

Codimension-2 brane cosmology

Antonios Papazoglou

ICG Portsmouth

Ch. Charmousis, A.P., arXiv:0804.2121[hep-th]

Ch. Charmousis, G. Kofinas, A.P., arXiv:0907.1640 [hep-th]

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Outline

- 1 Motivations
 - Why modify gravity at large scales?
 - Self-acceleration and self-tuning
- 2 Codimension-2 braneworlds
 - Codimension-2 distributions
 - Explicit vacua and their self-properties
 - Types of boundary conditions
- 3 Codimension-2 cosmology
 - FRW-like equations
 - Self-tuning scenario
 - Self-accelerating cosmology

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Gravity at small and large scales

- GR in excellent agreement with precise gravity tests
(Damour, PDG 2007)

Equivalence principle: 10^{-13} accuracy

Dynamics of gravitational field: 10^{-5} accuracy

- However, the large scale observations puzzling

Composition of the Universe

$$\Omega_b \approx 0.05 \quad , \quad \Omega_{DM} \approx .25 \quad , \quad \Omega_{DE} \approx 0.7$$

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Composition of the Universe

$$\Omega_b \approx 0.05 \quad , \quad \Omega_{DM} \approx .25 \quad , \quad \Omega_{DE} \approx 0.7$$

- The last component the most puzzling for both
 - ★ high energy physics
 - ★ cosmology

Cosmological constant problem(s)

- Field theoretic prediction far from observable value

$$\Lambda_{\text{natural}} \sim M_{\text{Pl}}^4 \quad \dots \quad \Lambda_{\text{obs}} \sim 10^{-120} M_{\text{Pl}}^4$$

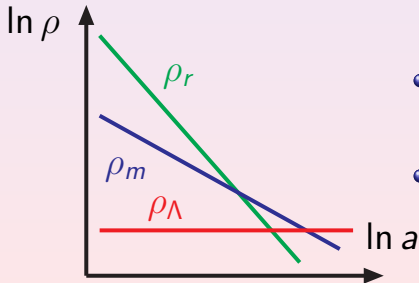
- No known symmetry which could enforce a vanishing vacuum energy and remain consistent with other observations

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- Cosmological constant finally **dominates**
- Why are we so close to the point that ρ_m and ρ_Λ intersect?

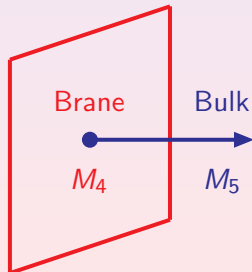
$$\Omega_\Lambda \sim \Omega_{DM} \sim 10^{-1} \Omega_b \quad !!!$$

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- The current acceleration is not due to exotic matter, but due to **different dynamics**
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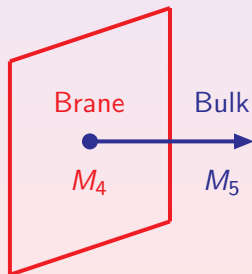


- Matter quantum loops induce an $M_4^2 \int \sqrt{g_4} R_4$ term on the brane
- There is a self-accelerating branch with (Deffayet, 2000)

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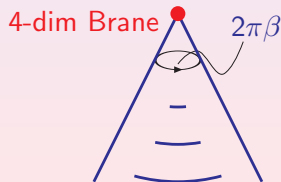
(however, with a ghost ...)

Self-tuning

- Mechanism to make vacuum energy **not gravitate**
- In brane worlds: transfer curvature from the brane to bulk
- Self-tuning: Curvature of the brane **insensitive to its vacuum energy** with no brane-bulk fine tuning

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- Mechanism to make vacuum energy **not gravitate**
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- Of special interest are the codimension-2 branes
(Chen, Luty, Ponton, 2000)



- In 2D sources do not curve the space but only introduce a deficit/excess angle
$$\text{Circ} \equiv 2\pi\beta = 2\pi - 4Gm$$

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Are codimension-2 distributions well defined?

- Distributional description for non-trivial codimension-2 sources **fails** in general (Geroch, Traschen, 1987)
- Sources should be regularised and the result will depend on the precise regularisation we pick
- Exception: pure vacuum energy branes

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- Exception: pure vacuum energy branes
- If we **increase the differential complexity** of the theory, the above can be avoided!
- Proposal: extension of bulk GR to **Lovelock gravity** (Bostock, Gregory, Navarro, Santiago, 2004)

Lovelock gravity

- Lovelock's theorem (1973): In $D = 4$ GR is the **unique** tensor theory which is
 - 1 Symmetric, $G_{\mu\nu} \equiv R_{\mu\nu} - \frac{1}{2}g_{\mu\nu}R$
 - 2 Divergence free, $\nabla^\mu G_{\mu\nu} = 0$
 - 3 Depends up to 2nd order derivatives of $g_{\mu\nu}$

Lovelock gravity

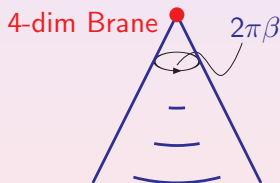
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- In $D > 4$ the theory is **extendable** to (curvature)² etc. additions. E.g. in $D = 6$ the unique theory is

$$\mathcal{S} = \frac{1}{16\pi G_6} \int d^6x \sqrt{-g} \left[R + \frac{\alpha}{6} (R^2 - 4R_{MN}^2 + R_{MKNL}^2) \right]$$

Gauss-Bonnet additional term

Brane gravity in a Gauss-Bonnet bulk

- Distributional description of general codim-2 sources **possible!**
- Brane junction conditions dictate an induced Einstein equation

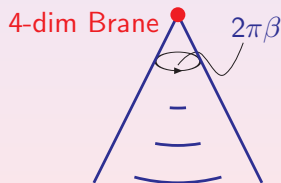


$$G_{\mu\nu} = 8\pi G_N T_{\mu\nu}$$

$$G_N = \frac{3G_6}{4\pi\alpha(1-\beta)}$$

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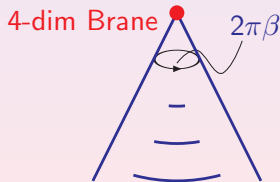
$$G_{\mu\nu} = 8\pi G_N T_{\mu\nu} + \frac{3}{2\alpha} g_{\mu\nu}$$

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$$G_{\mu\nu} = 8\pi G_N T_{\mu\nu} + \frac{3}{2\alpha} g_{\mu\nu} - W_{\mu\nu}$$



$$G_N = \frac{3G_6}{4\pi\alpha(1-\beta)}$$

$W_{\mu\nu} = "K_{\mu\nu}^2" \text{ terms}$

$K_{\mu\nu}$: brane extrinsic curvature

Maximally symmetric solutions

- There exist **exact** brane solutions for $K_{\mu\nu} = 0$:
 Double Wick rotated Gauss-Bonnet black holes

$$ds^2 = r^2 \left(-dt^2 + e^{2Ht} d\vec{x}^2 \right) + \frac{dr^2}{V(r)} + c^2 V(r) d\theta^2$$

- "Potential" $V(r)$ depends on α , Λ_6 and μ ("black hole" mass)

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- "Potential" $V(r)$ depends on α , Λ_6 and μ ("black hole" mass)
- Brane junction condition

$$H^2 = \underbrace{-\frac{1}{2\alpha}}_{\text{Geometric contribution}} + \underbrace{\frac{8\pi G_N T}{3}}_{\text{Brane content}}$$

Self-accelerating and self-tuning examples

- Self-acceleration exists for solutions with $T = 0$, e.g.

$$\mu = 0 \quad , \quad \alpha < 0 \quad , \quad \Lambda_6 = \frac{5}{2|\alpha|} \quad (\text{Born - Infeld limit})$$

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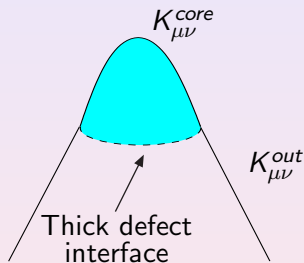
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- Then β is not constrained and $H^2 = \frac{1}{2|\alpha|}$
- Self-tuning exists for solutions with $H \approx 0$, when β self-tunes to cancel T , e.g.

$$\Lambda_6 = 0 \quad , \quad \alpha < 0 \quad , \quad \mu|\alpha|^{-3/2} \gg 1$$

- Then α is not constrained and G_N changes

Types of boundary conditions



- In general $K_{\mu\nu}^{core} \neq K_{\mu\nu}^{out}$
- Singularities structure is mixed
codimension-2 and **codimension-1**
(Kanno, Soda, 2004)
- If ones want to focus to pure
codimension-2, $K_{\mu\nu}$ should be **continuous**
(topological boundary conditions of
Charmousis, Zegers, 2005)
- $W_{\mu\nu}$ term not depending on β
(different from Bostock et al)

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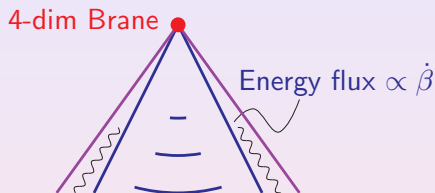
FRW-like equations on codimension-2 branes

- Brane Einstein equations + leading order of bulk equations on the brane: 3 **independent** equations

$$\begin{aligned}
 H^2 &= \frac{8\pi G_N}{3}\rho - \frac{\kappa}{a^2} - \frac{1}{2\alpha} + A^2 \\
 \frac{\ddot{a}}{a} &= -\frac{4\pi G_N}{3}(\rho + 3P) - \frac{1}{2\alpha} + \frac{3P}{\rho}A^2 \\
 \dot{\rho} + 3H(\rho + P) + \frac{\rho \dot{\beta}}{\beta(1-\beta)} &= 0
 \end{aligned}$$

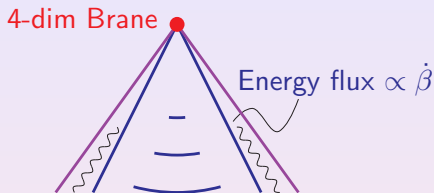
- New degrees of freedom $\beta(t)$ and $K_i^j = A(t) \delta_i^j$
- One free function - the system is not closed!

Self-tuning scenario



- Vacuum energy can be (in principle) relaxed when β changes
- Bulk boundary conditions necessary to determine $\beta(t)$
- $G_N \propto (1 - \beta)^{-1}$ is time dependent

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- $G_N \propto (1 - \beta)^{-1}$ is time dependent
- In the following $\beta = \text{constant}$
[This is a necessity if one regularises the codim-2 brane by a ring codim-1 brane]

Self-accelerating cosmology

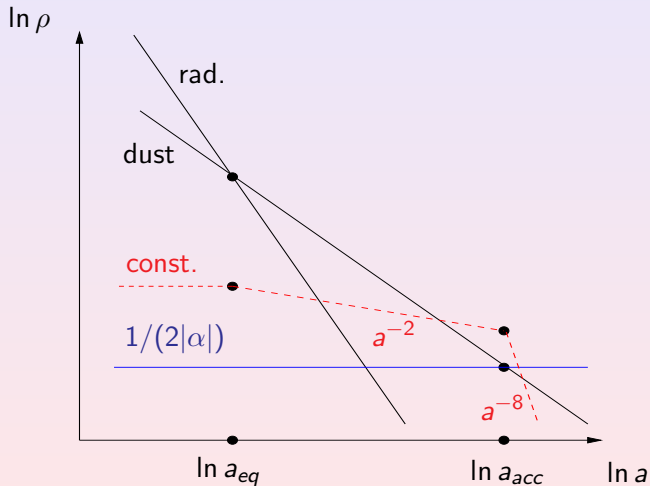
- For $\beta = \text{constant}$ the system is closed
- Friedmann equation for tensionless branes with $P = w \rho$

$$H^2 = \frac{8\pi G_N}{3} \rho - \frac{\kappa}{a^2} - \frac{1}{2\alpha} + \underbrace{C^2 \rho^{\frac{2(1-3w)}{3(1+w)}}}_{\propto a^{-2(1-3w)}}$$

- Late geometrical acceleration if $\alpha < 0$
- Correction term behaves like a **mirage fluid**

w	ρ	Correction
-1	const.	a^{-8}
0	a^{-3}	a^{-2}
1/3	a^{-4}	const.

A possible evolution



Conclusions

- Distributional description of codim-2 possible in 6D Lovelock gravity
- There exist self-accelerating and self-tuning solutions
- Friedmann equation receives corrections
 - geometrically induced vacuum energy
 - mirage fluid depending on w

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- There exist self-accelerating and self-tuning solutions
- Friedmann equation receives corrections
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 - mirage fluid depending on w
- Stability of these vacua should be studied (ghosts?)
- Challenge for phenomenology to constrain C^2
⇒ not-trivial signature of codimension-2