



A Minimal Explanation of Flavour Anomalies:

$R_{D^{(*)}}$, $R_{K^{(*)}}$, $(g-2)_\mu$ and CAA

Sokratis Trifinopoulos

Corfu

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[D. Marzocca,
ST] 2104.05730

B-Physics Anomalies: neutral-current

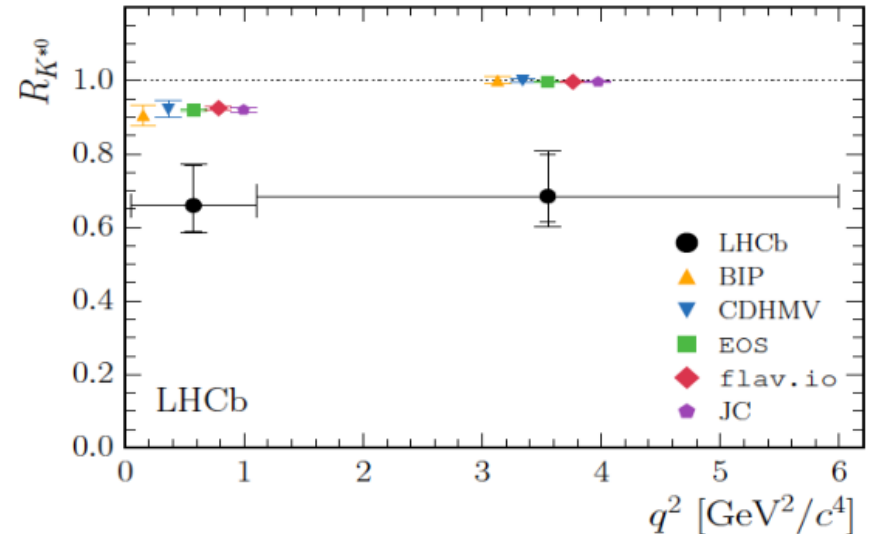
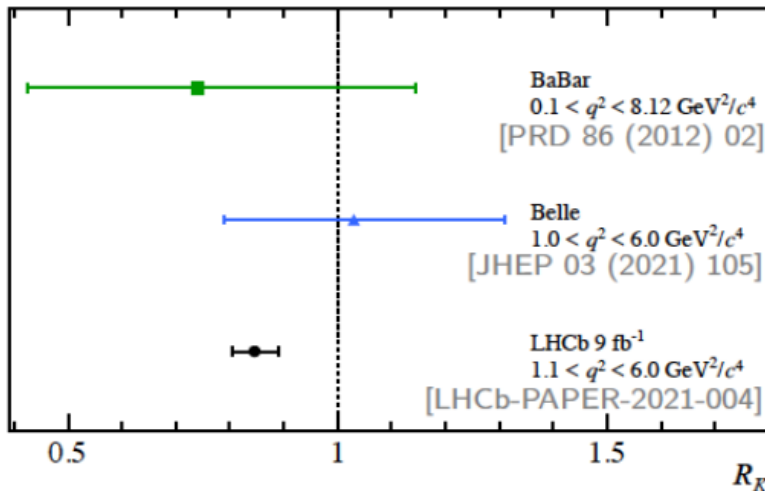
THE GOOD

- $b \rightarrow s\ell\ell$: Deficit of the charged-current transition in μ vs. e

$$R_{K^{(*)}} = \frac{\mathcal{B}(B \rightarrow K^{(*)} \mu \bar{\mu})}{\mathcal{B}(B \rightarrow K^{(*)} e \bar{e})}$$



+ $B_s \rightarrow \mu^+ \mu^-$, angular distributions of $B \rightarrow K^* \mu^+ \mu^-$, various BRs





***B*-Physics Anomalies: neutral-current**

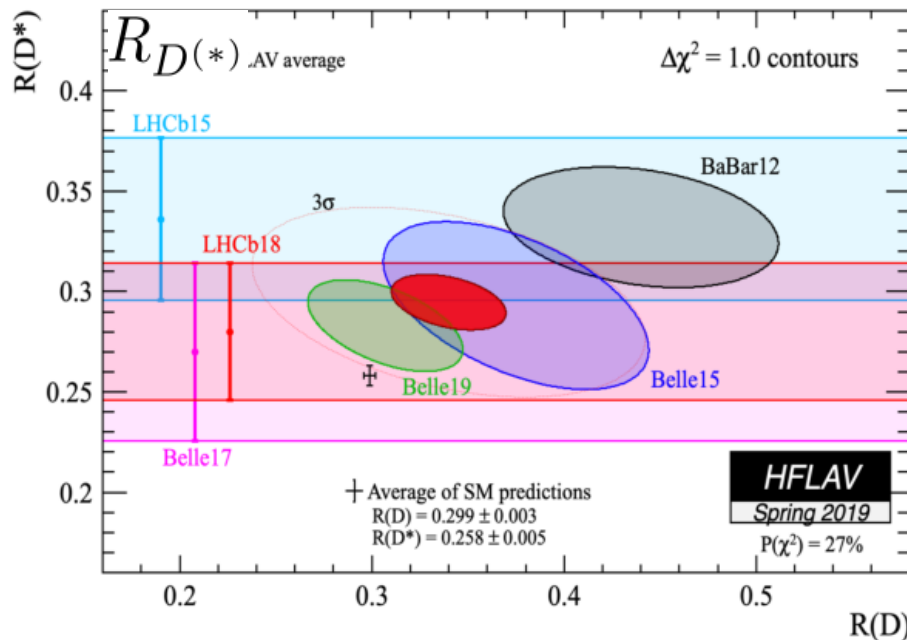
- R_K survived the latest update by the LHCb even after employing the full dataset of Run I and II independently at 3.1σ ! [LHCb] 2103.11769
- Fits obtained by varying one or two relevant NP WCs at a time yield pulls at a staggering 5σ level! [Altmannshofer et al] 2103.13370
A recent analysis evaluated the global significant using a conservative & unbiased method and found a **global 3.9σ** (still remarkable!)
[Lancieri et al] 2104.05631
- The anomalies implies violation of **L**epton **F**lavor **U**niversality (**LFUV**), which is an accidental symmetry of the SM.
↪ (if confirmed) definite NP signal of short-distance origin

B-Physics Anomalies: charged-current

THE (MAYBE NOT SO) GOOD

- $b \rightarrow c\tau\nu$: Enhancement of the charged-current transition in τ vs. ℓ

$$R_{D^{(*)}} = \frac{\mathcal{B}(B \rightarrow D^{(*)}\tau\bar{\nu})}{\mathcal{B}(B \rightarrow D^{(*)}\ell\bar{\nu})} \approx 3\sigma$$

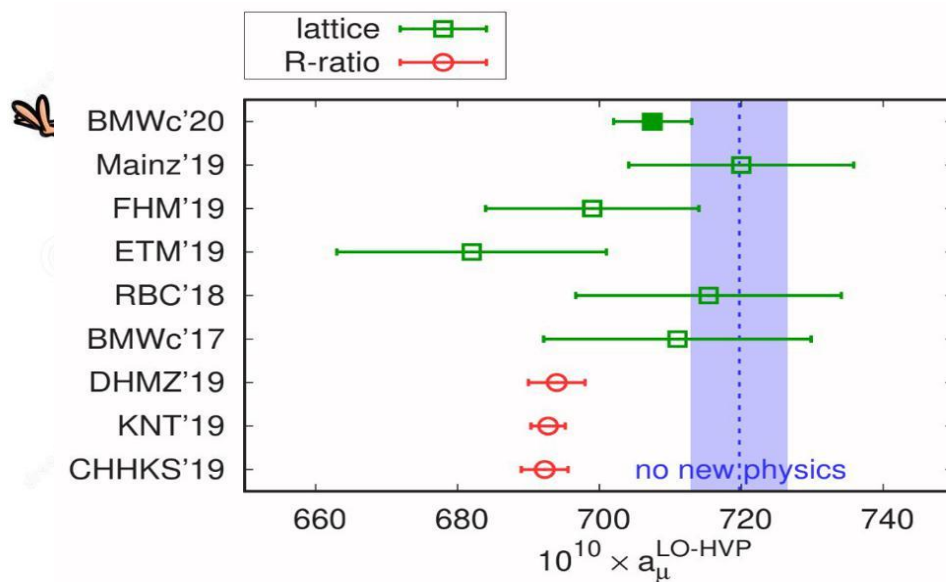


- The results from BaBar (2012) drive mainly the CL (without it the pull is 2.2σ)
- R_D agrees with the SM in the latest Belle update, but it has never been measured by LHCb → in anticipation of an analysis?!

Anomalous magnetic moment of the muon

THE (NOT SO) BAD

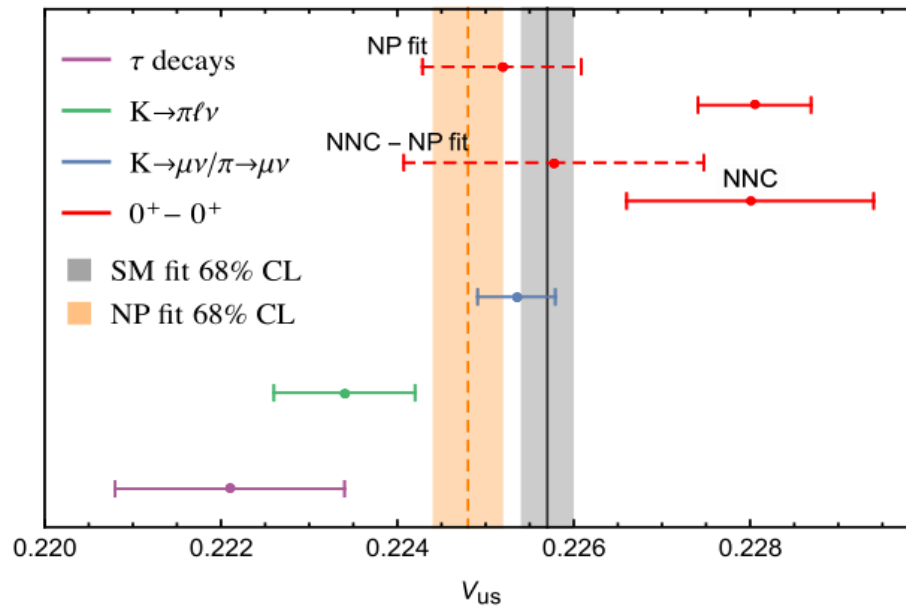
- $(g-2)_\mu$: Deviation in the muon anomalous magnetic moment
- The longstanding BNL result was recently updated by FNAL, confirming the previous trend and increasing the significance to 4.2σ !!! [Fermilab] 2104.03281
- **BUT**, recent lattice calculations yield SM prediction (almost) in agreement with experiment.



The Cabbibo-Angle Anomaly

AND THE UGLY

- **CAA:** Discrepancies between different determinations of V_{us} :



[Coutinho et al]
1912.08823

- Depending on the input from nuclear β decays we obtain 3-5 σ !

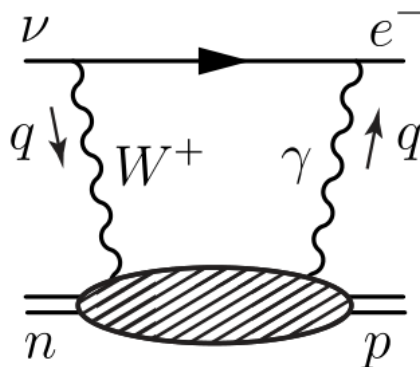
[Grossman et al] 1911.07821

Excursion: nuclear β decays

- Superaligned β decays are Fermi transitions ($S=0, \Delta J=0$) between isobaric analogue states with no parity change ($\Delta\pi=0$).

- Masterformula: $|V_{ud}|^2 = \frac{2984.432(3) \text{ s}}{\mathcal{F}t(1 + \Delta_R^V)}$
 - Universal contribution: EW corrections
 - “corrected” half-life: factoring out nucleus-dependent parts

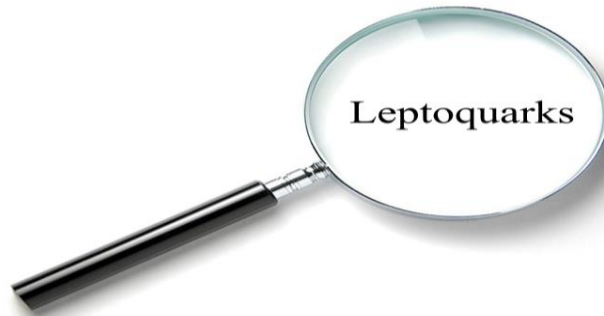
- What's **new**?



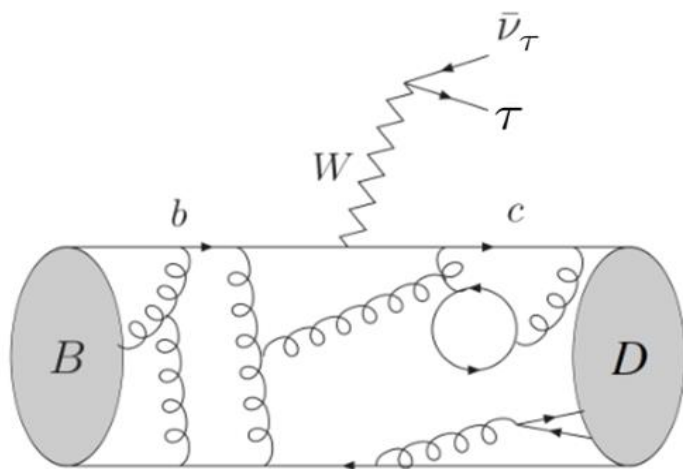
The uncertainty of Δ_R is dominated by the hadronic contribution to the $W\gamma$ box.
New analyses using dispersion relations and hybrid lattice QCD result in a shift of $|V_{ud}|$.

[Seng et al] 1812.03352,
2107.14708

[Czarnecki et al] 1907.06737

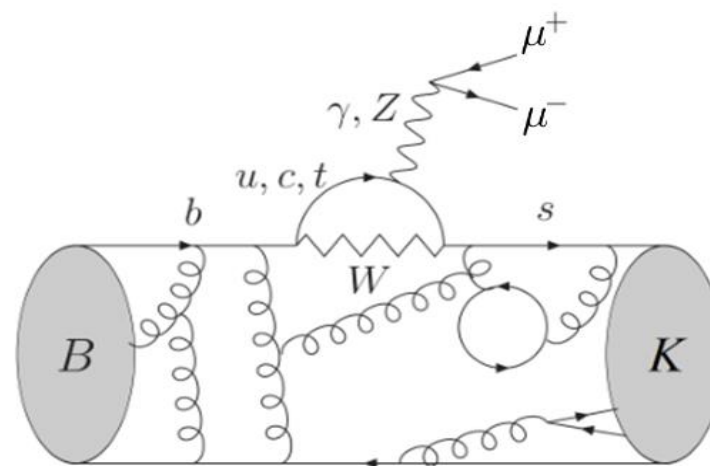


Following the trails...



- Tree-level process
- Mild CKM suppression

(presumably) tree-level NP effect induced by semi-leptonic 4-fermion operators \leftrightarrow Leptoquarks!



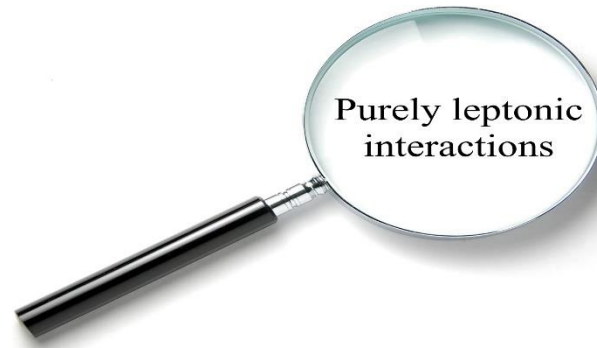
- Loop-level process
- CKM & GIM suppression

NP in (short-distance) WC of:

$$O_9^\mu = \bar{b}_L \gamma^\mu s_L \bar{\mu} \gamma_\mu \mu$$

$$O_{10}^\mu = \bar{b}_L \gamma^\mu s_L \bar{\mu} \gamma_\mu \gamma_5 \mu$$

BONUS: $(g-2)_\mu$ with LQ chiral enhanced loops



Following the trails...

- The tension between V_{us}^{CKM} and V_{us}^{β} can be **alleviated** by introducing a constructive interference in $\mu \rightarrow e\nu\nu$. In particular, one obtains:

$$V_{us}^{\beta} \equiv \sqrt{1 - (V_{ud}^{\beta})^2 - |V_{ub}|^2} \simeq V_{us}^{\text{CKM}} \left[1 + \left(\frac{V_{ud}^{\text{CKM}}}{V_{us}^{\text{CKM}}} \right)^2 \delta(\mu \rightarrow e\nu\nu) \right]$$

where $\delta(\ell \rightarrow \ell'\nu\nu) \equiv \mathcal{A}(\ell \rightarrow \ell'\nu\nu)_{\text{NP}} / \mathcal{A}(\ell \rightarrow \ell'\nu\nu)_{\text{SM}}$. [Crivellin et al] 2012.09845

- **Link** between CAA and the B -Physics anomalies:

$$\frac{g_{\tau}}{g_e} = \left| \frac{1 + \delta(\tau \rightarrow \mu\nu\nu)}{1 + \delta(\mu \rightarrow e\nu\nu)} \right| \approx 2\sigma$$

[Feruglio et al] 1606.00524

Interplay with LQ anticipated already by the EFT analysis

The Model

- **NP particle content:** $S_1 \sim (\bar{\mathbf{3}}, \mathbf{1})_{1/3}$, $\phi^+ \sim (\mathbf{1}, \mathbf{1})_1$.
- The SM Lagrangian is augmented by the following Yukawa-type terms

$$\mathcal{L}_{S1+\phi} = \frac{1}{2} \lambda_{\alpha\beta} \bar{\ell}_\alpha^c \epsilon \ell_\beta \phi^+ + \lambda_{i\alpha}^{1L} \bar{q}_i^c \epsilon \ell_\alpha S_1 + \lambda_{i\alpha}^{1R} \bar{u}_i^c e_\alpha S_1 + \text{h.c.}$$

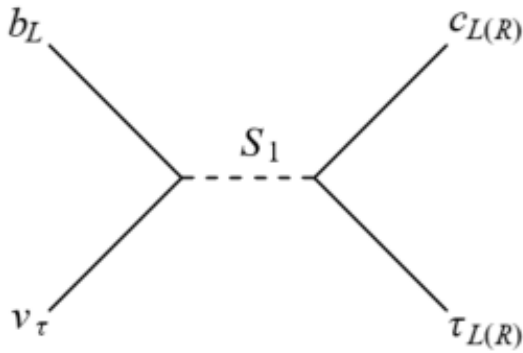
- We employ the following couplings:

$$\lambda \approx \begin{pmatrix} 0 & \lambda_{12} & \cancel{\lambda_{13}} \\ -\lambda_{12} & 0 & \lambda_{23} \\ \cancel{\lambda_{13}} & -\lambda_{23} & 0 \end{pmatrix}, \quad \lambda^{1L} \approx \begin{pmatrix} 0 & 0 & 0 \\ 0 & \lambda_{s\mu}^{1L} & \lambda_{s\tau}^{1L} \\ 0 & \lambda_{b\mu}^{1L} & \lambda_{b\tau}^{1L} \end{pmatrix}, \quad \lambda^{1R} \approx \begin{pmatrix} 0 & 0 & 0 \\ 0 & \lambda_{c\mu}^{1R} & \lambda_{c\tau}^{1R} \\ 0 & \lambda_{t\mu}^{1R} & \lambda_{t\tau}^{1R} \end{pmatrix}.$$

$\mu \rightarrow e\gamma$

$\tau \rightarrow \mu\gamma$

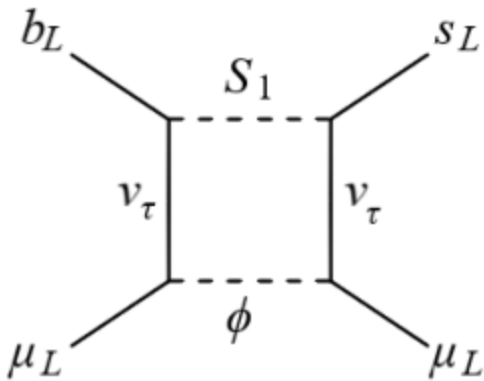
Solution to the Anomalies (diagrams)



➤ $b \rightarrow c \tau \nu$:

$$R_D \approx 0.299 - 0.235 \frac{\lambda_{b\tau}^{1L} \lambda_{c\tau}^{1R}}{m_1^2} (1 + 0.05 \log m_1^2)$$

$$R_{D^*} \approx 0.258 - 0.088 \frac{\lambda_{b\tau}^{1L} \lambda_{c\tau}^{1R}}{m_1^2} (1 + 0.02 \log m_1^2)$$



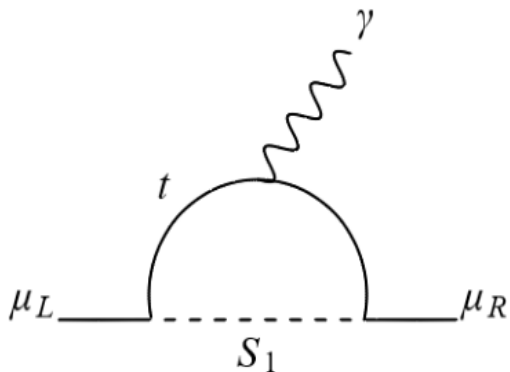
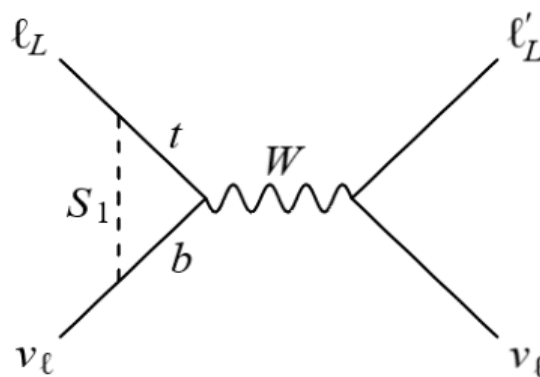
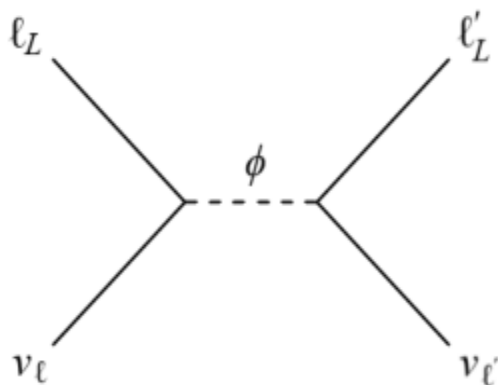
➤ $b \rightarrow s \ell \ell$: $\Delta C_{9,10}^\mu = (C_{LR} \pm C_{LL}) / (2\mathcal{N}_{sb})$

$$C_{LL} \approx -\lambda_{b\tau}^{1L} \lambda_{s\tau}^{1L*} \left(\frac{|\lambda_{b\mu}^{1L}|^2}{64\pi^2 M_1^2} + \frac{|\lambda_{\mu\tau}|^2 \log M_1^2 / M_\phi^2}{64\pi^2 (M_\phi^2 - M_1^2)} \right)$$

$$C_{LR} \approx -\frac{|\lambda_{c\mu}^{1R}|^2 \lambda_{b\tau}^{1L} \lambda_{s\tau}^{1L*}}{64\pi^2 M_1^2}$$

Solution to the Anomalies (diagrams)

➤ CAA: $\delta(\mu \rightarrow e\nu\nu) \approx \frac{v^2|\lambda_{12}|^2}{4M_\phi^2} + \frac{3m_t^2|\lambda_{b\mu}^{1L}|^2}{32\pi^2 M_1^2} \left(\frac{1}{2} - \log \frac{M_1^2}{m_t^2} \right)$

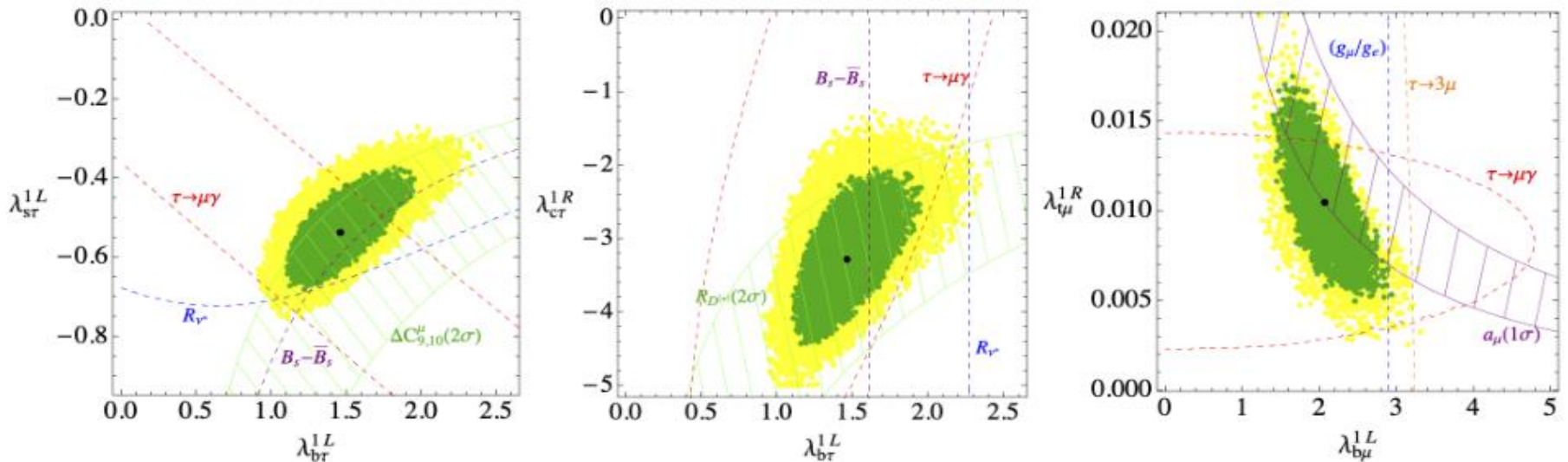


➤ $(g-2)_\mu$:

$$\Delta a_\mu \approx \frac{m_\mu m_t \lambda_{b\mu}^{1L} \lambda_{t\mu}^{1R}}{4\pi^2 M_1^2} \left(\log M_1^2 / m_t^2 - \frac{7}{4} \right)$$

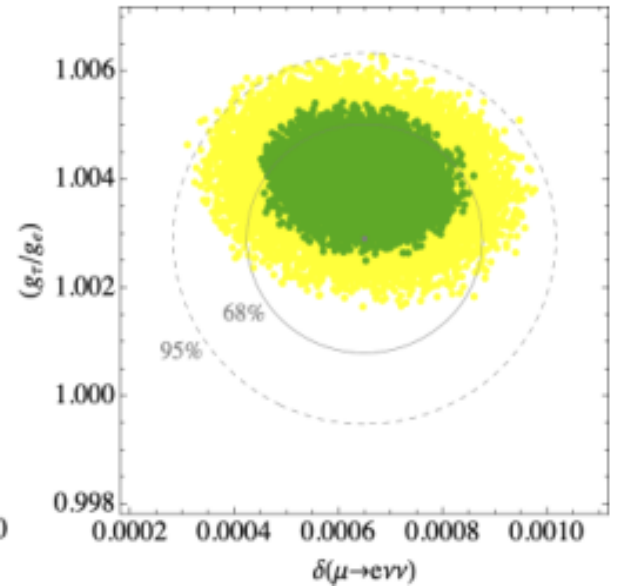
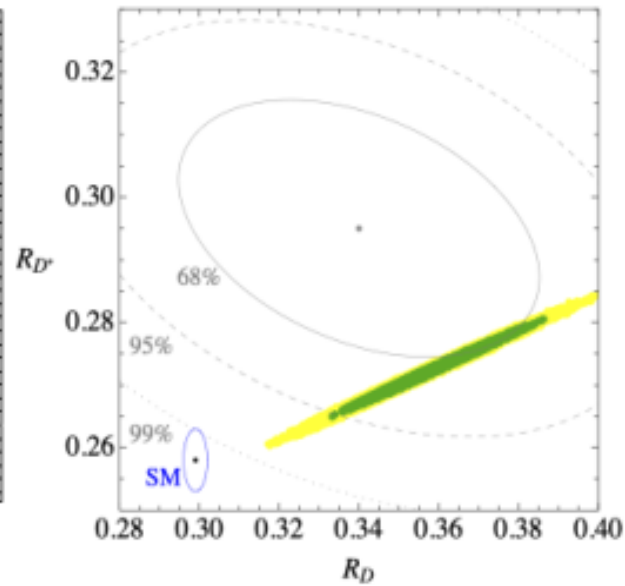
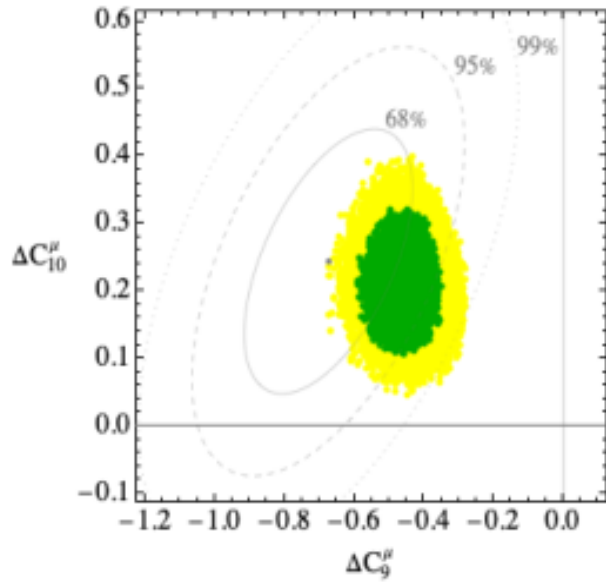
Model parameter space (at 2σ exclusion)

- We perform a χ^2 minimization with fixed masses, $M_1 = M_\phi = 5.5\text{TeV}$.



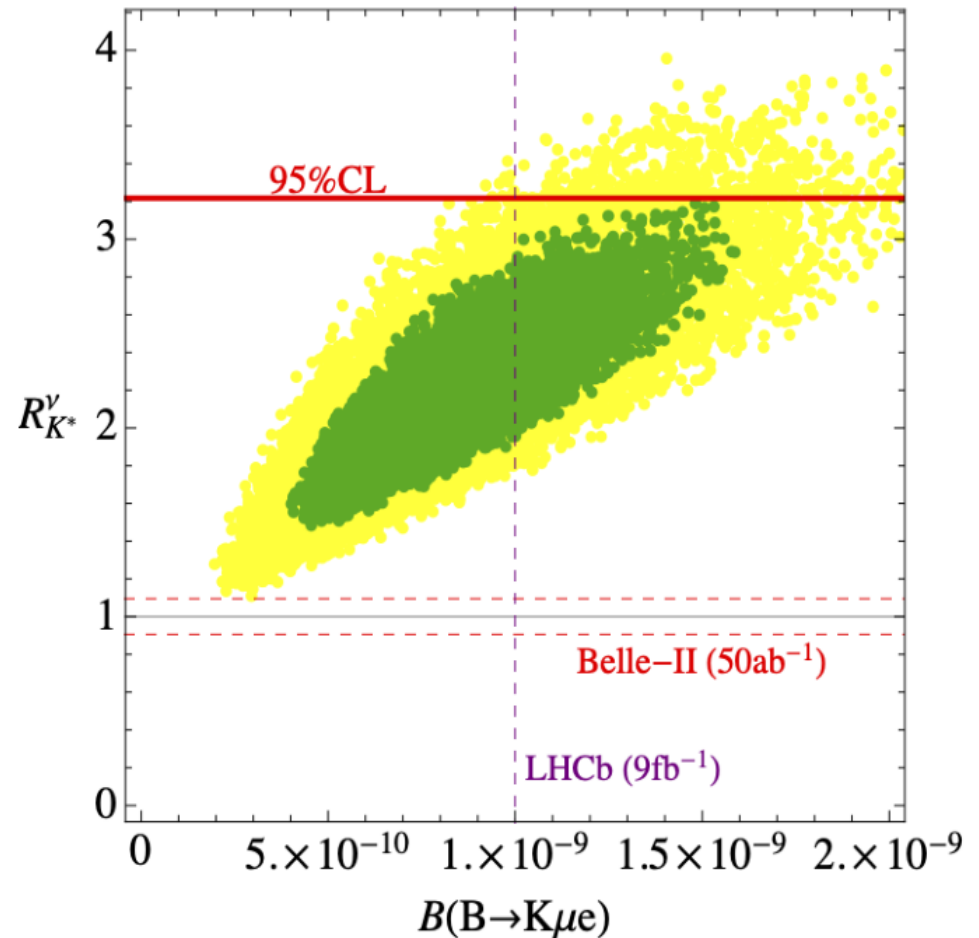
Solution to the Anomalies (fit)

$$\chi_{\text{SM}}^2 - \chi_{\text{best-fit}}^2 = 82$$



Future Prospects I

- The preferred region is **close** to B_s -mixing and $B \rightarrow K^{(*)}\nu\nu$ present bounds due to S_1 contributions.
- Unique **signatures**: $B \rightarrow K\mu e$ and $B_s \rightarrow \mu e$ (similar box that fixes C_9 & C_{10})
- **No** effects to $b \rightarrow s\tau\tau$.





Future Prospects II

- **LFU** ratios in τ decays (φ^+ at tree-level + S_1 at one-loop) :
$$\frac{g_\tau}{g_e} \frac{g_\tau}{g_\mu} \text{ and } \frac{g_\mu}{g_e} .$$
- **LFV** τ decays: $\tau \rightarrow \mu\gamma$, $\tau \rightarrow 3\mu$ and $\tau \rightarrow \mu ee$ (via S_1 at one-loop)
 $\tau \rightarrow e\gamma$, $\tau \rightarrow 3e$ and $\tau \rightarrow e\mu\mu$ (via φ^+ at one-loop) } Belle-II & LHCb
- Masses **too large** for direct searches at LHC (maybe effects in high-energy tails of Drell-Yan). **Ideal machines**: FCC-hh for S_1 and a muon collider for φ^+ .



Model-independent Prospects

- **Final verdict** on *B-Physics* anomalies: Belle-II with $\sim 20\text{ab}^{-1}$
- $(g-2)_\mu$ runs 2&3 already completed \rightarrow statistical error will soon shrink further!
- **CAA:** (*theory*) advance understanding of the EW radiative corrections

(*experiment*) improvement on $K \rightarrow \pi \ell \nu$ (LHCb, NA62, TREK etc.) &
 $K \rightarrow \mu \nu / K \rightarrow e \nu$ (J-PARC E36)

Additionally, complementary observables can become competitive e.g.
pion β decays $\pi^\pm \rightarrow \pi^0 e^\pm \nu_e$, neutron lifetime etc.

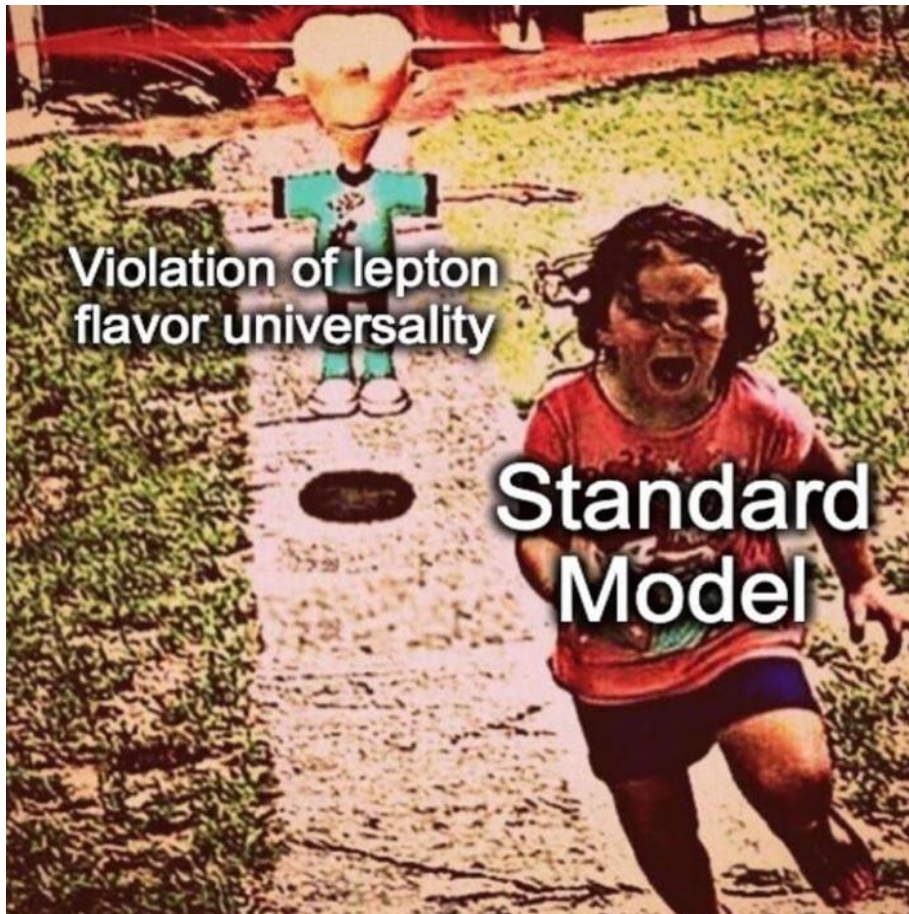


Conclusions

- **Model-building invitation:** Good reasons to consider all four major ($>3\sigma$) deviations from the SM in Flavour physics under the same NP (LFUV) framework.
- In this work we provide the **first combined solution**, respecting the rich set of phenomenological constraints.
- The particle content of the model ($S_1 + \varphi^+$) is (most probably) **the most minimal** one that can achieve this. Additionally, it features **both** LQ and purely leptonic interactions (=fresh new *market* direction vs the LQ *monopolies*).
- Many new and old interesting **signatures** (or model **killers**..)



Thank you!!!!



QUESTIONS ???

@largememecollider



Backup slides

Constraints:

Observable	Experimental value
R_D	0.34 ± 0.029 [56]
R_{D^*}	0.295 ± 0.013 [56]
ΔC_9^μ	-0.675 ± 0.16 [20]
ΔC_{10}^μ	0.244 ± 0.13 [20]
Δa_μ	$(2.51 \pm 0.59) \times 10^{-9}$ [27, 28]
$\delta(\mu \rightarrow e\nu\nu)$	$(6.5 \pm 1.5) \times 10^{-4}$ [41]
$R_D^{\mu/e}$	0.978 ± 0.035 [57, 58]
$\mathcal{B}(B_c \rightarrow \tau\nu)$	< 0.1 [59]
$R_{K^{(*)}}^\nu$	< 2.7 [60]
$C_{B_s}^1$	$< 2.01 \times 10^{-5} \text{ TeV}^{-2}$ [61]
$ \text{Re}(C_D^1) $	$< 3.57 \times 10^{-7} \text{ TeV}^{-2}$ [61]
$ \text{Im}(C_D^1) $	$< 2.23 \times 10^{-8} \text{ TeV}^{-2}$ [61]
$\frac{g_\tau}{g_e}$	1.0058 ± 0.0030 [56]
$\frac{g_\tau}{g_\mu}$	1.0022 ± 0.0030 [56]
$\frac{g_\mu}{g_e}$	1.0036 ± 0.0028 [56]
$\delta g_{\tau L}^Z$	$(-0.11 \pm 0.61) \times 10^{-3}$ [62]
$\delta g_{\tau R}^Z$	$(0.66 \pm 0.65) \times 10^{-3}$ [62]
$\delta g_{\mu L}^Z$	$(0.3 \pm 1.1) \times 10^{-3}$ [62]
$\delta g_{\mu R}^Z$	$(0.2 \pm 1.3) \times 10^{-3}$ [62]
$\mathcal{B}(\tau \rightarrow \mu\gamma)$	$< 4.4 \times 10^{-8}$ [63]
$\mathcal{B}(\tau \rightarrow 3\mu)$	$< 2.1 \times 10^{-8}$ [63]

➤ Best-fit point:

$$\lambda_{e\mu} = 1.35, \quad \lambda_{\mu\tau} = 3.17, \\ \lambda_{b\tau}^{1L} = 1.46, \quad \lambda_{s\tau}^{1L} = -0.54, \quad \lambda_{b\mu}^{1L} = 2.07, \\ \lambda_{c\tau}^{1R} = -3.28, \quad \lambda_{t\mu}^{1R} = 0.01, \quad \lambda_{c\mu}^{1R} = 2.35.$$



(Wild) speculations

- The LQ S_1 and φ^+ share the quantum numbers of the right-handed sbottom and stau, respectively.
- The NP couplings λ and λ^{IL} would correspond then to the couplings λ and λ' of the **RPV superpotential**. The right-handed couplings λ^{IR} can only be generated from non-holomorphic RPV terms. [Trifinopoulos] 1904.12940
- **Hint** towards a RPV scenario with lighter 3rd generation superpartners?
- λ and λ^{IL} **follow** the $U(2)$ flavour symmetry that respects $SU(5)$ unification. λ^{IR} is more **peculiar**, but since it originates from a hidden sector maybe there is a flavourful dynamical way to reproduce the observed structure. [Csaki et al] 1309.5957