Exploring mixed lepton-quark interactions in non-resonant leptoquark production at the LHC

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Work done in collaboration with: Antonio P. Morais, Roman Pasechnik, António Onofre and Werner Porod

Based on arXiv:2206.01674 and arXiv:2306.15460;

https://github.com/Mrazi09/LQ_collider_project

https://github.com/Mrazi09/Leptoquark-project

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Despite its successes, the SM either **fails** or is **insufficient** in tackling various observations, with **anomalous results** popping up in recent years ...

- <u>LFV</u>: Various anomalous results in *B* meson decays $(B_s \to \phi \mu^+ \mu^-, B \to K \mu^+ \mu^-, P'_5)$;
- a_μ: Tension with the SM [Phys.Rev.Lett. 126, 141801 (2021)]. Although lattice results indicate consistency [Nature 593 (2021) 7857, 51-55].
- M_W : Tension primarily driven by the new CDF result [Science 376 (2022) 170–176].
- $U_{PMNS} + m_{\nu}$ SM can not account for neutrino physics. Requires non-renormalisable 5D Weinberg operator [Phys. Lett. B 91 (1980) 51];

As shown in our previous work [2206.01674 [hep-ph]], a minimal LQ model is sufficient to address all of these open questions.

The model:

- One SU(2) singlet: $S \sim (\bar{3}, 1, 1/3)^{-}$;
- One SU(2) doublet: $R \sim (\bar{3}, 2, 1/6)^-$;
- Yukawas that generate proton decay $\rightarrow y_1 \bar{Q}^c Q S^\dagger + y_2 \bar{d} u S + {\rm h.c.}$

$$\mathbb{Z}_2$$
 symmetry: $\mathbb{P}_B = (-1)^{3B+2S} \implies S^-, R^-, q^+_{L,R}, \ell^-_{L,R}, H^-$

Group charges and particles are emergent in a gauge-flavour unification model based on $E_6\times SU(3)_F$ [Eur. Phys. J. C 80, 1162 (2020)].

- Yukawa Lagrangian: $\mathcal{L}_Y = \Theta_{ij} \bar{Q}_j^c L_i S + \Omega_{ij} \bar{L}_i d_i R^{\dagger} + \Upsilon_{ij} \bar{u}_j^c e_i S + h.c.$
- Relevant scalar potential:

$$\begin{split} V \supset & -\mu^2 |H|^2 + \mu_S^2 |S|^2 + \mu_R^2 |R|^2 + \lambda (H^{\dagger}H)^2 + g_{HR}(H^{\dagger}H)(R^{\dagger}R) + \\ & g'_{HR}(H^{\dagger}R)(R^{\dagger}H) + g_{HS}(H^{\dagger}H)(S^{\dagger}S) + \left(a_1 R S H^{\dagger} + \text{h.c.}\right) \,. \end{split}$$

• After EWSB, two physical (1/3)e appear in the spectrum

$$M_{LQ^{1/3}}^{2} = \begin{bmatrix} \mu_{S}^{2} + \frac{g_{HS}v^{2}}{2} & \frac{va_{1}}{\sqrt{2}} \\ \frac{va_{1}}{\sqrt{2}} & \mu_{R}^{2} + \frac{(g_{HR} + g'_{HR})v^{2}}{2} \end{bmatrix}$$

Trade physical masses, $m_{S_1^{1/3}}$ and $m_{S_2^{1/3}}$, for the μ terms

$$\mu_{S}^{2} = \frac{1}{2} (m_{S_{1}^{1/3}}^{2} + m_{S_{2}^{1/3}}^{2} - g_{HS}v^{2} + \sqrt{(m_{S_{1}^{1/3}}^{2} - m_{S_{2}^{1/3}}^{2})^{2} - 2a_{1}^{2}v^{2}}),$$

$$\mu_{R}^{2} = \frac{1}{2} (m_{S_{1}^{1/3}}^{2} + m_{S_{2}^{1/3}}^{2} - (g_{HR} + g'_{HR})v^{2} - \sqrt{(m_{S_{1}^{1/3}}^{2} - m_{S_{2}^{1/3}}^{2})^{2} - 2a_{1}^{2}v^{2}}).$$

Mixing controlled by the relation $\sin(2\theta)=(\sqrt{2}va_1)(m_{S_1^{1/3}}^2-m_{S_2^{1/3}}^2)^{-1}$

• A single (2/3)e state appears and does not mix with $m_{S^{2/3}}^2=\mu_R^2+\frac{g_{HR}v^2}{2}$.



- Limits based on the combination of various channels (pair prod., single prod., recast) [JHEP 01 (2019) 132];
- Most searches tend towards diagonal channels. Current constraints are not strong for off-diagonal contributions. Particularly for 1st and 3rd generation couplings; **Substantial phase-space left unprobed.**
- $\bullet\,$ We focus our analysis for the μc coupling and study its implications at future runs of the LHC.



- Vast region of the parameter space in reach of both run-III and HL-LHC;
- Cross-section driven primarily by $S_1^{1/3}$ due to small mixing $a_1 \sim \mathcal{O}(1 \text{ GeV}) \Rightarrow \sin(2\theta) \ll 1$;
- Various benchmarks are defined for different mass scales of the lightest (1/3e) LQ.



- Drell-Yan *t*-channel type process with two muons in the final state;
- LQ mass cannot be reconstructed. Although, kinematic distributions affected by it.
- Events generated at LO in MadGraph, with Delphes detector simulation.
 Backgrounds reweighted based on higher-order corrections;
- Selection criteria
 - $p_T(\mu^{\pm}) > 25 \text{ GeV}$
 - $|\eta(\mu^{\pm})| < 2.5$



- Dimension-full distributions offer the greater discriminating power. Angular distributions closely follow SM backgrounds;
- The signal tends to populate the high-energy regions, as opposed to the SM expectation;

M_{LQ} (GeV)	$M(\mu^+, \mu^-)$ $\mathcal{L} = 300 \text{ fb}^{-1}$	$M(\mu^+,\mu^-)$ $\mathcal{L} = 3000 \text{ fb}^{-1}$	$E(\mu^+)$ $\mathcal{L} = 300 \text{ fb}^{-1}$	$\begin{array}{c} E(\mu^+) \\ \mathcal{L} = 3000 \text{ fb}^{-1} \end{array}$	Combined $(300, 3000) \text{ fb}^{-1}$
1.5 TeV $(\mathcal{B}1)$	1.75σ	5.20σ	0.891σ	2.72σ	$(3.06\sigma, 9.72\sigma)$
$2.5 \text{ TeV} (\mathcal{B}2)$	0.573σ	2.28σ	0.744σ	2.33σ	$(1.15\sigma, 4.97\sigma)$
$3.5 \text{ TeV} (\mathcal{B}3)$	0.128σ	0.912σ	0.225σ	1.04σ	(0.288 <i>σ</i> , 1.97 <i>σ</i>)



- The combination of the various distributions can extend the discovery threshold of our model's LQs;
- $M/\Upsilon < 4.17 \text{ TeV}$ @ 95% CL for run-III;
- $M/\Upsilon < 6.55$ TeV @ 95% 2 CL for HL-LHC.

To conclude ...

- Discussed a simple economical model and how it solves current open questions;
- Showed that there are regions of parameter space that can be probed at LHC (run-III and/or HL-LHC)



• Flavour off-diagonal channels also possible. LQ Yukawa couplings with top more easily measured in lepton colliders.

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Thank you for your attention

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