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Regularized Codimension-2 Brane Cosmology

In collaboration with E. PAPANTONOPOULOS and A. PAPAZOGLOU hep-th/0611311 and 0707.1396 [hep-th]

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28 Sept. 2007

OUTLINE

- I. Setup and Static Solution of the 6D warped flux compactified model.
- II. Cosmology of the 4-Brane.
- III. Conclusions.

I. SETUP AND STATIC SOLUTION

M. Peloso, L. Sorbo, G. Tasinato (2006) E. Papantonopoulos, A. Papazoglou, VZ (2006)



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I. SETUP AND STATIC SOLUTION

➢ Metric Solution:

$$ds_{6}^{2} = z(r,\alpha)^{2} \eta_{\mu\nu} dx^{\mu} dx^{\nu} + R_{i}^{2} \left[\frac{dr^{2}}{f(r,\alpha)} + c_{i}^{2} f(r,\alpha) d\varphi^{2} \right],$$

$$F_{r\varphi} = -c_{i} R_{i} M^{2} S(\alpha) \frac{1}{z(r,\alpha)^{4}}$$

with
$$z(r,\alpha) = \frac{1}{2} \Big[(1-\alpha)r + (1+\alpha) \Big],$$

 $f(r,\alpha) = \frac{1}{5(1-\alpha)^2} \Big[-z(r)^2 + \frac{1-\alpha^8}{1-\alpha^3} \frac{1}{z(r)^3} - \alpha^3 \frac{1-\alpha^5}{1-\alpha^3} \frac{1}{z(r)^6} \Big],$
 $S(\alpha) = \sqrt{\frac{3}{5}\alpha^3 \frac{1-\alpha^5}{1-\alpha^3}}.$

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I. SETUP AND STATIC SOLUTION

✓ The cap is smooth at r = +1 as long as

$$c_c = \frac{1}{X_+}$$
 with $X_+ = \frac{20(1-\alpha)(1-\alpha^3)}{5+3\alpha^8-8\alpha^3}$

✓ The conical singularity at r = -1 is supported by a codim-2 brane with tension

$$T = 2\pi M^4 \left(1 - c_o X_- \right) \text{ with } X_- = \frac{3 + 5\alpha^8 - 8\alpha^5}{20\alpha^4 \left(1 - \alpha \right) \left(1 - \alpha^3 \right)}$$

✓ Restriction condition from quantum numbers

$$n = \frac{N}{2} \frac{2}{\left(1 - \alpha^3\right)} \left[\frac{5\left(1 - \alpha^3\right)}{8\left(1 - \alpha^5\right)} - \alpha^3 \right] \implies \alpha \text{ takes discrete values}$$

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II. <u>COSMOLOGY OF THE 4-BRANE</u>

E. Papantonopoulos, A. Papazoglou, VZ (2007)

- ➤ In general: inclusion of matter contribution ⇒ evolution of both brane and bulk. But as a first step we want to avoid time dependent bulk solution,
- \Rightarrow Approach: Bulk remains static while the brane matter merely makes the brane to move between the static bulk and the static cup with position R(t). Here the brane is not merely a probe brane and we will use Junction Conditions (we include backreaction of the brane energy density).
- Brane coordinates: $\sigma^{\hat{\mu}} = (t, x^i, \varphi).$
- Brane embedding X^{M} : $X^{i} = x^{i}$, $X^{r} = R(t)$ and $X^{\varphi} = \varphi$, Outer bulk section: $X^{0}_{(out)} = t$, Inner cap section: $X^{0}_{(in)} = T(t)$.

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- ✓ We put *matter* on the brane (perfect fluid), 6 parameters: ρ , P, \hat{P} , l, L and \hat{L} .
- \Rightarrow The brane *MOVES* in the *static* bulk.
- \Rightarrow This movement *induces cosmology* on the brane with the need of *warping* in the bulk.
- \checkmark Induced metric on the brane

$$ds_{(5)}^{2} = -d\tau^{2} + a^{2}(\tau)d\vec{x}^{2} + b^{2}(\tau)d\varphi^{2}.$$

 \checkmark Relations derived from the continuity of the induced metric.

$$c_0 R_0 = c_c R_c, \quad \dot{T}^2 \left(1 - \beta_+^2 \frac{\dot{R}^2}{\dot{T}^2} \frac{R_0^2}{fz^2} \right) = \left(1 - \dot{R}^2 \frac{R_0^2}{fz^2} \right)$$

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> Hubble parameters for the two scale factors:

$$H_{a} = \frac{1}{a} \frac{da}{d\tau} = \frac{z'}{z} \frac{\dot{R}}{\sqrt{1 - \dot{R}^{2}} \frac{R_{0}^{2}}{fz^{2}}}, \qquad H_{b} = \frac{1}{b} \frac{db}{d\tau} = \frac{f'}{2fz} \frac{\dot{R}}{\sqrt{1 - \dot{R}^{2}} \frac{R_{0}^{2}}{fz^{2}}}$$
The two Hubble rates are related:
$$H_{a} = \frac{zf'}{z} \frac{H}{\sqrt{1 - \dot{R}^{2}} \frac{R_{0}^{2}}{fz^{2}}}$$

 \Rightarrow The two Hubble rates are related:

$$H_b = \frac{zf'}{2fz'}H_a$$

Close to the would-be conical singularity we have: f' < 0

If
$$H_a > 0 \implies H_b < 0$$

(If 4D space expands then internal space shrinks)

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> Junction conditions:

$$\left\{ \hat{K}_{\hat{\mu}\hat{\nu}} \right\} = -\frac{1}{M^4} t_{\hat{\mu}\hat{\nu}}^{(br)}, \quad \text{where } \left\{ H \right\} = H^{in} + H^{out}$$
$$\left\{ n_M F_N^M \partial_{\hat{\kappa}} X^N \right\} = -\frac{\delta S_{br}}{\delta a^{\hat{\kappa}}}.$$

- ***** RESULTS:
- ✓ 2 Junction conditions for the gauge field (the (τ) and (*i*) components) give for the coupling: l = L = 0,
- ✓ The (φ) component of the gauge field junction and the (φφ) component of the metric junction determine \hat{P}, \hat{L} ,
- ✓ The 2 remaining metric junction conditions give the Friedmann equation and the acceleration equation.

> The Friedmann Equation:

$$H_a^2 = C_1(a)\rho^2 + \frac{C_2(a)}{\rho^2} + C_3(a),$$

where $C_i(a)$ are expressed as known functions of f and z.

> Expansion around the static case $\rho_m^{(4)} \ll \rho_0$: where the effective 4D $\rho_m^{(4)} = \int d\varphi \sqrt{g_{\varphi\varphi}} \rho_m = \frac{2\pi\beta_+}{X_+} R_0 \sqrt{f} \rho_m$

$$H_{a}^{2} = \frac{8\pi}{3} G_{eff}(a) \rho_{m}^{(4)} + \Lambda_{eff}(a) + O(\rho_{m}^{(4)2})$$

 Λ_{eff} : Mirage matter contribution

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II. COSMOLOGY OF THE 4-BRANE



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\succ Early Times Expansion of the Friedmann Equation: $\rho_m^{(4)} \gg \rho_0$:

 $H_a^2 = C_4(a)\rho_m^2$, 5D Freidmann Law

> We can also derive the energy continuity equation:

$$\frac{d\rho_{tot}}{d\tau} + 3(\rho + P)H_a + (\rho + \hat{P})\frac{zf'}{2fz'}H_a = -W(a, H_a)$$

Significant Energy Flow from brane to Bulk !!!

III. <u>CONCLUSIONS</u>

- First step for getting cosmology in a regularized codimension-2 brane.
- ➤ We use the simplest scenario of a static bulk and a moving brane ⇒ unsatisfactory cosmology!
- D. Langlois & M. Minamitsuji [0707.1426 hep-th]: Time-dependent Scalar Field.

Need of *time-dependent bulk solution*. (work in progress...)

K. i. Maeda & H. Nishino, Phys. Lett. B 154 (1985) 358
C. P. Burgess & al. [hep-th/0608083]
T. Kobayashi & M. Minamitsuji [0705.3500 hep-th]
E. J. Copeland & O. Seto [0705.4169 hep-th]

IV. EXTRA SLIDES

• Brane matter (Energy Momentum Tensor):

$$t_{\hat{\mu}}^{\hat{\nu}(br)} = -\frac{2}{\sqrt{-\gamma_{+}}} \frac{\delta S_{br}}{\delta \gamma_{+\nu}^{\mu}} = \text{diag}\left(-\rho, P, P, P, \hat{P}\right)$$

with $\rho = \rho_{0} + \rho_{m} = \lambda + \frac{\nu^{2} \left(n - eA_{\phi}^{+}\right)^{2}}{2c_{0}^{2}R_{0}^{2}f(r_{c})} + \rho_{m}$
 $P = P_{0} + P_{m} = -\rho_{0} + P_{m}$
 $\hat{P} = \hat{P}_{0} + \hat{P}_{m} = -\lambda + \frac{\nu^{2} \left(n - eA_{\phi}^{+}\right)^{2}}{2c_{0}^{2}R_{0}^{2}f(r_{c})} + \hat{P}_{m}$

• Coupling of the bulk gauge field to the brane matter :

$$\frac{\delta S_{br}}{\delta a^{\hat{\kappa}}} = (l, L, L, \hat{L}, \hat{L}) \quad \text{with} \quad \begin{array}{l} l = l_m \\ L = L_m \\ \hat{L} = \hat{L}_0 + \hat{L}_m = ev^2 \left(n - eA_{\varphi}^+\right) + \hat{L}_m \end{array}$$

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➤ Induced metric on the brane:

Need of warping

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