

Scattering amplitudes in N=6 Chern-Simons-matter theories

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Outline

- **AdS/CFT and ABJM**
- **Gluon Scattering at strong coupling**
- **Super Wilson Loops/Scattering in N=4**
- **Results in N=6 Chern-Simons matter theories**
- **Conclusions**

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AdS/CFT

- IIB Strings on AdS₅×S⁵

- $\sqrt{\lambda} = R^2$

- Perturbation in powers of $\frac{1}{\sqrt{\lambda}}$

- N=4 SYM field Theory

- $\lambda = g_{YM}^2 N$

- Perturbation in powers of $\frac{1}{\lambda}$

- IIA Strings on AdS₄×CP³

- $\sqrt{\lambda} = R^2$

- N=6 Chern-Simons matter theory

- $\lambda = \frac{N}{k}$



ABJM Component Action

$$S_{\text{kin}} = \frac{k}{4\pi} \int d^3x \text{tr} \left[A_\alpha \epsilon^{\alpha\beta\gamma} \partial_\beta A_\gamma - \hat{A}_\alpha \epsilon^{\alpha\beta\gamma} \partial_\beta \hat{A}_\gamma + Y_A^\dagger \partial_\mu \partial^\mu Y^A + i\psi^{\dagger B} \not{\partial} \psi_B \right]$$

$$\begin{aligned}
 S_{\text{int}} = \frac{k}{4\pi} \int d^3x \text{tr} & \left[\frac{2}{3} i \epsilon^{\alpha\beta\gamma} (A_\alpha A_\beta A_\gamma - \hat{A}_\alpha \hat{A}_\beta \hat{A}_\gamma) \right. \\
 & - i A_\mu Y^A \overleftrightarrow{\partial}^\mu Y_A^\dagger - i \hat{A}_\mu Y_A^\dagger \overleftrightarrow{\partial}^\mu Y^A + 2 Y_A^\dagger A_\mu Y^A \hat{A}^\mu \\
 & - \hat{A}_\mu \hat{A}_\mu Y_A^\dagger Y^A - A_\mu A_\mu Y^A Y_A^\dagger \\
 & - \psi^{\dagger B} \not{A} \psi_B + \hat{A}_\mu \psi^{\dagger B} \gamma^\mu \psi_B \\
 & + \frac{1}{12} Y^A Y_B^\dagger Y^C Y_D^\dagger Y^E Y_F^\dagger (\delta_A^B \delta_C^D \delta_E^F + \delta_A^F \delta_C^B \delta_E^D - 6 \delta_A^B \delta_C^F \delta_E^D + 4 \delta_A^D \delta_C^F \delta_E^B) \\
 & - \frac{i}{2} (Y_A^\dagger Y^B \psi^{\dagger C} \psi_D - \psi_D \psi^{\dagger C} Y^B Y_A^\dagger) (\delta_B^A \delta_C^D - 2 \delta_C^A \delta_B^D) \\
 & \left. + \frac{i}{2} \epsilon^{ABCD} Y_A^\dagger \psi_B Y_C^\dagger \psi_D - \frac{i}{2} \epsilon_{ABCD} Y^A \psi^{\dagger B} Y^C \psi^{\dagger D} \right]
 \end{aligned}$$



ABJM Component Action

$$S_{\text{kin}} = \frac{k}{4\pi} \int d^3x \text{tr} \left[A_\alpha \epsilon^{\alpha\beta\gamma} \partial_\beta A_\gamma - \hat{A}_\alpha \epsilon^{\alpha\beta\gamma} \partial_\beta \hat{A}_\gamma + Y_A^\dagger \partial_\mu \partial^\mu Y^A + i\psi^{\dagger B} \not{\partial} \psi_B \right]$$

CS

$$S_{\text{int}} = \frac{k}{4\pi} \int d^3x \text{tr} \left[\frac{2}{3} i \epsilon^{\alpha\beta\gamma} (A_\alpha A_\beta A_\gamma - \hat{A}_\alpha \hat{A}_\beta \hat{A}_\gamma) \right. \\ \left. - i A_\mu Y^A \overleftrightarrow{\partial}^\mu Y_A^\dagger - i \hat{A}_\mu Y_A^\dagger \overleftrightarrow{\partial}^\mu Y^A + 2 Y_A^\dagger A_\mu Y^A \hat{A}^\mu \right. \\ \left. - \hat{A}_\mu \hat{A}_\mu Y_A^\dagger Y^A - A_\mu A_\mu Y^A Y_A^\dagger \right. \\ \left. - \psi^{\dagger B} \not{A} \psi_B + \hat{A}_\mu \psi^{\dagger B} \gamma^\mu \psi_B \right] \quad \left. \vphantom{S_{\text{int}}} \right\} \text{m.c.}$$

ϕ^6 -potential

$$\left\{ + \frac{1}{12} Y^A Y_B^\dagger Y^C Y_D^\dagger Y^E Y_F^\dagger (\delta_A^B \delta_C^D \delta_E^F + \delta_A^F \delta_C^B \delta_E^D - 6 \delta_A^B \delta_C^F \delta_E^D + 4 \delta_A^D \delta_C^F \delta_E^B) \right. \\ \left. - \frac{i}{2} (Y_A^\dagger Y^B \psi^{\dagger C} \psi_D - \psi_D \psi^{\dagger C} Y^B Y_A^\dagger) (\delta_B^A \delta_C^D - 2 \delta_C^A \delta_B^D) \right. \\ \left. + \frac{i}{2} \epsilon^{ABCD} Y_A^\dagger \psi_B Y_C^\dagger \psi_D - \frac{i}{2} \epsilon_{ABCD} Y^A \psi^{\dagger B} Y^C \psi^{\dagger D} \right\} \quad \left. \vphantom{\phi^6} \right\} \text{3d Yukawa}$$



N=2 Action



N=2 Action

$$S_{CS} = \frac{k}{4\pi} \int d^3x d^4\theta \int_0^1 dt \operatorname{tr} \left[V \bar{D}^\alpha \left(e^{-tV} D_\alpha e^{tV} \right) - V \leftrightarrow \hat{V} \right]$$

$$S_{mat} = \int d^3x d^4\theta \operatorname{tr} \left(\bar{A}^A e^{\hat{V}} \mathcal{A}_A e^{-V} + \bar{B}_A e^V \mathcal{B}^A e^{-\hat{V}} \right)$$

$$S_{pot} = \frac{2\pi i}{k} \int d^3x d^2\theta \epsilon_{AC} \epsilon^{BD} \operatorname{tr} (\mathcal{B}^A \mathcal{A}_D \mathcal{B}^C \mathcal{A}_B) + h.c.$$

\mathcal{B}^A \mathcal{A}_D

Chiral Scalars

V \hat{V}

Gauge Vectors



N=2 Action

$$S_{CS} = \frac{k}{4\pi} \int d^3x d^4\theta \int_0^1 dt \operatorname{tr} \left[V \bar{D}^\alpha \left(e^{-tV} D_\alpha e^{tV} \right) - V \leftrightarrow \hat{V} \right]$$

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$$S_{pot} = \frac{2\pi i}{k} \int d^3x d^2\theta \epsilon_{AC} \epsilon^{BD} \operatorname{tr} (\mathcal{B}^A \mathcal{A}_D \mathcal{B}^C \mathcal{A}_B) + h.c.$$

\mathcal{B}^A \mathcal{A}_D

Chiral Scalars

V \hat{V}

Gauge Vectors

't hooft coupling: $\hat{\lambda} = \frac{N}{k}$

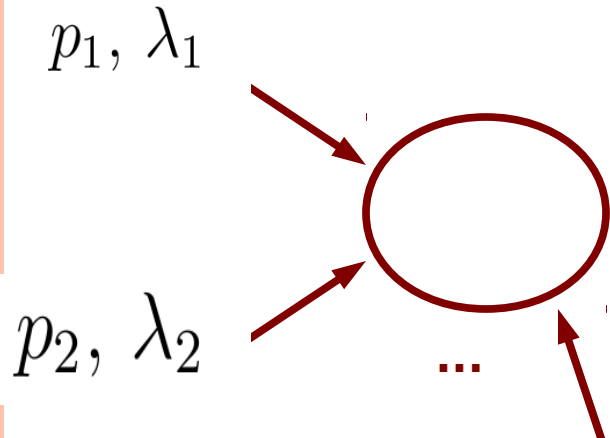


AdS/CFT Perspective



MHV Gluon-Scattering in N=4

Alday, Maldacena
0705.0303

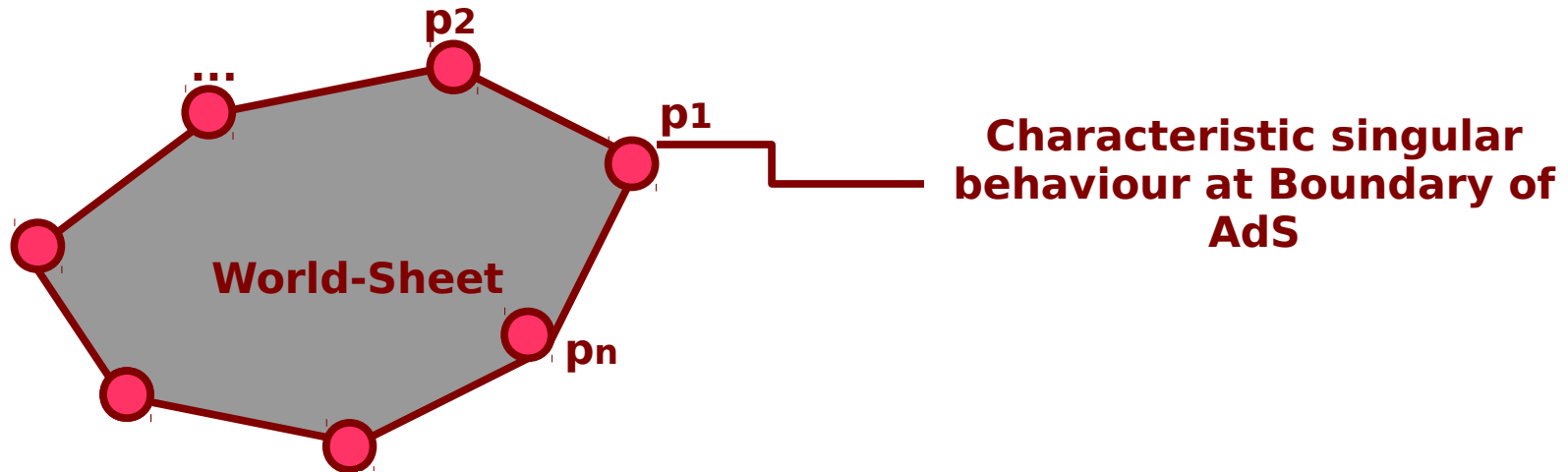


$$\frac{A}{A^{\text{tree}}} = e^{IR_{\text{div}} + f(\lambda)I(s_i) + \text{Rem}(s_i)}$$

$$f(\lambda) = \lambda + \mathcal{O}(\lambda^2)$$



Stringy dual picture



$$S_p = \frac{1}{2\pi} \int d^2\sigma \sqrt{-h} h^{ab} g_{\mu\nu} \partial_a \mathcal{X}^\mu \partial_b \mathcal{X}^\nu$$

$$g_{\mu\nu} d\mathcal{X}^\mu d\mathcal{X}^\nu = R^2 \left(\frac{dx_\mu dx^\mu + dr^2}{r^2} \right)$$



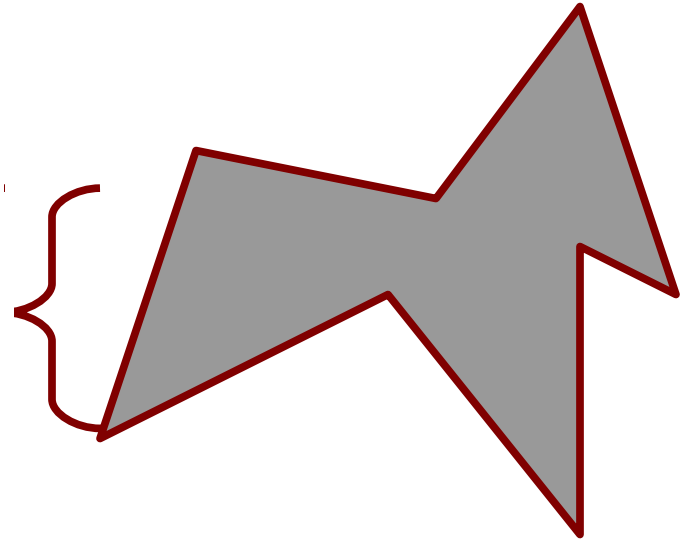
Stringy dual picture – T duality

$$\partial_a y^\mu = \frac{i}{r^2} \epsilon_{\alpha\beta} \partial_\beta x^\mu \quad z = \frac{R^2}{r}$$

$$p_i^\mu = \frac{\Delta_i y^\mu}{2\pi}$$

$$g_{\mu\nu} d\mathcal{Y}^\mu d\mathcal{Y}^\nu = R^2 \left(\frac{dy_\mu dy^\mu + dz^2}{z^2} \right)$$

$\Delta_i y^\mu$



$$\mathcal{A} \sim e^{-\frac{R^2}{2\pi} \text{Area}} = e^{-\frac{\sqrt{\lambda}}{2\pi} \text{Area}}$$



Wilson-Loops/Scattering amplitudes

$$W[c_n] = \frac{1}{N^2} \langle \text{Tr} \mathcal{P} \exp \left(ig \int_{C_n} dx_\mu A^\mu(x) \right) \rangle$$

Segments

$$\Delta_i y^\mu$$



Momenta

$$p_i^\mu$$

Light-like

$$(\Delta_i y)^2 = 0$$



On-Shell

$$p_i^2 = 0$$

Closed Polygon

$$\sum_i \Delta_i y^\mu = 0$$



Momentum conservation

$$\sum_i p_i^\mu = 0$$

Conformal Symmetry



Dual Conformal Symmetry



Wilson-Loops

Henn, Plefka,
Wiegandt
arXiv:1004.0226

$$\langle W_4 \rangle_{\text{ABJM}} = \frac{1}{2N} \left\{ \text{Tr} e^{i \int_{\gamma} A_{\mu} dz^{\mu}} + \text{Tr} e^{i \int_{\gamma} \hat{A}_{\mu} dz^{\mu}} \right\}$$

- **1-Loop Vanishing (n=4 analytic, n=6 numerical)**



- **2-Loop n=4**

(analytical) →

**Bicocca-
Celoria
1103.3675**

$$\langle W_4 \rangle_{\text{ABJM}}^{(2)} = \left(\frac{N}{k} \right)^2 \left[-\frac{(x_{13}^2 \mu'^2)^{2\epsilon}}{(2\epsilon)^2} - \frac{(x_{24}^2 \mu'^2)^{2\epsilon}}{(2\epsilon)^2} + \frac{1}{2} \ln^2 \left(\frac{x_{13}^2}{x_{24}^2} \right) + C \right]$$

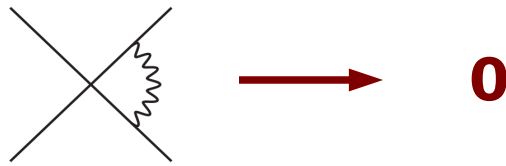
$$C = \frac{\pi^2}{2} + 2 \ln 2 + 5 \ln^2 2 - \frac{a_6}{4} \quad \mu'^2 = 8\pi e^{\gamma_E} \mu^2$$



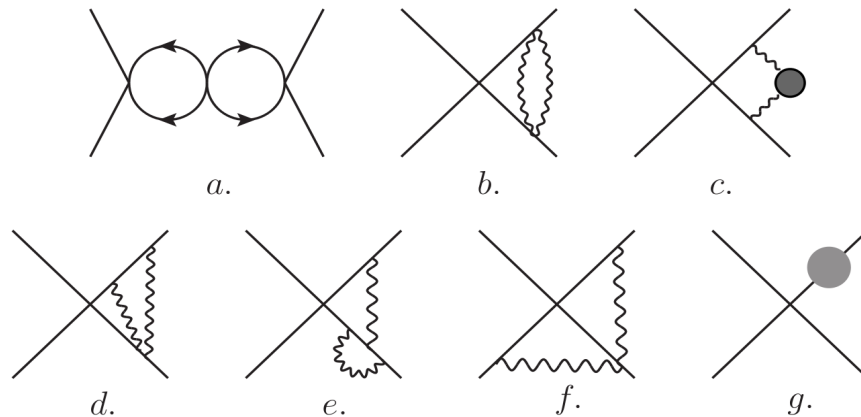
Results: Amplitudes

Bicocca-
Celoria
1107.3139

- **1-loop n=4 chiral amplitude** ($A^i B_j A^k B_l$)



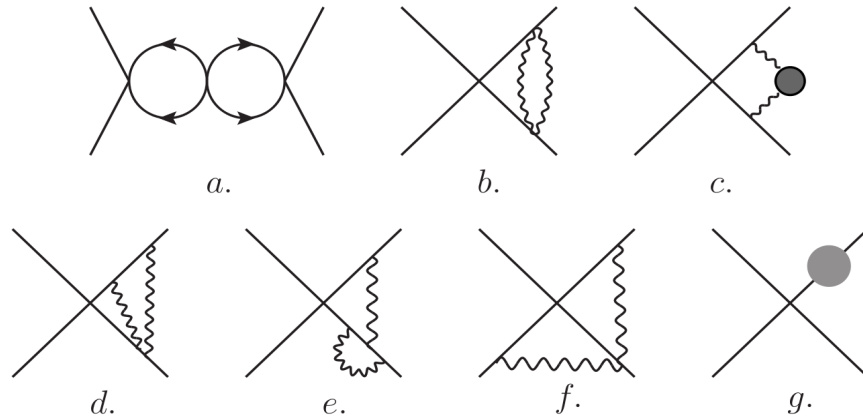
- **2-loop n=4 chiral amplitude**



Results: Amplitudes

Bicocca-
Cleoria
1107.3139

2-loop n=4 chiral amplitude $(A^i B_j A^k B_l)$



$$\mathcal{M}^{(2)} \equiv \frac{\mathcal{A}_4^{(2\text{loops})}}{\mathcal{A}_4^{\text{tree}}} = \lambda^2 \left[-\frac{(s/\mu'^2)^{-2\epsilon}}{(2\epsilon)^2} - \frac{(t/\mu'^2)^{-2\epsilon}}{(2\epsilon)^2} + \frac{1}{2} \log^2 \left(\frac{s}{t} \right) + K_1 + \mathcal{O}(\epsilon) \right]$$

$$K_1 = 4\zeta_2 + 3 \log^2 2$$

$$s = -x_{13}^2, t = -x_{24}^2$$

$$1/\mu_{WL}'^2 = \mu'^2$$

$$\left. \begin{array}{l} s = -x_{13}^2, t = -x_{24}^2 \\ 1/\mu_{WL}'^2 = \mu'^2 \end{array} \right\} \rightarrow \mathcal{M}^{(2)} \equiv \langle W_4 \rangle^{(2)} + \mathbf{Const}$$



Comments

- **Our Result permits to conjecture a BDS-like ansatz**

$$\frac{\mathcal{A}_4}{\mathcal{A}_4^{tree}} = e^{Div + \frac{f_{CS}(\lambda)}{8} \left(\ln^2\left(\frac{s}{t}\right) + \frac{4\pi^2}{3} \right) + C(\lambda)} \quad f_{CS}(\lambda) = \frac{1}{2} f_{\mathcal{N}=4}(\lambda) \Big|_{\frac{\sqrt{\lambda}}{4\pi} \rightarrow h(\lambda)}$$

- **We derived two-loop WL/Amplitude match without assumptions.**

Assuming dual conformal invariance:

**Chen, Huang,
arXiv:1107.2710**

- **We need to go to higher-loops and more scattered particles to check WL/Correlator duality for N=6**



More particles and superamplitudes



N=4 Superamplitudes

- Maximally helicity violating gluons

$$\mathcal{A}(+ + \dots + +) = 0, \quad \mathcal{A}(+ + \dots - \dots + +) = 0$$

$$\mathcal{A}(+ + \dots i^- \dots j^- \dots + +) = i \frac{\langle ij \rangle^4}{\langle 12 \rangle \langle 23 \rangle \dots \langle n1 \rangle}$$

$$p_{\alpha\dot{\alpha}} = p^\mu (\gamma_\mu)_{\alpha\dot{\alpha}} = \lambda_\alpha \bar{\lambda}_{\dot{\alpha}} \quad \langle 12 \rangle = \lambda_1^\alpha \lambda_{2\alpha} \quad [12] = \bar{\lambda}_1^{\dot{\alpha}} \bar{\lambda}_{2\dot{\alpha}}$$

(Parke-Taylor 1986)



N=4 Superamplitudes

- **MHV gluons and superpartners**

$$\begin{aligned}\Phi(p, \eta) = & G^+(p) + \eta^A \Gamma_A(p) + \frac{1}{2} \eta^A \eta^B S_{AB}(p) \\ & + \frac{1}{3!} \eta^A \eta^B \eta^C \epsilon_{ABCD} \bar{\Gamma}^D(p) + \frac{1}{4!} \eta^A \eta^B \eta^C \eta^D \epsilon_{ABCD} G^-(p)\end{aligned}$$

$$\mathcal{A}_n^{\text{MHV};0} = i \frac{\delta^{(4)} \left(\sum_{i=1}^n \lambda_i^\alpha \tilde{\lambda}_i^{\dot{\alpha}} \right) \delta^{(8)} \left(\sum_{i=1}^n \lambda_i^\alpha \eta_i^A \right)}{\langle 12 \rangle \langle 23 \rangle \dots \langle n1 \rangle}$$

Drummond et.al. 0807.1095

$$p_i^{\alpha\dot{\alpha}} = \lambda_i^\alpha \tilde{\lambda}_i^{\dot{\alpha}} = x_i^{\alpha\dot{\alpha}} - x_{i+1}^{\alpha\dot{\alpha}},$$

$$\lambda_i^\alpha \eta_i^A = \theta_i^{\alpha A} - \theta_{i+1}^{\alpha A}$$



N=4 Superamplitudes

- **Full Superamplitude**

$$\mathcal{A}_n = \mathcal{A}_n^{\text{MHV}} + \mathcal{A}_n^{\text{NMHV}} + \dots + \mathcal{A}_n^{\text{N}^{n-4}\text{MHV}}$$

$$\mathcal{A}_n^{\text{N}^k\text{MHV}} = i(2\pi)^4 \frac{\delta^{(4)}(p^{\dot{\alpha}\alpha}) \delta^{(8)}(q_\alpha^A)}{\langle 12 \rangle \langle 23 \rangle \dots \langle n1 \rangle} \hat{\mathcal{A}}_{n;k}(\lambda, \tilde{\lambda}, \eta; a)$$

$$\ln \hat{\mathcal{A}}_{n;0} = \ln W_n + O(\varepsilon), \quad W_n = \frac{1}{N_c} \left\langle \text{tr} P \exp \left(ig \int_{C_n} dx \cdot A(x) \right) \right\rangle$$



N=4 Superamplitudes

- **Super Wilson-loop**

$$\mathcal{W}_n = \frac{1}{N_c} \left\langle \text{tr } P \exp \left(ig \int_{\mathcal{C}_n} dx^\mu \mathcal{A}_\mu(x, \theta) + ig \int_{\mathcal{C}_n} d\theta^{\alpha A} \mathcal{F}_{\alpha A}(x, \theta) \right) \right\rangle$$

$$\mathcal{W}_n = \mathcal{W}_{n;0}(x_i; a) + \mathcal{W}_{n;1}(x_i, \theta_i^A; a) + \dots + \mathcal{W}_{n;n}(x_i, \theta_i^A; a)$$

$$\mathcal{W}_{n;k}(x, \theta; a) = a^k \hat{\mathcal{A}}_{n;k}(\lambda, \tilde{\lambda}, \eta; a)$$



ABJM Superamplitudes

- **On-Shell Superfields**

$$\begin{aligned}\Phi(\Lambda) &= \phi^4(\lambda) + \eta^A \psi_A(\lambda) + \frac{1}{2} \epsilon_{ABC} \eta^A \eta^B \phi^C(\lambda) + \frac{1}{3!} \eta^A \eta^B \eta^C \epsilon_{CBA} \psi_4(\lambda) \\ \bar{\Phi}(\Lambda) &= \bar{\psi}^4(\lambda) + \eta^A \bar{\phi}_A(\lambda) + \frac{1}{2} \epsilon_{ABC} \eta^A \eta^B \bar{\psi}^C(\lambda) + \frac{1}{3!} \eta^A \eta^B \eta^C \epsilon_{CBA} \bar{\phi}_4(\lambda)\end{aligned}$$

- **N=6 Superamplitudes**

$$\mathcal{A}_n(\Lambda_1, \dots, \Lambda_n) = \delta^{(3)}(P) \delta^{(6)}(Q) \sum_{k=1}^K f_{n,k} F_{n,k}$$



ABJM Superamplitudes

- **Four-Point superamplitude**

$$\mathcal{A}_{4;l} = \delta^3(P)\delta^6(Q) f_l(\lambda)$$

$$\mathcal{A}_{4;0} = \frac{\delta^3(P)\delta^6(Q)}{\langle 21 \rangle \langle 14 \rangle} \quad \mathcal{A}_{4;1} = 0$$

$$\mathcal{A}_{4;2} = \mathcal{A}_{4;0} \lambda^2 \left[-\frac{(s/\mu^2)^{-2\epsilon}}{(2\epsilon)^2} - \frac{(t/\mu^2)^{-2\epsilon}}{(2\epsilon)^2} + \frac{1}{2} \log^2(s/t) + C \right]$$



ABJM Superamplitudes

Bargheer et.al.
1003.6120

- **Six Point superamplitude**

$$\mathcal{A}_6^{(l)} = \delta^3(P) \delta^6(Q) [f^{+(l)}(\lambda) \delta^3(\alpha) + f^{-(l)}(\lambda) \delta^3(\beta)]$$

- **Constraints on** $f^{+(0)}(\lambda_i)$ and $f^{-(0)}(\lambda_i)$
- **Tree Level:**

$$\mathcal{A}_6^{(0)} = -\frac{\delta^3(P) \delta^6(Q)}{p_{123}^2} \left[\frac{(\epsilon_{ijk} \langle j k \rangle \eta_i^I + i \epsilon_{\bar{i}\bar{j}\bar{k}} \langle \bar{j} \bar{k} \rangle \eta_{\bar{i}}^I)^3}{(\langle 1 | p_{23} | 4 \rangle - i \langle 2 3 \rangle \langle 5 6 \rangle) (\langle 3 | p_{12} | 6 \rangle - i \langle 1 2 \rangle \langle 4 5 \rangle)} + \frac{(\epsilon_{ijk} \langle j k \rangle \eta_i^I - i \epsilon_{\bar{i}\bar{j}\bar{k}} \langle \bar{j} \bar{k} \rangle \eta_{\bar{i}}^I)^3}{(\langle 1 | p_{23} | 4 \rangle + i \langle 2 3 \rangle \langle 5 6 \rangle) (\langle 3 | p_{12} | 6 \rangle + i \langle 1 2 \rangle \langle 4 5 \rangle)} \right]$$

ABJM Superamplitudes

Bicocca- Celoria 1204.4407
Bargheer et.al 1204.4406

• 1-loop

$$\mathcal{A}_6^{(1)} = i \mathcal{C}(P) \frac{\delta^3(P) \delta^6(Q)}{p_{123}^2} \left[\frac{(\epsilon_{ijk} \langle j k \rangle \eta_i^I + i \epsilon_{\bar{i}\bar{j}\bar{k}} \langle \bar{j} \bar{k} \rangle \eta_{\bar{i}}^I)^3}{(\langle 1 | p_{23} | 4 \rangle - i \langle 2 3 \rangle \langle 5 6 \rangle) (\langle 3 | p_{12} | 6 \rangle - i \langle 1 2 \rangle \langle 4 5 \rangle)} - \frac{(\epsilon_{ijk} \langle j k \rangle \eta_i^I - i \epsilon_{\bar{i}\bar{j}\bar{k}} \langle \bar{j} \bar{k} \rangle \eta_{\bar{i}}^I)^3}{(\langle 1 | p_{23} | 4 \rangle + i \langle 2 3 \rangle \langle 5 6 \rangle) (\langle 3 | p_{12} | 6 \rangle + i \langle 1 2 \rangle \langle 4 5 \rangle)} \right]$$

$$\mathcal{C}(P) = \frac{\pi \lambda}{2} \left[\frac{\langle 12 \rangle}{\sqrt{\langle 12 \rangle^2}} \frac{\langle 34 \rangle}{\sqrt{\langle 34 \rangle^2}} \frac{\langle 56 \rangle}{\sqrt{\langle 56 \rangle^2}} + \frac{\langle 23 \rangle}{\sqrt{\langle 23 \rangle^2}} \frac{\langle 45 \rangle}{\sqrt{\langle 45 \rangle^2}} \frac{\langle 61 \rangle}{\sqrt{\langle 61 \rangle^2}} \right]$$

$$f^{+(1)} = -i \mathcal{C}(P) f^{+(0)} \quad \text{and} \quad f^{-(1)} = i \mathcal{C}(P) f^{-(0)}$$



ABJM Superamplitudes

- **The f functions at tree level → SCI, DSCI**

$$f^{\pm 1} \sim \mathcal{C}(P) f^{\pm 0}$$

- **3D** $\frac{d}{dx} Sg(x) = 2 \delta(x) \quad \rightarrow \quad \frac{\partial}{\partial \lambda_i^\alpha} Sg(\langle ij \rangle) = 2 \lambda_{j\alpha} \delta(\langle ij \rangle)$
- **4D** $\frac{\partial}{\partial \bar{z}} \frac{1}{z} = \pi \delta^2(z) \quad \rightarrow \quad \frac{\partial}{\partial \bar{\lambda}_i^{\dot{\alpha}}} \frac{1}{\langle ij \rangle} = \pi \bar{\lambda}_{j\dot{\alpha}} \delta^2(\langle ij \rangle)$



Conclusions

- **Four- and Six-points amplitudes in ABJM**
- **Similarities and differences with N=4 SYM**
- **4 points match with bosonic Wilson-Loop**
- **Found 3D-Analogous “holomorphic” anomaly**
- **Super Wilson-Loop is needed**



Thank you!

