□
$H$	0	□	1	1
$N$	1/2	1	1	□
$N'$	1/2	1	□	1

$$V = \lambda_{S1} (\text{Tr} S^\dagger S)^2 + \lambda_{S2} \text{Tr}(S^\dagger S S^\dagger S) + \lambda_H (H^\dagger H)^2 + \lambda_{HS} (H^\dagger H) \text{Tr}(S S^\dagger)$$

# Gauged Higgs UV Fixed Point

Pelaggi, Sannino, Strumia, Vigiani, 1701.01453

Fields	Gauge symmetries		Global symmetries	
	Spin	$SU(N_c)$	$U(N_F)_L$	$U(N_F)_R$
$\psi$	1/2	□	□	1
$\bar{\psi}$	1/2	□	1	□
$S$	0	1	□	□
$H$	0	□	1	1
$N$	1/2	1	1	□
$N'$	1/2	1	□	1

$$V = \lambda_{S1} (\text{Tr} S^\dagger S)^2 + \lambda_{S2} \text{Tr}(S^\dagger S S^\dagger S) + \lambda_H (H^\dagger H)^2 + \lambda_{HS} (H^\dagger H) \text{Tr}(S S^\dagger)$$

$$\mathcal{L}_Y = y S_{ij} \psi_i \bar{\psi}_j + y' S_{ij}^* N_i N'_j + \tilde{y} H \bar{\psi}_i N_i + \tilde{y}' H^* \psi_i N'_i + \text{h.c.}$$

# Gauged Higgs UV Fixed Point

Pelaggi, Sannino, Strumia, Vigiani, 1701.01453

Fields	Gauge symmetries		Global symmetries	
	Spin	$SU(N_c)$	$U(N_F)_L$	$U(N_F)_R$
$\psi$	1/2	□	□	1
$\bar{\psi}$	1/2	□	1	□
$S$	0	1	□	□
$H$	0	□	1	1
$N$	1/2	1	1	□
$N'$	1/2	1	□	1

$$V = \lambda_{S1} (\text{Tr} S^\dagger S)^2 + \lambda_{S2} \text{Tr}(S^\dagger S S^\dagger S) + \lambda_H (H^\dagger H)^2 + \lambda_{HS} (H^\dagger H) \text{Tr}(S S^\dagger)$$

$$\mathcal{L}_Y = y S_{ij} \psi_i \bar{\psi}_j + y' S_{ij}^* N_i N'_j + \boxed{\tilde{y} H \bar{\psi}_i N_i + \tilde{y}' H^* \psi_i N'_i + \text{h.c.}}$$

# Gauged Higgs UV Fixed Point

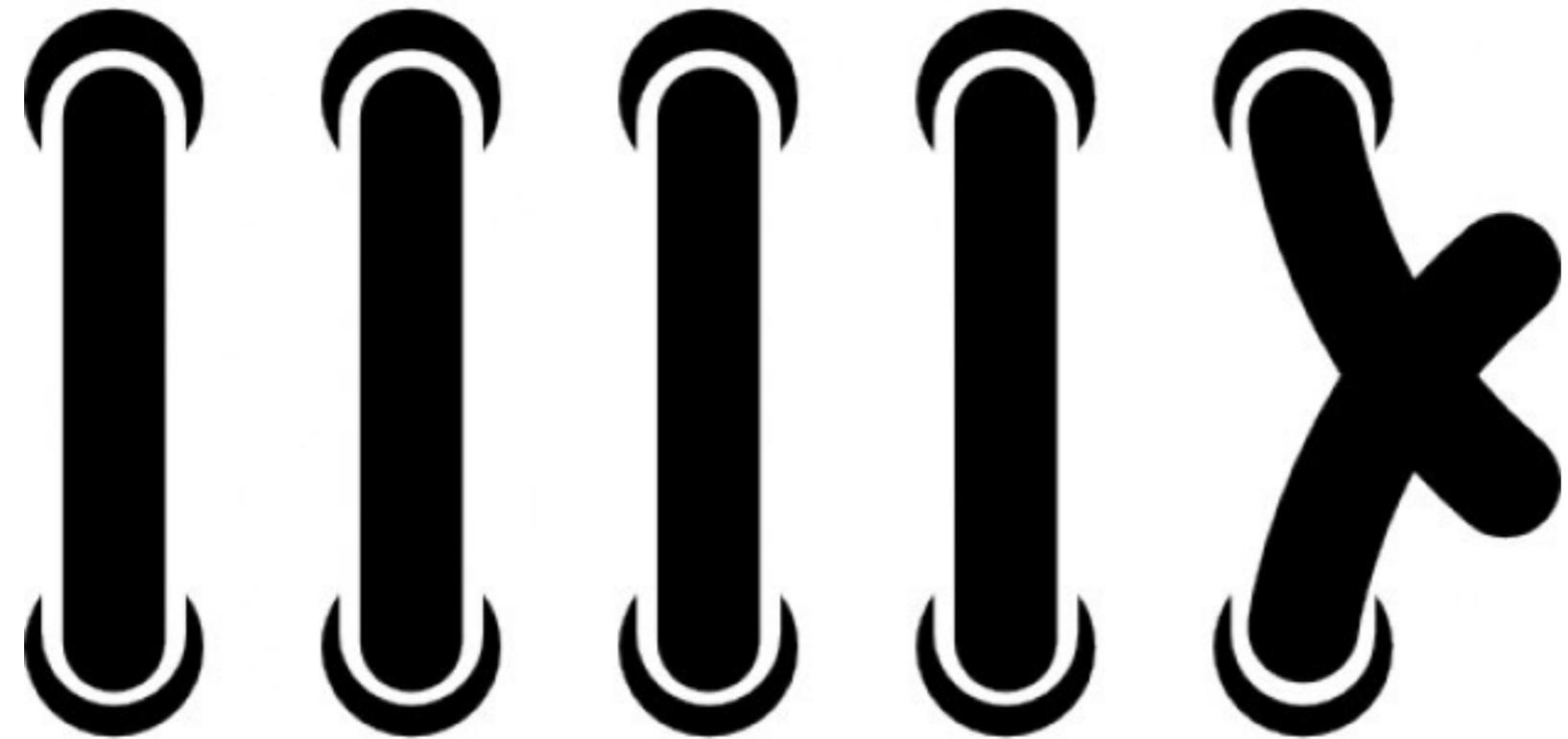
Pelaggi, Sannino, Strumia, Vigiani, 1701.01453

Fields	Gauge symmetries		Global symmetries	
	Spin	$SU(N_c)$	$U(N_F)_L$	$U(N_F)_R$
$\psi$	1/2	□	□	1
$\bar{\psi}$	1/2	□	1	□
$S$	0	1	□	□
$H$	0	□	1	1
$N$	1/2	1	1	□
$N'$	1/2	1	□	1

$$V = \lambda_{S1} (\text{Tr} S^\dagger S)^2 + \lambda_{S2} \text{Tr}(S^\dagger S S^\dagger S) + \lambda_H (H^\dagger H)^2 + \lambda_{HS} (H^\dagger H) \text{Tr}(S S^\dagger)$$

$$\mathcal{L}_Y = y S_{ij} \psi_i \bar{\psi}_j + y' S_{ij}^* N_i N'_j + \boxed{\tilde{y} H \bar{\psi}_i N_i + \tilde{y}' H^* \psi_i N'_i + \text{h.c.}}$$

*Controllably safe in all couplings*



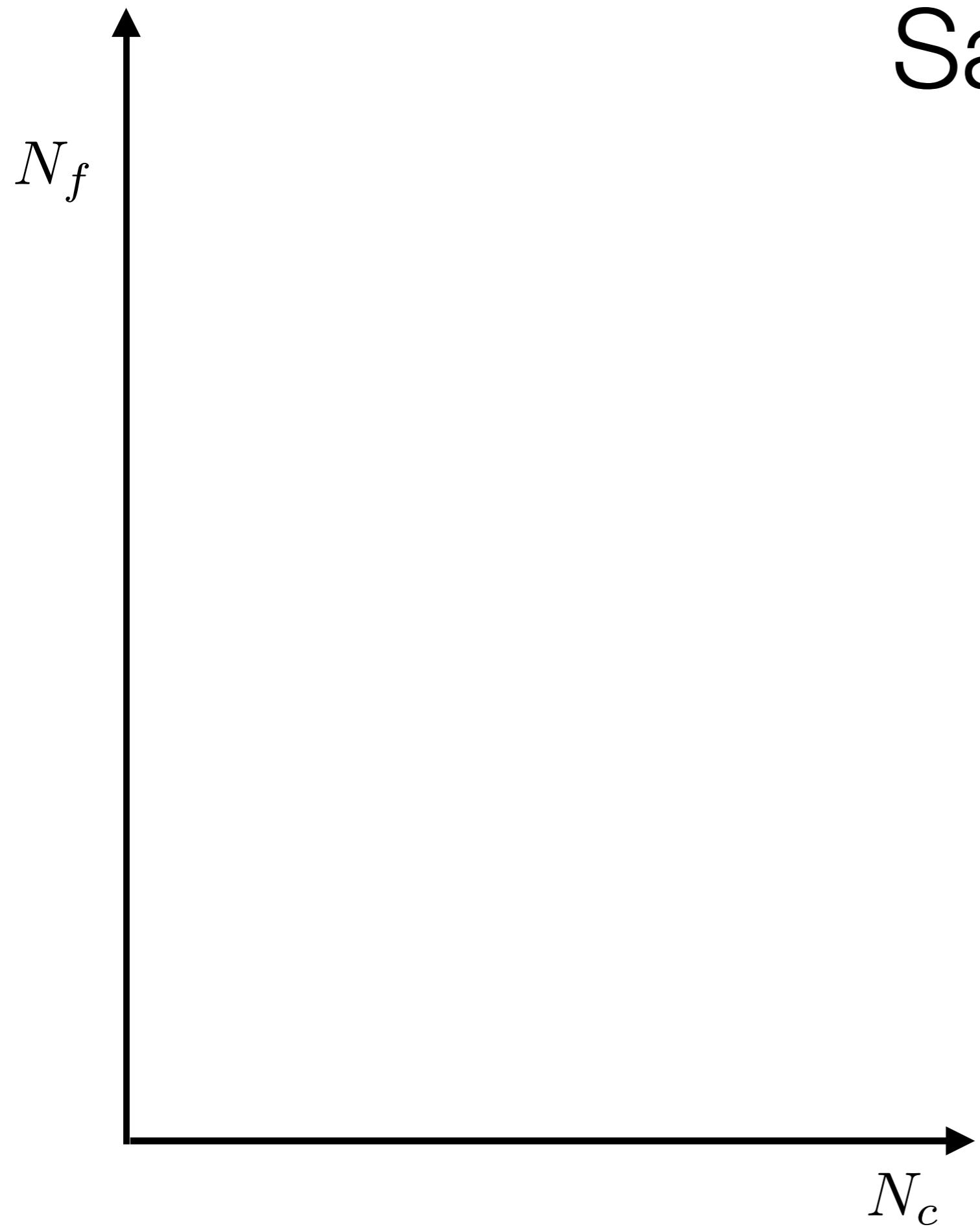
*Higgs as shoelace*

# Conformal Window 2.0: Large Nf Story

Sannino, ERG 2016, Heidelberg

Antipin and Sannino to appear

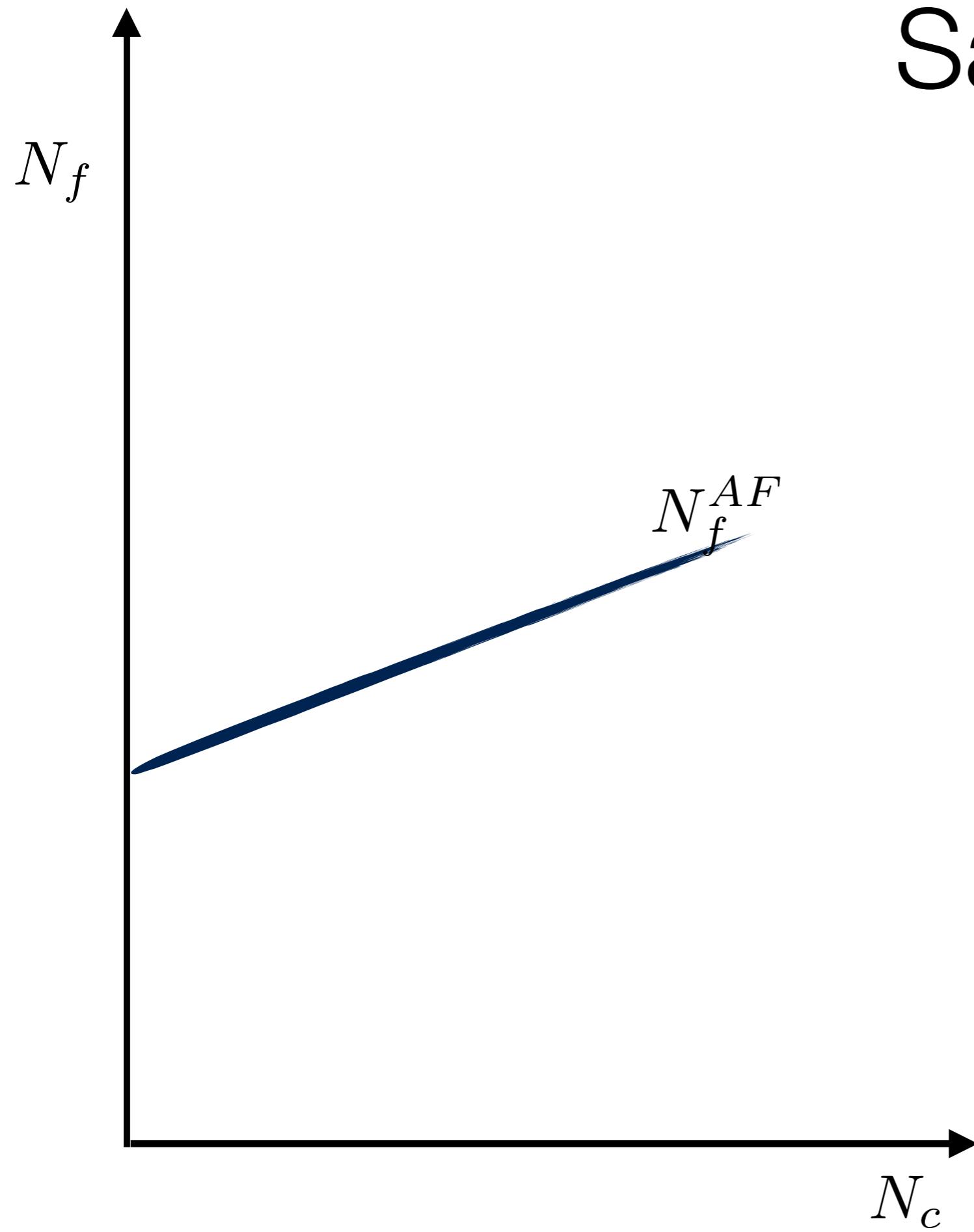
# Safe QCD



Sannino, ERG 2016, Heidelberg

Antipin and Sannino, to appear  
Pica and Sannino 1011.5917, PRD

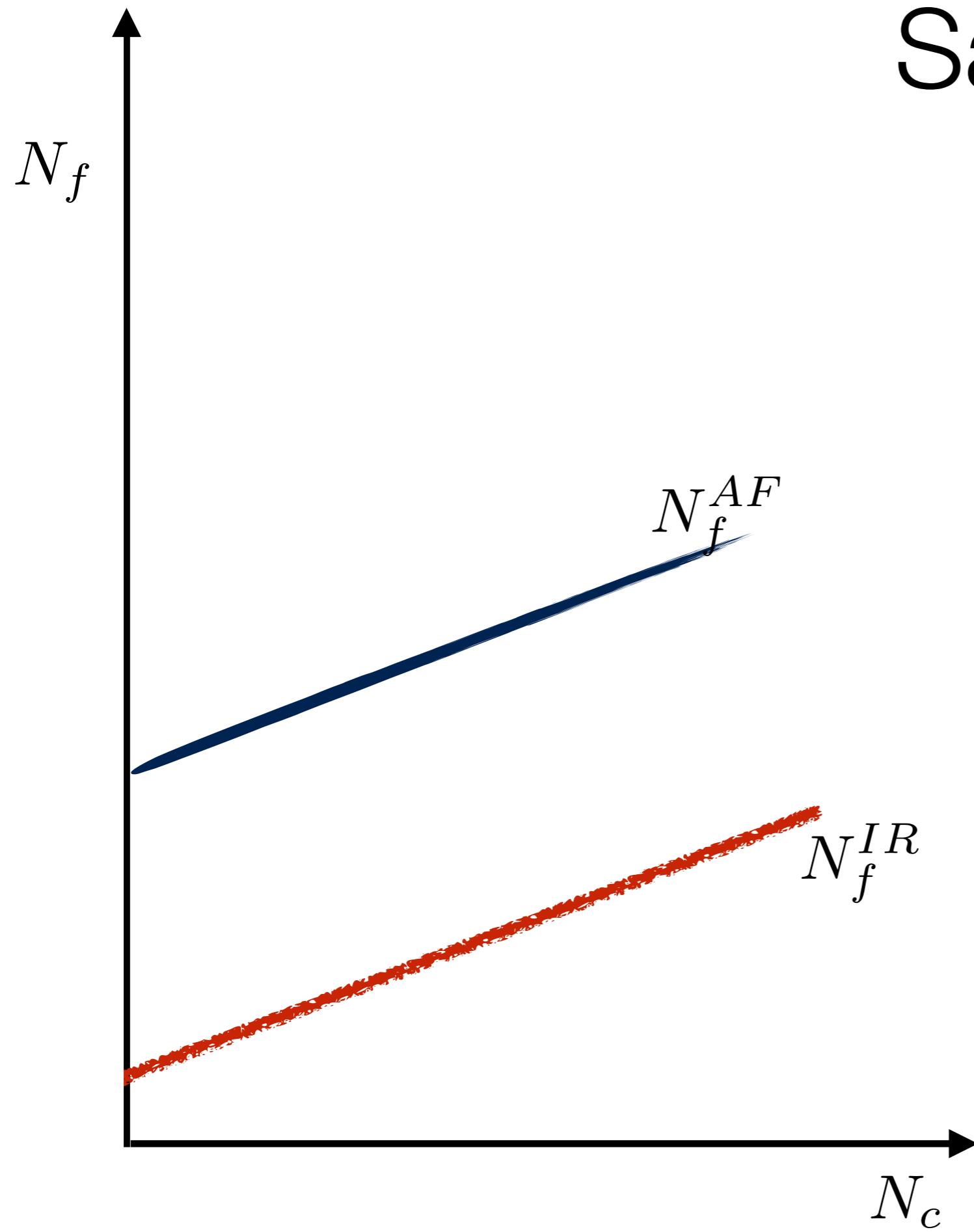
# Safe QCD



Sannino, ERG 2016, Heidelberg

Antipin and Sannino, to appear  
Pica and Sannino 1011.5917, PRD

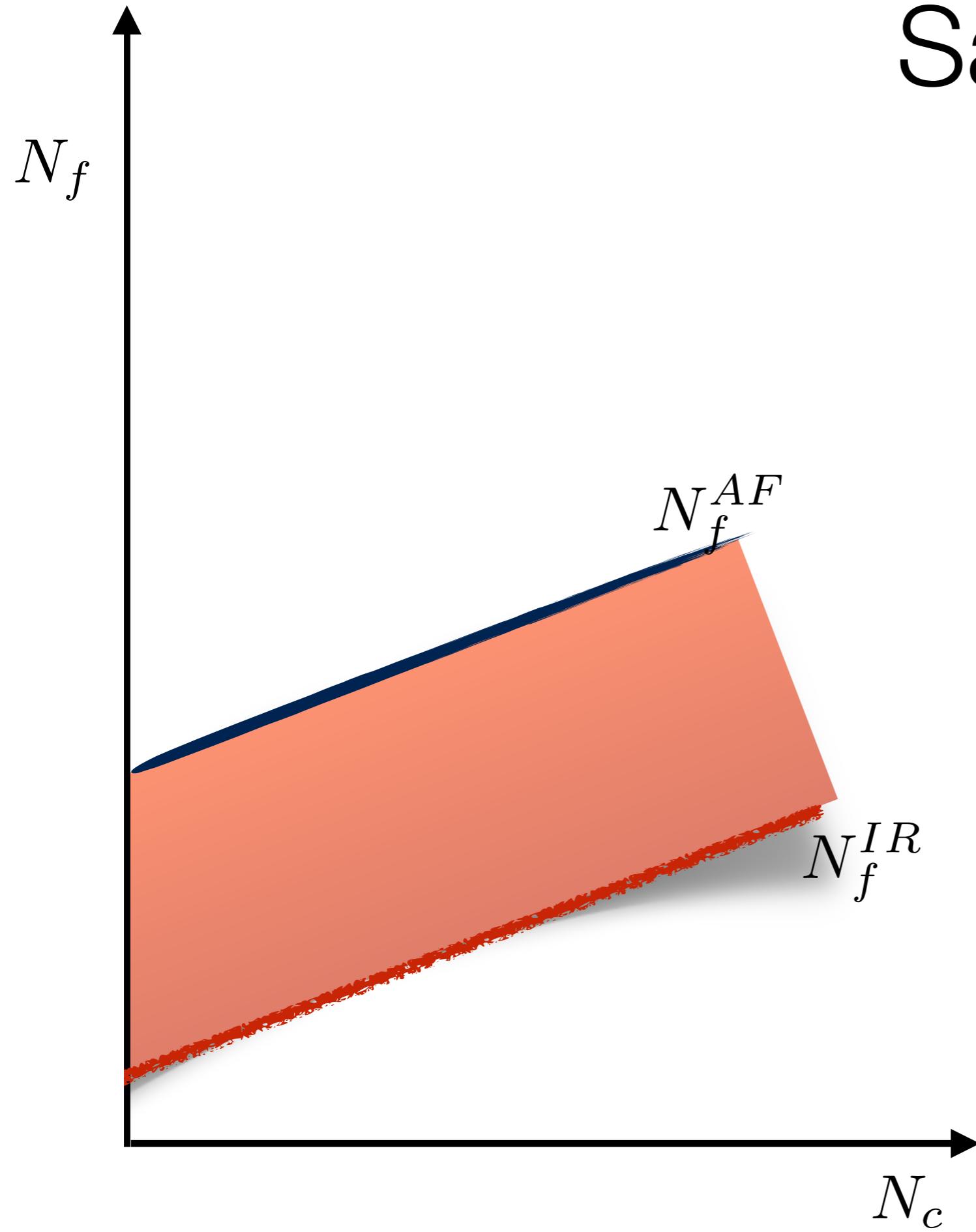
# Safe QCD



Sannino, ERG 2016, Heidelberg

Antipin and Sannino, to appear  
Pica and Sannino 1011.5917, PRD

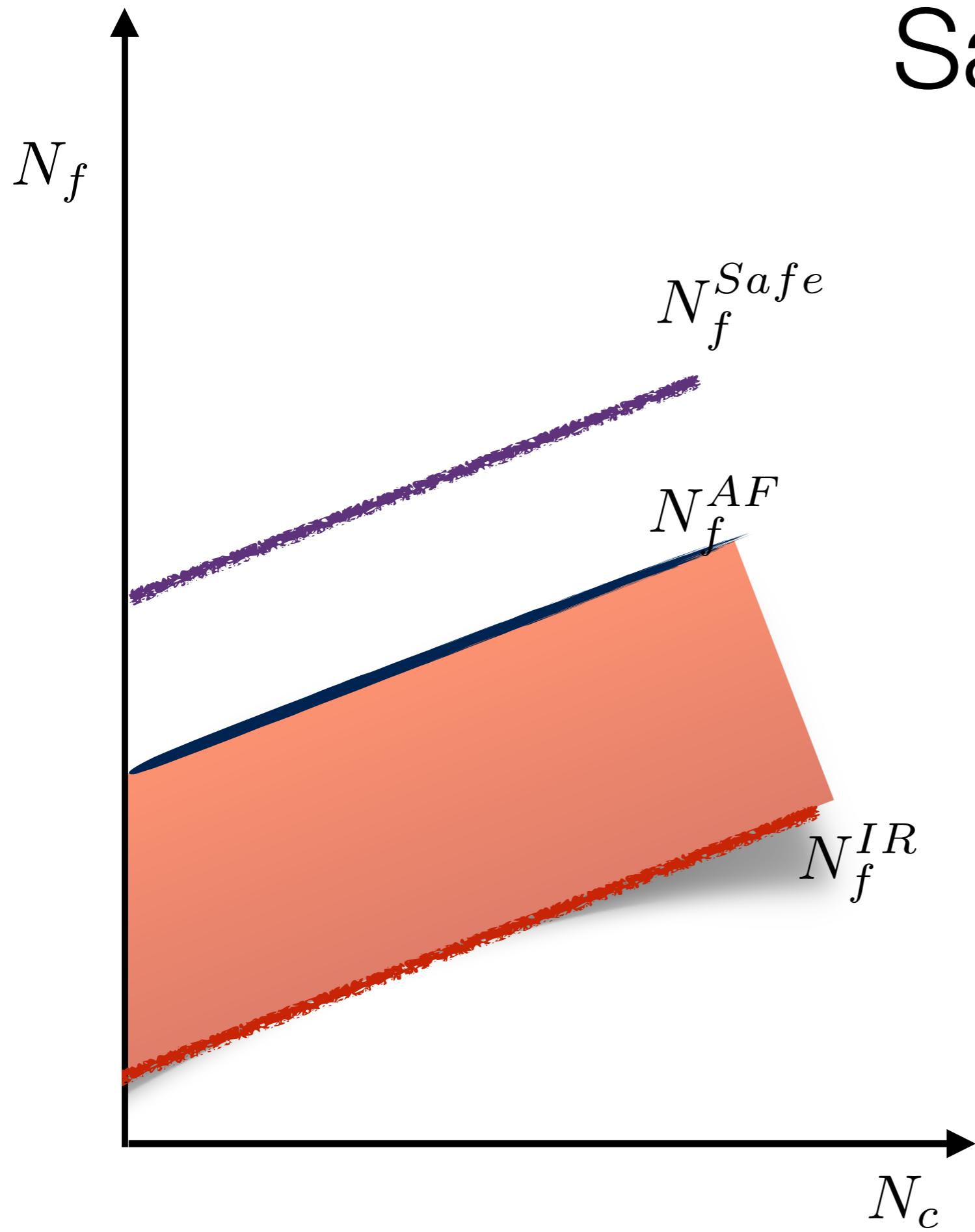
# Safe QCD



Sannino, ERG 2016, Heidelberg

Antipin and Sannino, to appear  
Pica and Sannino 1011.5917, PRD

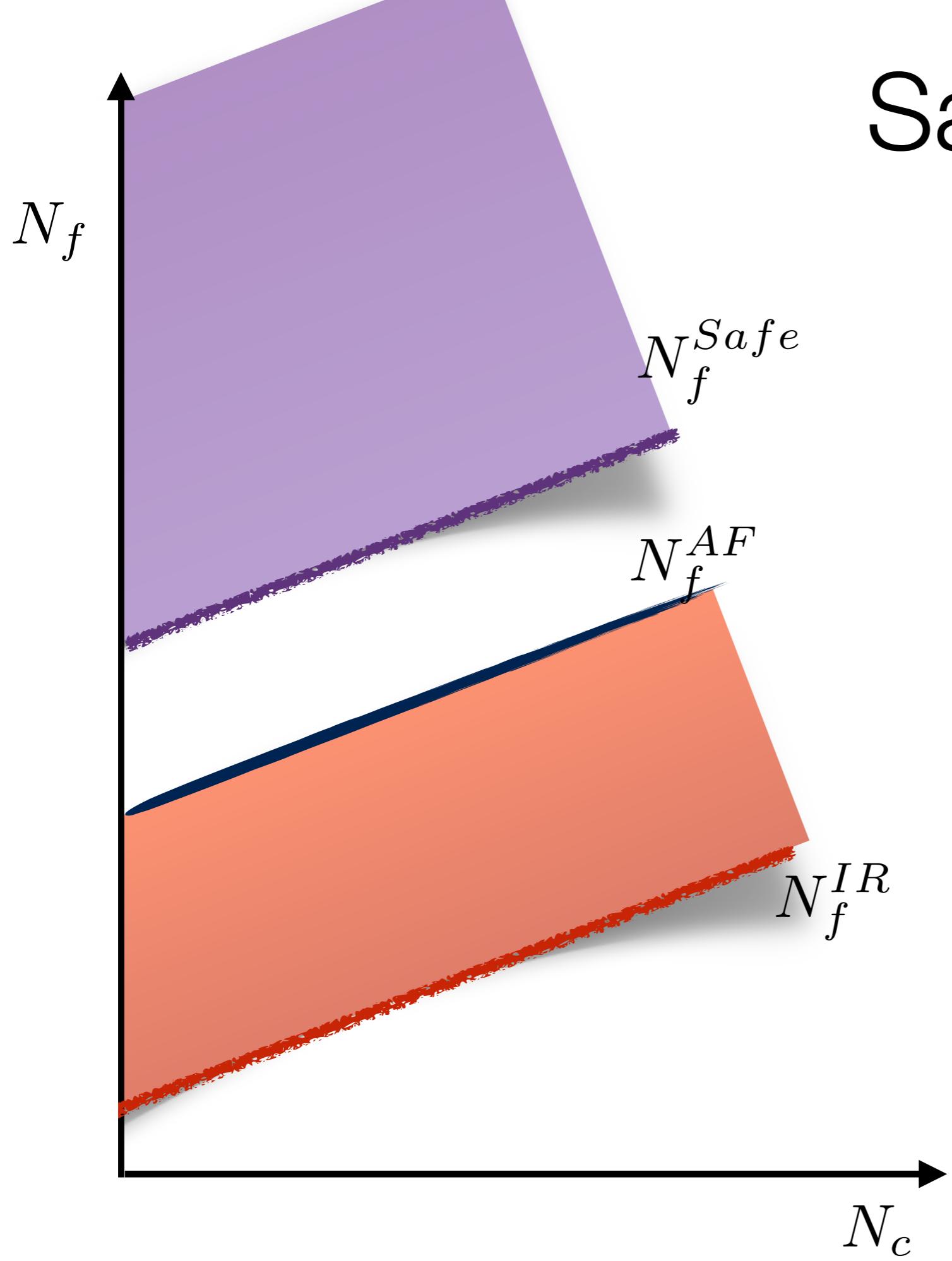
# Safe QCD



Sannino, ERG 2016, Heidelberg

Antipin and Sannino, to appear  
Pica and Sannino 1011.5917, PRD

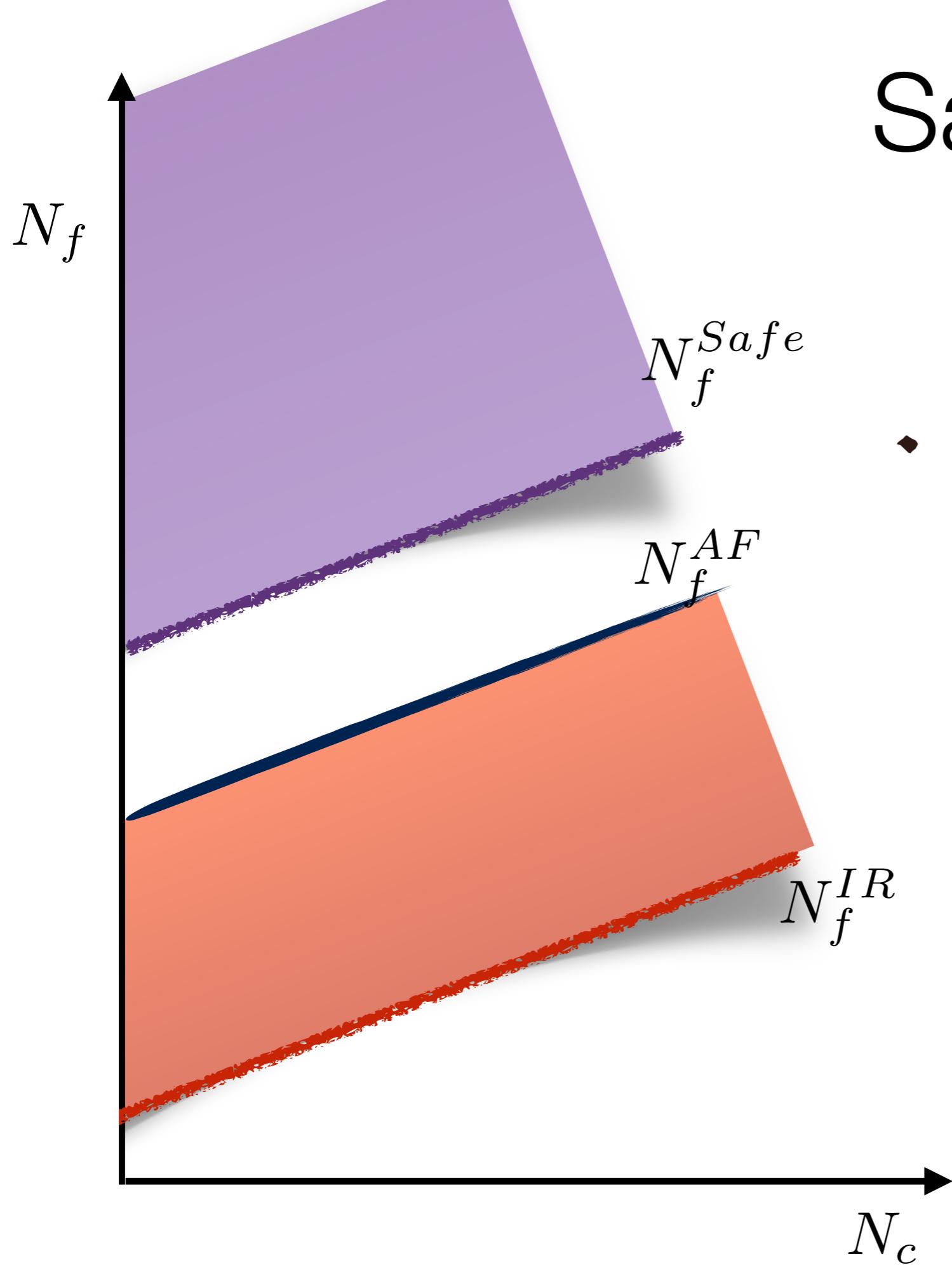
# Safe QCD



Sannino, ERG 2016, Heidelberg

Antipin and Sannino, to appear  
Pica and Sannino 1011.5917, PRD

# Safe QCD

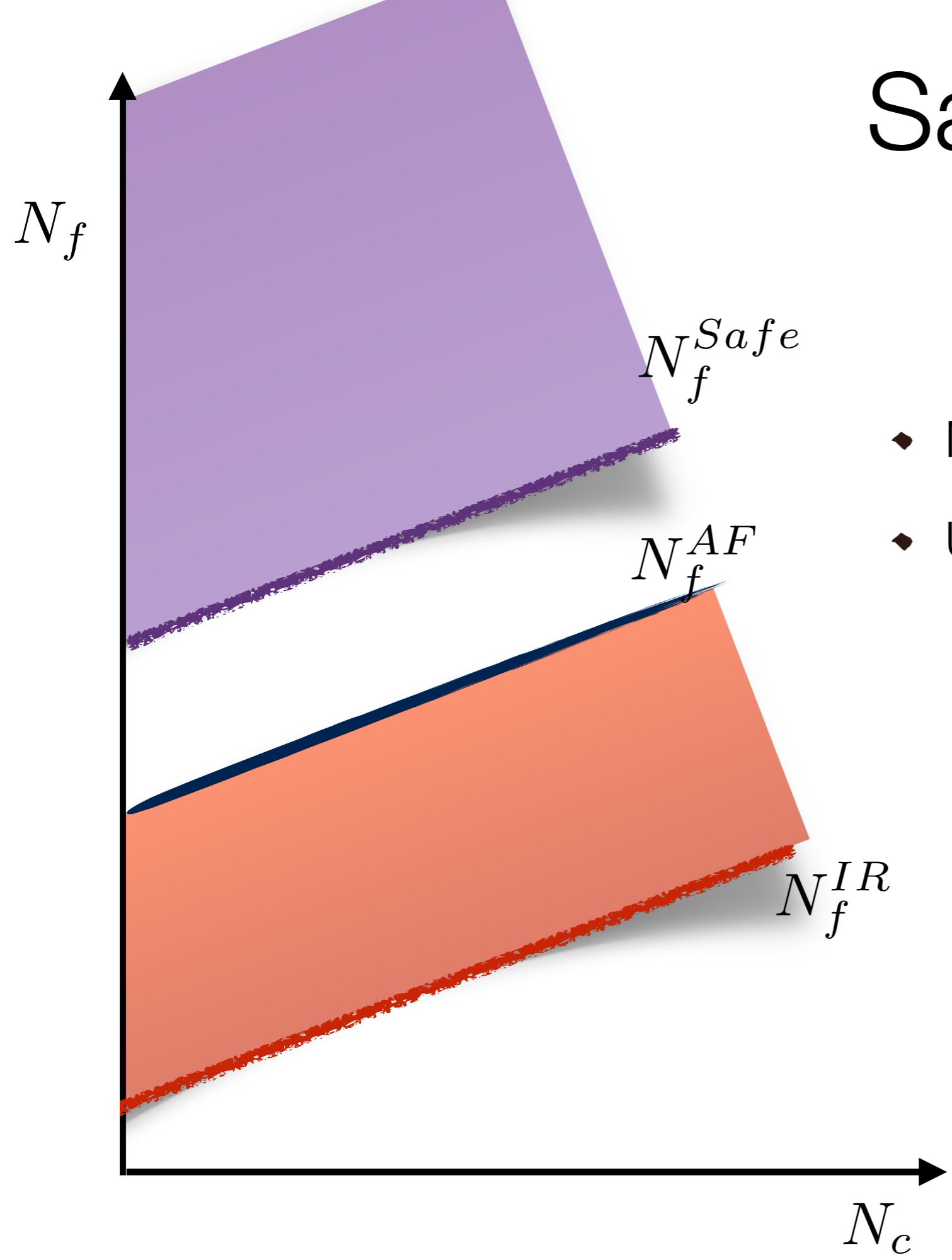


- ◆ Must exist a critical Safe  $N_f$

Sannino, ERG 2016, Heidelberg

Antipin and Sannino, to appear  
Pica and Sannino 1011.5917, PRD

# Safe QCD

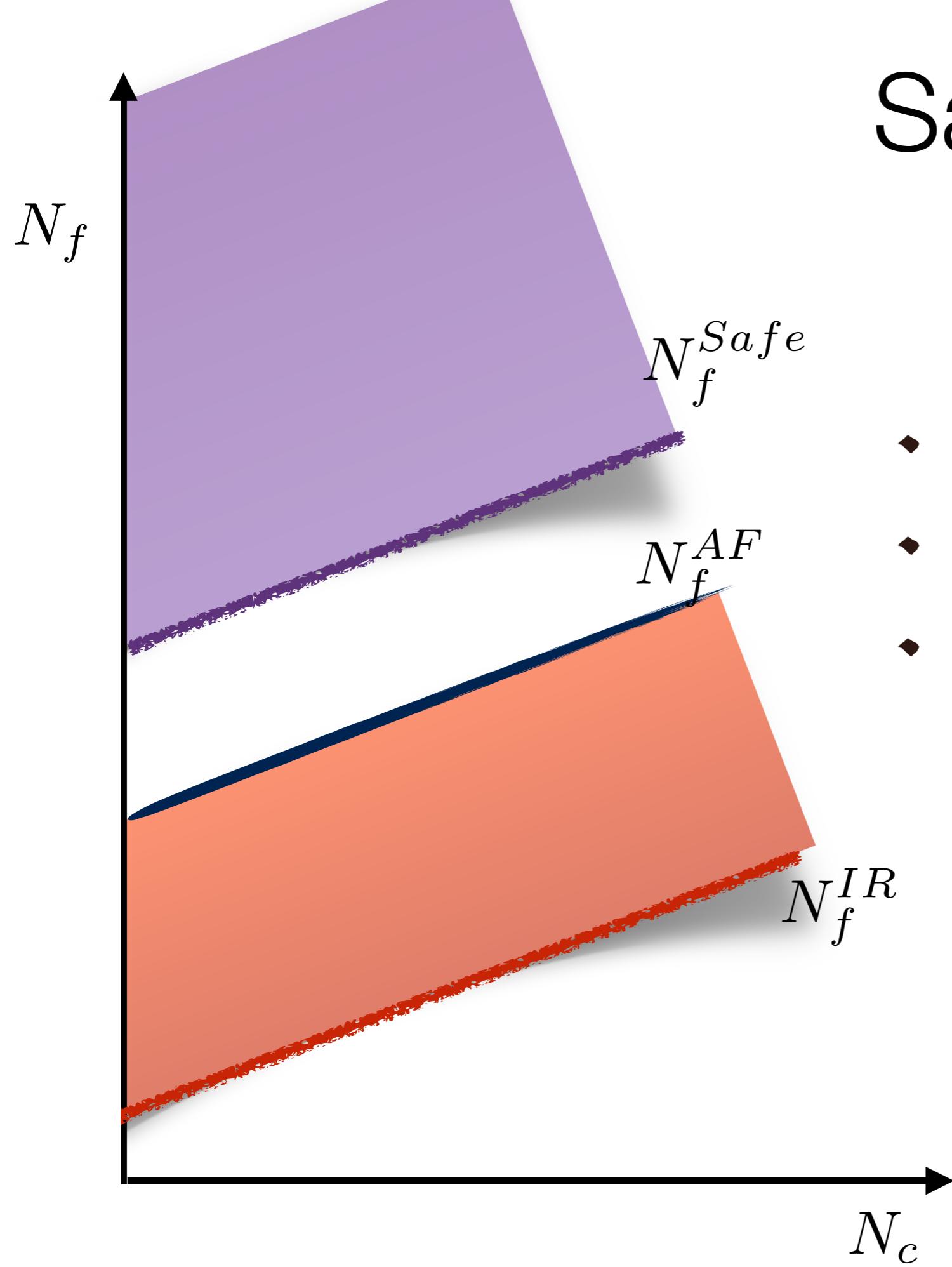


- ◆ Must exist a critical Safe  $N_f$
- ◆ Unsafe region in  $N_f$ - $N_c$

Sannino, ERG 2016, Heidelberg

Antipin and Sannino, to appear  
Pica and Sannino 1011.5917, PRD

# Safe QCD



- ◆ Must exist a critical Safe  $N_f$
- ◆ Unsafe region in  $N_f$ - $N_c$
- ◆ Continuous (Walking) transition?

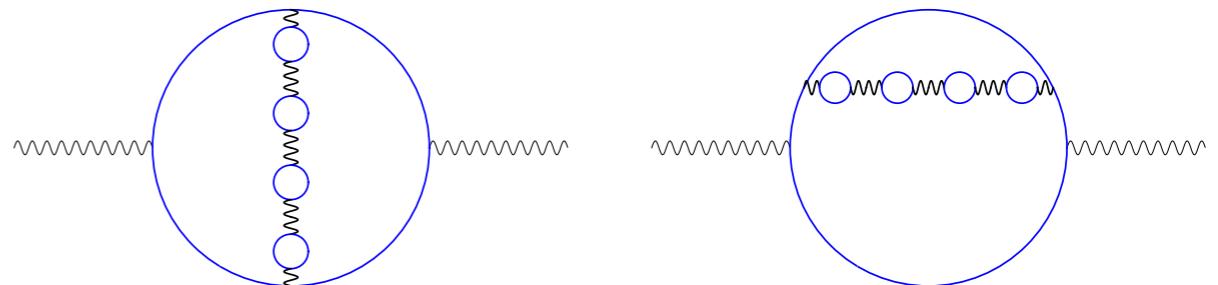
Sannino, ERG 2016, Heidelberg

Antipin and Sannino, to appear  
Pica and Sannino 1011.5917, PRD

# Large Nf (Standard Model)

# Large Nf (Standard Model)

$$\frac{\partial \alpha_i}{\partial \ln \mu} = \beta_{\alpha_i} = \beta_{\alpha_i}^{\text{SM}} + \beta_{\alpha_i}^{\text{extra}}$$

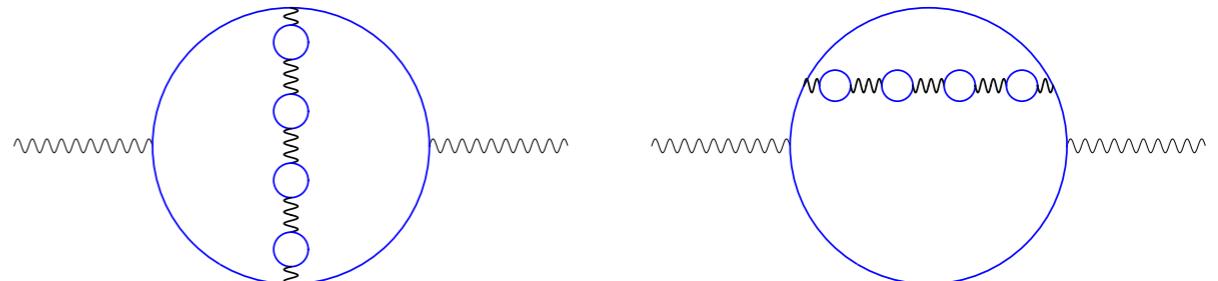


Antipin and Sannino, to appear  
Palanques-Mestre, Pascual, Commun. Math. Phys. 84  
Gracey, PLB, 96, Holdom PLB 2011  
Pica and Sannino 1011.5917, PRD

# Large Nf (Standard Model)

$$\frac{\partial \alpha_i}{\partial \ln \mu} = \beta_{\alpha_i} = \beta_{\alpha_i}^{\text{SM}} + \beta_{\alpha_i}^{\text{extra}}$$

$$\beta_{\alpha_i}^{\text{extra}} = \frac{\alpha_i^2}{2\pi} \Delta b_i + \frac{\alpha_i^2}{3\pi} F_i(\Delta b_i \frac{\alpha_i}{4\pi})$$



Antipin and Sannino, to appear  
Palanques-Mestre, Pascual, Commun. Math. Phys. 84  
Gracey, PLB, 96, Holdom PLB 2011  
Pica and Sannino 1011.5917, PRD

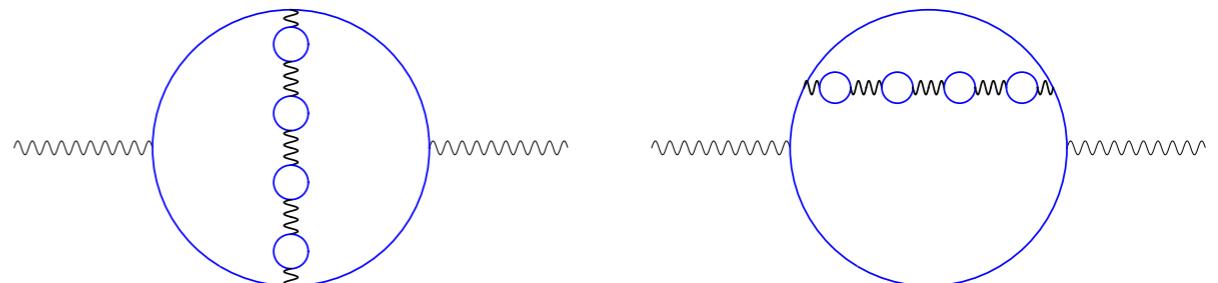
# Large Nf (Standard Model)

$$\frac{\partial \alpha_i}{\partial \ln \mu} = \beta_{\alpha_i} = \beta_{\alpha_i}^{\text{SM}} + \beta_{\alpha_i}^{\text{extra}}$$

$$\beta_{\alpha_i}^{\text{extra}} = \frac{\alpha_i^2}{2\pi} \Delta b_i + \frac{\alpha_i^2}{3\pi} F_i(\Delta b_i \frac{\alpha_i}{4\pi})$$

$$\Delta b_Y = \frac{4}{3} Y^2 N_F D_{R_2} D_{R_3} \quad \Delta b_2 = \frac{2}{3} N_F D_{R_3}$$

$$\Delta b_3 = \frac{2}{3} N_F D_{R_2}$$



Antipin and Sannino, to appear  
Palanques-Mestre, Pascual, Commun. Math. Phys. 84  
Gracey, PLB, 96, Holdom PLB 2011  
Pica and Sannino 1011.5917, PRD

# Large Nf (Standard Model)

$$\frac{\partial \alpha_i}{\partial \ln \mu} = \beta_{\alpha_i} = \beta_{\alpha_i}^{\text{SM}} + \beta_{\alpha_i}^{\text{extra}}$$

$$\beta_{\alpha_i}^{\text{extra}} = \frac{\alpha_i^2}{2\pi} \Delta b_i + \frac{\alpha_i^2}{3\pi} F_i(\Delta b_i \frac{\alpha_i}{4\pi})$$

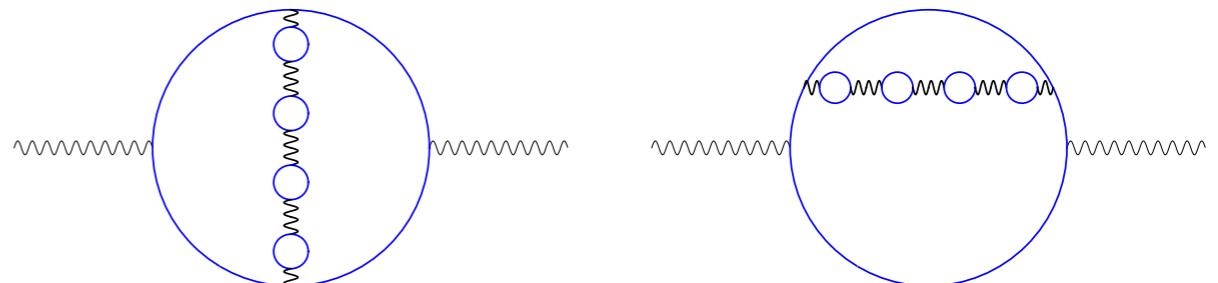
$$\Delta b_Y = \frac{4}{3} Y^2 N_F D_{R_2} D_{R_3} \quad \Delta b_2 = \frac{2}{3} N_F D_{R_3}$$

$$\Delta b_3 = \frac{2}{3} N_F D_{R_2}$$

$$A \equiv \Delta b \frac{\alpha}{4\pi}$$

$$F_1(A) \equiv 2 \int_0^A I_1(x) dx$$

$$F_n(A) \equiv \int_0^A I_1(x) I_n(x) \quad \text{for } n = 2, 3$$



Antipin and Sannino, to appear  
 Palanques-Mestre, Pascual, Commun. Math. Phys. 84  
 Gracey, PLB, 96, Holdom PLB 2011  
 Pica and Sannino 1011.5917, PRD

# Large Nf (Standard Model)

$$\frac{\partial \alpha_i}{\partial \ln \mu} = \beta_{\alpha_i} = \beta_{\alpha_i}^{\text{SM}} + \beta_{\alpha_i}^{\text{extra}}$$

$$\beta_{\alpha_i}^{\text{extra}} = \frac{\alpha_i^2}{2\pi} \Delta b_i + \frac{\alpha_i^2}{3\pi} F_i(\Delta b_i \frac{\alpha_i}{4\pi})$$

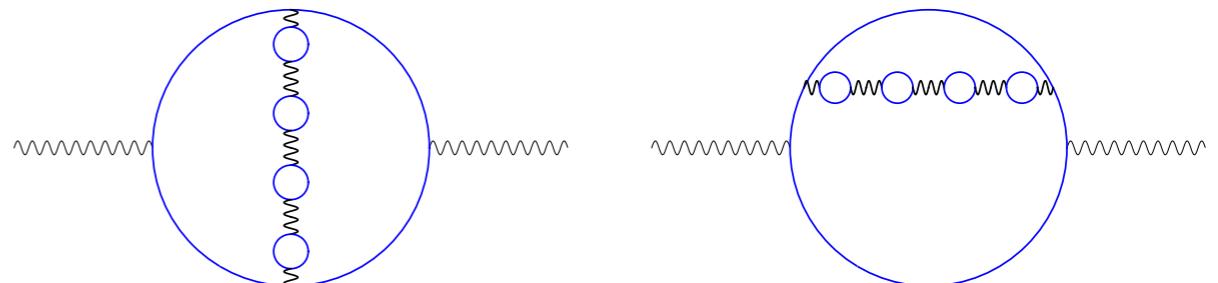
$$\Delta b_Y = \frac{4}{3} Y^2 N_F D_{R_2} D_{R_3} \quad \Delta b_2 = \frac{2}{3} N_F D_{R_3}$$

$$\Delta b_3 = \frac{2}{3} N_F D_{R_2}$$

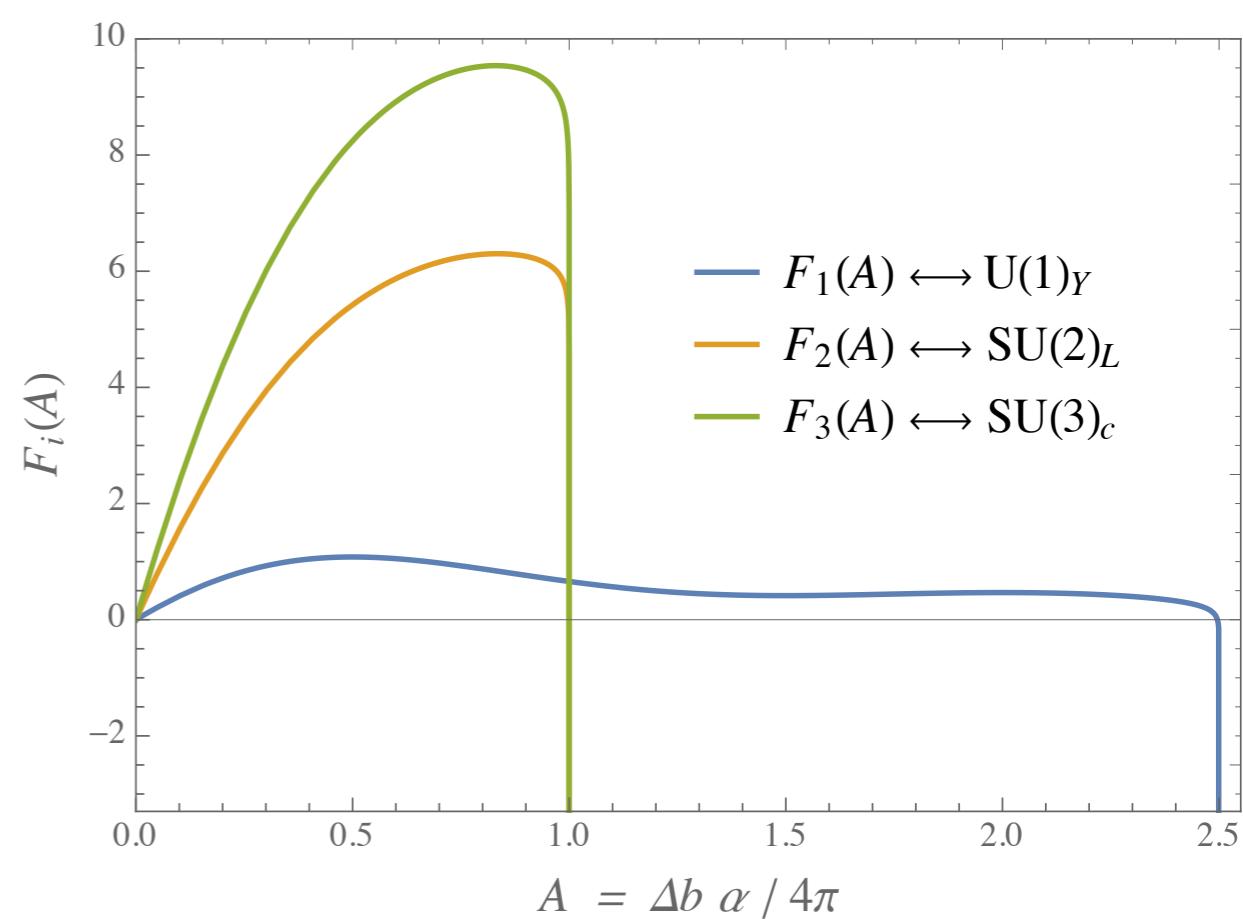
$$A \equiv \Delta b \frac{\alpha}{4\pi}$$

$$F_1(A) \equiv 2 \int_0^A I_1(x) dx$$

$$F_n(A) \equiv \int_0^A I_1(x) I_n(x) \quad \text{for } n = 2, 3$$



Antipin and Sannino, to appear  
 Palanques-Mestre, Pascual, Commun. Math. Phys. 84  
 Gracey, PLB, 96, Holdom PLB 2011  
 Pica and Sannino 1011.5917, PRD



# Large Nf (Standard Model)

$$\frac{\partial \alpha_i}{\partial \ln \mu} = \beta_{\alpha_i} = \beta_{\alpha_i}^{\text{SM}} + \beta_{\alpha_i}^{\text{extra}}$$

$$\beta_{\alpha_i}^{\text{extra}} = \frac{\alpha_i^2}{2\pi} \Delta b_i + \frac{\alpha_i^2}{3\pi} F_i(\Delta b_i \frac{\alpha_i}{4\pi})$$

$$\Delta b_Y = \frac{4}{3} Y^2 N_F D_{R_2} D_{R_3} \quad \Delta b_2 = \frac{2}{3} N_F D_{R_3}$$

$$\Delta b_3 = \frac{2}{3} N_F D_{R_2}$$

$$A \equiv \Delta b \frac{\alpha}{4\pi}$$

$$F_1(A) \equiv 2 \int_0^A I_1(x) dx$$

$$F_n(A) \equiv \int_0^A I_1(x) I_n(x) \quad \text{for } n = 2, 3$$

Abel, Sannino

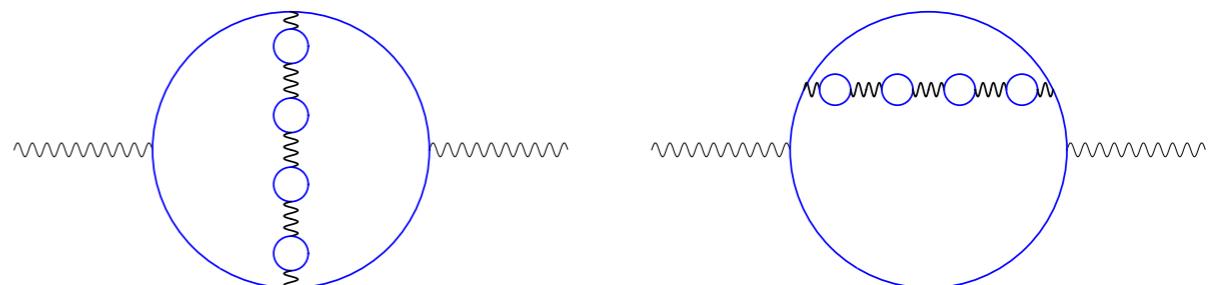
1707.06638

Pelaggi, Plascencia, Salvio, Sannino, Smirnov, Strumia

1708.00437

Mann, Meffe, Wang, Sannino, Steele, Zhang

1707.02942

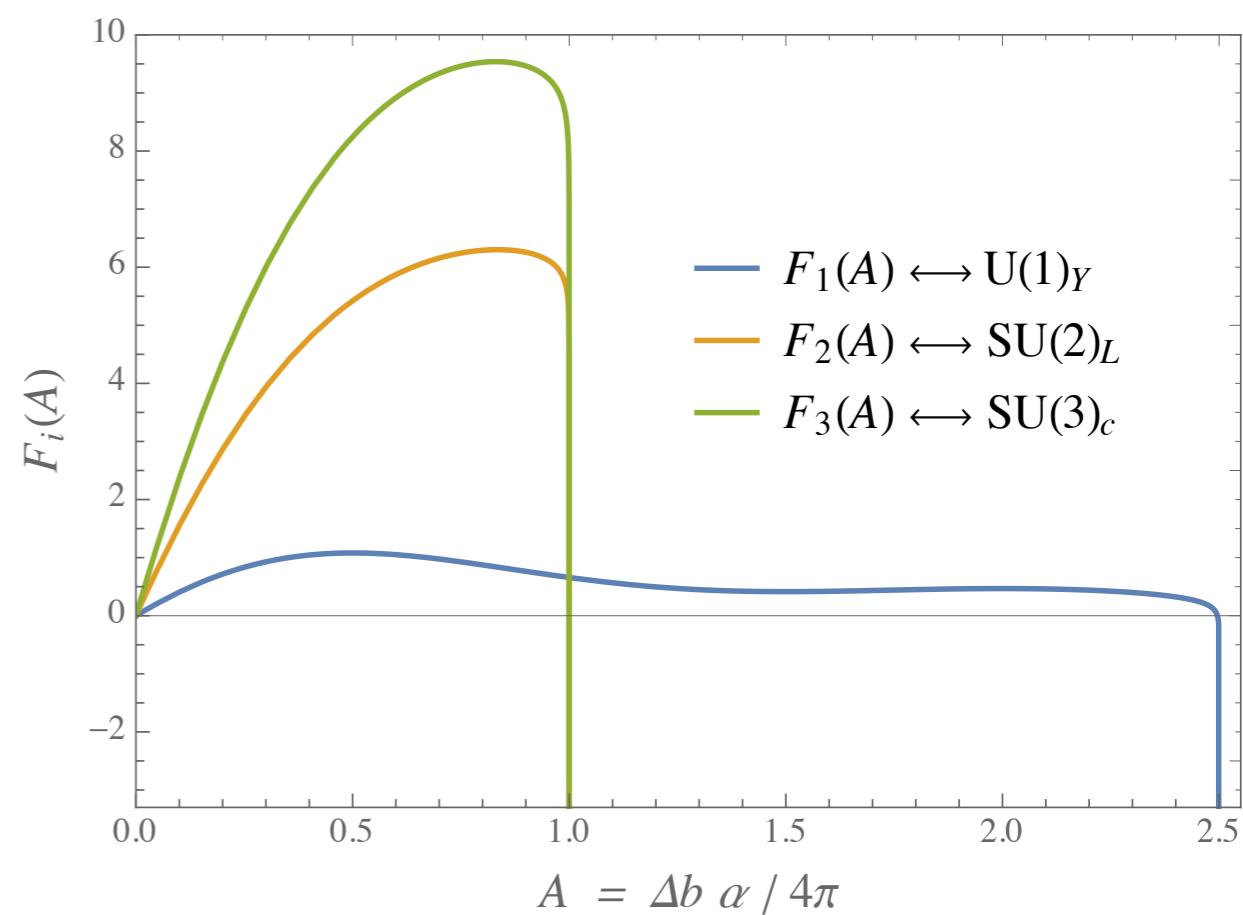


Antipin and Sannino, to appear

Palanques-Mestre, Pascual, Commun. Math. Phys. 84

Gracey, PLB, 96, Holdom PLB 2011

Pica and Sannino 1011.5917, PRD



# Paths to a Safe Standard Model

Large  $N_f$  resummation via vector-like fermions  
[Mann et al. 1707.02942, Pelaggi et al. 1708.00437]

(Perturbative) safety via dynamical breaking,  
[Abel, Sannino 1707.06638, Bond et al. 1702.01727 ]

New paths not yet explored

# Large Nf Safe Standard Model

QCD can be either free [1708.00437] or safe [1707.02942]

Remaining couplings can be treated perturbatively [1707.06638]

# Large Nf Safe Standard Model

QCD can be either free [1708.00437] or safe [1707.02942]

Remaining couplings can be treated perturbatively [1707.06638]

$$\beta_{\lambda_H}^{(1)}|_{\text{SM}} = (4\pi)^2 \frac{d\lambda_H}{d \ln \bar{\mu}} = 24\lambda_H^2 + \lambda_H (12y_t^2 - 9g_2^2 - 3g_Y^2) + \frac{9g_2^4}{8} + \frac{3g_Y^4}{8} + \frac{3g_2^2 g_Y^2}{4} - 6y_t^4$$

# Large Nf Safe Standard Model

QCD can be either free [1708.00437] or safe [1707.02942]

Remaining couplings can be treated perturbatively [1707.06638]

$$\beta_{\lambda_H}^{(1)}|_{\text{SM}} = (4\pi)^2 \frac{d\lambda_H}{d \ln \bar{\mu}} = 24\lambda_H^2 + \lambda_H (12y_t^2 - 9g_2^2 - 3g_Y^2) + \frac{9g_2^4}{8} + \frac{3g_Y^4}{8} + \frac{3g_2^2 g_Y^2}{4} - 6y_t^4$$

Example(s) from 1707.02942

# Large Nf Safe Standard Model

QCD can be either free [1708.00437] or safe [1707.02942]

Remaining couplings can be treated perturbatively [1707.06638]

$$\beta_{\lambda_H}^{(1)}|_{\text{SM}} = (4\pi)^2 \frac{d\lambda_H}{d \ln \bar{\mu}} = 24\lambda_H^2 + \lambda_H (12y_t^2 - 9g_2^2 - 3g_Y^2) + \frac{9g_2^4}{8} + \frac{3g_Y^4}{8} + \frac{3g_2^2 g_Y^2}{4} - 6y_t^4$$

Example(s) from 1707.02942

Add vector-like leptons & RH neutrinos

# Large Nf Safe Standard Model

QCD can be either free [1708.00437] or safe [1707.02942]

Remaining couplings can be treated perturbatively [1707.06638]

$$\beta_{\lambda_H}^{(1)}|_{\text{SM}} = (4\pi)^2 \frac{d\lambda_H}{d \ln \bar{\mu}} = 24\lambda_H^2 + \lambda_H (12y_t^2 - 9g_2^2 - 3g_Y^2) + \frac{9g_2^4}{8} + \frac{3g_Y^4}{8} + \frac{3g_2^2 g_Y^2}{4} - 6y_t^4$$

Example(s) from 1707.02942

Add vector-like leptons & RH neutrinos

$$\Delta b_2 = \Delta b_Y = \frac{2}{3}N_F, \quad \Delta b_3 = 0$$

# Large Nf Safe Standard Model

QCD can be either free [1708.00437] or safe [1707.02942]

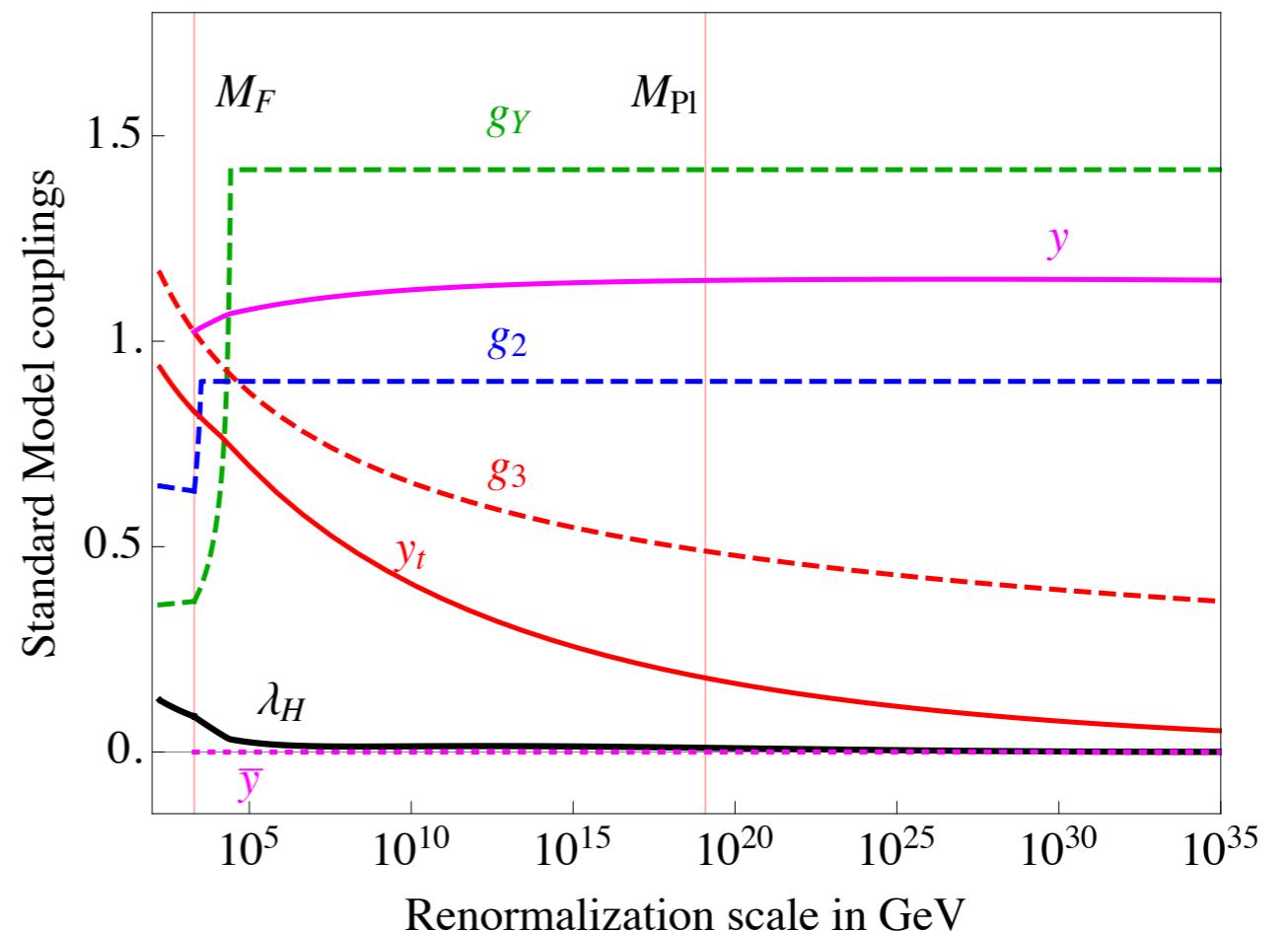
Remaining couplings can be treated perturbatively [1707.06638]

$$\beta_{\lambda_H}^{(1)}|_{\text{SM}} = (4\pi)^2 \frac{d\lambda_H}{d \ln \bar{\mu}} = 24\lambda_H^2 + \lambda_H (12y_t^2 - 9g_2^2 - 3g_Y^2) + \frac{9g_2^4}{8} + \frac{3g_Y^4}{8} + \frac{3g_2^2 g_Y^2}{4} - 6y_t^4$$

Example(s) from 1707.02942

Add vector-like leptons & RH neutrinos

$$\Delta b_2 = \Delta b_Y = \frac{2}{3}N_F, \quad \Delta b_3 = 0$$



# Large Nf Safe Standard Model

QCD can be either free [1708.00437] or safe [1707.02942]

Remaining couplings can be treated perturbatively [1707.06638]

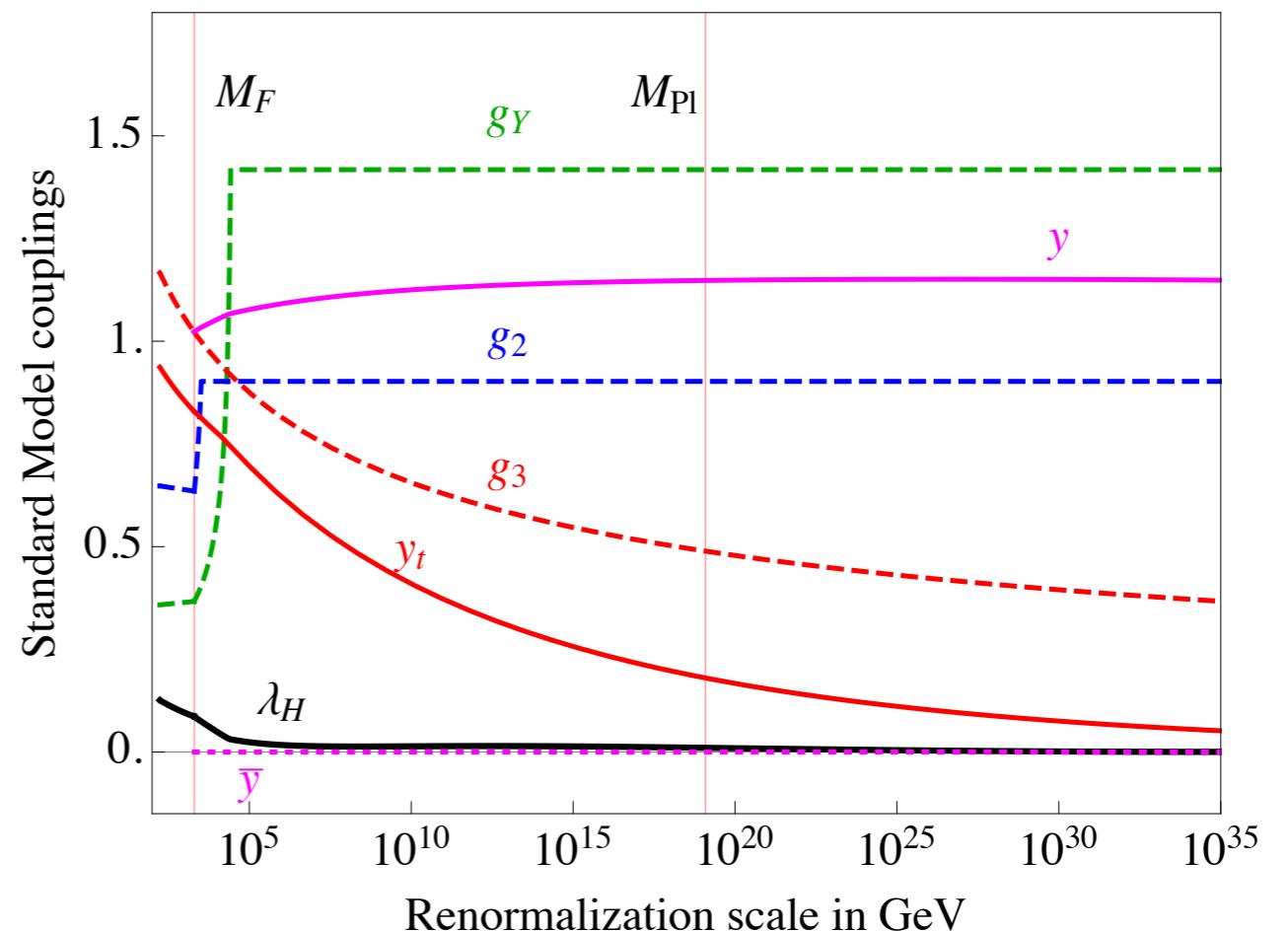
$$\beta_{\lambda_H}^{(1)}|_{\text{SM}} = (4\pi)^2 \frac{d\lambda_H}{d \ln \bar{\mu}} = 24\lambda_H^2 + \lambda_H (12y_t^2 - 9g_2^2 - 3g_Y^2) + \frac{9g_2^4}{8} + \frac{3g_Y^4}{8} + \frac{3g_2^2 g_Y^2}{4} - 6y_t^4$$

Example(s) from 1707.02942

Add vector-like leptons & RH neutrinos

$$\Delta b_2 = \Delta b_Y = \frac{2}{3}N_F, \quad \Delta b_3 = 0$$

Higgs fine tuning can be tamed



# Large Nf Safe Standard Model

QCD can be either free [1708.00437] or safe [1707.02942]

Remaining couplings can be treated perturbatively [1707.06638]

$$\beta_{\lambda_H}^{(1)}|_{\text{SM}} = (4\pi)^2 \frac{d\lambda_H}{d \ln \bar{\mu}} = 24\lambda_H^2 + \lambda_H (12y_t^2 - 9g_2^2 - 3g_Y^2) + \frac{9g_2^4}{8} + \frac{3g_Y^4}{8} + \frac{3g_2^2 g_Y^2}{4} - 6y_t^4$$

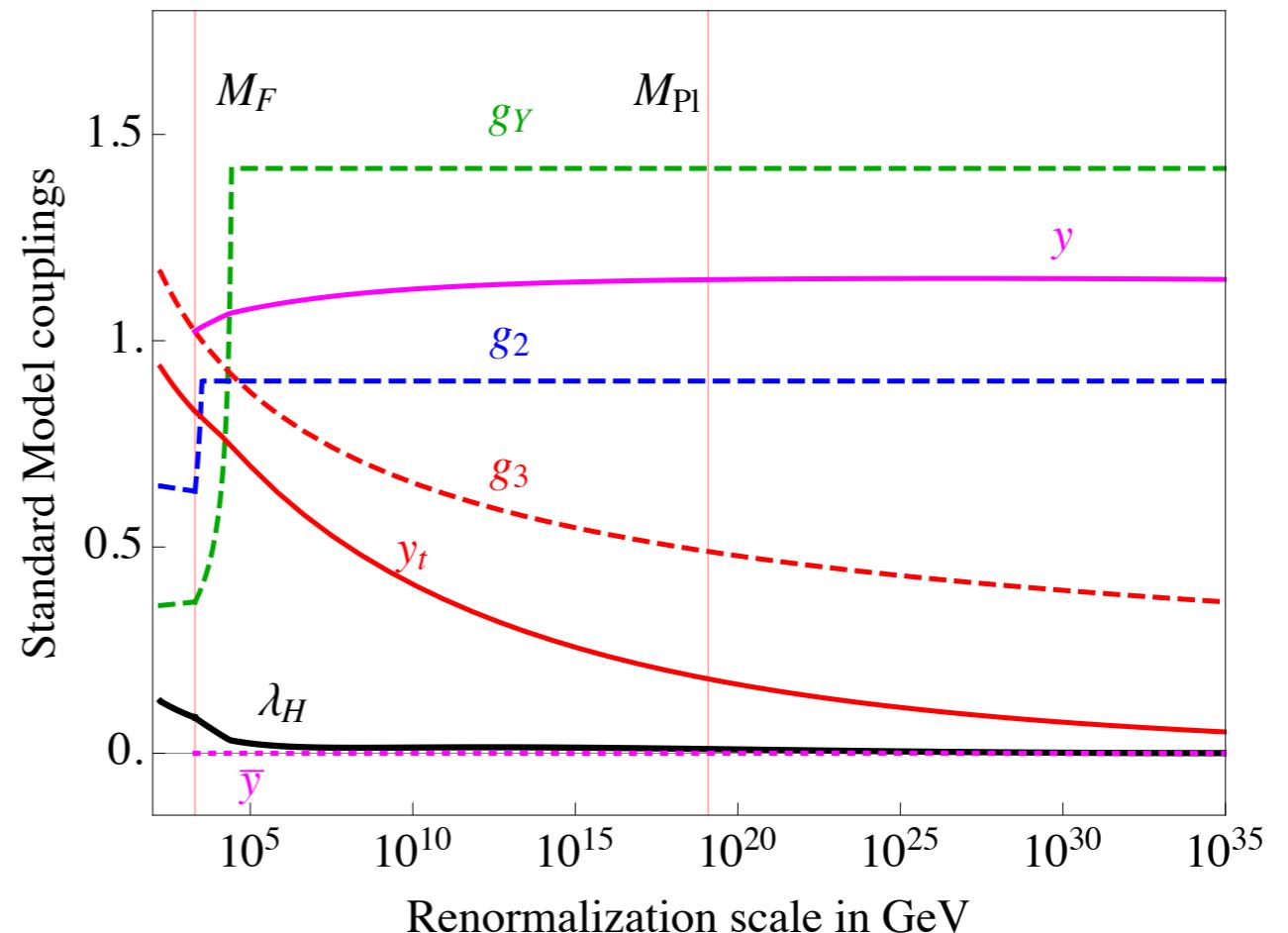
Example(s) from 1707.02942

Add vector-like leptons & RH neutrinos

$$\Delta b_2 = \Delta b_Y = \frac{2}{3}N_F, \quad \Delta b_3 = 0$$

Higgs fine tuning can be tamed

Large theory space at our disposal



# Supersymmetric (un)safety

Intriligator and Sannino, 1508.07413, JHEP

Bajc and Sannino, 1610.09681, JHEP

# Supersymmetric (un)safety

Intriligator and Sannino, 1508.07413, JHEP

Bajc and Sannino, 1610.09681, JHEP

*Exact results beyond perturbation theory*

# a-theorem

# a-theorem

For any super CFT besides positivity we have, following Cardy

# a-theorem

For any super CFT besides positivity we have, following Cardy

$$\Delta a \equiv a_{UV} - a_{IR} > 0$$

# a-theorem

For any super CFT besides positivity we have, following Cardy

$$\Delta a \equiv a_{UV} - a_{IR} > 0$$

$$\Delta a = a_{UV} - a_{IR} = \pm \frac{1}{9} \sum_i |r_i| [(3R_i - 2)^2(3R_i - 5)] > 0$$

# a-theorem

For any super CFT besides positivity we have, following Cardy

$$\Delta a \equiv a_{UV} - a_{IR} > 0$$

$$\Delta a = a_{UV} - a_{IR} = \pm \frac{1}{9} \sum_i |r_i| [(3R_i - 2)^2(3R_i - 5)] > 0$$

$r_i$  = dim. of matter rep.

# a-theorem

For any super CFT besides positivity we have, following Cardy

$$\Delta a \equiv a_{UV} - a_{IR} > 0$$

$$\Delta a = a_{UV} - a_{IR} = \pm \frac{1}{9} \sum_i |r_i| [(3R_i - 2)^2(3R_i - 5)] > 0$$

$r_i$  = dim. of matter rep.

+(-) for asymptotic safety (freedom)

# a-theorem

For any super CFT besides positivity we have, following Cardy

$$\Delta a \equiv a_{UV} - a_{IR} > 0$$

$$\Delta a = a_{UV} - a_{IR} = \pm \frac{1}{9} \sum_i |r_i| [(3R_i - 2)^2(3R_i - 5)] > 0$$

$r_i$  = dim. of matter rep.

+(-) for asymptotic safety (freedom)

# a-theorem

For any super CFT besides positivity we have, following Cardy

$$\Delta a \equiv a_{UV} - a_{IR} > 0$$

$$\Delta a = a_{UV} - a_{IR} = \pm \frac{1}{9} \sum_i |r_i| [(3R_i - 2)^2(3R_i - 5)] > 0$$

$r_i$  = dim. of matter rep.

+(-) for asymptotic safety (freedom)

Stronger constraint for asymp. safety, since at least one large  $R > 5/3$

# SQCD with H

Fields	$[SU(N_c)]$	$SU_L(N_f)$	$SU_R(N_f)$	$U_V(1)$	$U(1)_R$
$W_\alpha$	Adj	1	1	0	1
$Q$	$\square$	$\bar{\square}$	1	1	$1 - \frac{N_c}{N_f}$
$\tilde{Q}$	$\bar{\square}$	1	$\square$	-1	$1 - \frac{N_c}{N_f}$
$H$	1	$\square$	$\bar{\square}$	0	$2\frac{N_c}{N_f}$

# SQCD with $\mathbb{H}$

Fields	$[SU(N_c)]$	$SU_L(N_f)$	$SU_R(N_f)$	$U_V(1)$	$U(1)_R$	AF is lost
$W_\alpha$	Adj	1	1	0	1	
$Q$	$\square$	$\bar{\square}$	1	1	$1 - \frac{N_c}{N_f}$	$N_f > 3N_c$
$\tilde{Q}$	$\bar{\square}$	1	$\square$	-1	$1 - \frac{N_c}{N_f}$	
$H$	1	$\square$	$\bar{\square}$	0	$2\frac{N_c}{N_f}$	

# SQCD with H

Fields	$[SU(N_c)]$	$SU_L(N_f)$	$SU_R(N_f)$	$U_V(1)$	$U(1)_R$	AF is lost
$W_\alpha$	Adj	1	1	0	1	
$Q$	$\square$	$\bar{\square}$	1	1	$1 - \frac{N_c}{N_f}$	$N_f > 3N_c$
$\tilde{Q}$	$\bar{\square}$	1	$\square$	-1	$1 - \frac{N_c}{N_f}$	
$H$	1	$\square$	$\bar{\square}$	0	$2\frac{N_c}{N_f}$	$W = y \operatorname{Tr} Q H \tilde{Q}$

# SQCD with H

Fields	$[SU(N_c)]$	$SU_L(N_f)$	$SU_R(N_f)$	$U_V(1)$	$U(1)_R$
$W_\alpha$	Adj	1	1	0	1
$Q$	□	□	1	1	$1 - \frac{N_c}{N_f}$
$\tilde{Q}$	□	1	□	-1	$1 - \frac{N_c}{N_f}$
$H$	1	□	□	0	$2\frac{N_c}{N_f}$

AF is lost

$$N_f > 3N_c$$

$$W = y \operatorname{Tr} Q H \tilde{Q}$$

$$\beta(\alpha_g) \approx -2\alpha_g^2 \left[ 3 - \frac{N_f}{N_c} + \left( 6 - 4\frac{N_f}{N_c} \right) \alpha_g + 2\frac{N_f^2}{N_c^2} \alpha_y + \mathcal{O}(\alpha^2) \right]$$

$$\beta(\alpha_y) \approx 2\alpha_y \left[ \left( 2\frac{N_f}{N_c} + 1 \right) \alpha_y - 2\alpha_g + \mathcal{O}(\alpha^2) \right]$$

# SQCD with H

Fields	$[SU(N_c)]$	$SU_L(N_f)$	$SU_R(N_f)$	$U_V(1)$	$U(1)_R$
$W_\alpha$	Adj	1	1	0	1
$Q$	□	□	1	1	$1 - \frac{N_c}{N_f}$
$\tilde{Q}$	□	1	□	-1	$1 - \frac{N_c}{N_f}$
$H$	1	□	□	0	$2\frac{N_c}{N_f}$

AF is lost

$$N_f > 3N_c$$

$$W = y \operatorname{Tr} Q H \tilde{Q}$$

$$\beta(\alpha_g) \approx -2\alpha_g^2 \left[ 3 - \frac{N_f}{N_c} + \left( 6 - 4\frac{N_f}{N_c} \right) \alpha_g + 2\frac{N_f^2}{N_c^2} \alpha_y + \mathcal{O}(\alpha^2) \right]$$

$$\beta(\alpha_y) \approx 2\alpha_y \left[ \left( 2\frac{N_f}{N_c} + 1 \right) \alpha_y - 2\alpha_g + \mathcal{O}(\alpha^2) \right]$$

No perturbative UV fixed point

$$\beta(\alpha_g) \approx 2\alpha_g^2 \left[ \epsilon + \frac{6}{7}\alpha_g \right]$$

# SQCD with H

# SQCD with H

Assume a nonperturbative fixed point, however

# SQCD with H

Assume a nonperturbative fixed point, however

$$D(H) = \frac{3}{2}R(H) = 3\frac{N_c}{N_f} < 1 \quad \text{for} \quad N_f > 3N_c$$

# SQCD with H

Assume a nonperturbative fixed point, however

$$D(H) = \frac{3}{2}R(H) = 3\frac{N_c}{N_f} < 1 \quad \text{for} \quad N_f > 3N_c$$

Violates the unitarity bound

$$D(\mathcal{O}) \geq 1$$

# SQCD with H

Assume a nonperturbative fixed point, however

$$D(H) = \frac{3}{2}R(H) = 3\frac{N_c}{N_f} < 1 \quad \text{for} \quad N_f > 3N_c$$

Violates the unitarity bound

$$D(\mathcal{O}) \geq 1$$

Potential loophole: H is free and decouples at the fixed point

# SQCD with H

Assume a nonperturbative fixed point, however

$$D(H) = \frac{3}{2} R(H) = 3 \frac{N_c}{N_f} < 1 \quad \text{for} \quad N_f > 3N_c$$

Violates the unitarity bound

$$D(\mathcal{O}) \geq 1$$

Potential loophole: H is free and decouples at the fixed point

Check if SQCD without H has an UV fixed point

SQCD

# SQCD

Unitarity bound is not sufficient

# SQCD

Unitarity bound is not sufficient

$$\mathcal{M} = Q\tilde{Q} \quad D_{SCFT}(\mathcal{M}) = \frac{3}{2}R_{SCFT}(\mathcal{M}) = 3\frac{N_f - N_c}{N_f}$$

# SQCD

Unitarity bound is not sufficient

$$\mathcal{M} = Q\tilde{Q}$$

$$D_{SCFT}(\mathcal{M}) = \frac{3}{2}R_{SCFT}(\mathcal{M}) = 3\frac{N_f - N_c}{N_f}$$

$$\mathcal{B} = Q^{N_c}$$

$$D_{SCFT}(\mathcal{B}) = D_{SCFT}(\tilde{\mathcal{B}}) = \frac{3}{2}R_{SCFT}(\mathcal{B}) = \frac{3}{2}N_c\frac{N_f - N_c}{N_f}$$

# SQCD

Unitarity bound is not sufficient

$$\mathcal{M} = Q\tilde{Q}$$

$$D_{SCFT}(\mathcal{M}) = \frac{3}{2}R_{SCFT}(\mathcal{M}) = 3\frac{N_f - N_c}{N_f}$$

$$\mathcal{B} = Q^{N_c}$$

$$D_{SCFT}(\mathcal{B}) = D_{SCFT}(\tilde{\mathcal{B}}) = \frac{3}{2}R_{SCFT}(\mathcal{B}) = \frac{3}{2}N_c\frac{N_f - N_c}{N_f}$$

Can be ruled out via a-theorem      $a(R) = 3\text{Tr}U(1)_R^3 - \text{Tr}U(1)_R$

# SQCD

Unitarity bound is not sufficient

$$\mathcal{M} = Q\tilde{Q}$$

$$D_{SCFT}(\mathcal{M}) = \frac{3}{2}R_{SCFT}(\mathcal{M}) = 3\frac{N_f - N_c}{N_f}$$

$$\mathcal{B} = Q^{N_c}$$

$$D_{SCFT}(\mathcal{B}) = D_{SCFT}(\tilde{\mathcal{B}}) = \frac{3}{2}R_{SCFT}(\mathcal{B}) = \frac{3}{2}N_c\frac{N_f - N_c}{N_f}$$

Can be ruled out via a-theorem  $a(R) = 3\text{Tr}U(1)_R^3 - \text{Tr}U(1)_R$

$$a_{\text{UV-safe}} - a_{\text{IR-safe}} < 0$$

# SQCD

Unitarity bound is not sufficient

$$\mathcal{M} = Q\tilde{Q}$$

$$D_{SCFT}(\mathcal{M}) = \frac{3}{2}R_{SCFT}(\mathcal{M}) = 3\frac{N_f - N_c}{N_f}$$

$$\mathcal{B} = Q^{N_c}$$

$$D_{SCFT}(\mathcal{B}) = D_{SCFT}(\tilde{\mathcal{B}}) = \frac{3}{2}R_{SCFT}(\mathcal{B}) = \frac{3}{2}N_c\frac{N_f - N_c}{N_f}$$

Can be ruled out via a-theorem  $a(R) = 3\text{Tr}U(1)_R^3 - \text{Tr}U(1)_R$

$$a_{\text{UV-safe}} - a_{\text{IR-safe}} < 0$$

*Non-abelian SQED with(out) H cannot be asymptotically safe*

# SQCD

Unitarity bound is not sufficient

$$\mathcal{M} = Q\tilde{Q}$$

$$D_{SCFT}(\mathcal{M}) = \frac{3}{2}R_{SCFT}(\mathcal{M}) = 3\frac{N_f - N_c}{N_f}$$

$$\mathcal{B} = Q^{N_c}$$

$$D_{SCFT}(\mathcal{B}) = D_{SCFT}(\tilde{\mathcal{B}}) = \frac{3}{2}R_{SCFT}(\mathcal{B}) = \frac{3}{2}N_c\frac{N_f - N_c}{N_f}$$

Can be ruled out via a-theorem  $a(R) = 3\text{Tr}U(1)_R^3 - \text{Tr}U(1)_R$

$$a_{\text{UV-safe}} - a_{\text{IR-safe}} < 0$$

*Non-abelian SQED with(out) H cannot be asymptotically safe*

*Generalisation to several susy theories using a-maximisation\**

# Unsafe SUSY GUTs

Bajc and Sannino, 1610.09681, JHEP

# Unsafe SUSY GUTs

Bajc and Sannino, 1610.09681, JHEP

*Exact results*

Gaining R parity... but

# Gaining R parity... but

R-symmetry from SO(10) Cartan subalgebra generator B-L

# Gaining R parity... but

R-symmetry from SO(10) Cartan subalgebra generator B-L

$$R = (-1)^{3(B-L)+2S} = M(-1)^{2S} \quad \text{with} \quad M = (-1)^{3(B-L)}$$

# Gaining R parity... but

R-symmetry from SO(10) Cartan subalgebra generator B-L

$$R = (-1)^{3(B-L)+2S} = M(-1)^{2S} \quad \text{with} \quad M = (-1)^{3(B-L)}$$

M = matter parity

# Gaining R parity... but

R-symmetry from SO(10) Cartan subalgebra generator B-L

$$R = (-1)^{3(B-L)+2S} = M(-1)^{2S} \quad \text{with} \quad M = (-1)^{3(B-L)}$$

M = matter parity

Elegant breaking of SO(10) preserving R-parity:

# Gaining R parity... but

R-symmetry from SO(10) Cartan subalgebra generator B-L

$$R = (-1)^{3(B-L)+2S} = M(-1)^{2S} \quad \text{with} \quad M = (-1)^{3(B-L)}$$

M = matter parity

Elegant breaking of SO(10) preserving R-parity:

Introduce  $126 + 126^*$  Higgs in SO(10)

# Gaining R parity... but

R-symmetry from SO(10) Cartan subalgebra generator B-L

$$R = (-1)^{3(B-L)+2S} = M(-1)^{2S} \quad \text{with} \quad M = (-1)^{3(B-L)}$$

M = matter parity

Elegant breaking of SO(10) preserving R-parity:

Introduce  $126 + 126^*$  Higgs in SO(10)

$126(126^*)$  SM and SU(5) singlet has  $B-L=-2(2)$  preserving R-parity

asympt. freedom is lost

asympt. freedom is lost

$$W_{Yukawa} = 16_a \left( Y_{10}^{ab} 10 + Y_{126}^{ab} \overline{126} + Y_{120}^{ab} 120 \right) 16_b$$

asympt. freedom is lost

$$W_{Yukawa} = 16_a \left( Y_{10}^{ab} 10 + Y_{126}^{ab} \overline{126} + Y_{120}^{ab} 120 \right) 16_b$$

a, b run over generations

# asympt. freedom is lost

$$W_{Yukawa} = 16_a \left( Y_{10}^{ab} 10 + Y_{126}^{ab} \overline{126} + Y_{120}^{ab} 120 \right) 16_b$$

a, b run over generations

To fully break SO(10) to SM add 210 of SO(10)

# asympt. freedom is lost

$$W_{Yukawa} = 16_a \left( Y_{10}^{ab} 10 + Y_{126}^{ab} \overline{126} + Y_{120}^{ab} 120 \right) 16_b$$

a, b run over generations

To fully break SO(10) to SM add 210 of SO(10)

In summary:  $3 \times 16 + 126 + 126^* + 10 + 210$  contributes

# asympt. freedom is lost

$$W_{Yukawa} = 16_a \left( Y_{10}^{ab} 10 + Y_{126}^{ab} \overline{126} + Y_{120}^{ab} 120 \right) 16_b$$

a, b run over generations

To fully break SO(10) to SM add 210 of SO(10)

In summary:  $3 \times 16 + 126 + 126^* + 10 + 210$  contributes

$$\beta_{1-loop} = +109$$

# asymptotic freedom is lost

$$W_{Yukawa} = 16_a \left( Y_{10}^{ab} 10 + Y_{126}^{ab} \overline{126} + Y_{120}^{ab} 120 \right) 16_b$$

a, b run over generations

To fully break SO(10) to SM add 210 of SO(10)

In summary:  $3 \times 16 + 126 + 126^* + 10 + 210$  contributes

$$\beta_{1-loop} = +109$$

*Asymptotic freedom is badly lost!*

# Exact results

# Exact results

***Minimal SO(10) without super potential***

# Exact results

***Minimal SO(10) without super potential***

$3 \times 16 + 126 + 126^* + 10 + 210$  ***is unsafe.***

# Exact results

## ***Minimal SO(10) without super potential***

$3 \times 16 + 126 + 126^* + 10 + 210$  ***is unsafe.***

*Exotic examples exist requiring thousands of generations!*

# Exact results

## ***Minimal SO(10) without super potential***

$3 \times 16 + 126 + 126^* + 10 + 210$  ***is unsafe.***

*Exotic examples exist requiring thousands of generations!*

## ***Minimal SO(10) with general 3-linear super potential***

# Exact results

## ***Minimal SO(10) without super potential***

$3 \times 16 + 126 + 126^* + 10 + 210$  ***is unsafe.***

*Exotic examples exist requiring thousands of generations!*

## ***Minimal SO(10) with general 3-linear super potential***

$$W = y_1 210^3 + y_2 210 126 \overline{126} + y_3 210 126 10 + y_4 210 \overline{126} 10 + \sum_{a,b=1,2,3} 16_a 16_b (y_{5,ab} 10 + y_{6,ab} \overline{126})$$

# Exact results

## ***Minimal SO(10) without super potential***

$3 \times 16 + 126 + 126^* + 10 + 210$  ***is unsafe.***

*Exotic examples exist requiring thousands of generations!*

## ***Minimal SO(10) with general 3-linear super potential***

$$W = y_1 210^3 + y_2 210 126 \overline{126} + y_3 210 126 10 + y_4 210 \overline{126} 10 + \sum_{a,b=1,2,3} 16_a 16_b (y_{5,ab} 10 + y_{6,ab} \overline{126})$$

- All trilinear present then:  $R=2/3$  for all fields and no NSVZ UV fixed point
- Eliminate one 16 from super potential passes the constraints

# Exact results

## ***Minimal SO(10) without super potential***

$3 \times 16 + 126 + 126^* + 10 + 210$  ***is unsafe.***

*Exotic examples exist requiring thousands of generations!*

## ***Minimal SO(10) with general 3-linear super potential***

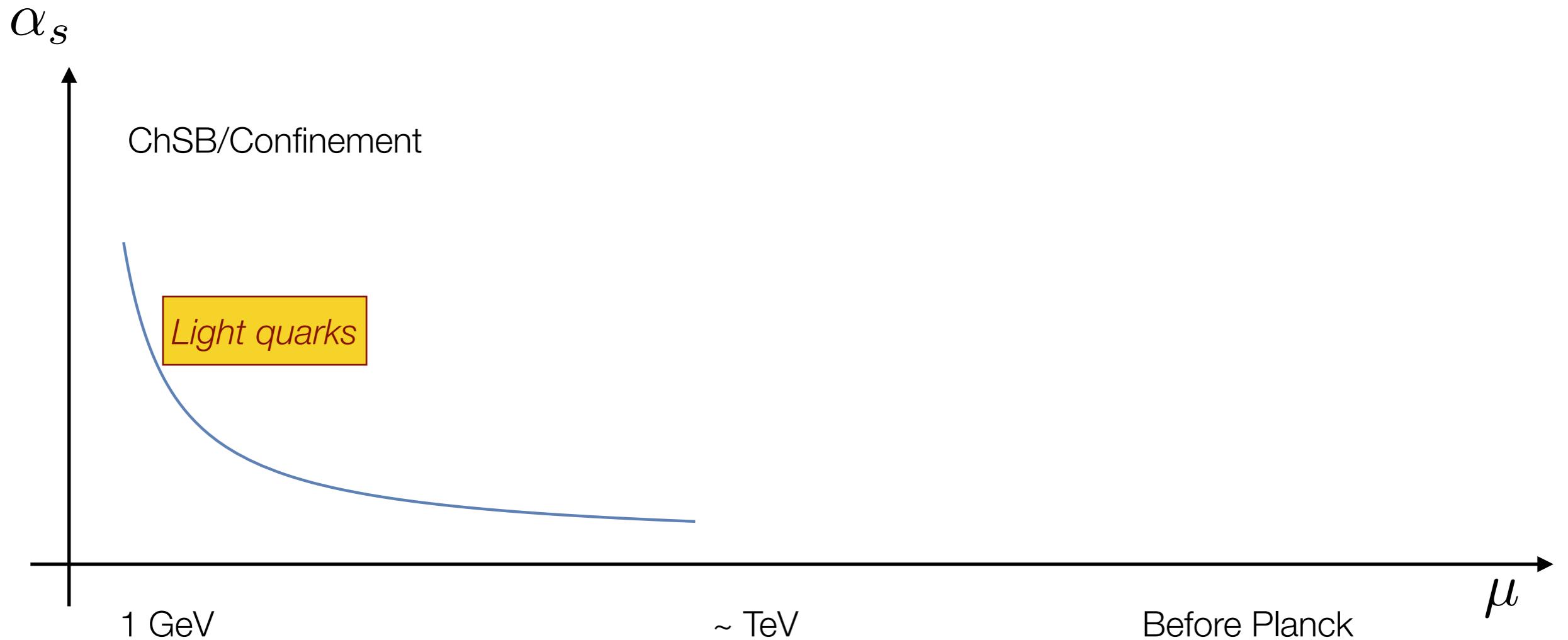
$$W = y_1 210^3 + y_2 210 126 \overline{126} + y_3 210 126 10 + y_4 210 \overline{126} 10 + \sum_{a,b=1,2,3} 16_a 16_b (y_{5,ab} 10 + y_{6,ab} \overline{126})$$

- All trilinear present then:  $R=2/3$  for all fields and no NSVZ UV fixed point
- Eliminate one 16 from super potential passes the constraints

*Super GUTs with R-charge are challenging!*

# Safe QCD

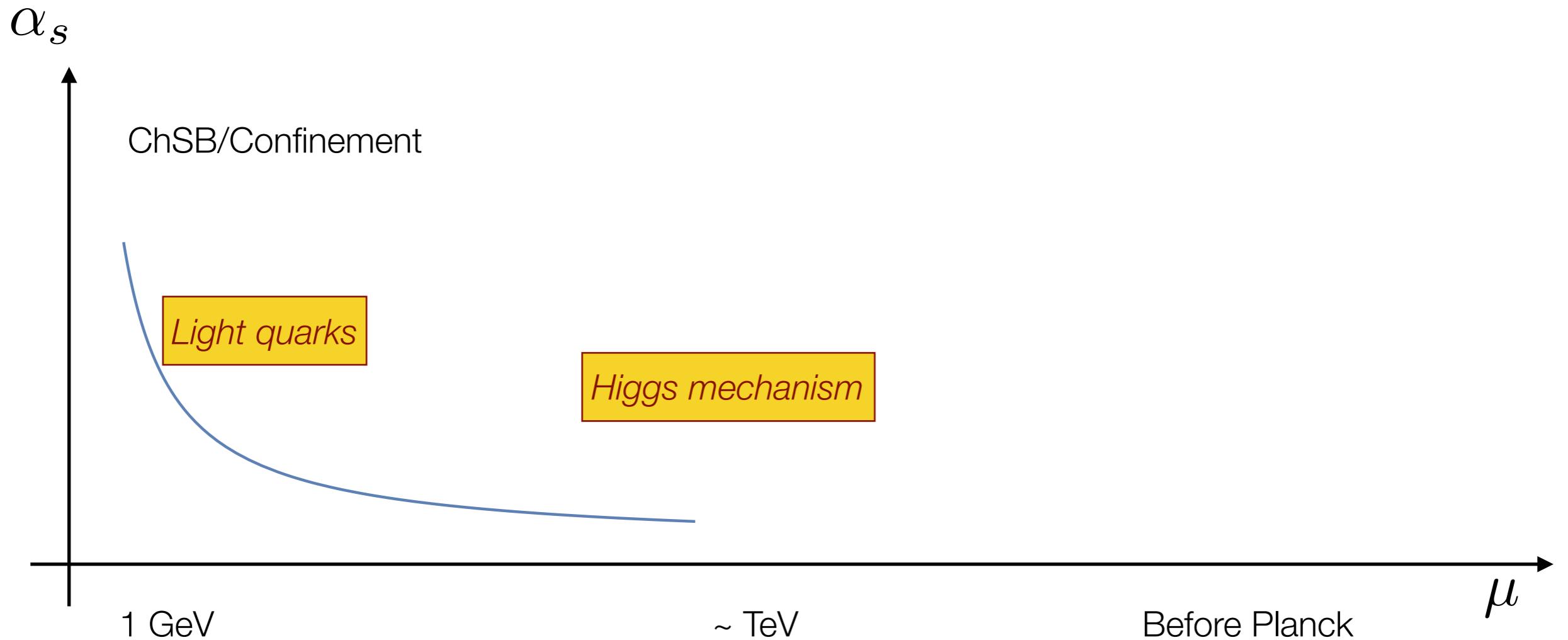
Sannino, 1511.09022



Pica & Sannino, 1011.5917 PRD

# Safe QCD

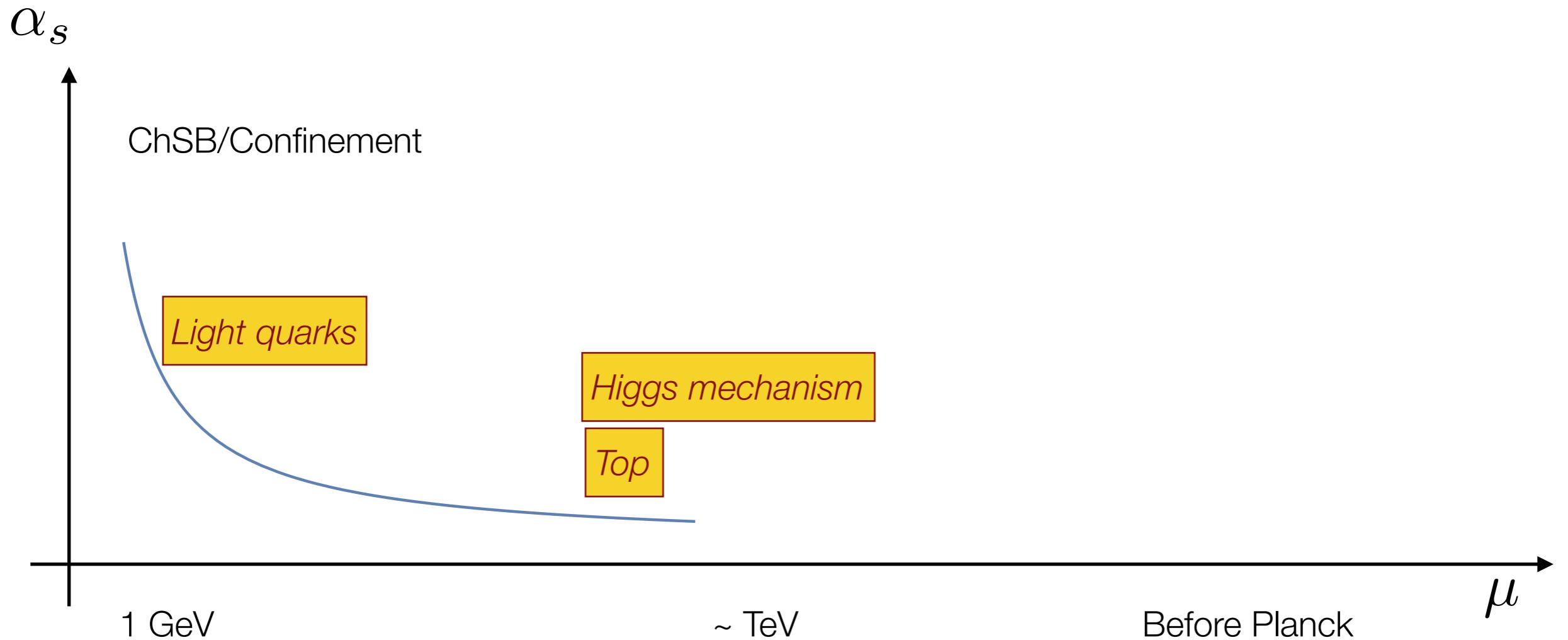
Sannino, 1511.09022



Pica & Sannino, 1011.5917 PRD

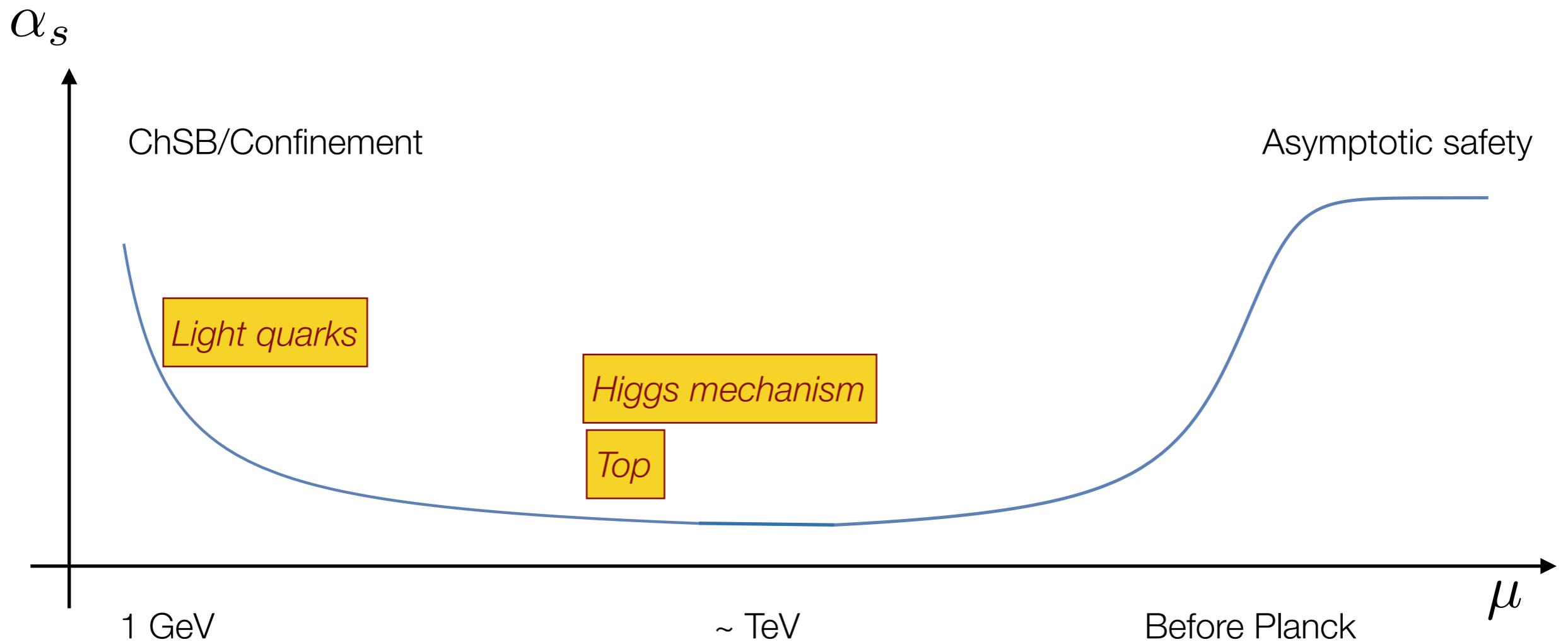
# Safe QCD

Sannino, 1511.09022



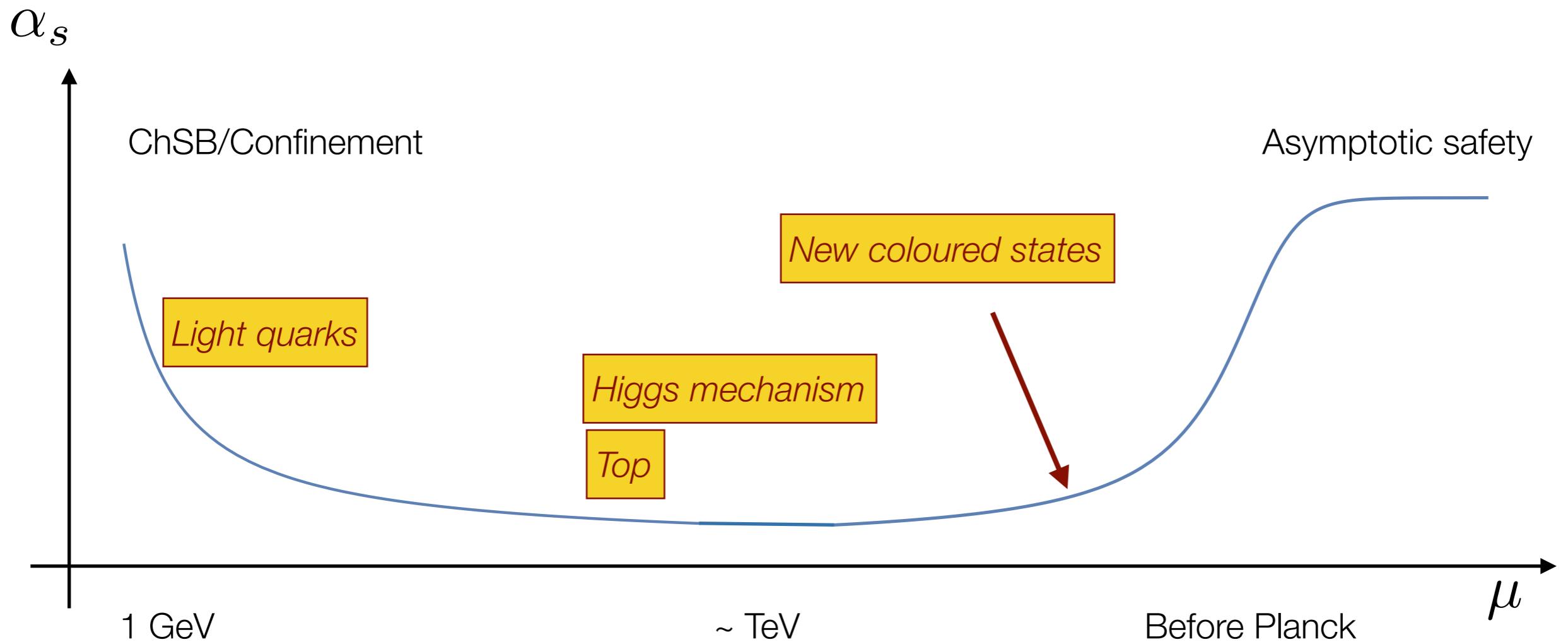
# Safe QCD

Sannino, 1511.09022



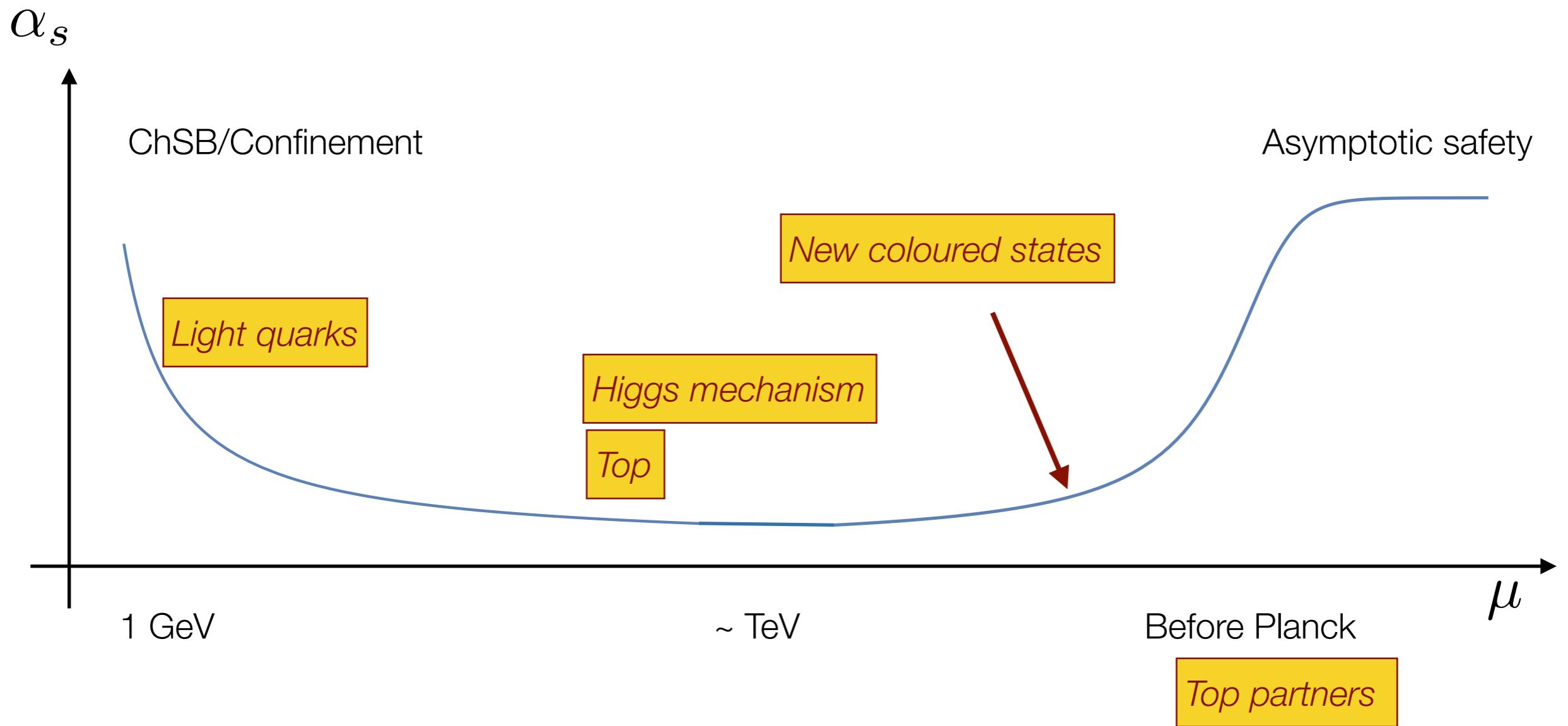
# Safe QCD

Sannino, 1511.09022



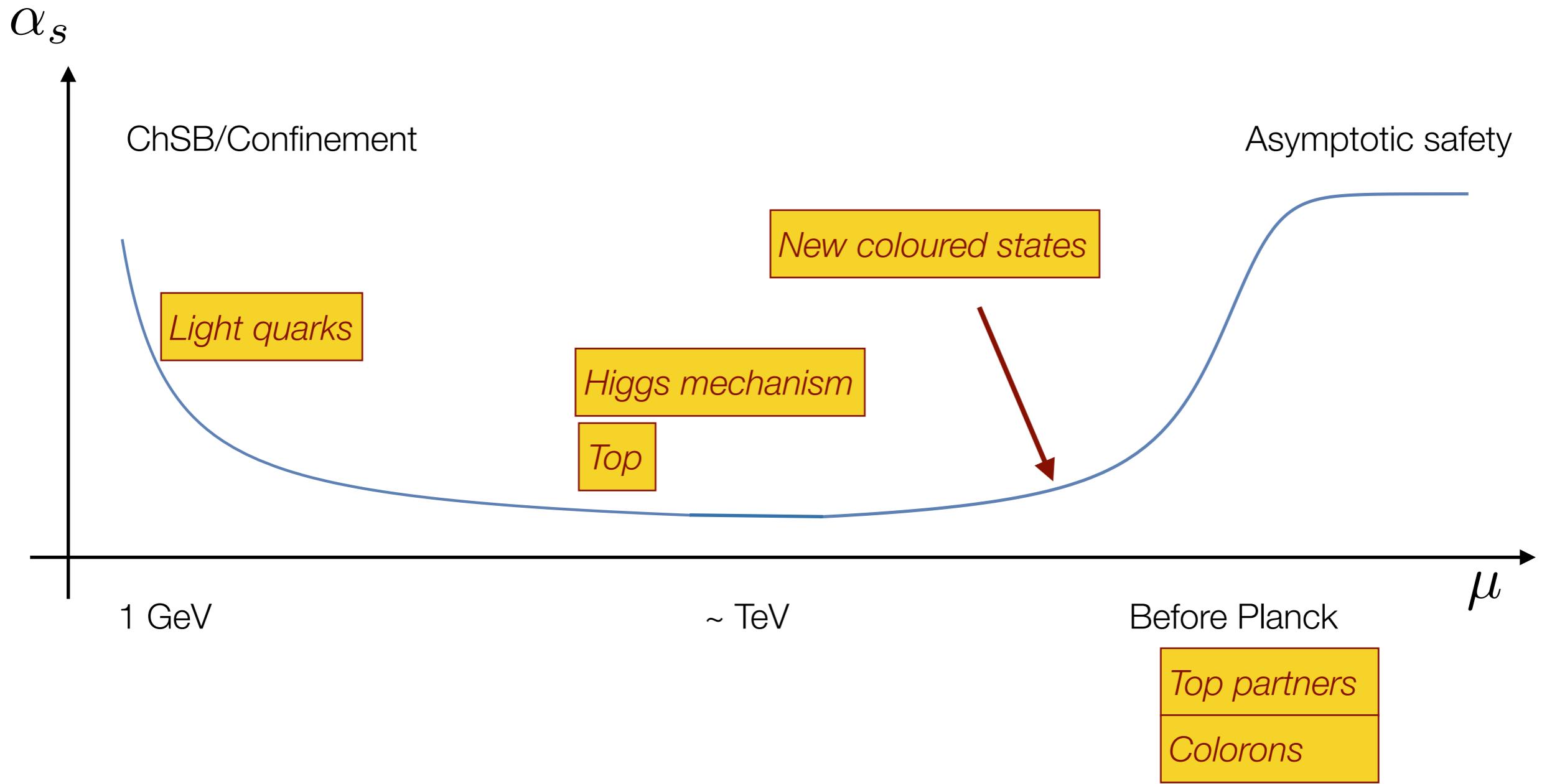
# Safe QCD

Sannino, 1511.09022



# Safe QCD

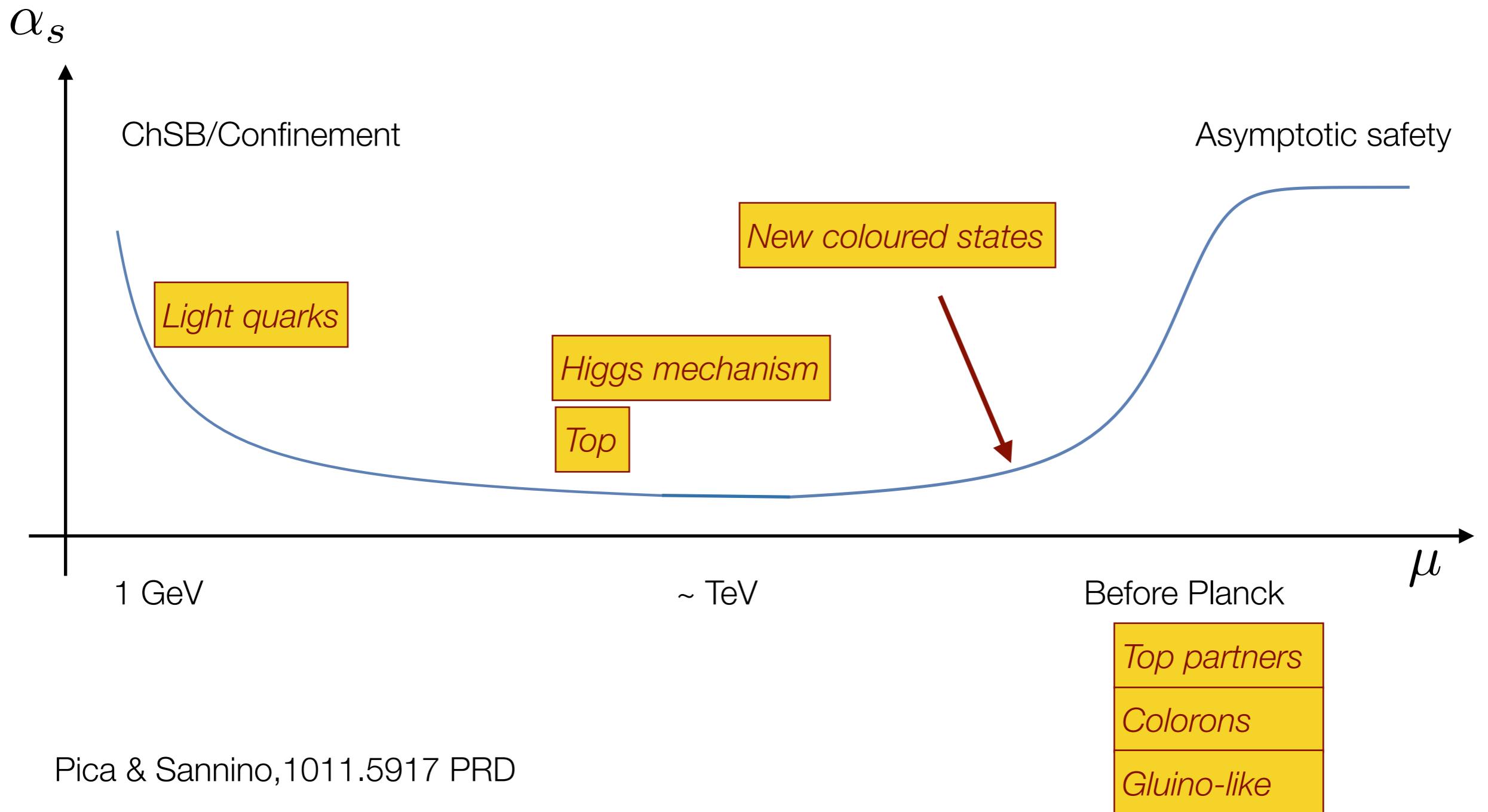
Sannino, 1511.09022



Pica & Sannino, 1011.5917 PRD

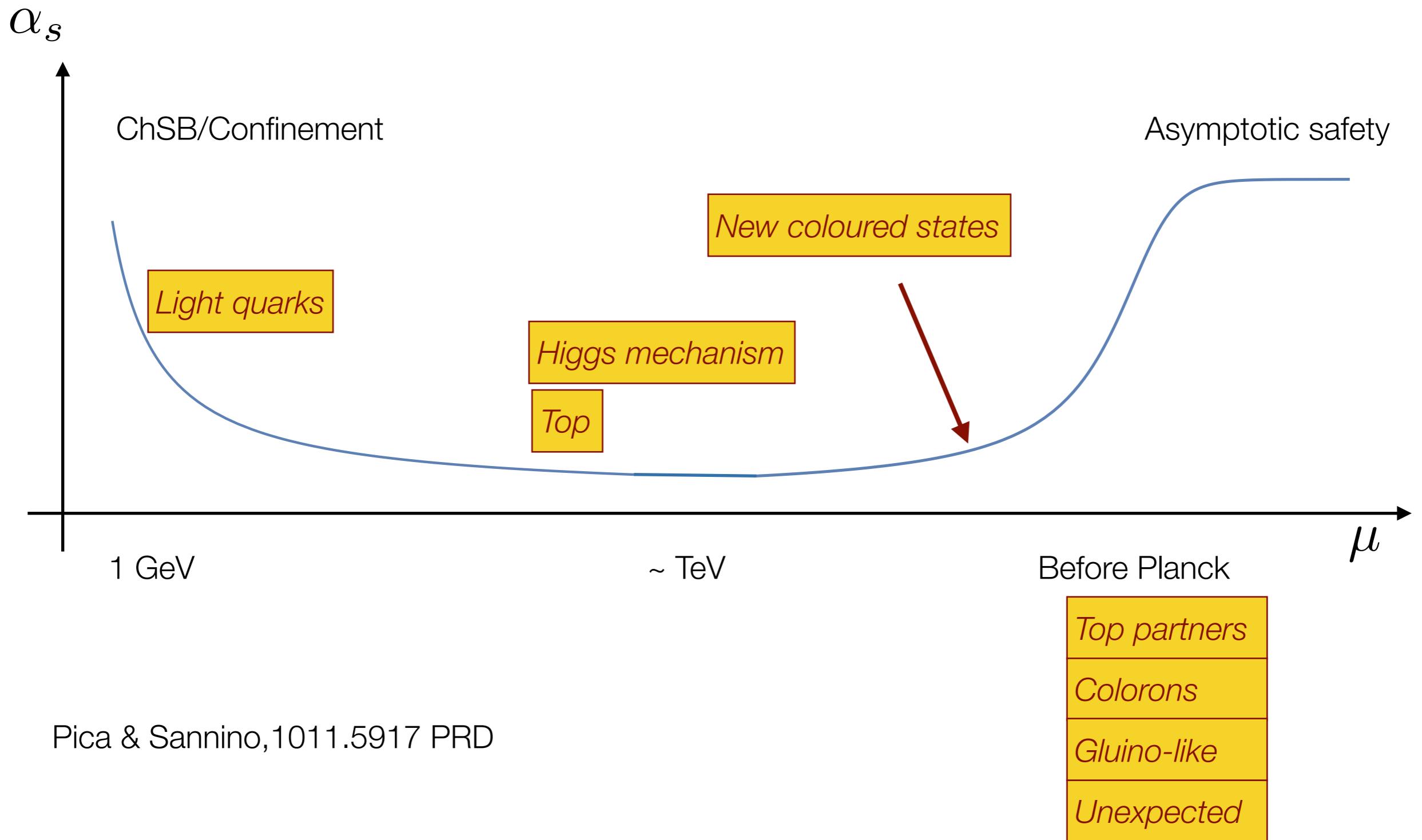
# Safe QCD

Sannino, 1511.09022



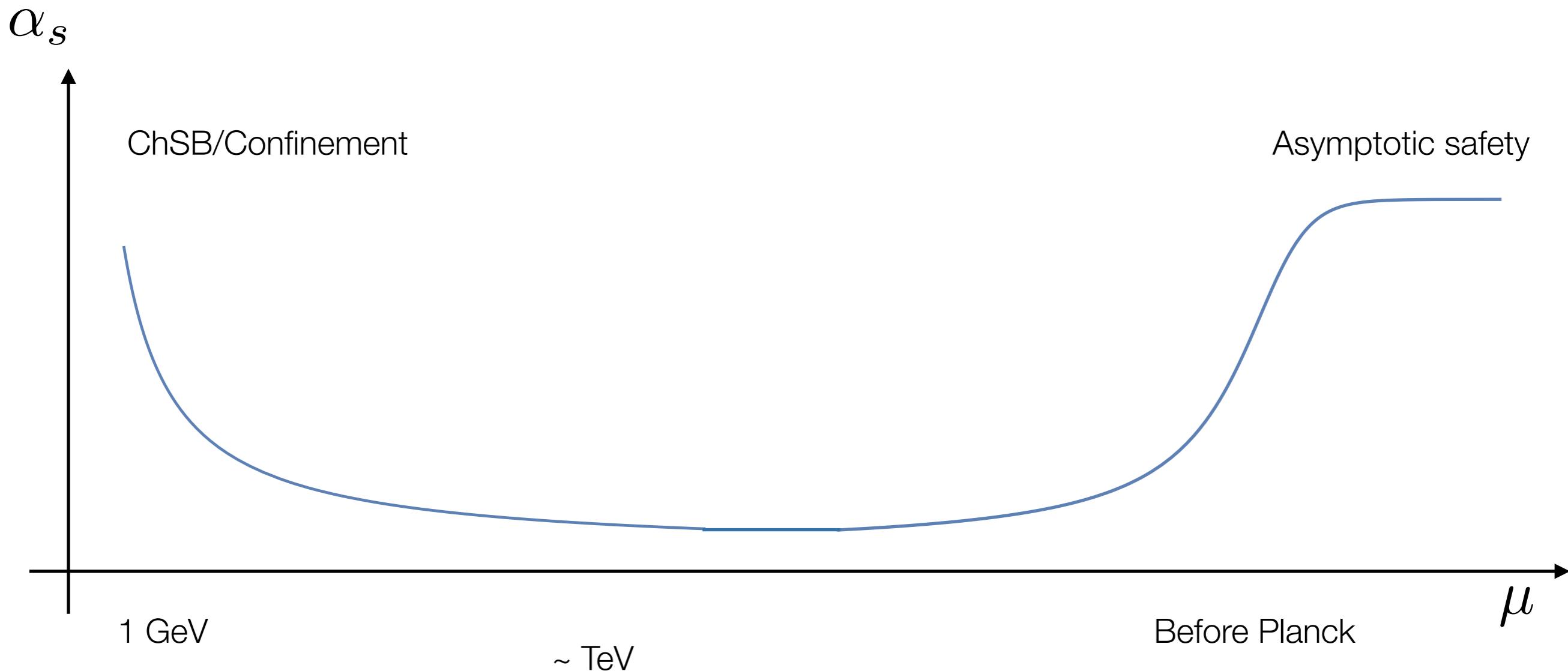
# Safe QCD

Sannino, 1511.09022



# Testing safe QCD scenarios

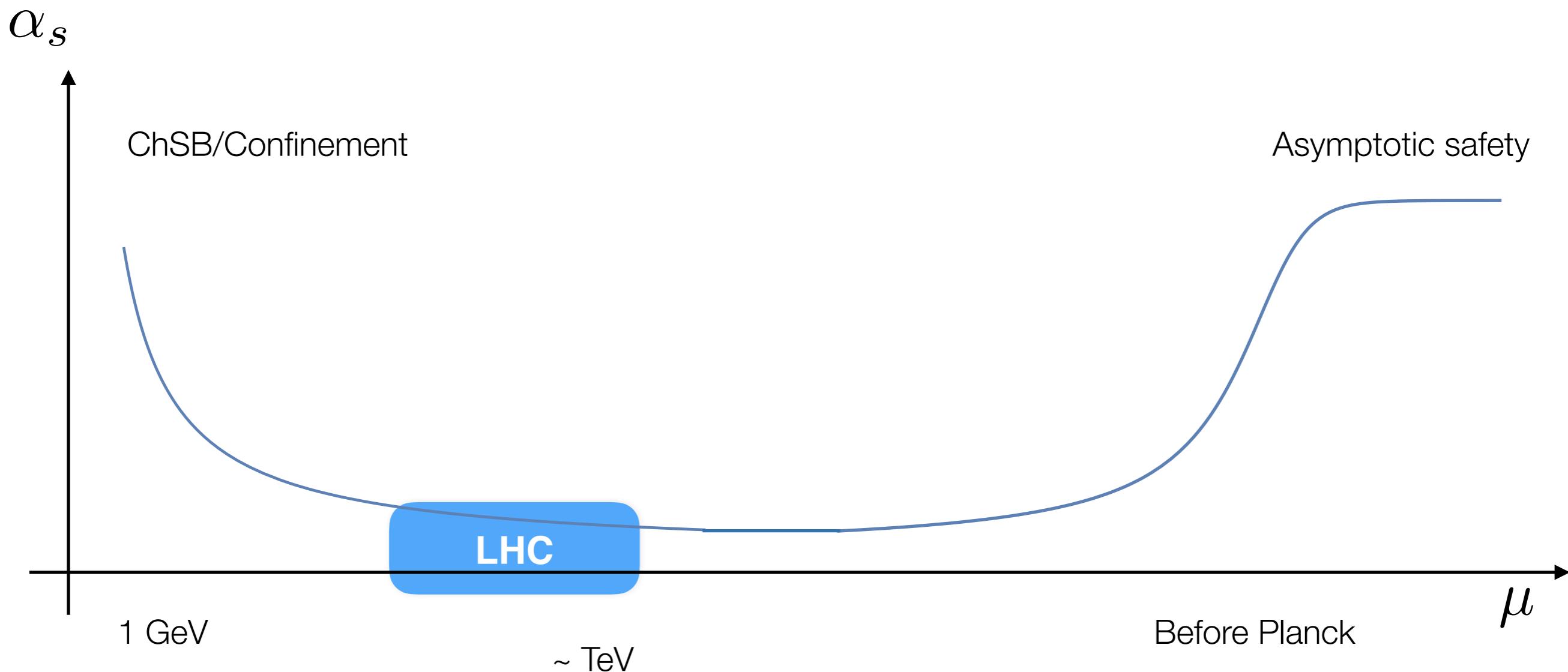
Sannino, 1511.09022



*Asymptotic freedom is not a must for UV complete theories*

# Testing safe QCD scenarios

Sannino, 1511.09022



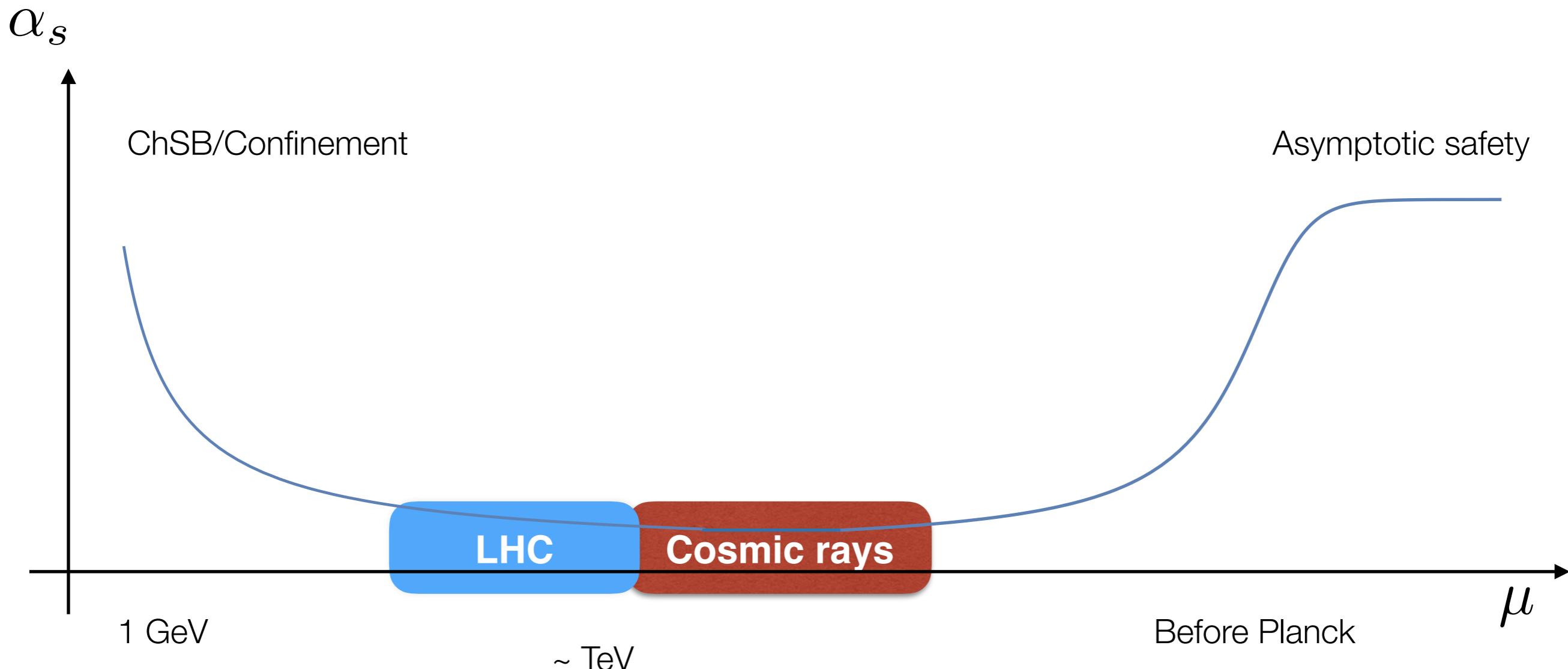
*Asymptotic freedom is not a must for UV complete theories*

Model independent tests of new coloured states at the LHC

Becciolini, Gillioz, Nardecchia, Sannino, Spannowsky 1403.7411, PRD

# Testing safe QCD scenarios

Sannino, 1511.09022



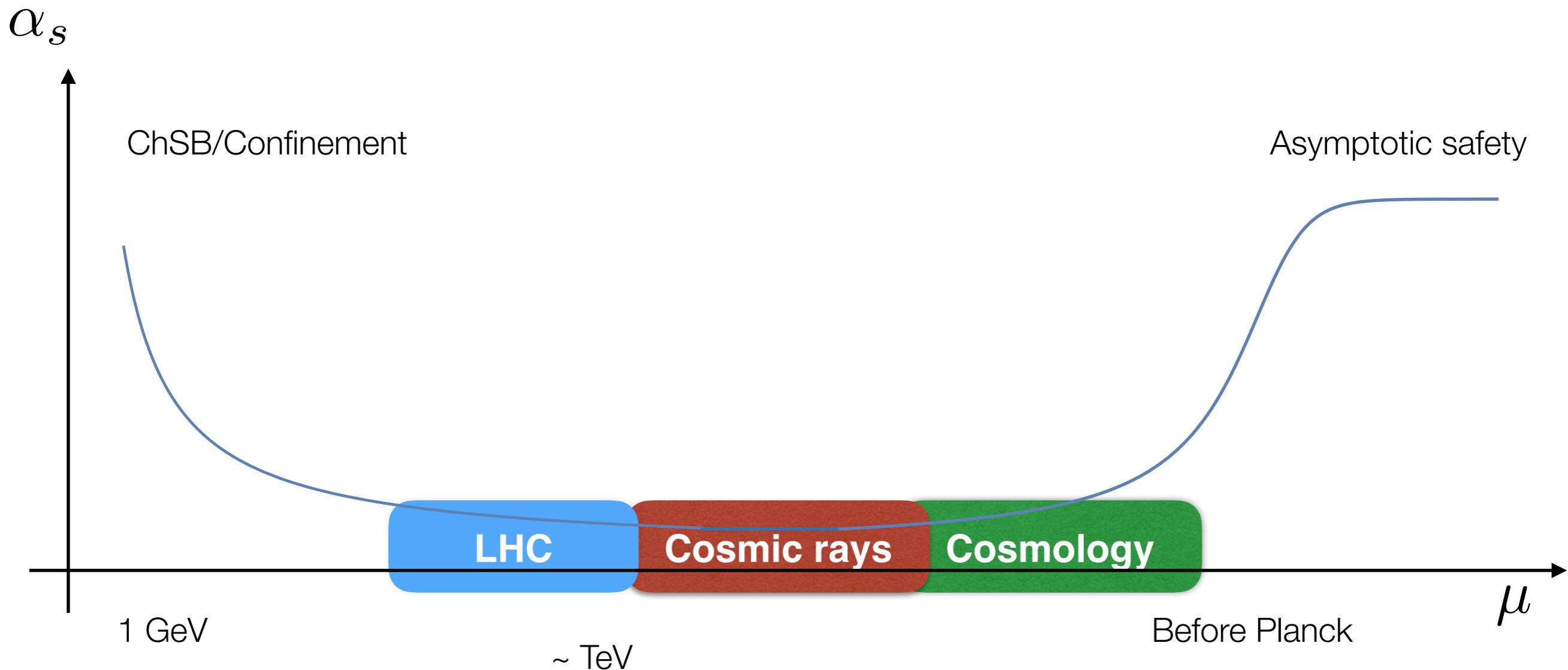
*Asymptotic freedom is not a must for UV complete theories*

Model independent tests of new coloured states at the LHC

Becciolini, Gillioz, Nardecchia, Sannino, Spannowsky 1403.7411, PRD

# Testing safe QCD scenarios

Sannino, 1511.09022



*Asymptotic freedom is not a must for UV complete theories*

Model independent tests of new coloured states at the LHC

Becciolini, Gillioz, Nardecchia, Sannino, Spannowsky 1403.7411, PRD

# Testing safe QCD scenarios

Sannino, 1511.09022

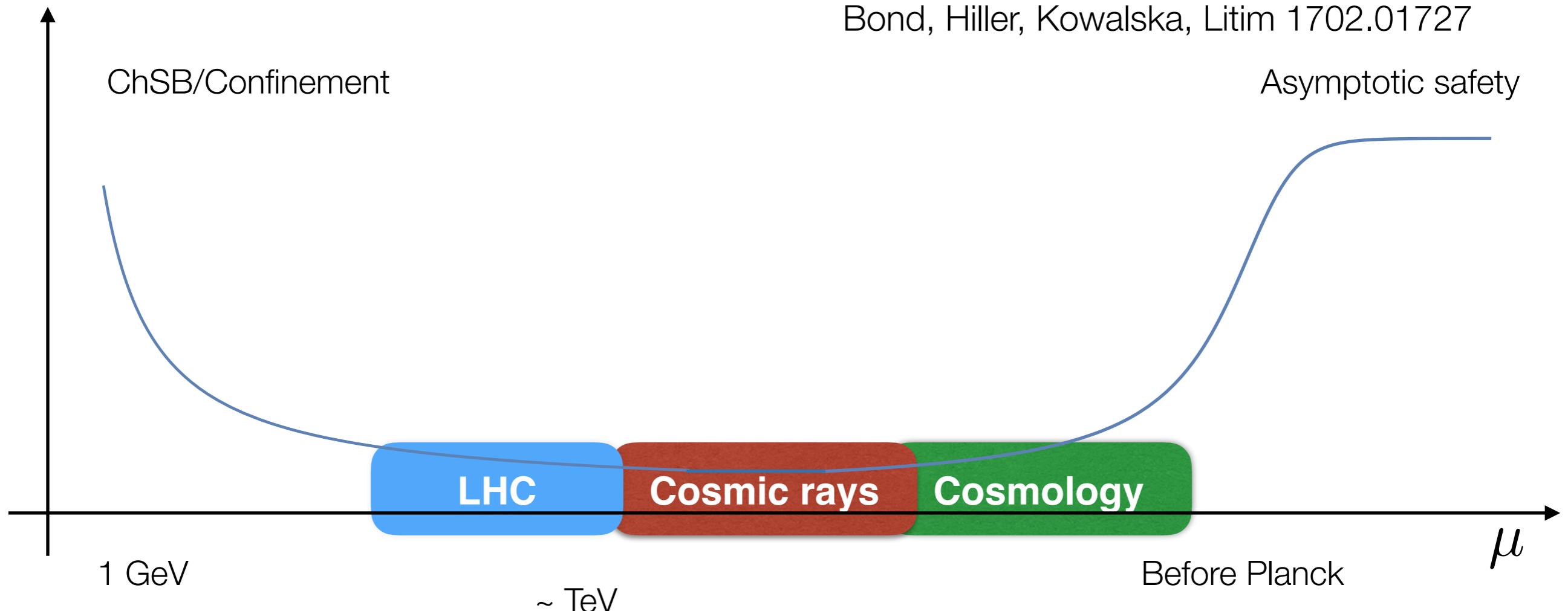
$\alpha_s$

Pelaggi, Sannino, Strumia, Vigiani 1701.01453

Bond, Hiller, Kowalska, Litim 1702.01727

Asymptotic safety

ChSB/Confinement



*Asymptotic freedom is not a must for UV complete theories*

Model independent tests of new coloured states at the LHC

Becciolini, Gillioz, Nardecchia, Sannino, Spannowsky 1403.7411, PRD

# What next?

# What next?

- ◆ Explore different paths for a safe extension of the SM
- ◆ Extend the number of (super) safe theories
- ◆ Cosmological and particle physics consequences
- ◆ Radiative symmetry breaking [Abel and Sannino 1704.00700]
- ◆ Interplay with gravity

Safe Dark Matter

