Resolving Cosmological Tensions with Unparticles

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Outline

- Introduction
- Beyond the scalar field paradigm
- Banks-Zaks Cosmology
- Emergent Unparticles DE Model and the tensions
- Phenomenological fluid approach (WiP)
- Conclusions

COSMOLOGY CRASH SLIDE

Consider the FLRW metric with no spatial curvature. Define the critical density

$$\rho_c = \frac{3H_0^2}{8\pi G}, \quad \Omega_{0i} = \frac{\rho}{\rho_c}, \quad w_i = \frac{p_i}{\rho_i}, \quad EOS$$
$$H(z)^2 = H_0^2 \sum \Omega_{0i} (1+z)^{3+3w_i}, \quad \dot{H} = -H_0^2 \sum \frac{3+3w_i}{2} \Omega_{0i} (1+z)^{3+3w_i}$$

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For z<3400 (equality of matter and radiation), equations simplify further:

$$H(z)^{2} \simeq H_{0}^{2} \left[\Omega_{m0} (1+z)^{3} + 1 - \Omega_{m0} \right], \quad \dot{H} \simeq -\frac{3}{2} H_{0}^{2} \Omega_{m0} (1+z)^{3}$$

• LCDM model ~1% accuracy (CMB) $H_0 = 67.4 \pm 0.5$, $\Omega_{m0} = 0.315 \pm 0.007$ $H_0, \Omega_b, \Omega_c, \tau, \sigma_8, n_s; \quad \Omega_{m0} = \Omega_b + \Omega_c; \quad \Omega_{m0} + \Omega_{\Lambda 0} = 1$

Fundamental Problems in Cosmology

- Big Bang Singularity Artymowski, IBD, Kumar JCAP 2019 +WIP.
- Early Universe Inflation/Bounce? Which model? Connection with SM Artymowski, IBD, Thattarampilly JCAP 2020, IBD, Thattarampilly 2308.00256
- Nature of present acceleration CC? Dark Energy? Modified gravity? Which model?
- Nature of Dark Matter

Scalar Fields!

- The Hubble Tension
- The LSS (S_8) Tension

 H_0 is the most important cosmological measurement and is relevant for fundamental physics as well.



 $H_0 = 67.4 \pm 0.3$ km/sec/Mpc $H_0 = 67.4 \pm 1.2$ km/sec/Mpc

74.8+-3.1 76.5+-4.0

THE LSS S8 TENSION

- Measurements of matter fluctuations on large scales is given by $S_8=\sigma_8\sqrt{\frac{\Omega_{m0}}{0.3}}$

Most attempts to reduce one tension result in increasing the other!

• σ_8 are the linear matter fluctuations smoothed over 8h Mpc

2-3 SIGMA!

- Ω_{m0} is the relative matter energy density.
- CMB (Planck measurement) $S_8 = 0.834 \pm 0.016$
- (Also KiDs ...) DES $S_8 = 0.776 \pm 0.017$

SCALAR FIELDS IN COSMOLOGY

- Daily practice in theoretical physics T(t,x)
- A single DOF, with a flat potential everyone can do that!
- Abundant in String Theory and extensions of the SM
- Fine-tuned models? organizing principle?
- Field Theory is much richer confinement, strong interactions, topological defects, conformal symmetry...
- The Swampland Conjecture in QG potentials are steep (see however, IBD, PRD 2018 "Draining the Swampland")

$$\Delta \phi \lesssim 1, \quad \frac{V'}{V} \gtrsim 1 \quad OR \quad \frac{V''}{V} \gtrsim 1$$



PROBLEMS WITH PRESENT ACCELERATION (CC/DE)

- CC is the simplest parametrization of the observed acceleration
- The CC is related to zero point vacuum fluctuations of fields and respects Lorentz inv., for a cutoff M, we expect $\rho_{\Lambda} \sim M^4$
- For any fundamental energy scale M, we find a CC much smaller than expected *The Cosmological Constant Problem* $\rho_{\Lambda}^{obs.} \sim 10^{-10} erg/cm^3$, $M_{QCD} \sim 10^{36} erg/cm^3$, $M_{EW} \sim 10^{47} erg/cm^3$, $M_{pl} \sim 10^{110} erg/cm^3$.

WHY NOW?

- Energy Density of matter ~ a⁻³, radiation ~ a⁻⁴, CC ~a⁰.
- Completely different scaling which is why we had a radiation dominated universe followed by after domination.

. Why do we have today
$${\rho_m\over
ho_\Lambda}\sim 1?$$

 Requires fine-tuned initial conditions in the early universe. Especially for scalar fields.

BEYOND SCALAR FIELDS IN COSMOLOGY

- Daily practice in theoretical physics T(t,x)
- A single DOF, with a flat potential everyone can do that! Emergent single DOF with an equation of state - everyone can do that!
- Abundant in String Theory and extensions of the SM
- Fine-tuned models? organizing principle?
- Field Theory is much richer confinement, strong interactions, topological defects, conformal symmetry...
- The Swampland Conjecture in QG potentials are steep



BANKS-ZAKS COSMOLOGY

- Consider a sector with conformal symmetry (SU(3) with N_f massless fermions) weakly coupled to the SM (suppressed by $\Lambda_{\mathcal{U}}$).
- At high temperature ($T \gg \Lambda_{\mathcal{U}}$) conformal symmetry is restored and the sector behaves like radiation.
- At low temperature the coupling to SM breaks the symmetry "unparticles" with anomalous scaling T^{δ} .

BANKS-ZAKS COSMOLOGY -UNPARTICLES

• Theory with conformal symmetry, slightly displaced from its conformal fixed point. β function vanishes in the conformal limit.

Very general, based on dimensional analysis, any broken CFT $heta_{\mu}^{\mu}$ ~

$$\mathcal{L}\left(\frac{T}{\Lambda_{\mathcal{U}}}\right)^{44}$$

- The thermal average gives: $\theta^{\mu}_{\mu} =
ho - 3p \propto T^{\delta}$

•
$$\rho = \sigma T^4 + BT^{4+\delta}$$

$$p = \frac{1}{3}\sigma T^4 + \frac{B}{\delta + 3}T^{4+\delta}$$

. The equation of state is not (nearly) constant anymore $w \equiv \frac{p}{\rho} = \frac{1}{3} \frac{\sigma + \frac{3B}{3+\delta}T^{\delta}}{\sigma + BT^{\delta}}$

BANKS-ZAKS COSMOLOGY -CONSEQUENCES W(T)

- Naturally behaves as different fluids at different epochs
- Can temporarily violate the Null Energy Condition (NEC) - No Big Bang singularity without QG or noncanonical Lagrangians
- Can have a limiting temperature an effective CC!

UNPARTICLES AS DARK ENERGY-UDE

@ high T, w=1/3 -radiation. @ low T, limiting temperature T_c.
 Dim-less temperature y=T/T_c

$$-3 < \delta < 0, \quad B < 0 \Rightarrow \quad T > T_c = \left[\frac{4(\delta+3)}{3(\delta+4)} \left(-\frac{\sigma}{B}\right)\right]^{\frac{1}{\delta}}$$

- The dynamical evolution starts from high T and asymptotes to T_{c.}
- Unparticles start as radiation and as they asymptote to T_c they behave as a CC.
- Deviations can only come from higher loop corrections of the beta function.

Artymowski, IBD, Kumar PRD 2021, 2022, IBD, Kumar 2023



Ν



UDE- PROS.

- No fine-tuning of initial conditions. Radiation and CC behavior are predictions.
- No "Swampland conjectures", no scalar fields, no modified gravity.
- B is fixed by present day DE density. We are very close to the critical temperature $y_0-1 \lesssim 10^{-4}$





UDE - PREDICTIONS

• Special redshift dependence of w: $w_u \simeq -1 + 4(\delta + 4)(y_0 - 1)(1 + z)^3$



3/4

- Contributes to N_{eff}, consistency condition, current limits $\Delta N_{eff} \lesssim 0.19$:

$$w_u(z) \simeq -1 + 0.58 (1+z)^3 \left(-1 - \frac{4}{\delta}\right)^{1/4} \left[\frac{\Omega_{r0}}{\Omega_{u0}} \Delta N_{eff}\right]$$

• Perturbation observables (γ , $f\sigma_8$) - as LCDM to 0.1%

$$\delta'_{u} = -(1+w_{u})\left(\theta_{u}-3\Phi'\right) - \frac{a'}{a}\left(\frac{\delta p_{u}}{\delta \rho_{u}}-w_{u}\right)\delta_{u},$$

$$\theta'_{u} = -\frac{a'}{a}\left(1-3w_{u}\right)\theta_{u} - \frac{w'_{u}}{1+w_{u}}\theta_{u} + \frac{\frac{\delta p_{u}}{\delta \rho_{u}}}{1+w_{u}}k^{2}\delta_{u} + k^{2}\Phi$$



UDE AND LCDM - LIKELIHOOD ANALYSIS

- Perform likelihood analysis of UDE and compare to LCDM
- Consider various data sets, each time removing one that is causing the tension.
- For UDE fix $\delta = -3$, data is insensitive to the exact value.
- Flat priors for the different parameters $H_0, \Omega_b, \Omega_c, \tau, \sigma_8, n_s; \quad y_0 = 1 + 10^{x_0}, \quad x_0 \in [-6, -3]$
- Model mostly changes the Early Universe z>1.

Planck Only

SN Only



Planck+Pantheon+SHoES



Hubble is shifted owards SN value! _		LCDM
	H ₀	$68.00^{+0.92}_{-0.70}(68.46)$
S ₈ is reduced	σ_8	$0.8082^{+0.0079}_{-0.0088}(0.8045)$
wards DES value!	S_8	$0.817^{+0.017}_{-0.023}(0.8045)$
	Ω_m	$0.3069^{+0.0088}_{-0.013}(0.3005)$

UDE

 $71.29^{+0.96}_{-1.1}(71.18)$

 $0.8132^{+0.0093}_{-0.0078}\,(0.813)$

 $0.803^{+0.018}_{-0.014}\,(0.806)$

 $0.2926^{+0.0087}_{-0.0077}\,(0.295)$

 $\Delta \chi^2 \simeq -11 !$

to

Planck+Lensing+BAO+DES+Pantheon+SHoES



Hubble Towards	e is shifted s SN value! to	S ₈ is reduced owards DES value!	
$\Delta \chi^2 = -7.24!$			
	LCDM	UDE	
H_0	$68.28^{+0.71}_{-0.31} (68.47)$	$70.87^{+0.61}_{-0.79}(70.76)$	
σ_8	$0.8071 \pm 0.0059 (0.8069)$	$0.8142 \pm 0.0073 (0.8140)$	
S_8	$0.8111^{+0.0080}_{-0.015} \left(0.8075 ight)$	$0.808^{+0.010}_{-0.0078}(0.8093)$	
Ω_m	$0.3030^{+0.0038}_{-0.0091} (0.3004)$	$0.2956^{+0.0049}_{-0.0041}(0.2965)$	

Improvement in the likelihood persists for all other combinations of data sets by at least $|\Delta\chi^2| > 2.1$

Both tensions are reduced to less than 1 STD w.r.t SHoES and DES!



Clear improvement in concordance and likelihood!

PHENOMENOLOGICAL FLUID WITH "TRACKER MECHANISM"

- Could unparticle behavior be more robust?
- Any DE model needs a "tracker mechanism" DE "tracks" the matter or radiation to avoid fine tuning.
- Possibility of measuring the spatial curvature. *Leonard et al. 2016* -need theoretical insights.
- Could some of the tensions be just parametrization issues?
- Parametrize a relatively sudden transition.

PHENO. FLUID DE

- Theoretical priors:
- 1) Fulfill the Null Energy Condition w>=-1
- 2) Today w~-1
- 3) At early times w=1/3 (or 0) for tracker mechanism
- Speed of sound $c_s^2 = 1$ or 0 or open $0 \le c_s^2 \le 1$ and adiabatic speed of sound:

$$c_a^2 = w_{DE}(1 - \frac{\pi}{4}) + \frac{\pi}{12}$$

• at transition redshift.

$$w_{DE}(a) = -1 + \frac{4/3}{1 + (a/a_t)^n}$$

- Slight reduction in H0 and S8 tension. H0~69, S8~0.81.
- Less statistically significant.
- $\Delta \chi^2 \simeq -2.7$
- n>3 preferred
- In all cases, zt~30.



BEYOND SCALAR FIELD SUMMARY

- Within UDE cosmological concordance is largely restored for both Hubble and S₈ tension
- Pheno. Fluid approach not as successful. Still $z_t \sim 30$
- Going beyond weakly coupled scalar fields opens a new model space with different problems and opportunities.
- Generic arguments about conformal symmetry and dim. analysis.
- Useful for fundamental problems Big Bang singularity, Swampland, CC... and for practical ones - Hubble tension, S₈ tension, N_{eff}, ...
- Highly predictive, consistency condition detected within a decade or bound consistency approaching LCDM.
- Future directions interactions with CMB, N_{eff}, growth of fluctuations

BACKUP: HUBBLE TENSION REVISITED

- SNIa: do not assume a model except isotropic redshift.
- Consider only low redshift SNIa, z << 1. $d_L(z) = \frac{1+z}{H_0} \int_0^z \frac{dz'}{\sqrt{\Omega_m (1+z')^3 + 1 \Omega_m}} \approx \frac{z}{H_0} \left[1 + \left(1 \frac{1+z}{2} + \frac{1+$
- <u>Measure</u> Hubble $H_0^{SN,obs.} = 73 \pm 1 \, km/sec/Mpc$
- Use Hubble and high redshift z~1 for matter density Ω_{m0} , CC etc.
- **<u>CMB</u>**: Take all possible data. Assume a model (like LCDM)
- Infer the model parameters from a likelihood analysis



BACKUP 2: EARLY UNIVERSE

- Consider unparticles+fluid. $H(z)^{2} = H_{0}^{2} \sum_{i} \Omega_{0i}(1+z)^{3+3w_{i}}, \quad \dot{H} = -H_{0}^{2} \sum_{i} \frac{3+3w_{i}}{2} \Omega_{0i}(1+z)^{3+3w_{i}}$
- Violates NEC near the Bounce.
- New stable solutions- de Sitter Bounce, standard Bounce, cyclic universe.
- Analysis of the different phases



Calculation of the primordial spectra and stability of the cyclic/bounce scenarios.

Artymowski, IBD, Kumar JCAP 2019 +WiP



Unparticles +Fluid



Unparticles only