# Effective Dynamics from Spinfoam Cosmology

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In collaboration with Etera Livine (paper in preparation, soon to appear)

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Effective dynamics from SF cosmology

### Introduction

- Extraction of effective (classical) dynamics from Spin foams
  - transition amplitudes between spin

networks,  $\psi_{in} \rightarrow \psi_{out}$ 

 Fixed graph: \u03c6<sub>in</sub> and \u03c6<sub>out</sub> same underlying graph



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finite graph  $\rightarrow$  finite number of dof  $\rightarrow$  minisuperspaces ??

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finite graph  $\rightarrow$  finite number of dof  $\rightarrow$  minisuperspaces ??

• Simple setting  $\rightarrow$  homogenous and isotropic configuration. Trying to modelize effective FRW models

# Spinfoam cosmology

Bianchi, Rovelli, Vidotto

# Spinfoam cosmology approach

#### Kinematics

- We choose a suitable graph and attach to it appropriate classical data: 3-geometry
- We reduce to the homogeneous and isotropic sector
  - $\blacktriangleright$  Finite graph  $\rightarrow$  finite region of the homogeneous geometry
- We define suitable coherent spin networks peaked on above symmetric configurations

# Spinfoam cosmology approach

#### **Dynamics**

- We calculate the transition amplitude between two such states at first order in a vertex expansion by using the Spinfoam ansatz (evaluation of the boundary spin network on the identity)
  - Renormalization/ coarse grain procedure in GFT would give an expansion in effective contributions. We assume such a one-vertex contribution is the leading order.
- We look for symmetries of the transition amplitude
  - $\rightarrow$  discrete diffeomorphisms

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- We look for symmetries of the transition amplitude
  - $\rightarrow$  discrete diffeomorphisms
    - $\blacktriangleright \text{ Homogeneity} \rightarrow \text{Hamiltonian constraint}$

(effective) FRW model ??

### Loop gravity on a fixed graph

- In LQG, phase space of gravity parametrized by holonomy-flux variables  $(g_e, X_e) \longrightarrow$  classical data attached to graphs  $\Gamma$
- Quantization of the holonomy-flux algebra: spin networks
- Given  $\Gamma$ , quantum states are gauge-invariant functions  $\psi_{\Gamma}(g_e)$  of the group elements  $g_e$  living on the edges  $e \in \Gamma$ 
  - ▶ Spin network basis: irreps. of SU(2) attached to edges, intertwiners [SU(2) invariant tensors] attached to vertices

#### Recent approach

- Phase space of loop gravity parameterized by spinors
  - Twisted geometries
     Freidel, Speziale, Livine, Tambornino
- Spin networks are the quantization of classical spinor networks
  - U(N) formalism for intertwiners

Girelli, Livine, Freidel, Dupuis...

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Effective dynamics from SF cosmology

### Choice of graph: two-vertex graph with N edges



- Two spinors  $z_i, w_i \in \mathbf{C}^2$  per edge, attached to  $\alpha$  and  $\beta$  resp.
- Poisson bracket:  $\{z_a, \bar{z}_b\} = -i\delta_{ab} = \{w_a, \bar{w}_b\}$ , a, b = 0, 1
- Closure contraints: SU(2) invariance in every vertex
- Matching constraints: U(1) invariance in every edge

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• 
$$|z\rangle\langle z| = \frac{1}{2}\left(\langle z|z\rangle\mathbb{I} + \vec{V}(z)\cdot\vec{\sigma}\right), \quad |z][z| = \frac{1}{2}\left(\langle z|z\rangle\mathbb{I} - \vec{V}(z)\cdot\vec{\sigma}\right)$$
  
•  $z_i$  defined by  $\vec{V}(z_i)$  up to a global phase

### 2-vertex spinor network: Geometrical interpretation



•  $ec{V}(z_i) = \langle z_i | ec{\sigma} | z_i 
angle$  vector normal to the i-th face

•  $|ec{V}(z_i)| = \langle z_i | z_i 
angle$  (twice) the area of the i-th face

• Closure constraints:  $C_{\alpha} \equiv \sum_{i} \vec{V}(z_i) = 0$ , analogously for  $\beta$ .

•  $\vec{V}(z_i)$  generators of SU(2) algebra  $\rightarrow$  fluxes



• Matching constraints:  $\mathcal{M}_i \equiv |ec{V}(z_i)| - |ec{V}(w_i)| = 0$ ,  $\forall i$ 

 $z_{\Lambda}$ 

 $w_i$ 

- U(N)-action on the set of spinors:  $z_k \to (Uz)_k$ ,  $U \in U(N)$ Commutes with closure constraint
- Scalar products between spinors are SU(2)-invariant

• 
$$E_{ij}^v = \langle z_i^v | z_j^v \rangle$$

★ They generate a U(N)-algebra

- $\blacktriangleright \ F^v_{ij} = [z^v_i | z^v_j \rangle, \quad \bar{F}^v_{ij} = \langle z^v_j | z^v_i ]$ 
  - ★ They close algebra with the above observables

### Symmetry reduction

• Matching constraints: invariance under  $U(1)^N$ 

$$\forall i, \quad \mathcal{M}_i = E_{ii}^{\alpha} - E_{ii}^{\beta} = 0 \longleftrightarrow \langle z_i | z_i \rangle = \langle w_i | w_i \rangle$$

• Symmetry reduction: imposing a larger symmetry

ightarrow invariance under U(N) Borja, Diaz-Polo, Garay, Livine  $orall i, j, \qquad \mathcal{E}_{ij} \equiv E^{lpha}_{ij} - E^{eta}_{ji} = 0, \qquad \mathcal{E}_{ii} = \mathcal{M}_i$ 

- Polyhedra dual to  $\alpha$  and  $\beta$  are identical  $\rightarrow$  homogeneity
- $\forall i \quad |w_i] = e^{i\phi} |z_i\rangle$ . Same phase for all edges  $\rightarrow$  isotropy
- Reduced phase space:  $\{A, \phi\} = 1$ ,  $A \equiv \frac{1}{2} \sum_{i} \langle z_i | z_i \rangle$ 

  - $\star~\phi$  matches the angle  $\xi$  parameterizing twisted geometries
  - ! Individual areas  $A_k = \langle z_k | z_k \rangle$  not imposed to be equal

### Quantum representation

- SU(2)-invariant observables:  $\hat{E}_{ij}$ ,  $\hat{F}_{ij}$ ,  $\hat{F}_{ij}^{\dagger}$
- Each space of N-valent intertwiners at fixed total area J  $\mathcal{R}^J = \bigoplus_{J=\sum_i j_i} \operatorname{Inv}_{SU(2)} \otimes_i V^i$ , carries an irrep. of U(N)•  $\hat{E}_{ij} : \mathcal{R}^J \to \mathcal{R}^J$  generator U(N)-action
- Whole space of N-valent intertwiners: H<sub>N</sub> = ⊕<sub>J</sub> R<sup>J</sup>
   → Fock structure
  - $\begin{array}{ll} \hat{F}_{ij}: \mathcal{R}^J \to \mathcal{R}^{J-1} & \quad \text{annhilation operator} \\ \hat{F}_{ij}^{\dagger}: \mathcal{R}^J \to \mathcal{R}^{J+1} & \quad \text{creation operator} \end{array}$
- $\bullet\,$  Quantum matching constraints in every edge  $\rightarrow\,$  spin networks

## ${\cal U}(N)$ coherent intertwiners

• Eigenstates of the annihilation operators  $\hat{F}_{ij}$ 

$$\begin{split} ||\{z_i\}\rangle &= \sum_J \frac{1}{J!(J+1)!} \left( \sum_{ij} [z_i|z_j\rangle \hat{F}_{ij}^{\dagger} \right)^J |0\rangle \\ &= \sum_{\{j_e\}} \frac{1}{\sqrt{\prod_e (2j_e)!}} ||\{j_e, z_e\}\rangle \\ &= \int dgg \triangleright (\mathsf{HO's-coherent state}) \end{split}$$

- U(N)-coherent:  $\hat{U}||\{z_i\}\rangle = ||\{(Uz)_i\}\rangle$
- Scalar product and norm explicitly known

$$\langle \{z_i\} || \{z_i\} \rangle = \sum_J \frac{A(z_i)^{2J}}{J!(J+1)!} = \frac{I_1(2A(z_i))}{A(z_i)}, \quad A(z_i) \equiv \frac{1}{2} \sum_i \langle z_i | z_i \rangle$$

### Basic transition amplitude

- SU(2)-BF theory
- One-vertex transition amplitude



• Symmetric configuration:  $|w_i] = e^{i\phi} |z_i\rangle \longrightarrow ||\{w_i\}\rangle = ||\{e^{i\phi}z_i\}\rangle$ 

### Basic transition amplitude

$$\mathcal{A}_{\sigma} = \langle \{e^{i\phi}z_i\} || \{z_i\} \rangle \langle \{e^{i\phi'}z_f\} || \{z_f\} \rangle$$
$$= \frac{e^{i\phi}I_1(2e^{-i\phi A})}{A} \frac{e^{i\phi'}I_1(2e^{-i\phi' A'})}{A'} = \psi_{in}(A,\phi)\psi_{out}(A',\phi')$$

Symmetries of  $\mathcal{A}_{\sigma}$ 

• 
$$\hat{C}\psi(A,\phi) = 0$$
,  $\hat{C} = A^2\partial_A^2 - 2A\partial_A + 2$ 

- Differential equation explicitly known
- No need to take the large area (spin) limit

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• 
$$\{A, \phi\} = 1 \quad \rightarrow \quad C = -A^2 \phi^2 - i2A\phi + 2$$

► Link with FRW (work in progress) (using SL(2, C) SF model Dupuis, Freidel, Livine Speziale)

### Outlook

- Using coherent intertwiners based on spinors it is possible to compute exactly basic transition amplitudes
- Within a simple setting, explicit realization of the recursion relations on boundary data (Bonzom, Freidel, Livine)
  - $\rightarrow$  symmetries of the transition amplitude
  - $\rightarrow$  effective classical dynamics
- Goal: to play around with the boundary data and the SF bulk to modelize specific models with physical interest (cosmology)