

# Dark Matter and Dark Energy interaction models

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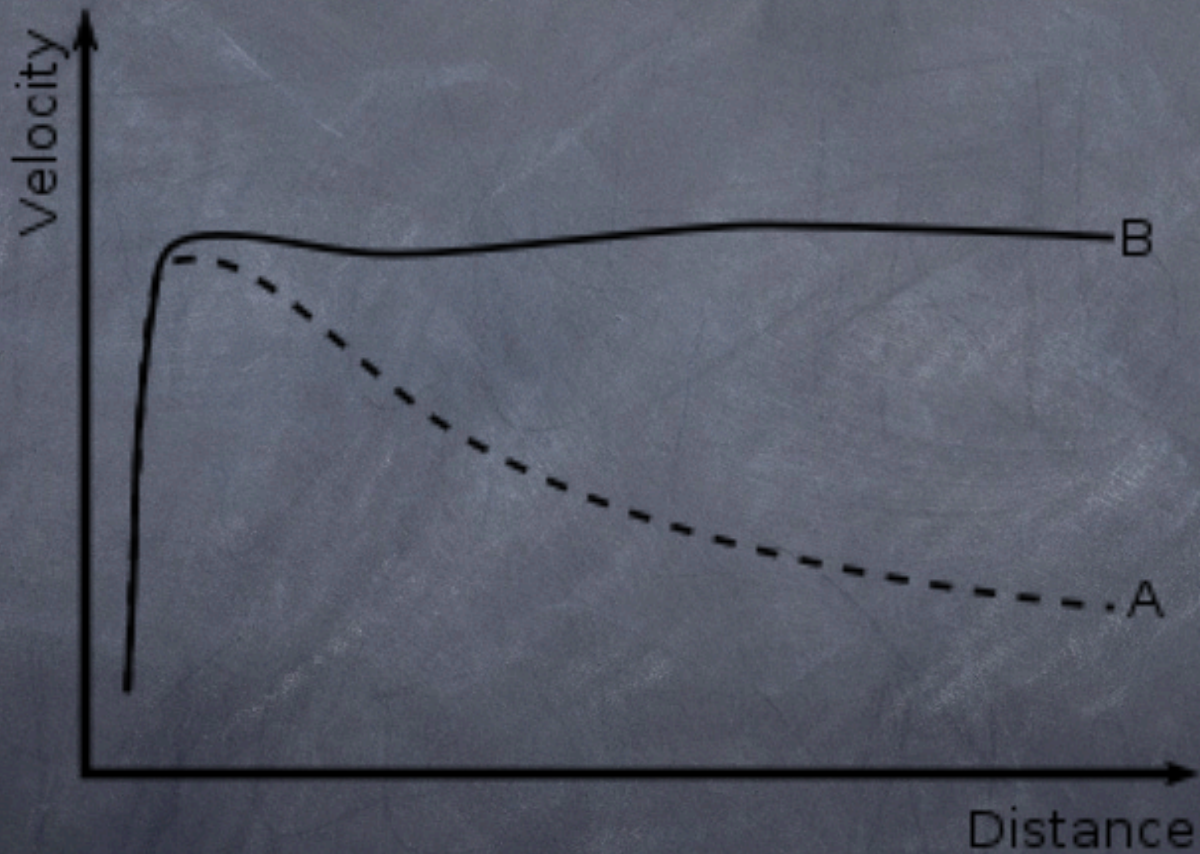
# Dark Matter

- Since Zwicky (in the thirties Coma Cluster)

$$V^2 = \frac{GM}{R}$$



# Dark Matter



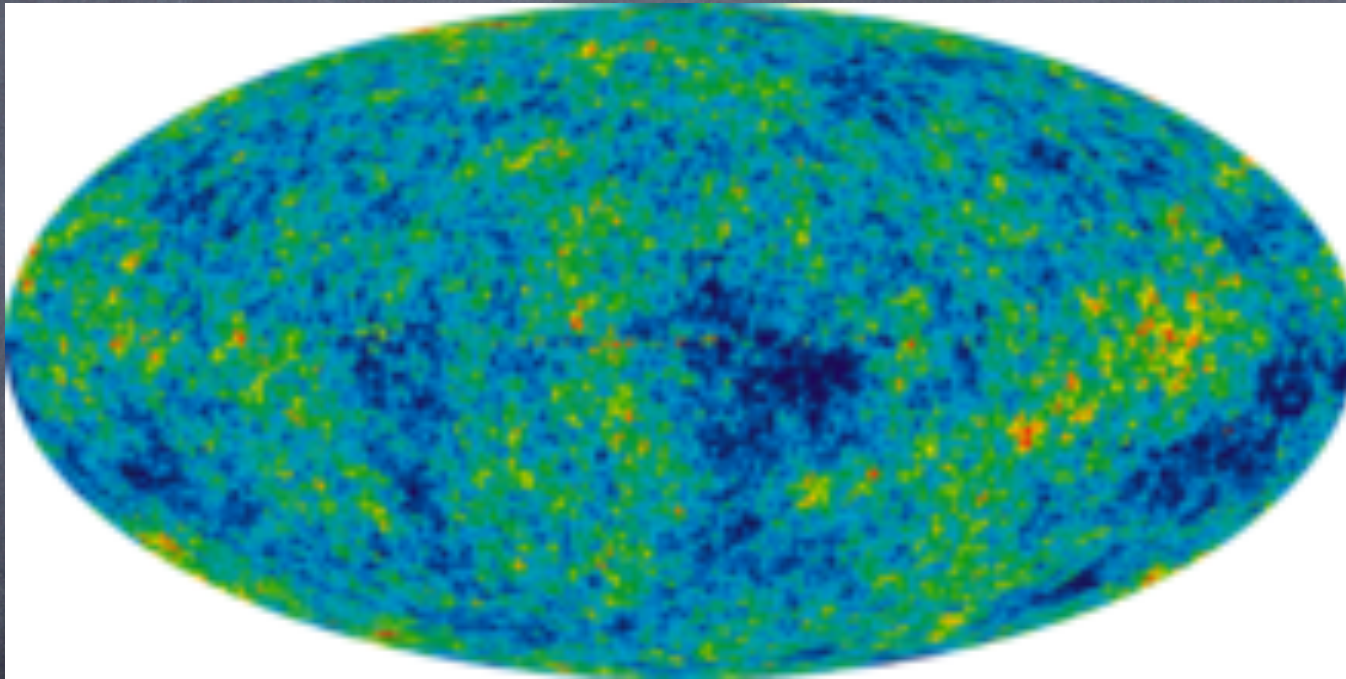


# Dark Energy

- Today's accelerated phase
- IA Supernovae (white dwarfs implosion): standard candles.
- CMB spectrum

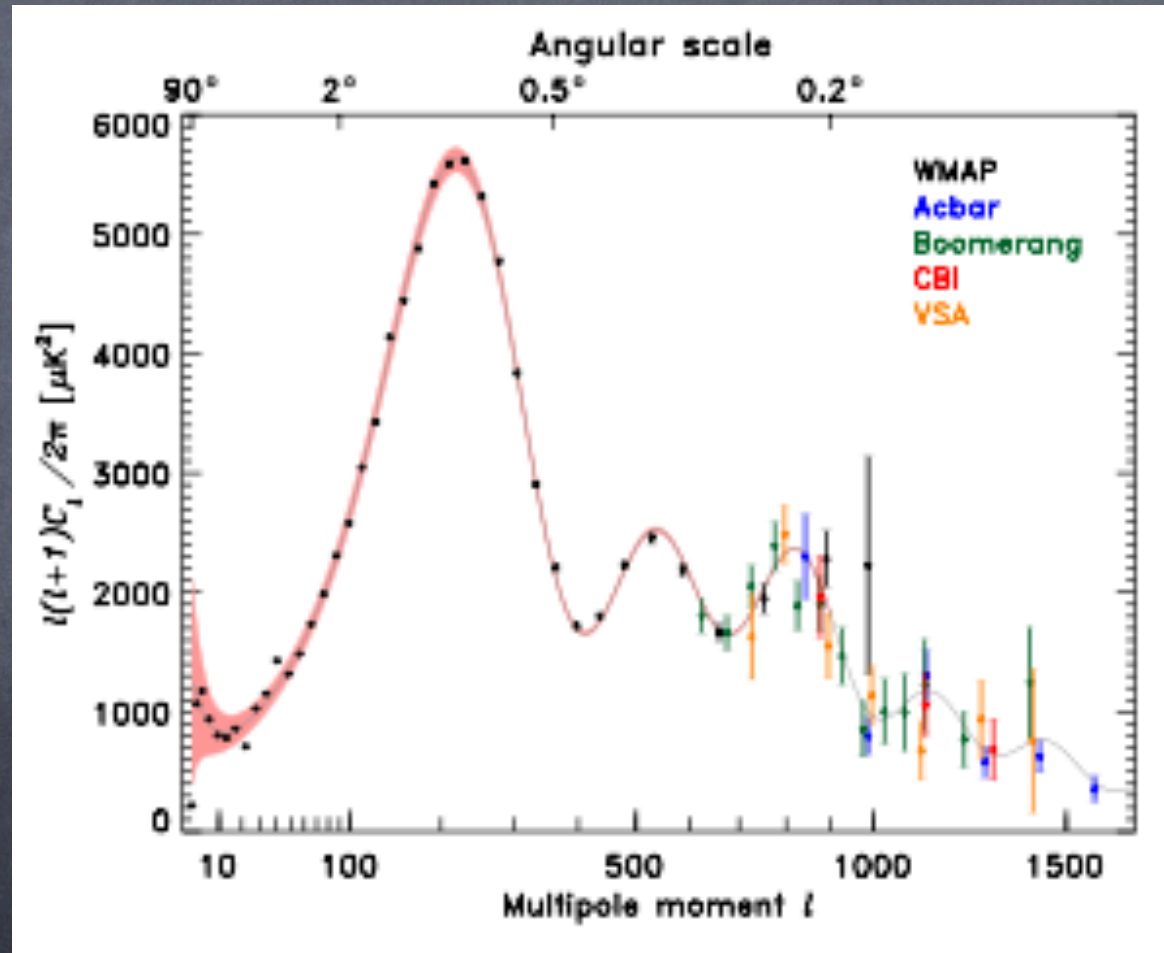


# CMB satellite WMAP



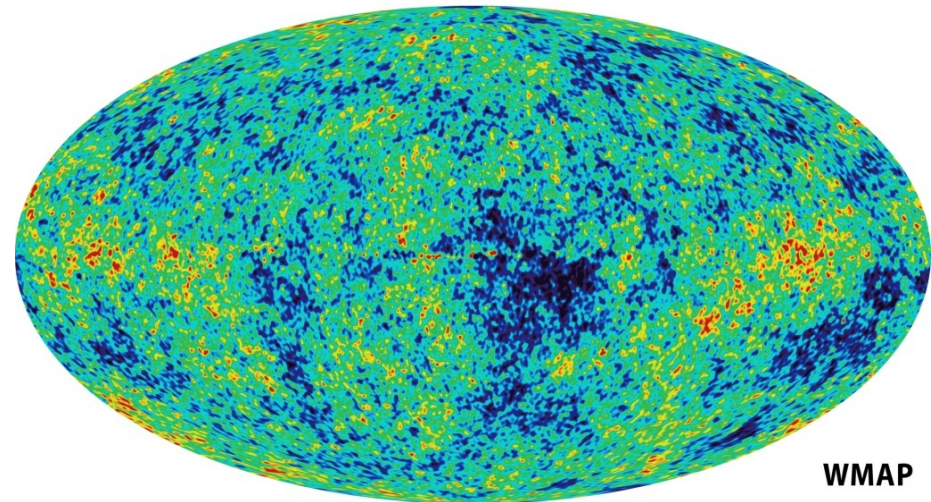
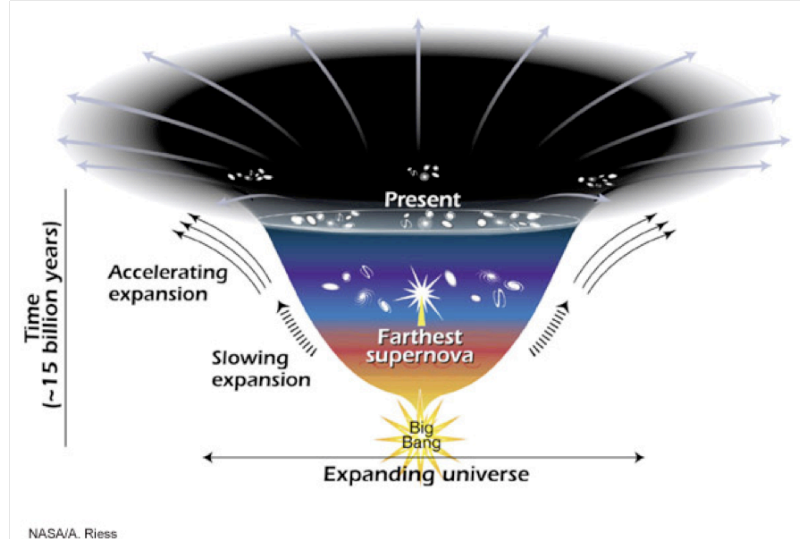
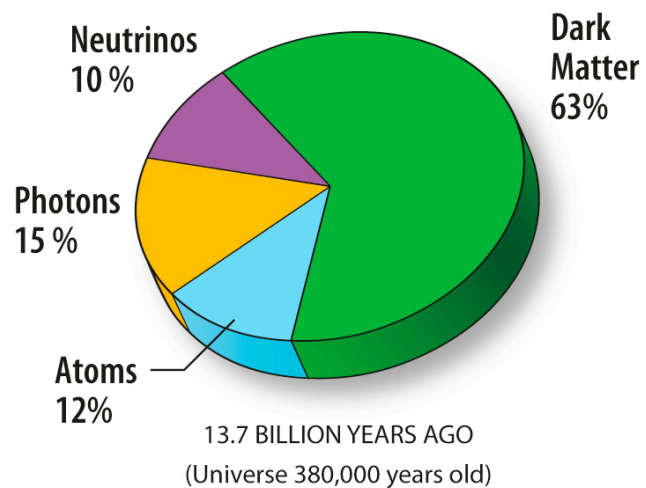
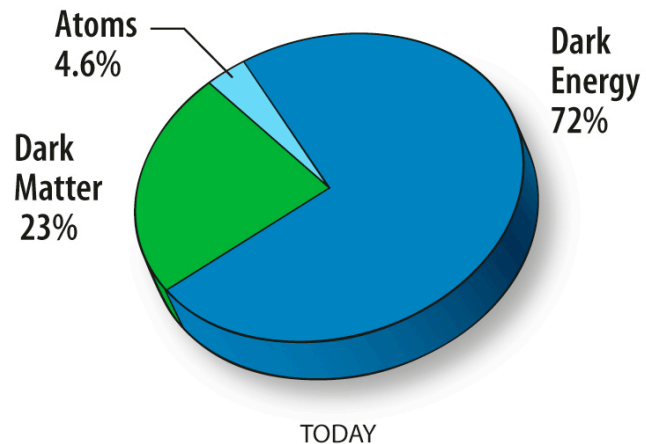


# CMB





# The standard "model" for cosmology



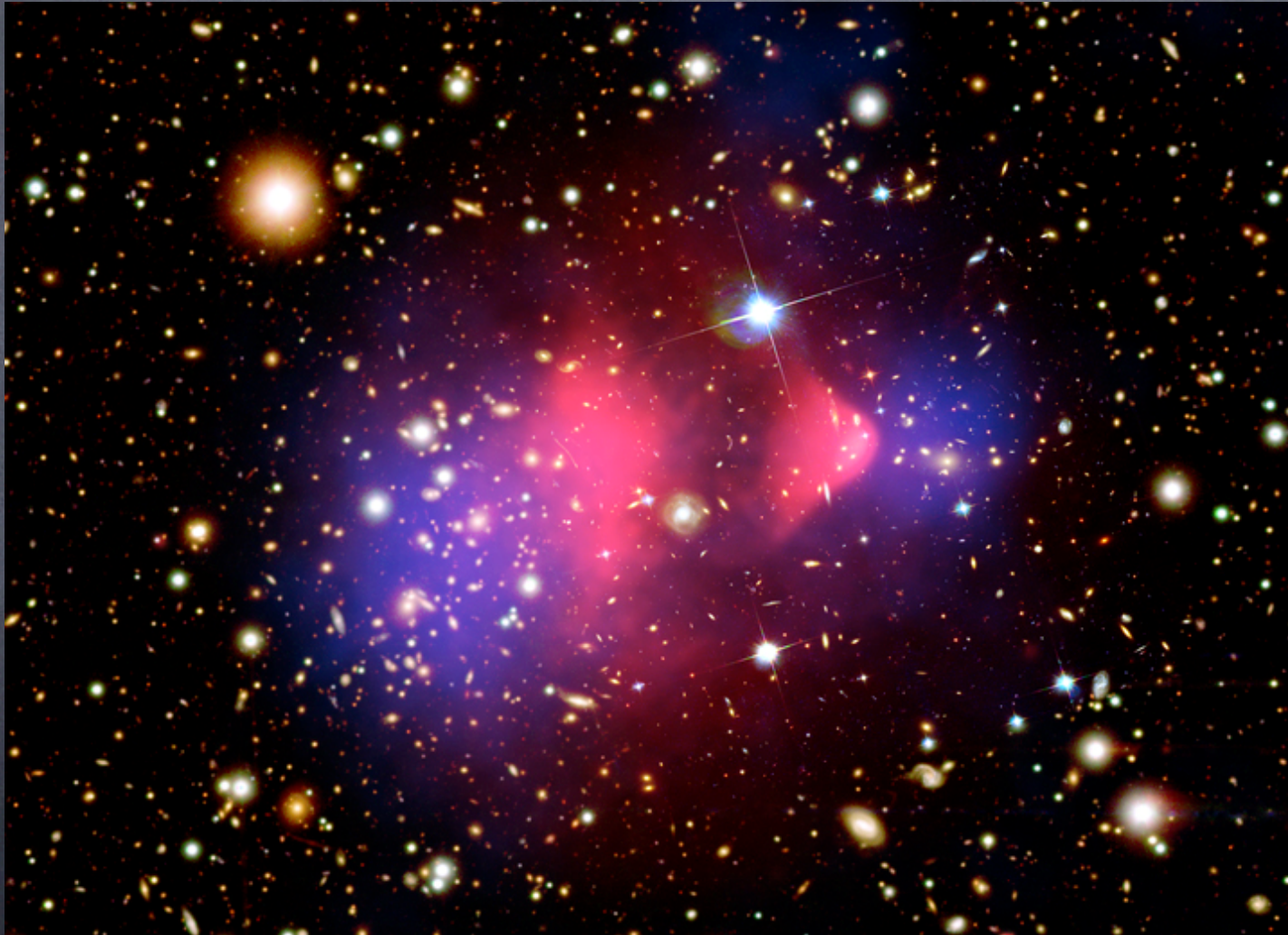


 All observations consistent with DM and DE, where:

- ✓ Total Energy density given by the critical value.
- ✓ Clustering Matter (baryons plus DM) represent 1/3 of total Energy.
- ✓ Strange ( **very** strange) object responsible for 2/3 of Universe energy content!



# Bullet Cluster





# Essence of Dark Energy

- ★ What is it?
- ★ Simple?
- ★ A new Particle/Field?
- ★ Cosmological Constant?



# How to get acceleration?

Friedmann Equation

$$\frac{\ddot{a}}{a} = -\frac{4\pi G}{3} (\rho + 3p)$$

Thus, to have acceleration:

$$p < -\frac{\rho}{3} \Leftrightarrow \omega < -\frac{1}{3}$$



# How to get acceleration?

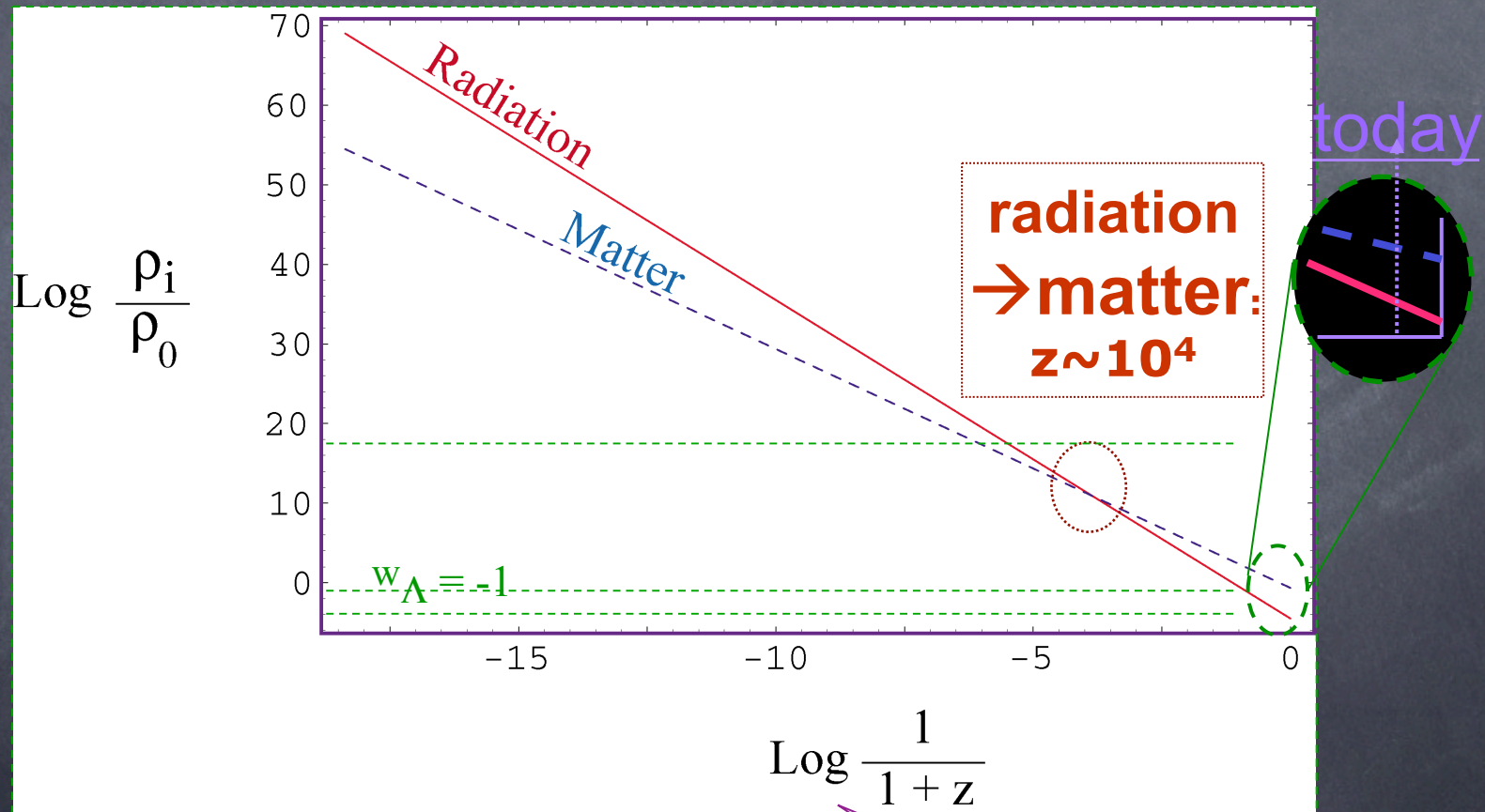
The first hypothesis was the cosmological constant. Is a simple solution, providing  $\omega = -1$ . However, there are some problems :

— The energy density consistent with the observations is such that the cosmological constant must be of the order of  $1(\text{MeV})^4$ , 120 orders of magnitude less than the value obtained from field theory.

— The fact that the cosmological constant is important at this stage of evolution of the Universe is a mystery.



# The Coincidence Problem

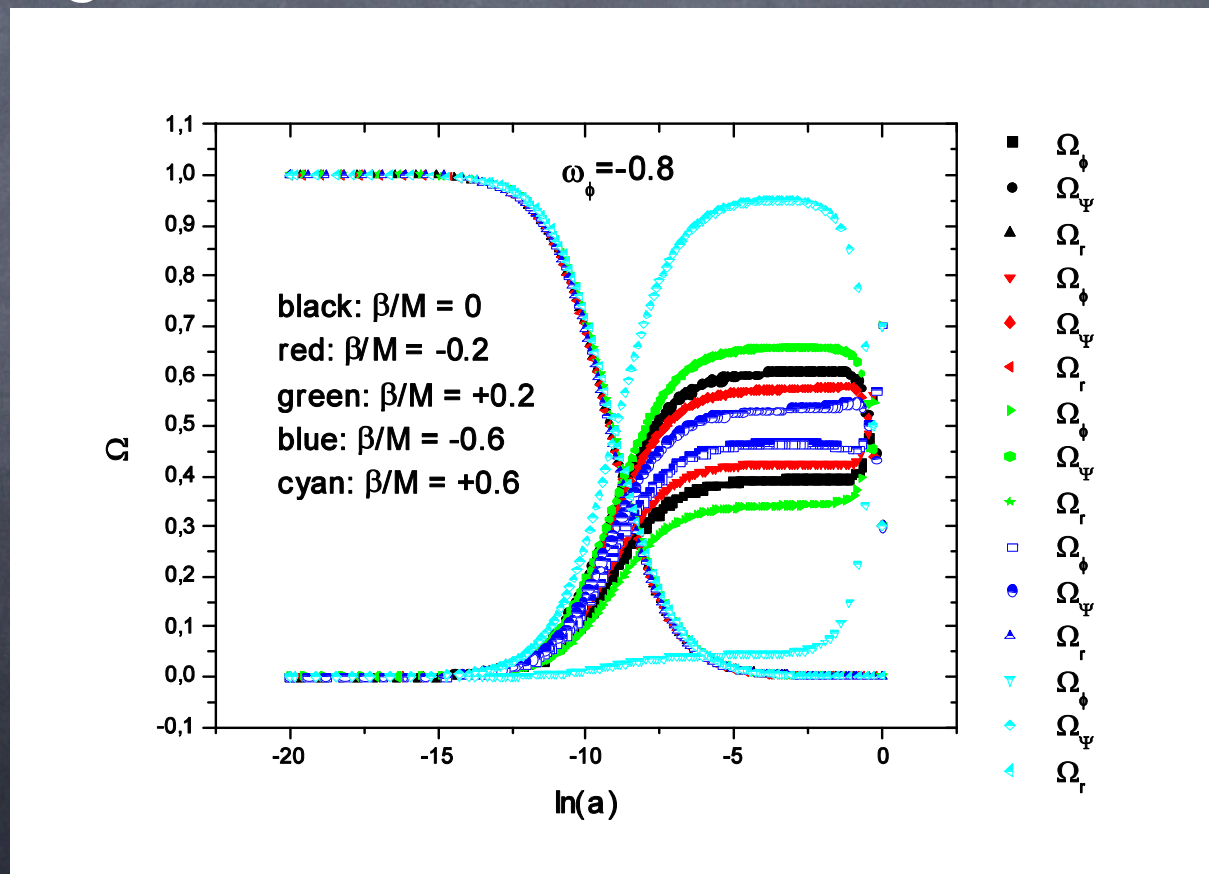


$$1+z = a_0/a$$



# The Coincidence Problem and the interaction DE/DM

- Why DE and DM have the same order of magnitude today (Micheletti, E.A., Wang)?





# The Coincidence Problem and the interaction DE/DM

- \* The decay of DE into DM alleviates the coincidence problem.

- \* Question: These ideas can be tested with observations?



# Interaction DE/DM: Preliminary

- Question of the suppression of lower multipoles of CMB power spectrum.
- Cosmic Holography (E.A., Wang, et al).
- Holographic idea (DE proportional to the square of the scale factor) is good for various non trivial models (**without** cosmological constant).



# Interacting Holographic Model

- The simplest interacting model for DE/DM is given by two fluids interacting

$$\dot{\rho}_{DM} + 3H\rho_{DM} = Q$$

$$\dot{\rho}_{DE} + 3H(1 + \omega)\rho_{DE} = -Q$$



# Interacting Holographic Model

- The simplest interaction is proportional to the density of dark energy

$$Q = 3b^2 H \rho$$



# Interacting Holographic Model

- Comparing with the supernovae data lead us to the following results (Wang, E.A. et al)

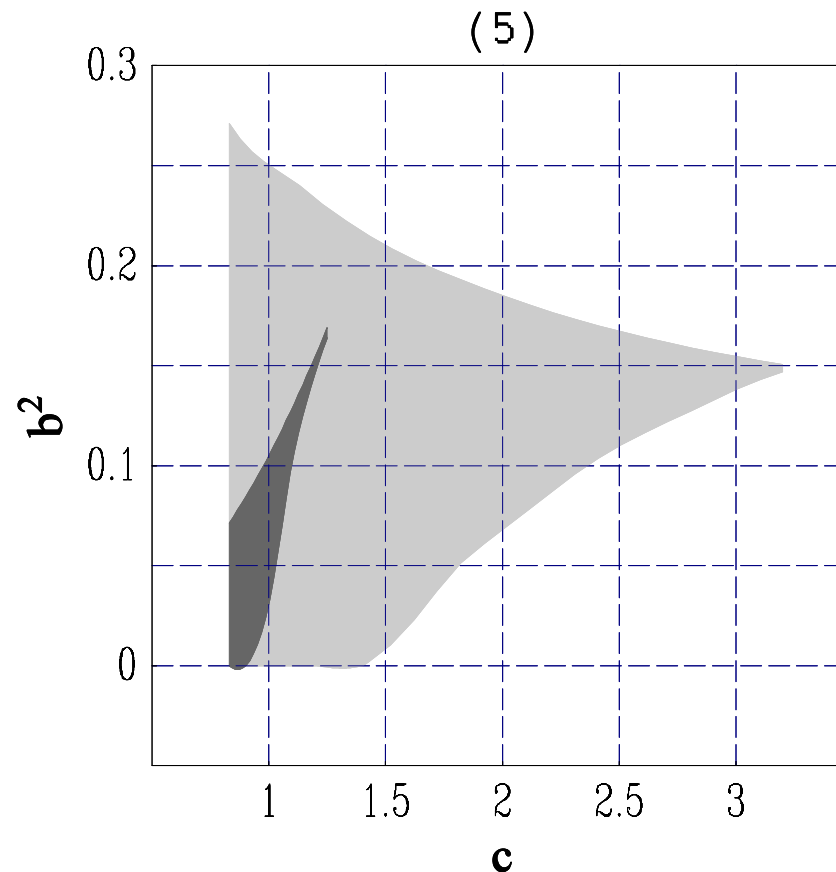
$$b^2 = 0.00^{+0.12}_{-0.00}$$

$$c = 0.40^{+0.75}_{-0.05}$$



# Interacting Holographic Model

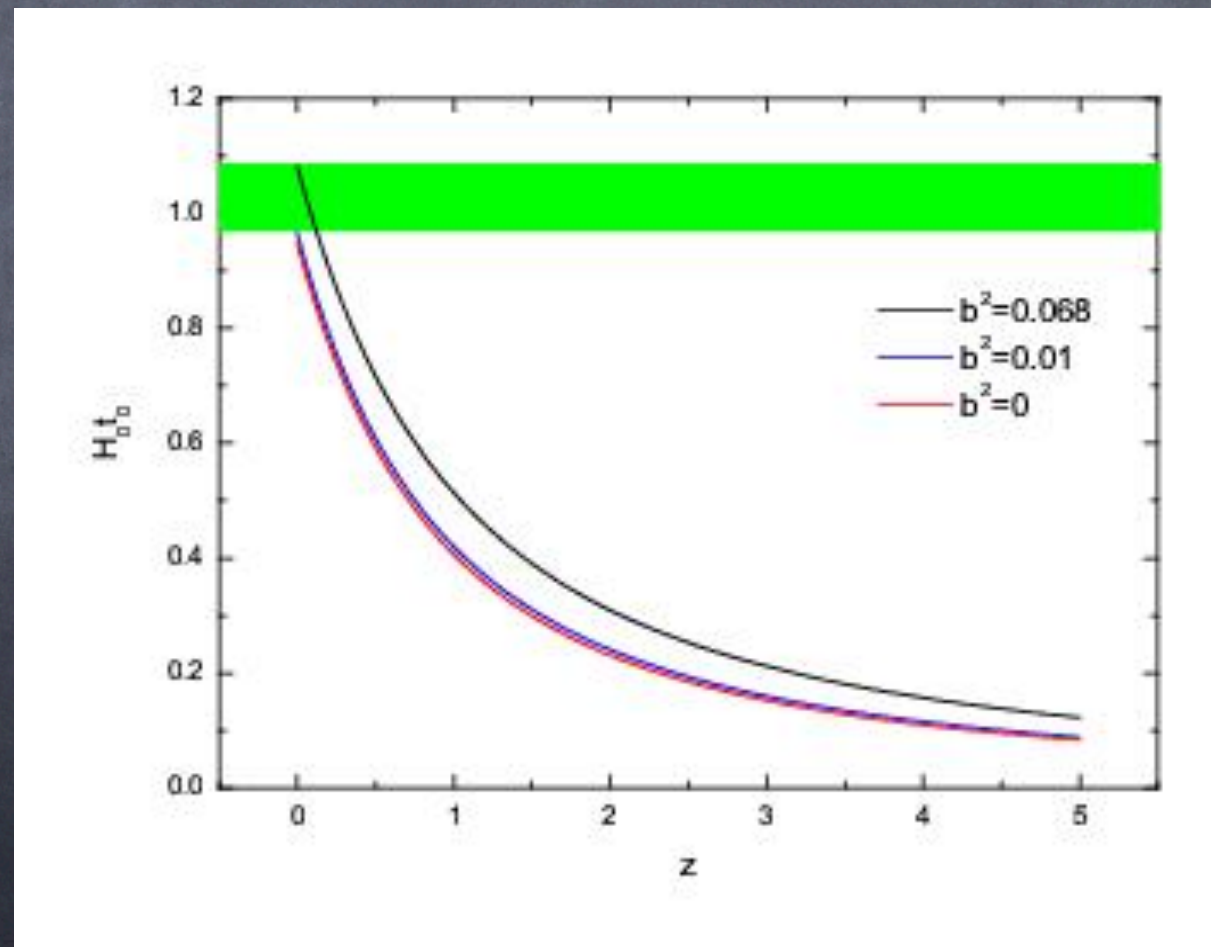
- Regarding to the equation of state, the comparison with observational data tell us





# Further on Phenomenology

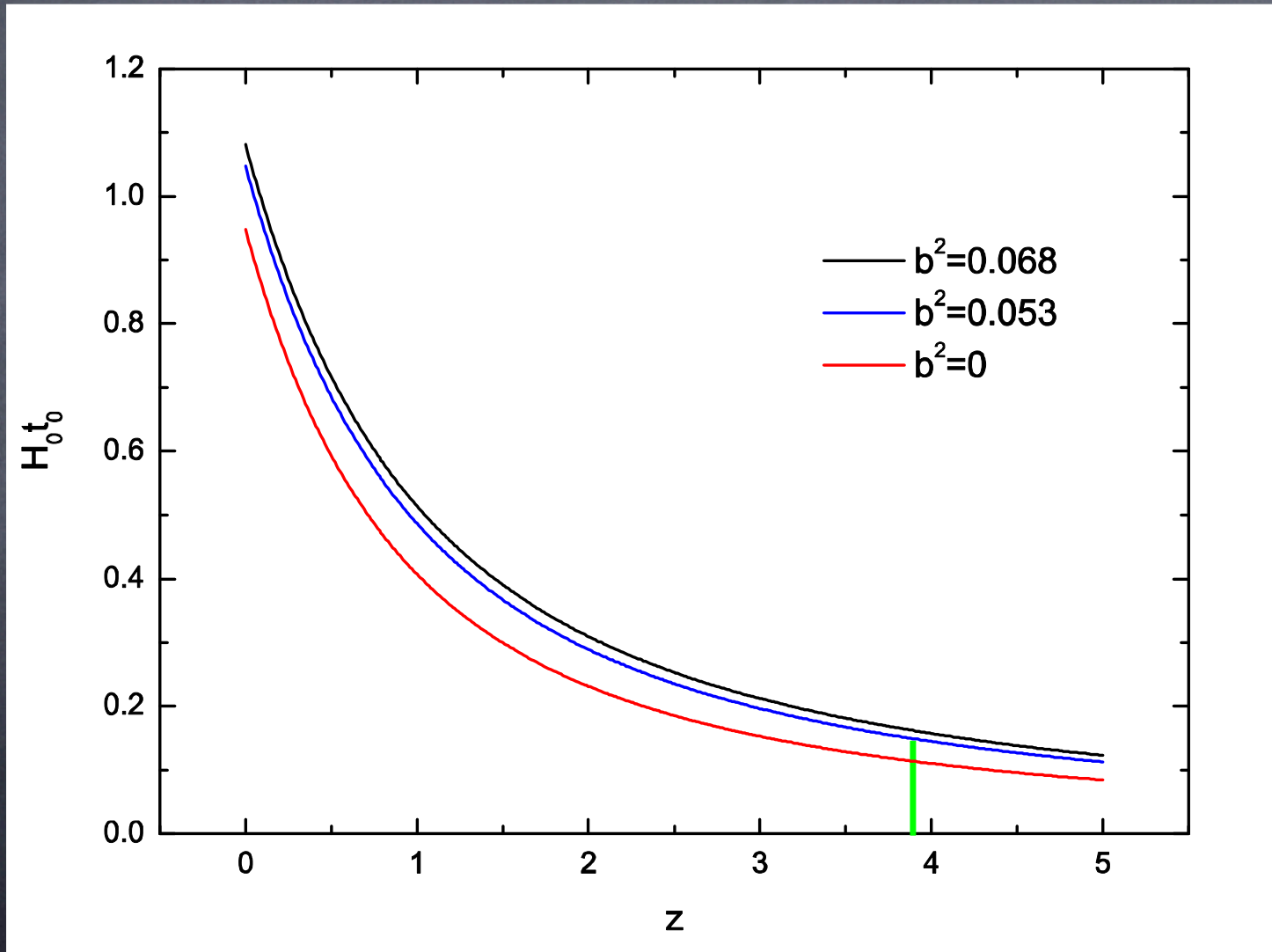
- Age of Universe





# Further on Phenomenology

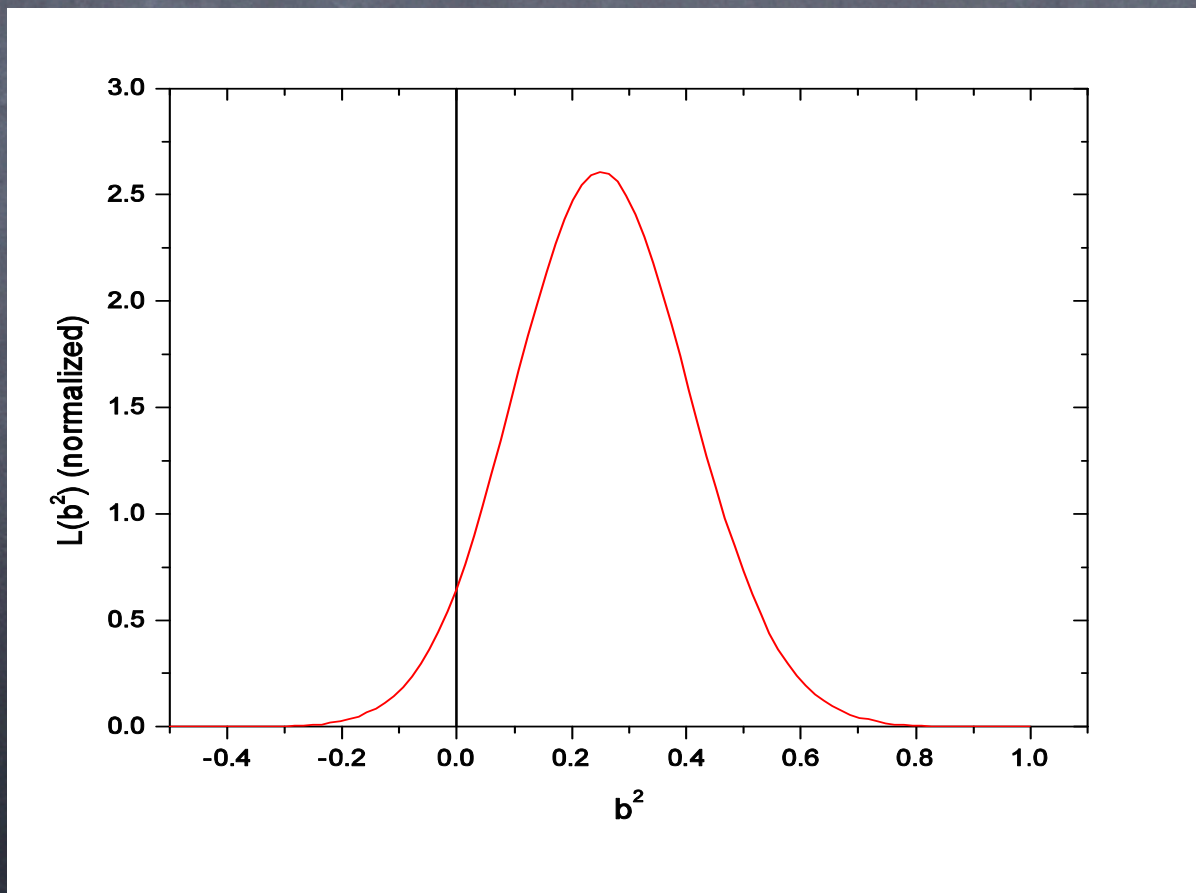
## ● Age of Quasar APM 08279+5255





# Further on Phenomenology

- Lower multipoles in CMB spectrum





# Further on Phenomenology

- Order of magnitude estimate for  $b^2$  (with some dependence on the DE equation of state) :

$$\star 1\sigma : \quad 0.10 < b^2 < 0.2$$

$$\star 2\sigma : \quad -0.10 < b^2 < 0.24$$

Probability of  $b^2$  positive: 95%



# Thermodynamics?

- The interaction must be DE  $\rightarrow$  DM in order to preserve the second law of thermodynamics (Wang, Pavon)



# Galaxy clusters and the Virial theorem

- Only baryonic matter is observed.
- Suppose that Dark Matter follows the pattern of baryonic matter.
- The interaction with Dark Energy works as an effective potential.
- Then, the Virial theorem has a correction (o.

Bertolami et al).



# Galaxy clusters and the Virial theorem

- Virial theorem :

$$2K + U = 0$$

- General Relativity: Layser-Irvine equation



# Galaxy clusters and the Virial theorem

- Layser-Irvine equation with interaction (Virial Theorem)

$$M_{vir} = \frac{(1 - \frac{\zeta}{2})}{(1 - 2\zeta)} M_X$$

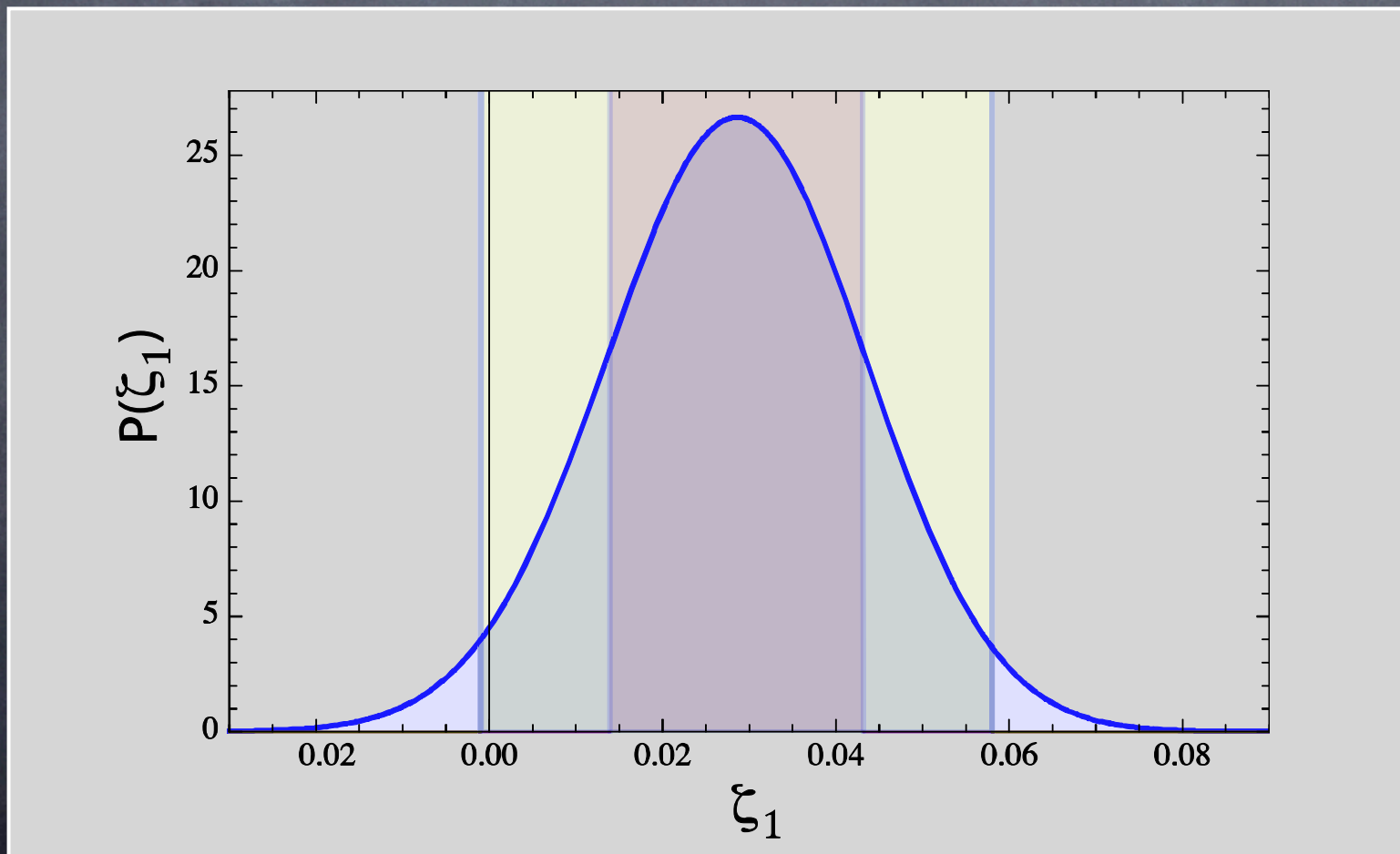
$$M_{vir} = \frac{(1 - \frac{\zeta}{2})}{(1 - 2\zeta)} M_{WL}$$

$$M_X = M_{WL}$$



# Galaxy clusters and the Virial theorem

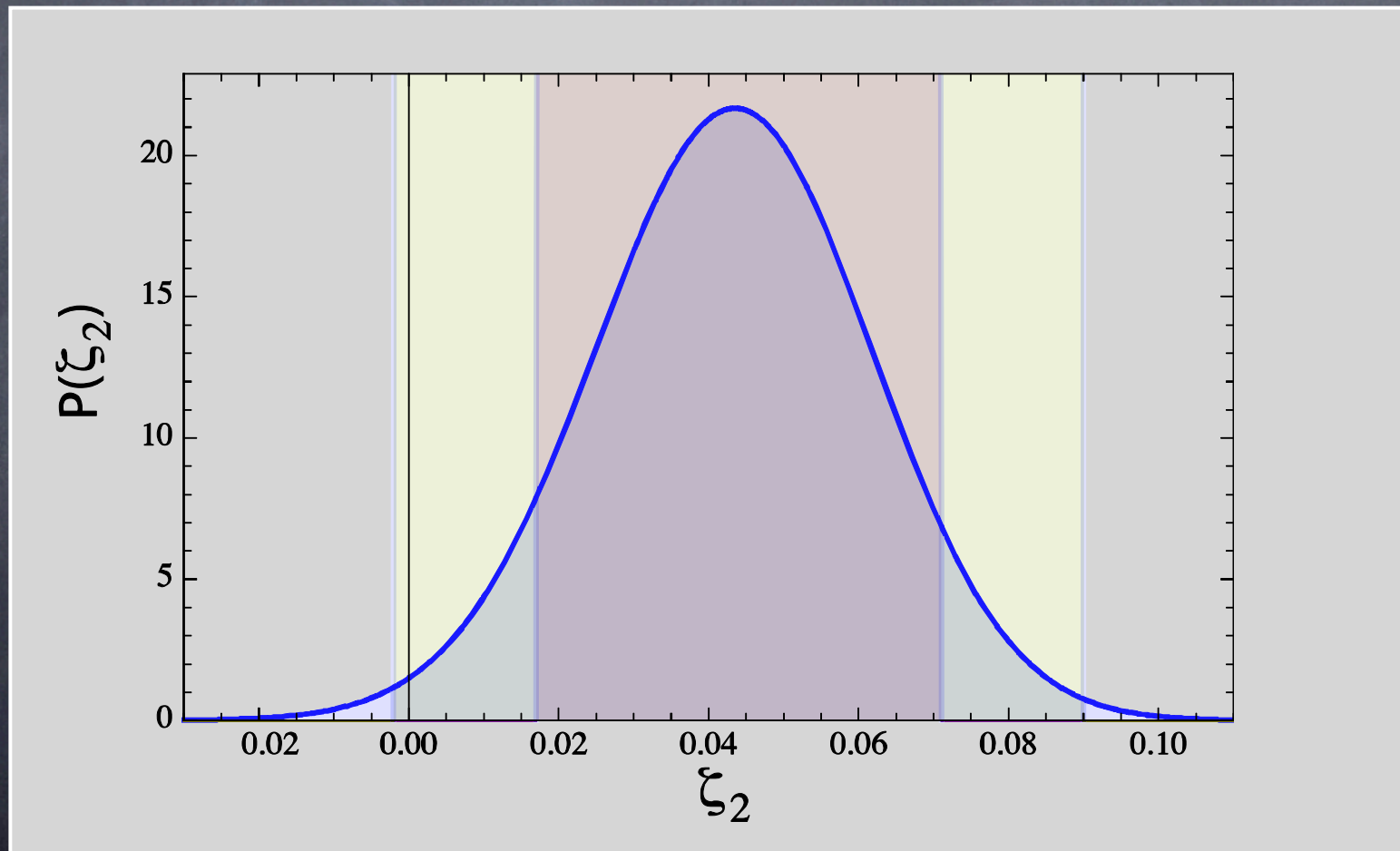
👁 Results (E.A., Abramo, Sodre, Wang, Campos)





# Galaxy clusters and the Virial theorem

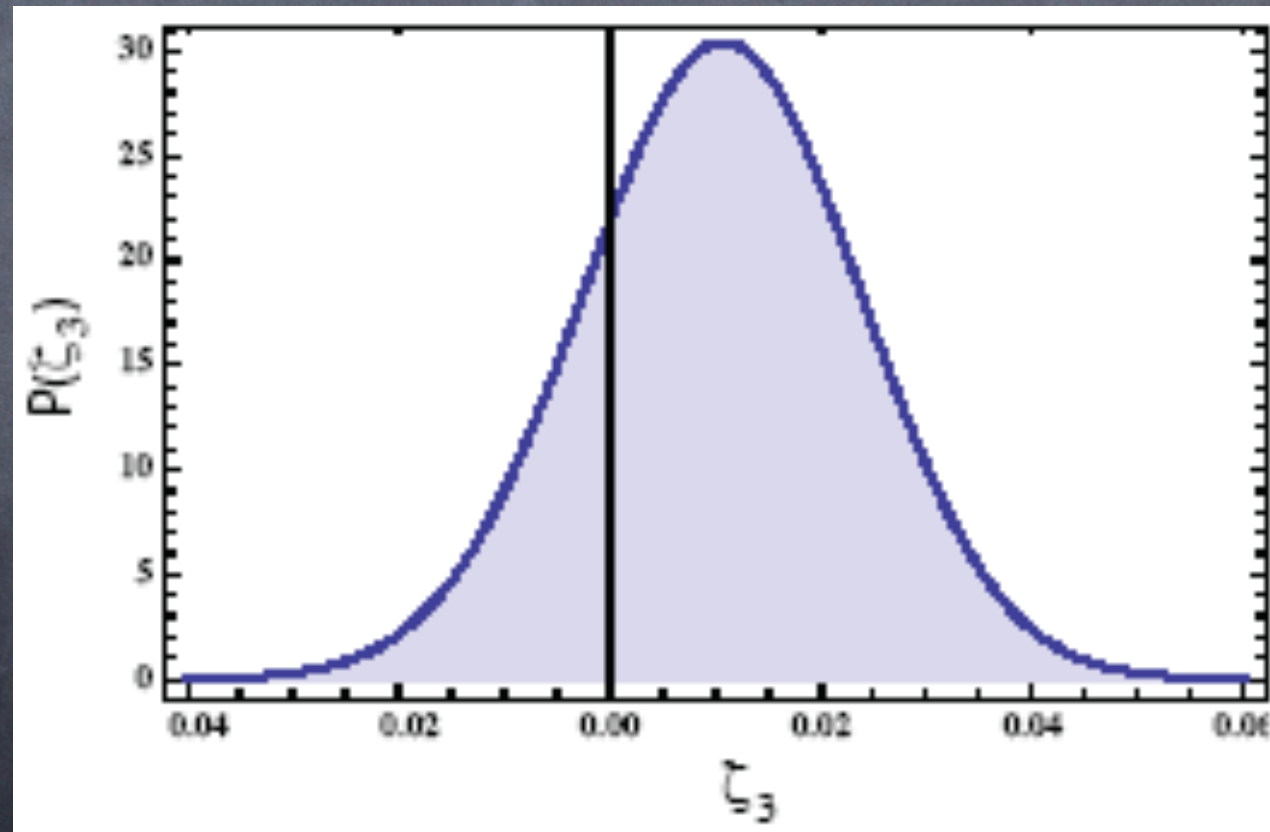
## Results





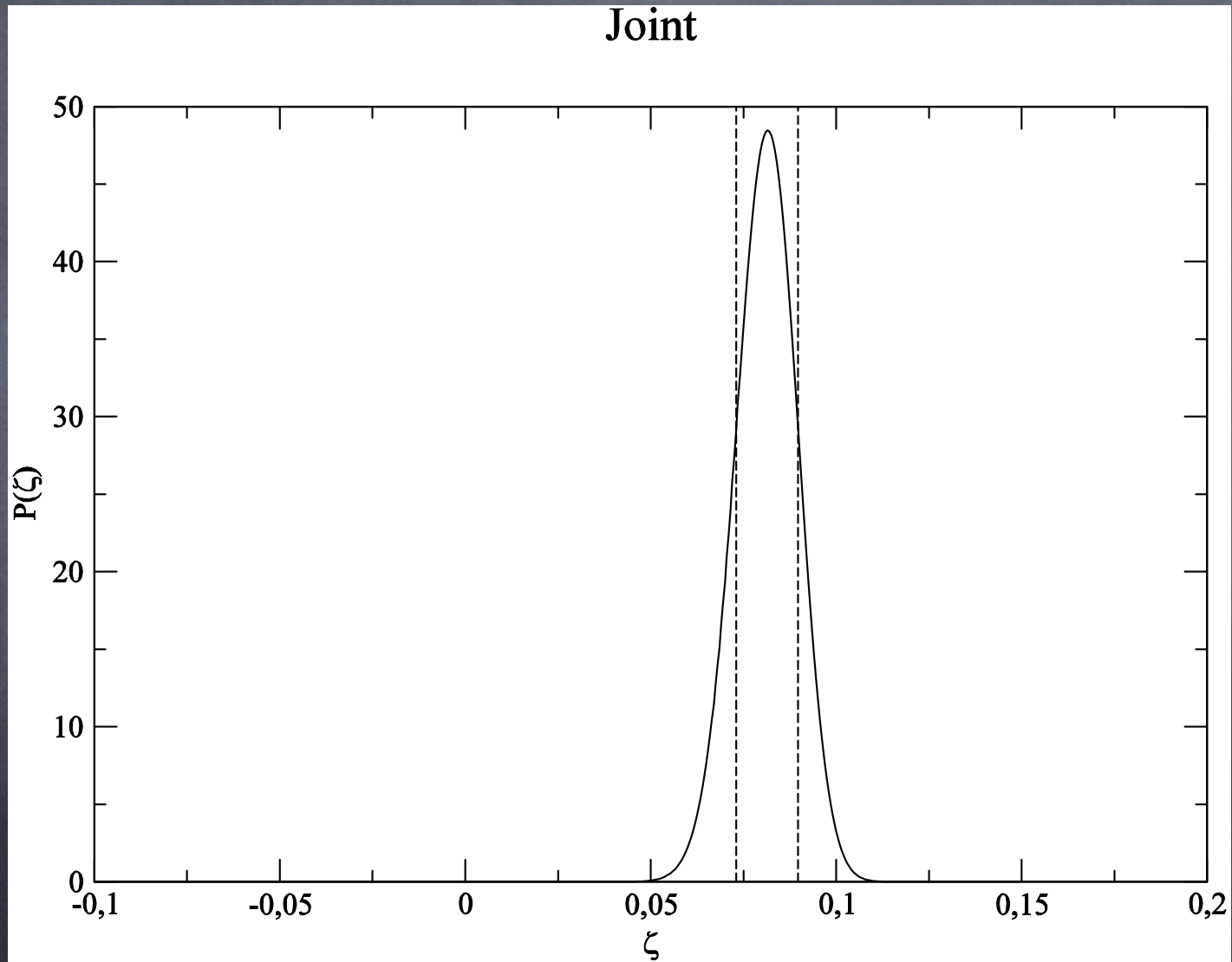
# Signature in Galaxy Clusters: an apparent correction of the Virial Theorem

## Results





# Clusters ( $\sim 100$ )





# Perturbations

- The interaction  $\zeta\rho_{dm}$  gives rise to very strong non perturbative effects (He, Wang, E.A.).
- They come from classical solutions of equations of motion of order  $\frac{1}{\zeta}$ .
- So, the perturbative methods fail.
- The interaction  $\zeta\rho_{de}$  behaves differently.
- Alternative methods are needed.



# Modeling the interaction

- Suppose DE/DM are interacting parts (E.A., Wang)•
- Suppose that there is a “Thompson scattering amplitude”.
- Existence of galaxies with  $M_{dm}$  of order  $10M_{baryons}$  leads to a coupling constant  $b^2 \sim 3\alpha_{QED}$ .
- According to previous calculations.



# Modeling the interaction

- Suppose DM and baryons are at equilibrium after the Big Bang.
- In the balance between the expansion rate and the self annihilation into ordinary matter, we have a relation between the DM interaction constant and the fine structure constant



# Theoretical model to DE/DM

(Micheletti, Wang, E.A.)

For  $g_{\mu\nu} = \text{diag}(1, -a(t)^2, -a(t)^2, -a(t)^2)$  one finds

$$\ddot{\varphi} = -(1 - \alpha\dot{\varphi}^2) \left[ \frac{1}{\alpha} \frac{d \ln V(\varphi)}{d\varphi} + 3H\dot{\varphi} - \frac{\beta \bar{\Psi} \Psi}{\alpha V(\varphi)} \sqrt{1 - \alpha\dot{\varphi}^2} \right],$$

with  $H = \frac{\dot{a}}{a}$ . We also find

$$\frac{d(a^3 \Psi^\dagger \Psi)}{dt} = 0,$$
$$\frac{d(a^3 \bar{\Psi} \Psi)}{dt} = 0.$$

From the latter,  $\bar{\Psi} \Psi = \frac{\bar{\Psi}_0 \Psi_0 a_0^3}{a^3}$ .



# Theoretical model to DE/DM

Moreover,

$$\rho_\varphi = \frac{V(\varphi)}{\sqrt{1 - \alpha\dot{\varphi}^2}},$$

$$P_\varphi = -V(\varphi)\sqrt{1 - \alpha\dot{\varphi}^2},$$

$$\rho_\Psi = M^*\bar{\Psi}\Psi$$

$$P_\Psi = 0,$$

where  $M^* \equiv M - \beta\varphi$ . Note that  $\omega_\varphi \equiv P_\varphi/\rho_\varphi = -(1 - \alpha\dot{\varphi}^2)$ . Using time derivatives and equations of motion, we also find

$$\begin{aligned}\dot{\rho}_\varphi + 3H\rho_\varphi(\omega_\varphi + 1) &= \beta\dot{\varphi}\frac{\bar{\Psi}_0\Psi_0a_0^3}{a^3} \\ \dot{\rho}_\Psi + 3H\rho_\Psi &= -\beta\dot{\varphi}\frac{\bar{\Psi}_0\Psi_0a_0^3}{a^3}.\end{aligned}$$



# Theoretical model to DE/DM

Similar to eq. for the interaction of Dark Matter and Dark Energy. The right hand side does not contain the Hubble parameter  $H$  explicitly, but it does contain the time derivative of the scalar fields, which should behave as the inverse cosmological time, replacing thus the Hubble parameter.

Friedmann equation for a flat universe reads

$$H^2 = \frac{1}{3M_{pl}^2} \left[ M^* \frac{\bar{\Psi}_0 \Psi_0 a_0^3}{a^3} + \frac{V(\varphi)}{\sqrt{1 - \alpha \dot{\varphi}^2}} \right],$$

where  $M_{pl}^2 = (8\pi G)^{-1}$  e  $H = \frac{\dot{a}}{a}$ .

Some analytic solutions in the pure bosonic case have been found elsewhere for the potential

$$V(\varphi) = \frac{m^{4+n}}{\varphi^n}, n > 0.$$



# Theoretical model to DE/DM

$$\frac{d\omega}{dz} = \frac{4\omega\sqrt{\omega+1}}{H_0 E(z)\phi(1+z)} - \frac{6\omega(\omega+1)}{1+z}$$

$$+ \frac{2\left(\frac{\beta}{M}\right)\frac{\Omega_{\Psi 0}}{\sqrt{\alpha}}}{\left(1 - \Omega_{\Psi 0}\right)\left(1 - \frac{\beta}{M}\frac{\phi_0}{\sqrt{\alpha}}\right)} \frac{\omega\sqrt{\left|\frac{\omega}{\omega_0}\right|(\omega+1)}\left(\frac{\phi}{\phi_0}\right)^2}{H_0 E(z)}(1+z)^2,$$

$$\frac{d\phi}{dz} = -\frac{\sqrt{\omega+1}}{H_0 E(z)(1+z)},$$

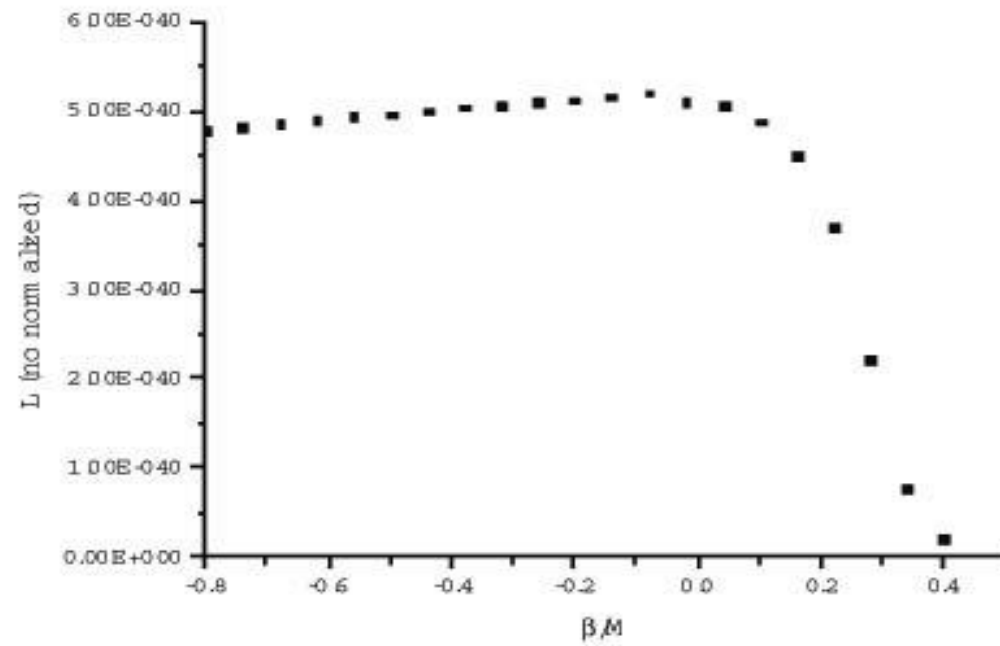
$$\frac{dt}{dz} = -\frac{1}{(1+z)H_0 E(z)},$$

where  $H_0 = 2.133h \times 10^{-39} \text{ MeV}$  is the value of the Hubble parameter,  $\Omega_{\Psi 0} = \frac{\rho_{\Psi}}{3M_{pl}^2 H_0^2}$  and

$$E(z) = \sqrt{\frac{1 - \frac{\beta}{M}\frac{\phi}{\sqrt{\alpha}}}{1 - \frac{\beta}{M}\frac{\phi_0}{\sqrt{\alpha}}} \Omega_{\Psi 0} (1+z)^3 + \left(\frac{\phi_0}{\phi}\right)^2 \sqrt{\left|\frac{\omega_0}{\omega}\right|} (1 - \Omega_{\Psi 0})}.$$

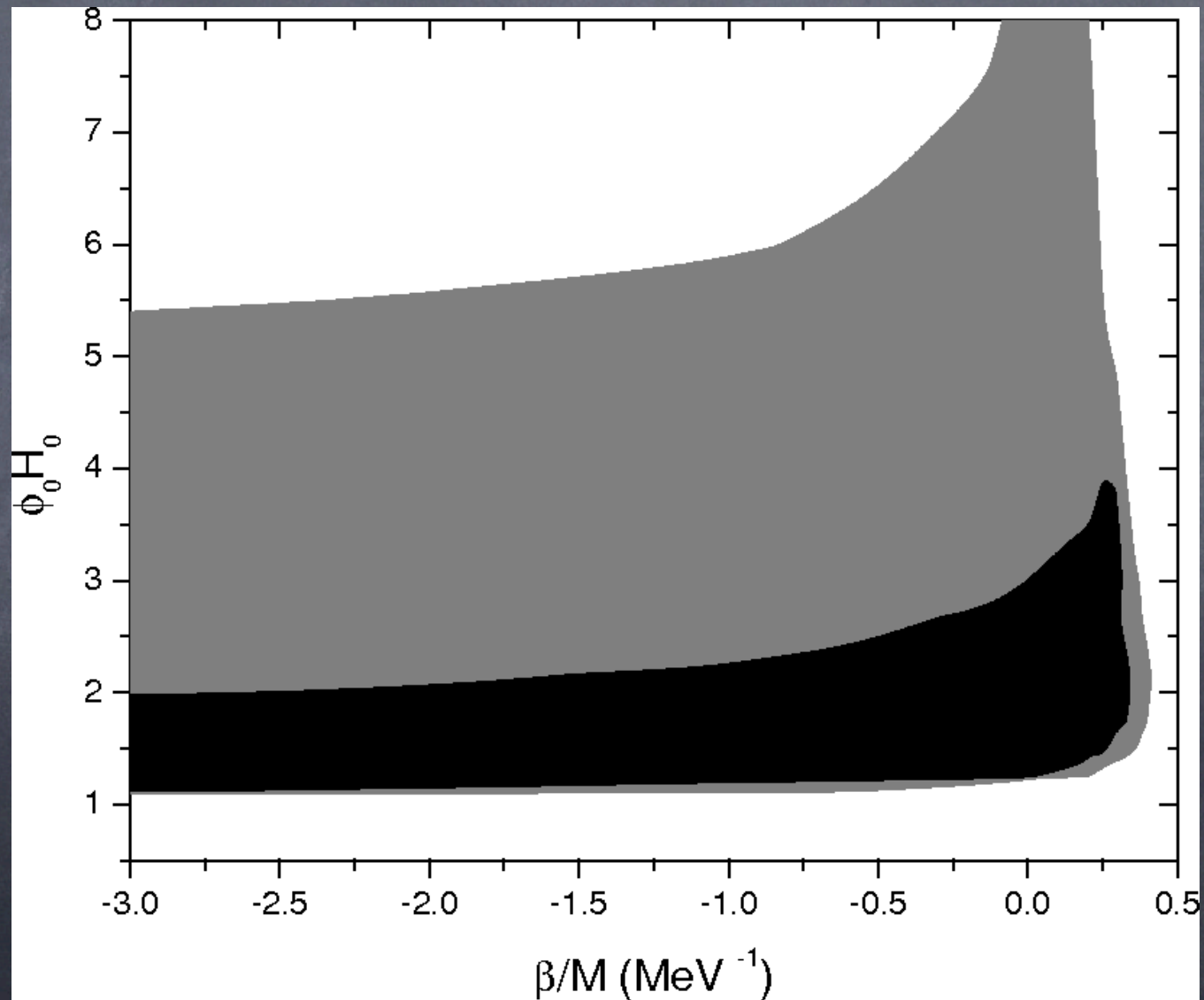


# Likelihood $\beta$



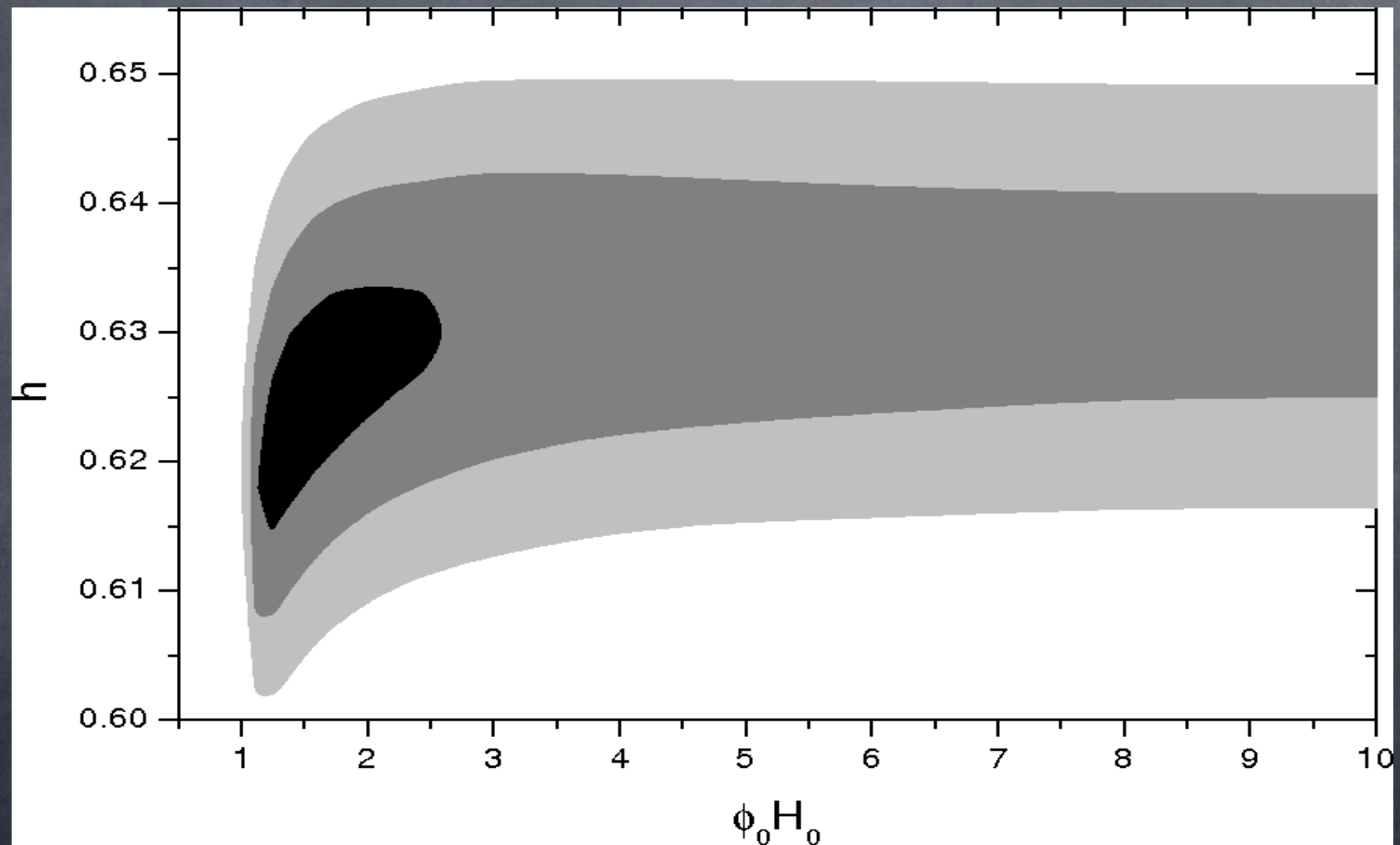


# 2D-Likelihood $\Phi_0\beta$



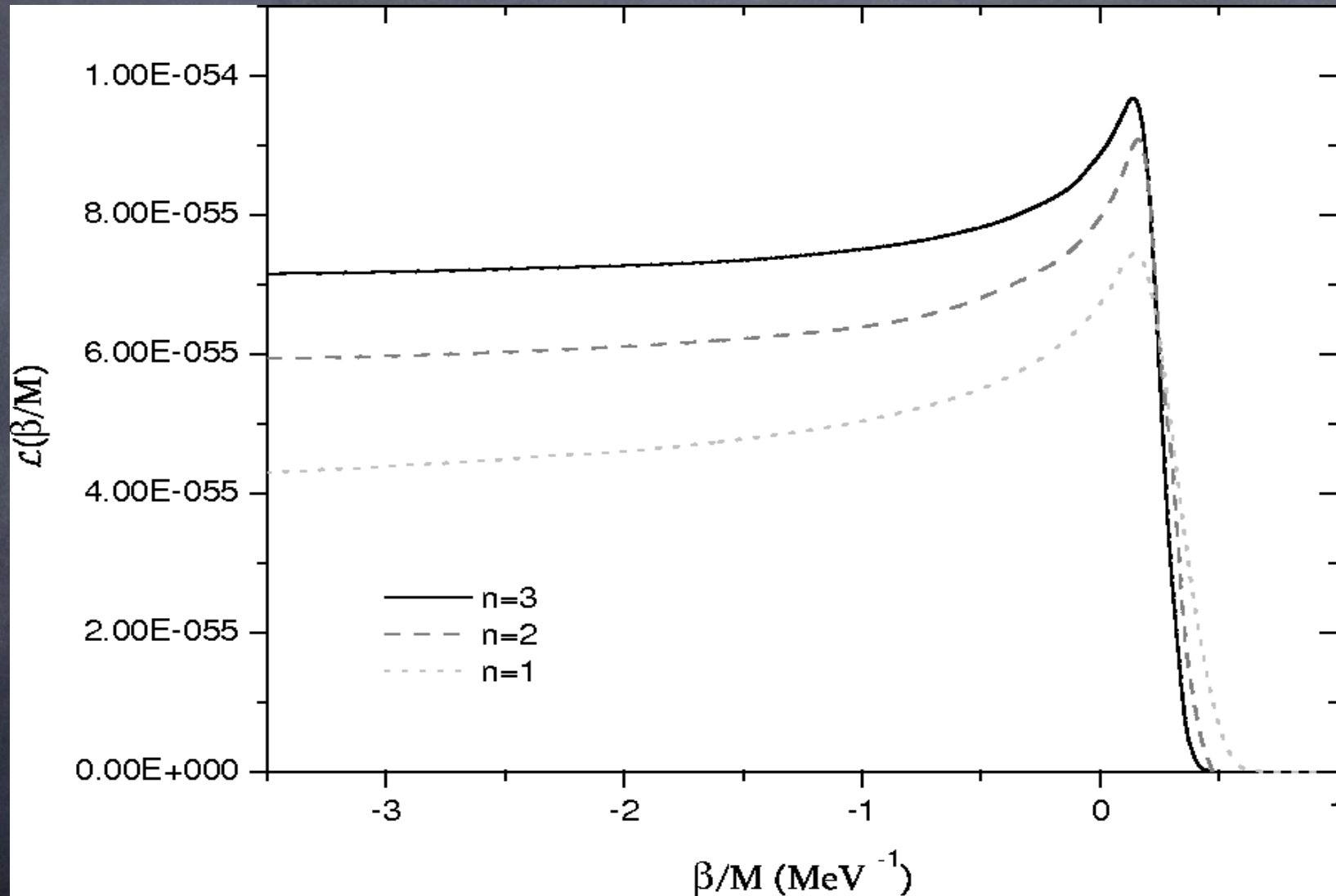


# 2D-Likelihood $\Phi_0 H_0$



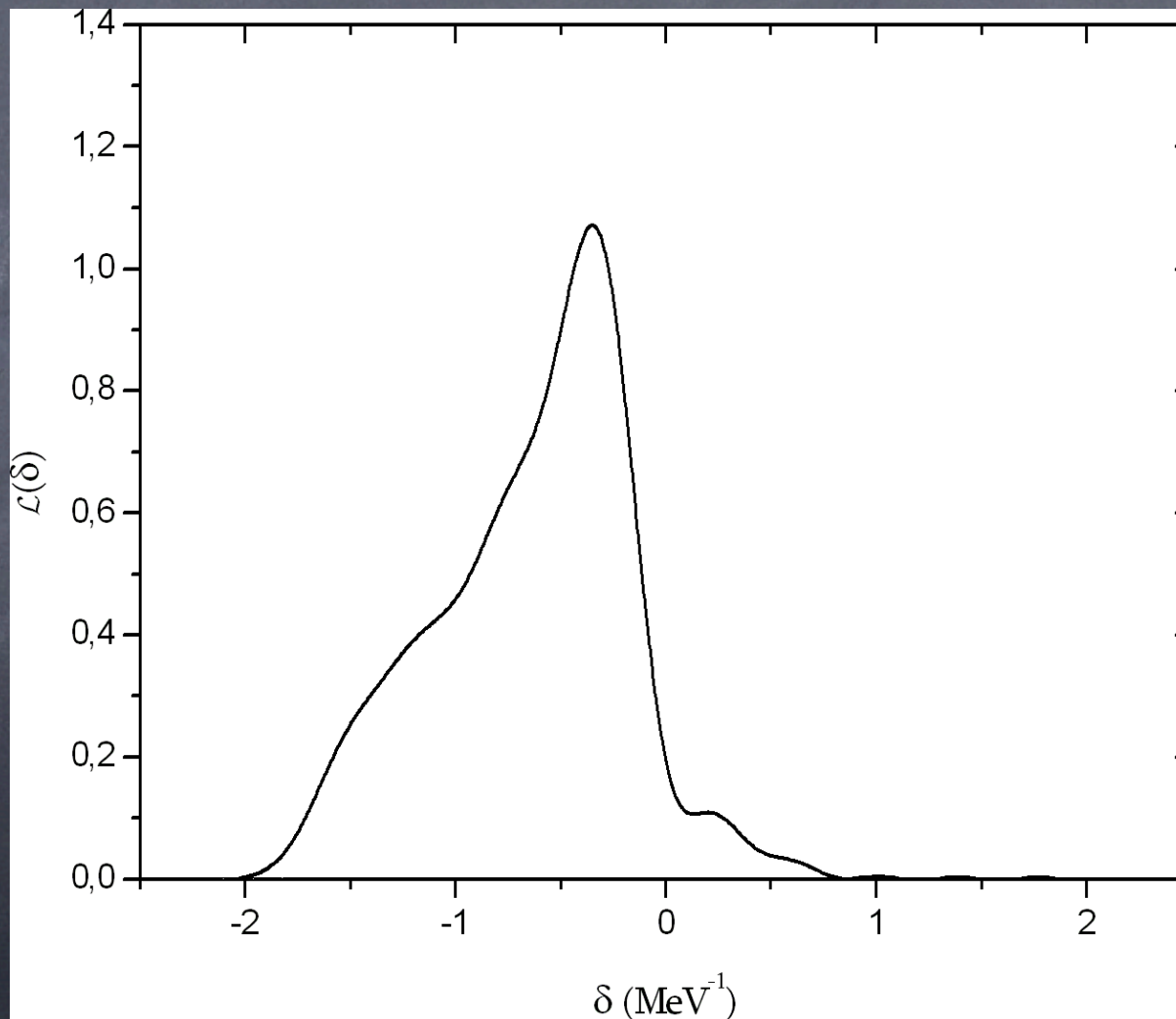


# Likelihood $n = 1, 2, 3$





# Likelihood $\delta$ (effective interaction)





# Theoretical model to DE/DM

- Are there standard model particles that constitute a natural candidate for DM?
- **Neutralinos**: candidate for lighter supersymmetric particle (LSP), would be stable, heavy enough, it interacts very weakly with ordinary matter.
- **Gravitinos** : candidate for LSP, stable, it interacts almost exclusively by gravitation.



# Theoretical model to DE/DM

- Are there standard model particles that constitute a natural candidate for DE?
- Such particle must be very stable ( $\tau$  of order of 10 billions years) interacting very weakly.
- The standard model can have multiple scalar particles, the stability condition is quite strong. There are uncertainties.



We compute now the likelihood of the function  $\delta \equiv \frac{\frac{\beta}{M\sqrt{\alpha}}}{1 - \frac{\beta}{M\sqrt{\alpha}}\phi_0}$

The likelihood of  $\delta$  is determined from the likelihood of  $\frac{\beta}{M\sqrt{\alpha}}$  and  $\phi_0$  according to

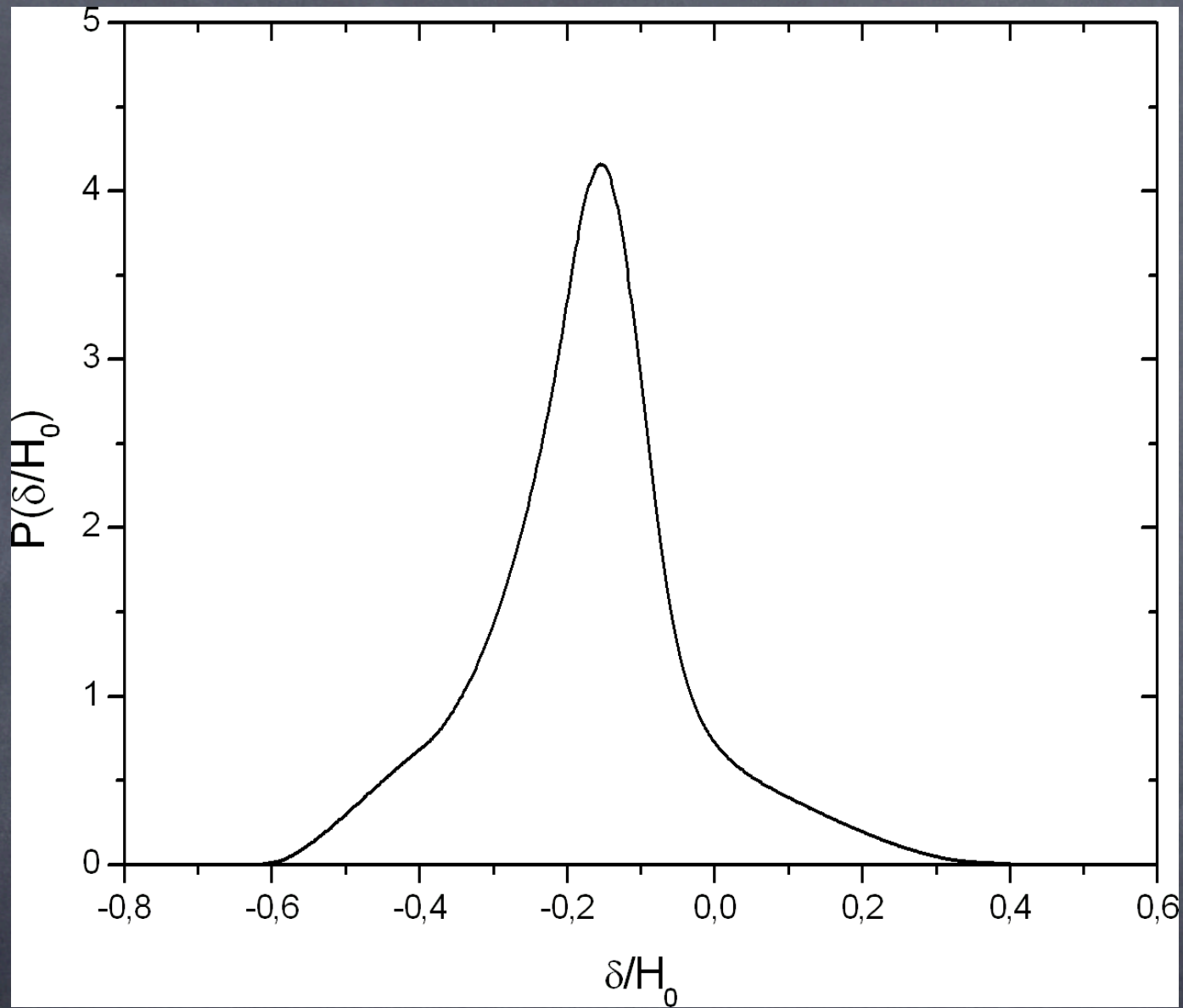
$$P(u') = \int \int \delta(u' - \delta) P\left(\frac{\beta}{M\sqrt{\alpha}}, \phi_0\right) d\left(\frac{\beta}{M\sqrt{\alpha}}\right) d\phi_0, ,$$

where

$$P\left(\frac{\beta}{M\sqrt{\alpha}}, \phi_0\right) \sim \text{prior}\left(\frac{\beta}{M\sqrt{\alpha}}, \phi_0\right) \mathcal{L}\left(\frac{\beta}{M\sqrt{\alpha}}, \phi_0\right)$$

for some prior we choose for  $\frac{\beta}{M\sqrt{\alpha}}$  and  $\phi_0$ .





The likelihood function



## Calculation of the Decay Rate of the Dark energy

We know that the decay rate (per unit time and volume) of a scalar particle in a meta stable vacuum is given by the formula (Coleman 77):

$$\frac{\Gamma}{V} = \frac{S(\phi_b)^2}{(2\pi\hbar)^2} \cdot e^{-\left(\frac{S_e}{\hbar} - \frac{S_\Lambda}{\hbar}\right)} \cdot \left( \frac{\det'(-\partial_\mu\partial_\mu + V''(\bar{\phi}(\rho)))}{\det(-\partial_\mu\partial_\mu + V''(2m/3\lambda))} \right)^{-\frac{1}{2}}$$



We will calculate the decay rate of dark energy scalar field described by the Wess-Zumino potential (E.A. Graeff, Wang)

$$U = |2m\phi - 3\lambda\phi^2|^2$$

with a term of supersymmetry breaking equal to

$$\xi\phi = \epsilon = 10^{-47} GeV^4$$

This potential corresponds to that of dark energy, which currently is at the metastable minimum corresponding to the cosmological constant, that has the value

$$10^{-47} GeV^4$$

We wish to calculate the decay's time of the metastable minimum to the true stable minimum, which corresponds to a decay of dark energy in dark matter. The dominant term of the formula of the decay's rate is the term  $e^{-S_e/\hbar}$ . So, we compute calculate it first.



The classical solution follows the equation

$$\frac{\delta S_e}{\delta \phi} = \left( -\frac{\partial^2 \phi}{\partial \rho^2} + V'(\phi) \right) = 0$$

$S_e$  = euclidean action

Due to the  $O(4)$  invariance of the field  $\phi$ , this equation can be written as:

$$\frac{\partial^2 \phi}{\partial \rho^2} + \frac{3}{\rho} \cdot \frac{\partial}{\partial \rho} \phi - V'(\phi) = 0$$



This equation of motion can be interpreted similarly to the equation of motion of a particle at position  $\phi$  moving in time  $\rho$ , subject to a force  $= -V'$ , and a frictional force with a coefficient depending on  $\rho$ . The classic solution of minimal action is the oscillatory solution that satisfies this equation, which in analogy corresponds to a particle oscillating between the two maxima in the potential  $-V(\phi)$ . Due to the invariance  $O(4)$  we can interpret the solution as a bubble of true vacuum separated by a wall of the false vacuum.

$$m > 10^{-13} \text{ GeV}$$



# Conclusions

- At level of  $1\sigma$  there are good reasons to believe that DE and DM are interacting fields.
- At level of  $2\sigma$  the results are still reliable.
- There are some solid indications from galaxy clusters
- There would be models of particles/fields describing DE/DM interaction ? (we are working)



# Conclusions

- The QFT can give us some convincing models, however, more constraints are needed. The models are still very poor, lacking a strong phenomenology, and the results so far are still compatible with no interaction.
- The problem comes from the DE sector: the quintessence described is somewhat sensitive to various parameters.



# Conclusions

- QFT models could describe correctly the interaction DE/DM? (work in progress)
- Neutralinos and gravitinos seem to be good candidates to DM.
- Indeed, fermions seems to be good candidates to DM. What about DE?
- What happens to String theory claims about anthropic principle and all that?