# On Lorentz violation in superluminal theories

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- Motivation: modified gravity theories
- A toy model with superluminality
- Instabilities and Lorentz violation
- Pathologies in stress-energy tensor
- Conclusions

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# dRGT Massive Gravity

Features

- Free of Boulware-Deser ghosts
- Allows self-accelerated cosmological solution

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# dRGT Massive Gravity

## Features

- Free of Boulware-Deser ghosts
- Allows self-accelerated cosmological solution

# Galileon theory

#### Features

- Ghost-free
- Self-accelerated cosmological solution
- Galilean Genesis as alternative to inflation

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# Modified gravity theories: problems

Presence of superluminal modes on non-trivial backgrounds.

Key moment: number of derivatives in self-coupling should be greater than 2; the sign of the coupling should be minus.

Adams et al., 2006

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Example: Galileon theory

$$\mathcal{L} = \partial_\mu \pi \ \partial^\mu \pi + rac{c}{\Lambda^4} \partial^2_\mu \pi (\partial_
u \pi)^2 + ... \ ; \quad c < 0$$

Superluminality itself looks not so bad...

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- Possible problems with causality (a principal possibility to make a closed time-like curves)
- Problems with UV-completion (scattering amplitudes do not satisfy S-matrix analyticity axioms)

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• Problems with UV-completion (scattering amplitudes do not satisfy S-matrix analyticity axioms)

Superluminal effective theories could not be UV-completed to local QFT or unitary string theory!

What is the particular mechanism of Lorentz breaking? Look for a physical difference between static and boosted classical solutions.

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What is the particular mechanism of Lorentz breaking?

Look for a physical difference between static and boosted classical solutions.

Toy model with superluminality in  $D{=}2$ 

$$S = \int d^2 x \left[ \frac{1}{2} (\partial_\mu \phi \, \partial^\mu \phi) - \frac{m^2}{2} \phi^2 - \frac{1}{4\Lambda^2} (\partial_\mu \phi \, \partial^\mu \phi)^2 + M \delta(x) \phi \right]$$

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Consider classical perturbations on a background of a static solution. It is enough to consider EOM without non-linear  $\phi$  term.

$$\phi'' - m^2 \phi = -M\delta(x) \Rightarrow \tilde{\phi} = \frac{M}{2m} e^{-m|x|}$$

Consider quadratic action for classical perturbations  $\phi = \tilde{\phi} + \xi$ . Collect terms up to the order  $\frac{1}{\Lambda^2}$ 

$$S_{\xi} = \int d^2 x \left[ \frac{1}{2} Z^{\mu\nu} \partial_{\mu} \xi \, \partial_{\nu} \xi - \frac{m^2}{2} \xi^2 \right];$$

$$Z^{\mu\nu} = \begin{pmatrix} 1 + \frac{M^2}{4\Lambda^2} e^{-2m|x|} & 0\\ 0 & -1 - \frac{3M^2}{4\Lambda^2} e^{-2m|x|} \end{pmatrix}$$

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Dispersion relation

$$\omega_{\pm} = \pm \sqrt{-rac{Z^{11}}{Z^{00}}} \, k \quad \Rightarrow \quad \left| rac{d\omega}{dk} 
ight| > 1$$

Group velocity of the excitations exceeds the speed of light.

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# Looking for a classical instability

In a static frame

 $\det Z^{\mu
u}_{static} < 0 \quad \Rightarrow ext{the solution } ilde{\phi} ext{ is classically stable}$ 

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#### Looking for a classical instability

#### In a static frame

det  $Z_{static}^{\mu\nu} < 0 \implies$  the solution  $\tilde{\phi}$  is classically stable In a boosted frame  $(x \rightarrow \gamma | x + \beta t |)$ 

$$Z_{boost}^{\mu\nu} = \begin{pmatrix} 1 - \frac{M^2}{4\Lambda^2} \gamma^2 (3\beta^2 - 1) e^{-2m\gamma|\mathbf{x} + \beta t|} & \frac{M^2}{2\Lambda^2} \gamma^2 \beta e^{-2m\gamma|\mathbf{x} + \beta t|} \\ \frac{M^2}{2\Lambda^2} \gamma^2 \beta e^{-2m\gamma|\mathbf{x} + \beta t|} & -1 - \frac{M^2}{4\Lambda^2} \gamma^2 (3 - \beta^2) e^{-2m\gamma|\mathbf{x} + \beta t|} \end{pmatrix}$$

However, det 
$$Z_{boost}^{\mu\nu} = \det \Lambda_{\rho}^{\mu} \Lambda_{\sigma}^{\nu} Z_{static}^{\rho\sigma} = \det Z_{static}^{\mu\nu}$$
  
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The solution  $\tilde{\phi}$  is classically stable in any frame.

# Looking for a quantum instability

As  $Z^{00}$  change its sign in a boosted frame  $\Rightarrow$  Ghosts appear!

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#### Looking for a quantum instability

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Looking for UV dynamics; we can take m = 0 and regard boosted background as a constant

$$S_{\xi} = \int d^2 x \left[ rac{1}{2} Z^{\mu
u} \partial_{\mu} \xi \ \partial_{
u} \xi 
ight]; \ Z^{\mu
u} = egin{pmatrix} Z^{00} & Z^{01} \ Z^{01} & Z^{11} \end{pmatrix} = ext{const}$$

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Add interaction with matter scalar field  $\chi$ .

$$S_{\chi} = \int d^2 x \left[ rac{1}{2} (\partial_{\mu} \chi \, \partial^{\mu} \chi) - rac{m_{\chi}^2}{2} \chi^2 + rac{g}{2} \xi \chi \chi 
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ight]$$

Triples  $\xi \chi \chi$  created from vacuum as  $Z^{00} < 0$ Required hierarchy of parameters:

$$m_{\chi} < m\gamma \ll k \ll \Lambda$$
;  $m \ll M$ ;  $g \frac{M}{m} \ll m_{\chi}^2$ 

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Estimations for background decay rate and energy loss rate Integral saturated in IR

$$rac{dE}{dt}\sim \Gamma\sim rac{g^2}{2\pi m^2\gamma^2}\lesssim rac{m_\chi^4}{\Lambda^2}$$

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Small, but non-zero quantity  $\Rightarrow$  Lorentz invariance is violated!

Naive estimation for 4D

$$S_{\chi} = \int d^2 x \left[ rac{1}{2} (\partial_{\mu} \chi \, \partial^{\mu} \chi) - rac{m_{\chi}^2}{2} \chi^2 + rac{g'}{2} \xi (\partial_{\mu} \chi \partial^{\mu} \chi) 
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Estimation for background decay rate Integral saturated in UV

$$\Gamma \sim {{g'}^2\over 8\pi} \Lambda^2$$

Estimation for energy loss rate

$$rac{dE}{dt}\sim rac{{g'}^2}{8\pi}rac{\Lambda^3}{m\gamma}$$

This can be very high and probably could destroy the classical solution.

Expect some features of  $T^{\tilde{\phi}}_{\mu\nu}(\xi)$  as  $\gamma \to \gamma_{critical}$ Below  $\gamma_{critical}$ ,  $T^{\tilde{\phi}^{boosted}}_{\mu\nu}(\xi) = T^{\tilde{\phi}^{static}}_{\mu\nu}(\xi)$ 

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Looking for accelerated solution  $\tilde{\phi}$  with velocity  $\beta=t/\tau\simeq 1$ 

$$Z^{\mu
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Solution of EOM  $\partial_{\mu}[Z^{\mu\nu}\partial_{\nu}\xi] = 0$  with ansatz  $\xi = \xi(t) e^{ikx}$ 

$$\xi(t,x) = c_1 e^{-ik(t-x)} + c_2 e^{-ik(t-x-\alpha \log \frac{\alpha+t}{\alpha-t})}$$

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We quantize this excitations considering  $Z^{\mu\nu}$  as "curved" background. Then we calculate renormalized  $T_{\mu\nu}$ 

$$\langle T_{\mu\nu}^{\text{ren.}} \rangle = \lim_{x \to x'} \mathcal{D}_{\mu\nu} [G_{\text{ex.}}^{(1)}(x, x') - G_{\text{pert.}}^{(1)}(x, x')]; \quad G_{\text{ex.}}^{(1)}(x, x') = \langle [\xi(x), \xi(x')] \rangle$$

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Using Schwinger-DeWitt technique we get

$$G_{ren.}^{(1)} = \lim_{x \to x'} [G_{ex.}^{(1)}(x, x') - G_{pert.}^{(1)}(x, x')] = \log \left[\frac{(\alpha^2 - t'^2)^2}{(\alpha^2 + t'^2)\alpha^2}\right] + \text{regular terms}$$

Indication for pathology as  $t' \rightarrow \alpha$ , although renormalization of  $T_{\mu\nu}$  is needed

Work in progress...

- Physical solution in the simple superluminal quantum field theory shows different behaviour in static and boosted frames. I.e. Lorentz invariance is broken.
  - Theory is pathological and we don't need to consider it.
  - As the theory is non-Lorentzian we need to complete a Lagrangian with other Lorentz-braking terms.
  - It is a spontaneous breaking of Lorentz invariance, so we need to consider a phase with a broken symmetry.

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  - As the theory is non-Lorentzian we need to complete a Lagrangian with other Lorentz-braking terms.
  - It is a spontaneous breaking of Lorentz invariance, so we need to consider a phase with a broken symmetry.
- This analysis seem to be suitable for other theories that have superluminal propagation on classical backgrounds.

# Thank you!