



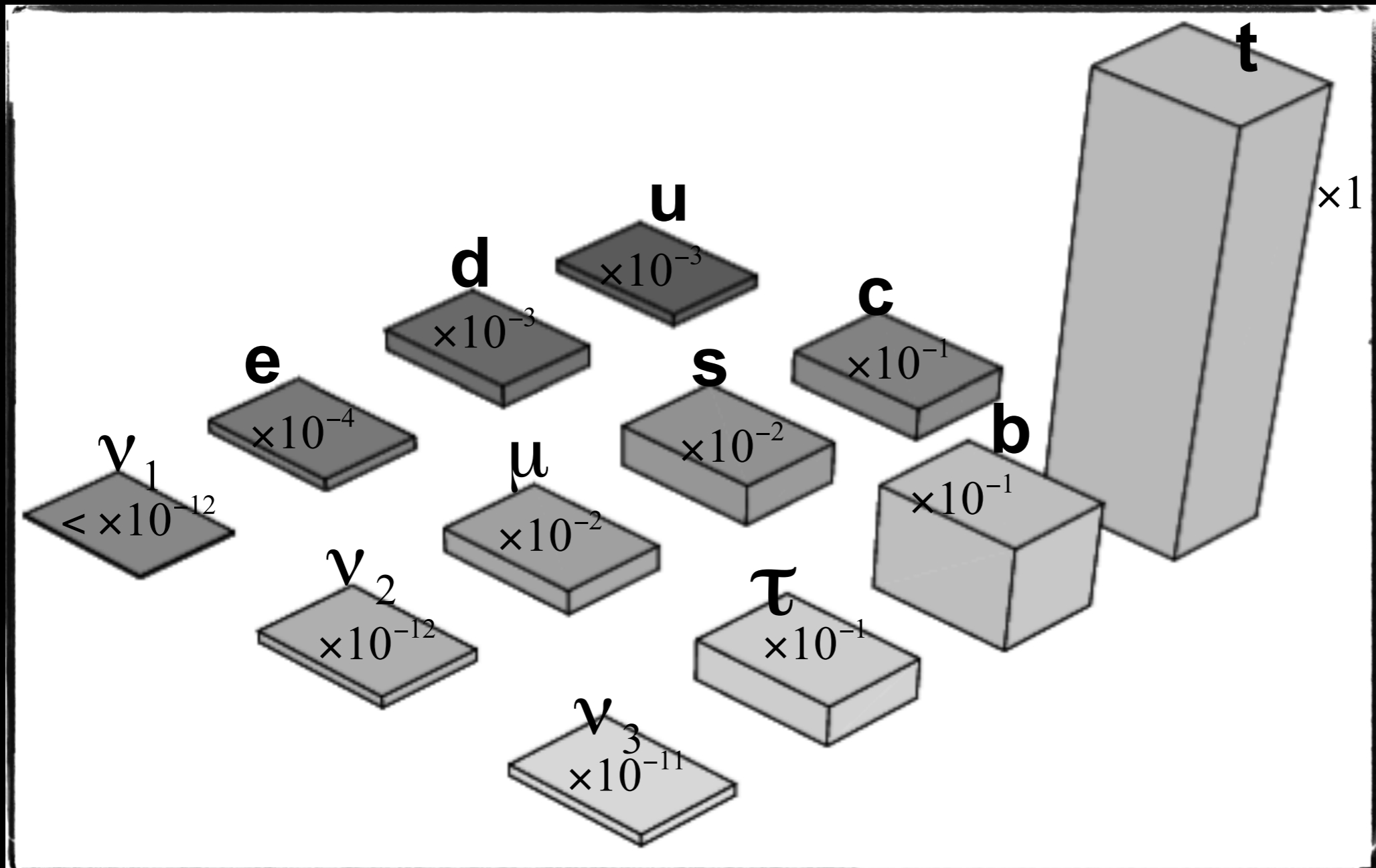
# Twin Pati-Salam Theory of Flavour with a TeV scale vector leptoquark

Based on [arXiv:2106.03876 [hep-ph]]

Steve King, 2nd September 2021

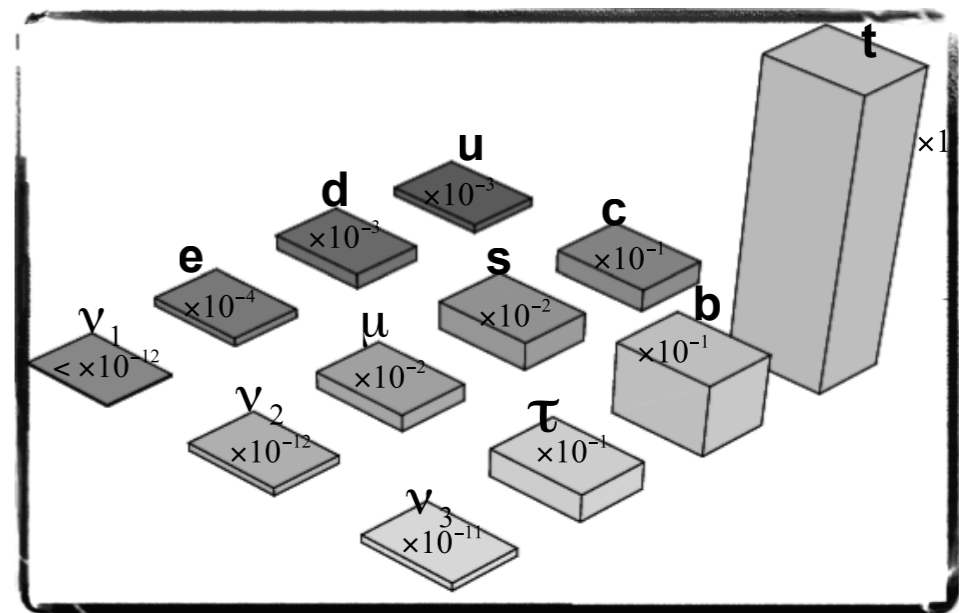
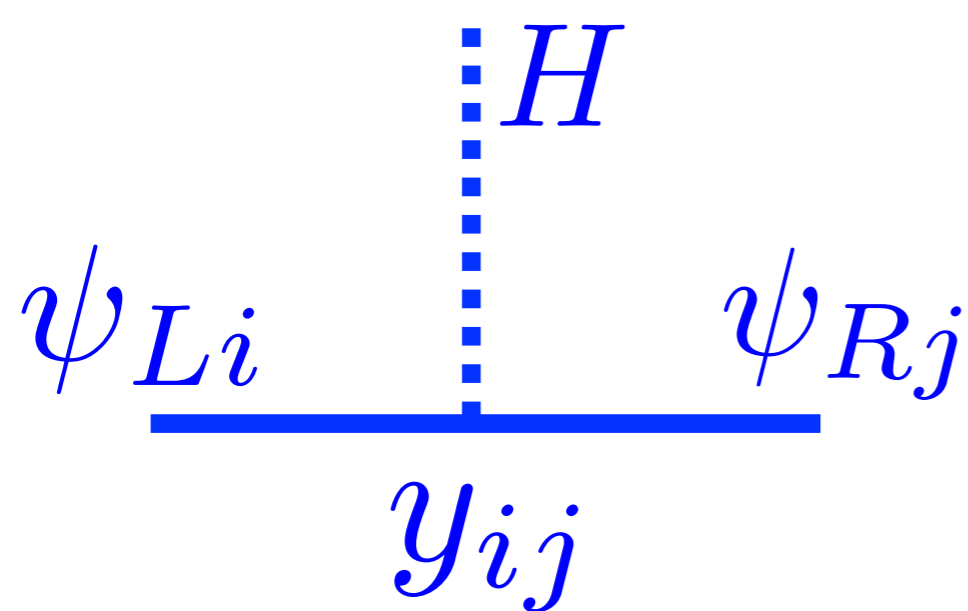


# Flavour Problem



# SM Yukawa couplings

$$y_{ij} H \bar{\psi}_{Li} \psi_{Rj}$$



Other flavour problems related to Lepton Flavour Universality (LFU)...

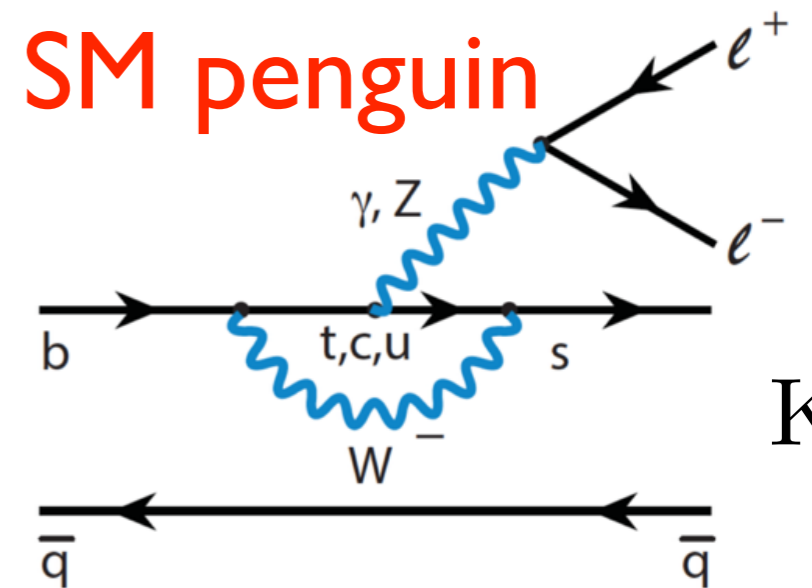
## Machine finds tantalising hints of new physics

By Pallab Ghosh  
Science correspondent

23 March

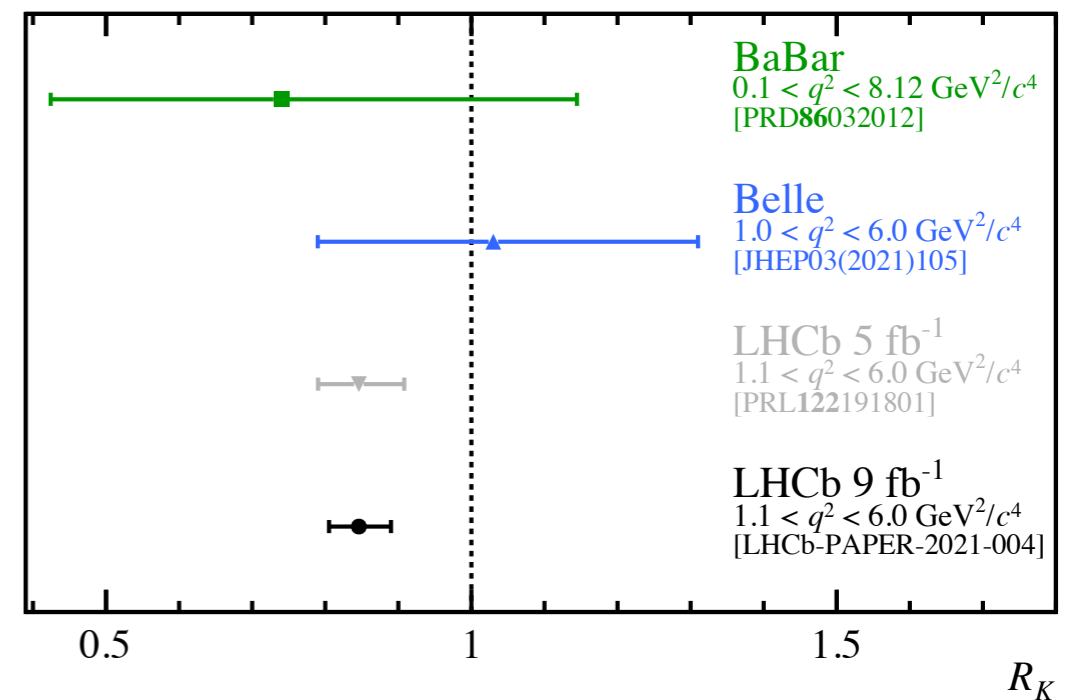
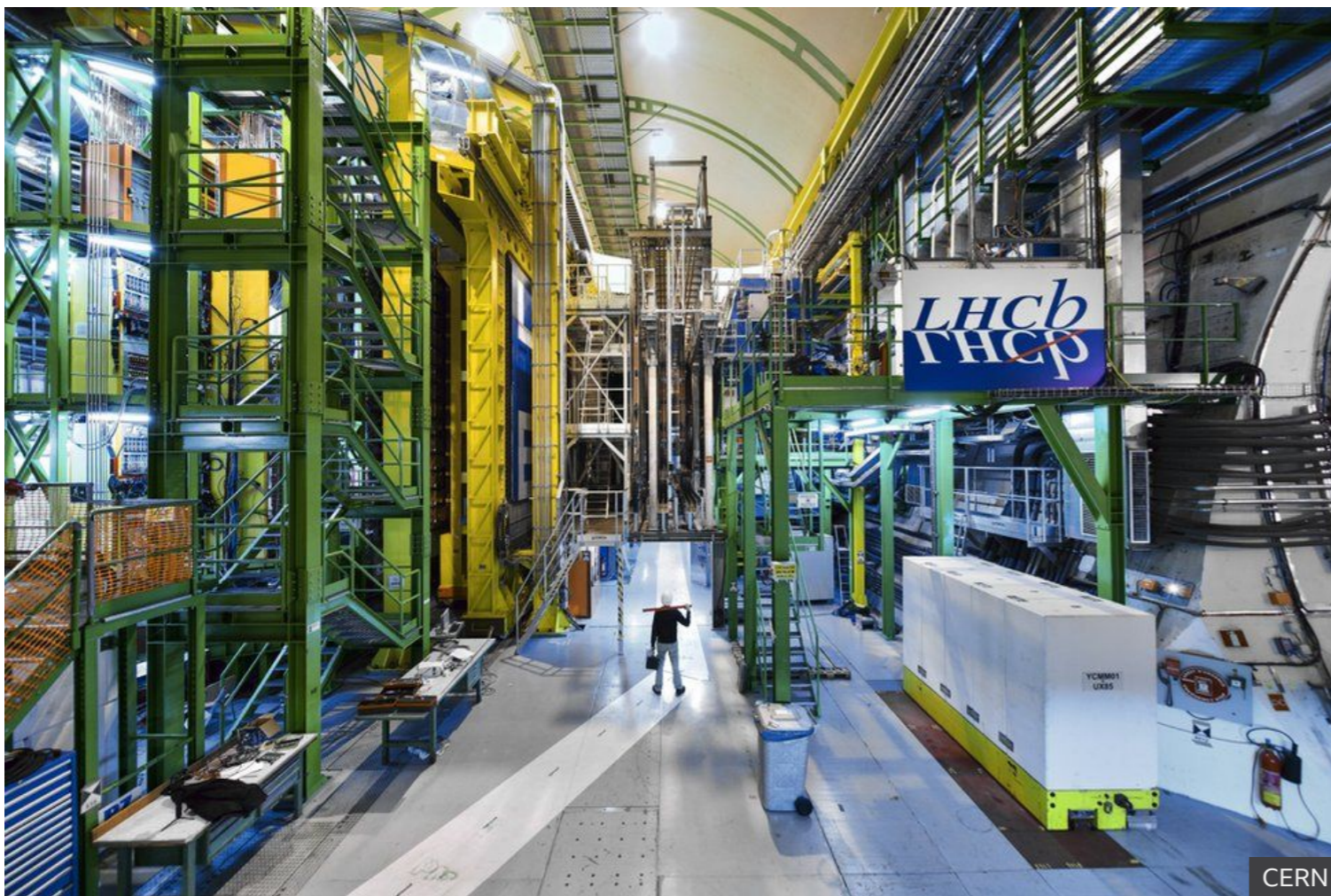


$$R_K = \frac{\int_{1.1 \text{ GeV}^2}^{6.0 \text{ GeV}^2} \frac{d\mathcal{B}(B^+ \rightarrow K^+ \mu^+ \mu^-)}{dq^2} dq^2}{\int_{1.1 \text{ GeV}^2}^{6.0 \text{ GeV}^2} \frac{d\mathcal{B}(B^+ \rightarrow K^+ e^+ e^-)}{dq^2} dq^2}$$



$$R_{K^{(*)}} := \frac{\mathcal{B}(B \rightarrow K^{(*)} \mu^+ \mu^-)}{\mathcal{B}(B \rightarrow K^{(*)} e^+ e^-)} \stackrel{\text{SM}}{\approx} 1$$

Actually use double ratio - see Monica talk



$$R_K = 0.846^{+0.042}_{-0.039} (\text{stat})^{+0.013}_{-0.012} (\text{syst})$$

3.1σ departure from LFU!

R.Aaij et al. [LHCb], [arXiv:2103.11769]

Physicists have uncovered a potential flaw in a theory that explains how the building blocks of the Universe behave.

# Possible operators for $R_K, R_{K^*}$

$$\mathcal{L}_{b \rightarrow s \mu \mu}^{\text{NP}} \supset \frac{4G_F}{\sqrt{2}} V_{tb} V_{ts}^* (\delta C_9^\mu O_9^\mu + \delta C_{10}^\mu O_{10}^\mu) + \text{h.c.}$$

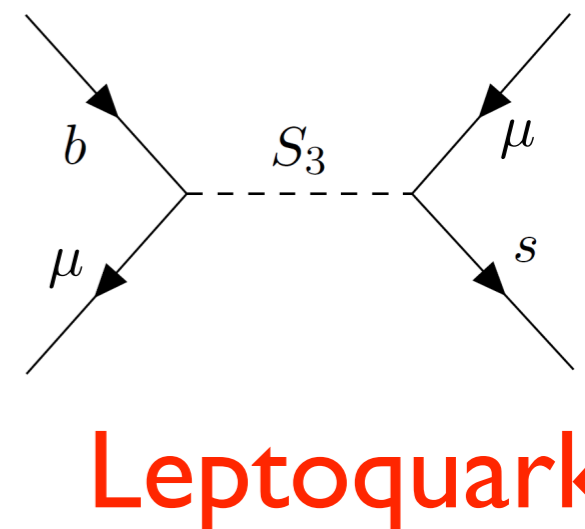
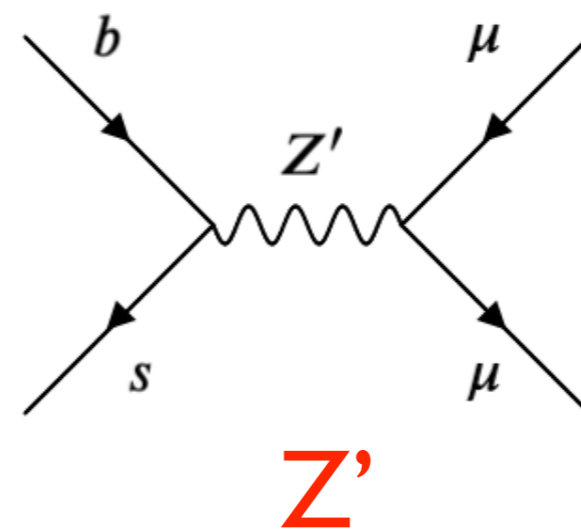
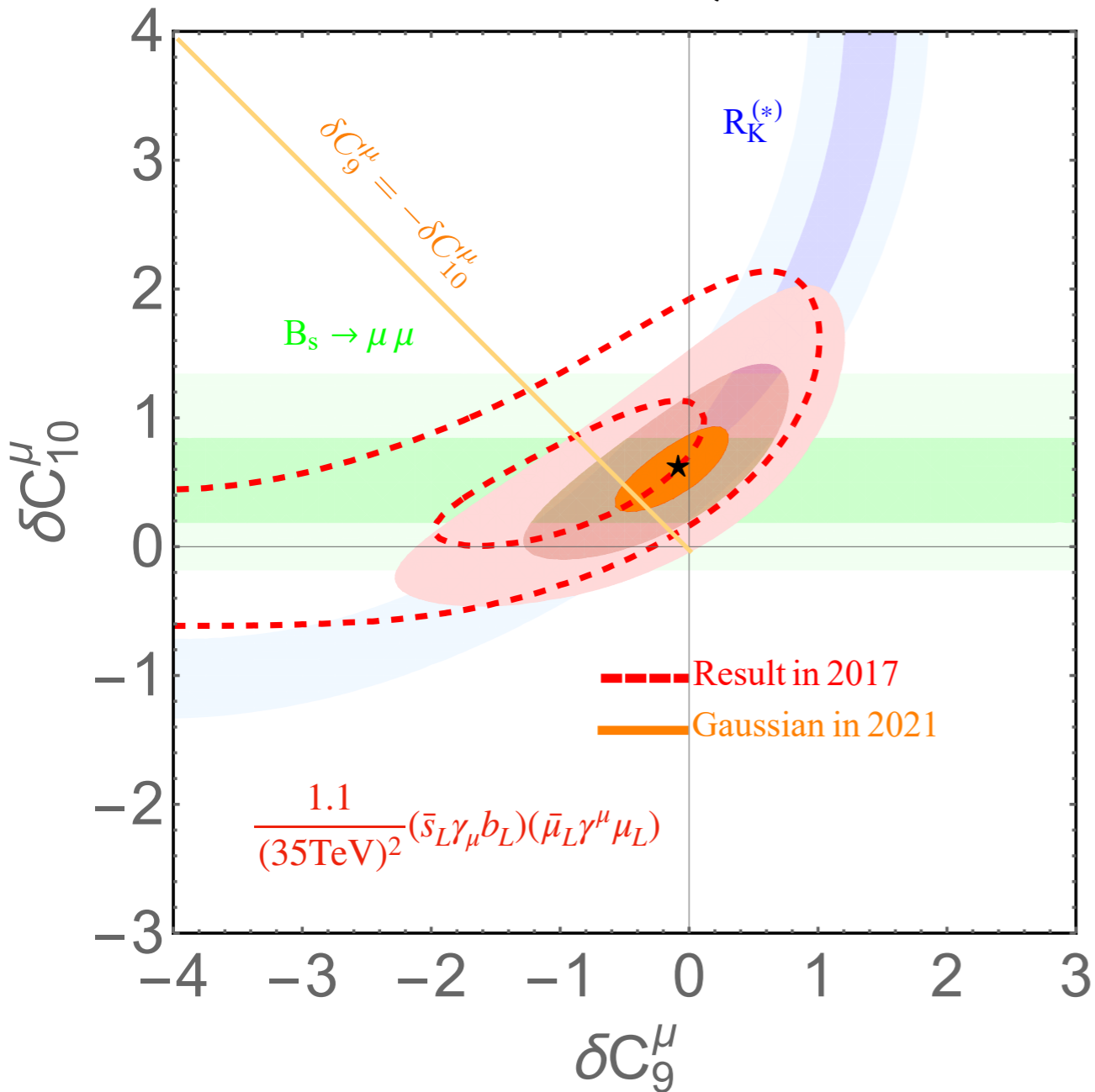
L.S.Geng, B.Grinstein, S.Jager, S.Y.Li,  
J.M.Camalich and R.X.Shi, 2103.12738

$$O_9^\mu = \frac{\alpha}{4\pi} (\bar{s}_L \gamma_\mu b_L) (\bar{\mu} \gamma^\mu \mu),$$

$$O_{10}^\mu = \frac{\alpha}{4\pi} (\bar{s}_L \gamma_\mu b_L) (\bar{\mu} \gamma^\mu \gamma_5 \mu).$$

LH muons

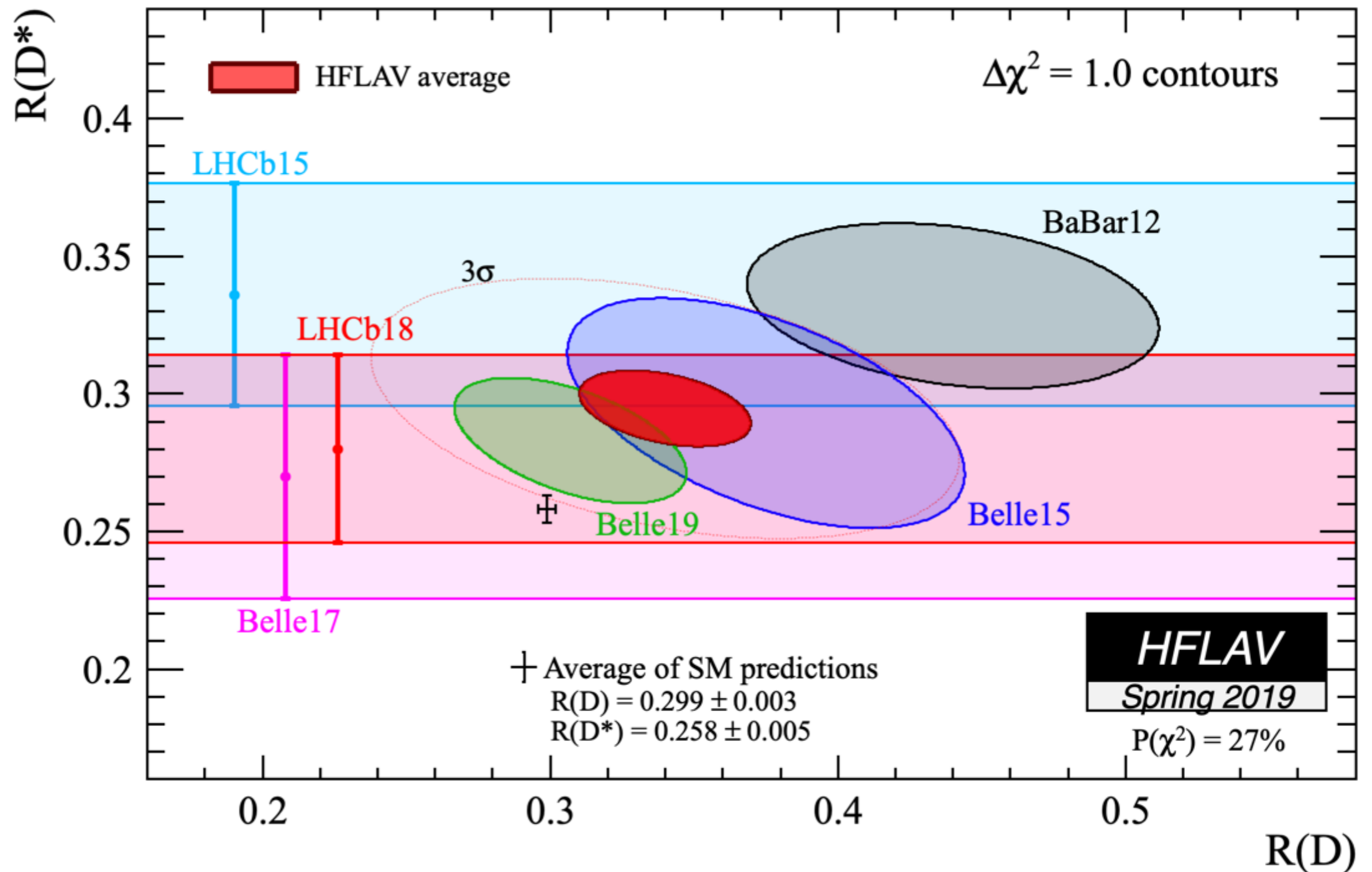
$$\text{Re}(\delta C_9^\mu) = -\text{Re}(\delta C_{10}^\mu)$$



# $b \rightarrow c \tau \nu$ Measurements

# $R_D, R_{D^*}$

$$R(D^{(*)}) = B \rightarrow D^{(*)} \tau \nu / B \rightarrow D^{(*)} \ell \nu$$

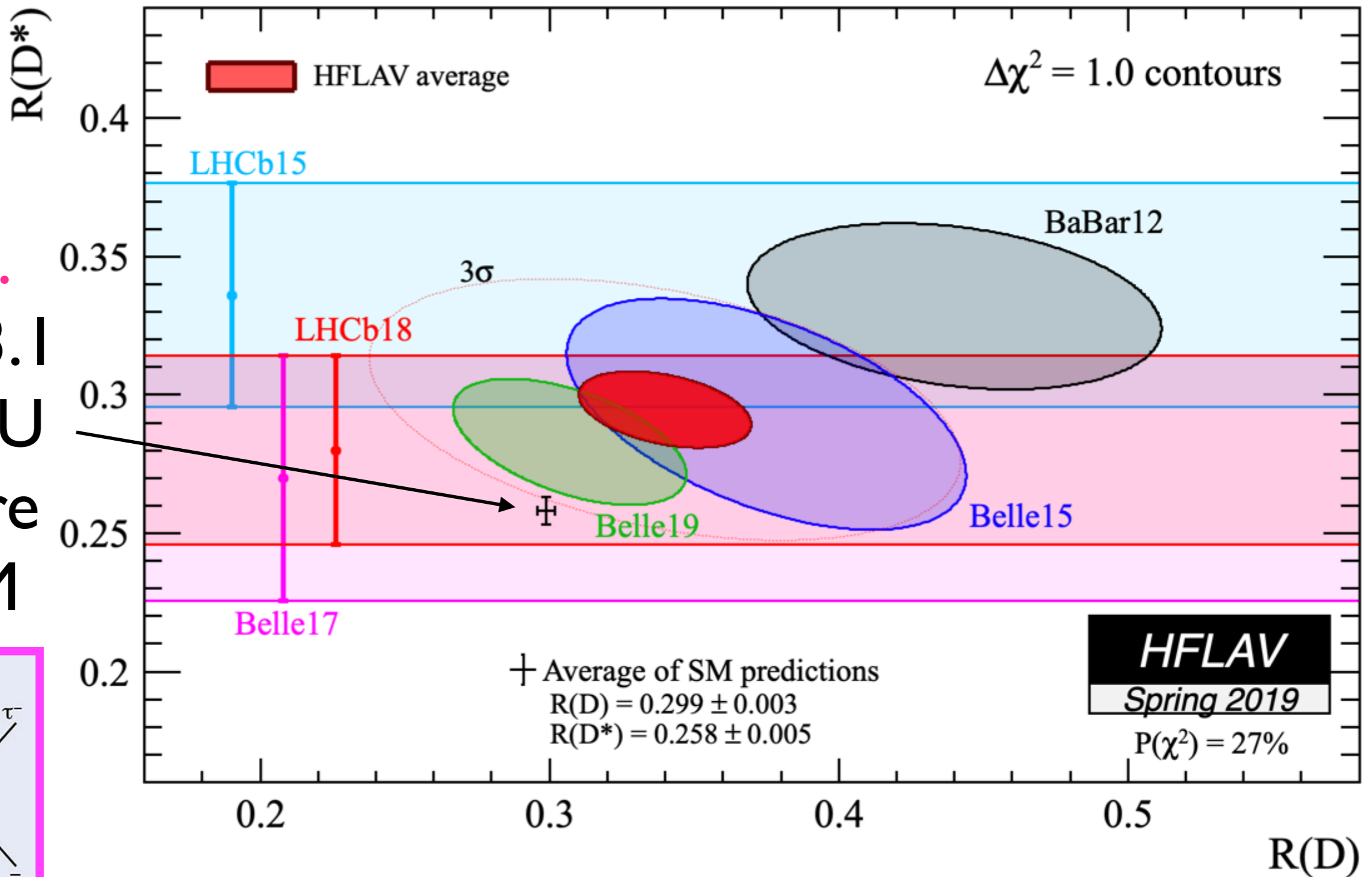
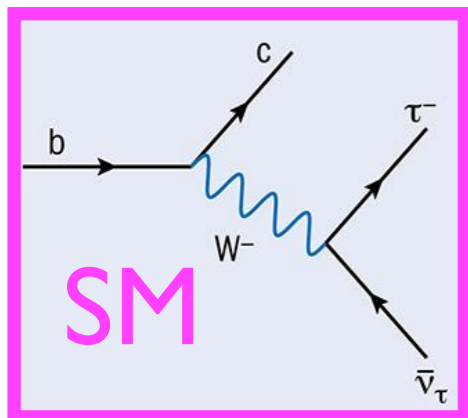


# $b \rightarrow c \tau \nu$ Measurements

# $R_D, R_{D^*}$

$$R(D^{(*)}) = B \rightarrow D^{(*)} \tau \nu / B \rightarrow D^{(*)} \ell \nu$$

Not BBC  
News...  
But still 3.1  
sigma LFU  
departure  
from SM



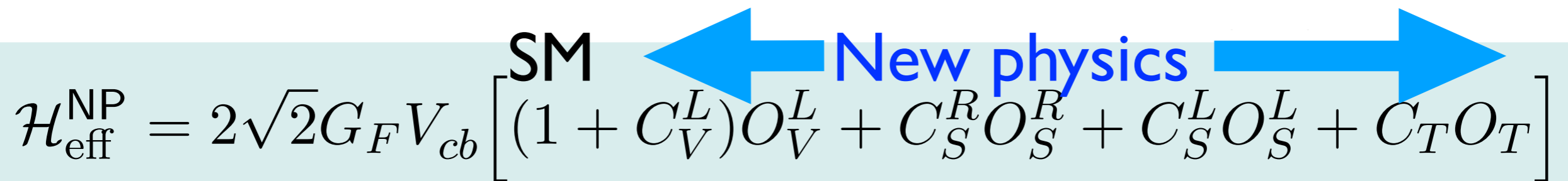
# Possible operators for $R_D, R_D^*$

$$O_V^L = (\bar{c}\gamma^\mu P_L b) (\bar{\tau}\gamma_\mu P_L \nu_\tau)$$

$$O_S^R = (\bar{c}P_R b) (\bar{\tau}P_L \nu_\tau)$$

$$O_T = (\bar{c}\sigma^{\mu\nu} P_L b) (\bar{\tau}\sigma_{\mu\nu} P_L \nu_\tau)$$

$$O_S^L = (\bar{c}P_L b) (\bar{\tau}P_L \nu_\tau)$$



$\mathcal{H}_{\text{eff}}^{\text{NP}} = 2\sqrt{2}G_F V_{cb} \left[ \overset{\text{SM}}{(1 + C_V^L)O_V^L} + C_S^R O_S^R + C_S^L O_S^L + C_T O_T \right]$



# Possible operators for $R_D, R_{D^*}$

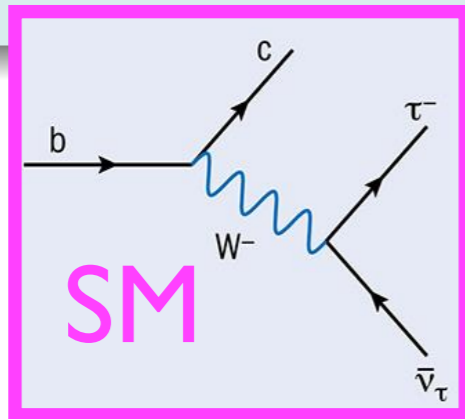
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$$\mathcal{H}_{\text{eff}}^{\text{NP}} = 2\sqrt{2}G_F V_{cb} \left[ \text{SM} \left( (1 + C_V^L) O_V^L + C_S^R O_S^R + C_S^L O_S^L + C_T O_T \right) \right]$$



W

W'

Charged Higgs



# Possible operators for $R_D, R_{D^*}$

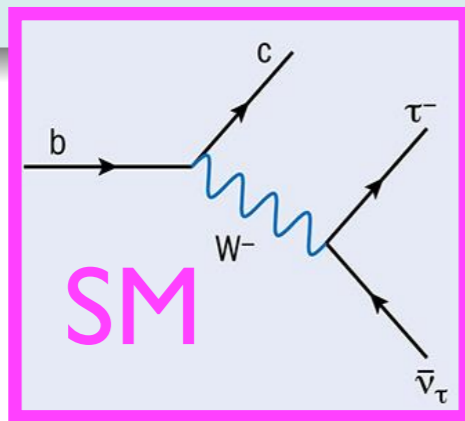
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$$\mathcal{H}_{\text{eff}}^{\text{NP}} = 2\sqrt{2}G_F V_{cb} \left[ \overset{\text{SM}}{(1 + C_V^L)O_V^L} + C_S^R O_S^R + C_S^L O_S^L + C_T O_T \right]$$



W

W'

Charged Higgs

Leptoquarks can contribute to all

- Leptoquark scalar S1=(3,1,1/3)
- Leptoquark scalar S2=(3,2,7/6)
- Leptoquark scalar S3=(3,3,1/3)
- Leptoquark vector U1=(3,1,2/3)
- Leptoquark vector U3=(3,3,2/3)

# Possible operators for $R_D, R_D^*$

$$O_V^L = (\bar{c}\gamma^\mu P_L b) (\bar{\tau}\gamma_\mu P_L \nu_\tau)$$

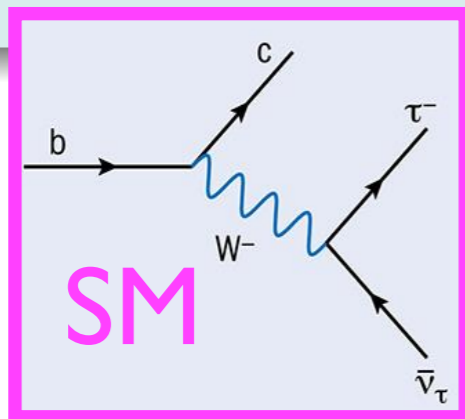
$$O_S^R = (\bar{c}P_R b) (\bar{\tau}P_L \nu_\tau)$$

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$$O_S^L = (\bar{c}P_L b) (\bar{\tau}P_L \nu_\tau)$$

$$\mathcal{H}_{\text{eff}}^{\text{NP}} = 2\sqrt{2}G_F V_{cb} \left[ \text{SM} \left( (1 + C_V^L) O_V^L + C_S^R O_S^R + C_S^L O_S^L + C_T O_T \right) \right]$$

← New physics →



W

W'

Charged Higgs

Leptoquarks can contribute to all

$U_1$  can explain both  $R_K$  and  $R_D$

Leptoquark scalar  $S1=(3,1,1/3)$

Leptoquark scalar  $S2=(3,2,7/6)$

Leptoquark scalar  $S3=(3,3,1/3)$

Leptoquark vector  $U1=(3,1,2/3)$

Leptoquark vector  $U3=(3,3,2/3)$

# Vector Leptoquark $U_1$

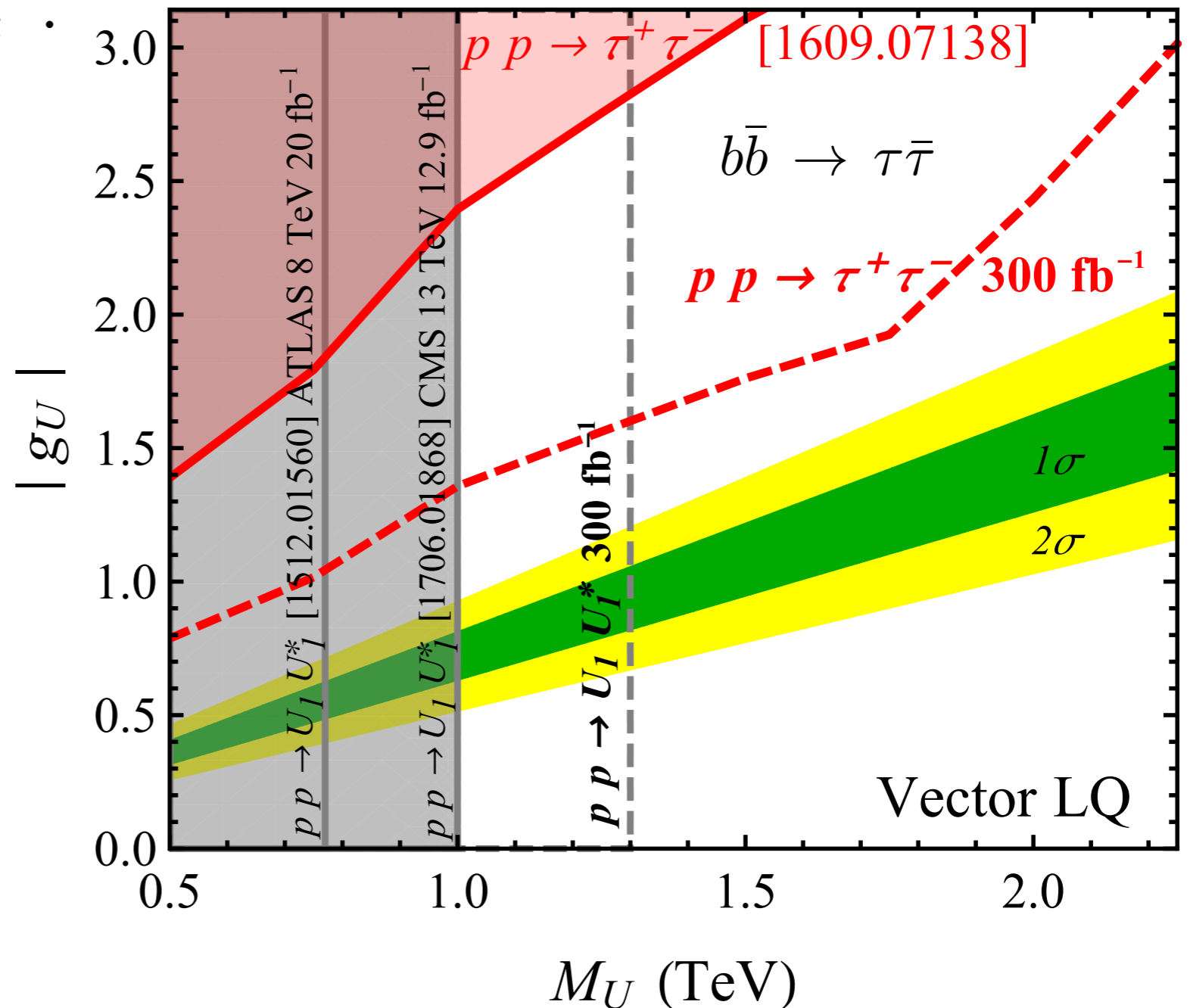
$$\mathcal{L}_U = -\frac{1}{2} U_{1,\mu\nu}^\dagger U^{1,\mu\nu} + M_U^2 U_{1,\mu}^\dagger U_1^\mu + g_U (J_U^\mu U_{1,\mu} + \text{h.c.})$$

$$J_U^\mu \equiv \beta_{i\alpha} \bar{Q}_i \gamma^\mu L_\alpha \cdot$$

LHC bound

$$M_U \gtrsim 1 \text{ TeV}$$

$$g_U \lesssim 2.5$$

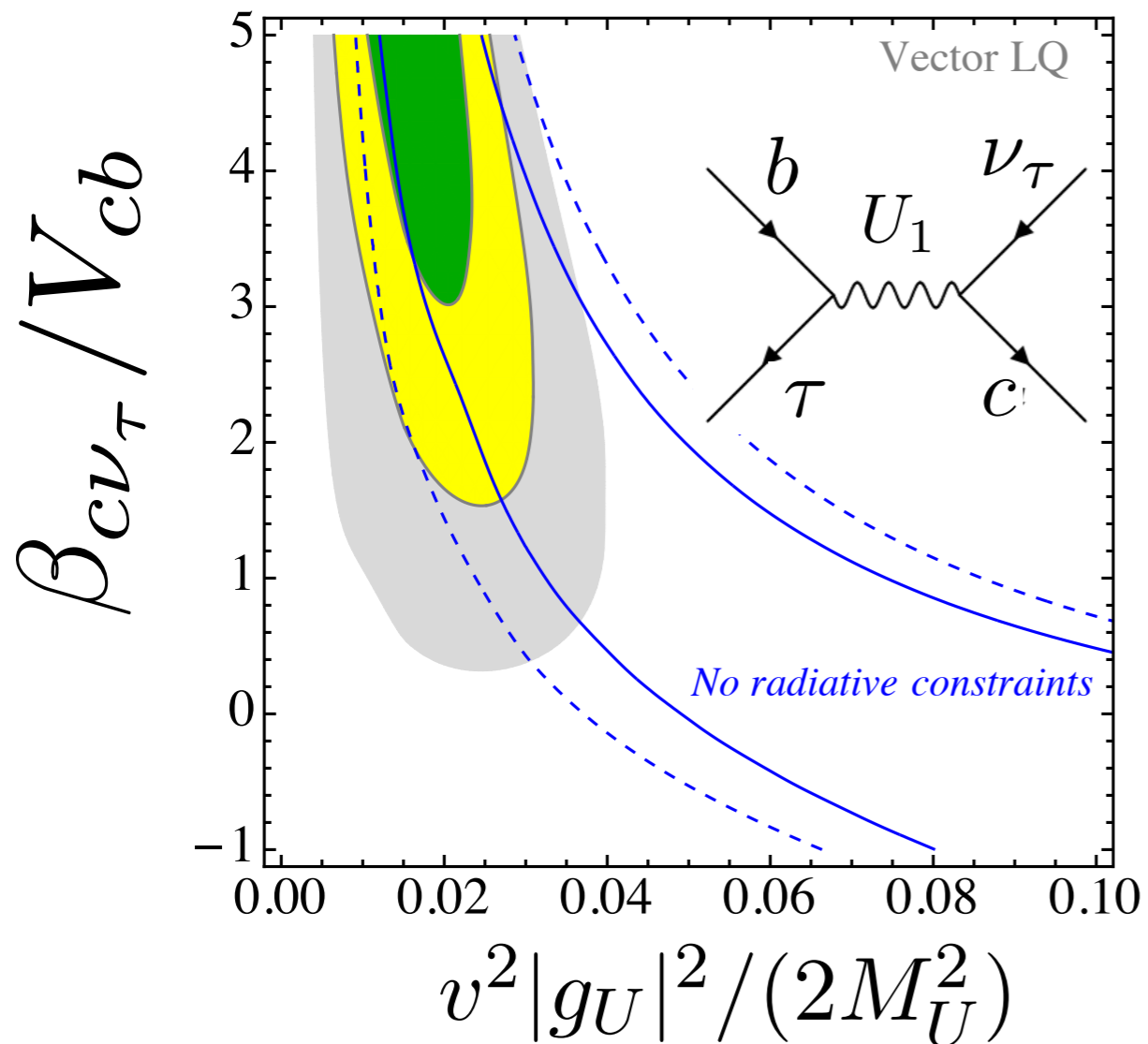


# Vector Leptoquark $U_1$

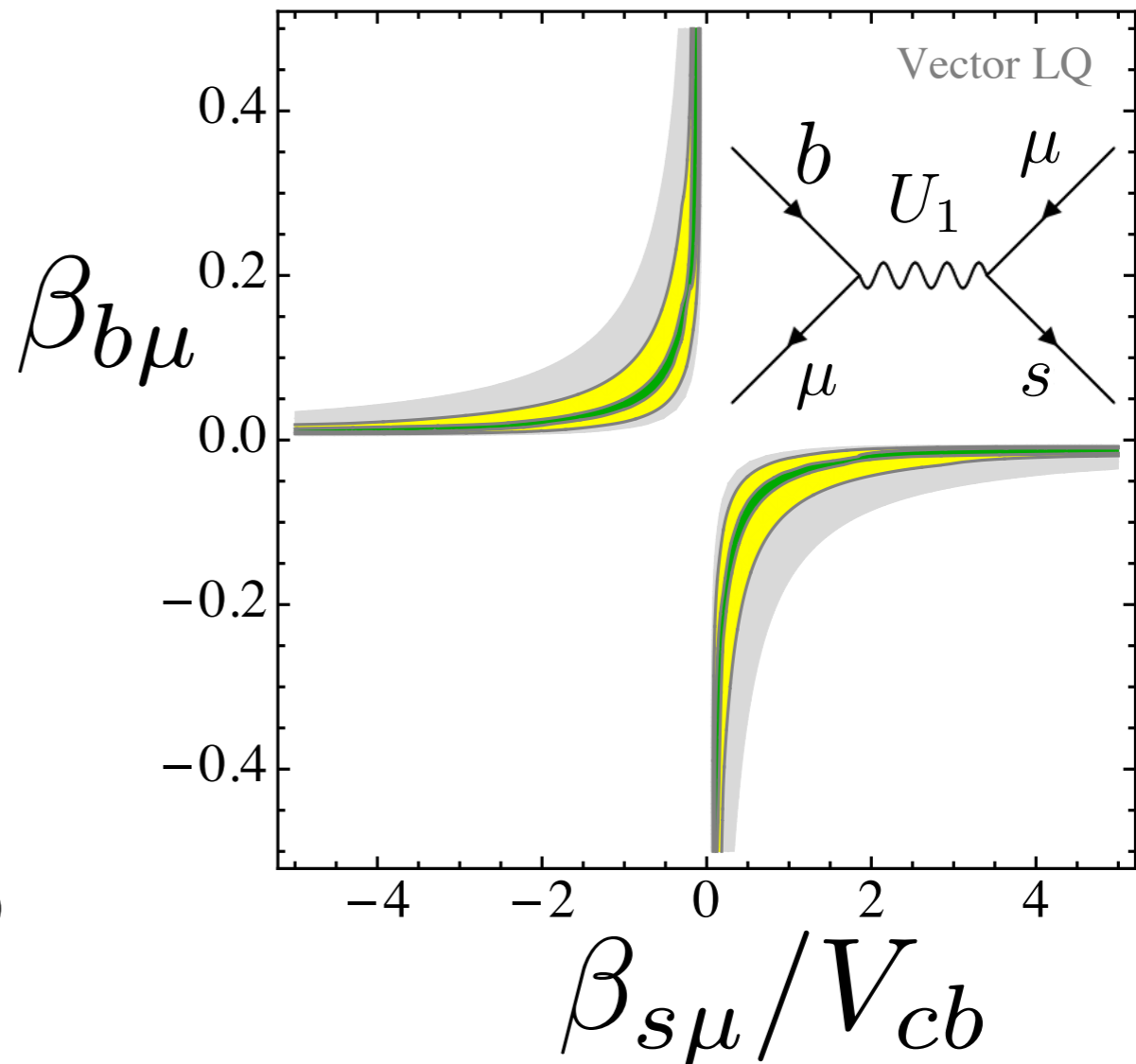
$$\mathcal{L}_U = -\frac{1}{2}U_{1,\mu\nu}^\dagger U^{1,\mu\nu} + M_U^2 U_{1,\mu}^\dagger U_1^\mu + g_U (J_U^\mu U_{1,\mu} + \text{h.c.})$$

$$J_U^\mu \equiv \beta_{i\alpha} \bar{Q}_i \gamma^\mu L_\alpha .$$

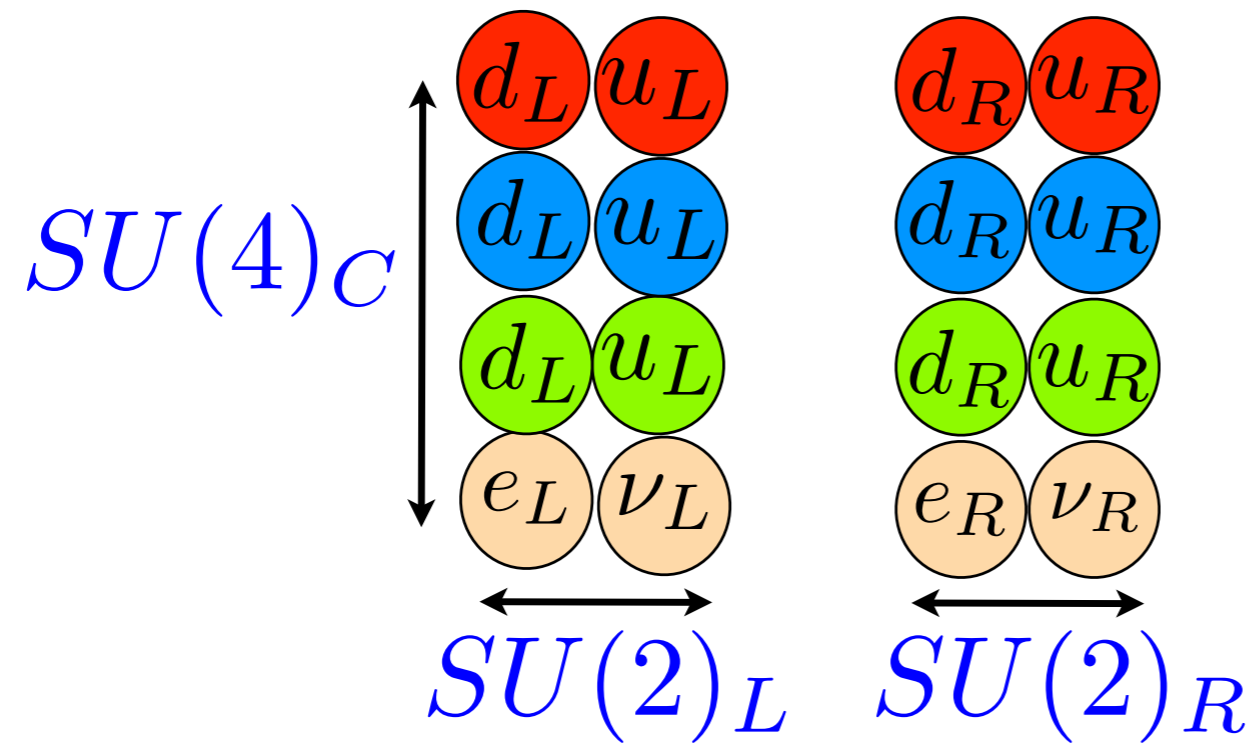
Fit to  $R_D$



Fit to  $R_K$

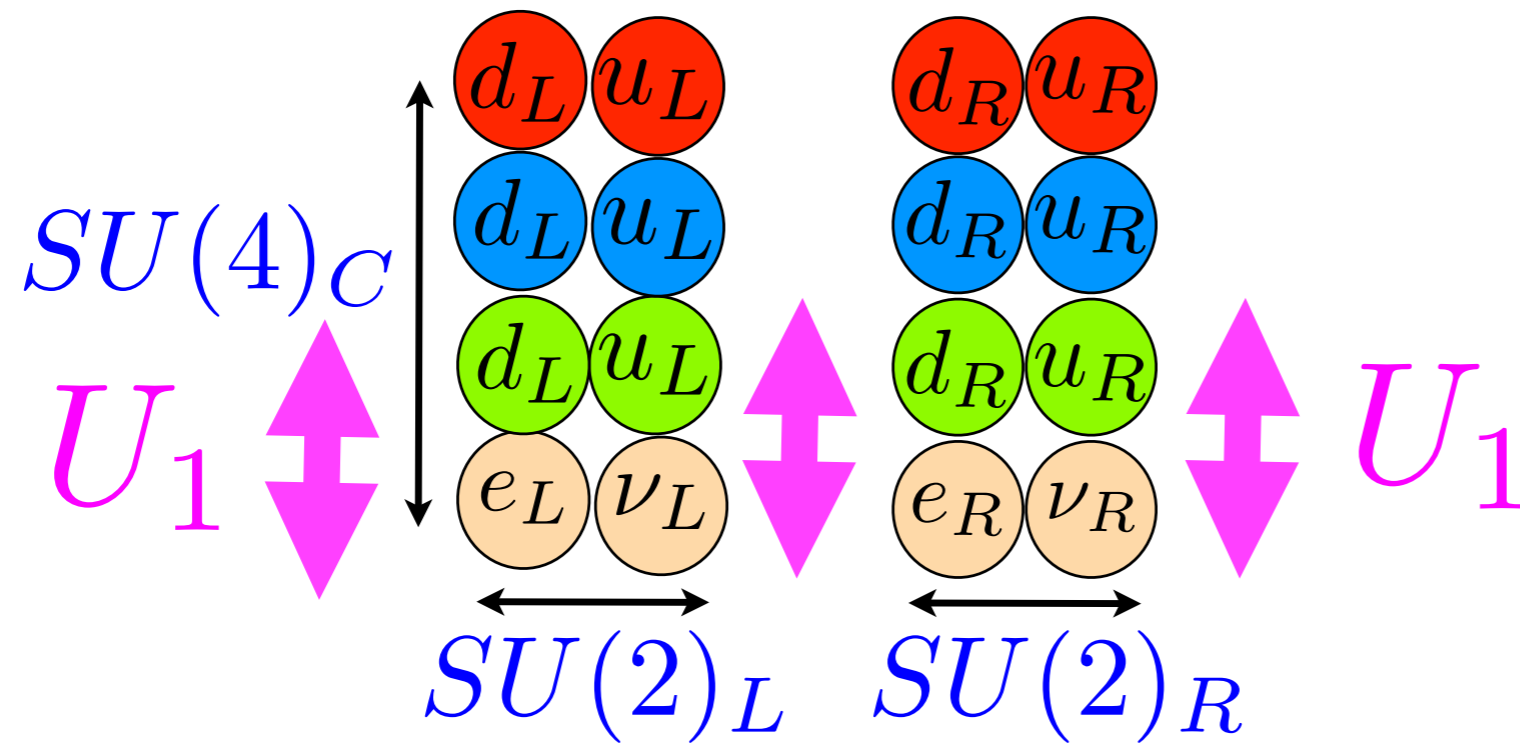


# Vector Leptoquark $U_1$ from Pati-Salam



Quark-lepton  
unification  
“leptons are  
fourth colour”

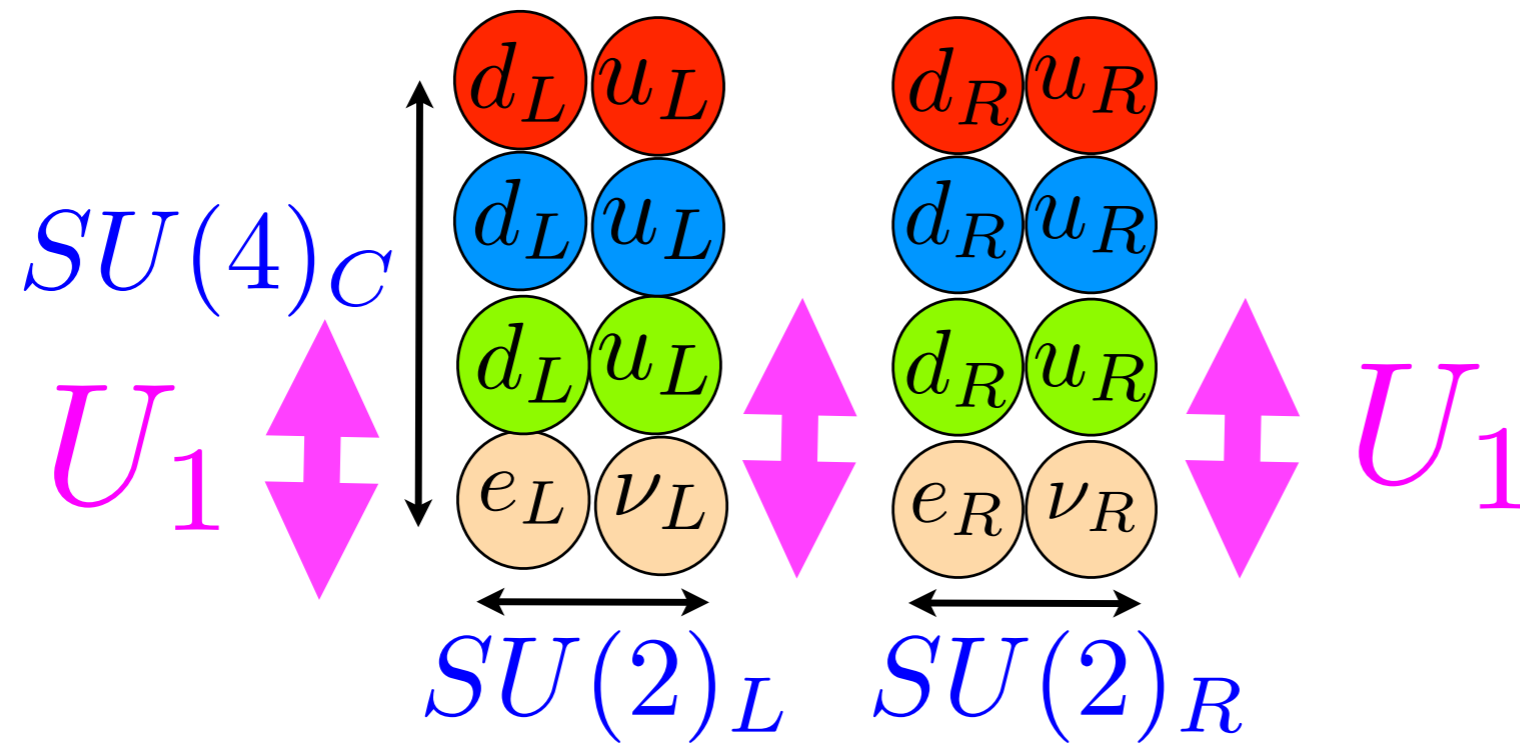
# Vector Leptoquark $U_1$ from Pati-Salam



Quark-lepton  
unification  
“leptons are  
fourth colour”

$U_1$  can explain both  $R_K$  and  $R_D$  for  $\sim 1$  TeV mass

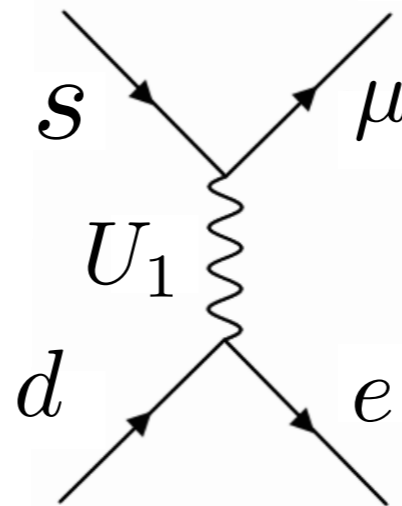
# Vector Leptoquark $U_1$ from Pati-Salam



Quark-lepton  
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$U_1$  can explain both  $R_K$  and  $R_D$  for 1 TeV mass

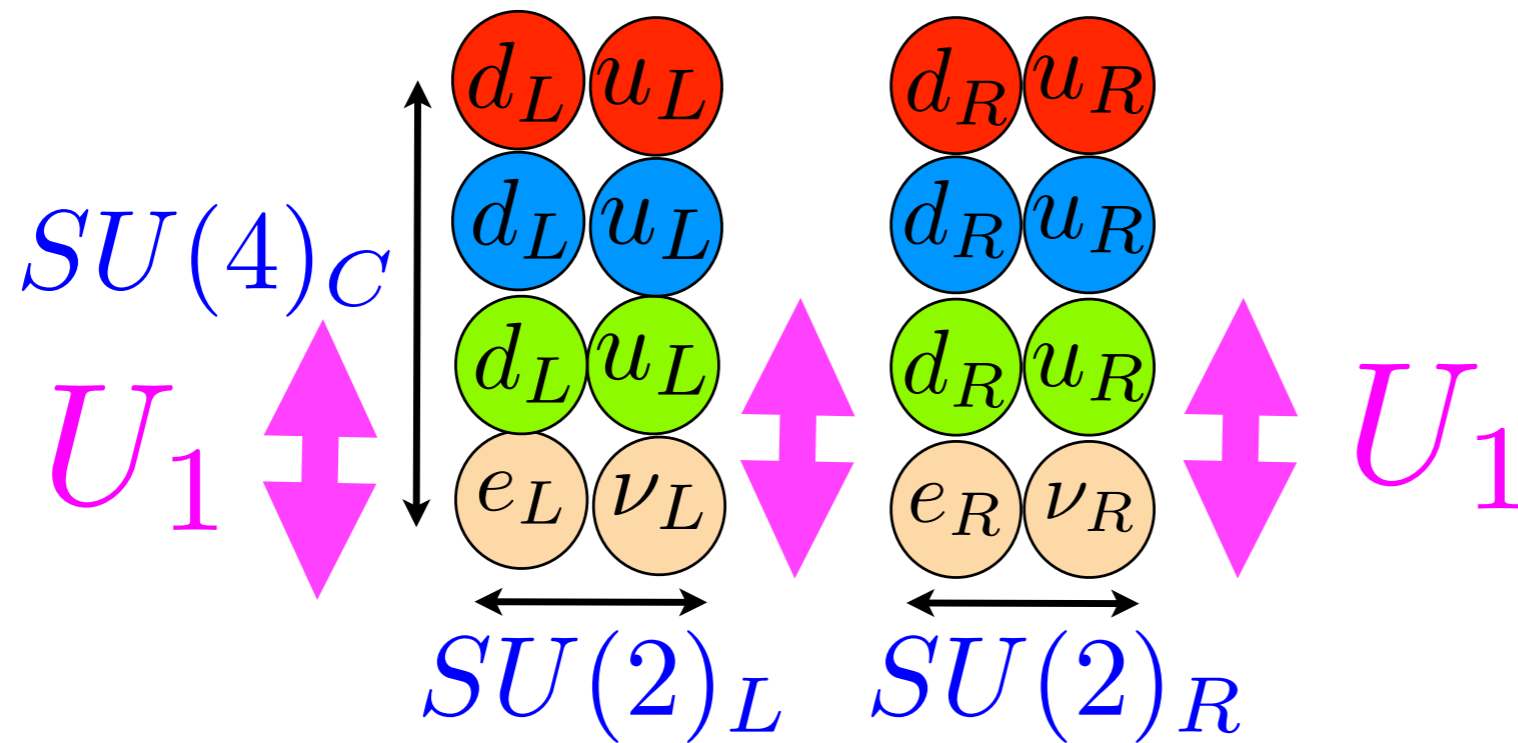
But TeV mass not allowed due to  $K_L \rightarrow \mu e$



$$M_{U_1} \gtrsim 1 \text{ PeV}$$



# Vector Leptoquark $U_1$ from Pati-Salam

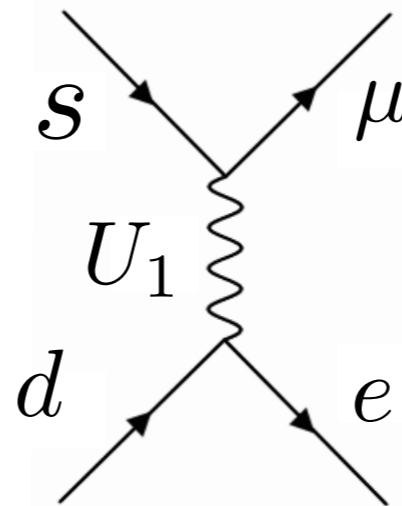


Quark-lepton  
unification  
“leptons are  
fourth colour”

$U_1$  can explain both  $R_K$  and  $R_D$  for 1 TeV mass

But TeV mass not allowed due to  $K_L \rightarrow \mu e$

Difficult to reconcile in a  
UV complete PS model



$$M_{U_1} \gtrsim 1 \text{ PeV}$$

- $SU(4) \times SU(3)' \times SU(2)_L \times U(1)_Y$  + Vector-like fermions  
L. Di Luzio, A. Greljo, M. Nardecchia, arXiv:1708.08450
- $SU(4) \times U(2)_L \times SU(2)_R$  + Vector-like fermions  
L. Calibbi, AC, T. Li, arXiv:1709.00692
- $SU(4) \times SU(4) \times SU(4)$   
M. Bordone, C. Cornella, J. Fuentes-Martin, G. Isidori, arXiv:1712.01368
- $SU(4) \times SU(2)_L \times SU(2)_R$  including scalar LQs and light right-handed neutrinos  
J. Heeck, D. Teresi, arXiv:1808.07492
- $SU(8)$  might even explain  $\varepsilon'/\varepsilon$   
S. Matsuzaki, K. Nishiwaki and K. Yamamoto, arXiv:1806.02312
- $SU(4) \times SU(2)_L \times SU(2)_R$  in RS background  
M. Blanke, AC, arXiv:1801.07256

# Possible UV completions

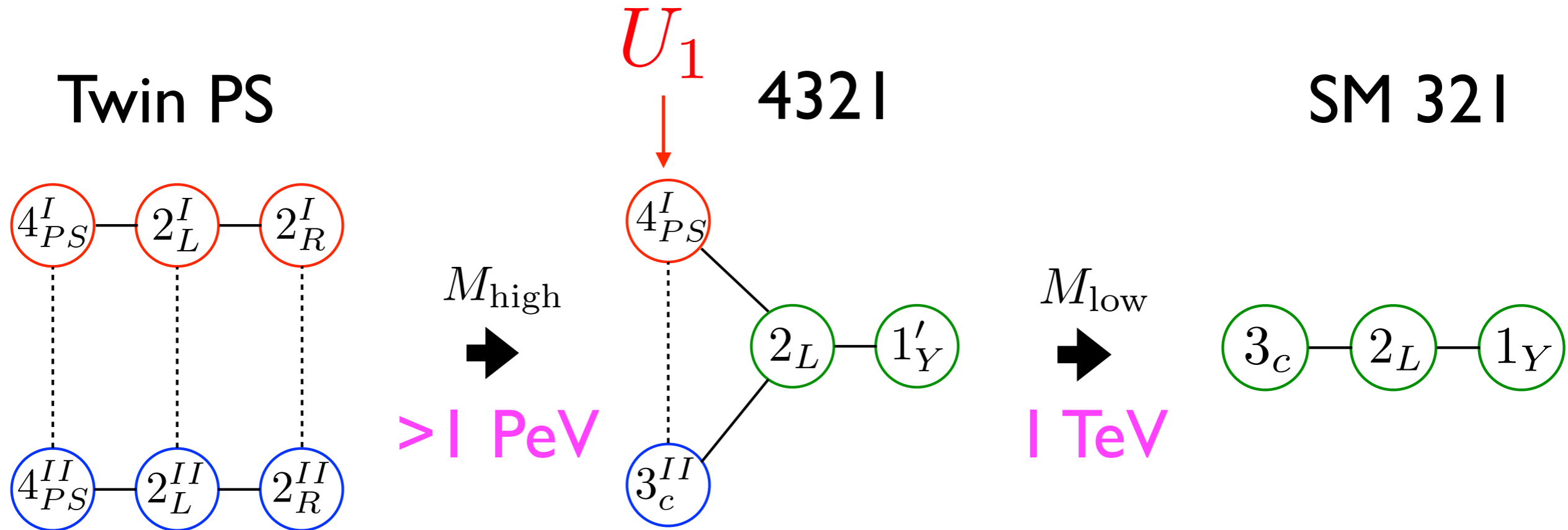
Andreas Crivellin

- $SU(4) \times SU(3)' \times SU(2)_L \times U(1)_Y$  + Vector-like fermions  
L. Di Luzio, A. Greljo, M. Nardecchia, arXiv:1708.08450
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M. Bordone, C. Cornella, J. F. Marchesoni, G. Isidori, arXiv:1712.01368
- $SU(4) \times SU(2)_L \times SU(2)_R$  + Vector-like fermions + scalar LQs and light right-handed neutrinos  
J. Heer, arXiv:1808.07492
- $SU(4) \times SU(2)_L \times SU(2)_R$  + Vector-like fermions + scalar LQs and light right-handed neutrinos even explain  $\epsilon'/\epsilon$   
S. M. Barr, K. Nishiwaki and K. Yamamoto, arXiv:1806.02312
- $SU(4) \times SU(2)_L \times SU(2)_R$  in RS background  
M. Blanke, AC, arXiv:1801.07256

None of them attempt to explain the origin of quark and lepton masses

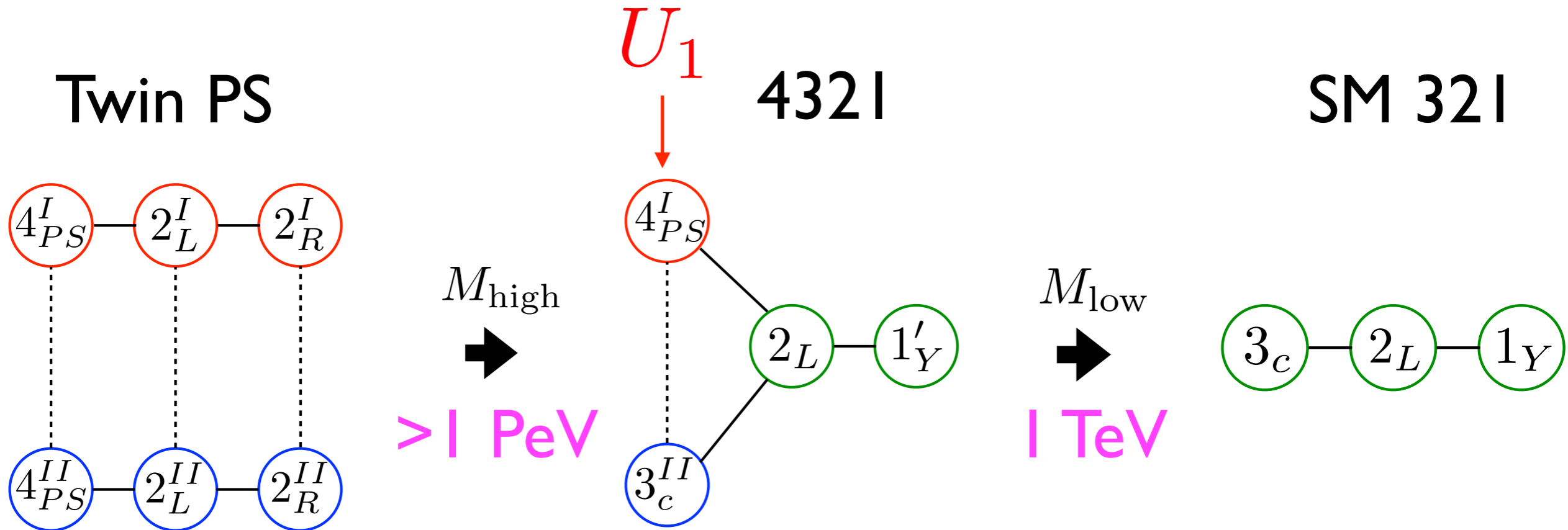
# Twin Pati-Salam

Each circle represents a gauge group



# Twin Pati-Salam

Each circle represents a gauge group



An attempt to explain the origin of fermion masses  
as well as RK and RD simultaneously

(N.B. Due to time limitations we ignore first family and neutrino masses in this talk - see paper for full model)

Field	$SU(4)_I^{PS}$	$SU(2)_L^I$	$SU(2)_R^I$	$SU(4)_{PS}^{II}$	$SU(2)_L^{II}$	$SU(2)_R^{II}$
$\psi_{1,2,3}$	<b>1</b>	<b>1</b>	<b>1</b>	<b>4</b>	<b>2</b>	<b>1</b>
$\psi_{1,2,3}^c$	<b>1</b>	<b>1</b>	<b>1</b>	$\bar{4}$	<b>1</b>	$\bar{2}$

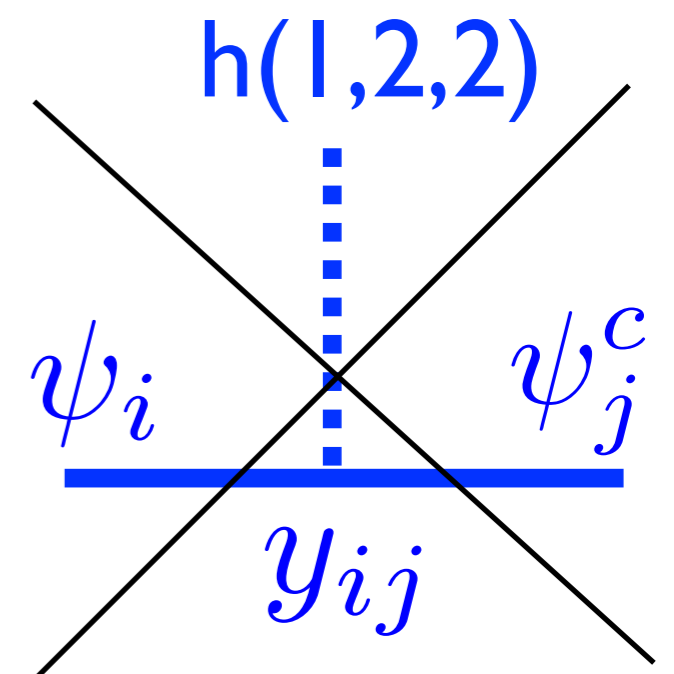
3 chiral families under second group do not couple to  $U_1$

Field	$SU(4)_I^{PS}$	$SU(2)_L^I$	$SU(2)_R^I$	$SU(4)_{PS}^{II}$	$SU(2)_L^{II}$	$SU(2)_R^{II}$
$\psi_{1,2,3}$	<b>1</b>	<b>1</b>	<b>1</b>	<b>4</b>	<b>2</b>	<b>1</b>
$\psi_{1,2,3}^c$	<b>1</b>	<b>1</b>	<b>1</b>	$\bar{4}$	<b>1</b>	$\bar{2}$
$\psi_4$	<b>4</b>	<b>2</b>	<b>1</b>	<b>1</b>	<b>1</b>	<b>1</b>
$\bar{\psi}_4$	$\bar{4}$	$\bar{2}$	<b>1</b>	<b>1</b>	<b>1</b>	<b>1</b>
$\psi_4^c$	$\bar{4}$	<b>1</b>	$\bar{2}$	<b>1</b>	<b>1</b>	<b>1</b>
$\bar{\psi}_4^c$	<b>4</b>	<b>1</b>	<b>2</b>	<b>1</b>	<b>1</b>	<b>1</b>

4th vector-like family under first group couple to  $U_1$

Field	$SU(4)_{PS}^I$	$SU(2)_L^I$	$SU(2)_R^I$	$SU(4)_{PS}^{II}$	$SU(2)_L^{II}$	$SU(2)_R^{II}$
$\psi_{1,2,3}$	1	1	1	4	2	1
$\psi_{1,2,3}^c$	1	1	1	$\bar{4}$	1	$\bar{2}$
$\psi_4$	4	2	1	1	1	1
$\bar{\psi}_4$	$\bar{4}$	$\bar{2}$	1	1	1	1
$\psi_4^c$	$\bar{4}$	1	$\bar{2}$	1	1	1
$\bar{\psi}_4^c$	4	1	2	1	1	1
$\phi$	4	2	1	$\bar{4}$	$\bar{2}$	1
$\bar{\phi}$	$\bar{4}$	1	$\bar{2}$	4	1	2
$H$	$\bar{4}$	$\bar{2}$	1	4	1	2
$\bar{H}$	4	1	2	$\bar{4}$	$\bar{2}$	1

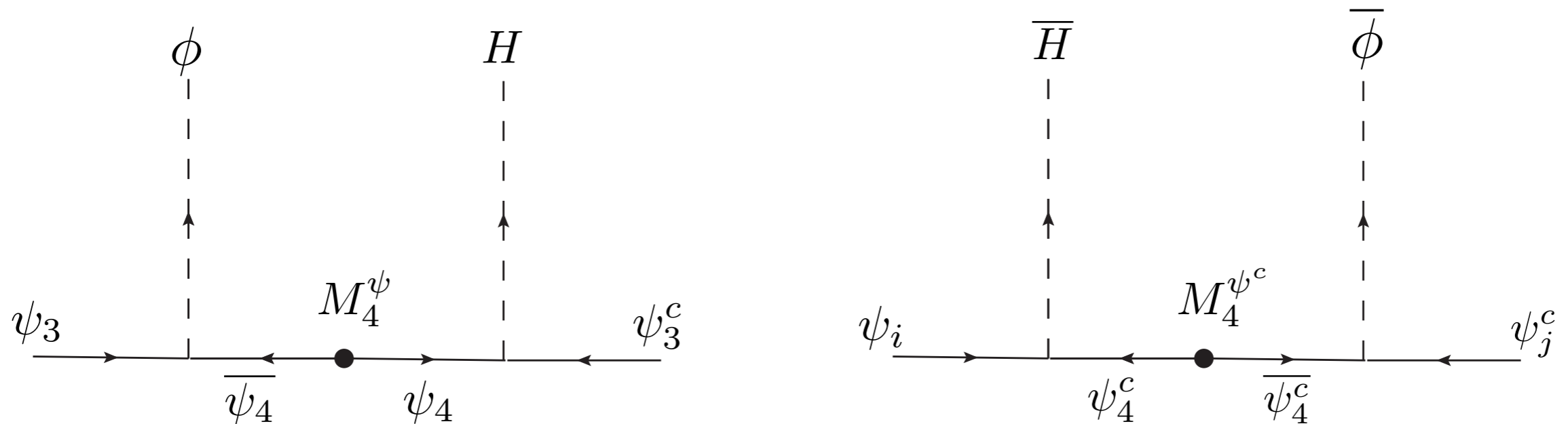
- Higgs transform under both groups and generate the mixing
- No Higgs  $h(1,2,2)$  under second group so no Yukawa coupling





# Fermion Masses

Generated by mixing with vector-like family

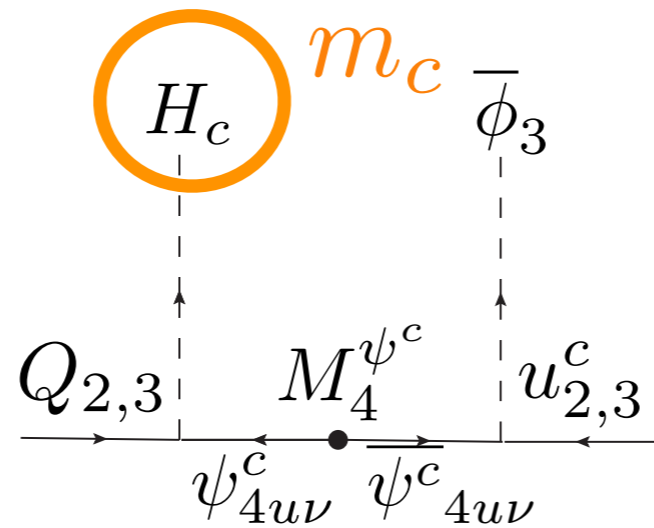
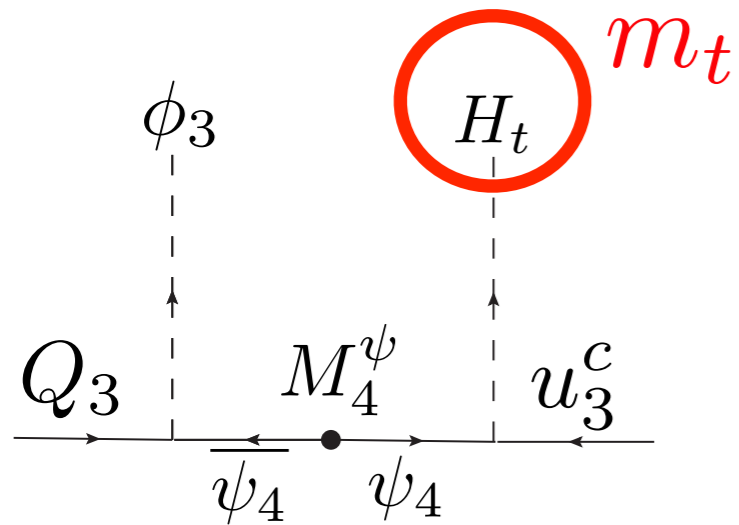


$$\mathcal{L}_{4eff}^{Yuk} = \begin{pmatrix} & \psi_1^c & \psi_2^c & \psi_3^c \\ \psi_1 | & 0 & 0 & 0 \\ \psi_2 | & 0 & 0 & 0 \\ \psi_3 | & 0 & 0 & x_{34}^\psi y_{43}^\psi \end{pmatrix} \left( \frac{\phi}{M_4^\psi} H \right) \sim 1 + \begin{pmatrix} & \psi_1^c & \psi_2^c & \psi_3^c \\ \psi_1 | & 0 & 0 & 0 \\ \psi_2 | & 0 & y_{24}^\psi x_{42}^{\psi^c} & y_{24}^\psi x_{43}^{\psi^c} \\ \psi_3 | & 0 & y_{34}^\psi x_{42}^{\psi^c} & y_{34}^\psi x_{43}^{\psi^c} \end{pmatrix} \left( \frac{\bar{\phi}}{M_4^{\psi^c}} \bar{H} \right) \ll 1$$

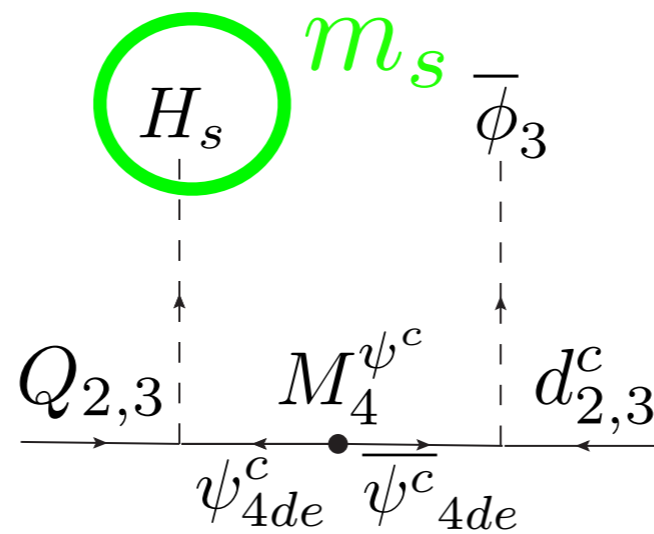
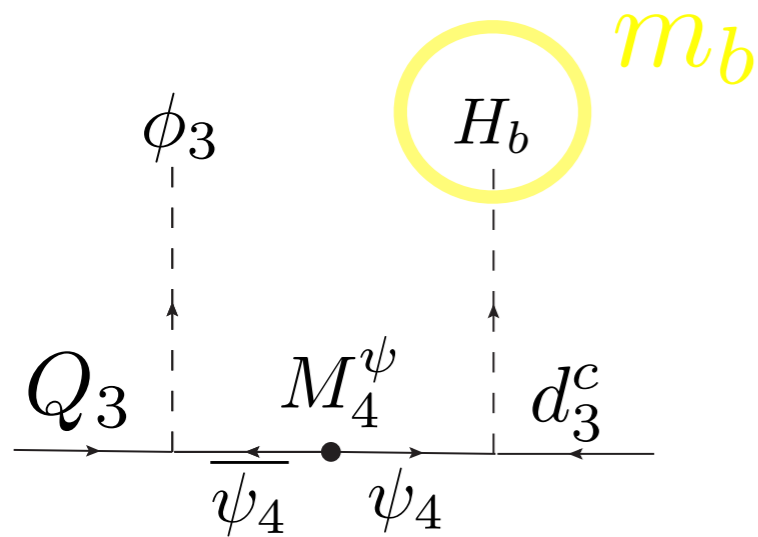
3rd family

23 block

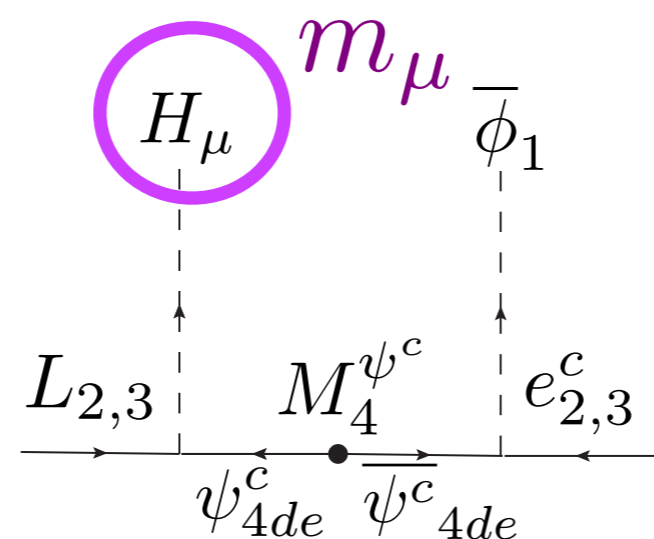
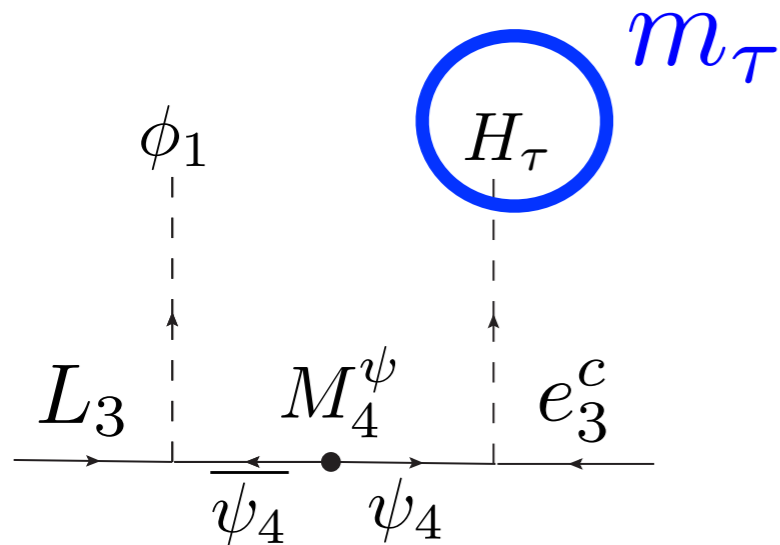
# Under 4321 have a personal Higgs for each mass



$$M_u \sim \begin{pmatrix} 0 & \\ m_c & m_c \\ m_c & m_t \end{pmatrix}$$



$$M_d \sim \begin{pmatrix} 0 & \\ m_s & m_s \\ m_s & m_b \end{pmatrix}$$



$$M_e \sim \begin{pmatrix} 0 & \\ m_\mu & m_\mu \\ m_\mu & m_\tau \end{pmatrix}$$

# 4321 phenomenology

$$\frac{g_4}{\sqrt{2}} (\bar{Q}_{L4} \gamma^\mu L_{L4} + \text{H.c.}) U_{1\mu} \quad \text{Vector Leptoquark}$$

# 4321 phenomenology

no couplings to chiral quarks and leptons

$$\frac{g_4}{\sqrt{2}} (\bar{Q}_{L4} \gamma^\mu L_{L4} + \text{H.c.}) U_{1\mu} \quad \text{Vector Leptoquark}$$

# 4321 phenomenology

no couplings to chiral quarks and leptons

$$\frac{g_4}{\sqrt{2}} (\bar{Q}_{L4} \gamma^\mu L_{L4} + \text{H.c.}) U_{1\mu} \quad \text{Vector Leptoquark}$$

$$+ \frac{g_4 g_s}{g_3} \left( \bar{Q}_{L4} \gamma^\mu T^a Q_{L4} - \frac{g_3^2}{g_4^2} \bar{Q}_{Li} \gamma^\mu T^a Q_{Li} \right) g_\mu^{Ia} \quad \text{coloron} \quad Z'$$

$$+ \frac{\sqrt{3} g_4 g_Y}{\sqrt{2} g_1} \left( \frac{1}{6} \bar{Q}_{L4} \gamma^\mu Q_{L4} - \frac{1}{2} \bar{L}_{L4} \gamma^\mu L_{L4} - \frac{g_1^2}{9g_4^2} \bar{Q}_{Li} \gamma^\mu Q_{Li} + \frac{g_1^2}{3g_4^2} \bar{L}_{Li} \gamma^\mu L_{Li} \right) Z'_\mu$$

# 4321 phenomenology

no couplings to chiral quarks and leptons

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suppressed couplings to chiral quarks and leptons

# 4321 phenomenology

no couplings to chiral quarks and leptons

$$\begin{aligned}
 & \frac{g_4}{\sqrt{2}} (\bar{Q}_{L4} \gamma^\mu L_{L4} + \text{H.c.}) U_{1\mu} \quad \text{Vector Leptoquark} \\
 & + \frac{g_4 g_s}{g_3} \left( \bar{Q}_{L4} \gamma^\mu T^a Q_{L4} - \frac{g_3^2}{g_4^2} \bar{Q}_{Li} \gamma^\mu T^a Q_{Li} \right) g_\mu^{\prime a} \quad \text{coloron} \quad Z' \\
 & + \frac{\sqrt{3} g_4 g_Y}{\sqrt{2} g_1} \left( \frac{1}{6} \bar{Q}_{L4} \gamma^\mu Q_{L4} - \frac{1}{2} \bar{L}_{L4} \gamma^\mu L_{L4} - \frac{g_1^2}{9g_4^2} \bar{Q}_{Li} \gamma^\mu Q_{Li} + \frac{g_1^2}{3g_4^2} \bar{L}_{Li} \gamma^\mu L_{Li} \right) Z'_\mu
 \end{aligned}$$

suppressed couplings to chiral quarks and leptons

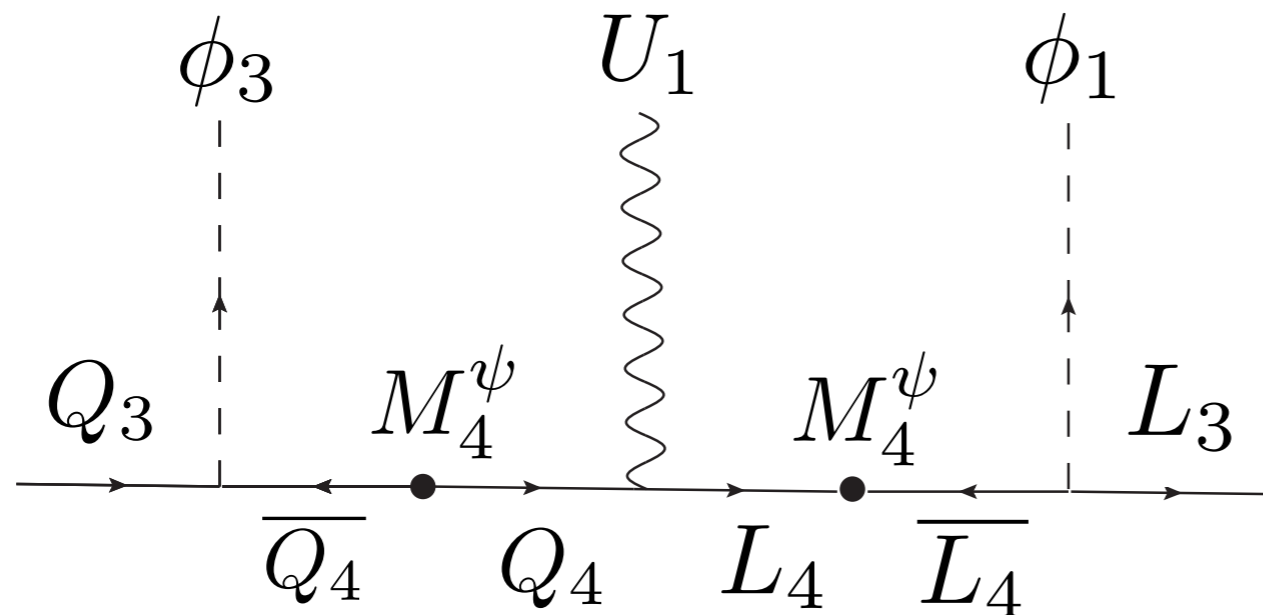
Typical Benchmark

$$g_4 \approx 3, \quad g_3 \approx g_s \approx 1, \quad g_1 \approx g_Y \approx 0.36,$$

$$M_{Z'} \approx 1.4 \text{ TeV}, \quad M_{U_1} \approx 1.6 \text{ TeV}, \quad M_{g'} \approx 2.0 \text{ TeV}$$

TeV masses

# Effective Vector Leptoquark $U_1$ couplings



Mixing generates  
couplings to  
chiral quarks and  
leptons

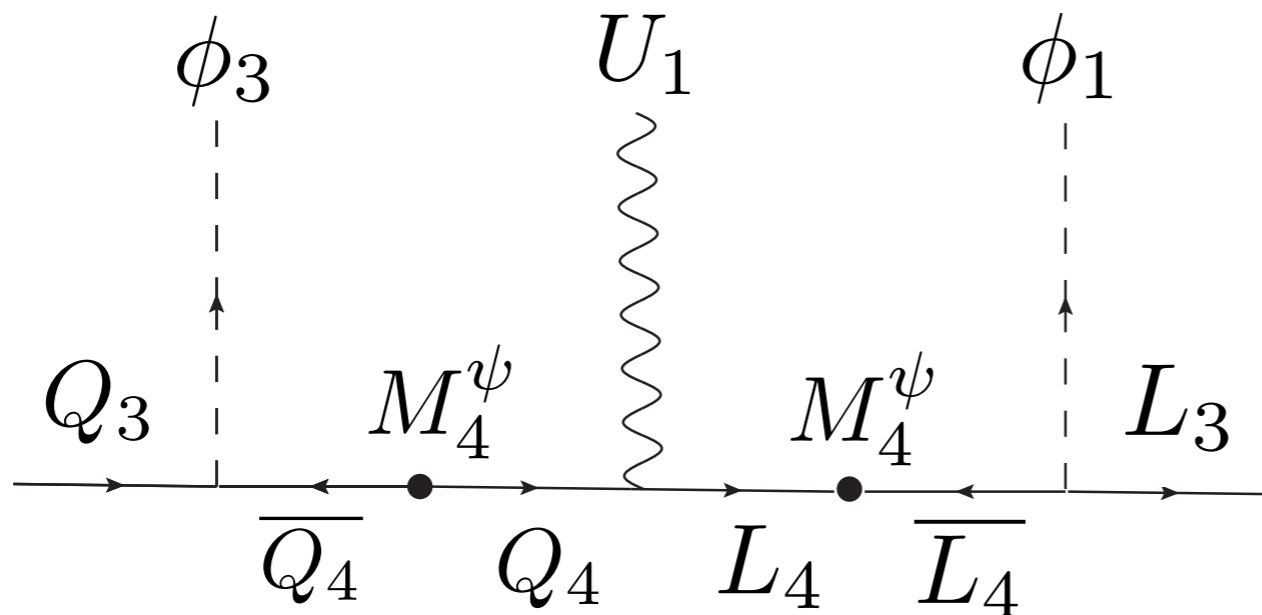
Third family only

$$\frac{g_4}{\sqrt{2}} \frac{x_{34}^\psi \langle \phi_1 \rangle}{M_4^\psi} \frac{x_{34}^\psi \langle \phi_3 \rangle}{M_4^\psi} \bar{Q}_{L3} \gamma^\mu L_{L3} U_{1\mu} \equiv g_U \bar{Q}_{L3} \gamma^\mu L_{L3} U_{1\mu}$$

After diagonalising mass matrices  $\rightarrow g_U \beta_{i\alpha} \bar{Q}_{Li} \gamma^\mu L_{L\alpha} U_{1\mu}$



# Effective Vector Leptoquark $U_1$ couplings



Mixing generates  
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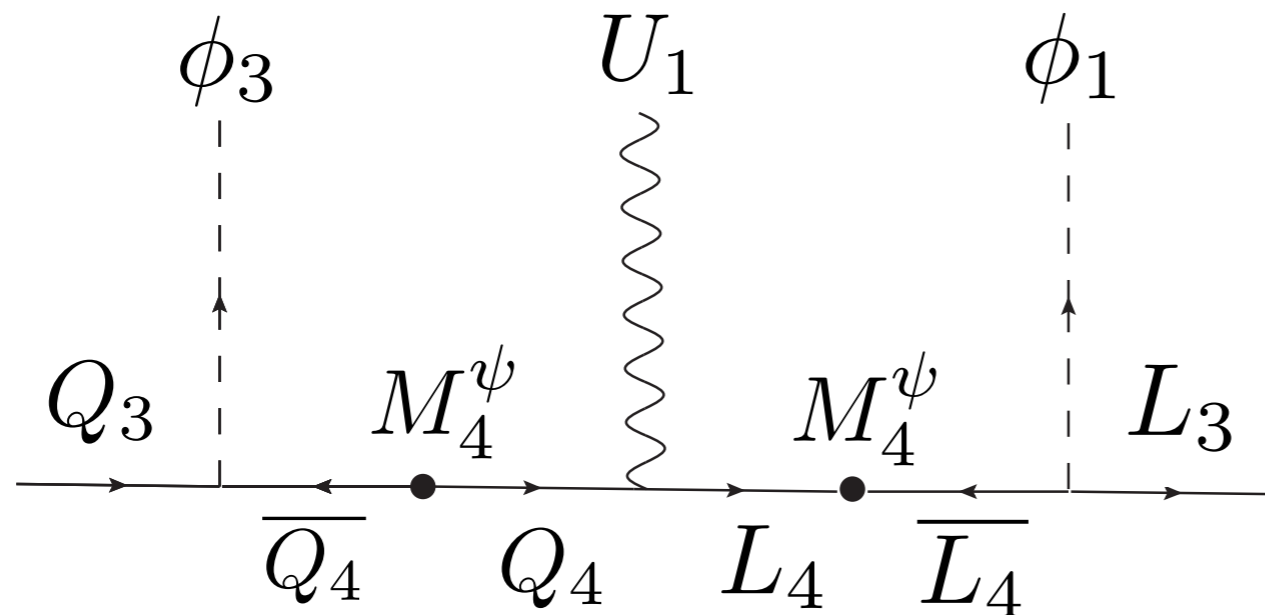
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After diagonalising mass matrices  $\rightarrow g_U \beta_{i\alpha} \bar{Q}_{Li} \gamma^\mu L_{L\alpha} U_{1\mu}$

$$\beta_{b\tau} \sim 1 \quad \beta_{c\nu_\tau} \sim \frac{m_c}{m_t} \quad \beta_{b\mu} \sim \frac{m_\mu}{m_\tau} \quad \beta_{s\mu} \sim \frac{m_\mu}{m_\tau} \frac{m_s}{m_b}$$

# Effective Vector Leptoquark $U_1$ couplings



Mixing generates couplings to chiral quarks and leptons

Third family only

$$\frac{g_4}{\sqrt{2}} \frac{x_{34}^\psi \langle \phi_1 \rangle}{M_4^\psi} \frac{x_{34}^\psi \langle \phi_3 \rangle}{M_4^\psi} \bar{Q}_{L3} \gamma^\mu L_{L3} U_{1\mu} \equiv g_U \bar{Q}_{L3} \gamma^\mu L_{L3} U_{1\mu}$$

After diagonalising mass matrices  $\rightarrow g_U \beta_{i\alpha} \bar{Q}_{Li} \gamma^\mu L_{L\alpha} U_{1\mu}$

$$\beta_{b\tau} \sim 1 \quad \begin{array}{c} \text{RD} \\ \begin{array}{ccc} b & U_1 & \nu_\tau \\ \tau & & c \end{array} \end{array} \quad \beta_{c\nu_\tau} \sim \frac{m_c}{m_t} \quad \beta_{b\mu} \sim \frac{m_\mu}{m_\tau} \quad \begin{array}{c} \text{RK} \\ \begin{array}{ccc} b & U_1 & \mu \\ \mu & & s \end{array} \end{array} \quad \beta_{s\mu} \sim \frac{m_\mu}{m_\tau} \frac{m_s}{m_b}$$

Too small c.f.  $R_D$  best fit value  $\beta_{c\nu_\tau} \sim 4V_{cb} \sim 0.16$

# Conclusion

- Fermion mass and mixing is a great mystery of SM, which may be related to recent B physics anomalies
- $R_K$  and  $R_D$  anomalies may be simultaneously explained by TeV scale vector leptoquark  $U_1$  from Pati-Salam
- Pati-Salam usually broken above PeV to avoid FCNCs so this requires non-trivial UV completion
- The twin PS model yields a TeV scale vector leptoquark **and** explains fermion masses (usually not addressed)
- But the predicted mass matrices do not yield large enough couplings to explain  $R_D$  given the current data