

Supersymmetrizing Massive Gravity

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Outline

- Motivation
- Forming matter action
- Coupling to Supergravity
- Main results
- Summary

MOTIVATION

- Forming massive gravity via Higgs mechanism (by Chamseddine and Mukhanov)
- Four scalar fields with global Lorentz symmetry coupled to gravity
- Scalar fields taking a vacuum expectation value breaking diffeomorphism invariance spontaneously
- Result massive graviton

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FOUR CHIRAL SUPERFIELDS

- Consider four $N = 1$ chiral superfields with global Lorentz symmetry

$$\bar{D}_{\dot{\alpha}} \Phi^A(x, \theta, \bar{\theta}) = 0$$

- Given by

$$\begin{aligned} \Phi_A = & \varphi_A + i(\theta\sigma^\mu\bar{\theta})\partial_\mu\varphi_A - \frac{1}{4}\theta\theta\bar{\theta}\bar{\theta}\partial_\mu\partial^\mu\varphi_A + \sqrt{2}\theta\psi_A \\ & - \frac{i}{\sqrt{2}}\theta\theta(\partial_\mu\psi_A\sigma^\mu\bar{\theta}) + \theta\theta F_A \end{aligned}$$

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BASIC FIELD

- In bosonic case: $H_{AB} = g^{\mu\nu} \partial_\mu \phi_A \partial_\nu \phi_B$
- Our case: Define H_{ABC} as the basic field

$$H_{ABC} = D^\alpha \phi_A (\sigma_B)_{\alpha\dot{\alpha}} \bar{D}^{\dot{\alpha}} \phi_C^* = D\phi_A \sigma_B \bar{D}\phi_C^* \quad (1)$$

- Work instead with $\bar{H}_{ABC} = H_{ABC} - DX_A \sigma_B \bar{D}X_C^*$ to avoid including higher order terms

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FORMING OUR ACTION

- Using this \bar{H}_{ABC} construct D-type terms to form the matter multiplets ($\bar{H}_{ABC}\bar{H}_{BCA},\dots$)
- Couple to the Supergravity Lagrangian using the rules of tensor calculus.

$$L_{S.G} = -\frac{e}{2\kappa^2}R(e, w) - \frac{e}{3}|u|^2 + \frac{e}{3}A_\mu A^\mu - \frac{1}{2}\bar{\phi}_\mu R^\mu \quad (2)$$

- Lagrangian field content:
 - spin-2 field, $e_{a\mu}$
 - spin-3/2 field, ϕ_μ
 - auxiliary fields u, A_μ

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COUPLING VECTOR MULTIPLY TO SUPERGRAVITY

- Vector multiplet

$$V = (C, \xi, M, V_\mu, \lambda, D) \quad (3)$$

- D-type multiplet action formula

$$\begin{aligned}
 e^{-1} L_D = & D + \frac{i\kappa}{2} \bar{\phi}_\mu \gamma^5 \gamma^\mu \lambda - \frac{\kappa}{3} (uM^* + u^*M) + \frac{i\kappa^2}{8} \epsilon^{\mu\nu\rho\sigma} \bar{\phi}_\mu \gamma_\nu \phi_\rho \bar{\xi} \phi_\sigma \\
 & + \frac{2}{3} \kappa V_\mu \left(A^\mu + \frac{3}{8} i e^{-1} \epsilon^{\mu\rho\sigma\tau} \bar{\phi}_\rho \gamma_\tau \phi_\sigma \right) - i \frac{\kappa}{3} e^{-1} \bar{\xi} \gamma^5 \gamma_\mu R^\mu \\
 & - \frac{2}{3} \kappa^2 C e^{-1} L_{S.G.} + e^{-1} L_{S.G.} \quad (4)
 \end{aligned}$$

[P. Nath and R. Arnowitt and A. Chamseddine: Applied N=1 Supergravity, World Scientific Publishing, (1984)]

[E. Cremmer, S. Ferrara, L. Girardello and A. Van Proeyen, Nuclear Physics B, 212, (1983), 413]

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WHAT WE WANT?

- Expanding the fields around the vacuum solution

$$\varphi^A = x^A + \chi^A, \quad e_A^\mu = \delta_A^\mu + \bar{e}_A^\mu, \quad (5)$$

- An action with the following required conditions

- A Fierz-Pauli term for the vierbeins ($\bar{e}_A^\mu \bar{e}_\mu^A - \bar{e}^2$)
- No linear vierbein term
- Maxwell form for the χ_A fields

$$I (\partial_\mu \chi_A \partial^\mu \chi^{A*} - \partial_A \chi^A \partial_B \chi^{B*}) \quad (6)$$

where I is a constant

- Ghost free: no terms like

$$\partial_\mu \chi_A \partial^\mu \chi^A, \quad \text{or} \quad \partial_A \chi^A \partial_B \chi^B \quad (7)$$

- Massive gravitinos

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MATTER ACTION

- First four conditions well satisfied with three D-type terms:
 - $\bar{H}_{ABC}\bar{H}_{BCA}$
 - $\bar{H}_{ABB}\bar{H}_{CCA}$
 - $\bar{H}_{AB}\bar{H}_{AB}^*$ where $H_{AB} = D\Phi_A D\Phi_B$
- Gravitino massive by adding two F-type terms and coupling them to supergravity:
 - $\bar{D}^2 (D\Phi_A \sigma^{AB} D\Phi_B)$
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COUPLING CHIRAL MULTIPLY TO SUPERGRAVITY

- Chiral multiplet (F-type)

$$F = (z, X_L, h) \quad (8)$$

- Complex scalar field z
- Left-handed Weyl spinors X_L
- Complex auxiliary field h

- Action formula

$$e^{-1} L_F = h + \kappa UZ + \kappa \bar{\phi}_\mu \gamma^\mu \chi + i \kappa^2 \bar{\phi}_\mu \gamma^{\mu\nu} \phi_{\nu R} Z + h.c. \quad (9)$$

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- Global supersymmetry promoted to a local one using the rules of tensor calculus
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MAIN RESULTS

- Global Lorentz index A identified with the space-time Lorentz index
 - Scalar fields $\chi_A \rightarrow$ vectors
 - Chiral spinors $\psi_A \rightarrow$ spin-3/2 Rarita-Schwinger fields
- Spectrum of the model in the broken phase:
 - A massive spin-2 field
 - Two massive spin-3/2 fields with different masses
 - A massive vector

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DEGREES OF FREEDOM: BEFORE COUPLING

- Counting degrees of freedom: supergravity coupled to a $N = 1$ supersymmetry model similar to the Wess-Zumino model
- Before the coupling, supergravity contains
 - one massless spin-2 graviton (two bosonic dof)
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DEGREES OF FREEDOM: BEFORE COUPLING

- $N = 1$ supersymmetry model:
 - Four spin-0 particles, φ^A , with only six degrees of freedom (3 times 2) since φ^0 decouples due to Fierz-Pauli choice
 - Six fermionic degrees of freedom forming a multiplet
- Overall eight fermionic degrees of freedom and eight bosonic degrees of freedom

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DEGREES OF FREEDOM: AFTER COUPLING

- After coupling to supergravity: $N = 1$ massive representation
- Bosonic degrees of freedom
 - A massive spin-2 particle, with five degrees of freedom
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- Fermionic degrees of freedom
 - Two massive spin-3/2 particles, ϕ_μ and ψ_A with four degrees of freedom each
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- Same number of degrees of freedom as before coupling

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SUMMARY

- Constructed a supersymmetric theory of massive gravity
- Coupled a pure Supergravity multiplet to a matter multiplet
- Left with two massive spin-3/2 particles, ϕ_μ and ψ_A with completely different masses
- Similar to the $N = 2$ supersymmetry in which we have two gravitinos, but there they have the same mass.

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- Supersymmetry completely broken exactly at the same scale as the diffeomorphism breaking
- Before diffeomorphism breaking:
spin-1/2 and not spin-3/2 \rightarrow two gravitinos with different masses
- ϕ_μ : genuine gravitino
- ψ_A : identified with a gravitino after the breaking

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