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THE LARGE SCALE STRUCTURE OF WARM DARK MATTER

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OUTLINE

- free-streaming
- the halo model
- simulations dark matter only
- degeneracies: WDM+b & WDM+ ν

ACDM, PRESENTTIME

73% Dark Energy or Λ
23% Cold Dark Matter
4.5% Baryonic matter
negligible radiation and other standard model matter

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WARM DARK MATTER PARTICLE

- DM: elementary particles
- DM: here assume thermal or sterile v
- We have: $\Omega_{\rm dm}(m_{\rm dm},g^*(T_{\rm d}),\langle\sigma v\rangle)$
- choose a DM model → get rid of one variable

CONSTRAINTS

• I measurement: Dark Matter abundance from CMB experiments (WMAP7, assuming standard Λ CDM):

 $\Omega_{\rm dm}=0.227\pm0.014$

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• I measurement: Dark Matter abundance from CMB experiments (WMAP7, assuming standard Λ CDM):

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- I constraint: Choose a heavy DM particle á la WIMP

 → constraints from DM high-density regions
- I constraint: Choose a light particle (m ~ I keV)
 - → Warm Dark Matter
 - → constraints from Large Scale Structure

FREE-STREAMING

- early relativistic particles \Rightarrow free-streaming
- .: the linear matter power spectrum small scale suppression
- How does this translate to non-linearities?



Smith & Markovic (PRD; 2011)

in CDM, assume:
 all of density field ← haloes

$$\rho(\boldsymbol{x}) = \sum_{i=1}^{N} M_i u_{\rm h}(|\boldsymbol{x} - \boldsymbol{x}_{0,i}|, M_i)$$

Smith & Markovic (PRD; 2011)

in CDM, assume:
 all of density field ← haloes

$$ho(m{x}) = \sum_{i=1}^{N} M_i \, u_{
m h}(|m{x} - m{x}_{0,i}|, M_i) +
ho_{
m s}(m{x})$$

in WDM, split:
 clumped + smooth

Smith & Markovic (PRD; 2011)

- volume average
- get background values

$$\langle \rho(\boldsymbol{x}) \rangle = \langle \rho_{\rm s}(\boldsymbol{x}) \rangle + \langle \rho_{\rm c}(\boldsymbol{x}) \rangle$$

define a "clumped fraction"

Smith & Markovic (PRD; 2011)

- volume average
- get background values

 $\bar{\rho} = \bar{\rho}_{\rm s} + \bar{\rho}_{\rm c}$

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- get background values

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define a "clumped fraction"

mass function / (no. density of haloes per dM)

$$f = \bar{\rho}_{\rm c} / \bar{\rho} = \frac{1}{\bar{\rho}} \int_{M_{\rm cut}}^{\infty} dM M \frac{dn}{dM}$$

Smith & Markovic (PRD; 2011)

$$\langle \rho_i(\boldsymbol{x}) \rho_j(\boldsymbol{x} + \boldsymbol{r}) \rangle = \bar{\rho}_i \bar{\rho}_j \left(1 + \langle \delta_i(\boldsymbol{x}) \delta_j(\boldsymbol{x} + \boldsymbol{r}) \rangle \right)$$

 $\rho_i(\boldsymbol{x}) \equiv \bar{\rho}_i \left(1 + \delta_i(\boldsymbol{x}) \right)$

Smith & Markovic (PRD; 2011)

• we want statistic of clustering, so as always:

 $\begin{array}{l} \langle \rho_i(\boldsymbol{x})\rho_j(\boldsymbol{x}+\boldsymbol{r})\rangle = \bar{\rho}_i\bar{\rho}_j\left(1+\xi_{ij}(\boldsymbol{r})\right) \\ \rho_i(\boldsymbol{x}) \equiv \bar{\rho}_i\left(1+\delta_i(\boldsymbol{x})\right) \\ \xi_{ij}(|\boldsymbol{r}|) \equiv \xi_{ij}(\boldsymbol{r}) \equiv \langle \delta_i(\boldsymbol{x})\delta_j(\boldsymbol{x}+\boldsymbol{r})\rangle \\ + \text{ statistical} \\ \text{ isotropy } \& \\ \text{ homogeneity} \end{array}$

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Smith & Markovic (PRD; 2011)

$$\langle
ho_i(\boldsymbol{x})
ho_j(\boldsymbol{x}+\boldsymbol{r})
angle = ar{
ho}_iar{
ho}_j(1+\xi_{ij}(r))$$

= $\langle
ho_{\mathrm{s}}(\boldsymbol{x})
ho_{\mathrm{s}}(\boldsymbol{x}+\boldsymbol{r})
angle + \langle
ho_{\mathrm{c}}(\boldsymbol{x})
ho_{\mathrm{c}}(\boldsymbol{x}+\boldsymbol{r})
angle + 2 \langle
ho_{\mathrm{c}}(\boldsymbol{x})
ho_{\mathrm{s}}(\boldsymbol{x}+\boldsymbol{r})
angle$

Smith & Markovic (PRD; 2011)

$$\langle \rho_i(\boldsymbol{x}) \rho_j(\boldsymbol{x} + \boldsymbol{r}) \rangle = \bar{\rho}_i \bar{\rho}_j (1 + \xi_{ij}(r))$$

= $\bar{\rho}_s^2 (1 + \xi_{ss}(r)) + \bar{\rho}_c^2 (1 + \xi_{cc}(r)) + 2\bar{\rho}_c \bar{\rho}_s (1 + \xi_{cs}(r))$

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 $\xi(r) = (1 - f)^2 \xi_{\rm ss}(r) + 2f(1 - f)\xi_{\rm sc}(r) + f^2 \xi_{\rm cc}(r)$

 \downarrow

Smith & Markovic (PRD; 2011)

• we want statistic of clustering, so as always:

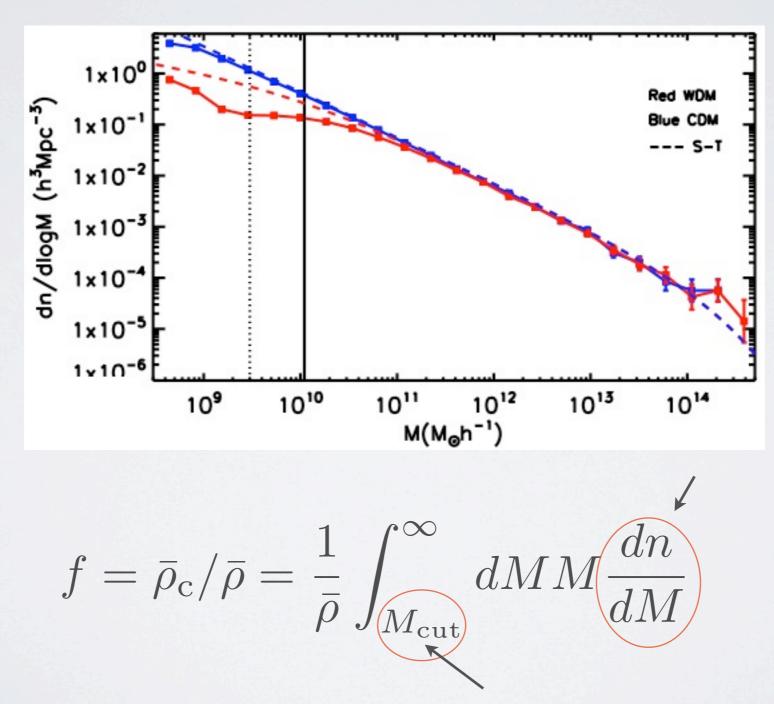
$$\langle \rho_i(\boldsymbol{x}) \rho_j(\boldsymbol{x} + \boldsymbol{r}) \rangle = \bar{\rho}_i \bar{\rho}_j (1 + \xi_{ij}(r))$$

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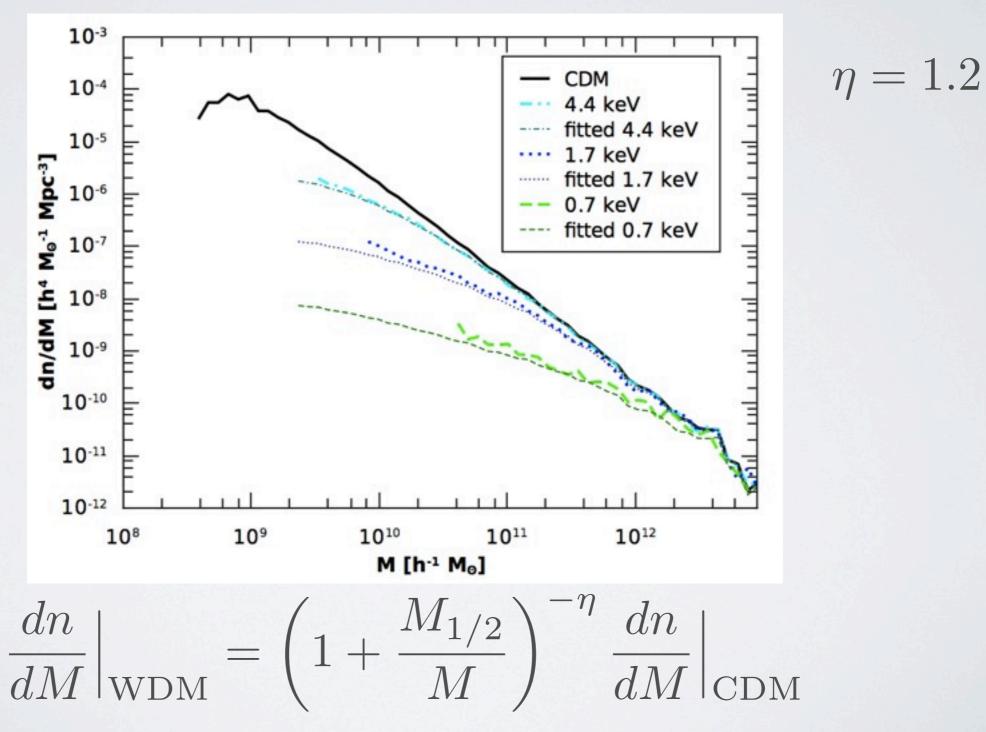
 $P(k) = (1 - f)^2 P_{\rm ss}(k) + 2(1 - f) f P_{\rm sc}(k) + f^2 P_{\rm cc}(k)$ $P_{\rm cc}^{\rm 2h}(k) + P_{\rm cc}^{\rm 1h}(k)$

 \downarrow

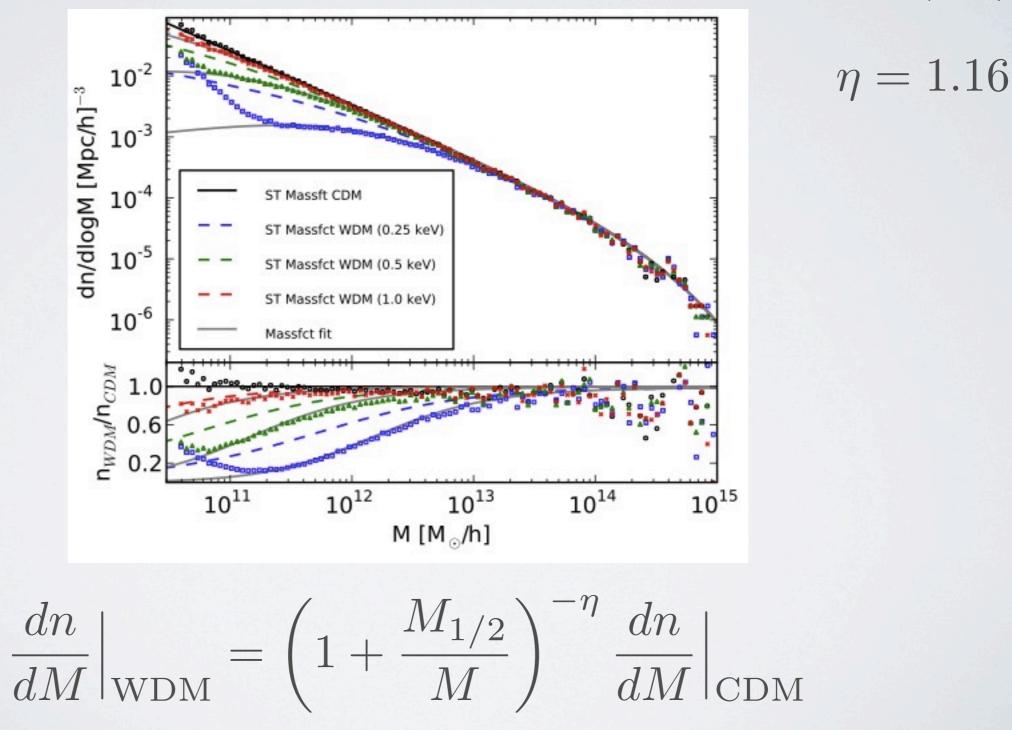
Zavala et al. (2009)



Dunstan, Abazajian, Polisensky & Ricotti (2011)



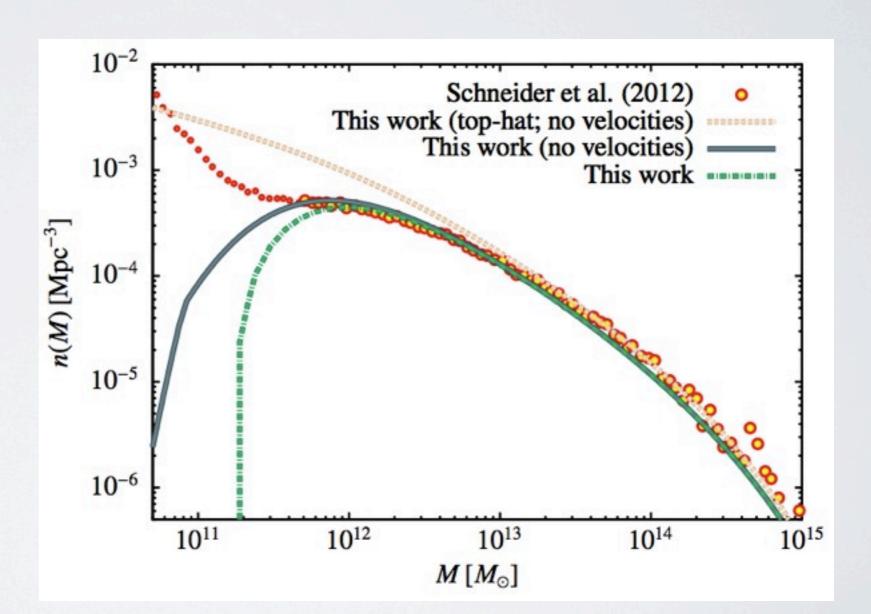
Schneider, Smith, Maccio & Moore (2012)

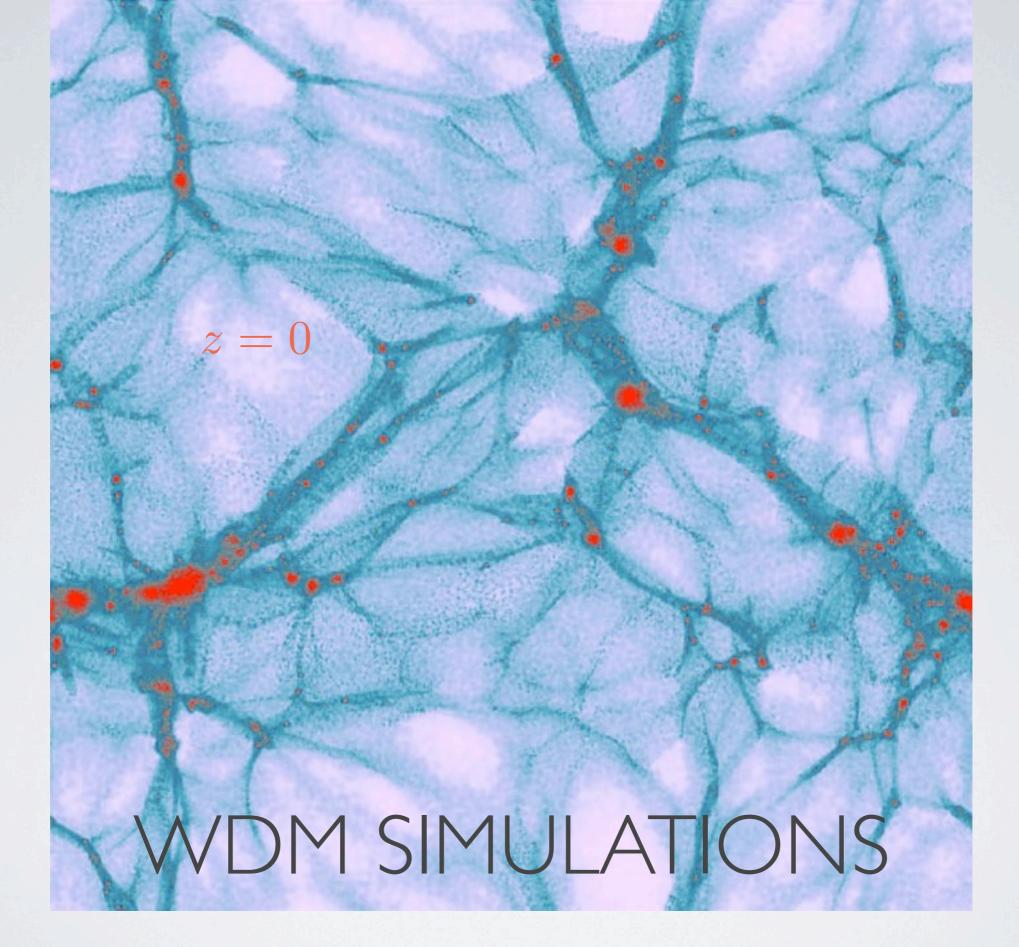


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Benson et al. (2012)

- arXiv, Monday
- merger trees
- no top-down
- smooth accretion!

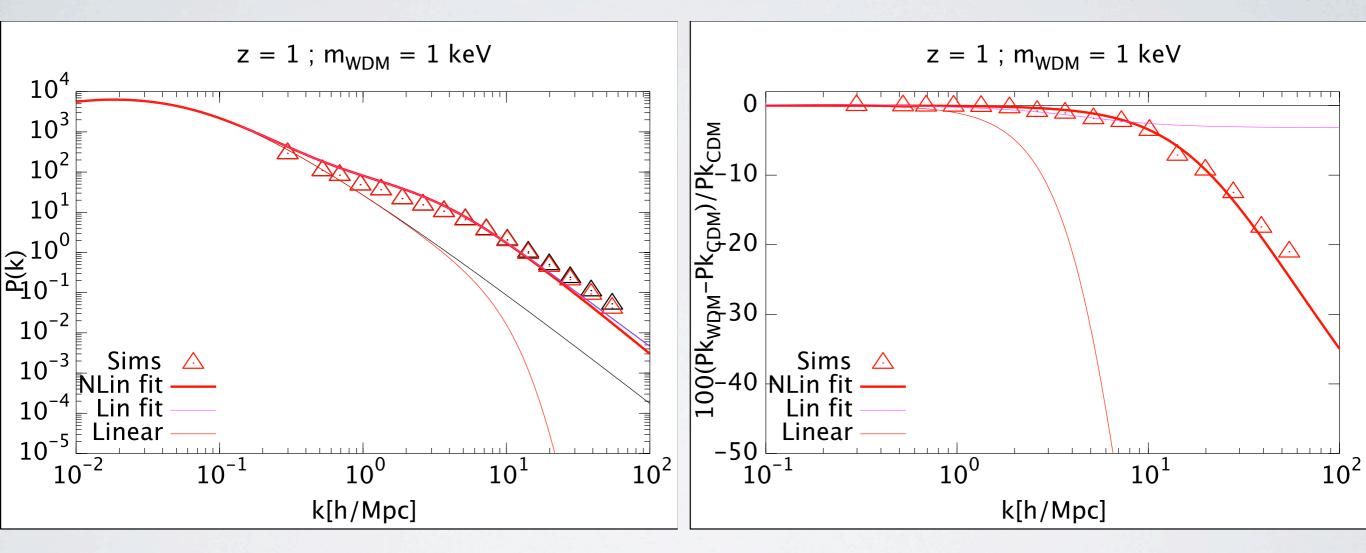




Viel, Markovic, Baldi & Weller (MNRAS; 2012)

WDM SIMULATIONS

Viel, Markovic, Baldi & Weller (MNRAS; 2012)

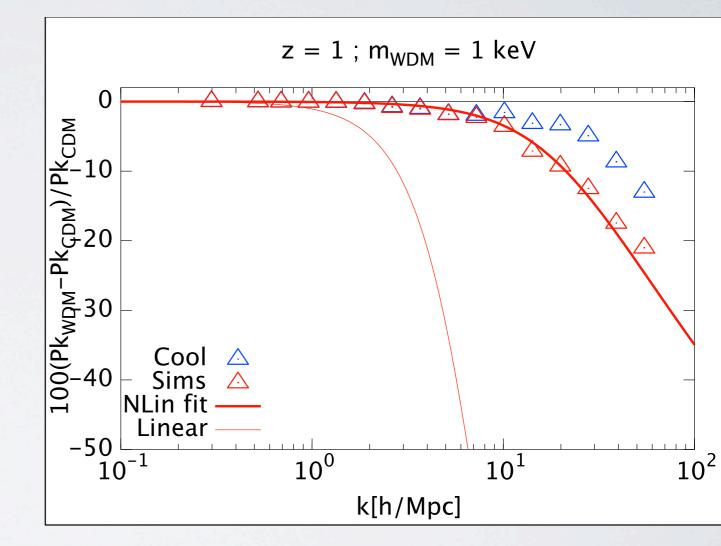


 $T_{\rm nlin}^2(k) \equiv P_{\rm WDM}(k) / P_{\Lambda {\rm CDM}}(k) = (1 + (\alpha \, k)^{\nu l})^{-s/\nu}$

Depends on particle mass

WDM+b

Viel, Markovic, Baldi & Weller (MNRAS; 2012)



- blue: SF, winds
- cooling erases WDM effect

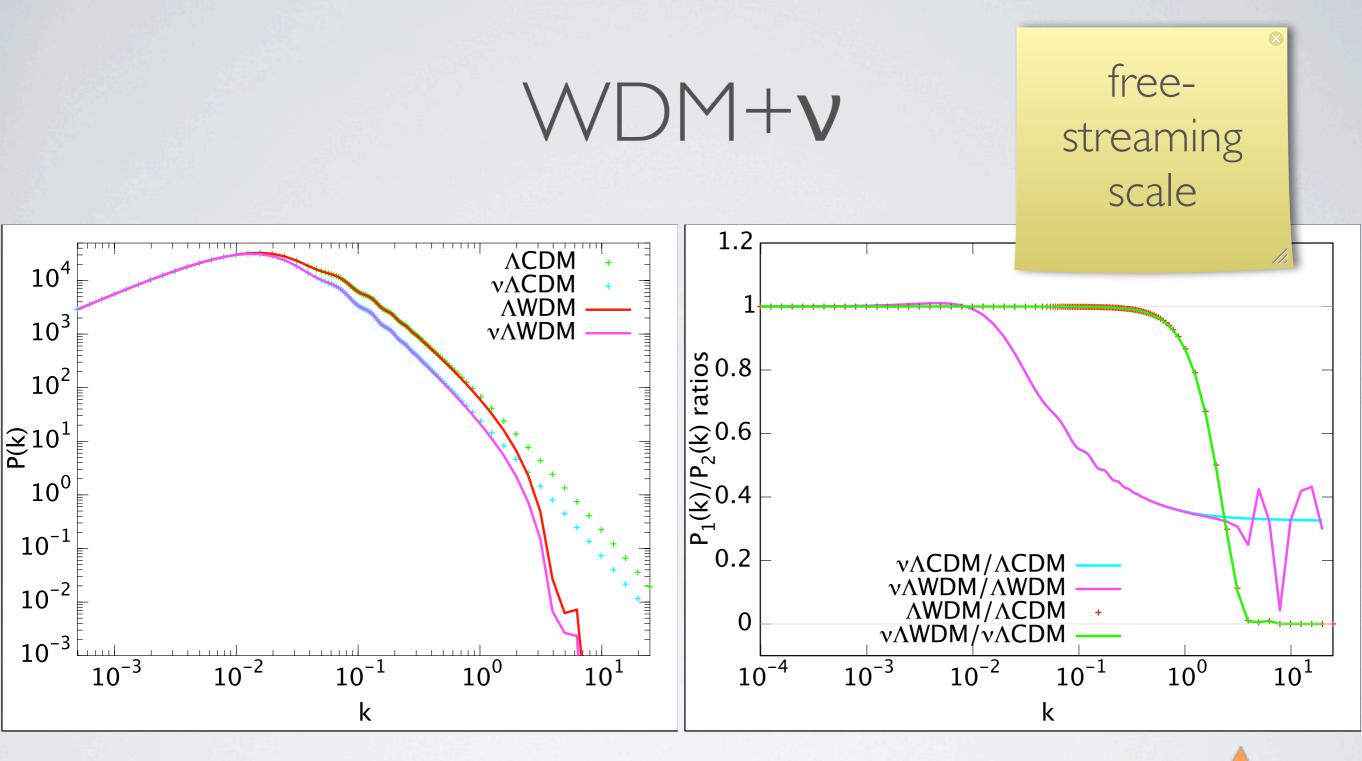


WDM+v

Markovic, Abdalla, Lahav & Weller (in prep)

- neutrinos = HDM
- CLASS code (Lesgourgues, 2011)
- halofit (CDM-based)
- ... new, calibrated halo model

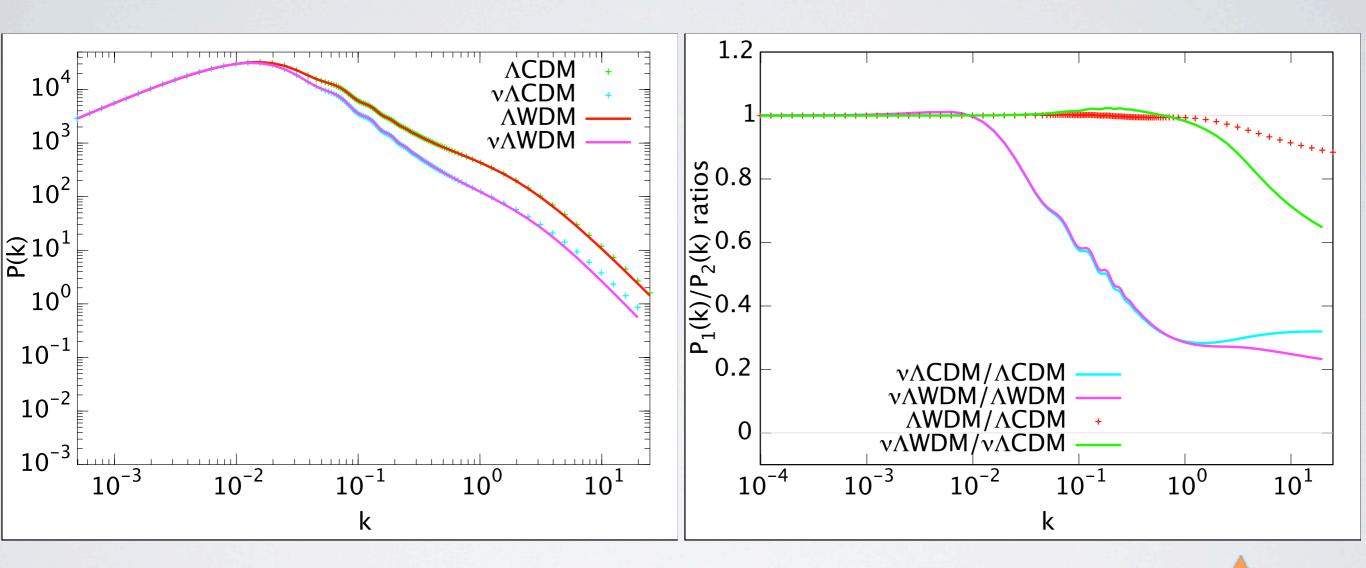




Linear



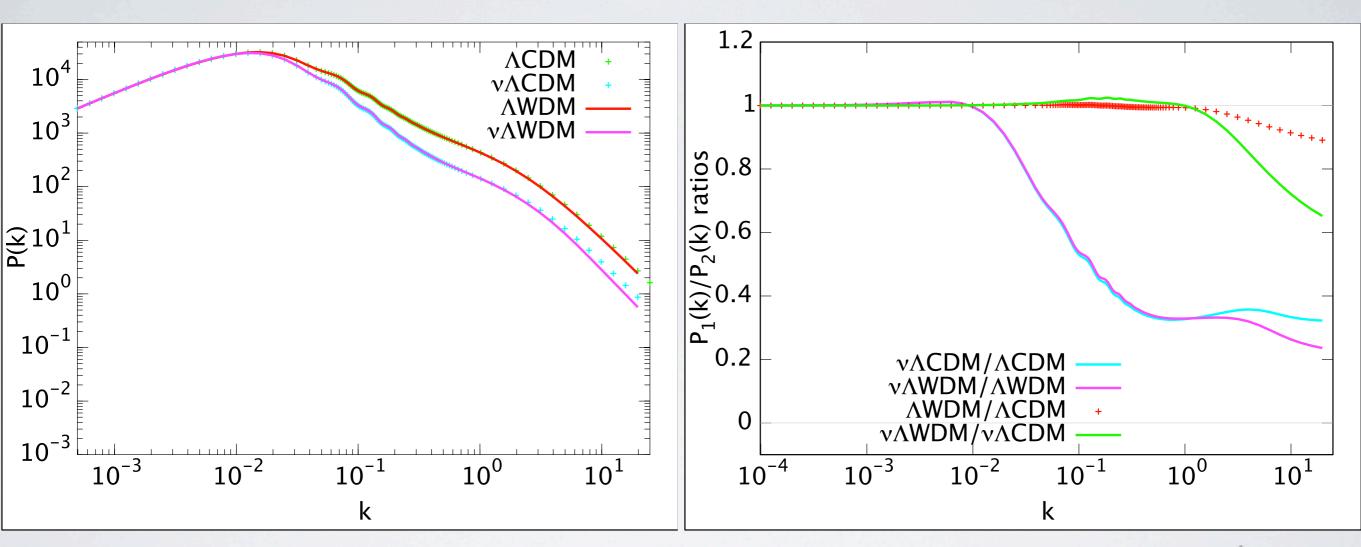
WDM+v



halofit



WDM+v



halofit +correction



TL;DR

- WDM is a generalisation of CDM.
- We know good prescriptions to calculate non-linear corrections in Λ WDM.
- BUT we need to know baryonic effects!

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