ASPECT OF NON-MINIMAL DERIVATIVE COUPLING

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INTRODUCTION

- There are many models have been proposed to describe the accelerated expansion of the universe.
- The derivative coupling of a scalar filed to the curvature can be considered as a generalized scalar-tensor.
- It can be extending a scalar-tensor theories by allowing the non-minimal couplings between the derivatives of the scalar fields and the Einstein tensor which giving more interesting behaviors.
- The power-law expansion, a x t^a, is a model that can remove flatness and horizon problems in a simplest inflationary model and it is also considering in specific gravity or dark energy model i.e. f(R) and f(G) gravities. It is also considering in a context of scalar field cosmology, phantom scalar field cosmology.



INTRODUCTION

- In this work, it is also assumed flat FLRW universe filled with a perfect fluid and scalar field.
- We use a power-law expansion to constrain the couplings constant in non-minimal derivative scalar field which there is non-minimal coupling between derivative of the scalar field and the Einstein tensor.
- We use the observational data from a combined WMAP9 data set (WMAP9+eCMB+BAO+H₀) to estimate (approximately) the value of the cosmological parameters and use them to find the value of the coupling constants.
- The signature using in this work is most positive and we set a speed of light c = 1.



Consider the theory of gravity with the action

$$S = \int d^4x \sqrt{-g} \left\{ \frac{R}{8\pi G} - \left[\varepsilon g_{\mu\nu} + \kappa G_{\mu\nu} \right] \phi^{\mu} \phi^{\nu} - 2V(\phi) \right\} + S_m \qquad (1)$$

• In the flat FLRW universe, we can write

$$3H^2 = 4\pi G\dot{\phi}^2 \left(\varepsilon - 9\kappa H^2\right) + 8\pi GV + 8\pi G\rho \tag{2}$$

$$2\dot{H} + 3H^{2} = -4\pi G \left[\varepsilon + \kappa \left(2\dot{H} + 3H^{2} + 4H\ddot{\phi}\dot{\phi}^{-1} \right) \right] + 8\pi G V - 8\pi G p \quad (3)$$

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$$\left(\varepsilon - 3\kappa H^2\right)\ddot{\phi} + 3H\left(\varepsilon - 3\kappa H^2 - 2\kappa\dot{H}\right)\dot{\phi} = -V_{,\phi} \quad (4)$$

where V_{ϕ} (= $dV/d\phi$) is a derivative with respect to scalar field ϕ .

 From above equations, we can extract the energy density and pressure of scalar field as follow,

$$\rho_{\phi} = \frac{1}{2} \dot{\phi}^2 \left(\varepsilon - 9\kappa H^2 \right) + V \tag{5}$$

$$p_{\phi} = \frac{1}{2}\dot{\phi}^2 \left(\varepsilon + 2\kappa\dot{H} + 3\kappa H^2 + 4\kappa H\ddot{\phi}\dot{\phi}^{-1}\right) - V \tag{6}$$

Considering the continuity equation of scalar filed

$$\dot{\rho}_{\phi} = -3H\rho_{\phi}(1+w_{\phi}) \tag{7}$$

• Substituting the energy density from (5) on the LHS of (7)

$$\dot{\rho}_{\phi} = -3H \left[\frac{1}{2} \dot{\phi}^2 \left(\varepsilon - 9\kappa H^2 \right) + V \right] (1 + w_{\phi}) \qquad (8)$$



- Comparing with the time derivative of the energy density from (5),

$$\dot{\rho}_{\phi} = \left(\varepsilon - 9\kappa H^2\right) \dot{\phi} \ddot{\phi} - 9\kappa H \dot{H} \dot{\phi}^2 + V_{,\phi} \dot{\phi} \qquad (9)$$

• We can obtain the equation of state (eos) parameter form

$$w_{\phi} = -\frac{\left(\varepsilon - 9\kappa H^{2}\right)\dot{\phi}\ddot{\phi} - 9\kappa H\dot{H}\dot{\phi}^{2} + V_{,\phi}\dot{\phi}}{3H\left(\frac{1}{2}\dot{\phi}^{2}\left(\varepsilon - 9\kappa H^{2}\right) + V\right)} - 1$$
⁽¹⁰⁾

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- By using Klein-Gordon (KG) equation (4), we have

$$\ddot{\phi} = -3H\dot{\phi} + \frac{6\kappa H\dot{H}}{\varepsilon - 3\kappa H^2}\dot{\phi} - \frac{V_{,\phi}}{\varepsilon - 3\kappa H^2} \tag{11}$$

• Then substituting back to obtain the eos parameter (10) becomes

$$w_{\phi} = \frac{\frac{1}{2}\dot{\phi}^{2}\left(\varepsilon - 9\kappa H^{2}\right)\left[2\kappa\dot{H}\left(\frac{\varepsilon + 9\kappa H^{2}}{\left(\varepsilon - 9\kappa H^{2}\right)\left(\varepsilon - 3\kappa H^{2}\right)}\right) + 1\right] - \frac{2\kappa H\dot{\phi}V_{,\phi}}{\left(\varepsilon - 3\kappa H^{2}\right)} - V$$

$$\frac{1}{2}\dot{\phi}^{2}\left(\varepsilon - 9\kappa H^{2}\right) + V$$
(12)

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• From the eos parameter, we obtain the scalar field pressure in a form

$$p_{\phi} = \frac{1}{2}\dot{\phi}^{2}\left(\varepsilon - 9\kappa H^{2}\right) \left[2\kappa\dot{H}\left(\frac{\varepsilon + 9\kappa H^{2}}{\left(\varepsilon - 9\kappa H^{2}\right)\left(\varepsilon - 3\kappa H^{2}\right)}\right) + 1\right] - \frac{2\kappa H\dot{\phi}V_{,\phi}}{\left(\varepsilon - 3\kappa H^{2}\right)} - V$$
(13)

- Taking time derivative to Friedmann equation (2), we have

$$\dot{H} = -\frac{4\pi G}{3H} \dot{\phi} \Big[3\kappa H \dot{H} \dot{\phi} - (\varepsilon - 9\kappa H^2) \ddot{\phi} + V_{,\phi} \Big] + \frac{4\pi G}{3H} \dot{\rho} \qquad (14)$$



• By using a continuity of pressure-less matter and using (11), we obtain

$$\dot{\phi}^{2} = \frac{\left(\varepsilon - 9\kappa H^{2}\right)V_{,\phi}\dot{\phi} + \frac{\left(\varepsilon - 3\kappa H^{2}\right)}{3H}V_{,\phi} + \frac{\left(\varepsilon - 3\kappa H^{2}\right)}{4\pi G}\left(-\dot{H} + 4\pi G\rho\right)}{3\kappa\dot{H}\left(\varepsilon - 3\kappa H^{2}\right) + \left(\varepsilon - 3\kappa H^{2} - 2\kappa\dot{H}\right)\left(\varepsilon - 9\kappa H^{2}\right)}$$
(15)



THE CONSTANT POTENTIAL

• We are interesting in the case of the scalar field's potential plays a role of a cosmological constant in a form

$$V \equiv \frac{\Lambda}{8\pi G} \tag{16}$$

where $\Lambda \ge 0$ and $V_{\phi} = 0$ in this case.

• Then the kinetic term of scalar field (15) becomes

$$\dot{\phi}^{2} = \frac{\frac{1}{4\pi G} \left(\varepsilon - 3\kappa H^{2}\right) \left(-\dot{H} + 4\pi G\rho\right)}{3\kappa \dot{H} \left(\varepsilon - 3\kappa H^{2}\right) + \left(\varepsilon - 3\kappa H^{2} - 2\kappa \dot{H}\right) \left(\varepsilon - 9\kappa H^{2}\right)}$$
(17)

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THE CONSTANT POTENTIAL

• The eos parameter becomes

$$\left\{ \left(-\dot{H} + 4\pi G\rho \right) \left[2\kappa \dot{H} \left(\varepsilon + 9\kappa H^2 \right) + \left(\varepsilon - 9\kappa H^2 \right) \left(\varepsilon - 3\kappa H^2 \right) \right] \right. \\
 w_{\phi} = \frac{-3\kappa \dot{H} \left(\varepsilon - 3\kappa H^2 \right) \Lambda - \left(\varepsilon - 3\kappa H^2 - 2\kappa \dot{H} \right) \left(\varepsilon - 9\kappa H^2 \right) \Lambda \right\}}{\left\{ \left(-\dot{H} + 4\pi G\rho \right) \left(\varepsilon - 3\kappa H^2 \right) \left(\varepsilon - 9\kappa H^2 \right) + 3\kappa \dot{H} \left(\varepsilon - 3\kappa H^2 \right) \Lambda \right. \\
 + \left(\varepsilon - 3\kappa H^2 - 2\kappa \dot{H} \right) \left(\varepsilon - 9\kappa H^2 \right) \Lambda \right\}$$
(18)



Now we consider the power-law expansion model

$$a \equiv a_0 \left(\frac{t}{t_0}\right)^{\alpha} \tag{19}$$

where α is a constant in a range $0 \le \alpha \le \infty$, *a* is scale factor at a cosmic time *t* and a_0 is a scale factor at a present time t_0 and set to unity in FLRW universe.

• From a definition of Hubble parameter $H = \dot{a}/a$, then

$$H = \frac{\alpha}{t}, \quad \dot{H} = -\frac{\alpha}{t^2} \tag{20}$$

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Substituting these parameter into eos parameter (18),

$$\begin{split} & \left\{ \left(\frac{\alpha t^{3\alpha} + 4\pi G \rho_0 t_0^{3\alpha} t^2}{t^{2+3\alpha}} \right) \left[-2\alpha \kappa \left(\varepsilon t^2 + 9\kappa \alpha^2 \right) + \left(\varepsilon t^2 - 9\kappa \alpha^2 \right) \left(\varepsilon t^2 - 3\kappa \alpha^2 \right) \right] \right. \\ & w_{\phi} = \frac{+3\alpha \kappa \left(\varepsilon t^2 - 3\kappa \alpha^2 \right) \Lambda - \left(\varepsilon t^2 - 3\kappa \alpha^2 + 2\kappa \alpha \right) \left(\varepsilon t^2 - 9\kappa \alpha^2 \right) \Lambda \right\}}{\left\{ \left(\frac{\alpha t^{3\alpha} + 4\pi G \rho_0 t_0^{3\alpha} t^2}{t^{2+3\alpha}} \right) \left(\varepsilon t^2 - 3\kappa \alpha^2 \right) \left(\varepsilon t^2 - 9\kappa \alpha^2 \right) - 3\kappa \alpha^2 \left(\varepsilon t^2 - 3\kappa \alpha^2 \right) \Lambda \right. \\ & \left. + \left(\varepsilon t^2 - 3\kappa \alpha^2 + 2\kappa \alpha \right) \left(\varepsilon t^2 - 9\kappa \alpha^2 \right) \Lambda \right\} \end{split}$$

$$(21)$$

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• Consider a cosmic time $t \to t_0$ then $w_{\phi} \to w_{\phi,0}$, $\Lambda \to \Lambda_0$ and $\alpha \to \alpha_0$ and the eos parameter becomes

$$\mathcal{N}_{\phi,0} = \frac{\left\{A_{0}\left[-2\alpha_{0}t_{0}^{2}\varepsilon\kappa\left(1+6\alpha_{0}\right)-9\alpha_{0}^{3}\kappa^{2}\left(2-3\alpha_{0}\right)+\varepsilon^{2}t_{0}^{4}\right]+3\alpha_{0}\kappa\left(\varepsilon t_{0}^{2}-3\kappa\alpha_{0}^{2}\right)\Lambda_{0}\right\}}{\left\{A_{0}\left[\varepsilon^{2}t_{0}^{4}+2\varepsilon\kappa\alpha_{0}t_{0}^{2}\left(1-6\alpha_{0}\right)-9\alpha_{0}^{3}\kappa^{2}\left(2-3\alpha_{0}\right)\right]\Lambda_{0}\right\}}\right. + \left[\varepsilon^{2}t_{0}^{4}+2\varepsilon\kappa\alpha_{0}t_{0}^{2}\left(1-6\alpha_{0}\right)-9\kappa^{2}\alpha_{0}^{3}\left(2-3\alpha_{0}\right)\right]\Lambda_{0}\right\}}$$

$$(22)$$



where A_0 is constant and

$$A_0 = \left(\frac{\alpha_0}{t_0^2} + 4\pi G\rho_0\right) \tag{23}$$

- To derive the coupling constant, κ , we use the Λ CDM data set derived from a combined WMAP9 data set (WMAP9+eCMB+BAO+ H_0) as shown in Table-I.
- We are considering into 2 cases,
 - $w_{\phi 0} = w_{\text{DE}} = -1$ and
 - $w_{\phi 0} = w_{\text{DE}} = -1.073$ (derived from a combined WMAP9 data set)

Table-I: Cosmological parameters from a combined WMAP9 (WMAP9+eCMB+BAO+ H_0) and derived parameters (in a lower part)

Parameter	WMAP9+eCMB+BAO+H ₀	
t_0 (Gyr)	$\frac{13.772 \pm 0.059}{\left[\left(4.352 \pm 0.019 \right) \times 10^{17} s \right]}$	
H_0 (km/s/Mpc)	$\frac{69.32 \pm 0.80}{\left[\left(2.246 \pm 0.026 \right) \times 10^{-18} \text{s}^{-1} \right]}$	
$\Omega_{\mathrm{m,0}}$	$0.2865^{+0.0096}_{-0.0095}$	
$\Omega_{\Lambda,0}$	$0.7135^{+0.0095}_{-0.0096}$	
$w_{\rm DE,0}$	$-1.073_{-0.089}^{+0.090}$	
$\rho_{\rm m,0}({\rm kg/m^3/s^2})$	$(2.5868^{+0.1164}_{-0.1148}) \times 10^{-27}$	
$\Lambda_0 (s^{-2})$	$(1.0797^{+0.0268}_{-0.0269}) \times 10^{-35}$	
α_0	0.977 ± 0.016	



Using data from Table-I we can write the eos parameter equation as

$$w_{\phi,0} = \frac{\left(-1.2449 \times 10^{35}\right)\varepsilon^2 + 6.8152\varepsilon\kappa - \left(1.1774 \times 10^{-34}\right)\kappa^2}{\left(6.5013 \times 10^{35}\right)\varepsilon^2 - 41.1782\varepsilon\kappa + \left(3.5314 \times 10^{-34}\right)\kappa^2}$$
(24)

• In the case of canonical field, $\varepsilon = 1$, $w_{\phi,0}$ becomes

$$w_{\phi,0} = \frac{-1.2449 \times 10^{35} + 6.8152\kappa - (1.1774 \times 10^{-34})\kappa^2}{6.5013 \times 10^{35} - 41.1782\kappa + (3.5314 \times 10^{-34})\kappa^2}$$
(25)



• For $w_{\phi,0} = w_{DE} = -1$, we can solve for the coupling constants and derived

$$\kappa_{c1} = 1.7362 \times 10^{34}$$

$$\kappa_{c2} = 1.2861 \times 10^{35}$$
• For $w_{\phi,0} = w_{\rm DE} = -1.073$, the coupling constants are

$$\kappa_{c1}' = 1.7469 \times 10^{34}$$
(27)
 $\kappa_{c2}' = 1.2561 \times 10^{35}$



• In the case of phantom field, $\mathcal{E} = -1$, $w_{\phi,0}$ becomes

$$w_{\phi,0} = \frac{-1.2449 \times 10^{35} - 6.8152\kappa - (1.1774 \times 10^{-34})\kappa^2}{6.5013 \times 10^{35} + 41.1782\kappa + (3.5314 \times 10^{-34})\kappa^2}$$
(28)



• For $w_{\phi,0} = w_{DE} = -1$, we can solve for the coupling constants and derived

$$\kappa_{p1} = -1.2861 \times 10^{35}$$

$$\kappa_{p2} = -1.7362 \times 10^{34}$$
• For $w_{\phi 0} = w_{\rm DE} = -1.073$, the coupling constants are
$$\kappa'_{p1} = -1.2561 \times 10^{35}$$

$$\kappa'_{p2} = -1.7469 \times 10^{34}$$
(26)
(27)



CONCLUSIONS

- In this work, we using the non-minimal derivative coupling which there is a coupling between the derivative of scalar fields and the Einstein tensor.
- We assumed the flat FLRW universe filled with a perfect fluids and set c = 1.
- We used the power-law expansion to constraint the cosmological parameters and try to find the coupling constants.
- To derived the values of the coupling constants, we used the derived parameters from a combined WMAP9 data set (WMAP9+eCMB+BAO+ H_0).
- In the case of canonical field ($\varepsilon = 1$) and phantom field ($\varepsilon = -1$) we obtained the coupling constants with the same magnitude but opposite sign for the same value of the eos parameter, $w_{\phi,0} = -1$ and $w_{\phi,0} = -1.073$.
- The coupling constants are shown in Table-II.



CONCLUSIONS

Table-II: The coupling constants derived from the power-law combined with the non-minimal derivative coupling model by using the combined WMAP9 data set

	Canonical ($\varepsilon = 1$)	Phantom (ε = -1)
$w_{\phi,0} = -1$	$\kappa_{c1} = 1.7362 \times 10^{34}$	$\kappa_{p1} = -1.2861 \times 10^{35}$
	$\kappa_{c2} = 1.2861 \times 10^{35}$	$\kappa_{p2} = -1.7362 \times 10^{34}$
$w_{\phi,0} = -1.073$	$\kappa_{c1}' = 1.7469 \times 10^{34}$	$\kappa'_{p1} = -1.2561 \times 10^{35}$
	$\kappa_{c2}' = 1.2561 \times 10^{35}$	$\kappa'_{p2} = -1.7469 \times 10^{34}$



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