Seven hints that early-time new physics alone is not sufficient to solve the Hubble tension

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Dipartimento di Fisica





The trouble

Overall trend:

- "early-time" model-dependent measurements prefer low *H*₀
- "late-time" direct measurements prefer high *H*₀



Review by Di Valentino et al., CQG 38 (2021) 153001

Hubble tension "no-go theorem"

Solving the tension while providing a good fit to BAO data and Hubble flow SNeIa data seems to require lowering r_s by $\approx 7\%$



Knox & Millea, PRD 101 (2020) 043533

This would seem to require early-time (pre-recombination) new physics!

Hubble tension "no-go theorem"

Late-time guard rails: BAO and Hubble flow SNeIa are very unforgiving!



Bernal, Verde & Riess, JCAP 1610 (2016) 019

...yet, we still haven't been able to construct a model truly fixing H_0 (empirically, early-Universe new physics only seems to get to $H_0 \sim 70 -$ with Planck CMB data and without including local H_0 priors)

Is early-time new physics the end of the story?

My sociological worry: "the Hubble tension calls for early-time new physics" may have been uncritically elevated to the mantra "the Hubble tension calls **exclusively** for early-time new physics"

Seven hints

- Ages of the oldest astrophysical objects
- Baryon Acoustic Oscillations r_d-H₀ degeneracy slope
- Cosmic chronometers
- \mathcal{D} escending trends observed in a wide range of low-z datasets
- \mathcal{E} arly integrated Sachs-Wolfe effect and its restrictions on early-time new physics
- ${m {\cal F}}$ ractional matter density (Ω_m) constraints from uncalibrated cosmic standards
- \mathcal{G} alaxy power spectrum r_{d} and k_{eq} -based determinations of H_0

Why seven? (Why not?) Miller's law - see Miller, Psychol. Rev. 63 (1956) 81





Opinion

SV. Universe 9 (2023) 393

Seven Hints that Early-Time New Physics Alone Is Not Sufficient to Solve the Hubble Tension

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Seven hints

a) Just reducing the sound horizon will introduce other problems:

• \mathcal{B} aryon Acoustic Oscillations r_d - H_0 degeneracy slope

b) Analyses more-or-less independent of pre-recombination physics – some residual amount of post-recombination physics seems to be required:

- \mathcal{F} ractional matter density (Ω_m) constraints from uncalibrated cosmic standards
- Cosmic chronometers
- \mathcal{A}_{ges} of the oldest astrophysical objects
- \mathcal{D} escending trends observed in a wide range of low-z datasets

c) Early-time guard rails – introducing pre-recombination new physics and maintaining the level of early-time consistency of Λ CDM is difficult:

- \mathcal{E} arly integrated Sachs-Wolfe effect and its restrictions on early-time new physics
- \mathcal{G} alaxy power spectrum r_d and k_{eq} -based determinations of H_0

Historically (1960s-1998) high-z OAOs provided the first hints for the existence of dark energy ($\Omega \neq 1$, $\Omega_{\Lambda} > 0$)

A 3.5-Gyr-old galaxy at redshift 1.55

James Dunlop, John Peacock, Hyron Spinrad, Arjun Dey, Raul Jimenez, Daniel Stern & Rogier Windhorst

Nature 381, 581-584 (1996) Cite this article

The observational case for a low-density Universe with a non-zero cosmological constant

J. P. Ostriker & Paul J. Steinhardt

Nature 377, 600-602 (1995) Cite this article

What can OAOs do for cosmology in the 2020s?

$$t_U(z) = \int_z^\infty rac{dz'}{(1+z')H(z')} \propto rac{1}{H_0}$$

- \bullet OAOs cannot be older than the Universe \rightarrow upper limit on H_0
- $t_U(z)$ integral insensitive to early-time cosmology
- $\bullet \rightarrow \textbf{late-time } \land \textbf{CDM consistency test independent of early times!}$
- Ages of astrophysical objects at z > 0 hard to estimate robustly $\underline{\mathbb{M}}$

Usefulness in relation to the H_0 tension:

- Contradiction between OAOs upper limit on H_0 and local H_0 measurements could indicate the need for non-standard late-time ($z \lesssim 10$) physics, or non-standard local physics
- Conclusions completely independent of pre-recombination physics

Age-redshift diagram up to $z\sim 8$



Assume Λ CDM at late times, constrain H_0 , Ω_m , and incubation time τ_{in} Prior for τ_{in} following Jiménez et al., JCAP 1903 (2019) 043; Valcin et al., JCAP 2012 (2020) 022



SV, Pacucci & Loeb, JHEAp 36 (2022) 27 $H_0 < 73.2 \ (95\% \ {
m C.L.})$

CAVEAT – If the OAOs ages are reliable, possible explanations are:

- ACDM may not be the end of the story at $z \lesssim 10$
- Nothing wrong with Λ CDM at $z \lesssim 10$, need new physics on local scales
- Just a boring 2σ fluke or systematics?

Hint 1: OAOs

Cosmic triangles: current cosmological data within a given model are over-constrained, look at quantities beyond H_0 and r_d (e.g. Ω_m , t_U)



If high $t_U(z = 0)$ measured reliably and with small uncertainties, models with high H_0 and standard low-z physics disfavored

Hint 2: BAO r_d - H_0 degeneracy slope

CMB and BAO constrain respectively:

$$\theta_{\star} \equiv \frac{r_{\star}}{D(z_{\star})}, \qquad \theta_d(z_{\rm obs}) \equiv \frac{r_d}{D(z_{\rm obs})}$$

Two sound horizons closely related:

$$r_d \approx 1.0184 r_{\star}$$

Given ω_m , imposing $\theta_{\star} = \text{const}$ and $\theta_d(z_{\text{obs}}) = \text{const}$ defines degeneracy line in r_d - H_0 plane with very different slopes for CMB and BAO (steeper for CMB, because $z_{\star} \gg z_{\text{obs}}$)

Q: what happens if H_0 is raised while only lowering r_d ...?

Hint 2: BAO r_d - H_0 degeneracy slope

A: quickly run into trouble with BAO and/or WL data if ω_m is unchanged, but even changing ω_m cannot bring agreement with both!



Jedamzik, Pogosian & Zhao, Commun. Phys. 4 (2021) 123

Hint 2: BAO r_d - H_0 degeneracy slope

Lower $\omega_m \implies$ tension with BAO data Higher $\omega_m \implies$ tension with WL data (worsen S_8 tension)



- **ACDM**
- Early Dark Energy (EDE) model I [14]
- EDE model II [14]
- evolving scalar fields | [15]
- evolving scalar fields II [15]
- a new EDE model [18]
- interacting neutrino cosmology I [19]
- interacting neutrino cosmology II [19]
- neutrino sector radiation [22]
- ultralight scalar decay [24]
- decaving dark matter (DM) I [25]
- decaying DM II [25]
- DM dark radiation interaction [26]
- swampland & fading DM [29]
- primordial magnetic fields | [30]
- primordial magnetic fields II [30]
- non-standard recombination I [31]
- non-standard recombination II [31]
- early recombination [33]

Jedamzik, Pogosian & Zhao, Commun. Phys. 4 (2021) 123 New physics which only reduces r_s is not enough!

Hint 3: Cosmic chronometers

Take two ensembles of galaxies that formed around the same time and are separated by a small redshift interval Δz around z_{eff} : Jiménez & Loeb, ApJ 573 (2002) 37

$$rac{dt}{dz} = -rac{1}{(1+z)H(z)} \implies H(z_{
m eff}) = -rac{1}{1+z_{
m eff}}rac{\Delta z}{\Delta t}$$

Use massive, early-time, passively-evolving galaxies (evolving on a much longer timescale than their age differences)



Hint 3: Cosmic chronometers

- CCs are completely (cosmological) model-independent
- CCs can be used to infer cosmological/non-local value of H_0
- Analyzing CC requires no assumptions on early-Universe physics
- Contradiction between CCs value of H_0 (assuming Λ CDM) and local H_0 measurements could indicate the need for non-standard late-time ($z \leq 2$) physics beyond Λ CDM, or non-standard local physics



Hint 3: Cosmic chronometers

Early-time-independent consistency test of Λ CDM: assuming Λ CDM holds at late times, from CC alone infer $H_0 = 67.5 \pm 3.0$ (note: no systematics!)

- Central value in excellent agreement with Planck
- Almost 2σ "tension" with local Cepheid-calibrated SNeIa H_0
- Preference for low H_0 not driven by any specific datapoint
- If uncertainties decrease and central value doesn't move, will need new late-time ($z\lesssim 2$) physics and/or new local physics



Mathematically speaking, dynamical models (e.g. ACDM) break down if values of (constant) fitting parameters pick up time dependence

Integrate 1st Friedmann equation with $w_{eff}(z)$ prescribed (in FLRW):

$$H_0 = H(z) \exp\left[-\frac{3}{2} \int_0^z dz' \frac{1 + w_{\text{eff}}(z')}{1 + z'}\right]$$

 $H(z) \sim \text{data}$ $w_{\text{eff}}(z')$: prescribed model H_0 : inferred fitting parameter (here mathematically integration constant)

If input $w_{eff}(z)$ and data "disagree", H_0 picks up z-dependence and "runs" at all redshifts Krishnan et al., PRD 103 (2021) 103509

If H_0 tension physical and at least some late-time new physics involved, *z*-evolution of H_0 at intermediate z ($0 < z < z_*$) inevitable!

- Has such a z-evolution already been observed in current data?
- Has it been observed in independent datasets with a common trend?
- Are there mundane explanations for its size and direction?





Combination of (binned) low-z datasets: megamaser distances, CCs, isotropic BAO, *Pantheon* SNeIa (r_d treated as free parameter)



 $\sim 2.1\sigma$ significance, slope consistent with H0LiCOW, by construction independent of early-Universe physics

(Binned) Pantheon SNela



Dainotti et al., ApJ 912 (2021) 150

 $\sim 2\sigma$ significance, well fit by $H(z) = H_0(1+z)^{-\alpha}$, with $\alpha \sim 10^{-2}$

Similar trends (descending H_0 and/or increasing Ω_m) observed in many different dataset combinations:

- PantheonPlus+SH0ES SNela Jia, Hu & Wang, A&A 674 (2023) A45
- PantheonPlus SNela Malekjani et al., arXiv:2301.12725
- Pantheon SNela Horstmann, Pietschke & Schwarz, A&A 668 (2022) A34
- CC+Pantheon SNela+QSOs Ó Colgáin *et al.*, arXiv:2206.11447
- QSOs Risaliti & Lusso, Nat. Astron. 3 (2019) 272
- $f\sigma_8$ measurements: S_8 increasing with redshift Adil et al., arXiv:2303.06928
- …and others!

Question: could this be expected even within ΛCDM ? (naïve guess: at high z lose sensitivity to DE, so expect $\Omega_m \uparrow \implies H_0 \downarrow$)

Forecast with mock data matching expected sensitivity of DESI



Ó Colgáin, Sheikh-Jabbari & Solomon, PDU 40 (2023) 101216

- Slight trend actually in the opposite direction
- Trend seen at z smaller than those where one expects to see it
- Expected size in any case much smaller than what is observed

Taken seriously, descending trends indicate need for new late-time physics

Hint 5: Early ISW effect

Around recombination: Universe not fully matter dominated \implies residual decay of gravitational potentials \implies eISW effect sources anisotropies

$$\Theta = \int_{0}^{\eta_{0}} d\eta \left[\underbrace{\propto g(\Theta_{0} + \Psi)}_{\text{Sachs-Wolfe}} + \underbrace{\propto gv_{b} \frac{d}{d\eta}}_{\text{Doppler}} + \underbrace{\propto e^{-\tau}(\dot{\Psi} - \dot{\Phi})}_{\text{ISW}} + \underbrace{\propto (g\Pi + [\ddot{g}\Pi])}_{\text{Polarization}} \right] j_{\ell}(k\Delta\eta)$$
$$\Theta_{\ell}^{\text{ISW}}(k) = \underbrace{\int_{0}^{\eta_{m}} d\eta \, e^{-\tau} \left(\dot{\Psi} - \dot{\Phi}\right) j_{\ell}(k\Delta\eta)}_{\text{early ISW}} + \underbrace{\int_{\eta_{m}}^{\eta_{0}} d\eta \, e^{-\tau} \left(\dot{\Psi} - \dot{\Phi}\right) j_{\ell}(k\Delta\eta)}_{\text{late ISW}}$$

(A substantial amount of) New physics increasing H(z) around z_{eq}/z_{\star} should leave an imprint on the eISW effect!

Why is there no clear sign of early-time new physics in CMB data alone?

Hint 5: Early ISW effect

$$\Theta_{\ell}^{\mathsf{elSW}}(k) = A_{\mathsf{elSW}} \int_{0}^{\eta_{m}} d\eta \, e^{-\tau} \left(\dot{\Psi} - \dot{\Phi} \right) j_{\ell}(k\Delta\eta)$$

Consistency check: within ΛCDM , data consistent with $A_{elSW} = 1$?

Yes! $A_{eISW} = 0.988 \pm 0.027$ (other parameters stable to within $\lesssim 0.3\sigma$)





SV, PRD 104 (2021) 063524

Hint 5: Early ISW effect (EDE application)

High H_0 EDE fit to CMB requires increased $\omega_c \rightarrow$ worsens S_8 tension? Hill *et al.*, PRD 102 (2020) 043507; Ivanov *et al.*, PRD 102 (2020) 103502; D'Amico *et al.*, JCAP 2105 (2021) 072; see partial rebuttals in: Murgia *et al.*, PRD 103 (2021) 063502; Smith *et al.*, PRD 103 (2021) 123542

Editors' Suggestion

Early dark energy does not restore cosmological concordance

J. Colin Hill, Evan McDonough, Michael W. Toomey, and Stephon Alexander Phys. Rev. D **102**, 043507 – Published 5 August 2020

Parameter	ΛCDM	EDE (high ω_c)	EDE (low ω_c)
$100\omega_b$	2.253	2.253	2.253
ω_c	0.1177	0.1322	0.1177
$H_0 [{ m km/s/Mpc}]$	68.21	72.19	72.19
τ	0.085	0.072	0.072
$\ln(10^{10}A_s)$	3.0983	3.0978	3.0978
n_s	0.9686	0.9889	0.9889
$f_{\rm EDE}$	-	0.122	0.122
$\log_{10} z_c$	-	3.562	3.562
θ_i	-	2.83	2.83
n	-	3	3



Hint 5: Early ISW effect (EDE application)

Let's extract only eISW contribution to temperature anisotropies...

Low ω_c

High ω_c



Almost 20% eISW excess!

No more than \lesssim 3-5% eISW excess

Problem generic to models increasing pre-recombination H(z)

Hint 6: Ω_m constraints from uncalibrated cosmic standards

Beneficial to look at joint H_0 - Ω_m constraints rather than just projected H_0 constraints Lin, Mack & Hou, ApJL 904 (2020) L22 Can we determine Ω_m :

- At a level competitive with the CMB model-dependent value?
- Free from early-Universe assumptions (as with BAO+SNela)?

 $\Delta r H_0$ small & insensitive to early-Universe physics Lin, Chen & Mack, ApJ 920 (2021) 159

$$\Delta r H_0 \equiv (r_d - r_{\star}) H_0 = \int_{z_d}^{z_{\star}} dz \, \frac{c_s(z)}{E(z)} \qquad (z_d - z_{\star}) \sim 30$$

Combine θ_{\star} (CMB) and θ_d (BAO) in almost early Universe-independent way, with long lever arm to constrain Ω_m at level competitive with CMB: Early Universe Physics Insensitive Uncalibrated Cosmic Standards (UCS)

Hint 6: Ω_m constraints from uncalibrated cosmic standards

Data: θ_{\star} (*Planck*+ACT+WMAP), θ_d (eBOSS), CMB priors on z_{\star} and Δz_s , BBN prior on $\Omega_b h^2$ Parameters: Ω_m , \mathcal{M} , $r_d H_0$, h (weak dependence)



Lin, Chen & Mack, ApJ 920 (2021) 159

Purely geometrical, early Universe-independent value: $\Omega_m = 0.302 \pm 0.008$ For comparison $\Omega_m = 0.310 \pm 0.006$ in ACDM using full CMB information

Hint 6: Ω_m constraints from uncalibrated cosmic standards

Constraints not exactly along Ω_m direction, weak Ω_m -h degeneracy

$$\left(\frac{\Omega_m}{0.3}\right) \left(\frac{h}{0.7}\right)^{-0.08} = 1.0060 \pm 0.0258$$

Combine UCS with several early Universe-independent late-time, non-local measurements to infer H_0 in an early Universe-independent way

Methods	$H_0 ({\rm kms^{-1}Mpc^{-1}})$			n-σ from R21	
UCS and individual nonlocal observation	Without θ_{cmb}	With θ_{cmb}	Without θ_{cmb}	With θ_{cmb}	
Cosmic chronometers					
Current public data	69.1 ± 1.7	$\textbf{68.8} \pm \textbf{1.6}$	1.9σ	2.1σ	
Extra systematic	69.4 ± 2.3	$\textbf{69.2} \pm \textbf{2.1}$	1.4σ	1.6σ	
Extra systematic, conservative	69.3 ± 3.4	68.9 ± 3.3	1.1σ	1.2σ	
γ-ray optical depth	66.2 ± 3.5	66.1 ± 3.4	1.9σ	2.0σ	
Cosmic age					
$t_{\rm U} = 13.5 \pm 0.27 \text{ Gyr}$	70.2 ± 1.7	69.8 ± 1.5	1.4σ	1.7σ	
$t_U = 13.5 \pm 0.33$ Gyr	70.3 ± 2.1	69.8 ± 1.9	1.2σ	1.5σ	
CMBlens+DES+BBN	68.8 ± 2.4	$\textbf{68.6} \pm \textbf{2.0}$	1.6σ	1.9σ	
UCS and joint nonlocal observations ^a					
All nonlocal observations	69.1 ± 1.5	$\textbf{68.8} \pm \textbf{1.3}$	2.0σ	2.4σ	
Nonlocal observations without cosmic age	68.3 ± 1.9	68.1 ± 1.6	2.1 <i>σ</i>	2.5σ	
Nonlocal observations without LSS	69.1 ± 1.6	$\textbf{68.8} \pm \textbf{1.5}$	2.0σ	2.2σ	

Lin, Chen & Mack, ApJ 920 (2021) 159

Residual $\approx 2\sigma$ tension can have nothing to do with early-Universe physics: need late-time new physics and/or local new physics (systematics very unlikely given consistency among independent probes)

Hint 7: r_s - and k_{eq} -based constraints on H_0 from P(k)

Two scales in P(k), both standard rulers



Credits: Oliver Philcox

- $k_{\rm eq} = \sqrt{2\Omega_m H_0 z_{\rm eq}}$ (if no extra components with significant pressure support) sets peak and overall shape ($z_{\rm eq} \approx 3500$)
- r_d sets BAO frequency ($z_{\star} \approx 1100$)

Both can be used to infer H_0 : in the presence of a substantial amount of early-time new physics, no reason two values should agree!

Hint 7: r_s - and k_{eq} -based constraints on H_0 from P(k)

Can analyze P(k) data removing (most) r_d information (effectively marginalizing over r_d), similarly CMB lensing also sensitive to k_{eq}



Philcox et al., PRD 106 (2022) 063530

 $H_0 = 64.8^{+2.2}_{-2.5}$ (only k_{eq} info): agrees with Λ CDM r_d -based value of H_0 , disfavors significant amount of early-time new physics?

Hint 7: r_s - and k_{eq} -based constraints on H_0 from P(k)

Caveats:

- Current error bars still guite large
- r_d vs r_d-marginalized comparison model-dependent...
- ...and (Ω_m) prior-dependent



Smith. Poulin & Simon. arXiv:2208.12992

Future data should improve discriminatory power, but for the moment this is a consistency test of ΛCDM at best

Where to from here? Some scattered thoughts

- Empirically: early-time physics only seems to reach $H_0 \sim 70$ (no external priors)
- Idea: combine early-time and late-time (both non-local) and local new physics?
- Direction of late-time physics: lower $d_A(z)$ at z > 0 (phantom/interacting DE?)
- CMB+BAO/SNeIa actually can tolerate w as low as ~ -1.07 , H_0 responds as $\Delta H_0 \sim -20(1 + w)$, so this can help as much as $\Delta H_0 \sim 1.5$ sv, PRD 102 (2020) 023518
- If there is also some local new physics lowering local H₀, maybe don't need non-local H₀ to go all the way up to ~ 74 after all? (two can meet halfway)
- Early-time new physics probably still need to do the lion's share of the job...
- Early+late: can two models decouple, both "push" non-local *H*₀ up separately, combining their tension-solving virtues "in phase" / "constructively"?



Objection: wouldn't this violate Occam's razor?

My opinion \downarrow



Credits: Wiley Miller

Nature is under no obligation to look simple to us!

Phantom dark energy? Really?

The state of the dark energy equation of state circa 2023

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Escamilla et al., arXiv:2307.14802 (submitted to PRD)



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Where to from here? What about the S_8 tension?

Early times: a relatively successful early-time model (EDE and variants, $\Delta m_{e},...$) <u>Late times</u>: scattering-type new physics (at 1st order does not affect background but only perturbations) involving DM and/or DE \rightarrow decouple S_8 -solving effects from H_0 -solving ones, combine the two constructively?

Example: DE-baryon scattering

$$\dot{\theta}_{b} = -\mathcal{H}\theta_{b} + c_{s}^{2}k^{2}\delta_{b} + \frac{4\rho_{\gamma}}{3\rho_{b}}a_{ne}\sigma_{T}(\theta_{\gamma} - \theta_{b}) + (1 + w_{x})\frac{\rho_{x}}{\rho_{b}}a_{ne}\sigma_{xb}(\theta_{x} - \theta_{b})$$

$$\dot{\theta}_{x} = -\mathcal{H}(1 - 3c_{s}^{2})\theta_{x} + \frac{c_{s}^{2}k^{2}}{1 + w_{x}}\delta_{x} + a_{ne}\sigma_{xb}(\theta_{b} - \theta_{x})$$





Dark scattering (and S_8)

Lots of room for dark scattering



Concrete recent example explicitly discussing the S_8 tension

Sigma-8 tension is a drag

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Poulin et al., PRD 107 (2023) 123538

Simpson, PRD 82 (2010) 083505

Possible underlying Lagrangian: "Type 3" coupled DE models (scalar field derivative coupling to velocity)

Models of dark matter coupled to dark energy

A. Pourtsidou, C. Skordis, and E. J. Copeland Phys. Rev. D **88**, 083505 – Published 9 October 2013

See classification presented in Pourtsidou, Skordis & Copeland, PRD 88 (2013) 083505

Where to from here?

Pictorial representation of what I think could be a promising scenario



Conclusions

- Seve(ral/n) hints that early-time new physics alone cannot solve the Hubble tension
- My opinion: will probably need a combination of early-time and late-time (both non-local) and local new physics, non-local and local H_0 might not need to meet at \sim 74 but halfway
- "Decoupling" of early- and late-time tension-solving effects may also help S₈: I find scattering-type models particularly promising





Opinion

Seven Hints that Early-Time New Physics Alone Is Not Sufficient to Solve the Hubble Tension

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