



UNIVERSITY OF
OXFORD

Taming distributional celestial amplitudes

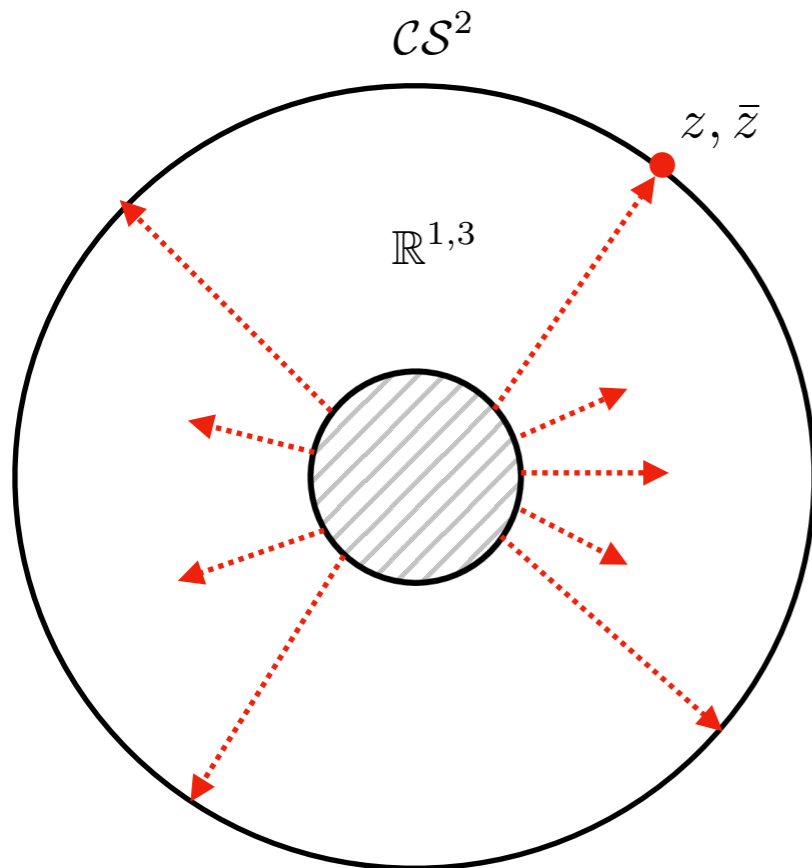
Corfu, 2021

Atul Sharma
Mathematical Institute
Oxford

Based on [Sharma '21]

Locality in the sky?

[Pasterski, Shao '17]



Momentum
Eigenstates

Boost
Eigenstates

$$|\omega, z, \bar{z}, \ell, \epsilon\rangle \longrightarrow |z, \bar{z}, h, \bar{h}, \epsilon\rangle = \int_0^\infty \frac{d\omega}{\omega} \omega^\Delta |\omega, z, \bar{z}, \ell, \epsilon\rangle$$

$$h = \frac{\Delta + \ell}{2}, \quad \bar{h} = \frac{\Delta - \ell}{2}$$

$$\epsilon = \pm 1$$

$$\mathcal{O}_{h, \bar{h}}^\epsilon(z, \bar{z})$$

Local?

Shadow
transform

$$\mathbf{S}[\mathcal{O}_{h, \bar{h}}^\epsilon](z, \bar{z}) = \int_{\mathbb{C}} \frac{d^2 w \mathcal{O}_{h, \bar{h}}^\epsilon(w, \bar{w})}{(w - z)^{2-2h} (\bar{w} - \bar{z})^{2-2\bar{h}}}$$

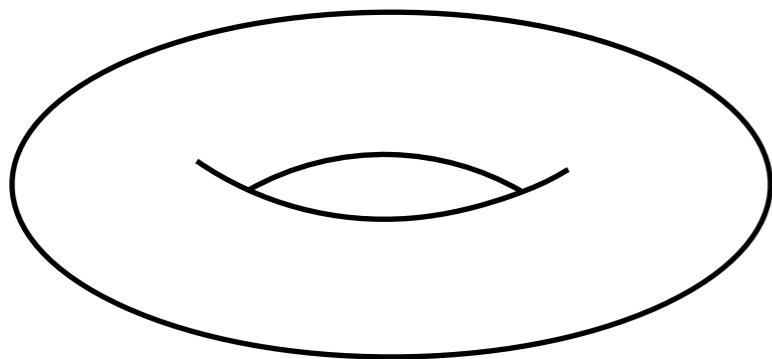
Bulk dual state

$$|z, \bar{z}, 1 - h, 1 - \bar{h}, \epsilon\rangle_{\mathbf{S}}$$

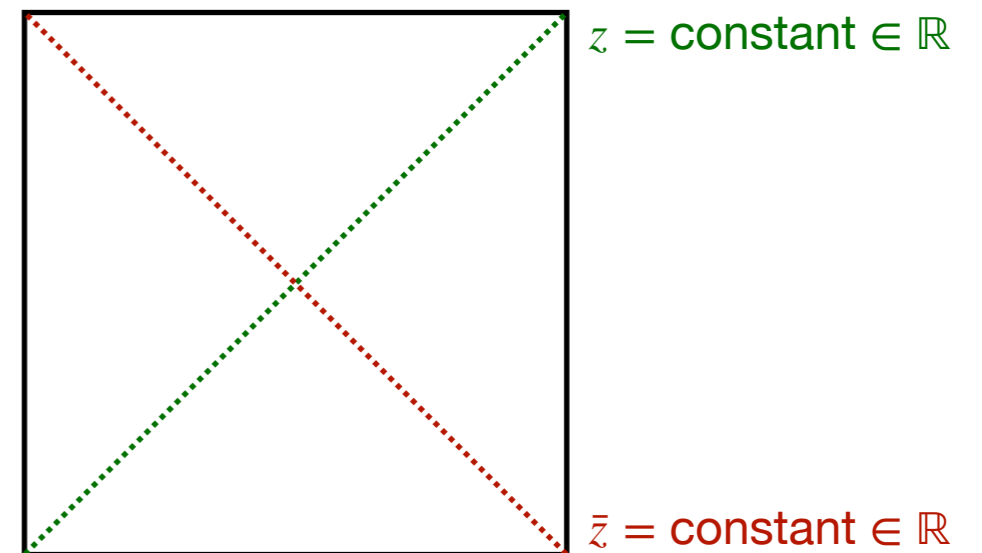
Light transforms

[Atanasov, Ball, Melton, Raclariu, Strominger '21]

Celestial torus
of $\mathbb{R}^{2,2}$



$$ds^2 = dz d\bar{z}$$



**Lights
transforms**

$$\mathbf{L}[\mathcal{O}_{h,\bar{h}}^\epsilon](z, \bar{z}) = \int_{\mathbb{R}} \frac{dw \mathcal{O}_{h,\bar{h}}^\epsilon(w, \bar{z})}{(w - z)^{2-2h}}$$

$$\bar{\mathbf{L}}[\mathcal{O}_{h,\bar{h}}^\epsilon](z, \bar{z}) = \int_{\mathbb{R}} \frac{d\bar{w} \mathcal{O}_{h,\bar{h}}^\epsilon(z, \bar{w})}{(\bar{w} - \bar{z})^{2-2\bar{h}}}$$

Bulk dual state

$$|z, \bar{z}, 1 - h, \bar{h}, \epsilon\rangle_{\mathbf{L}}$$

$$|z, \bar{z}, h, 1 - \bar{h}, \epsilon\rangle_{\bar{\mathbf{L}}}$$

Previously in celestial news

- Shadow transform \longrightarrow adjoint operation in 2d CCFT.

[Crawley, Miller, Narayanan, Strominger '21]

- Shadow transforms used for gluon conformal block expansions.

[Fan, Fotopoulos, Stieberger, Taylor, Zhu '21]

- Light transform exchanges observed in scalar conformal block expansions.

[Atanasov, Melton, Raclariu, Strominger '21]

- Light transform unravels symmetries of self-dual CCFT.

[Himwich, Pate, Singh '21] [Strominger '21] [Guevara, Himwich, Pate, Strominger '21]

Today...

- Light transformed celestial amplitudes
- Light transformed celestial OPE
- Connections with twistors

Ambidextrous choices

[Pasterski, Shao, Strominger '17]

$$\mathcal{A}_n(z_i, \bar{z}_i, \Delta_i, l_i, \epsilon_i) = \int_{\mathbb{R}_+^n} \prod_{j=1}^n \frac{d\omega_j}{\omega_j} \omega_j^{\Delta_j} \mathcal{A}_n(\epsilon_i \omega_i q_i, l_i) \delta^4\left(\sum_{j=1}^n \epsilon_j \omega_j q_j\right)$$

Celestial amplitude

Momentum space amplitude

$$q_{i\alpha\dot{\alpha}} = \begin{pmatrix} z_i \bar{z}_i & z_i \\ \bar{z}_i & 1 \end{pmatrix}_{\alpha\dot{\alpha}}$$

$$\epsilon_i = \pm 1$$

$$\mathcal{L}_n(z_i, \bar{z}_i, \Delta_i, l_i, \epsilon_i) = \int_{\mathbb{R}^n} \prod_{\bar{a}} \frac{d\bar{w}_{\bar{a}}}{(\bar{w}_{\bar{a}} - \bar{z}_{\bar{a}})^{2-2h_{\bar{a}}}} \prod_a \frac{dw_a}{(w_a - z_a)^{2-2h_a}} \times \mathcal{A}_n(z_{\bar{a}}, \bar{w}_{\bar{a}}, \Delta_{\bar{a}}, -, \epsilon_{\bar{a}}; w_a, \bar{z}_a, \Delta_a, +, \epsilon_a)$$

$a \longrightarrow$ Positive helicity
gluon/graviton

$\bar{a} \longrightarrow$ Negative helicity
gluon/graviton

Gluon examples

$$\mathcal{A}_n(i_{\Delta_i, l_i}^{\epsilon_i}) \equiv \mathcal{A}_n(z_i, \bar{z}_i, \Delta_i, l_i, \epsilon_i) \propto \delta(\beta), \quad \beta = i \sum_{j=1}^n (\Delta_j - 1)$$

2 points

$$\mathcal{A}_2(1_{\Delta_1, +1}^+, 2_{\Delta_2, -1}^-) = \delta(\beta) \delta(z_{12}) \delta(\bar{z}_{12})$$

$$z_{ij} = z_i - z_j$$

$$\begin{aligned} \Rightarrow \mathcal{L}_2(1_{\Delta_1, +1}^+, 2_{\Delta_2, -1}^-) &= \delta(\beta) \int_{\mathbb{R}^2} \frac{dw_1 d\bar{w}_2 \delta(w_1 - z_2) \delta(\bar{w}_1 - \bar{w}_2)}{(w_1 - z_1)^{1-\Delta_1} (\bar{w}_2 - \bar{z}_2)^{1-\Delta_2}} \\ &= \frac{\delta(\beta)}{z_{21}^{1-\Delta_1} \bar{z}_{12}^{1-\Delta_2}}. \end{aligned}$$

Gluon examples

3 points

$$\mathcal{A}_3(1_{\Delta_1, -1}^-, 2_{\Delta_2, -1}^-, 3_{\Delta_3, +1}^+) = \delta(\beta) \Theta\left(\frac{z_{13}}{z_{12}}\right) \Theta\left(\frac{z_{32}}{z_{12}}\right) \frac{\delta(\bar{z}_{13}) \delta(\bar{z}_{23})}{z_{12}^{-\Delta_3} z_{32}^{2-\Delta_1} z_{13}^{2-\Delta_2}}$$

$$\begin{aligned} \mathcal{L}_3(1_{\Delta_1, -1}^-, 2_{\Delta_2, -1}^-, 3_{\Delta_3, +1}^+) &= \delta(\beta) \int_{\mathbb{R}^3} \frac{d\bar{w}_1}{(\bar{w}_1 - \bar{z}_1)^{1-\Delta_1}} \frac{d\bar{w}_2}{(\bar{w}_2 - \bar{z}_2)^{1-\Delta_2}} \frac{dw_3}{(w_3 - z_3)^{1-\Delta_3}} \\ &\quad \times \Theta\left(\frac{z_1 - w_3}{z_1 - z_2}\right) \Theta\left(\frac{w_3 - z_2}{z_1 - z_2}\right) \frac{\delta(\bar{w}_1 - \bar{z}_3) \delta(\bar{w}_2 - \bar{z}_3)}{z_{12}^{-\Delta_3} (w_3 - z_2)^{2-\Delta_1} (z_1 - w_3)^{2-\Delta_2}} \\ &= \frac{\delta(\beta)}{z_{12}^{-\Delta_3} \bar{z}_{31}^{1-\Delta_1} \bar{z}_{32}^{1-\Delta_2}} \int_{z_2}^{z_1} \frac{\text{sgn}(z_{12}) dw_3}{(w_3 - z_3)^{1-\Delta_3} (w_3 - z_2)^{2-\Delta_1} (z_1 - w_3)^{2-\Delta_2}} \end{aligned}$$

$$\Rightarrow \mathcal{L}_3(1_{\Delta_1, -1}^-, 2_{\Delta_2, -1}^-, 3_{\Delta_3, +1}^+) = \frac{\delta(\beta) B(\Delta_1 - 1, \Delta_2 - 1) \text{sgn}(z_{12})}{z_{13}^{\Delta_1 - 1} z_{23}^{\Delta_2 - 1} \bar{z}_{31}^{1-\Delta_1} \bar{z}_{32}^{1-\Delta_2}}.$$

**Parity
conjugate**

$$\mathcal{L}_3(1_{\Delta_1, +1}^+, 2_{\Delta_2, +1}^+, 3_{\Delta_3, -1}^-) = \frac{\delta(\beta) B(\Delta_1 - 1, \Delta_2 - 1) \text{sgn}(\bar{z}_{12})}{z_{31}^{1-\Delta_1} z_{32}^{1-\Delta_2} \bar{z}_{13}^{\Delta_1 - 1} \bar{z}_{23}^{\Delta_2 - 1}}.$$

Observations

$$\mathcal{L}_2(1_{\Delta_1,+}^+, 2_{\Delta_2,-}^-) = \frac{\delta(\beta)}{z_{21}^{1-\Delta_1} \bar{z}_{12}^{1-\Delta_2}}.$$

$$\mathcal{L}_3(1_{\Delta_1,+}^+, 2_{\Delta_2,+}^+, 3_{\Delta_3,-}^-) = \frac{\delta(\beta) B(\Delta_1 - 1, \Delta_2 - 1) \text{sgn}(\bar{z}_{12})}{z_{31}^{1-\Delta_1} z_{32}^{1-\Delta_2} \bar{z}_{13}^{\Delta_1-1} \bar{z}_{23}^{\Delta_2-1}}.$$

- Ambidextrous prescription absorbs *all* momentum conserving delta functions.
- Brings amplitude to standard conformally covariant power law.
- Can read off light transform celestial OPE.

$$\begin{aligned} \mathbf{L}[O_{\Delta_1,+}^a](z_1, \bar{z}_1) \mathbf{L}[O_{\Delta_2,+}^b](z_2, \bar{z}_2) \\ \sim \text{sgn}(\bar{z}_{12}) B(\Delta_1 - 1, \Delta_2 - 1) f^{abc} \mathbf{L}[O_{\Delta_1+\Delta_2-1,+}^c](z_2, \bar{z}_2) + \dots \end{aligned}$$

Graviton examples

2 points

$$\mathcal{A}_2(1_{\Delta_1,+2}^+, 2_{\Delta_2,-2}^-) = \delta(\beta) \delta(z_{12}) \delta(\bar{z}_{12})$$

3 points

$$\mathcal{A}_3(1_{\Delta_1,+2}^+, 2_{\Delta_2,+2}^+, 3_{\Delta_3,-2}^-) = \delta(\beta + i) \Theta\left(\frac{\bar{z}_{13}}{\bar{z}_{12}}\right) \Theta\left(\frac{\bar{z}_{32}}{\bar{z}_{12}}\right) \frac{\delta(z_{13}) \delta(z_{23})}{\bar{z}_{12}^{-2-\Delta_3} \bar{z}_{32}^{2-\Delta_1} \bar{z}_{13}^{2-\Delta_2}}$$



$$\mathcal{L}_2(1_{\Delta_1,+2}^+, 2_{\Delta_2,-2}^-) = \frac{\delta(\beta)}{z_{21}^{-\Delta_1} \bar{z}_{12}^{-\Delta_2}}.$$

$$\mathcal{L}_3(1_{\Delta_1,+2}^+, 2_{\Delta_2,+2}^+, 3_{\Delta_3,-2}^-) = \frac{\delta(\beta + i) B(\Delta_1 - 1, \Delta_2 - 1) |\bar{z}_{12}|}{z_{31}^{-\Delta_1} z_{32}^{-\Delta_2} \bar{z}_{13}^{\Delta_1-1} \bar{z}_{23}^{\Delta_2-1}}.$$

**Celestial
OPE**

$$\mathbf{L}[G_{\Delta_1,+}](z_1, \bar{z}_1) \mathbf{L}[G_{\Delta_2,+}](z_2, \bar{z}_2) \sim |\bar{z}_{12}| B(\Delta_1 - 1, \Delta_2 - 1) \mathbf{L}[G_{\Delta_1+\Delta_2,+}](z_2, \bar{z}_2) + \dots$$

Comparing celestial OPE

[Pate, Raclariu, Strominger, Yuan '19] [Fan, Fotopoulos, Taylor '19]

For boost
eigenstates

$$O_{\Delta_1,+}^a(z_1, \bar{z}_1) O_{\Delta_2,+}^b(z_2, \bar{z}_2) \sim \frac{1}{z_{12}} B(\Delta_1 - 1, \Delta_2 - 1) f^{abc} O_{\Delta_1+\Delta_2-1,+}^c(z_2, \bar{z}_2) + \dots$$

$$G_{\Delta_1,+}(z_1, \bar{z}_1) G_{\Delta_2,+}(z_2, \bar{z}_2) \sim \frac{\bar{z}_{12}}{z_{12}} B(\Delta_1 - 1, \Delta_2 - 1) G_{\Delta_1+\Delta_2,+}(z_2, \bar{z}_2) + \dots$$

For light
transformed
states

$$\mathbf{L}[O_{\Delta_1,+}^a](z_1, \bar{z}_1) \mathbf{L}[O_{\Delta_2,+}^b](z_2, \bar{z}_2) \sim \text{sgn}(\bar{z}_{12}) B(\Delta_1 - 1, \Delta_2 - 1) f^{abc} \mathbf{L}[O_{\Delta_1+\Delta_2-1,+}^c](z_2, \bar{z}_2) + \dots$$

$$\mathbf{L}[G_{\Delta_1,+}](z_1, \bar{z}_1) \mathbf{L}[G_{\Delta_2,+}](z_2, \bar{z}_2) \sim |\bar{z}_{12}| B(\Delta_1 - 1, \Delta_2 - 1) \mathbf{L}[G_{\Delta_1+\Delta_2,+}](z_2, \bar{z}_2) + \dots$$

Non-singular OPE?



Twistor eigenstates

Spinor-helicity

$$p_{\alpha\dot{\alpha}} = \lambda_{\alpha} \bar{\lambda}_{\dot{\alpha}}, \quad \lambda_{\alpha}, \bar{\lambda}_{\dot{\alpha}} \in \mathbb{R}^2$$

Real and independent
in $\mathbb{R}^{2,2}$

Witten's
1/2 Fourier
transform

$$|Z, \ell\rangle_{\mathbb{T}} \equiv |\lambda, \mu, \ell\rangle_{\mathbb{T}} := \int_{\mathbb{R}^2} d^2\bar{\lambda} e^{i[\mu \bar{\lambda}]} |\lambda, \bar{\lambda}, \ell\rangle$$

↑
Twistor eigenstate

↑
Momentum eigenstate

$$\mu^{\dot{\alpha}} \xleftrightarrow{\text{Fourier conjugate}} \bar{\lambda}_{\dot{\alpha}}$$

[Witten '03]

[Mason, Skinner '09]

[Arkani-Hamed, Cachazo,

Cheung, Kaplan '09]

Little group

$$|r Z, \ell\rangle_{\mathbb{T}} = r^{-2\ell-2} |Z, \ell\rangle_{\mathbb{T}}, \quad r \in \mathbb{R}^*$$

Twistors

$$Z^A = (\mu^{\dot{\alpha}}, \lambda_{\alpha}) \in \mathbb{RP}^3$$

Primary twistor eigenstates

$$(\mu^{\dot{\alpha}}, \lambda_{\alpha}) \sim r (\mu^{\dot{\alpha}}, \lambda_{\alpha}) \xrightarrow{\text{Scale fixing}} \mu^{\dot{\alpha}} = \begin{pmatrix} \bar{z} \\ 1 \end{pmatrix}^{\dot{\alpha}}, \quad \lambda_{\alpha} = \epsilon \omega \begin{pmatrix} z \\ 1 \end{pmatrix}_{\alpha}$$

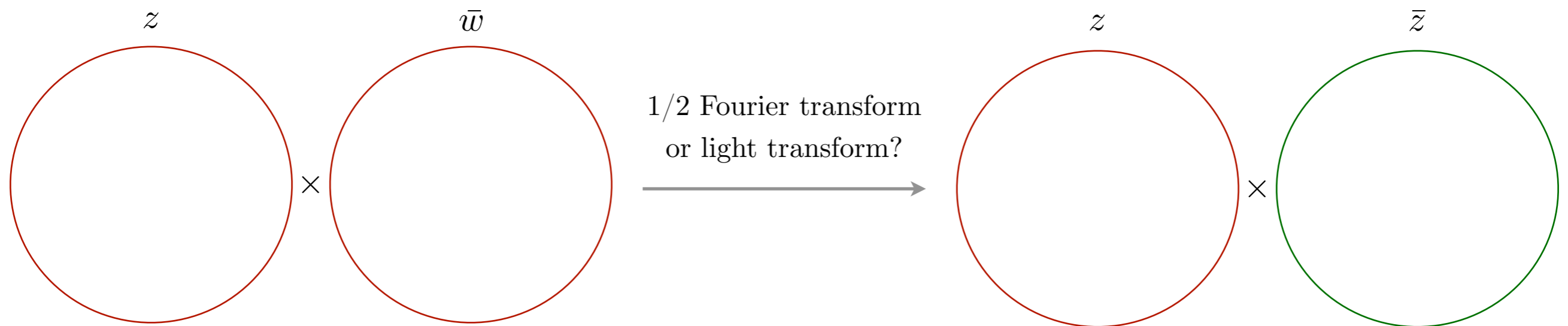
Conformal primary
twistor eigenstate

$$|z, \bar{z}, h, 1 - \bar{h}, \epsilon\rangle_{\mathbb{T}} := \int_0^{\infty} \frac{d\omega}{\omega} \omega^{2h} |\lambda, \mu, \ell\rangle_{\mathbb{T}} \stackrel{?}{=} |z, \bar{z}, h, 1 - \bar{h}, \epsilon\rangle_{\bar{\mathbb{L}}}$$

Point on twistor space
of AdS₃/Z slices

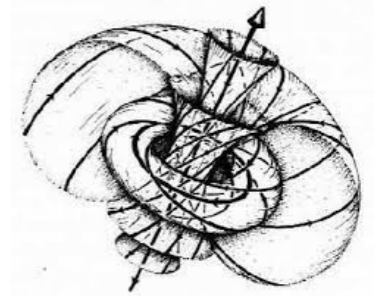
Transforms as conformal
primary with these weights

"1/2 Mellin
transform"

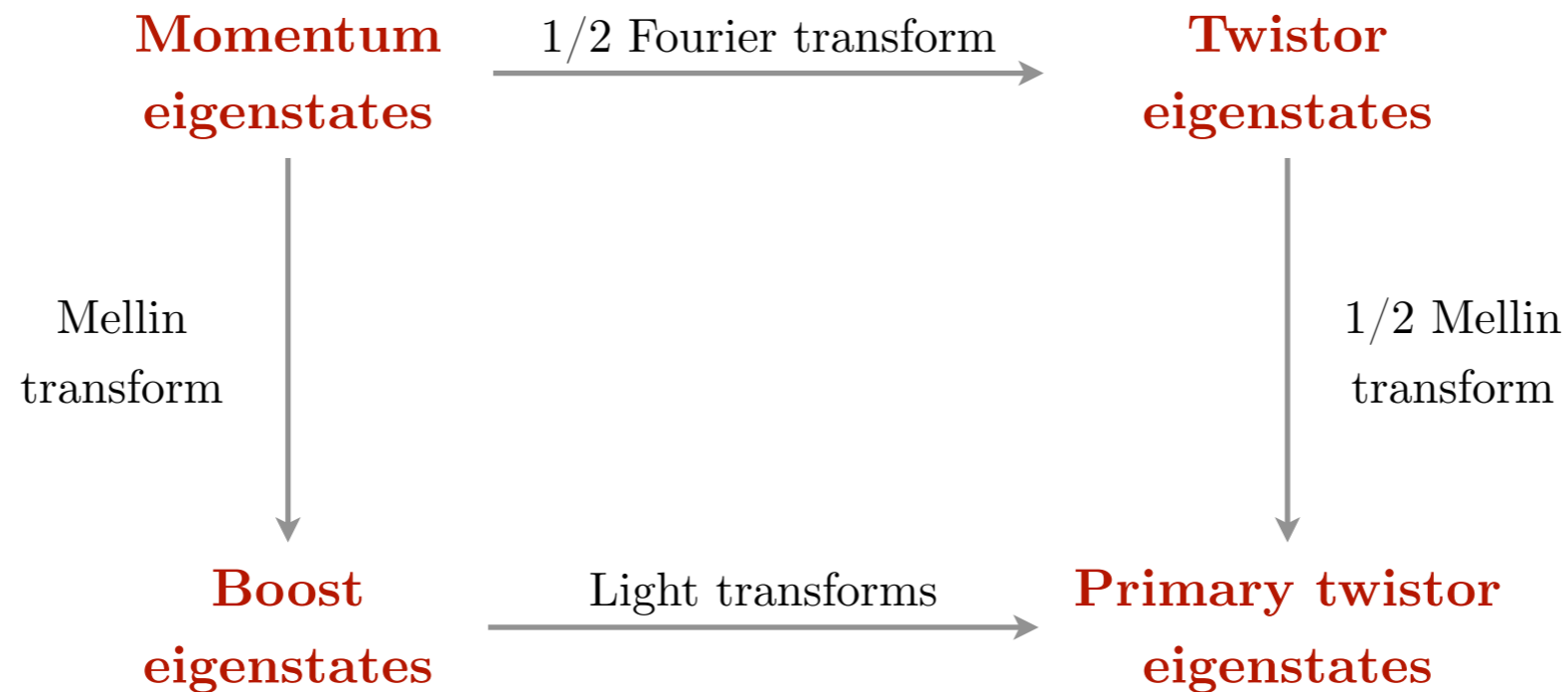


(Considered as an auxiliary
copy of the celestial torus.)

All roads lead to twistors...



$$|z, \bar{z}, h, 1 - \bar{h}, \epsilon\rangle_{\mathbb{T}} \\ = \epsilon^{-2\ell} \Gamma(2 - 2\bar{h}) \left(e^{\pi i(1 - \bar{h})} |z, \bar{z}, h, 1 - \bar{h}, \epsilon\rangle_{\bar{\mathbf{L}}} + e^{\pi i\bar{h}} |z, \bar{z}, h, 1 - \bar{h}, -\epsilon\rangle_{\bar{\mathbf{L}}} \right)$$



→ Original motivation for ambidextrous prescription.

Questions and speculations

- Light transformed boost eigenstates dual to local operators?
- In usual CFTs, correlators involving only light transformed local operators vanish. Reconcile with CCFT results? [Kravchuk, Simmons-Duffin '18]
- Maybe original boost eigenstates dual to *non-local* operators?
- Light transforms “refine” the transform to twistor eigenstates.
- Is celestial holography the origin of the success of twistor strings?
- Twistorial origins of $w_{1+\infty}$ symmetries in $(2, 2)$ signature in a light transform basis? [Strominger '21] (See Mason’s talk for story in Lorentzian signature.)