

LINEAR e^+e^- COLLIDERS – FUTURE HIGGS FACTORIES



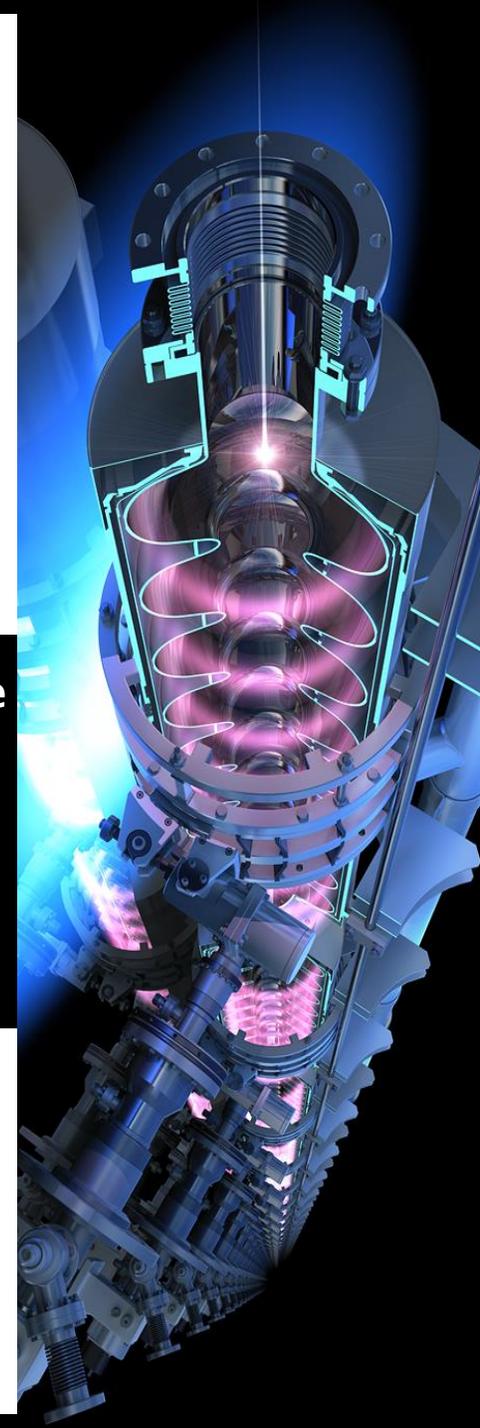
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Corfu Summer Institute

Workshop on the
Standard Model and
Beyond

29.08.-08.09. 2021

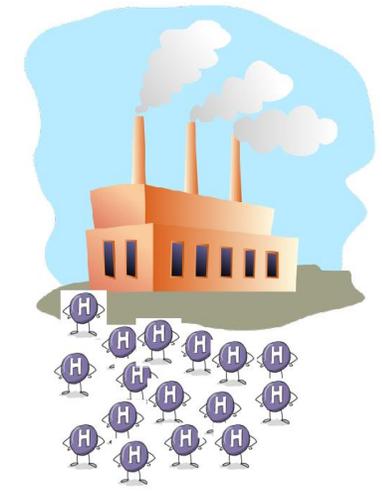


OVERVIEW

- WHY HIGGS FACTORIES?
- LINEAR COLLIDERS
- DETECTOR CONCEPTS & TECHNOLOGIES
- HIGGS PHYSICS AT LINEAR COLLIDERS
- OUTLOOK

WHY HIGGS FACTORIES?

- Higgs discovery ended era of reductionism and symmetries in particle physics [1]
- Never seen before fundamental scalar is discovered, unique (with exception of gravity) in its self-coupling
- Higgs discovery opened several important questions of nature of relativistic vacuum:
 - How can we accommodate it in energy density of the Universe?
 - Why the Higgs is not enormously massive (even Planckian)?



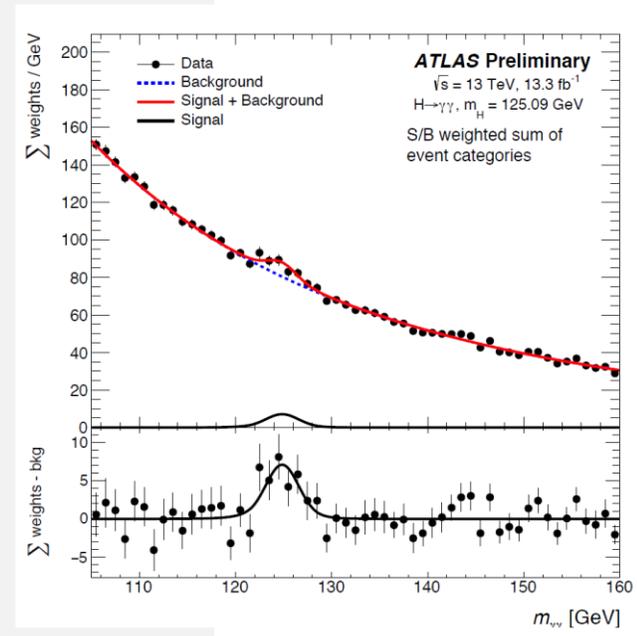
$$\left(\frac{E}{V} \right) = \int d^3k \frac{1}{2} \sqrt{k^2 + m_B^2} - \frac{1}{2} \sqrt{k^2 + m_F^2}$$

\downarrow $g^4 |k|^2$ \downarrow $\lambda^2 |k|^2$

$$= \Lambda_{UV}^4 + (g^2 - \lambda^2) \Lambda_{UV}^2 |k|^2 + \dots$$

\uparrow Cosmological Constant Problem \uparrow Hierarchy Problem

[1] WHY IS THERE A MACROSCOPIC UNIVERSE?



LIMITS, LIMITS,...

- No New Physics discovery at LHC
- With the LHC resolution to probe Higgs compositeness, the Higgs could be as elementary as pion. So, how pointlike is it?
- λ can be significantly enhanced in EW baryogenesis models. HL-LHC will probe λ with 50% uncertainty [2]

Higgs is Really New Physics!

* We've never seen anything like it

* Harbinger of profound New Principles at work in quantum vacuum

* MUST LOOK AT IT CLOSELY!

[1]

ATLAS SUSY Searches* - 95% CL Lower Limits
December 2017

ATLAS Preliminary
 $\sqrt{s} = 7, 8, 13 \text{ TeV}$

Model	e, μ, τ, γ	Jets	E_T^{miss}	$L \mathcal{B} [fb^{-1}]$	Mass limit	$\sqrt{s} = 7, 8 \text{ TeV}$	$\sqrt{s} = 13 \text{ TeV}$	Reference	
Inclusive Searches	$\tilde{q}\tilde{q} \rightarrow q\tilde{q}\tilde{q}^0$	0	2-6 jets	Yes	36.1	$\tilde{q}\tilde{q}$	1.57 TeV	$m(\tilde{q}^0) < 200 \text{ GeV}, m(1^{\text{st}} \text{ gen. } \tilde{q}) = m(2^{\text{nd}} \text{ gen. } \tilde{q})$	1712.0232
	$\tilde{q}\tilde{q} \rightarrow q\tilde{q}\tilde{q}^0$ (compressed)	mono-jet	1-3 jets	Yes	36.1	$\tilde{q}\tilde{q}$	710 GeV	$m(\tilde{q}) = m(\tilde{q}^0) + 5 \text{ GeV}$	1711.03301
	$\tilde{g}\tilde{g} \rightarrow g\tilde{g}\tilde{g}^0$	0	2-6 jets	Yes	36.1	$\tilde{g}\tilde{g}$	2.02 TeV	$m(\tilde{g}^0) < 200 \text{ GeV}$	1712.0232
	$\tilde{g}\tilde{g} \rightarrow g\tilde{g}\tilde{g}^0$	0	2-6 jets	Yes	36.1	$\tilde{g}\tilde{g}$	2.01 TeV	$m(\tilde{g}^0) < 200 \text{ GeV}, m(\tilde{g}^0) > 0.5(m(\tilde{g}^0) + m(\tilde{g}))$	1611.05791
	$\tilde{g}\tilde{g} \rightarrow g\tilde{g}\tilde{g}^0$	$ee, \mu\mu$	2 jets	Yes	14.7	$\tilde{g}\tilde{g}$	1.7 TeV	$m(\tilde{g}^0) < 300 \text{ GeV}$	1706.03731
	$\tilde{g}\tilde{g} \rightarrow g\tilde{g}\tilde{g}^0$	$3e, \mu$	4 jets	-	36.1	$\tilde{g}\tilde{g}$	1.87 TeV	$m(\tilde{g}^0) = 0 \text{ GeV}$	1708.02794
	$\tilde{g}\tilde{g} \rightarrow g\tilde{g}\tilde{g}^0$	0	7-11 jets	Yes	36.1	$\tilde{g}\tilde{g}$	1.8 TeV	$m(\tilde{g}^0) < 400 \text{ GeV}$	1607.05979
	GMSB (\tilde{L} NLSP)	$1-2 \tau + 0-1 \ell$	0-2 jets	Yes	3.2	$\tilde{g}\tilde{g}$	2.0 TeV	$\tau \tau (\text{NLSP}) < 0.1 \text{ mm}$	ATLAS-CONF-2017-080
	GGM (bino NLSP)	2 γ	2 jets	Yes	36.1	$\tilde{g}\tilde{g}$	2.05 TeV	$m(\tilde{g}^0) = 1700 \text{ GeV}, \tau \tau (\text{NLSP}) < 0.1 \text{ mm}, \mu \mu > 0$	ATLAS-CONF-2017-080
	GGM (higgsino-bino NLSP)	0	mono-jet	Yes	20.3	$\tilde{g}\tilde{g}$	2.05 TeV	$m(\tilde{g}^0) > 1.5 \times 10^{-4} \text{ eV}, m(\tilde{g}) = m(\tilde{g}^0) + 1.5 \text{ TeV}$	1502.01519
1 st gen. squarks	$\tilde{u}\tilde{u} \rightarrow u\tilde{u}\tilde{u}^0$	0	3 b	Yes	36.1	$\tilde{u}\tilde{u}$	1.92 TeV	$m(\tilde{u}^0) < 600 \text{ GeV}$	1711.01901
	$\tilde{u}\tilde{u} \rightarrow u\tilde{u}\tilde{u}^0$	0-1 e, μ	3 b	Yes	36.1	$\tilde{u}\tilde{u}$	1.97 TeV	$m(\tilde{u}^0) < 200 \text{ GeV}$	1711.01901
	$\tilde{d}\tilde{d} \rightarrow d\tilde{d}\tilde{d}^0$	0	2 b	Yes	36.1	$\tilde{d}\tilde{d}$	950 GeV	$m(\tilde{d}^0) < 420 \text{ GeV}$	1708.09266
	$\tilde{b}\tilde{b} \rightarrow b\tilde{b}\tilde{b}^0$	0	2 b	Yes	36.1	$\tilde{b}\tilde{b}$	275-700 GeV	$m(\tilde{b}^0) < 200 \text{ GeV}, m(\tilde{b}^0) = m(\tilde{b}^0) + 100 \text{ GeV}$	1706.03731
	$\tilde{t}\tilde{t} \rightarrow t\tilde{t}\tilde{t}^0$	2 e, μ (SS)	1 b	Yes	36.1	$\tilde{t}\tilde{t}$	200-720 GeV	$m(\tilde{t}^0) = 2m(\tilde{b}^0), m(\tilde{t}^0) = m(\tilde{t}^0) + 55 \text{ GeV}$	1209.2102, ATLAS-CONF-2016-077
	$\tilde{t}\tilde{t} \rightarrow t\tilde{t}\tilde{t}^0$	0-2 e, μ	1-2 b	Yes	4.7/13.3	$\tilde{t}\tilde{t}$	117-170 GeV	$m(\tilde{t}^0) = 1 \text{ GeV}$	1506.09616, 1709.04163, 1711.11520
	$\tilde{t}\tilde{t} \rightarrow t\tilde{t}\tilde{t}^0$	0-2 e, μ	0-3 jets+1-2 b	Yes	20.3/336.1	$\tilde{t}\tilde{t}$	90-198 GeV	$m(\tilde{t}^0) = m(\tilde{t}^0) + 5 \text{ GeV}$	1711.03301
	$\tilde{t}\tilde{t} \rightarrow t\tilde{t}\tilde{t}^0$	0	mono-jet	Yes	36.1	$\tilde{t}\tilde{t}$	90-430 GeV	$m(\tilde{t}^0) > 150 \text{ GeV}$	1403.5222
	$\tilde{t}\tilde{t} \rightarrow t\tilde{t}\tilde{t}^0$	$\tilde{t}\tilde{t}$ (natural GMSB)	2 e, μ (Z)	1 b	Yes	20.3	$\tilde{t}\tilde{t}$	150-600 GeV	1706.03986
	$\tilde{t}\tilde{t} \rightarrow t\tilde{t}\tilde{t}^0$	$3e, \mu$ (Z)	1 b	Yes	36.1	$\tilde{t}\tilde{t}$	290-750 GeV	$m(\tilde{t}^0) = 0 \text{ GeV}$	1706.03986
EW direct	$\tilde{g}\tilde{g} \rightarrow g\tilde{g}\tilde{g}^0$	2 e, μ	0	Yes	36.1	$\tilde{g}\tilde{g}$	90-500 GeV	$m(\tilde{g}^0) = 0$	ATLAS-CONF-2017-039
	$\tilde{g}\tilde{g} \rightarrow g\tilde{g}\tilde{g}^0$	2 e, μ	0	Yes	36.1	$\tilde{g}\tilde{g}$	750 GeV	$m(\tilde{g}^0) = 0, m(\tilde{g}^0) > 0.5(m(\tilde{g}^0) + m(\tilde{g}^0))$	ATLAS-CONF-2017-039
	$\tilde{g}\tilde{g} \rightarrow g\tilde{g}\tilde{g}^0$	2 τ	-	Yes	36.1	$\tilde{g}\tilde{g}$	760 GeV	$m(\tilde{g}^0) = 0, m(\tilde{g}^0) > 0.5(m(\tilde{g}^0) + m(\tilde{g}^0))$	1708.07875
	$\tilde{g}\tilde{g} \rightarrow g\tilde{g}\tilde{g}^0$	3 e, μ	0	Yes	36.1	$\tilde{g}\tilde{g}$	1.13 TeV	$m(\tilde{g}^0) = m(\tilde{g}^0), m(\tilde{g}^0) = 0, m(\tilde{g}^0) > 0.5(m(\tilde{g}^0) + m(\tilde{g}^0))$	ATLAS-CONF-2017-039
	$\tilde{g}\tilde{g} \rightarrow g\tilde{g}\tilde{g}^0$	2-3 e, μ	0-2 jets	Yes	36.1	$\tilde{g}\tilde{g}$	580 GeV	$m(\tilde{g}^0) = m(\tilde{g}^0), m(\tilde{g}^0) = 0, \ell$ decoupled	ATLAS-CONF-2017-039
	$\tilde{g}\tilde{g} \rightarrow g\tilde{g}\tilde{g}^0$	e, μ, γ	0-2 b	Yes	20.3	$\tilde{g}\tilde{g}$	270 GeV	$m(\tilde{g}^0) = m(\tilde{g}^0), m(\tilde{g}^0) = 0, \ell$ decoupled	1501.07110
	$\tilde{g}\tilde{g} \rightarrow g\tilde{g}\tilde{g}^0$	4 e, μ	0	Yes	20.3	$\tilde{g}\tilde{g}$	635 GeV	$m(\tilde{g}^0) = m(\tilde{g}^0), m(\tilde{g}^0) = 0, m(\tilde{g}^0) > 0.5(m(\tilde{g}^0) + m(\tilde{g}^0))$	1405.5086
	GGM (bino NLSP) weak prod. $\tilde{g}\tilde{g} \rightarrow \gamma G$	1 $e, \mu + \gamma$	-	Yes	20.3	$\tilde{g}\tilde{g}$	115-370 GeV	$m(\tilde{g}^0) = m(\tilde{g}^0), m(\tilde{g}^0) = 0, m(\tilde{g}^0) > 0.5(m(\tilde{g}^0) + m(\tilde{g}^0))$	1507.05493
	GGM (bino NLSP) weak prod. $\tilde{g}\tilde{g} \rightarrow \gamma G$	2 γ	-	Yes	36.1	$\tilde{g}\tilde{g}$	1.06 TeV	$\tau \tau < 1 \text{ mm}$	ATLAS-CONF-2017-080
	Long-lived particles	Direct $\tilde{L}\tilde{L}$ prod., long-lived \tilde{L}^0	Disapp. trk	1 jet	Yes	36.1	$\tilde{L}\tilde{L}$	460 GeV	$m(\tilde{L}^0) = m(\tilde{L}^0) = 160 \text{ MeV}, \tau(\tilde{L}^0) > 0.2 \text{ ns}$
Direct $\tilde{L}\tilde{L}$ prod., long-lived \tilde{L}^0		dE/dx trk	-	Yes	18.4	$\tilde{L}\tilde{L}$	495 GeV	$m(\tilde{L}^0) = m(\tilde{L}^0) = 160 \text{ MeV}, \tau(\tilde{L}^0) < 15 \text{ ns}$	1506.05332
Stable, stopped \tilde{R} -hadron		0	1-5 jets	Yes	27.9	$\tilde{R}\tilde{R}$	850 GeV	$m(\tilde{R}^0) = 100 \text{ GeV}, 10 \mu\text{s} < \tau < 1000 \text{ s}$	1310.6584
Stable \tilde{g} -hadron		trk	-	-	3.2	$\tilde{g}\tilde{g}$	1.58 TeV	$m(\tilde{g}^0) = 100 \text{ GeV}, \tau > 10 \text{ ns}$	1606.05129
Metastable \tilde{g} -hadron		dE/dx trk	-	-	3.2	$\tilde{g}\tilde{g}$	1.57 TeV	$m(\tilde{g}^0) = 100 \text{ GeV}, \tau > 10 \text{ ns}$	1604.04520
Metastable \tilde{g} -hadron, $\tilde{g} \rightarrow q\tilde{q}^0$		displ. vtx	-	Yes	32.8	$\tilde{g}\tilde{g}$	2.37 TeV	$m(\tilde{g}^0) = 100 \text{ GeV}, \tau > 0.17 \text{ ns}, m(\tilde{g}^0) = 100 \text{ GeV}$	1710.04901
GMSB, stable $\tilde{L}^0 \rightarrow \tau \tilde{L}^0 \rightarrow \tau \tilde{L}^0 + \tau \tilde{L}^0 + \tau \tilde{L}^0$		1-2 μ	-	Yes	19.1	$\tilde{L}\tilde{L}$	537 GeV	$10^{-9} \text{ s} < \tau < 50$	1411.6795
GMSB, $\tilde{L}^0 \rightarrow e \tilde{L}^0$, long-lived \tilde{L}^0		2 γ	-	Yes	20.3	$\tilde{L}\tilde{L}$	440 GeV	$1 < \tau < 10^3 \text{ s}$, SPSS model	1409.5542
$\tilde{g}\tilde{g} \rightarrow g\tilde{g}\tilde{g}^0$		displ. $e\ell/g\mu/\mu\mu$	-	-	20.3	$\tilde{g}\tilde{g}$	1.0 TeV	$7 < \tau < 10^3 \text{ s}$, $740 \text{ MeV}, m(\tilde{g}) = 1.3 \text{ TeV}$	1504.05182
RPV		LFV $\tilde{g}\tilde{g} \rightarrow \tau \tilde{g} \tau + X, \tilde{g} \rightarrow \tau \tilde{g} \tau / \mu \tau / \mu \mu$	$e\mu, e\tau, \mu\tau$	-	-	3.2	$\tilde{g}\tilde{g}$	1.9 TeV	$A_{13} = 0.11, A_{33} = 0.02, \mu = 0.07$
	Bilinear RPV GMSB	2 e, μ (SS)	0-3 b	Yes	20.3	$\tilde{g}\tilde{g}$	1.45 TeV	$m(\tilde{g}^0) = m(\tilde{g}^0), \tau < 1 \text{ mm}$	1404.2500
	$\tilde{L}\tilde{L} \rightarrow \tau \tilde{L}^0 \tau + \tau \tilde{L}^0 \tau + \tau \tilde{L}^0 \tau$	4 e, μ	-	Yes	13.3	$\tilde{L}\tilde{L}$	1.14 TeV	$m(\tilde{L}^0) = 400 \text{ GeV}, A_{13} = 0 (A = 1, 2)$	ATLAS-CONF-2016-075
	$\tilde{L}\tilde{L} \rightarrow \tau \tilde{L}^0 \tau + \tau \tilde{L}^0 \tau + \tau \tilde{L}^0 \tau$	3 $e, \mu + \tau$	-	Yes	20.3	$\tilde{L}\tilde{L}$	450 GeV	$m(\tilde{L}^0) = 0.2 \times m(\tilde{L}^0), A_{13} \neq 0$	1405.5086
	$\tilde{g}\tilde{g} \rightarrow g\tilde{g}\tilde{g}^0$	0	4-5 large-R jets	-	36.1	$\tilde{g}\tilde{g}$	1.875 TeV	$m(\tilde{g}^0) = 1075 \text{ GeV}$	SUSY-2016-122
	$\tilde{g}\tilde{g} \rightarrow g\tilde{g}\tilde{g}^0$	1 e, μ	8-10 jets/0-4 b	-	36.1	$\tilde{g}\tilde{g}$	2.1 TeV	$m(\tilde{g}^0) = 1 \text{ TeV}, A_{13} = 0$	1704.08493
	$\tilde{g}\tilde{g} \rightarrow g\tilde{g}\tilde{g}^0$	1 e, μ	8-10 jets/0-4 b	-	36.1	$\tilde{g}\tilde{g}$	1.65 TeV	$m(\tilde{g}^0) = 1 \text{ TeV}, A_{13} = 0$	1704.08493
	$\tilde{g}\tilde{g} \rightarrow g\tilde{g}\tilde{g}^0$	0	2 jets + 2 b	-	36.7	$\tilde{g}\tilde{g}$	100-470 GeV	$BR(\tilde{g} \rightarrow b\tilde{g}) > 20\%$	1710.07171
	$\tilde{g}\tilde{g} \rightarrow g\tilde{g}\tilde{g}^0$	2 e, μ	2 b	-	36.1	$\tilde{g}\tilde{g}$	480-610 GeV	$m(\tilde{g}^0) = 200 \text{ GeV}$	1710.05544
	Other	Scalar charm, $\tilde{c} \rightarrow c\tilde{c}^0$	0	2 c	Yes	20.3	$\tilde{c}\tilde{c}$	510 GeV	$m(\tilde{c}^0) < 200 \text{ GeV}$

*Only a selection of the available mass limits on new states or phenomena is shown. Many of the limits are based on simplified models, c.f. refs. for the assumptions made.

WHAT BRINGS US TO THE HIGGS FACTORIES

In the European PP Strategy Update 2020, Higgs factories are the highest priority future initiatives [3]

Several projects on the market ($\sim 10^6$ Higgs bosons)

All electron-positron colliders

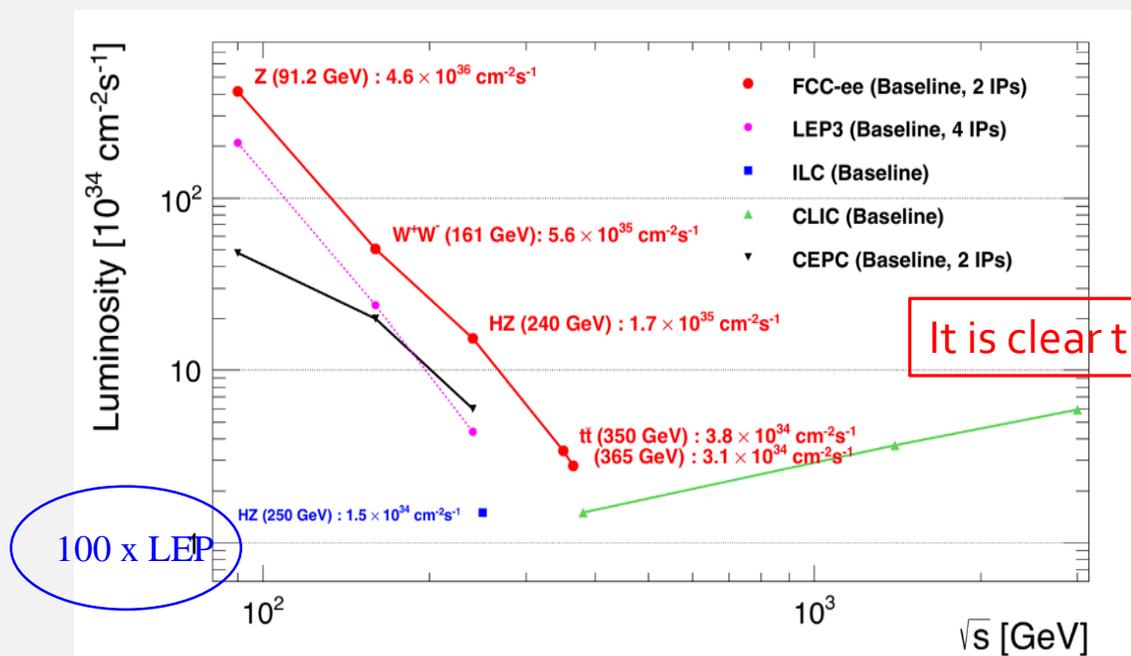
- Initial state well-defined
- High (TeV) center-of-mass energies - LCs
- Clean environment

→ High-precision measurements (dominated by statistical uncertainty)

Linear (ILC, CLIC) vs. circular (CEPC, FCCee)

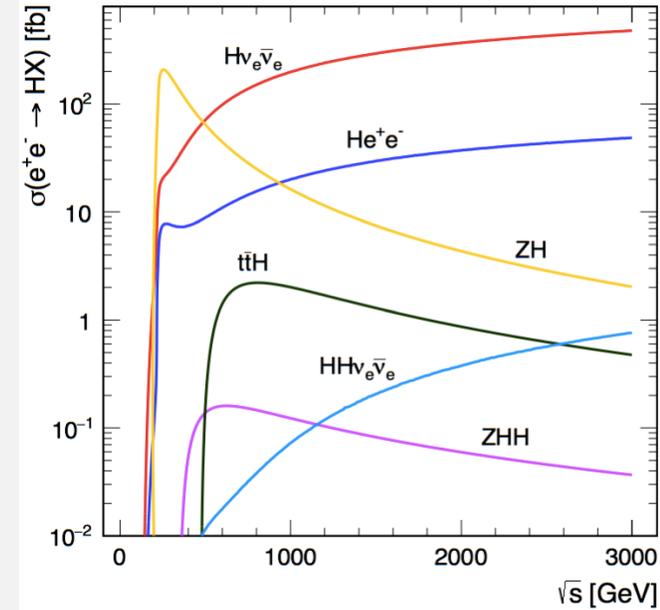
But, other aspects are also important:

- Extensibility of the physics span - flexibility to accommodate other options (pp, hh, ep, gamma gamma, plasma....)
- Flexibility to accommodate changes in scenario (i.e. unexpected HL-LHC discovery)
- Technological feasibility and cost
- Politics (it's a game of supremacy, unfortunately)

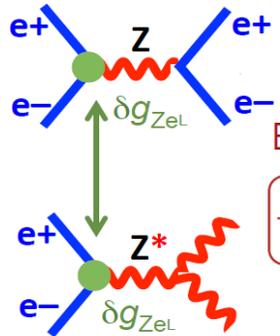


LINEAR COLLIDERS

- Comes as mature technological options developed for decade(s) – ‘ready to take’
- Staged, upgradable machines
 - Various Higgs production mechanisms accessible over the energy scale span
 - Less precise determination of an observable at high energy leads to the same precision on coupling as at low energy
- Beam polarization
 - Chiral nature of charge currents results in significant sensitivity of WW-fusion cross-section on polarization scheme ($\sim 2 \cdot \mathcal{L}$)
 - Provides new observables sensitive to New Physics
 - Helps characterization of newly discovered particles
- A few technical benefits
 - Triggerless
 - Power-pulsing



At low energy ($\sqrt{s}=m_Z$)



Effect grows as s
 $\left(\frac{3000}{91.2}\right)^2 \sim 1000$

At high energy ($\sqrt{s}=3\text{TeV}$)

Imagine measuring

$$\frac{d\sigma}{\sigma_{\text{SM}}} \Big|_{\sqrt{s}=m_Z} \sim 10^{-4} \Rightarrow \delta g_{ZeL} \sim 10^{-4}$$

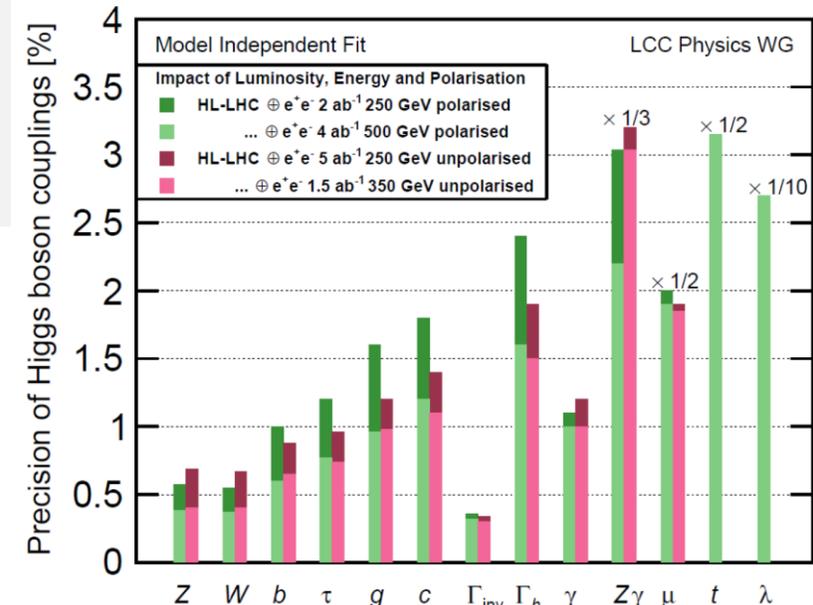
...equivalent to

$$\frac{d\sigma}{\sigma_{\text{SM}}} \Big|_{\sqrt{s}=3\text{TeV}} \sim 10\% \Rightarrow \delta g_{ZeL} \sim 10^{-4}$$

same precision!

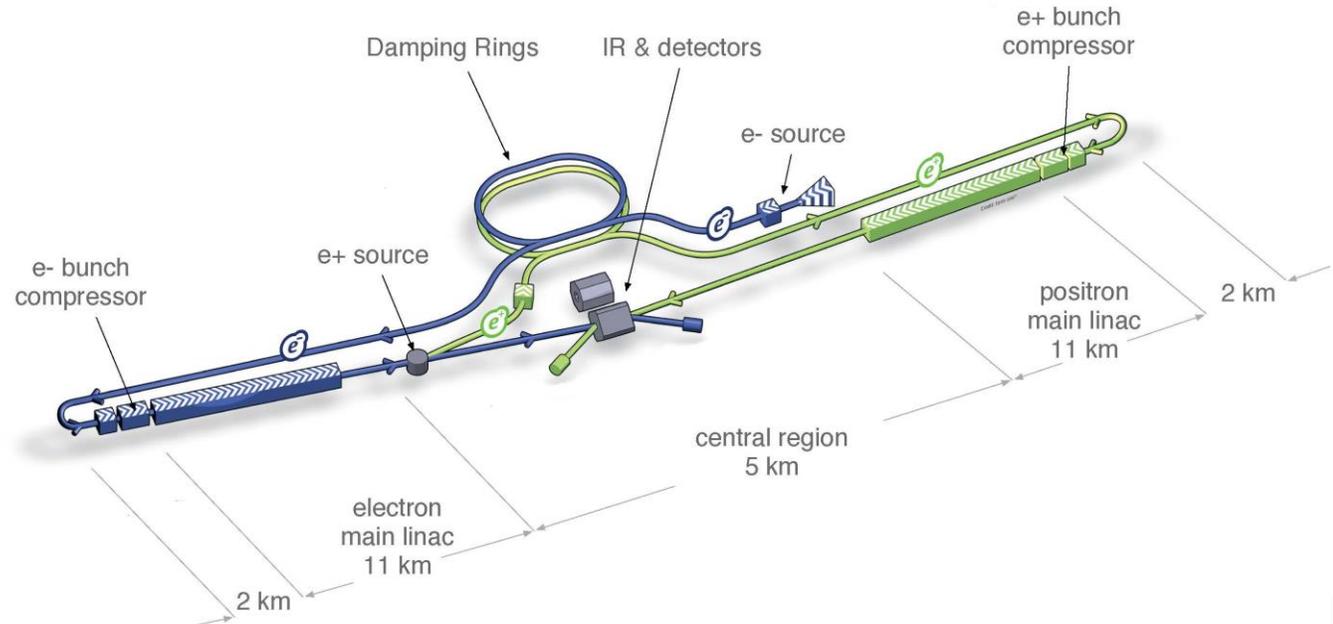
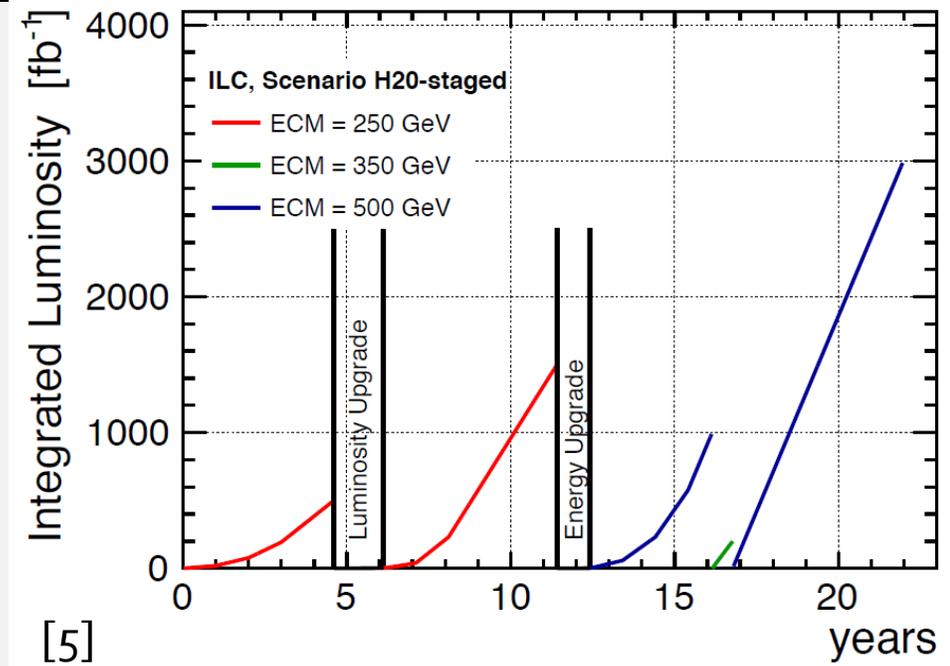
[4]

-> strongly benefit from high energies





- e^+e^- centre-of-mass energy
 - first stage: 250 GeV
 - tunable
 - upgrades: 500 GeV, 1 TeV
 - further options:
running at Z pole & WW threshold
- luminosity at 250 GeV:
 - $1.35 \times 10^{34} / \text{cm}^2 / \text{s}$
 - upgrade $2.7 \times 10^{34} / \text{cm}^2 / \text{s}$ (cheap)
 - upgrade $5.4 \times 10^{34} / \text{cm}^2 / \text{s}$ (expensive)
- beam polarisation
 - $P(e^-) \geq \pm 80\%$
 - $P(e^+) = \pm 30\%$,
at 500 GeV upgradable to 60%
- total length (250 GeV): 20.5 km

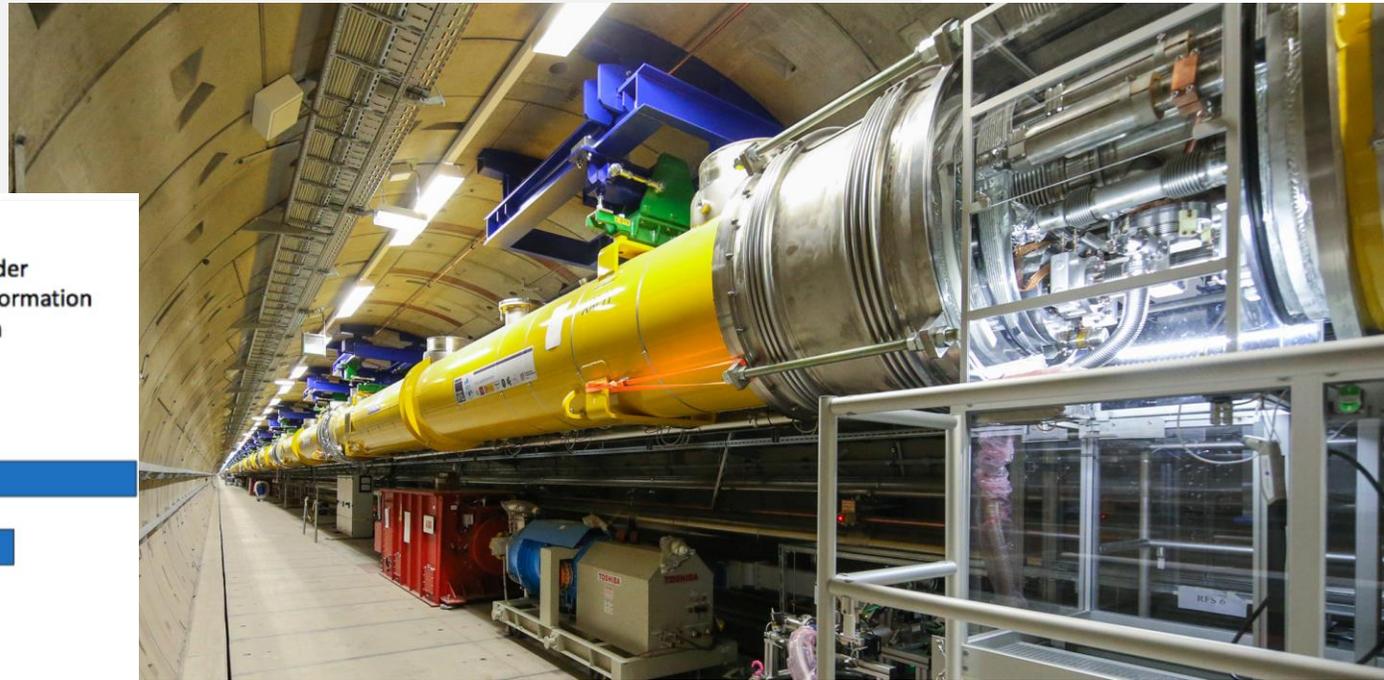
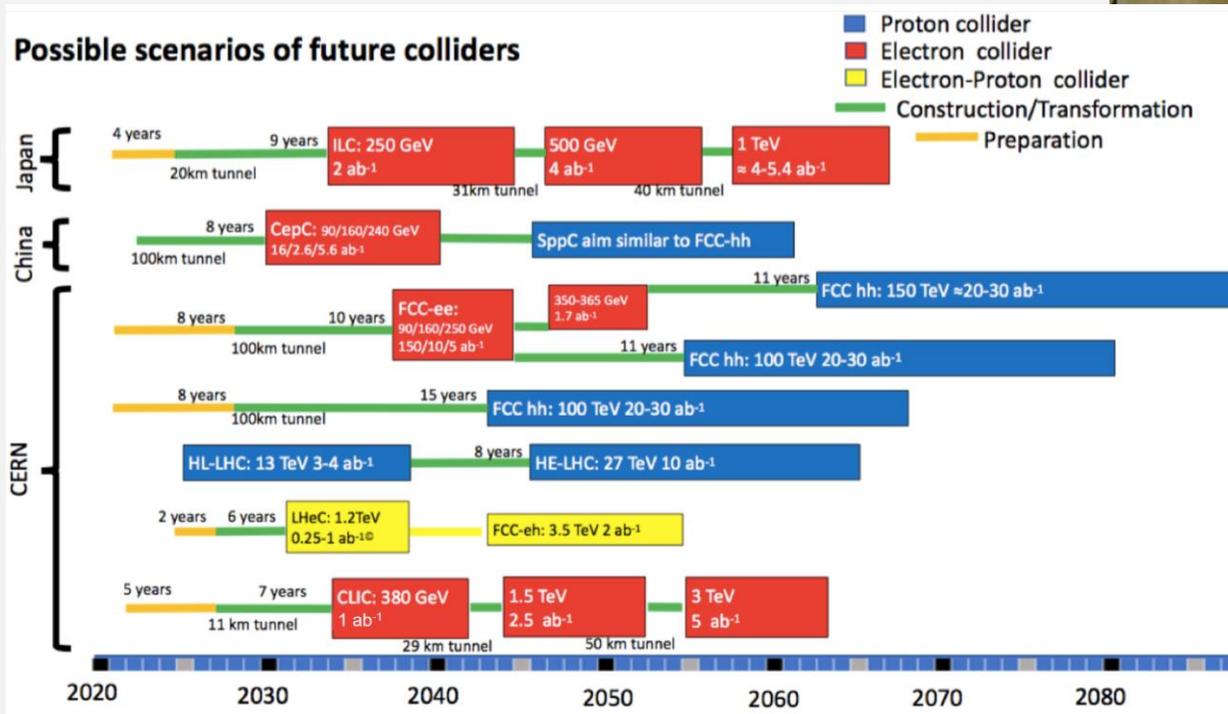




Timeline:

- Around 2000's – TESLA, NLC, JLC
- (2004) ILC based on 'cold' TESLA technology
- (2013) Technical Design Report [6]
- (2020) International Development Team (IDT)
- (?) Preparatory lab in Japan
- (2035) First collisions [7]

- Largest ever accelerator prototype (operating now as E-XFEL)
- Full industrialization of RF cavity production

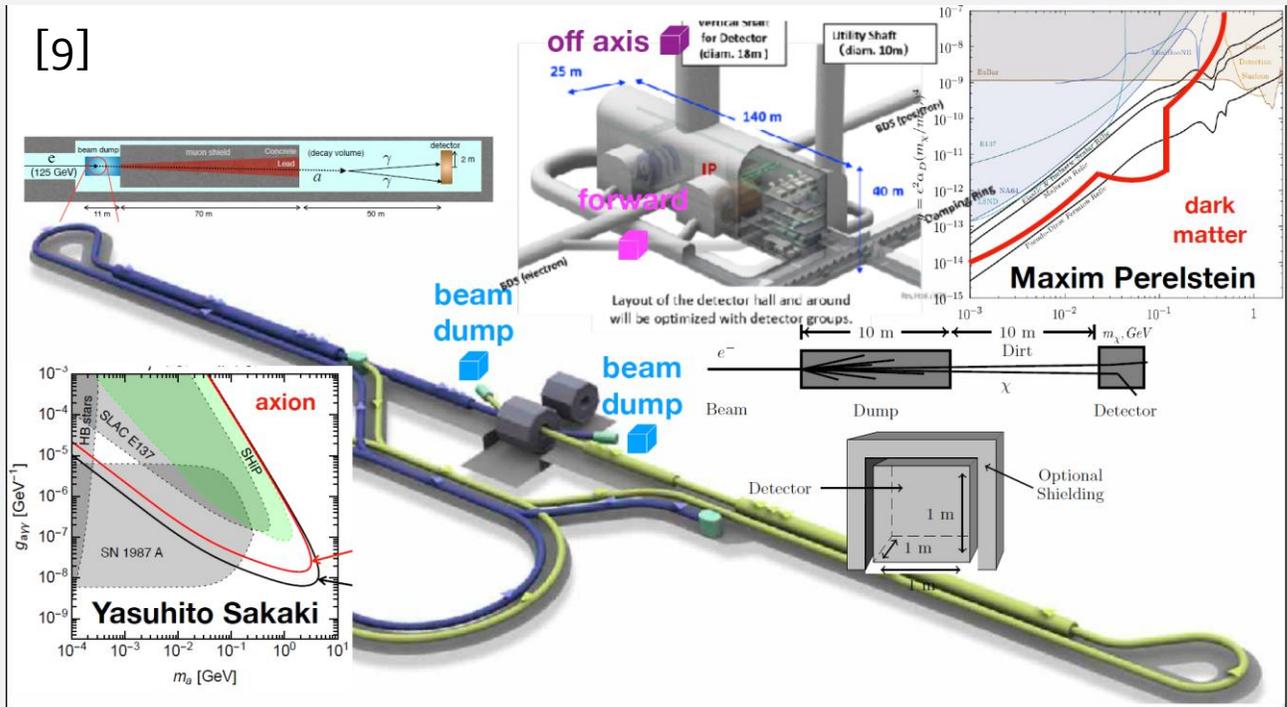


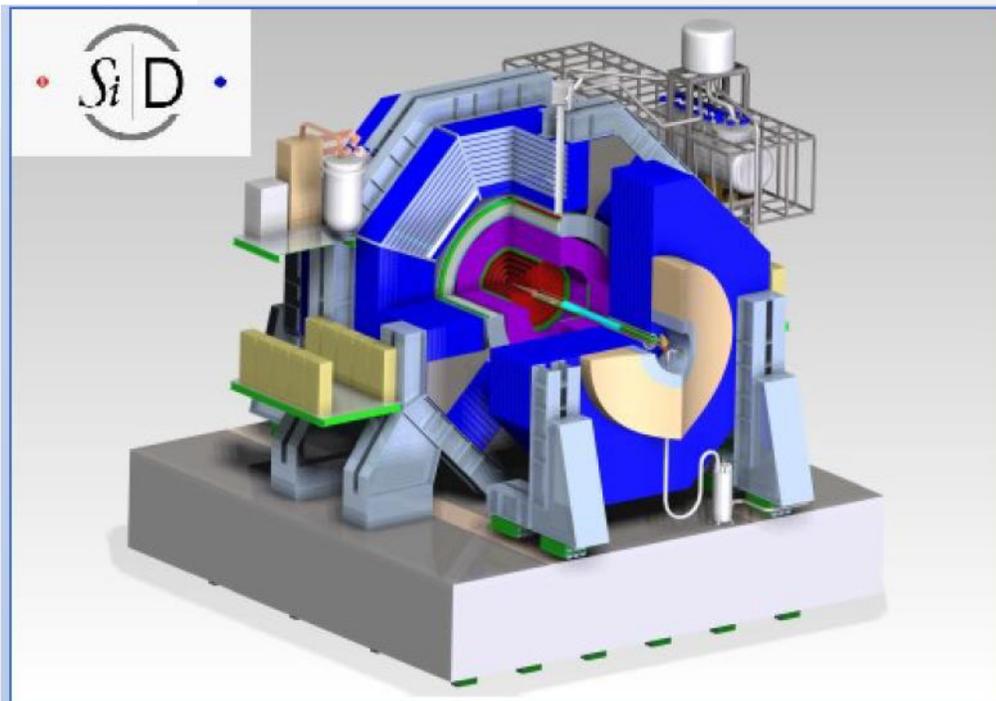


ILC comes with the collider program and rich auxiliary experiments

- At the LHC, experiments search for dark particles produced by pp collisions are placed in existing tunnels and caverns at CERN (FASER @ATLAS, MilliQan @CMS)
- Dark sector (ILC-BDX), fixed-target and beam dump experiments (ILCX)

Potential ILC site in Kitakami





SiD Detector

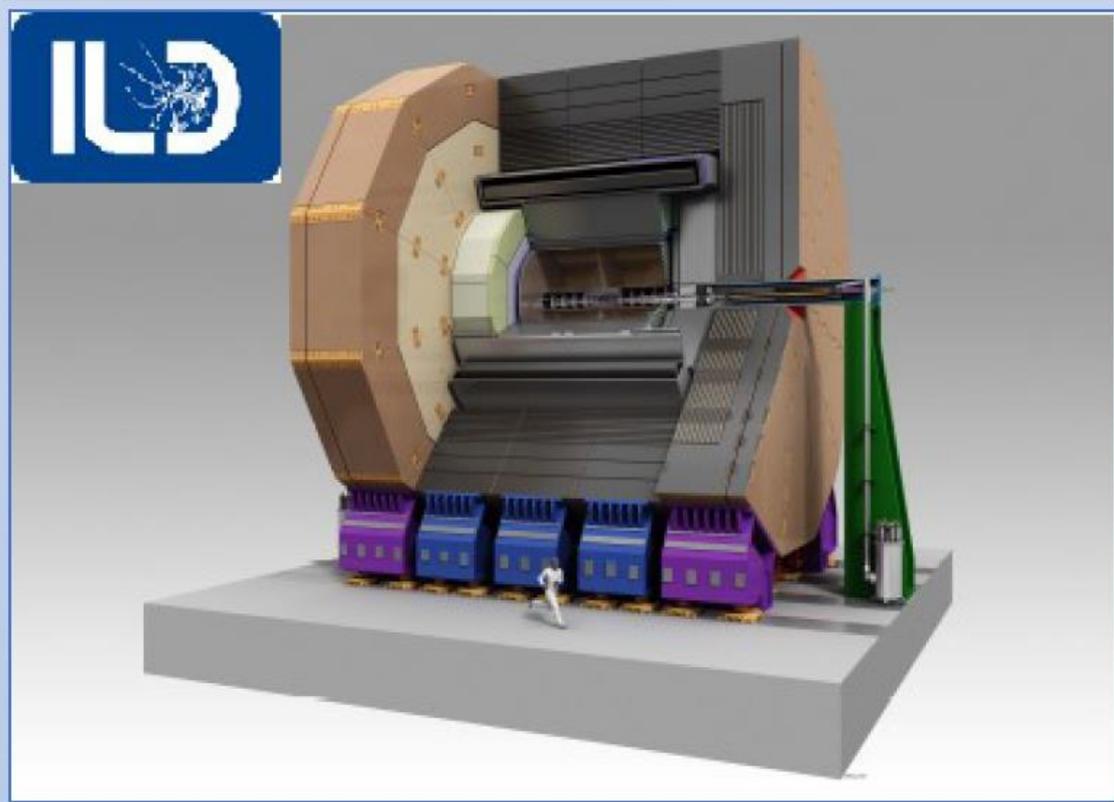
- 5 T field
- More compact
- All Si

Track momentum resolution: $\sigma_{1/p} < 5 \cdot 10^{-5} \text{ GeV}^{-1}$ CMS/40

Impact parameter resolution: $\sigma_d < 5 \mu\text{m} \oplus 10 \mu\text{m} \frac{1 \text{ GeV}}{p \sin^{3/2} \Theta}$ CMS/4

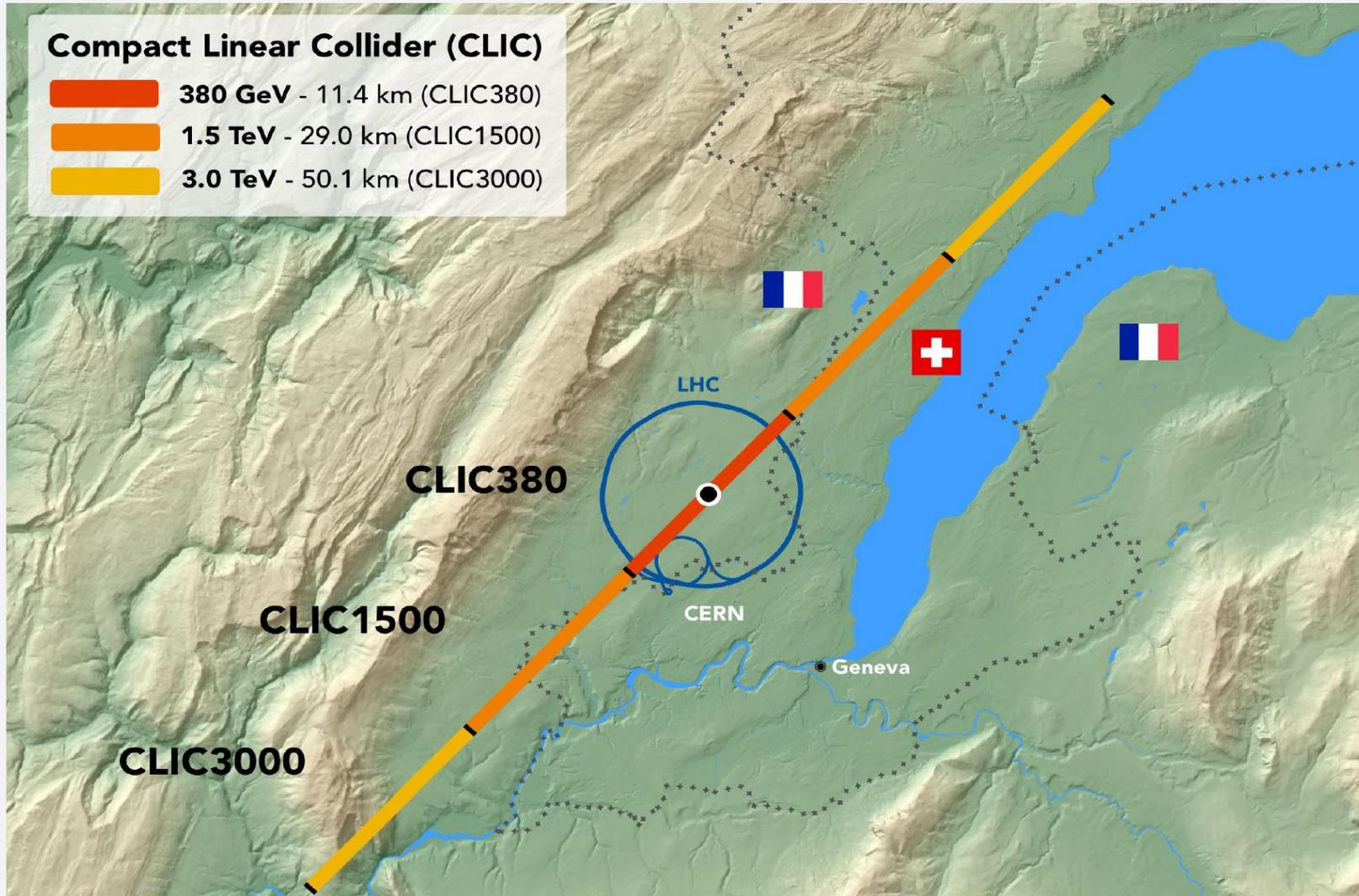
Jet energy resolution: $\sigma_E/E = 3 - 4\%$ (for highest jet energies) ATLAS/2

Hermecity: $\Theta_{min} = 5 \text{ mrad}$ ATLAS/3



ILD Detector

- 3.5 T field
- Optimized for CM energies 90 GeV - 1 TeV
- Si/gaseous tracking
- Particle flow calorimetry
- Mature design and available technologies

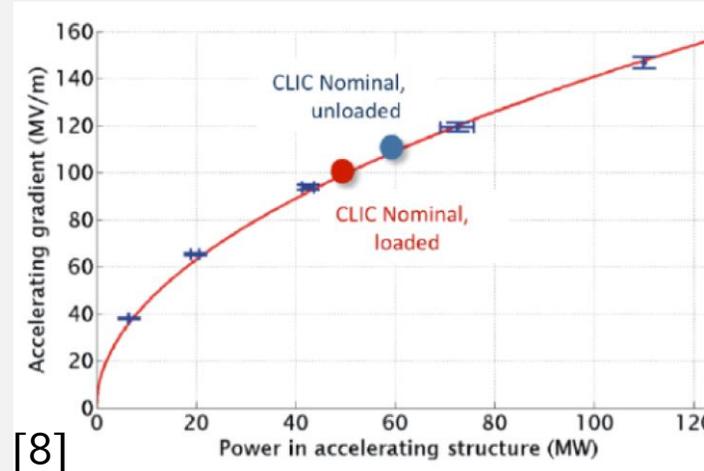




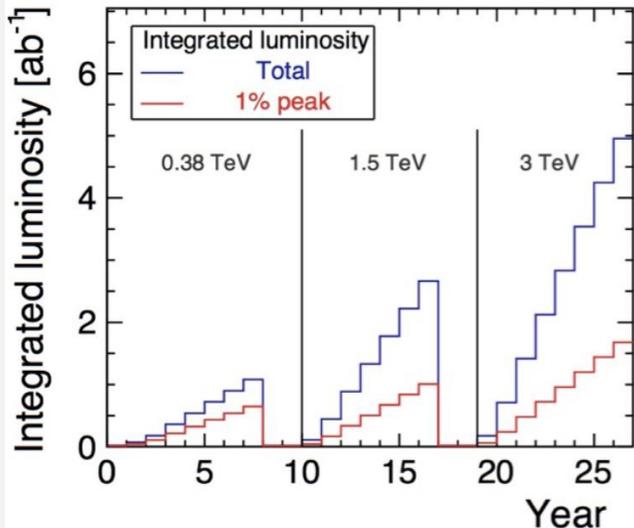
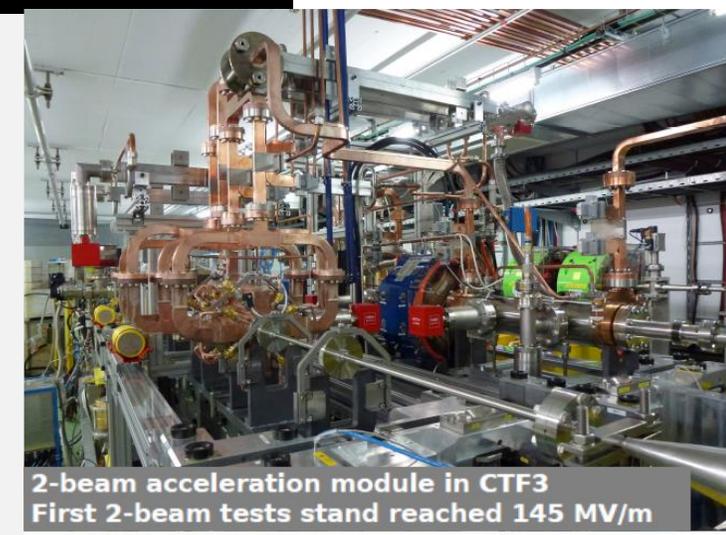
The only LC to go above 1 TeV

- CLIC Test Facility, CTF₃, at CERN now the 'CERN Linear Electron Accelerator for Research' facility, CLEAR

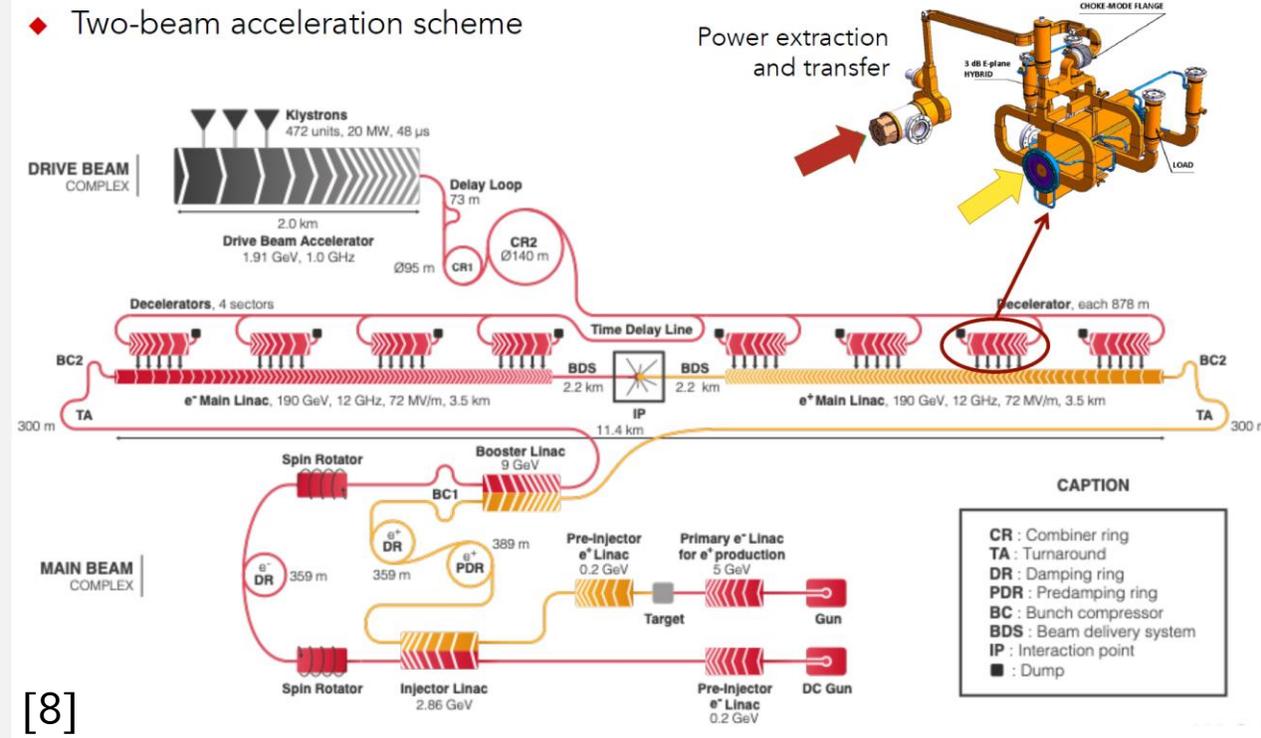
- Normal conductive high-current drive beam
- 380 GeV and 1.5 TeV – one drive-beam
- 3 TeV two drive-beam complexes
- 100 MV/m gradient in the main-beam cavities



[8]



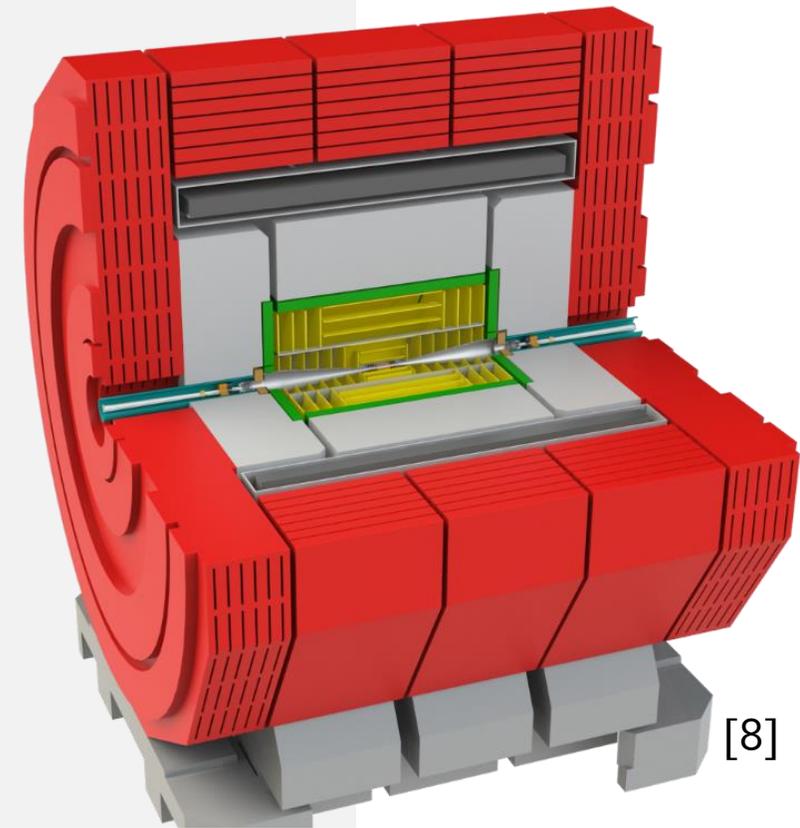
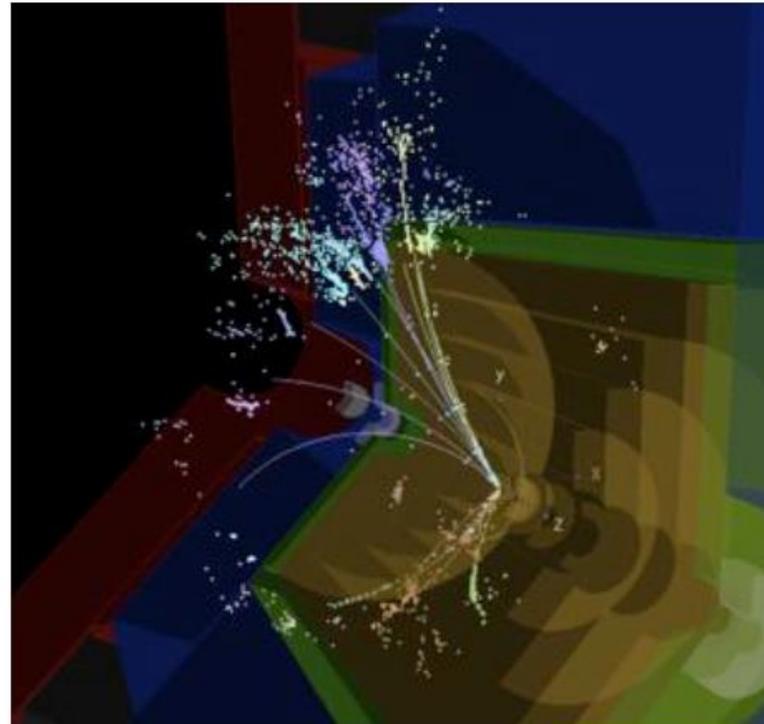
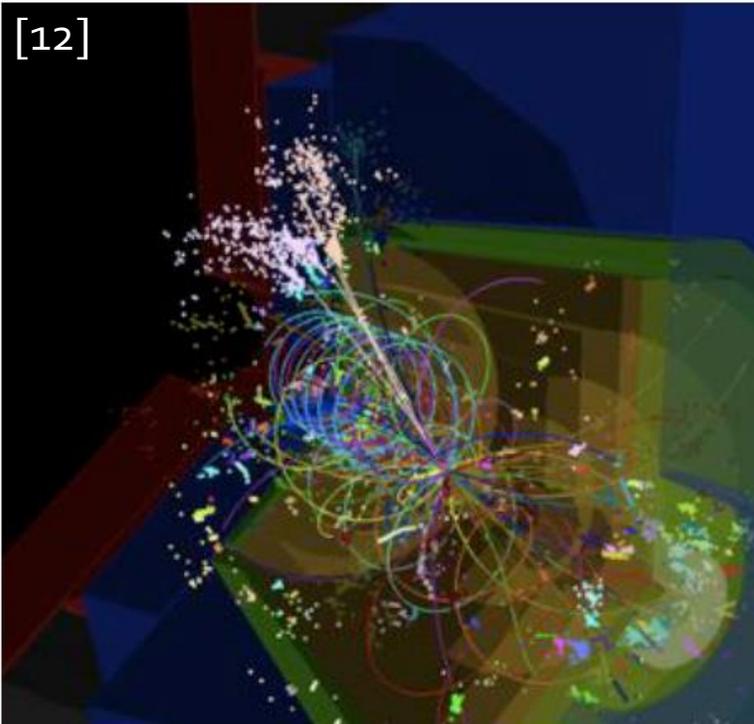
◆ Two-beam acceleration scheme



[8]

CLIC det

- 4 T field
- Ultra low-mass VTX
- All Si tracking
- Particle flow calorimetry
- Time-stamped readout (10 ns) due to pronounced Beamstrahlung background at higher energies



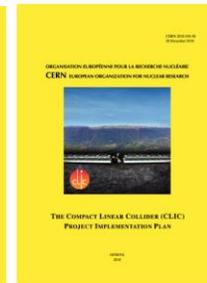
◆ 4 Yellow Reports 2018



Summary Report



Physics Potential



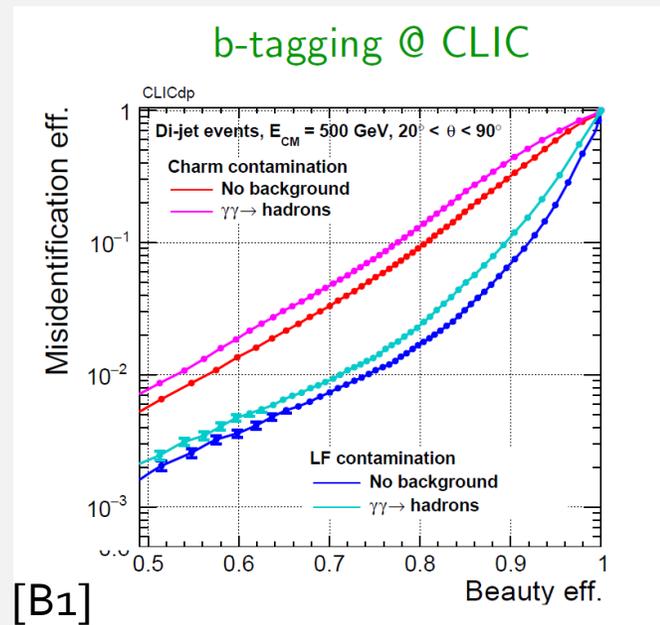
Project Implementation



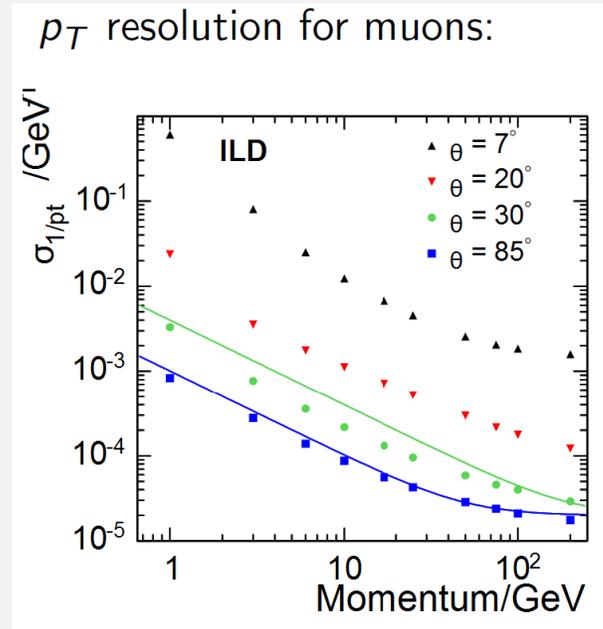
Detector Technologies

SIMILAR PERFORMANCE OF LC DETECTORS

c/b-tagging, Higgs branching ratios

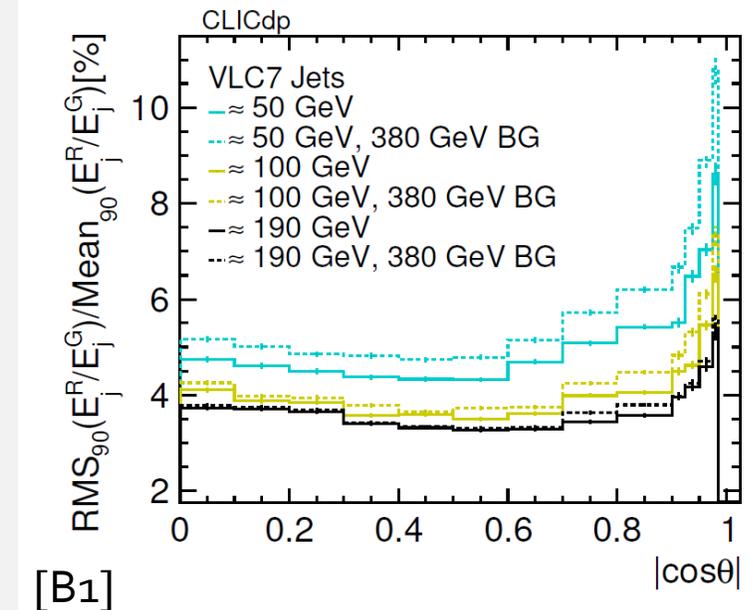


Higgs recoil mass, smuon endpoint,
Higgs coupling to muons



Separation of W/Z/H di-jets

3%–4% jet energy resolution gives $\sim 2.6 - 2.3\sigma$ W/Z separation



Particle Flow is the 'key word'. Only neutral particles ID (γ (30%), neutral hadrons (10%)) are left to calorimeters.

PHYSICS PROGRAMME AT A LC

- HIGGS PHYSICS

HIGGS COUPLINGS: model-independent measurements

κ -framework

EFT approach

HIGGS SELF-COUPLING

HIGGS AS A PROBE TO DARK SECTOR AND BSM IN GENERAL

CPV IN THE HIGGS SECTOR

High E

- t-PHYSICS

top-quark mass

electroweak couplings

rare decays

Low E

top Yukawa coupling

CP properties

BSM constraints

High E

- BSM

direct searches

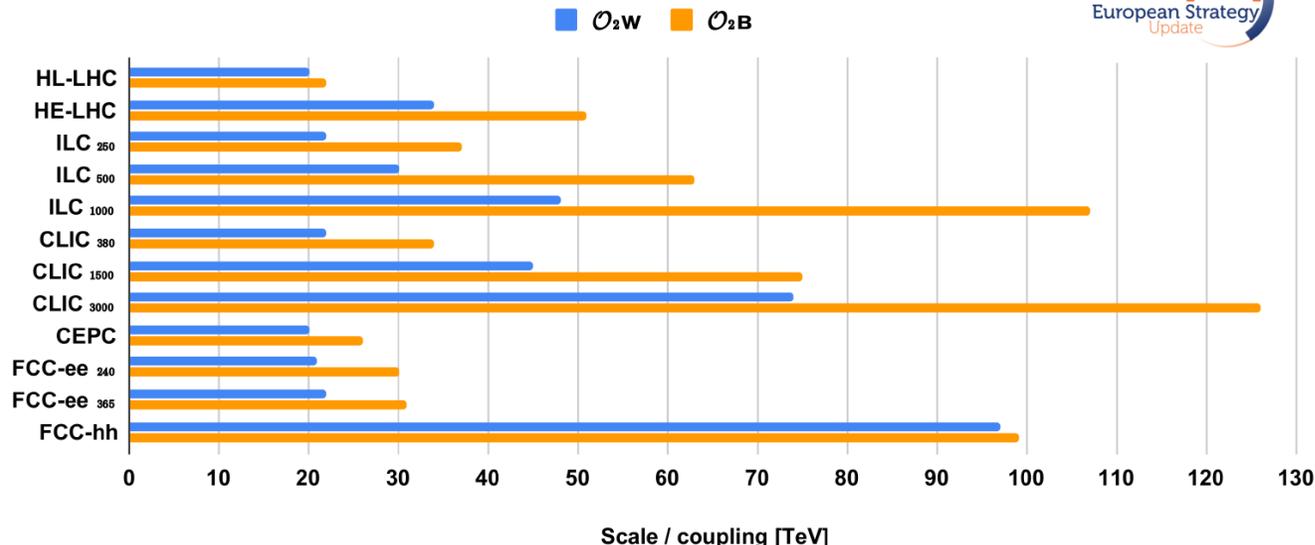
models with weak couplings or soft signatures

indirect searches

high sensitivity

High E

95% CL scale limits on 4-fermion contact interactions

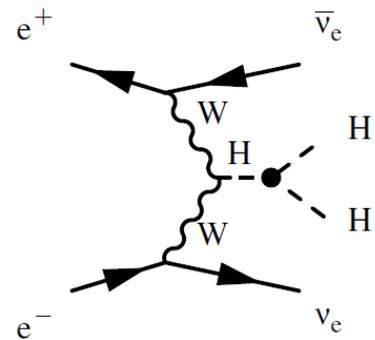
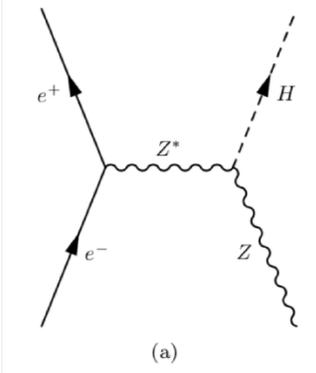
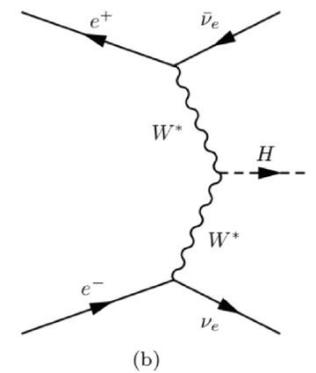
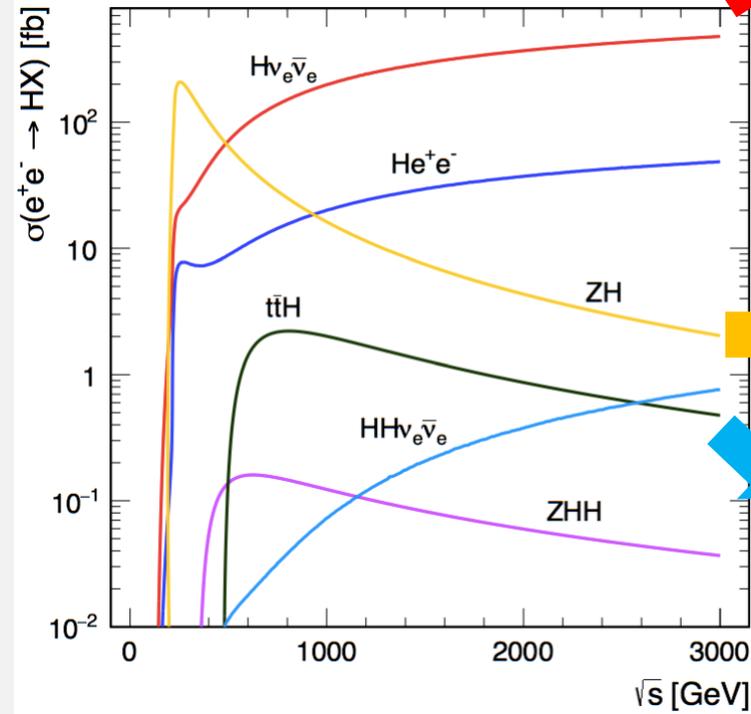


Due to staged realization of LCs, these are ideal machines to explore large physics span, with indirect access to the ~ 100 TeV scale

HIGGS PRODUCTION MECHANISMS AT LC

- **Higgsstrahlung** (ZH) is a unique feature of particle-antiparticle collisions (i.e. e^+e^- colliders)
 - It facilitates g_{HZZ} measurement in a **model-independent way** * (ZH cross-section)
 - Higgs invisible width can be determined from the recoil mass
 - Most of the Higgs couplings can be determined with a better precision than at HL-LHC only from ZH
-
- Linear colliders foreseen as staged machines benefit from additional statistics from **WW-fusion** (clear example is CLIC with $\sim 3M$ Higgs bosons at all stages)
 - **Double Higgs production** at higher energies enables self-coupling measurement

* Theory warning: level of accuracy $<1\%$ requires incorporation of loop-corrections \rightarrow loss of strict model-independence



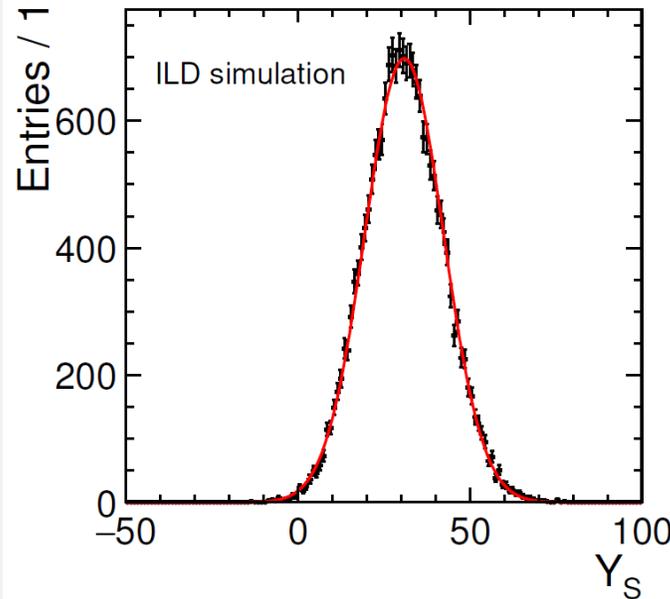
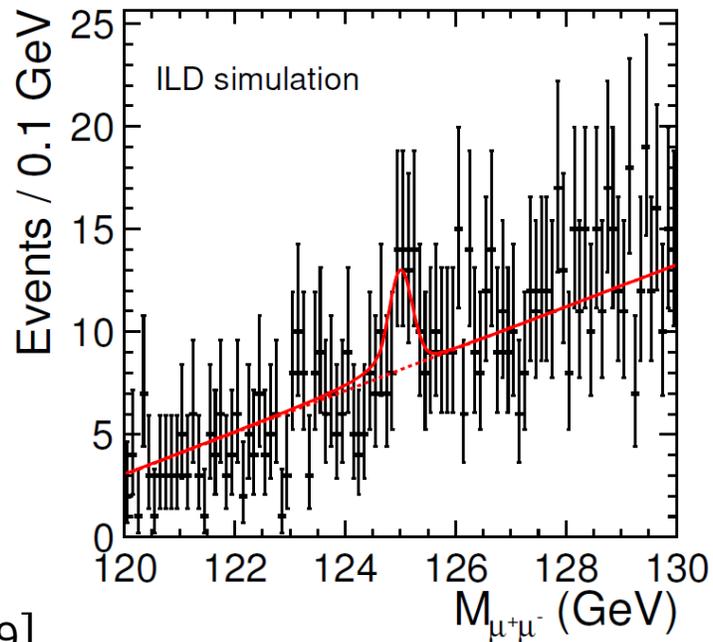
High-energy benefits, polarization, combination ➔ access to rare Higgs decays



- Clear advantage from rising cross-section for WW-fusion with energy
- ttH production, suitable i.e. for CPV study in the Higgs sector
- Multiple-Higgs production → self-coupling measurement
- Less precise determination of the observable at high energy leads to the same precision on coupling as at low energy

decays

Decay mode	Branching ratio
$H \rightarrow b\bar{b}$	56.1%
$H \rightarrow WW^*$	23.1%
$H \rightarrow gg$	8.5%
$H \rightarrow \tau^+\tau^-$	6.2%
$H \rightarrow c\bar{c}$	2.8%
$H \rightarrow ZZ^*$	2.9%
$H \rightarrow \gamma\gamma$	0.23%
$H \rightarrow Z\gamma$	0.16%
$H \rightarrow \mu^+\mu^-$	0.021%
Γ_H	4.2 MeV



$\sqrt{s} = 250$ GeV	$q\bar{q}H$	$\nu\bar{\nu}H$	ILC250	ILC250+500
L	34%	113%	23%	17%
R	36%	111%		
$\sqrt{s} = 500$ GeV	$q\bar{q}H$	$\nu\bar{\nu}H$	ILC500	
L	43%	37%	24%	
R	48%	106%		

[9]

HIGGS PHYSICS

Situation at LHC (HL-LHC, and pp in general)

- No absolute measurement of the production cross-section (like ZH at e+e- colliders)
- Higgs couplings come in combination:
$$\sigma(H) \times \text{BR}(H \rightarrow a + b) \sim \frac{\Gamma_{\text{prod}} \Gamma_{\text{decay}}}{\Gamma_{\text{tot}}}$$
- Only ratio of couplings can be directly determined (i.e. $g_{H\tau\tau}^2 / g_{HWW}^2$)

e+e- colliders

- Absolute measurement of the ZH cross-section
- Absolute measurement of the Higgs BRs
- Nearly model-independent determination of the Higgs total width and couplings
- High energy benefits of LCs: λ , CPV, BSM extensions of the Higgs sector

HIGGS COUPLINGS

How well do we need to know Higgs couplings?

- In many BSM models one expects only **% level deviations** from the SM couplings for BSM particles in the TeV range
- Higgs to EW bosons couplings are particularly sensitive to BSM; λ even more
- Example, 2HDM-type model in decoupling limit [B2]

$$\frac{g_{hVV}}{g_{h_{SM}VV}} \simeq 1 - 0.3\% \left(\frac{200 \text{ GeV}}{m_A} \right)^4$$

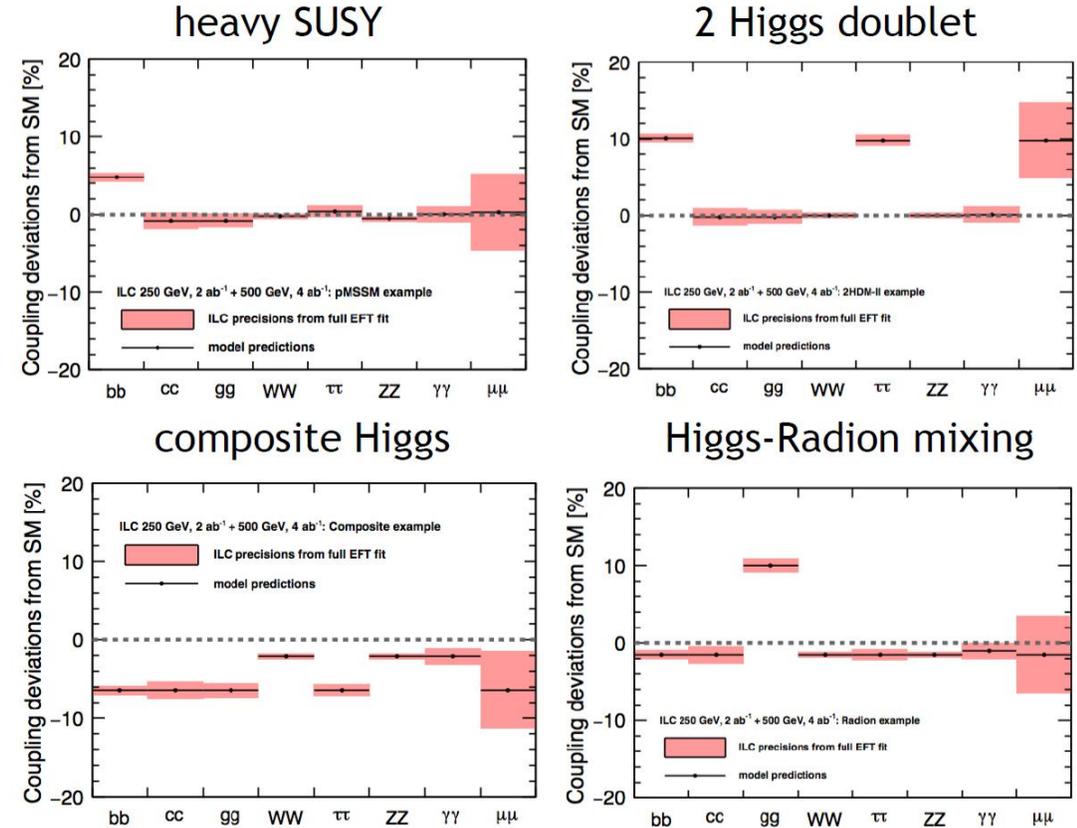
$$\frac{g_{htt}}{g_{h_{SM}tt}} = \frac{g_{hcc}}{g_{h_{SM}cc}} \simeq 1 - 1.7\% \left(\frac{200 \text{ GeV}}{m_A} \right)^2$$

$$\frac{g_{hbb}}{g_{h_{SM}bb}} = \frac{g_{h\tau\tau}}{g_{h_{SM}\tau\tau}} \simeq 1 + 40\% \left(\frac{200 \text{ GeV}}{m_A} \right)^2.$$

Percent order accuracy on Higgs couplings offers access to various BSM scenarios

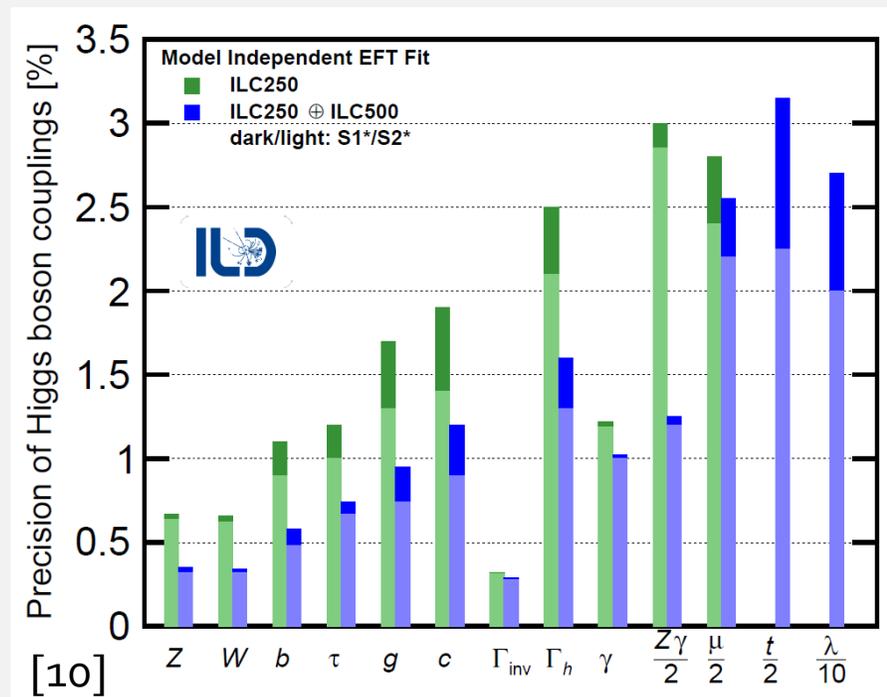
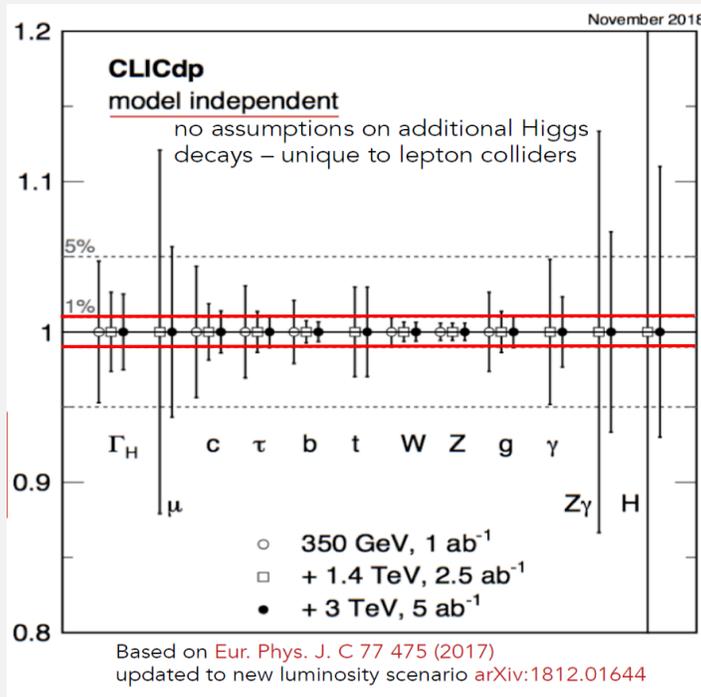
The models below are outside the HL-LHC reach

[T. Barklow et al. '17]



HIGGS COUPLINGS

Model independent approach*,
precision better than 1% for most couplings



Similar at circular colliders...

in %	FCC-ee 240 GeV	+FCC-ee 365 GeV	+HL- LHC
δg_{HZZ}	0.25	0.22	0.21
δg_{HWW}	1.3	0.47	0.44
δg_{Hbb}	1.4	0.68	0.58
δg_{Hcc}	1.8	1.23	1.20
δg_{Hgg}	1.7	1.03	0.83
$\delta g_{H\tau\tau}$	1.4	0.8	0.71
$\delta g_{H\mu\mu}$	9.6	8.6	3.4
$\delta g_{H\gamma\gamma}$	4.7	3.8	1.3
δg_{Htt}			3.3
$\delta \Gamma_H$	2.8	1.56	1.3

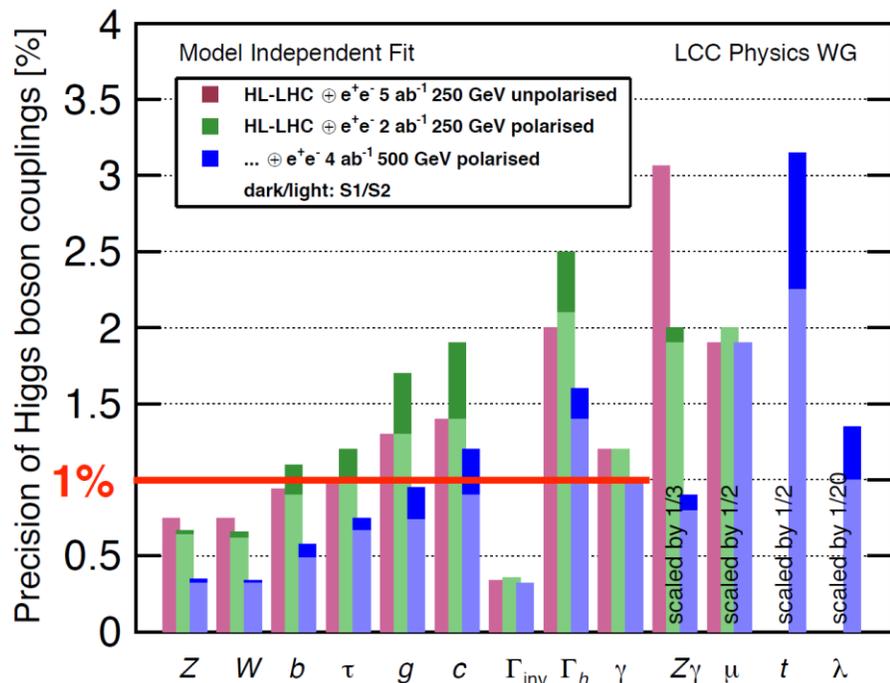
Statistical uncertainties are shown for 5 ab⁻¹@240 GeV and 1.5 ab⁻¹@365 GeV (from FCC-ee CDR)

COMBINATION WITH HL-LHC

To what extent future e+e- experiments are synergistic with the HL- LHC?

Evident synergy

- An example: ILC250 with 250 fb⁻¹
- Already the single measurement of the HZ cross section at ILC 250 yields a very large improvement of the HL-LHC accuracies



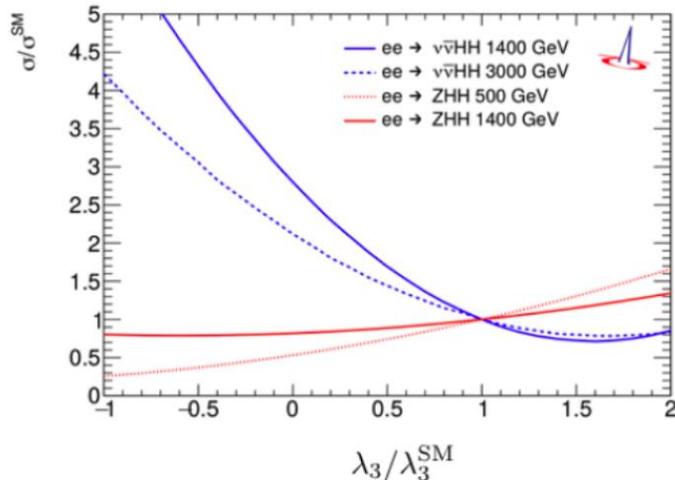
[11]

	Benchmark	HL-LHC	HL-LHC + CLIC		HL-LHC + FCC-ee	
			380 (4 ab ⁻¹)	380 (1 ab ⁻¹) + 1500 (2.5 ab ⁻¹)	240	365
$g_{HZZ}^{\text{eff}} [\%]$	SMEFT _{ND}	3.6	0.3	0.2	0.5	0.3
$g_{HWW}^{\text{eff}} [\%]$	SMEFT _{ND}	3.2	0.3	0.2	0.5	0.3
$g_{H\gamma\gamma}^{\text{eff}} [\%]$	SMEFT _{ND}	3.6	1.3	1.3	1.3	1.2
$g_{HZ\gamma}^{\text{eff}} [\%]$	SMEFT _{ND}	11.	9.3	4.6	9.8	9.3
$g_{Hgg}^{\text{eff}} [\%]$	SMEFT _{ND}	2.3	0.9	1.0	1.0	0.8
$g_{Htt}^{\text{eff}} [\%]$	SMEFT _{ND}	3.5	3.1	2.2	3.1	3.1
$g_{Hcc}^{\text{eff}} [\%]$	SMEFT _{ND}	—	2.1	1.8	1.4	1.2
$g_{Hbb}^{\text{eff}} [\%]$	SMEFT _{ND}	5.3	0.6	0.4	0.7	0.6
$g_{H\tau\tau}^{\text{eff}} [\%]$	SMEFT _{ND}	3.4	1.0	0.9	0.7	0.6
$g_{H\mu\mu}^{\text{eff}} [\%]$	SMEFT _{ND}	5.5	4.3	4.1	4.	3.8
$\delta g_{1Z} [\times 10^2]$	SMEFT _{ND}	0.66	0.027	0.013	0.085	0.036
$\delta \kappa_\gamma [\times 10^2]$	SMEFT _{ND}	3.2	0.032	0.044	0.086	0.049
$\lambda_Z [\times 10^2]$	SMEFT _{ND}	3.2	0.022	0.005	0.1	0.051

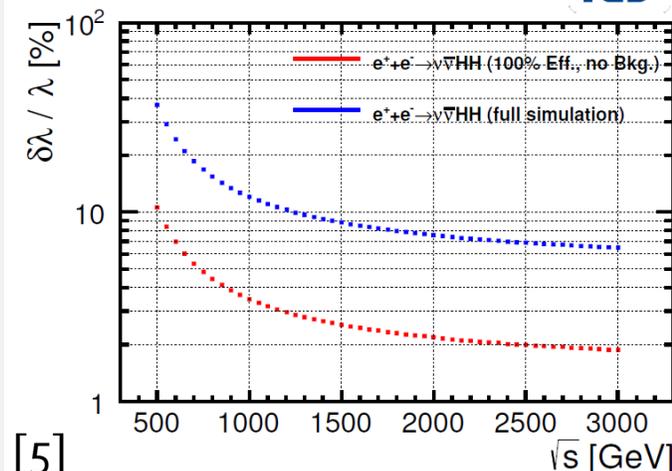
The same holds for CLIC (and FCCee, CEPC)

HIGGS SELF-COUPLING

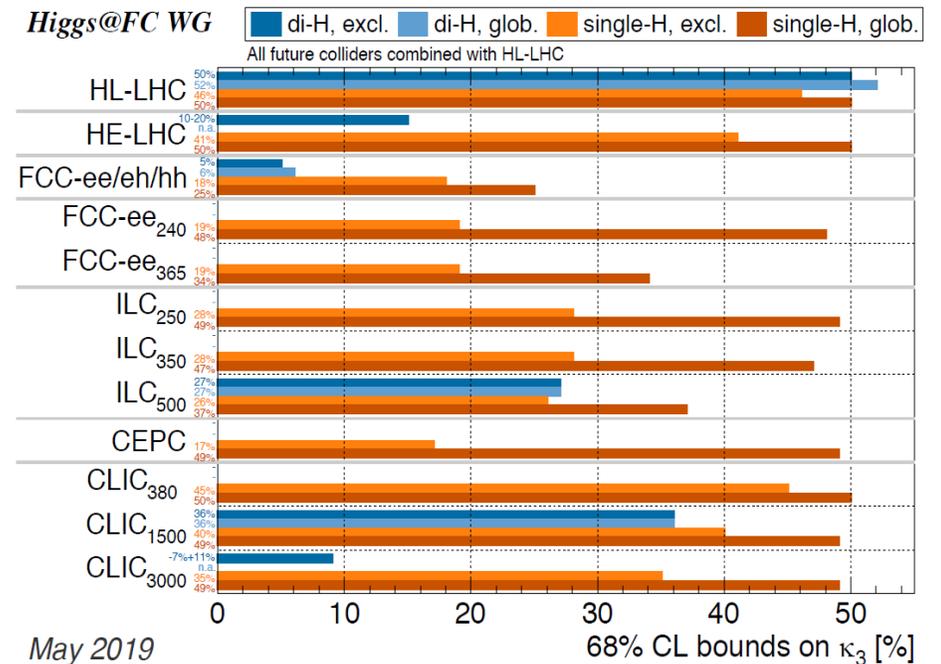
- High energy (>1 TeV) e+e- collider is superior in determination of the Higgs self-coupling
- High energy (double) Higgs production is the most sensitive to deviations of the Higgs self-coupling
- λ is determined from the total rate of HH events (ILD) or template fit of m_{HH} and BDT output (CLICdp)
- Polarization (i.e. -80%) almost doubles the HHvv rate



$$e^+e^- \rightarrow \nu\bar{\nu}HH$$



[5]

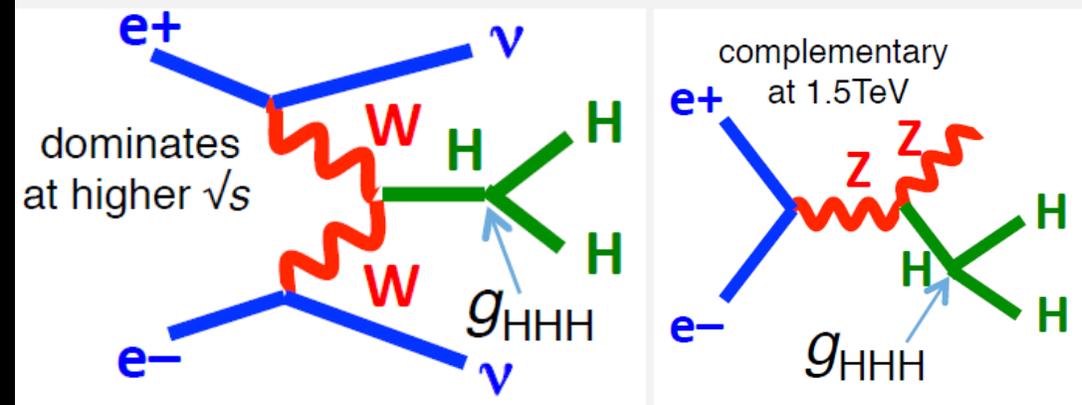


Low energy e+e- colliders (single Higgs production)

in combination with HL-LHC:

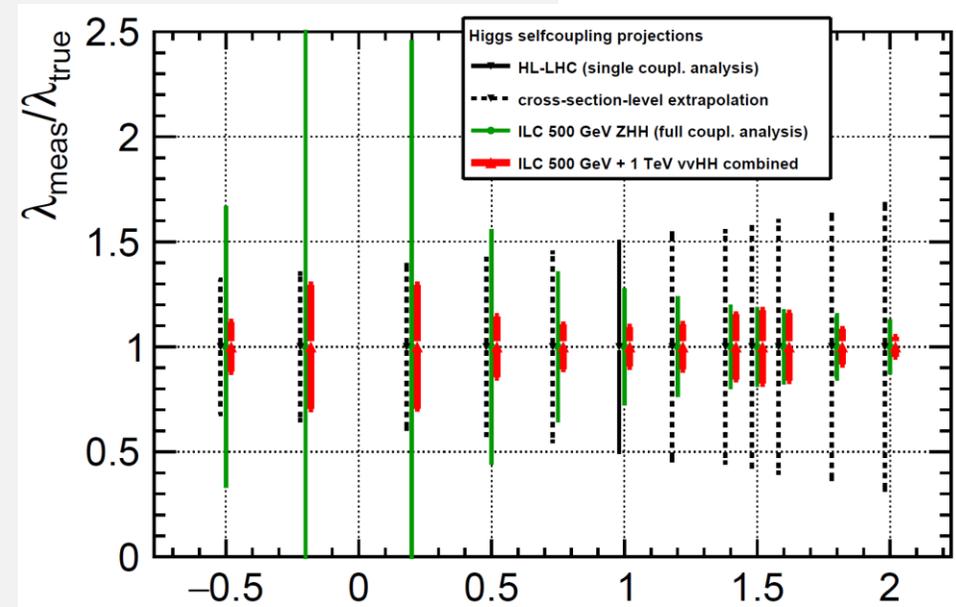
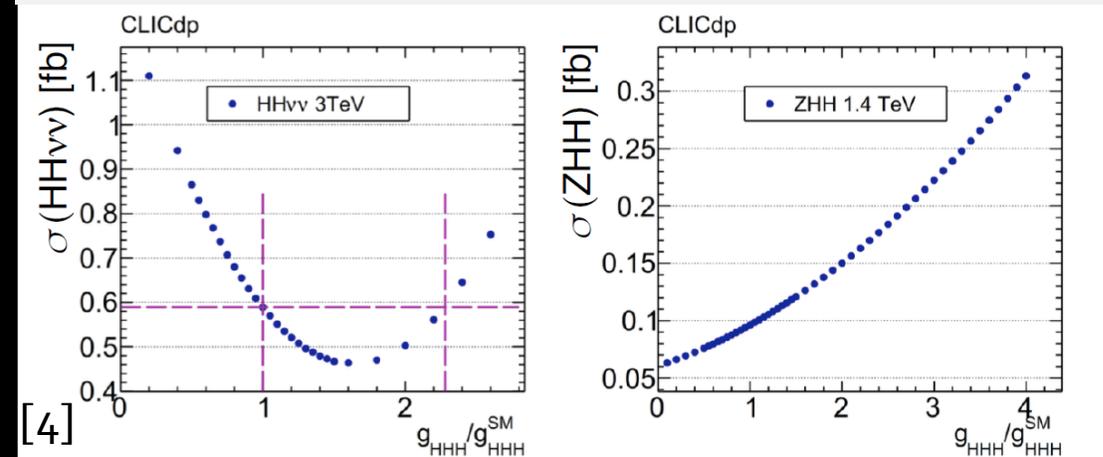
- ILC₂₅₀ and FCCee₃₆₅, $\pm 35\%$
- Double-Higgs production:
 - HL-LHC: $\sim \pm 50\%$
 - ILC₅₀₀ $\sim \pm 27\%$
 - CLIC₃₀₀₀ $\sim \pm 9\%$
 - FCC-hh $\sim \pm 5\%$

LC BENEFITS: STAGING, COMBINATIONS...



[7]	$\Delta\lambda_{hhh}/\lambda_{hhh}$
4 ab ⁻¹ at ILC500	27%
+8 ab ⁻¹ at ILC1000	10%

Higgs self-coupling

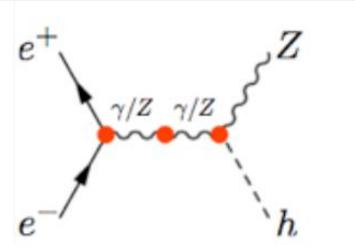


- Intermediate energy (1.4(5) TeV) at CLIC provides complementarity to 3 TeV option with ZHH production
- Different behavior of ZHH and double-Higgs production in WW-fusion, for non-SM values of triple Higgs couplings **resolves ambiguity from interference**
- Statistical uncertainty reduction in combination
- **Clear gain from high center-of-mass energies**

$\lambda_{\text{true}}/\lambda_{\text{SM}}$ [B4]

HIGGS AS A PROBE TO BSM

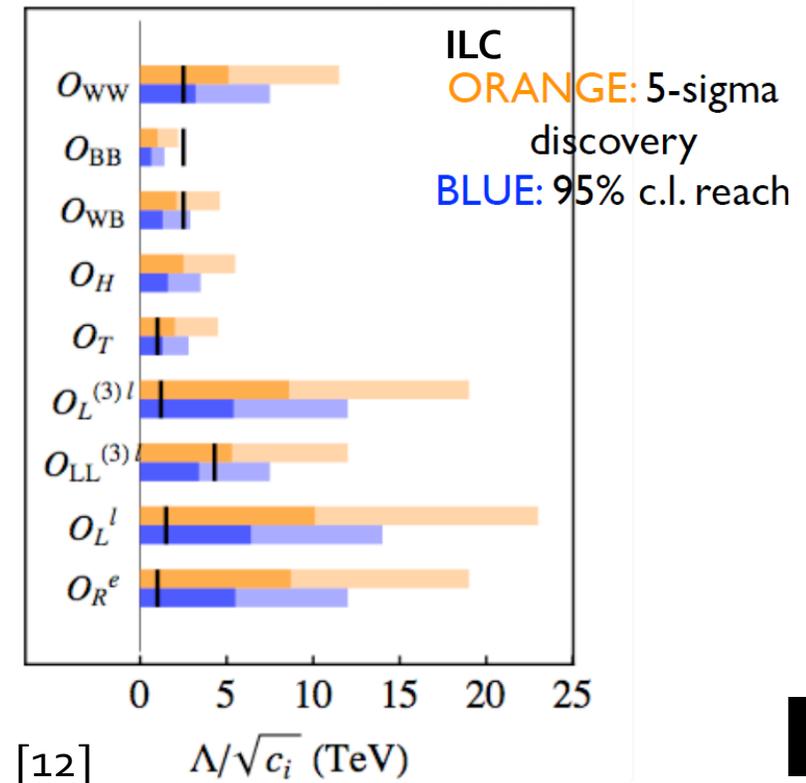
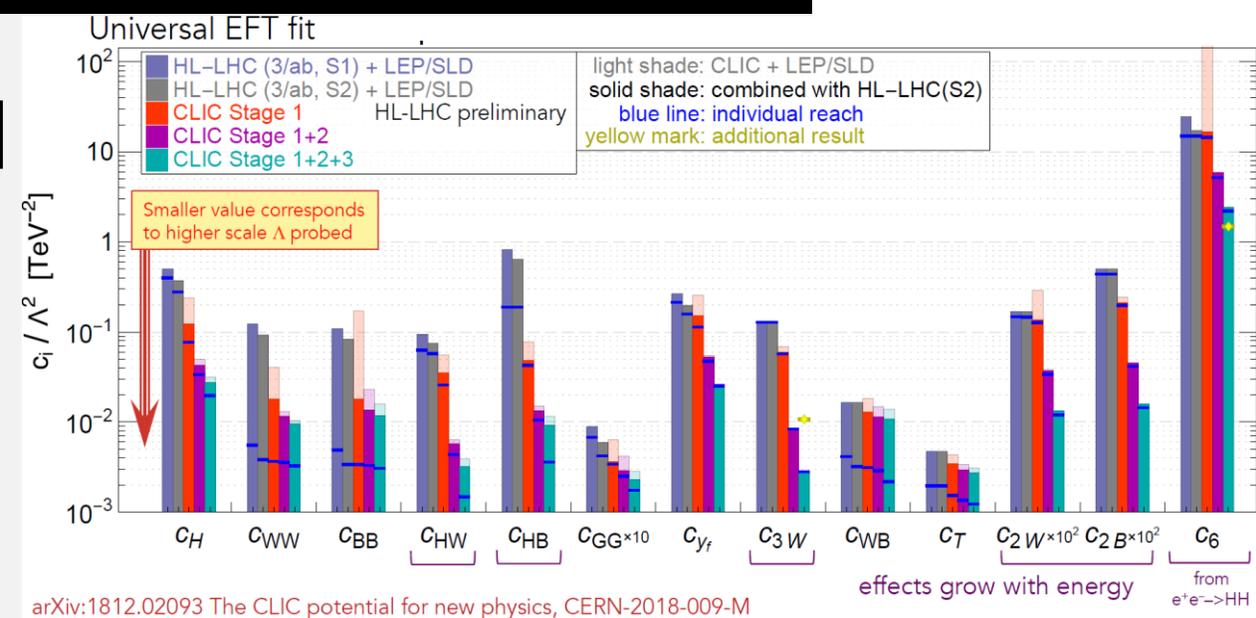
$$\mathcal{L}_{\text{pre-EWSB}} = \sum_i \frac{c_i}{\Lambda^2} \mathcal{O}_i$$



[12] $\delta\sigma/\sigma = 0.5\%/0.1\%$

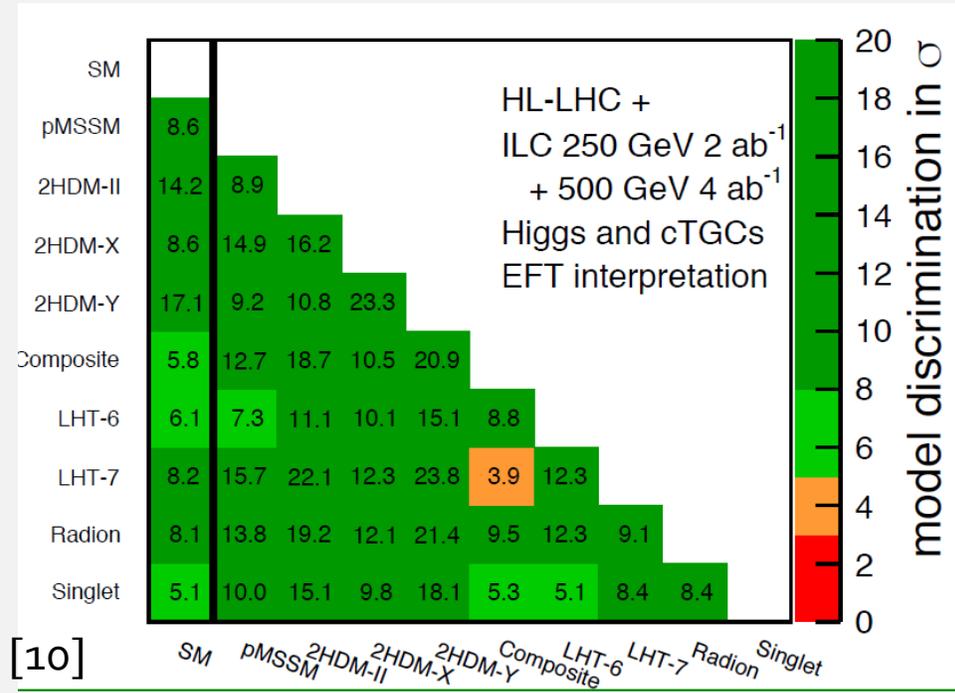
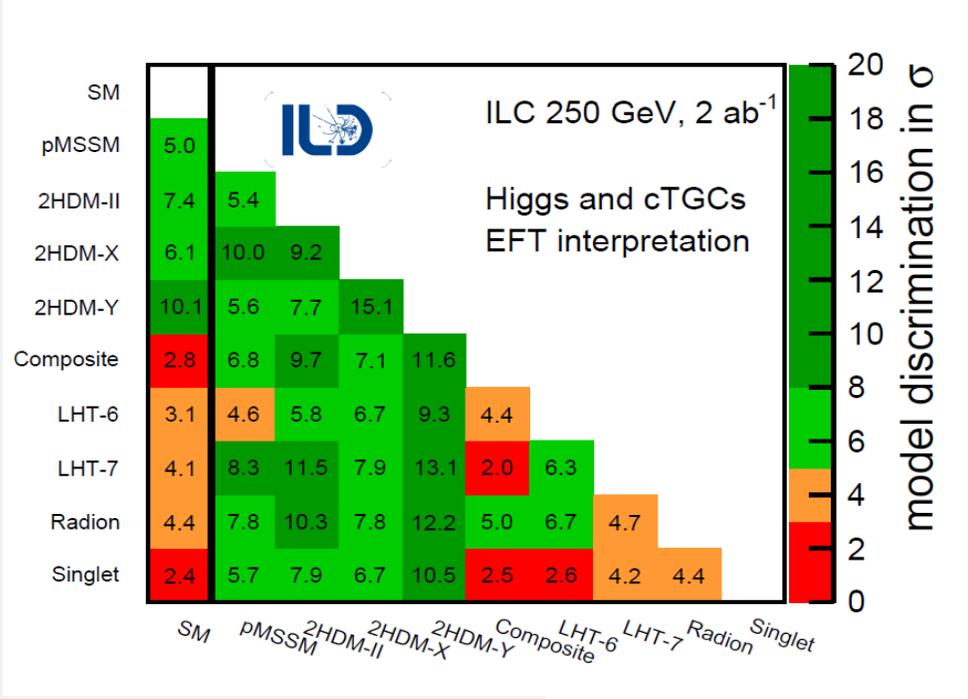
- BSM physics can manifest itself in the Higgs sector in several ways:
 - Contribution from the higher order operators (EFT approach)
 - Higgs compositeness
 - Extended Higgs sector
 - DM portal
 - CPV

High energy Higgs production is the most sensitive to contributions from the 6D operators in the EFT approach and thus can probe the highest New Physics scale Λ



HIGGS AS A PROBE TO BSM – EFT INTERPRETATIONS

Model	$b\bar{b}$	$c\bar{c}$	gg	WW	$\tau\tau$	ZZ	$\gamma\gamma$	$\mu\mu$
1 MSSM [36]	+4.8	-0.8	-0.8	-0.2	+0.4	-0.5	+0.1	+0.3
2 Type II 2HD [35]	+10.1	-0.2	-0.2	0.0	+9.8	0.0	+0.1	+9.8
3 Type X 2HD [35]	-0.2	-0.2	-0.2	0.0	+7.8	0.0	0.0	+7.8
4 Type Y 2HD [35]	+10.1	-0.2	-0.2	0.0	-0.2	0.0	0.1	-0.2
5 Composite Higgs [37]	-6.4	-6.4	-6.4	-2.1	-6.4	-2.1	-2.1	-6.4
6 Little Higgs w. T-parity [38]	0.0	0.0	-6.1	-2.5	0.0	-2.5	-1.5	0.0
7 Little Higgs w. T-parity [39]	-7.8	-4.6	-3.5	-1.5	-7.8	-1.5	-1.0	-7.8
8 Higgs-Radion [40]	-1.5	-1.5	+10.	-1.5	-1.5	-1.5	-1.0	-1.5
9 Higgs Singlet [41]	-3.5	-3.5	-3.5	-3.5	-3.5	-3.5	-3.5	-3.5

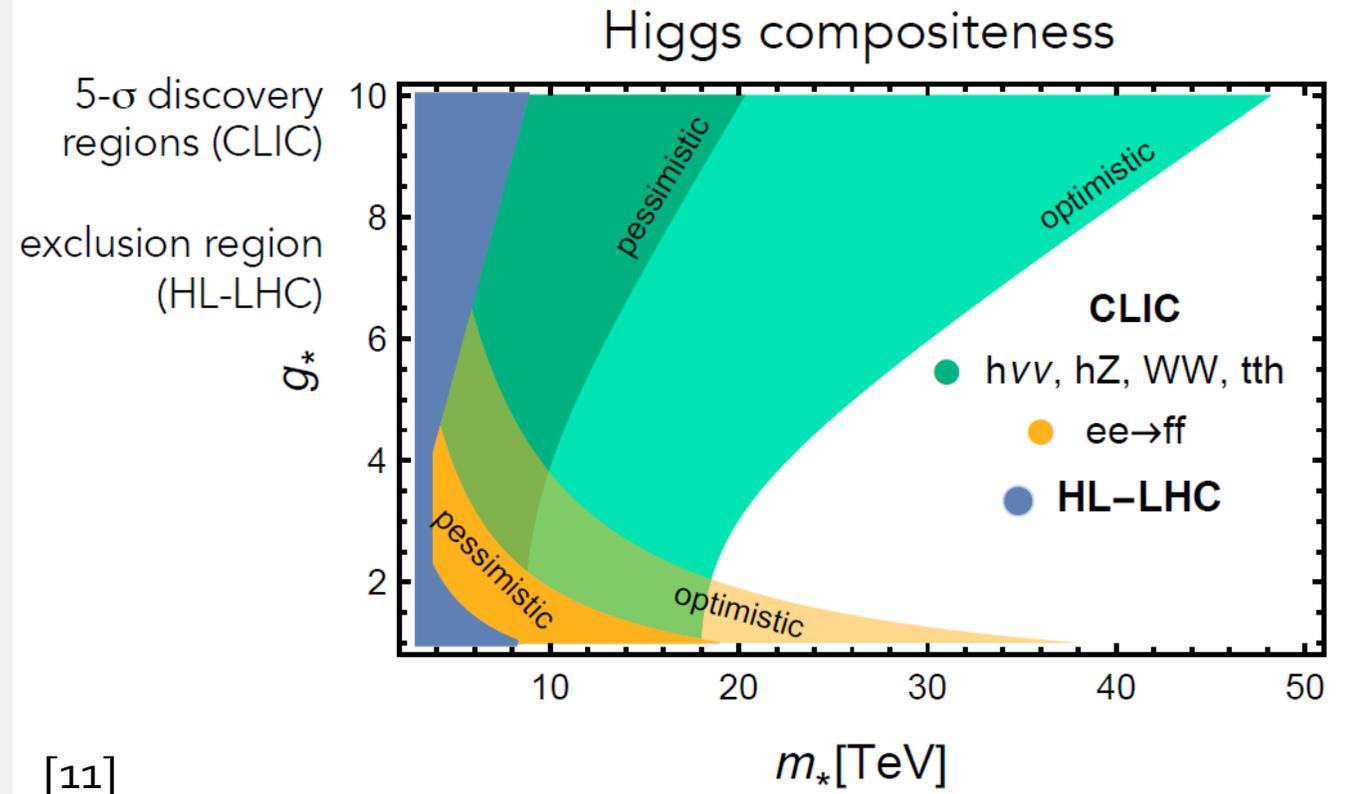


Above 5 σ model discrimination already with 250 GeV ILC

- Substantial improvement at higher energies (linear e+e- colliders):
 @ILC a factor 2 in Higgs couplings precision with 500 GeV polarized beams
 - Complementarity with HL-LHC

HIGGS AS A PROBE TO BSM

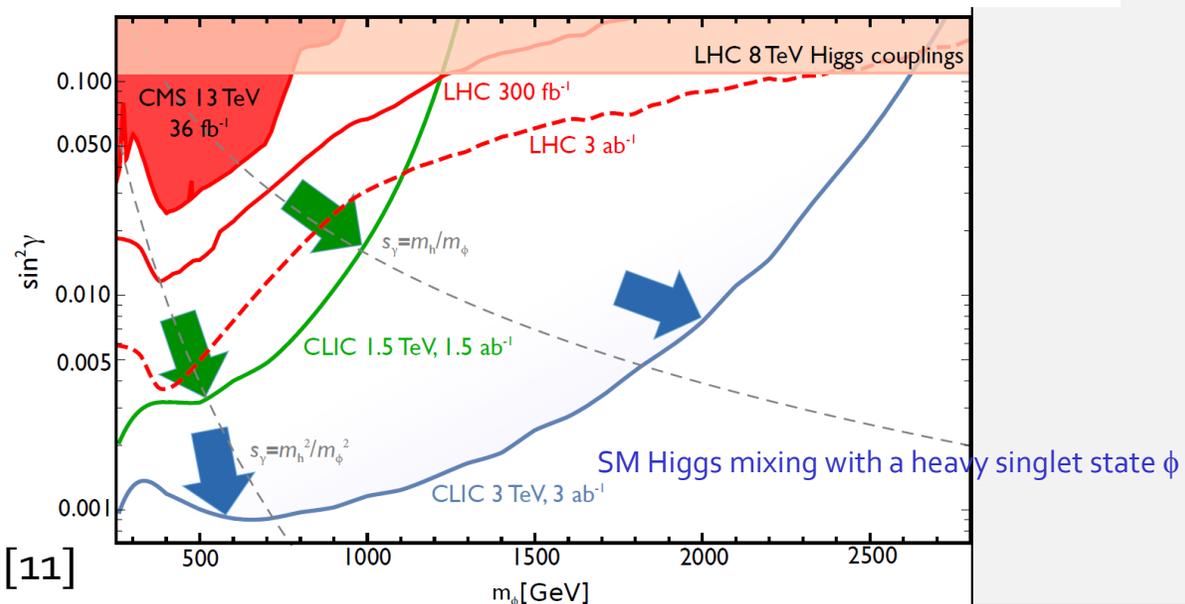
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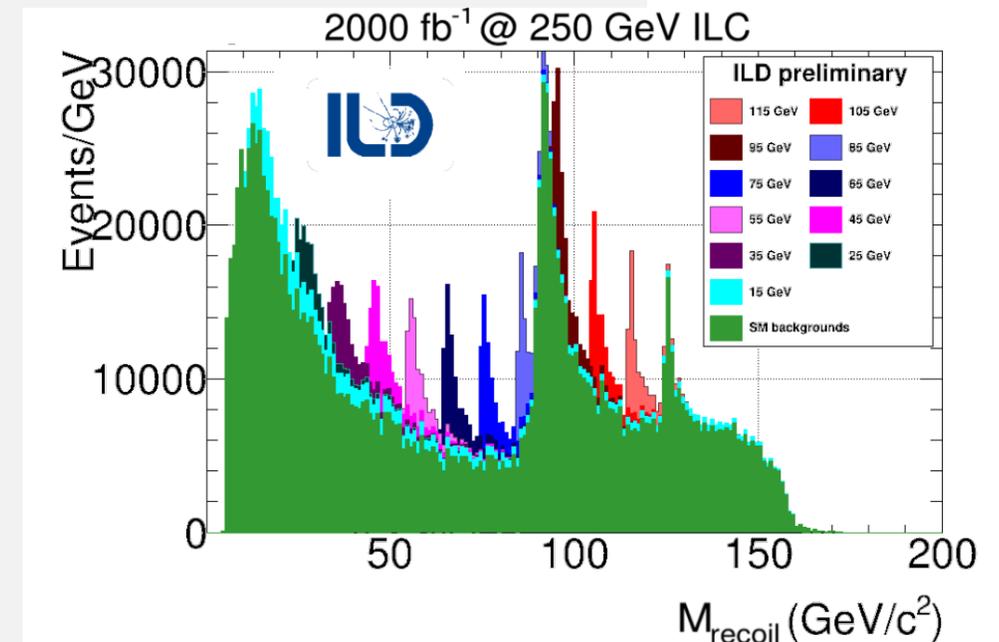
The scale of compositeness can be probed significantly higher from the high-energy collider kinematic limit

HIGGS AS A PROBE TO BSM

- BSM physics can manifest itself in the Higgs sector in several ways:
 - Contribution from the higher order operators (EFT approach)
 - Higgs compositeness
 - Extended Higgs sector
 - DM portal
 - CPV



- In majority of BSM models, SM Higgs comes with additional Higgses (2HDM, SUSY in general, compositeness,..etc.)
- Can be a lighter scalar than SM Higgs – it is important to be capable of probing such states at future colliders
- If SM Higgs is the lightest, other states are nearly mass-degenerated

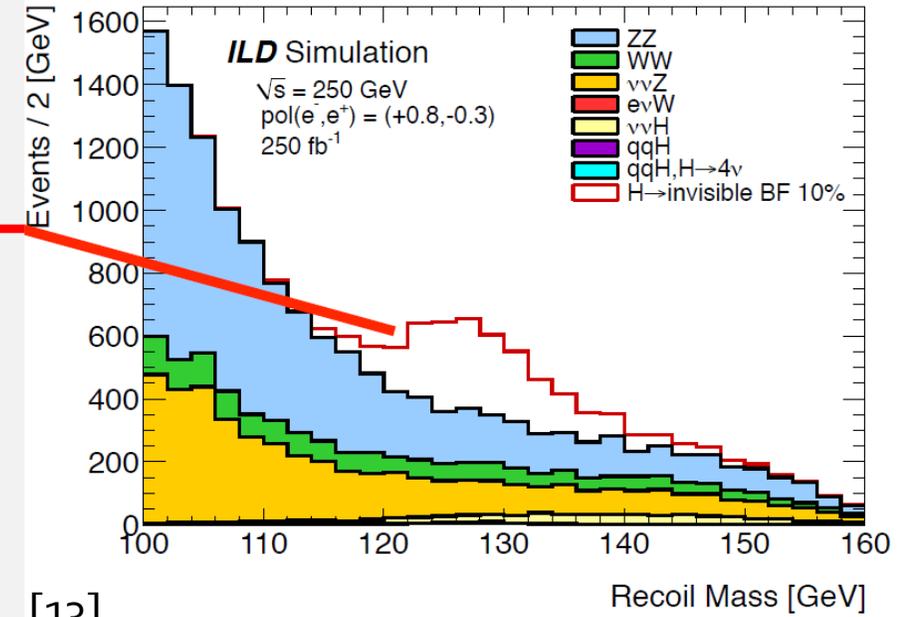


[5] $e^+e^- \rightarrow Z S^0 \rightarrow \mu^+\mu^- + inv$

HIGGS TO INVISIBLE

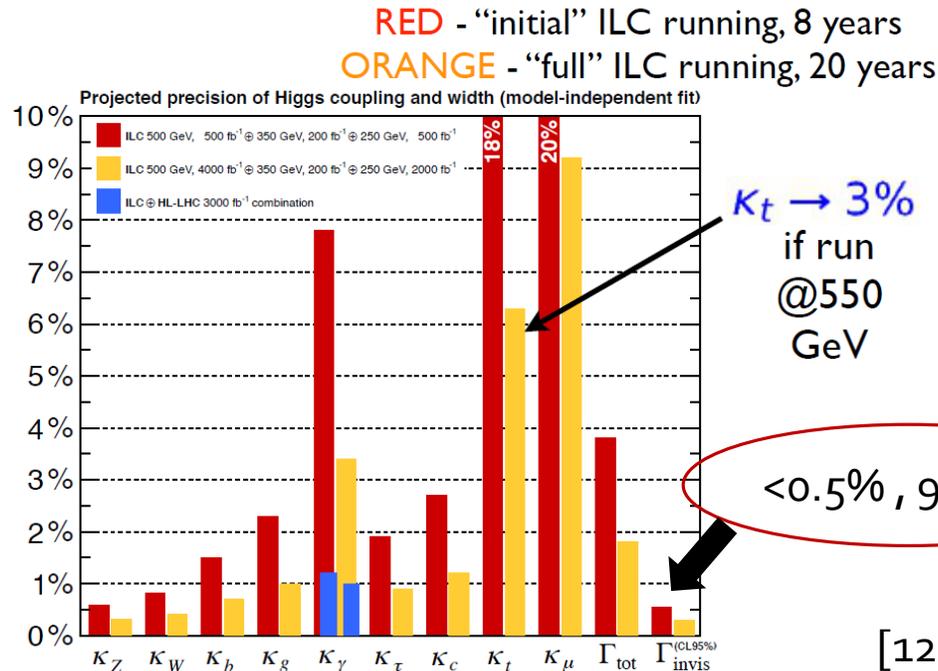
- BSM physics can manifest itself in the Higgs sector in several ways:
 - Contribution from the higher order operators (EFT approach)
 - Higgs compositeness
 - Extended Higgs sector
 - DM portal
 - CPV

$H \rightarrow inv.$



[13]

- Looking at the recoil mass under the condition that nothing observable is recoiling against the Z boson (only one Z per event)
- Access to DM connected to SM particles through a specific set of operators (portals)



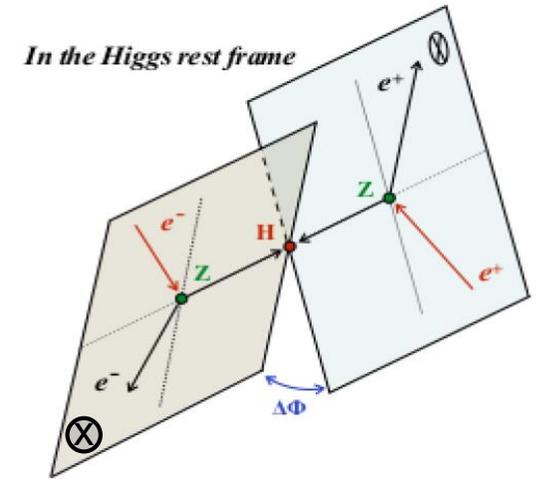
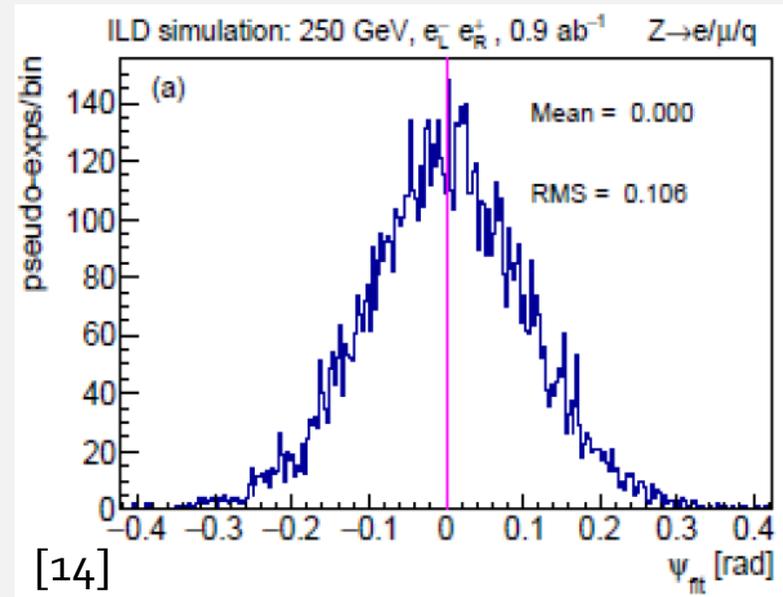
[12]

$$\frac{1}{2} \epsilon_Y F_{\mu\nu}^Y F'^{\mu\nu} \quad \epsilon_H |H|^2 |\Phi|^2 \quad \epsilon_a \frac{a}{f_a} F_{\mu\nu} \tilde{F}^{\mu\nu}$$

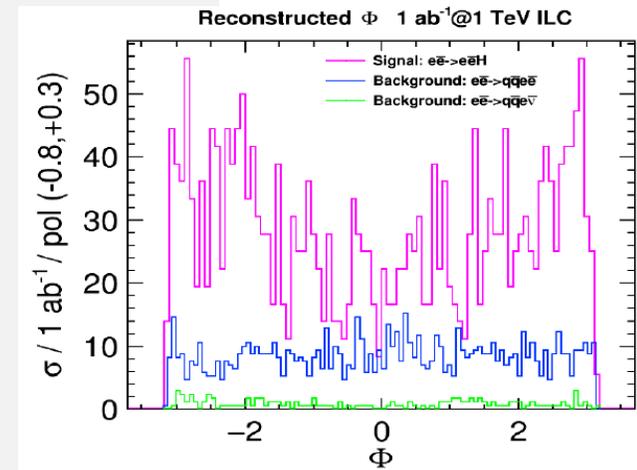
CPV IN THE HIGGS SECTOR

- More difficult than just a spin/parity determination: Higgs can be a mixture of different CP eigenstates

$$h = H \cdot \cos \psi + A \cdot \sin \psi$$
- Can be measured in Hff and HVV vertices, both in Higgs production and decays
- Hff (HVV) sensitive to CPV contributions at the tree (loop) level
- Only lose bounds (at present) on a quantum superposition of different CP states, while experimentally disfavored hypothesis on purely CP odd state



Collider	ψ_{CP}
HL-LHC	8°
HE-LHC	—
CEPC	—
FCC-ee ₂₄₀	10°
ILC ₂₅₀	4°



ILC₂₅₀ – benefit from polarization & combination ($H\tau\tau$)
 1 TeV – optimal for ZZ-fusion

LINEAR VS. CIRCULAR

- Precision wise, linear and circular colliders' precision is comparable when it comes to the Higgs couplings
- Due to high-energy access of high cross-section Higgs production mechanisms, LCs are superior in probing of the Higgs self-coupling
- Extensibility of the physics span (pp collisions, 100 TeV center-of mass energy) is a great advantage of circular colliders
- But, + a 100 TeV hadron-collider, comes at the moment with quite a few open issues:
 - Accelerator & detector technologies
 - Huge pile-up
 - Systematics control and theoretical uncertainties

[B. Heinemann '19]

Theoretical Uncertainties: production

Production at hadron colliders

- For HL-LHC uncertainties expected to be improved by factor 2 w.r.t. current
- HE-LHC: another factor of 2
- FCC-hh: well below 1%

Requires e.g.

- Improved PDFs
- Higher precision calculations
- Improved non-perturbative aspects
- ...

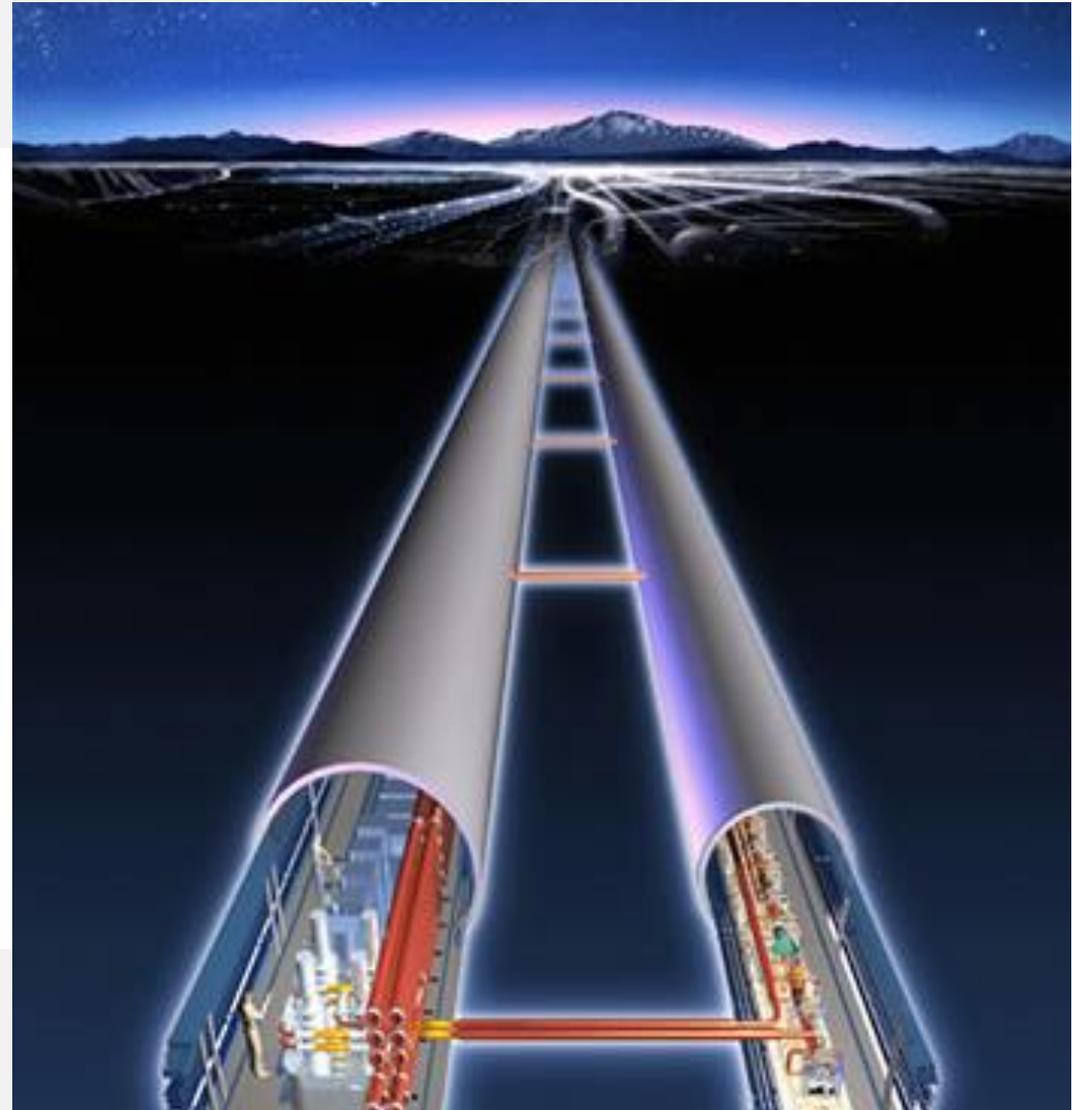
Note: this is related to the fact that FCC-hh is assumed to be realised only far in the future!

- Precision measurement of couplings at hadron colliders are limited by the systematic (theoretical) uncertainties
- This is also a reason for the fact that the Higgs coupling projections for HE-LHC show only relatively small improvements over HL-LHC
- FCC-hh projections, in particular when taken separately, depend on a drastic reduction of theory uncertainties [B2].

SUMMARY

- All future e+e- projects bring significant added value to the projected HL-LHC sensitivities in the Higgs sector...
- ... enabling discrimination of BSM models inaccessible at HL-LHC
- Already lowest energy phases brings sensitivity far beyond the projected HL-LHC precision
- Higher center of mass energies significantly extends physics span of a LC (Higgs self-coupling, BSM scenarios)
➡ upgrade is important – genuine advantage of a LC
- Additional enhancement from polarization (precision, model discrimination)

READY-TO-WEAR PROJECTS



EISA
European Institute for Sciences and Their Applications



Corfu Summer Institute
Workshop on the Standard
Model and Beyond
29.08.-08.09. 2021

THANK YOU

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Workshop on the Standard
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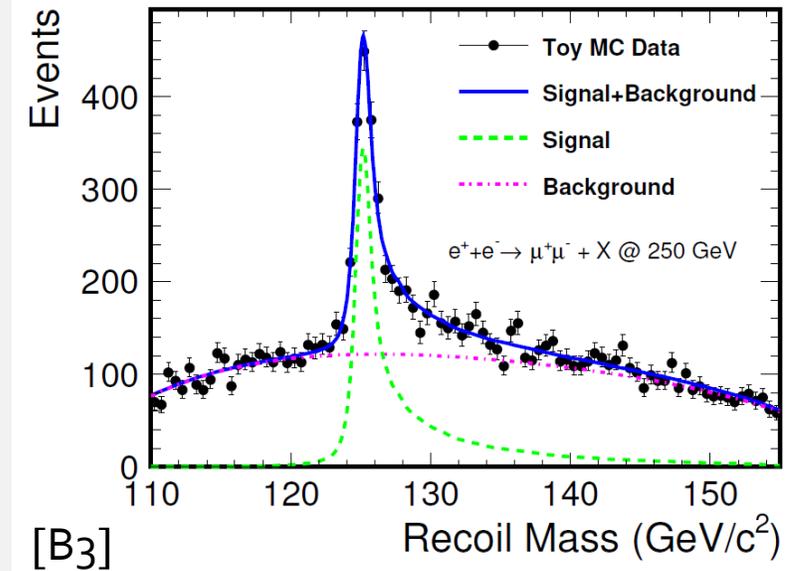
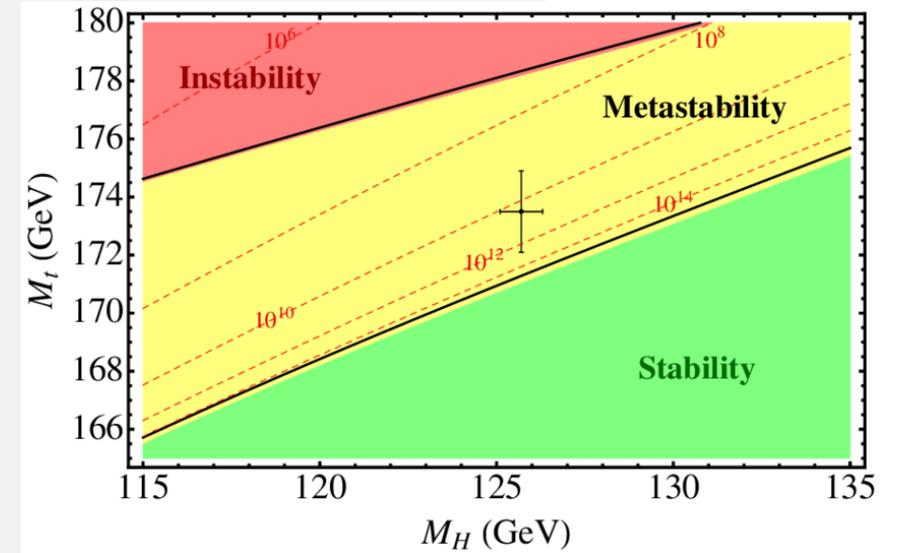
BACK UP

HIGGS MASS

- Which precision of the Higgs mass is needed?
 - Vacuum stability (at least several GeV)
 - Impact on $H \rightarrow ZZ^*$ width (a few tens of MeV)
- Current precision 160 MeV
- Comparable precision with HL-LHC

Collider Scenario	Strategy	δm_H (MeV)	$\delta(\Gamma_{ZZ^*})$ (%)
LHC Run-2	$m(ZZ), m(\gamma\gamma)$	160	1.9
HL-LHC	$m(ZZ)$	10-20	0.12-0.24
ILC ₂₅₀	ZH recoil	14	0.17
CLIC ₃₈₀	ZH recoil	78	1.3
CLIC ₁₅₀₀	$m(bb)$ in $H\nu\nu$	30 ¹⁵	0.56
CLIC ₃₀₀₀	$m(bb)$ in $H\nu\nu$	23	0.53
FCC-ee	ZH recoil	11	0.13
CEPC	ZH recoil	5.9	0.07

M. Cepeda, Higgs precision measurements at future colliders, IFT UAM-CSIC, Madrid, Spain, July 2019 and [2]



HIGGS WIDTH

- Being less than 5 MeV, Higgs decay width can not be *directly* measured at any proposed e+e-collider
- Can be determined from individual decays (quasi-direct measurement), i.e. $H \rightarrow WW$ decays in WW-fusion, $H \rightarrow ZZ$ in HZ)

$$\sigma(ee \rightarrow ZH) \cdot \text{BR}(H \rightarrow ZZ) \propto \frac{g_{HZ}^4}{\Gamma}$$

- In a combination of measurements:

$$\frac{\sigma(ee \rightarrow ZH) \cdot \text{BR}(H \rightarrow WW) \cdot \sigma(ee \rightarrow ZH) \cdot \text{BR}(H \rightarrow bb)}{\sigma(ee \rightarrow \nu\nu H) \cdot \text{BR}(H \rightarrow bb)}$$

$$\propto \frac{g_{HZ}^2 \cdot g_{HW}^2}{\Gamma} \cdot \frac{g_{HZ}^2 \cdot g_{Hb}^2}{\cancel{X}} \cdot \frac{\cancel{X}}{g_{HW}^2 \cdot g_{Hb}^2} = \frac{g_{HZ}^4}{\Gamma}$$

- The ultimate precision is reached in a global fit, (model-independent or in the LHC-style, so called κ -framework):

$$\Gamma_H = \frac{\Gamma_H^{\text{SM}} \cdot \kappa_H^2}{1 - (\text{BR}_{inv} + \text{BR}_{unt})}$$

- Or in a global (model-dependent) EFT fit (assumes the new physics scale $\Lambda \gg M_H$)

Statistical accuracy of 1-2%

Collider	$\delta\Gamma_H$ (%) from Ref.	Extraction technique standalone result	$\delta\Gamma_H$ (%) kappa-3 fit
ILC ₂₅₀	2.4	EFT fit [3]	2.4
ILC ₅₀₀	1.6	EFT fit [3, 11]	1.1
CLIC ₃₅₀	4.7	κ -framework [85]	2.6
CLIC ₁₅₀₀	2.6	κ -framework [85]	1.7
CLIC ₃₀₀₀	2.5	κ -framework [85]	1.6
CEPC	3.1	$\sigma(ZH, \nu\bar{\nu}H), \text{BR}(H \rightarrow Z, b\bar{b}, WW)$ [90]	1.8
FCC-ee ₂₄₀	2.7	κ -framework [1]	1.9
FCC-ee ₃₆₅	1.3	κ -framework [1]	1.2

[arXiv:1905.03764](https://arxiv.org/abs/1905.03764)

ILC and CLIC parameters

Property	unit	ILC at 500 GeV	ILC at 1 TeV	CLIC at 380 GeV	CLIC at 3 TeV
L	$\text{cm}^{-2}\text{s}^{-1}$	$1.8 \cdot 10^{34}$	$3.5 \cdot 10^{34}$	$1.5 \cdot 10^{34}$	$5.9 \cdot 10^{34}$
$L_{0.01}$	$\text{cm}^{-2}\text{s}^{-1}$	$1.0 \cdot 10^{34}$	$1.2 \cdot 10^{34}$	$0.9 \cdot 10^{34}$	$2.0 \cdot 10^{34}$
$L_{0.01}/L$	%	58	59	60	34
Repetition rate	Hz	5 Hz	4 Hz	50 Hz	50 Hz
Train duration	ns	727 μs	897 μs	178 ns	156 ns
BX / train		1312	2450	356	312
Bunch separation	ns	554 ns	366 ns	0.5 ns	0.5 ns
Duty cycle	%	0.36	0.36	0.00089	0.00078
σ_x/σ_y	nm	474/5.9	481/2.8	$\sim 150/3$	$\sim 40/1$
σ_z	μm	300	250	70	44