

Non-Abelian T-duality

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Disclaimers

1) I will not be talking about String theory.

2) “T-duality” is a bit of a misnomer.



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D4-D8 near-horizon

$AdS_6 \times S^4$ in massive IIA

Ferrara, Kehagias, Partouche, Zaffaroni (1998); Brandhuber, Oz (1999)

Recently...

Orbifold theories $\Leftrightarrow AdS_6 \times S^4 / \mathbb{Z}_n$ Bergman & Rodriguez-Gómez

Localization SCFT₅ Jafferis & Pufu; Kim² & Lee

Uniqueness in massive IIA Passias

“In fact this is the only known class of supersymmetric AdS_6 solutions.”



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I KNOW ONE MORE*.

*We all probably know one more. I am saying that I know one more than that!

with D. Rodriguez-Gómez, K. Sfetsos



NEW SUSY AdS_6 solution

$$d\hat{s}^2 = \frac{1}{4}W^2 \left[9ds^2(AdS_6) + 4d\theta^2 \right] + e^{-2A}dr^2 + \frac{r^2 e^{2A}}{r^2 + e^{4A}} ds^2(S^2),$$

$$\hat{B} = \frac{r^3}{r^2 + e^{4A}} \text{vol}(S^2),$$

$$e^{-2\hat{\phi}} = e^{-2\phi} e^{2A} (r^2 + e^{4A}),$$

$$F_1 = -G_1 - mrdr,$$

$$F_3 = \left[-\frac{r^3}{r^2 + e^{4A}} G_1 + \frac{mr^2 e^{4A}}{r^2 + e^{4A}} dr \right] \wedge \text{vol}(S^2),$$

where

$$W = (m \cos \theta)^{-1/6}, \quad e^\phi = \frac{2}{3(m \cos \theta)^{5/6}},$$

$$e^A = \frac{W \sin \theta}{2}, \quad G_1 = \frac{5}{8} (m \cos \theta)^{1/3} \sin^3 \theta d\theta.$$



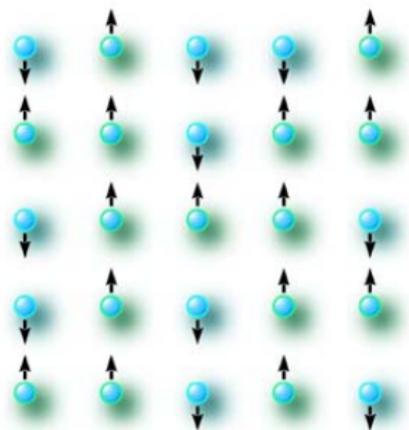
AdS_6 via $SU(2)$

Outline

- Abelian, non-Abelian dualities
- Review T-duality à la Buscher
- Status of non-Abelian T-duality
- Consistent KK reductions
- Return to D4-D8



Traditional (Abelian) Dualities



- Kramers-Wannier (1941)

- $\mathcal{F}_{\text{low}T} \Leftrightarrow \mathcal{F}_{\text{high}T}$

- Phase transition



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non-Abelian duality problem

Well Motivated: Nature is often far from linear

QCD, Heisenberg ferromagnets, FLRW metric

't Hooft, Mandelstam, Englert (1978-1980)

Technical problems: Ex. $D = 2$ $O(3)$ Heisenberg ferromagnet

$$H = -J \sum \cos (\Delta_{\mu} \Omega_i), \quad Z = \int D\Omega \exp \left(\frac{J}{kT} \sum \cos (\Delta_{\mu} \Omega_i) \right)$$

Expand in $Y_{l,m}$, integral gives constraints on four l, m of $Y_{l,m}$



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T-duality

NLSM

$$S = \frac{1}{4\pi\alpha'} \int d^2\zeta [\sqrt{h}(h^{\mu\nu}g_{ij} + i\epsilon^{\mu\nu}b_{ij})\partial_\mu x^i\partial_\nu x^j]$$

Gauge isometry $x^0 \equiv \theta$

$$S = \frac{1}{4\pi\alpha'} \int d^2\zeta [\sqrt{h}h^{\mu\nu}g_{00}(\partial_\mu\theta + A_\mu)(\partial_\nu\theta + A_\nu) + \dots + 2i\epsilon^{\mu\nu}\tilde{\theta}F_{\mu\nu}]$$

Integrate out $A, \tilde{\theta}$ dual coordinate

$$A_\mu = -\frac{1}{g_{00}}(g_{0\alpha}\partial_\mu x^\alpha + i\frac{\epsilon_\mu^\nu}{\sqrt{h}}(b_{0\alpha}\partial_\nu x^\alpha + \partial_\nu\tilde{\theta}))$$



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T-duality

- 1 Compare SM with T-dual SM - **Buscher (1988)**

$$\begin{aligned}\tilde{g}_{00} &= g_{00}^{-1}, & \tilde{g}_{0\alpha} &= -g_{00}^{-1}b_{0\alpha}, & \tilde{b}_{0\alpha} &= -g_{00}^{-1}g_{0\alpha}, \\ 2\tilde{\phi} &= 2\phi - \ln g_{00}, + \dots\end{aligned}$$

- 2 RR-flux: reduction to *unique* $\mathcal{N} = 2$ $D = 9$ supergravity
Bergshoeff, Hull & Ortin (1995)

- 3 KS rotation: $\tilde{\epsilon}_- = \epsilon_-$, $\tilde{\epsilon}_+ = \Omega\epsilon_+$, $\Omega = \sqrt{g_{00}^{-1}\Gamma_{11}\Gamma^0}$

$$\begin{aligned}\delta\tilde{\lambda}_- &= \delta\lambda_- - g_{00}^{-1}\Gamma^0\delta\psi_{-0}, \\ \delta\tilde{\lambda}_+ &= \Omega\left(\delta\lambda_+ - g_{00}^{-1}\Gamma^0\delta\psi_{+0}\right).\end{aligned}$$

Rotation on RR sector **Hassan (1999)**



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Status of non-Abelian T-duality

Old story Buscher (1988, 1992), De la Ossa, Quevedo; Roček, Verlinde
Gasperini, Ricci, Veneziano (1993), Alvarez, Alvarez-Gaume, Lozano (1994)

Gauge isometry $\partial x^m \rightarrow Dx^m \equiv \partial x^m + A^\alpha (T_\alpha)_n^m x^n$

Symmetry of type II sugra with $SU(2) \times SU(2)$ PCM Sfetsos, Thompson (2010)

i) $AdS_3 \times S^3 \times T^4$ ii) $AdS_5 \times S^5 \Rightarrow$ Gaiotto-Maldacena

Coset G/H manifolds, $S^4 \simeq SO(4)/SO(3)$, $S^5 \simeq SO(6)/SO(5)$
chirality unchanged Lozano, Ó C, Sfetsos, Thompson (2011)

Not traditional duality - one way, isometries (SUSY) break, non-compact



IIB reduction on S^3

Motivation: Can one emulate Bergshoeff, Hull, Ortin?

Natural ansatz with $SO(4)$ symmetry Itsios, Lozano, Ó C, Sfetsos (2012)

$$ds^2 = ds^2(M_7) + e^{2A} ds^2(S^3).$$

Fluxes

$$F_5 = G_2 \wedge \text{Vol}(S^3) - e^{-3A} \star_7 G_2,$$

$$F_3 = G_3 - m \text{Vol}(S^3),$$

$$F_1 = G_1.$$

Background B, Φ

$SO(4)$ singlets \Rightarrow consistent reduction to $D = 7$ theory (not SUGRA)



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Massive IIA on $\mathbb{R} \times S^2$

NS sector

$$ds^2 = ds^2(M_7) + e^{-2A} dr^2 + \frac{r^2 e^{2A}}{r^2 + e^{4A}} ds^2(S^2),$$

$$\tilde{B} = B + \frac{r^3}{r^2 + e^{4A}} \text{vol}(S^2),$$

$$e^{-2\tilde{\Phi}} = e^{-2\Phi} e^{2A} (r^2 + e^{4A}).$$

Flux ansatz

$$\hat{F}_0 = m,$$

$$\hat{F}_2 = \frac{mr^3}{r^2 + e^{4A}} \text{vol}(S^2) + r dr \wedge G_1 - G_2,$$

$$\hat{F}_4 = \frac{r^2 e^{4A}}{r^2 + e^{4A}} G_1 \wedge dr \wedge \text{vol}(S^2) - \frac{r^3}{r^2 + e^{4A}} G_2 \wedge \text{vol}(S^2) + r dr \wedge G_3 + e^{3A} \star_7 G_3.$$



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$D = 7$ Action

Field content: $G_1 = dC_0$, $G_2 = dC_1 - mB$, $G_3 = dC_2 - C_0H$

Diagonalisation & rescaling: $\Phi = 5\check{\Phi} + \frac{3}{2}A$

$$\begin{aligned}\mathcal{L}_{\text{Einstein}} = & R - 3(\partial A)^2 - 20(\partial\Phi)^2 - \frac{1}{2}e^{10\Phi+3A}(\partial C_0)^2 - \frac{1}{12}e^{-8\Phi}H^2 \\ & - \frac{1}{2}\left(m^2e^{14\Phi-3A} - 3e^{4\Phi-2A} + \frac{1}{2}e^{6\Phi-3A}G_2^2 + \frac{1}{6}e^{2\Phi+3A}G_3^2\right) \\ & + G_2 \wedge C_2 \wedge H\end{aligned}$$

Any solution uplifts to both IIB and massive IIA: pp-wave, LLM, Lifshitz

cf reduction to $D = 5$ Romans [Lu, Pope, Tran \(1999\)](#); [Gauntlett, Varela \(2007\)](#)



Rotation on the Killing spinor generic

$$\epsilon = e^X \tilde{\epsilon} = \exp \left(-\frac{1}{2} \tan^{-1} \left(\frac{e^{2A}}{r} \right) \Gamma^{\theta\phi} \sigma^3 \right) \tilde{\epsilon}$$

Variations between IIA and IIB match up to

$$\delta\psi_r = e^X \left[\frac{1}{2} \not{\partial} A \Gamma_r - \frac{e^{-A}}{4} \Gamma^{\theta\phi} \sigma^3 + \frac{e^\Phi}{8} \left(m e^{-3A} \Gamma^{\theta\phi} \Gamma_r (i\sigma^2) \right. \right. \\ \left. \left. - \mathcal{G}_1 (i\sigma^2) - \frac{e^{-3A}}{2} \mathcal{G}_2 \Gamma^{\theta\phi} \Gamma_r \sigma^1 - \frac{1}{6} \mathcal{G}_3 \sigma^1 \right) \right] \tilde{\epsilon}$$

Redefinitions $\tilde{\epsilon}_+ = \Gamma^r \epsilon_+$, $\tilde{\epsilon}_- = -\epsilon_-$;

$\delta\psi_r$ reproduces D3, D1-D5 projections of **Sfetsos, Thompson**



D4-D8

$AdS_6 \times S^4$

$$ds^2 = \frac{1}{4} W^2 \left[9 ds^2(AdS_6) + 4 \left(d\theta^2 + \sin^2 \theta ds^2(S^3) \right) \right],$$

$$F_4 = 5(m \cos \theta)^{1/3} \sin^3 \theta d\theta \wedge \text{vol}(S^3),$$

$$e^\phi = \frac{2}{3(m \cos \theta)^{5/6}}, \quad W = (m \cos \theta)^{-1/6}.$$

Warping breaks $SO(5)$ to $SO(4) \simeq SU(2)_R \times SU(2)_G$

Can do Abelian Hopf-fibre T-duality [Cvetic, Lu, Pope, Vazquez-Portiz \(1999\)](#)

More generally, also have freedom for $SU(2)$.



Future directions

Is there a dual CFT? *Rodríguez-Gómez, Ó C, Sfetsos (sometime soon)*

Classification of AdS_6 in IIB, more solutions?

NA T-duality turns on a B -field - applications?

Do thermodynamics change? - *1206.1367*

non-Abelian pre Big Bang cosmology? *Gasperini, Veneziano (1992)*

SUSY breaking models?



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