

Melting Domain Walls in the Early Universe as the origin of the NANOGrav signal and Dark Matter

Dmitry Gorbunov

Institute for Nuclear Research of RAS, Moscow

Workshop on Standard Model
and Beyond

Mon Repos, Corfu, Greece



(If) particles:

- 1 stable on cosmological time-scale
- 2 nonrelativistic long before RD/MD-transition (either Cold or Warm, free-streaming is limited $v_{RD/MD} \lesssim 10^{-3}$)
- 3 (almost) collisionless
- 4 (almost) electrically neutral
- 5 adiabatic matter perturbations

If were in thermal equilibrium:

$$M_X \gtrsim 1 \text{ keV}$$

If not:

$$\lambda = 2\pi/(M_X v_X), \text{ in a galaxy } v_X \sim 0.5 \cdot 10^{-3} \longrightarrow M_X \gtrsim 3 \cdot 10^{-22} \text{ eV}$$

for bosons

$$M_X \gtrsim 750 \text{ eV}$$

for fermions

Pauli blocking:

$$f(\mathbf{p}, \mathbf{x}) = \frac{\rho_X(\mathbf{x})}{M_X} \cdot \frac{1}{\left(\sqrt{2\pi} M_X v_X\right)^3} \cdot e^{-\frac{\mathbf{p}^2}{2M_X^2 v_X^2}} \Big|_{\mathbf{p}=0} \leq \frac{g_X}{(2\pi)^3}$$

Microscopic processes in the expanding Universe

A competition between scattering, decays, etc and expansion

for general processes one should solve kinetic equations

$$\frac{dn_{X_i}}{dt} + 3Hn_{X_i} = \sum (\text{production} - \text{destruction})$$

Boltzmann equation in a comoving volume: $\frac{d}{dt}(na^3) = a^3 \int \dots$

production:

$$\sigma(A+B \rightarrow X+C)n_A n_B, \quad \Gamma(D \rightarrow E+X)n_D \cdot M_D/E_D, \quad \text{etc}$$

desrtuction:

$$\sigma(A+X \rightarrow C+B)n_A n_X, \quad \Gamma(X \rightarrow F+G)n_X \cdot M_X/E_X, \quad \text{etc}$$

Fast direct and inverse processes, $\Gamma \gtrsim H$, are in equilibrium,

$\Sigma(\) = 0$ and thermalize particles

Freeze-out of nonrelativistic Dark Matter

Assumptions:

- ① no $X - \bar{X}$ asymmetry either $X = \bar{X}$ or $n_X = n_{\bar{X}}$
- ② @ $T \lesssim M_X$ in thermal equilibrium with plasma (e.g. neutrons)

$$n_X = n_{\bar{X}} = g_X \left(\frac{M_X T}{2\pi} \right)^{3/2} e^{-M_X/T}$$



freeze-out temperature T_f $H \equiv T^2/M_{\text{Pl}}^*$, $M_{\text{Pl}}^* = M_{\text{Pl}}/1.66\sqrt{g_*}$

$$n_X \langle \sigma_{\text{ann}} v \rangle = H(T_f) \longrightarrow T_f = \frac{M_X}{\ln \left(\frac{g_X M_X M_{\text{Pl}}^* \sigma_0}{(2\pi)^{3/2}} \right)} .$$

Bethe formula:

$$\text{s-wave: } \sigma_{\text{ann}} = \frac{\sigma_0}{v}$$

Weakly Interacting Massive Particles

density after freeze-out:

$$n_x(T_f) = \frac{T_f^2}{M_{Pl}^* \sigma_0}$$

present density:

$$n_x(T_0) = \left(\frac{a(T_f)}{a(T_0)} \right)^3 n_x(T_f) = \left(\frac{s_0}{s(T_f)} \right) n_x(T_f) \propto \frac{1}{T_f}$$

$X + \bar{X}$ contribution to critical density:

$$\begin{aligned} \Omega_x &= 2 \frac{M_x n_x(T_0)}{\rho_c} = 7.6 \frac{s_0 \ln \left(\frac{g_x M_{Pl}^* M_x \sigma_0}{(2\pi)^{3/2}} \right)}{\rho_c \sigma_0 M_{Pl} \sqrt{g_*(T_f)}} \\ &= 0.1 \cdot \left(\frac{(10 \text{ TeV})^{-2}}{\sigma_0} \right) \frac{10}{\sqrt{g_*(T_f)}} \ln \left(\frac{g_x M_{Pl}^* M_x \sigma_0}{(2\pi)^{3/2}} \right) \cdot \frac{1}{2h^2} \end{aligned}$$

Dark Matter: many well-motivated candidates

- WIMPs related to EW scale, SUSY
- sterile neutrinos active neutrino oscillations
- light scalar field string theory
- axion strong CP-problem
- gravitino local SUSY
- Heavy relics GUTs
- (Topological) defects GUTs
- Massive Astrophysical Compact Heavy Objects
- Primordial black hole (remnants) Phase transitions
exotic inflation, reheating

Multicomponent Dark Matter ?

γ , v , H, He

A simple example of scalar DM

most general renormalizable coupled to SM:

Z_2 -invariant Higgs (Φ) portal

$$\Delta\mathcal{L} = \frac{1}{2}g^{\mu\nu}\partial_\mu X\partial_\nu X - \frac{1}{2}M^2X^2 + g^2X^2\Phi^\dagger\Phi - \frac{\lambda}{4}X^4$$

Options:

- freeze-out: sufficiently large g^2

$$v\sigma_{XX \rightarrow hh} \times n_h \sim H \rightarrow \Omega_X \propto \frac{1}{\sigma_0}, \text{ with } \frac{g^4}{(4\pi\ldots)^2 M^2} = \sigma_0 \equiv \sigma v$$

- freeze-in: intermediate g^2

$$\dot{n}_X + 3Hn_X = \sigma_{hh \rightarrow XX} n_h^2 \rightarrow \frac{n_X}{s} = \# \int dT \frac{n_h^2}{sHT} \times \frac{g^4}{T^2} \sim g^4 \frac{M_{Pl}}{M} \rightarrow$$

$$\Omega_X \propto g^4 \rightarrow g^2 \approx 10^{-11} \quad \text{still natural...}$$

Freeze in via gravitational scatterings..?

any particles A in plasma

$$\sigma_{AA \rightarrow XX} \propto \frac{T^2}{M_{Pl}^4} \rightarrow \Omega_X \propto M_X \frac{T_i^3}{M_{Pl}^3} \dots$$

assuming $m \ll T_i$

called “unnatural” being dependent on the initial conditions

Free massive scalar field

$$g^2 = 0$$

$$\mathcal{L} = \frac{1}{2} g^{\mu\nu} \partial_\mu \phi \partial_\nu \phi - \frac{1}{2} m_\phi^2 \phi^2$$

Homogeneous scalar field in the expanding Universe

$$\ddot{\phi} + 3H\dot{\phi} + m_\phi^2 \phi = 0$$

Two-stage evolution:

$$m_\phi < H(t) \implies \phi = \phi_i = \text{const}$$

$$m_\phi > H(t) \implies p = \langle E_k \rangle - \langle E_p \rangle = 0, \quad \rho \sim m_\phi^2 \phi^2 \propto 1/a^3$$

- dust-like substance in the late Universe, $\Omega \propto m_\phi^{1/2} \phi_i^2$
depends on initial conditions
fuzzy DM

DM from oscillating scalar

$$0 \neq g^2 < 10^{-11}$$

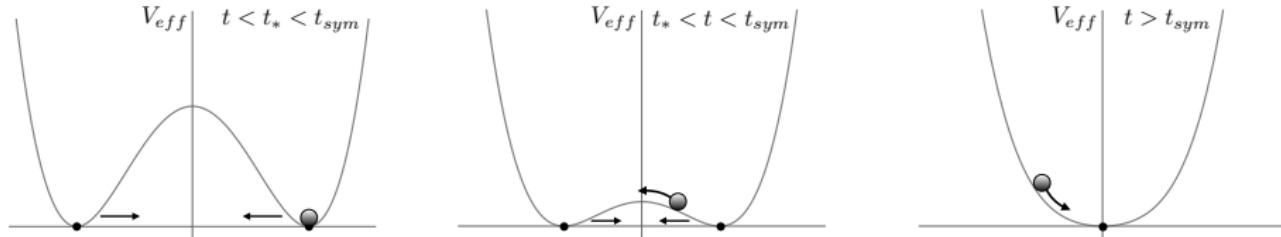
Z_2 -invariant Higgs (Φ) portal

$$\Delta\mathcal{L} = \frac{1}{2}g^{\mu\nu}\partial_\mu X\partial_\nu X - \frac{1}{2}M^2X^2 + g^2X^2\Phi^\dagger\Phi - \frac{\lambda}{4}X^4$$

Higgs particles in plasma change the potential:

$$g^2X^2\Phi^\dagger\Phi \rightarrow g^2X^2T^2/3$$

Z_2 symmetry is broken after reheating by the plasma contribution

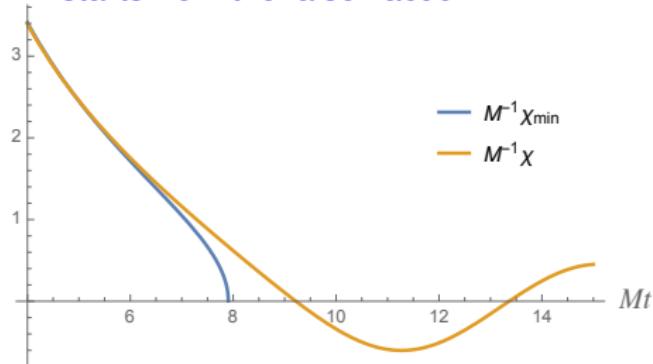


Temperature decrease restores Z_2

2004.03410

$$\Delta \mathcal{L} = \frac{1}{2} g^{\mu\nu} \partial_\mu X \partial_\nu X - \frac{1}{2} M^2 X^2 + g^2 X^2 T^2 / 3 - \frac{\lambda}{4} X^4$$

X starts from the false vacuum



at $g^2 T_*^2 \simeq M^2$ sign changes
and X starts to oscillate
gravitational misalignment

$$\rho_{DM}(t_*) = \frac{M^2 \cdot S_*^2}{2} \simeq \frac{(M^5 H_*)^{2/3}}{4\lambda}$$

And the correct amount of DM by classical oscillating field

$$p = \langle E_{kin} \rangle - \langle E_p \rangle = 0$$

$$g^2 \simeq 10^{-12} \times \left(\frac{\lambda}{10^{-6}} \right)^{6/5} \times \left(\frac{10^6 \text{ GeV}}{M} \right)^2$$

Dark Matter

with general setup:

white area

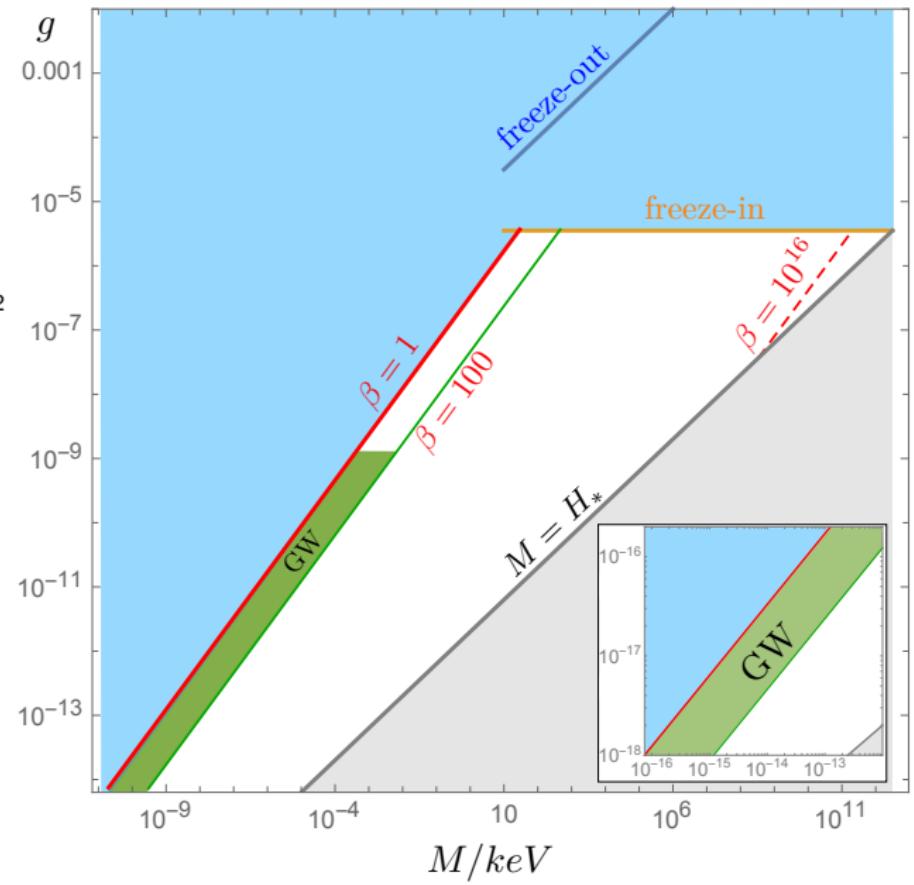
$$\beta \equiv \lambda/g^4 < 1$$

$$V = \frac{1}{2} M^2 X^2 + \frac{\lambda}{4} X^4 - \frac{g^2 T^2}{12} X^2$$

The inverse phase transition
may be accompanied by the
production of GW

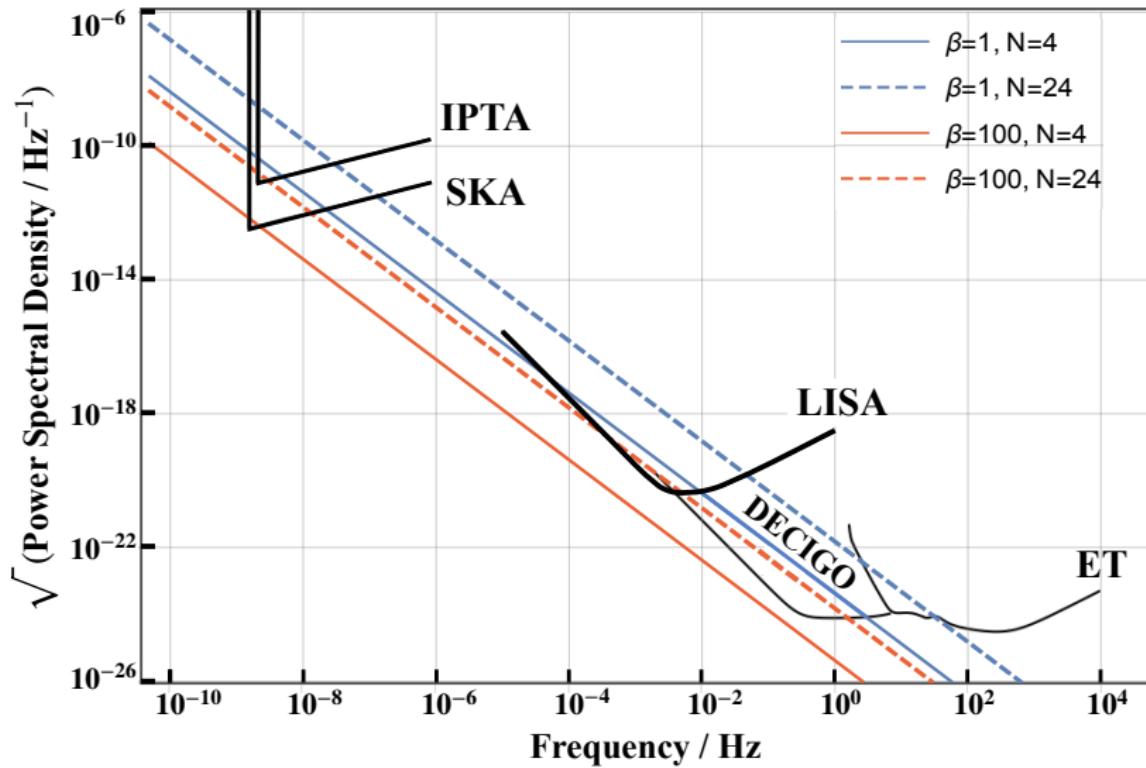
strong enough to be
detected by the present or
next generation
experiments

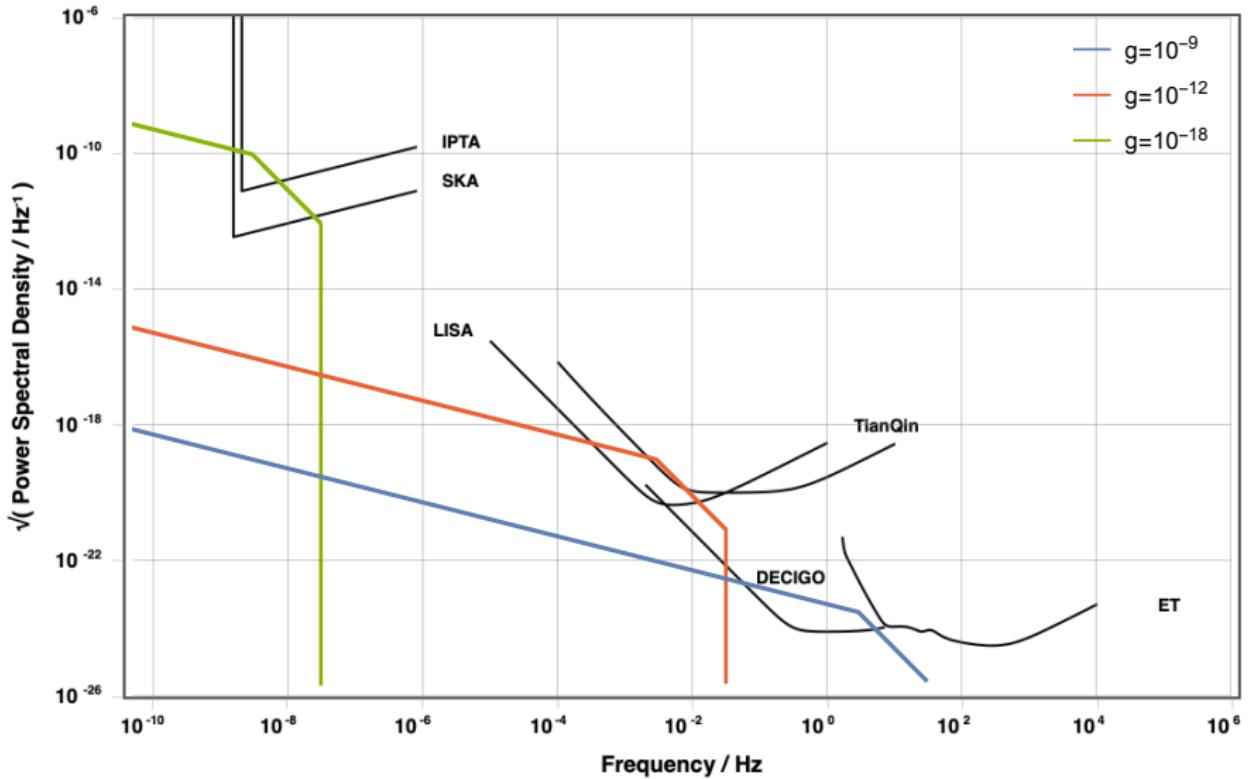
2104.13722



GW signals from dissipating DWs ...

2104.13722





strain: $\Omega_{GW} H_0^2 \equiv 2\pi^2 f^3 S/3$

2104.13722

Analog of a well-known effect in CMB:

by Rubakov, Sazhin and Veryaskin

From linear approximation to
the geodesic equation...
for tensor perturbations

$$\frac{\delta T}{T}(\mathbf{n}, \eta_0) = \frac{1}{2} \int_{\eta_r}^{\eta_0} d\eta n_i h_{ij}^{TT'} n_j ,$$

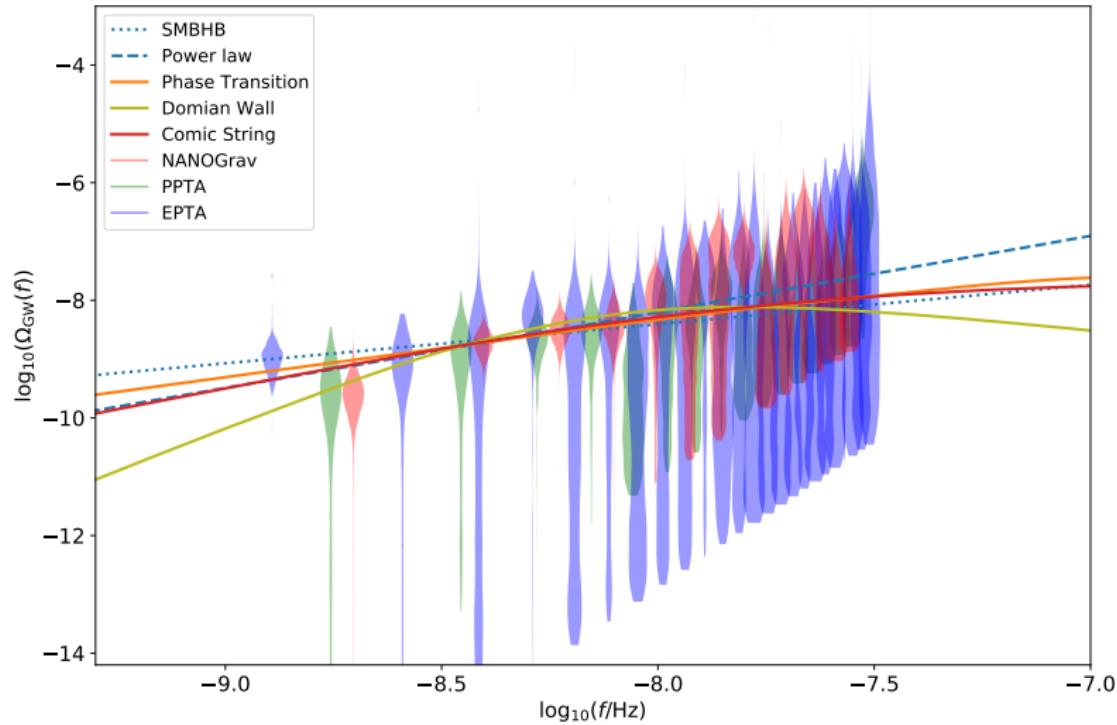
strongest limit on primordial
GW and Inflation scale:

$$h_{init} \sim H_{inf}/M_{Pl} \propto \sqrt{\rho_{infl}}$$

Pulsar frequencies are
different but stable, and we
observe its' time variation

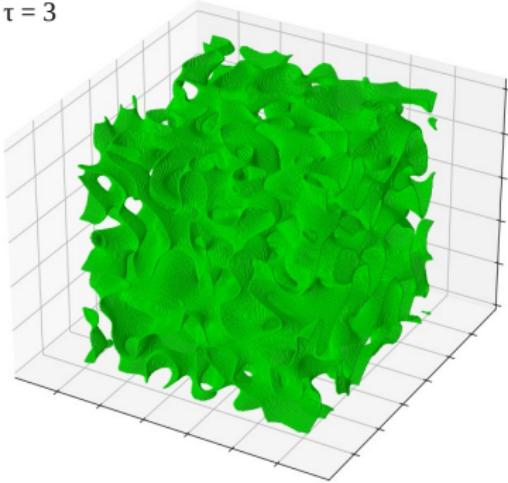
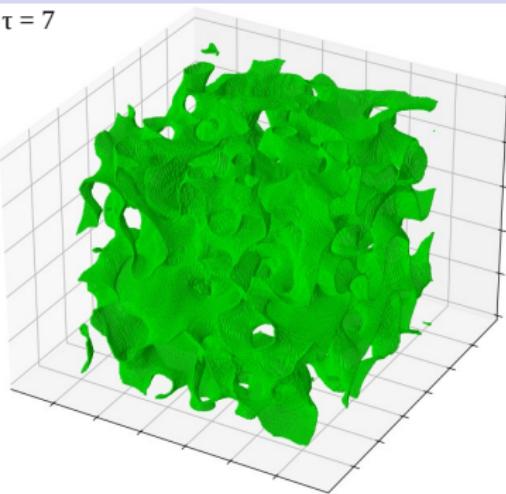
NANOGrav, T. Klein



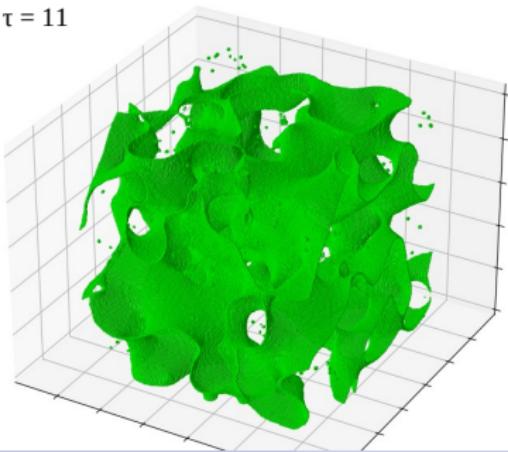
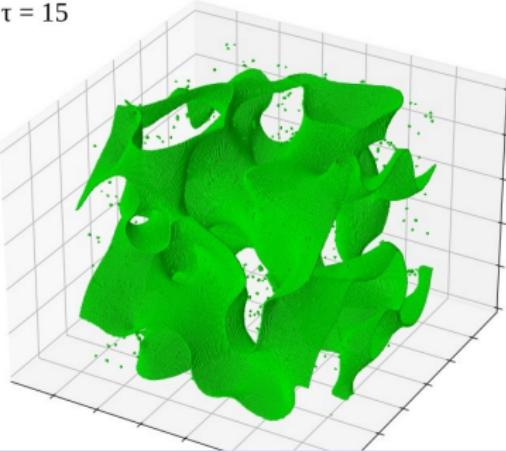


$$\Omega_{GW} h^2 = 6.3 \cdot 10^{-10} A^2 \left(\frac{f}{\text{yr}^{-1}} \right)^d, \quad A = \{ 6.4^{+4.2}_{-2.7}, 2.9^{+2.6}_{-1.8}, 3.1^{+1.3}_{-0.9} \}, \quad d = \{ 1.8 \pm 0.6, 0.81^{+0.63}_{-0.73}, 1.1 \pm 0.4 \}$$

2307.03141, 2406.02288

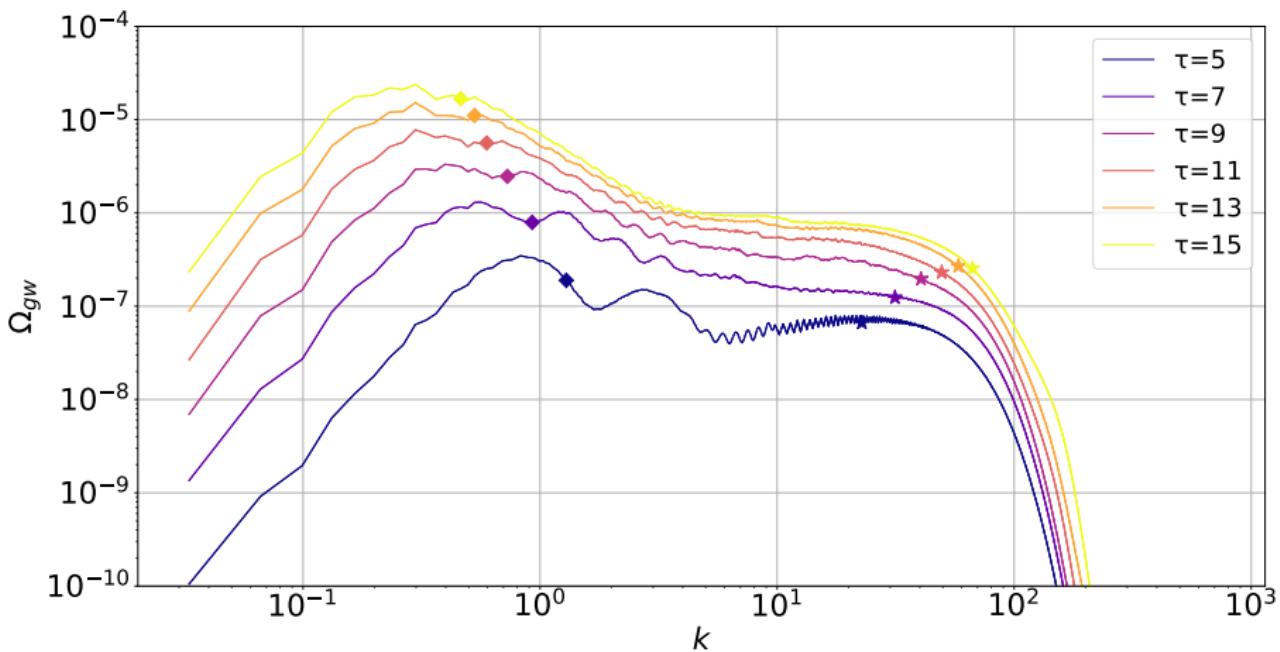
$\tau = 3$  $\tau = 7$ 

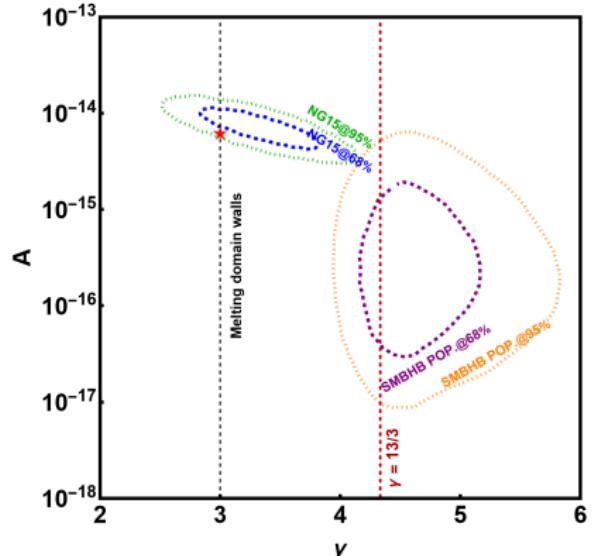
2406.17053

 $\tau = 11$  $\tau = 15$ 

GW spectra: simulation with CosmoLattice

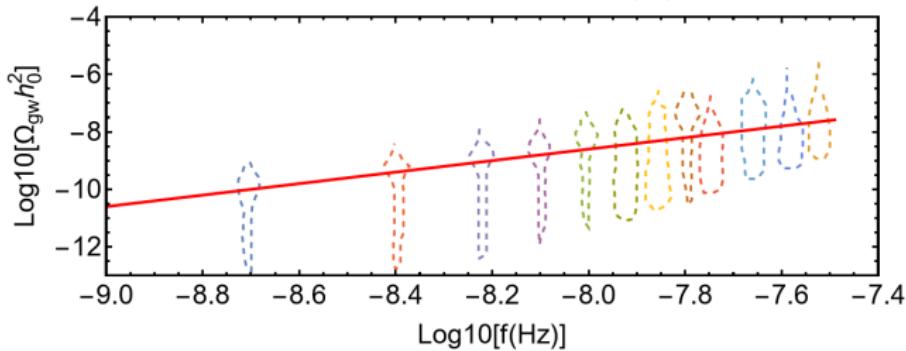
2406.17053

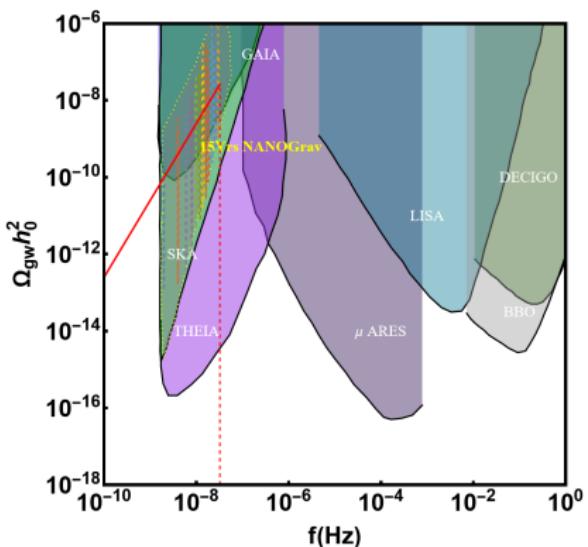
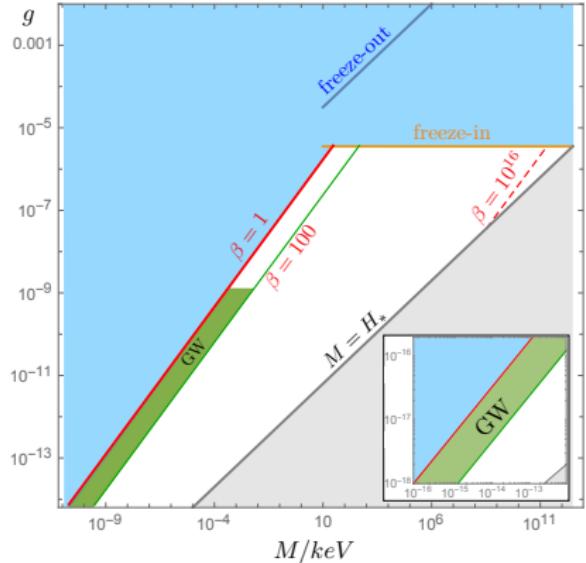




$$\gamma \equiv d + 1$$

2307.04582



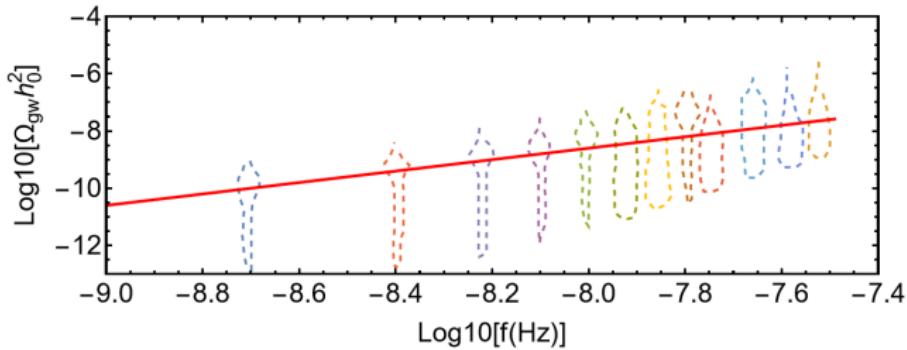


2104.13722

$$\beta \equiv \lambda/g^4$$

$$V = \frac{1}{2} M^2 X^2 + \frac{\lambda}{4} X^4 - \frac{g^2 T^2}{12} X^2$$

2307.04582



Conclusion

- What NANOGrav and others observe might be explained by the GW from melting domain walls
- they are expected in models with inverse phase transition
- which may induce light scalar dark matter production
- In realistic models its mass is of order 10^{-15} - 10^{-11} eV
- That predicts super-radiance instability of rotating black holes with astrophysical masses $10\text{-}10^5 M_\odot$

see e.g. 0905.4720

Might explain why the BH in binary system are not with $J \approx 1$

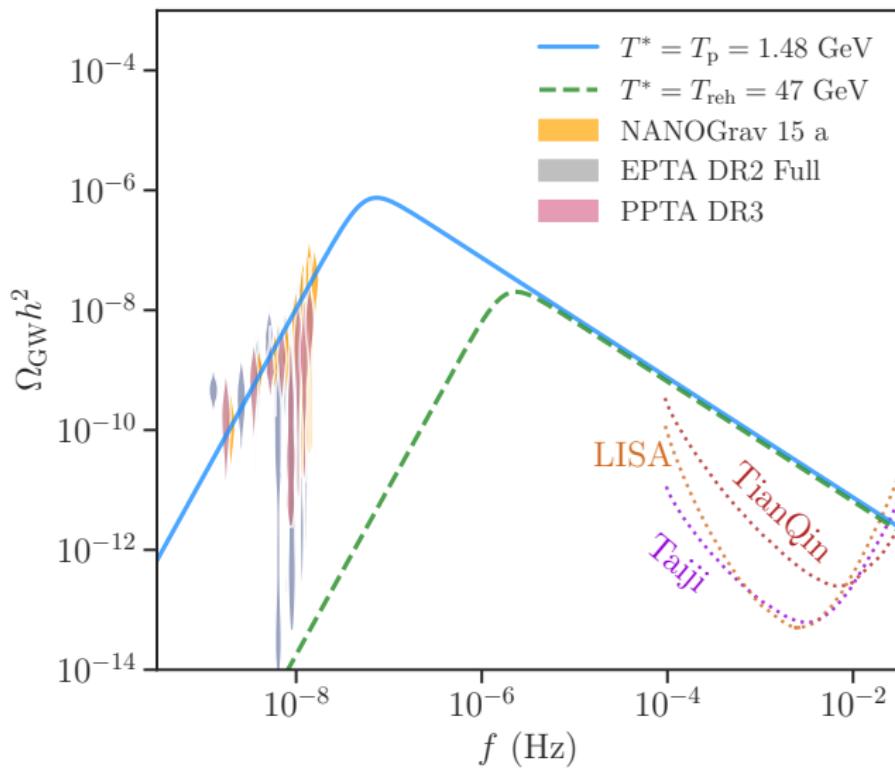
Large based on the results from

2004.03410, 2104.13722, 2112.12608, 2307.04582, 2406.17053

with Evgeny Babichev, Ivan Dankovsky, Sabir Ramazanov, Rome Samanta and Alex Vikman

Example: EW 1st order Phase Transition

2307.01072



$$\mathcal{L} = \frac{1}{2} g^{\mu\nu} \partial_\mu \phi \partial_\nu \phi - \frac{1}{2} m_\phi^2 \phi^2$$

Homogeneous scalar field in the expanding Universe

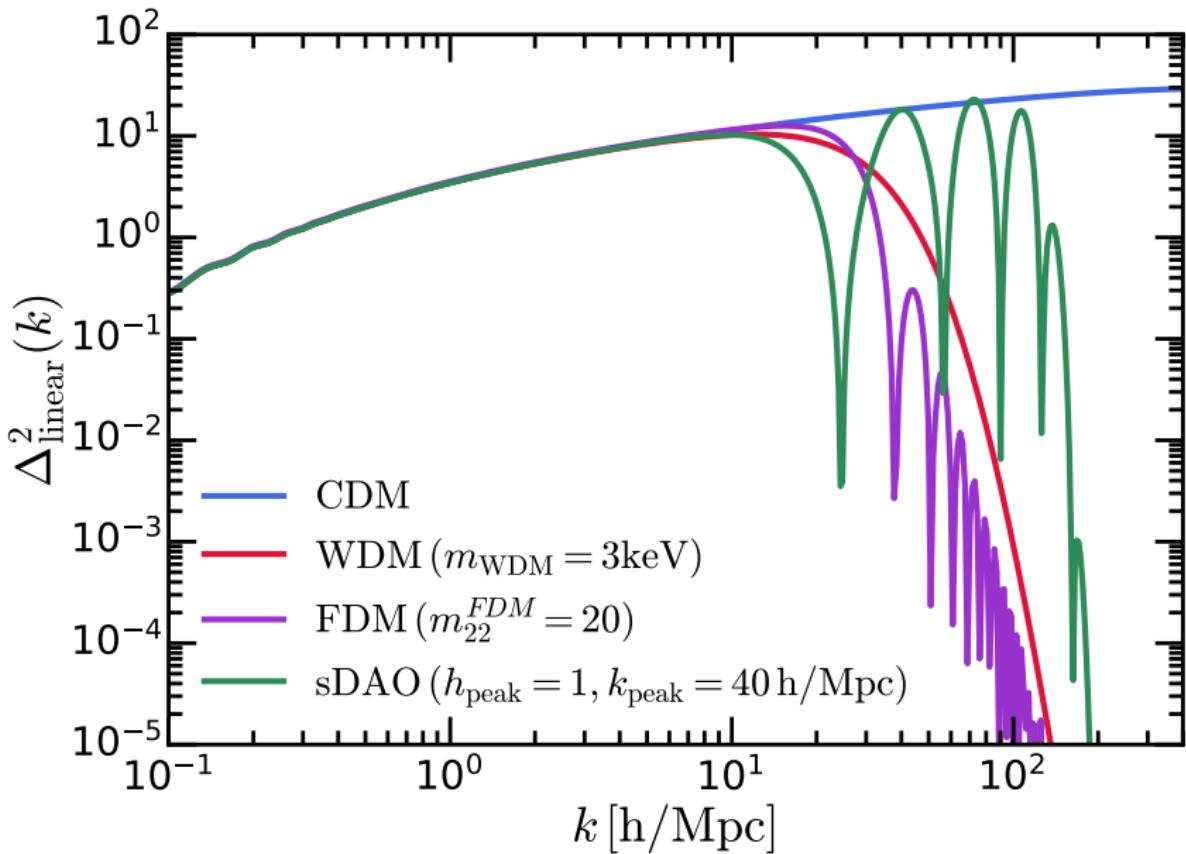
$$\ddot{\phi} + 3H\dot{\phi} + m_\phi^2 \phi = 0$$

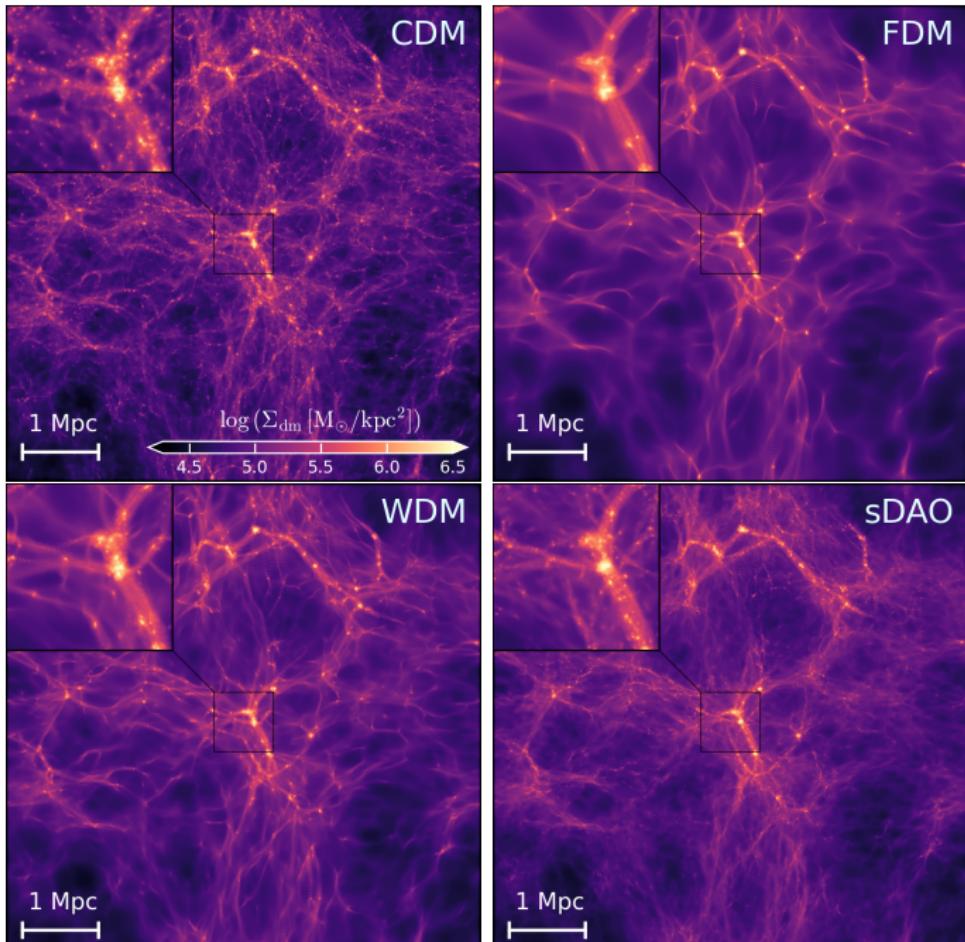
Two-stage evolution:

$$m_\phi < H(t) \implies \phi = \phi_i = \text{const}$$

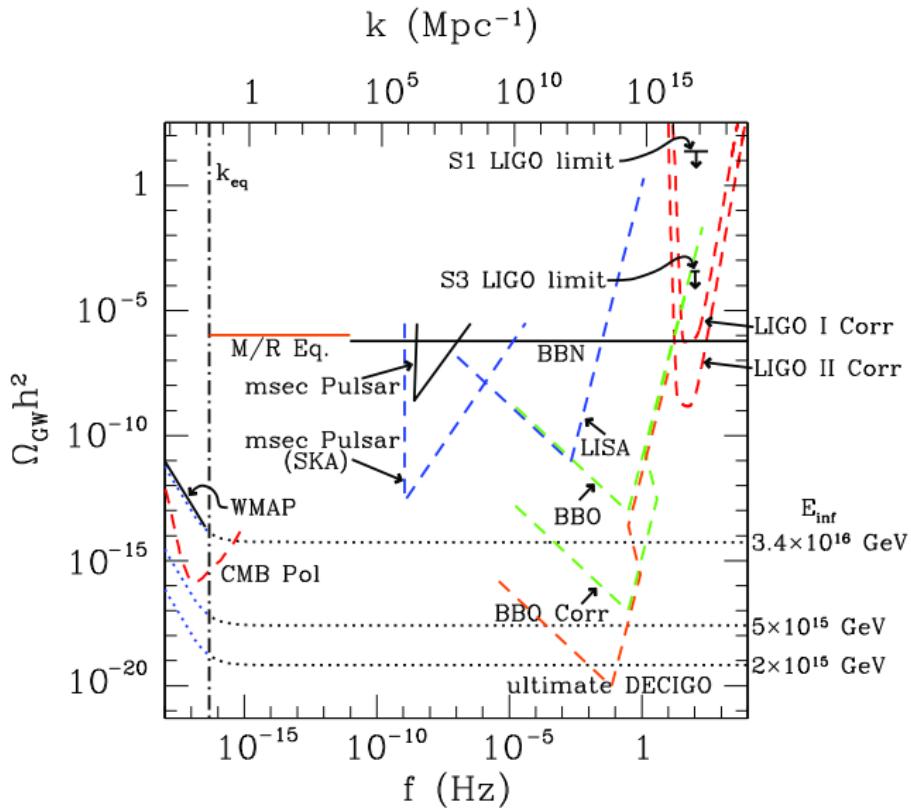
$$m_\phi > H(t) \implies p = \langle E_k \rangle - \langle E_p \rangle = 0, \quad \rho \sim m_\phi^2 \phi^2 \propto 1/a^3$$

- dust-like substance in the late Universe, $\Omega \propto m_\phi^{1/2} \phi_i^2$
at scales $l > 2\pi/m_\phi$ depends on initial conditions
perturbations are suppressed at $l > M_{Pl}^{1/2}/(m_\phi^{1/2} \rho^{1/4})$ fuzzy DM





Prospects in 2014



Conclusion

- What NANOGrav and others observe might be explained by the GW from melting domain walls
- they are expected in models with inverse phase transition
- which may induce light scalar dark matter production
- In realistic models it's mass is of order 10^{-16} - 10^{-12} eV
- That predicts super-radiance instability of rotating black holes with astrophysical masses

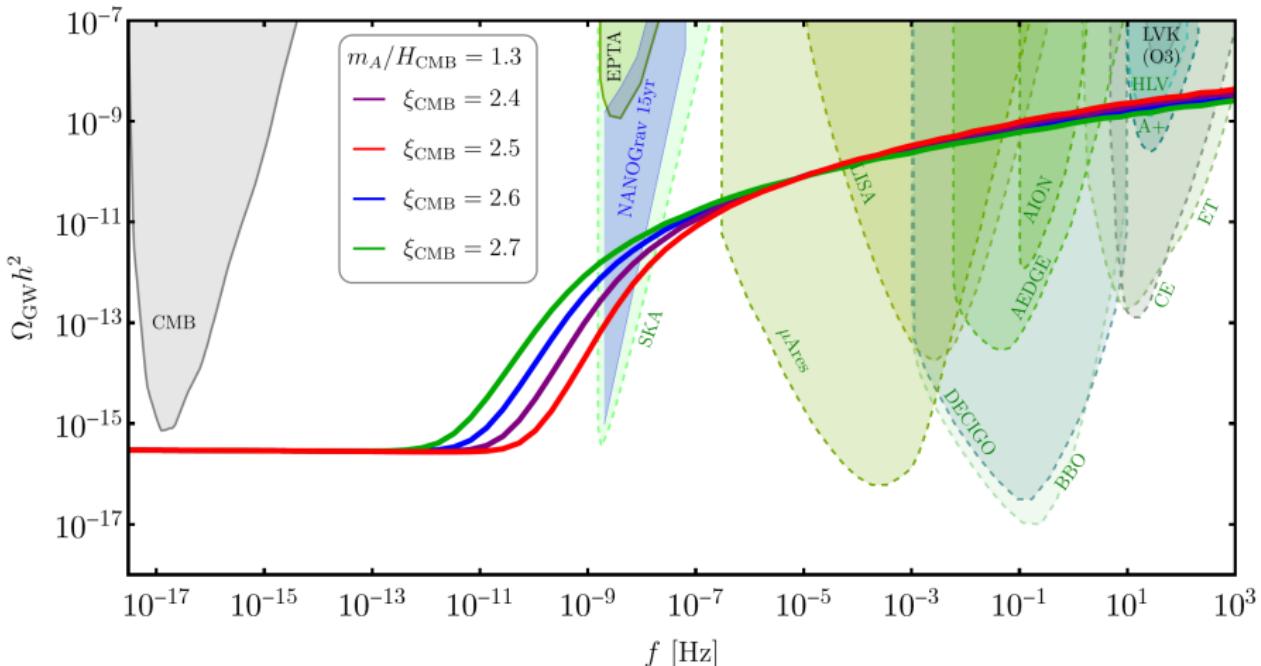
see e.g. 0905.4720

Other possible sources:

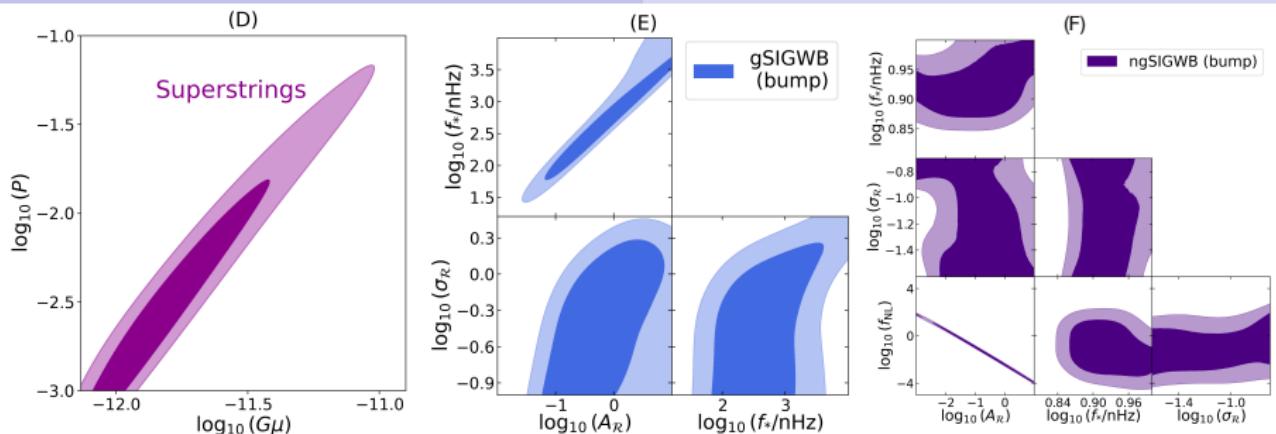
- scalar perturbations (but BH overproduction!)
- cosmic strings
- super heavy BH binaries (does not fit...)
- phase transitions
- violent reheating

$$\Omega_{GW} h^2 = 6.5^{+4.1}_{-2.8} \times 10^{-9} \text{ at } \nu \approx 1 \text{ yr}^{-1}$$

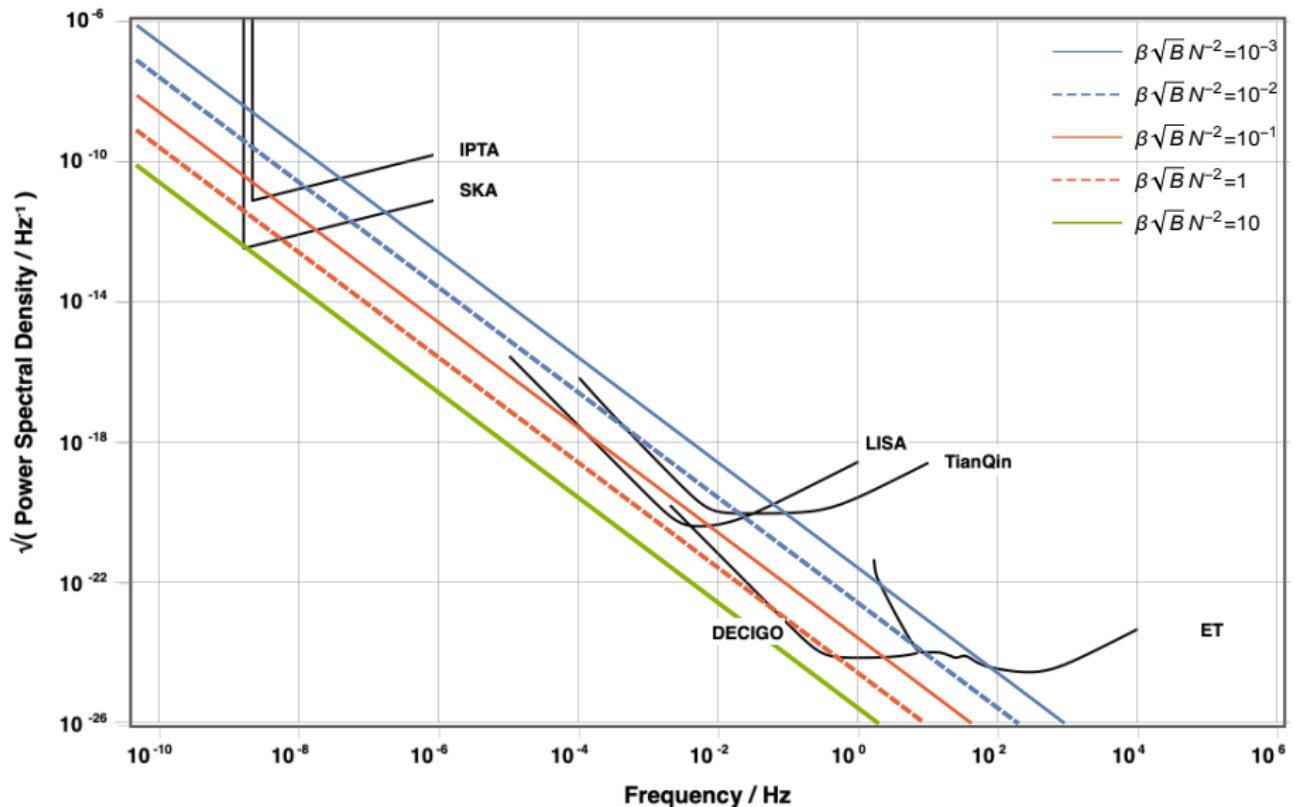
2306.17841



2307.01192



Template	BF_{NANO}	BF_{EPTA}	BF_{comb}	$\Delta \text{AIC}_{\text{comb}}$
PL($n_t = 0$)	0.02	0.62	0.0014	10.24
PL($n_t = 2/3$)	2.6	3.1	1.9	-5.20
PL(n_t, \mathcal{A}_*)	16	4.1	180	-12.25
field th. CS	0.18	1.3	0.019	5.22
super CS	29	3.3	58	-11.11
gSIGWB	110	14	1200	-13.53
ngSIGWB	0.00017	0.001	$5 \cdot 10^{-6}$	4.34
AA	36	2.7	130	-12.36
BPL	17	1.6	120	-0.19
PhTNR	8.5	8.1	150	-12.85
PhTRV	37	13	110	-12.34
MHD	26	3.9	210	-8.77



strain: $\Omega_{GW}H_0^2 \equiv 2\pi^2 f^3 S/3$

2104.13722