

Dark Energy from topology change at the foam level

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- The GB term appears in the effective low energy limit of string theory.
- The GB term is considered a topological invariant in 4D, thus it doesn't contribute to the equations of motion.
- Euclidean Quantum Gravity (EQG) predicts instantons, solutions at the foam level, of distinct topology from the background [Hawking, 1978].

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The main idea

- Take EQG topology changing instatons seriously.
- Calculate the variation of the GB term during topology change.

Topological significance of the GB term

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- The GB curvature polynomial

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$$\mathcal{G} = R^2 - 4R^{\mu\nu}R_{\mu\nu} + R^{\mu\nu\rho\sigma}R_{\mu\nu\rho\sigma}, \quad (1)$$

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- the essence of the theorem is that despite any deformations caused by smooth variations of the metric, the topology of the manifold remains constant, as it is characterized by the topological index χ

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- In Euclidean Quantum gravity, for the path intergral to converge:
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- Thus, the spacetime manifold is Euclideanized and its signature changes from $(-, +, +, +)$ to $(+, +, +, +)$
- then instatons, saddle point solutions, appear with different topology from the background.

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Table: Euler character χ of spacetime manifolds as it has been calculated in [Gibbons and Hawking, 1979]

Spacetime	χ
Minkowski	0
Extreme BHs	0
Self-dual Taub-NUT	1
Schwarchild and Kerr BHs	2
Nariai $S_2 \times S_2$	4
Euclidean Wormhole $S_1 \times S_3$	0

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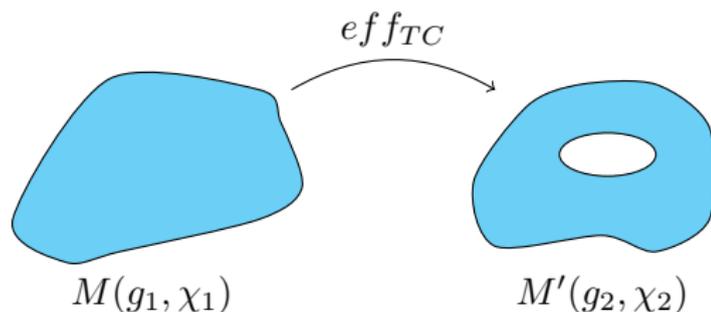
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$$\begin{aligned} M(g_1, \chi_1) &\xrightarrow{\text{eff}_{TC}} M'(g_2, \chi_2), \\ \text{eff}_{TC} : \delta h &\longrightarrow \delta\chi. \end{aligned} \tag{6}$$

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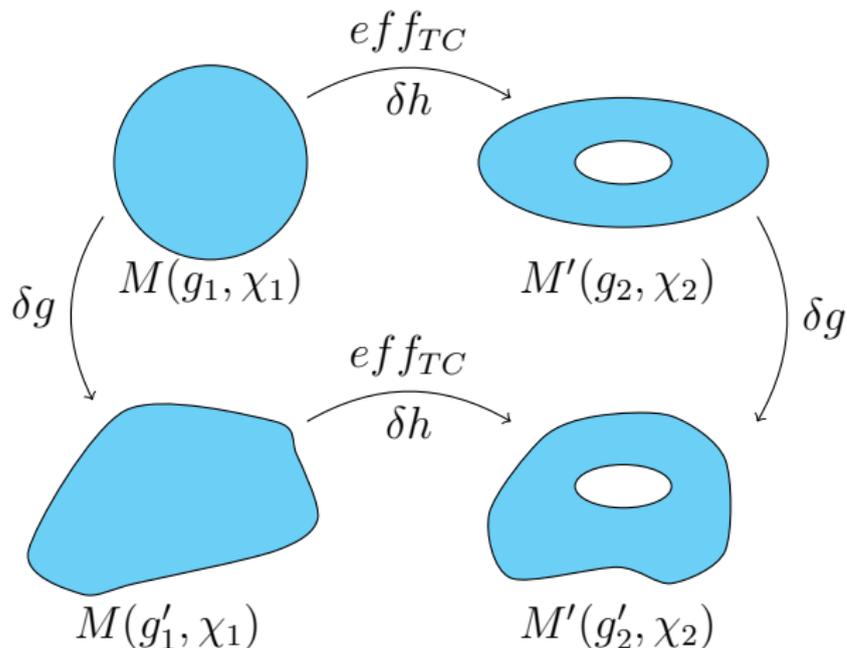
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- the extremization of the effective action by the tadpole condition

Background field approximation I

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The linear split of the metric

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Background independence

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Background independence

- Physical observables must be background independent and the action must be diffeomorphic invariant [Pawlowski and Reichert, 2021]

$$g(\tilde{g}, h) \rightarrow g(\tilde{g} + \delta\tilde{g}, h + \delta h) = g(\tilde{g}, h). \quad (8)$$

Background field approximation II

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- The tadpole condition provides a bridge for energy transfer between the quantum and the classical scale.

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$$S_{EH} = - \int d^4x \sqrt{g} R \quad (11)$$

$$S_{GB} = -\alpha \int d^4x \sqrt{g} (R^2 - 4R^{\mu\nu} R_{\mu\nu} + R^{\mu\nu\rho\sigma} R_{\mu\nu\rho\sigma}), \quad (12)$$

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- and Γ_{1L} corresponds to the gauge fixing and ghost terms

$$\Gamma_{1L} = \Gamma_{GF} + \Gamma_{Fgh}, \quad (13)$$

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$$\begin{aligned} S_0 &= - \int d^4x \sqrt{\tilde{g}} \tilde{R} \\ S_1 &= \int d^4x \sqrt{\tilde{g}} \left(\tilde{R}_{\mu\nu} - \frac{1}{2} \tilde{g}_{\mu\nu} \tilde{R} \right) h^{\mu\nu} \\ S_2 &= - \int d^4x \sqrt{\tilde{g}} \left\{ \frac{1}{4} h^{\mu\nu} \nabla^2 h_{\mu\nu} - \frac{1}{8} h \nabla^2 h \right. \\ &\quad + \frac{1}{2} \left(\nabla^\nu h_{\nu\mu} - \frac{1}{2} \nabla_\mu h \right)^2 + \frac{1}{2} h^{\mu\lambda} h^{\nu\sigma} \tilde{R}_{\mu\lambda\nu\sigma} \\ &\quad \left. + \frac{1}{2} (h^{\mu\lambda} h_{\lambda}^{\nu} - h h^{\mu\nu}) \tilde{R}_{\mu\nu} + \frac{1}{8} (h^2 - 2h^{\mu\nu} h_{\mu\nu}) \tilde{R} \right\}. \end{aligned} \tag{14}$$

Topological variation of the GB action I

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$$\begin{aligned} \delta_h S_{GB} &= \\ \delta_h \left(-\alpha \int_M d^4x \sqrt{g} (R^2 - 4R_{\mu\nu}R^{\mu\nu} + R_{\mu\nu\rho\sigma}R^{\mu\nu\rho\sigma}) \right) \\ &= -32\pi^2 \alpha \frac{\delta\chi}{\delta h^{\mu\nu}} \delta h^{\mu\nu}, \end{aligned} \quad (15)$$

Topological variation of the GB action II

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$$\begin{aligned}\delta_h S_{GB} &= -16\pi^2\alpha \frac{\partial\chi}{\partial V} \frac{\delta V}{\delta h^{\mu\nu}} \delta h^{\mu\nu} = -16\pi^2\alpha \frac{\partial\chi}{\partial V} \delta_h \left(\int_M d^4x \sqrt{g} \right) \\ &= -16\pi^2\alpha \frac{\partial\chi}{\partial V} \int_M d^4x \frac{\delta\sqrt{g}}{\delta h^{\mu\nu}} \delta h^{\mu\nu}.\end{aligned}\tag{16}$$

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$$\begin{aligned}\frac{\delta\sqrt{g}}{\delta h^{\mu\nu}} &= \frac{\delta}{\delta h^{\mu\nu}} \left(\sqrt{\tilde{g}} + \frac{1}{2}\sqrt{\tilde{g}}\tilde{g}_{\mu\nu}h^{\mu\nu} - \frac{1}{4}\sqrt{\tilde{g}}h^{\mu\nu}h_{\mu\nu} \right. \\ &\quad \left. + \frac{1}{8}\sqrt{\tilde{g}}(h_{\mu}^{\mu})^2 + \mathcal{O}(h^3) \right) \\ &= 0 + \frac{1}{2}\sqrt{\tilde{g}}\tilde{g}_{\mu\nu} + \mathcal{O}(h),\end{aligned}\tag{17}$$

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- Equivalently, one can express the topological variation of the Gauss-Bonnet term as

$$\frac{1}{\sqrt{\tilde{g}}} \frac{\delta S_{GB}}{\delta h^{\mu\nu}} = -16\pi^2 \alpha \frac{\partial \chi}{\partial V} \tilde{g}_{\mu\nu} + \mathcal{O}(h). \quad (20)$$

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$$\tilde{R}_{\mu\nu} - \frac{1}{2} \tilde{g}_{\mu\nu} \tilde{R} + \tilde{g}_{\mu\nu} \Lambda_{eff} = 8\pi G T_{\mu\nu}, \quad (22)$$

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- where the stress tensor originates from the one loop part of the effective action, containing the matter like correction terms in the right side of the equation.

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$$T^{\mu\nu} = -\frac{2}{\sqrt{g}}\frac{\delta}{\delta h^{\mu\nu}}\Gamma_{1L}|_{h=0}, \quad (23)$$

- and the dynamical topological term appears in the equation of motion as an effective cosmological constant

$$\Lambda_{eff} = -16\pi^2\alpha\frac{\partial\chi}{\partial V}. \quad (24)$$

Λ_{eff} interpretation

- Since $\delta\chi = 2$ corresponds to the formation of a Nariai instanton and $\delta\chi = -2$ to the formation of an Euclidean Wormhole, the term $\frac{\partial\chi}{\partial V}$ can be interpreted as the density of EW's or instantons per four-volume $\rho_w = \frac{N_w}{V}$

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- The upper bound for the density will be for one wormhole per Planck volume, that is $\rho_{Mw} = \frac{1}{l_p^4} \sim 10^{140}$

Bounds of Λ_{eff}

- There is a consensus over the various string models that the α constant of the GB term should be $\alpha = l_p^2$.
- If one inserts the observed present value of the cosmological constant $\Lambda_{obs} = 10^{-52} m^{-2}$ in (s.i) units to Eq. (25), then the model predicts $\rho_w = 10^{16}$ wormholes per cubic meter per second, which seems natural.
- The upper bound for the density will be for one wormhole per Planck volume, that is $\rho_{Mw} = \frac{1}{l_p^4} \sim 10^{140}$
- Thus, the range of the theoretical value of Λ_{eff} can be $0 \leq \Lambda_{eff} \leq \Lambda_M$ and the ratio $\frac{\Lambda_M}{\Lambda_{obs}} = 10^{124}$ of the model spans the 120 orders of magnitude, of the known discrepancy.

Sign of Λ_{eff}

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Since the sign of Λ_{eff} is determined by the type of instanton that is created, positive for Nariai and negative for EW's one can develop a spectrum of models mixing the two basic scenarios:

- 1 the creation of EWs provides a positive repulsive cosmological constant of topological origin, that can be the unique source of cosmic acceleration
- 2 the creation of Nariai instantons corresponds to a negative component of the net effective cosmological constant, absorbing up to 120 orders of vacuum energy, like a sponge absorbing water [Padmanabhan, 2003], thus providing a mechanism for the ending of the inflation period.

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- Ignatios Antoniadis and N. C. Tsamis. On the Cosmological Constant Problem. *Phys. Lett. B*, 144:55–60, 1984. doi: 10.1016/0370-2693(84)90175-8.
- Maximilian Becker and Martin Reuter. Background independent field quantization with sequences of gravity-coupled approximants. II. Metric fluctuations. *Phys. Rev. D*, 104(12):125008, 2021. doi: 10.1103/PhysRevD.104.125008.
- G. W. Gibbons. Topology change in classical and quantum gravity. 10 2011.
- G. W. Gibbons and S. W. Hawking. Classification of Gravitational Instanton Symmetries. *Commun. Math. Phys.*, 66:291–310, 1979. doi: 10.1007/BF01197189.

Bibliography II

- S.W. Hawking. Spacetime foam. *Nuclear Physics B*, 144(2): 349–362, 1978. ISSN 0550-3213. doi: [https://doi.org/10.1016/0550-3213\(78\)90375-9](https://doi.org/10.1016/0550-3213(78)90375-9). URL <https://www.sciencedirect.com/science/article/pii/0550321378903759>.
- T. Padmanabhan. Cosmological constant: The Weight of the vacuum. *Phys. Rept.*, 380:235–320, 2003. doi: 10.1016/S0370-1573(03)00120-0.
- Jan M. Pawłowski and Manuel Reichert. Quantum Gravity: A Fluctuating Point of View. *Front. in Phys.*, 8:551848, 2021. doi: 10.3389/fphy.2020.551848.
- Gerard 't Hooft and M. J. G. Veltman. One loop divergencies in the theory of gravitation. *Ann. Inst. H. Poincaré Phys. Theor. A*, 20:69–94, 1974.