

Ensemble Averages of Narain CFTs & Holography

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Jacob M. Leedom

Corfu Holography & Swampland, 09.09.2022



CLUSTER OF EXCELLENCE
QUANTUM UNIVERSE

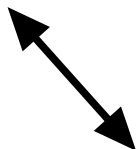
Holography

Swampland

Overview

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Quantum Gravity

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Holography

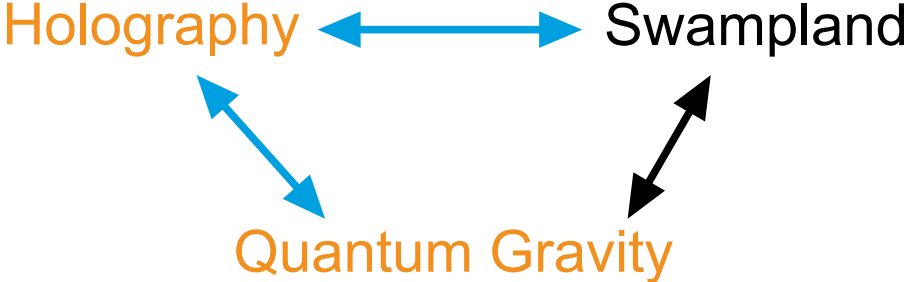
Swampland



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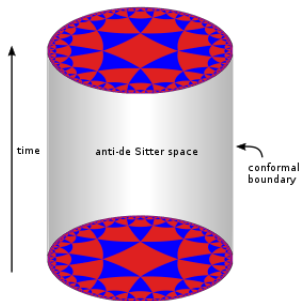
Overview

- > Peculiar Duality between AdS_3 & CFT_2



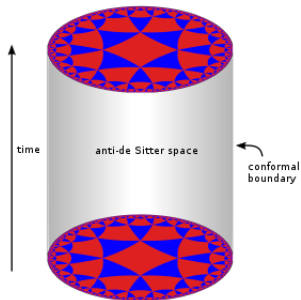
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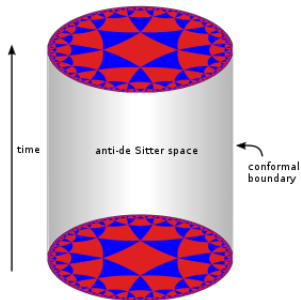


> Conventional Duality:

$$Z_{\text{Bulk}}[\phi(\text{boundary}) = J] = Z_{\text{CFT}}[J]$$

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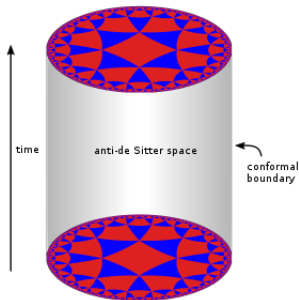


> Ensemble Duality:

$$\sum_{\text{AdS}_3 \text{ geometries}} Z_{\text{Bulk}}[\tau] = \int_{\text{moduli space}} [dm] Z_{\text{CFT}}[m; \tau]$$

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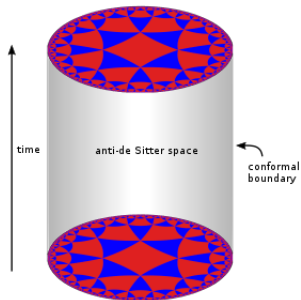
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You condense it with locusts and tape:
Still keeping one principal object in view—
To preserve its symmetrical shape."
—*Hunting of the Snark*, Lewis Carroll

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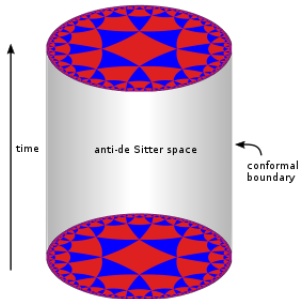
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$$\tau \rightarrow \tau_M = \frac{a\tau + b}{c\tau + d}$$
$$\begin{pmatrix} a & b \\ c & d \end{pmatrix} \in \text{SL}_2(\mathbb{Z})$$

AdS₃/CFT₂:



AdS₃/CFT₂: Ensemble of Even, Self-Dual CFTs



AdS₃/CFT₂: Ensemble of Even, Self-Dual CFTs

> Toy Illustration: 1 Compact Boson $\Rightarrow (c, \tilde{c}) = (1, 1)$

[2006.04839,2006.04855]

$$S_{CFT} = \frac{R^2}{2\pi} \int d^2\sigma \partial_a X \partial^a X$$

$$X \sim X + 2\pi$$

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$$\tau = \tau_1 + i\tau_2$$



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$$\text{T-Duality: } R \leftrightarrow 1/R$$

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$$ds^2 = G^{\alpha\beta} G^{\rho\sigma} (dG_{\alpha\rho} dG_{\beta\sigma} + dB_{\alpha\rho} dB_{\beta\sigma})$$



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- > Ensemble Average: Computable by Siegel-Weil Formula

$$\frac{1}{\text{Vol}(\mathcal{M}_{p,p})} \int_{\mathcal{M}_{(p,p)}} dm Z_{CFT}[m, \tau] = \frac{E_{p/2}(\tau)}{|\eta(\tau)|^{2p}}$$

$$E_s(\tau) = \sum_{(c,d)=1} \frac{1}{|c\tau + d|^{2s}}$$

AdS₃/CFT₂: U(1)^{2p} Chern-Simons & a Sum over Geometries



AdS₃/CFT₂: $U(1)^{2p}$ Chern-Simons & a Sum over Geometries

- > Boundary has $U(1)^{2p}$ current algebra & $U(1)^{2p}$ global symmetry
⇒ Bulk has $U(1)^{2p}$ gauge symmetry

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- > The level matrix of Chern-Simons is quantized, so a natural guess is:

$$S_{bulk} \supset \sum_{M,N} \frac{Q_{MN}}{2\pi} \int A^M \wedge dA^N$$



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Boundary current algebra (Brown-Henneaux-like photon modes) give:

$$Z_{ThAdS} = \frac{1}{|\eta(\tau)|^{2p}}$$

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- > Idea: what about other fillings?

AdS₃/CFT₂: U(1)^{2p} Chern-Simons & a Sum over Geometries

- > Described by so-called PSL₂(\mathbb{Z}) black holes M_(c,d)
 - solid tori defined by filling different cycles of boundary torus (genus 1 handlebodies)

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$$T^n = \begin{pmatrix} 1 & n \\ 0 & 1 \end{pmatrix}$$

- Thus geometries labeled by elements of $\Gamma_\infty \backslash \text{PSL}_2(\mathbb{Z}) \Rightarrow$ coprime integers (c, d)
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- > Sum over geometries \Rightarrow sum over images under $\Gamma_\infty \backslash \text{PSL}_2(\mathbb{Z})$:

$$\begin{aligned} Z(\tau) &= \sum_{g \in \Gamma_\infty \backslash \text{PSL}_2(\mathbb{Z})} \frac{1}{|\eta(g \cdot \tau)|^{2p}} \\ &= \frac{E_{p/2}(\tau)}{|\eta(\tau)|^{2p}} \end{aligned}$$

AdS₃/CFT₂: Duality Established

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- > Moduli Space M_(p,p)

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- > Moduli Space $\mathcal{M}_{(p,p)}$

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- > Average over moduli space:

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$$S_{CFT} = \frac{1}{2\pi} \int d^2\sigma \left(G_{\alpha\beta} \partial_a X^\alpha \partial^a X^\beta + i B_{\alpha\beta} \epsilon^{ab} \partial_a X^\alpha \partial_b X^\beta \right)$$

> Moduli Space $\mathcal{M}_{(p,p)}$

$$\text{O}(p, p; \mathbb{Z}) \backslash \text{O}(p, p; \mathbb{R}) / \text{O}(p) \times \text{O}(p)$$

> Average over moduli space:

AdS₃/CFT₂: Duality Established

- > U(1)^{2p} Chern-Simons Theory

$$S_{Bulk} = \sum_{MN} \frac{Q_{MN}}{2\pi} \int A^M \wedge dA^N$$

- > Geometries

M_(c,d) : PSL₂(Z) black holes

$$(c, d) \leftrightarrow \Gamma_\infty \backslash \text{PSL}_2(\mathbb{Z})$$

- > Sum Over Geometries:

$$\sum_{\text{AdS}_3 \text{ geometries}} Z_{\text{Bulk}}[\tau] = \int_{\text{moduli space}} [dm] Z_{\text{CFT}}[m; \tau]$$

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Number Theory Detour: The Siegel-Weil Formula



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$$\Theta_{Q,h}(m; \tau) = \sum_{\ell \in \Lambda} e^{i\pi\tau_1 Q[\ell+h] - \pi\tau_2 H[m; \ell+h]}$$

$$Q[\ell] \in 2\mathbb{Z}$$

$$H : H^T Q H = Q$$

$$h \in \mathcal{D} = \Lambda^* / \Lambda$$

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$$\sigma = p - q$$

$$\Theta_{Q,h}(m; \tau_M) = \frac{e^{-i\pi\sigma/4} c^{-\frac{p+q}{2}}}{|Q|^{1/2}} (c\tau + d)^{p/2} (c\bar{\tau} + d)^{q/2} \sum_{h' \in \mathcal{D}} \lambda_{hh'} \Theta_{Q,h'}(\tau)$$

$$\lambda_{hh'} = \sum_{g \in \Lambda/c\Lambda} \exp\left(\frac{i\pi}{\gamma} (aQ[g+h] - 2Q[h', g+h] + dQ[h'])\right)$$

$$\tau_M = \frac{a\tau + b}{c\tau + d} \quad \& \quad \begin{pmatrix} a & b \\ c & d \end{pmatrix} \in \mathrm{SL}_2(\mathbb{Z})$$

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$$\langle \Theta_{Q,h}(m; \tau) \rangle = \frac{1}{\text{Vol}(\mathcal{M}_Q)} \int_{\mathcal{M}_Q} [dm] \Theta_{Q,h}(m; \tau)$$

$$= \delta_{h \in \Lambda} + \sum_{\substack{(c,d)=1 \\ c>0}} \frac{\gamma_{Q,h}(c, d)}{(c\tau + d)^{\frac{p}{2}} (c\bar{\tau} + d)^{\frac{q}{2}}} := E_{Q,h}(\tau)$$

$$\gamma_{Q,h} := e^{i\pi\sigma/4} |Q|^{-\frac{1}{2}} c^{-\frac{p+q}{2}} \sum_{g \in \Lambda/c\Lambda} \exp\left(-i\pi \frac{d}{c} Q[g+h]\right)$$

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$$\gamma_{Q,h} = c^{-p} \sum_{n_i, w^i=0}^{c-1} \exp\left(-i\pi \frac{d}{c} n_i w^i\right) = 1$$

if $q = p$ and $|Q| = 1$

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$$E_{Q,h}(\tau) - \langle \Theta_{Q,h}(m; \tau) \rangle = 0 \text{ at cusps}$$

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$$\square_k = -\tau_2^2 (\partial_1^2 + \partial_2^2) + ik\tau_2 \partial_1$$

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$$\text{Minimum eigenvalue (square normalizable)} : \lambda_{min,k} = \frac{|k|}{2} \left(1 - \frac{|k|}{2} \right)$$

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no normalizable eigenfunction satisfies equation



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$$f_{Q,h}(\tau) = \tau_2^{(p+q)/4} (E_{Q,h}(\tau) - \langle \Theta_{Q,h}(m; \tau) \rangle) \Rightarrow \lambda < \lambda_{min}$$



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AdS₃/CFT₂: Even Q CFTs

- > Consider a CFT with $(c, \tilde{c}) = (p, q)$
- > Vertex Operators:

$$V_{k_L, k_R} = \exp(ik_L \cdot X_L(z) + ik_R \cdot X_R(\bar{z}))$$

- > To close OPE, need to define theory by picking out vertex operators that fill out a lattice Λ
- > For vertex operators to have integer spin, want an even quadratic form $Q[k] \in 2\mathbb{Z}$ of signature (p, q)
- > Also have Hamiltonian H that need not have integer values on Λ and depends on moduli
- > Moduli Space:

$$\mathcal{M}_Q = O(p, q, \mathbb{Z}) \backslash O(p, q, \mathbb{R}) / O(p, \mathbb{R}) \times O(q, \mathbb{R})$$

AdS₃/CFT₂: Even Q CFTs

- > The partition function is built from the Siegel-Narain theta functions:

$$Z_{Q,0}^{CFT}(m; \tau) = \frac{\Theta_Q(m; \tau)}{\eta^p(\tau)\bar{\eta}^q(\bar{\tau})}$$

- > this is not generally modular invariant and one needs to consider $Z_{Q,h}^{CFT}(m; \tau)$
- > Ensemble averages are direct application of Siegel-Weil Formula:

$$\langle Z_{Q,h}^{CFT}(m; \tau) \rangle = \frac{E_{Q,h}(\tau)}{\eta^p(\tau)\bar{\eta}^q(\bar{\tau})}$$

AdS₃/CFT₂: U(1)^{p+q} Chern-Simons Theory

> Averaged Partition functions can be written as

$$\langle Z_{Q,h}^{CFT}(m; \tau) \rangle = \sum_{g \in \Gamma_\infty \backslash \text{PSL}_2(\mathbb{Z})} e^{\frac{i\pi\sigma}{12} \Phi(g) - \frac{i\pi\sigma}{4}} \frac{\gamma_{Q,h}(c, d)}{\eta^p(g \cdot \tau) \bar{\eta}^q(g \cdot \bar{\tau})}$$

$\Phi(g)$ = Rademacher Phi

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$$\gamma_{Q,h}(c, d) \sim \lambda_{h,0}$$

Want elements of a transformation matrix in the bulk

AdS₃/CFT₂: U(1)^{p+q} Chern-Simons Theory

- > Canonical Quantization: CS States for general even Q labeled by elements of \mathcal{D} :

$$T|h\rangle = e^{i\pi Q[h]} e^{-\frac{i\pi\sigma}{12}} |h\rangle$$
$$S|h\rangle = \frac{1}{\sqrt{|Q|}} \sum_{h' \in \mathcal{D}} e^{-2\pi i Q[h, h']} |h'\rangle$$

- > The bulk analog of $\lambda_{h,0}$ is something like $\langle 0|U(g)|h\rangle$. In fact we have

$$\langle 0|U(g)|h\rangle^* = \langle h|U(g)^{-1}|0\rangle = e^{\frac{i\pi\Phi(g)}{12} - \frac{i\pi\sigma}{4}} \gamma_{Q,h}(c, d)$$

Partition functions of CS on Lens spaces! [\(see paper by Jeffery\)](#)



AdS₃/CFT₂: Duality for Even Q

$$\sum_{g \in \Gamma_\infty \backslash \mathrm{PSL}_2(\mathbb{Z})} \frac{\langle h | U(g)^{-1} | 0 \rangle}{\eta^p(g \cdot \tau) \bar{\eta}^q(g \cdot \bar{\tau})} = \frac{1}{\mathrm{Vol}(\mathcal{M}_Q)} \int_{\mathcal{M}_Q} \frac{\Theta_{Q,h}(m; \tau)}{\eta^p(\tau) \bar{\eta}^q(\bar{\tau})}$$

Global Symmetries, Orbifolds, & Ensembles

- > The Bulk theories have **global symmetries**, i.e. \mathbb{Z}_2 symmetry $A \rightarrow -A$



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- > One can gauge these symmetries to produce orbifolded dualities. Even, self dual case considered in [2103.15826,2105.12594]
- > General Q theories lead to novel Eisenstein series. The “gauge” sector partition function has generalized lattice theta functions :

$$\vartheta_{(\delta,\eta)}^I(\tau, \bar{\tau}; m) := \sum_{\ell \in I+\delta} e^{i\pi\tau Q_L(\ell) - i\pi\bar{\tau} Q_R(\ell)} e^{2i\pi Q_I(\ell, \eta)}$$

that average to “twisted” non-holomorphic Eisenstein series:

$$E_{(\delta,\eta)}^I(\tau, \bar{\tau}) = \delta_{\delta \in I} + \frac{1}{\sqrt{\det Q_I}} e^{-\frac{\pi i}{4}(p_I - q_I)} \sum_{(c,d)=1, c>0} c^{-\frac{p_I+q_I}{2}} (c\tau + d)^{-\frac{p_I}{2}} (c\bar{\tau} + d)^{-\frac{q_I}{2}} \mu_{(\delta,\eta)} \cdot M^{-1} \lambda_{0, -d\delta + c\eta}^I(M^{-1})$$



Global Symmetries, Orbifolds, & Ensembles

- > The Bulk theories have **global symmetries**, i.e. \mathbb{Z}_2 symmetry $A \rightarrow -A$
- > One can gauge these symmetries to produce orbifolded dualities. Even, self dual case considered in [2103.15826,2105.12594]
- > General Q theories lead to novel Eisenstein series. The “gauge” sector partition function has generalized lattice theta functions :

$$\vartheta_{(\delta,\eta)}^I(\tau, \bar{\tau}; m) := \sum_{\ell \in I+\delta} e^{i\pi\tau Q_L(\ell) - i\pi\bar{\tau} Q_R(\ell)} e^{2i\pi Q_I(\ell, \eta)}$$

that average to “twisted” non-holomorphic Eisenstein series:

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- > Global symmetries not in contradiction with [1810.05338+] - this is not Einstein gravity
- > However, this *is* holographic - lesson for quantum gravity?



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- > Must use *Maxwell-Chern-Simons* theory [\[arXiv:0403225\]](#)

$$S_{MCS} = \frac{1}{16\pi^2} \sum_{i,j} \int_M \left(-\frac{1}{2e^2} \lambda_{ij}^{-1} dA^i \wedge \star dA^j + 2\pi i Q_{ij} A^i \wedge dA^j \right)$$

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- > Connection to string theory: MCS arises in $AdS_3 \times K_7$ compactifications, where $K_7 = S^3 \times (S^1)^4$ or $S^3 \times S^3 \times S^1$



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- If $P(x)$ is a spherical polynomial with respect to Q , then $P(\partial X)$ is a primary operator and its one-point function & ensemble average are

$$\vartheta_{Q,P} = \sum_{\ell \in \Lambda} P(\ell) e^{iQ(\ell)\tau}, \quad \langle\langle \vartheta_{Q,P_m^\nu}(\tau) \rangle\rangle = C_k^{(\nu)} |T_m$$

- > Can also define average of orbifold twist operator correlation functions [\[2103.15826\]](#)



Things Left Unsaid

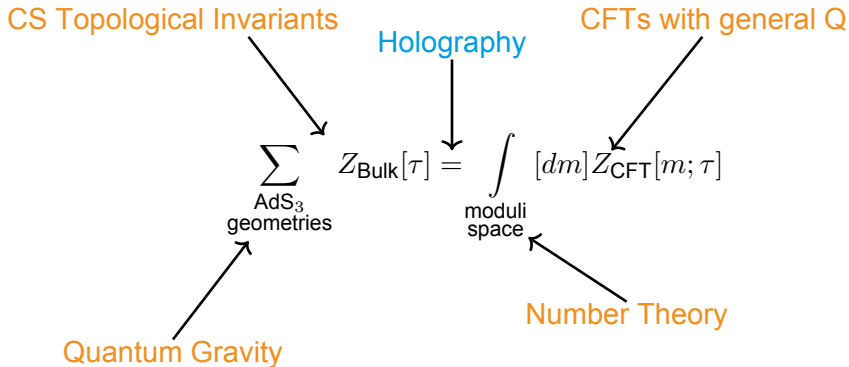
> Odd lattice CFTs & Spin Chern Simons

$$\Theta_{Q,h}^{\epsilon_1, \epsilon_2}(m; \tau) = \sum_{\ell \in \Lambda + h + \epsilon_1 W/2} e^{i\pi\tau Q_L[\ell] - i\pi\bar{\tau} Q_R[\ell]} (-1)^{\epsilon_2(W, \ell)}$$

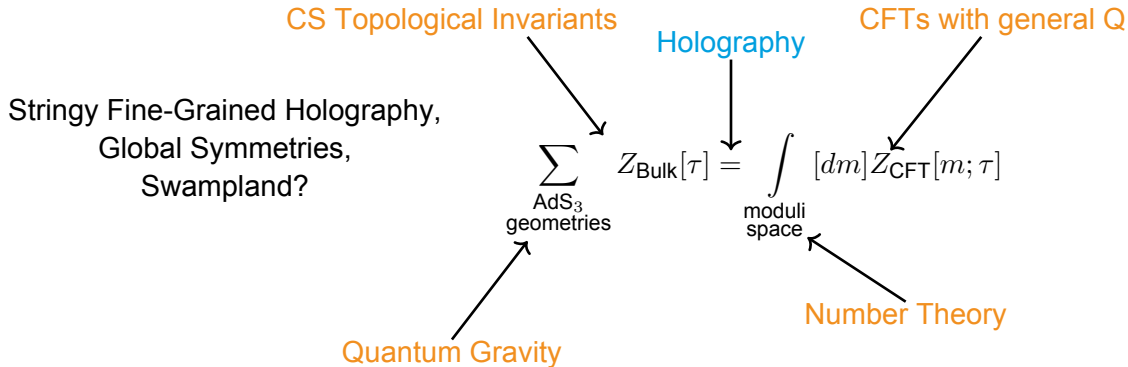
> Higher genus

$$E_{Q, \vec{h}}^g(\Omega) = \sum_{\gamma \in \Gamma_\infty \backslash \mathrm{Sp}(2g, \mathbb{Z})} \frac{\gamma_{\vec{h}}(C, D)}{\det(C\Omega + D)^{p/2} \det(C\bar{\Omega} + D)^{q/2}}$$

Conclusion



Conclusion




Thank you!

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