

# The supersymmetric anti-D3-brane action in KKLT

**Niccolò Cribiori**



TECHNISCHE  
UNIVERSITÄT  
WIEN

Technische Universität Wien

Conference on Recent Developments in Strings and Gravity  
Corfu, September 15th, 2019

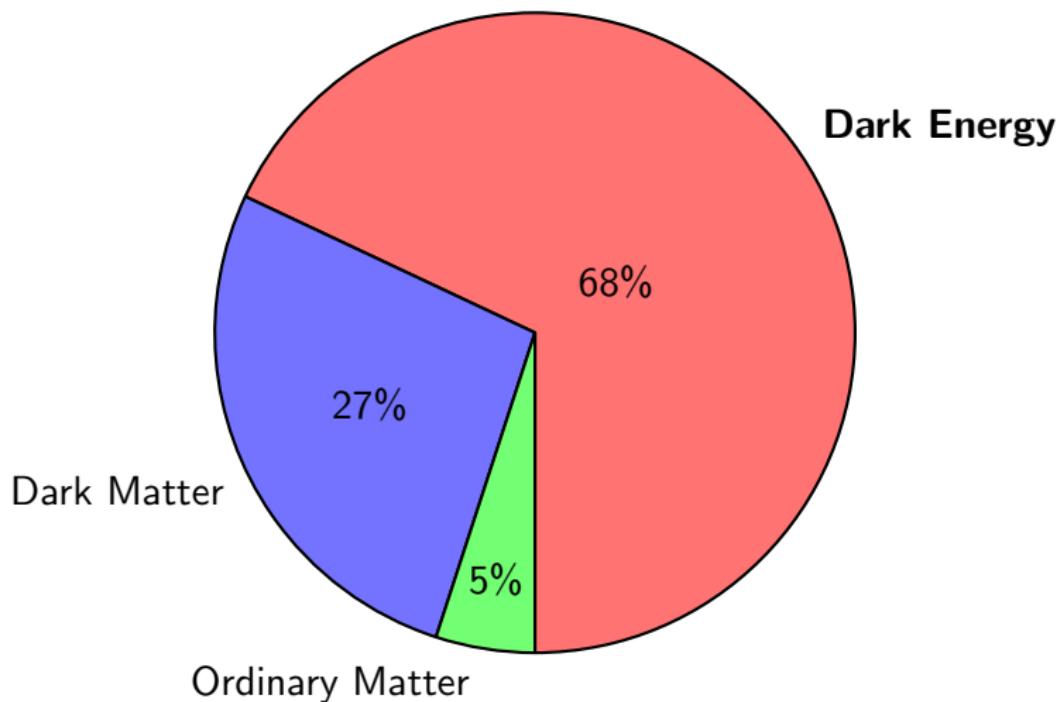
Based on 1906.07727, with C. Roupec, T. Wräse and Y. Yamada

# Introduction

# Our observed Universe

[Planck '18]

## $\Lambda$ CDM Cosmology



# What do we know about Dark Energy?

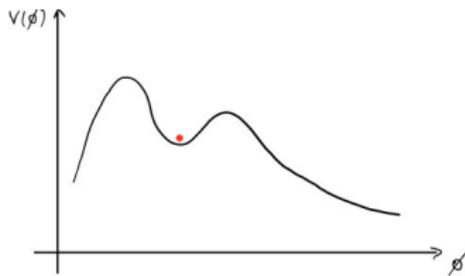
- ▶ 68% of the total energy of observable universe
- ▶ Measured energy density  $\mathcal{V} \sim 10^{-122} M_P^4 > 0$  [Planck '18]
- ▶ Assuming FRW universe and equation of state  $P = w\rho$ ,  
**accelerated expansion** occurs for

$$\frac{\ddot{a}}{a} = -\frac{1}{6}(\rho + 3P) = -\frac{1}{6}\rho(1 + 3w) > 0 \iff w < -\frac{1}{3}$$

Measured value  $w = -1.03 \pm 0.03$  [Planck '18]

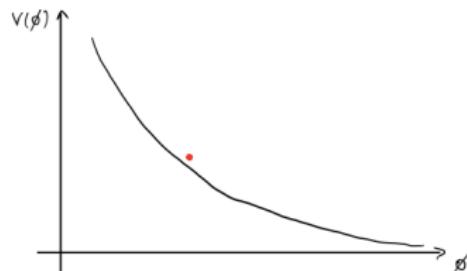
# How to explain Dark Energy?

## Accelerated expansion: possible explanations



constant vacuum energy  
de Sitter vacuum

$$w = -1$$



slowly changing in time  
quintessence

$$w > -1$$

At present they are both compatible with experiments

$$w = -1.03 \pm 0.03$$

[Planck '18]

## de Sitter vacuum

- ▶ Dark energy = positive cosmological constant

$$R_{\mu\nu} - \frac{1}{2}g_{\mu\nu}(R + 2\Lambda) = 0$$

- ▶ Understand it microscopically → string theory

$$\Lambda \sim \langle \mathcal{V}(\phi) \rangle$$

- ▶ Problems: moduli stabilization, corrections, ...

Moreover

- ▶ de Sitter breaks supersymmetry  
[Pilch, van Nieuwenhuizen, Sohnius '85]

$$\mathcal{V}_{\text{SUGRA}} = (\delta(\text{spin-1/2}))^2 - (\delta(\text{spin-3/2}))^2 > 0$$

## De Sitter and string theory

## Some bad news

- ▶ It is not straightforward to obtain dS from string theory.
- ▶ Is the difficulty accidental or is manifestation of some principle yet to be understood?
- ▶ It has been even conjectured that there is no dS!  
[Danielsson, Van Riet '18; Obied, Ooguri, Spodyneiko, Vafa '18; Ooguri, Palti, Shiu, Vafa '18]
- ▶ No-go theorems  
[Maldacena, Núñez '00; Hertzberg, Kachru, Taylor, Tegmark '07; Flauger, Paban, Robbins, Wrase '08; Wrase, Zagermann '10; Junghans, Zagermann '16,...]

## Some good news

- ▶ **Moduli stabilization:** compactifications with fluxes. Usually 10d SUGRA is trustable at large volume  $L/l_s \gg 1$  and weak coupling  $e^\phi \ll 1$
- ▶ **Supersymmetry breaking:** extended objects (D-branes)

$$\text{Brane susy-breaking} \quad \mathcal{V} \sim e^{\gamma\phi}$$

[Sugimoto '99; Antoniadis, Dudas, Sagnotti '99; Aldazabal, Uranga '99; Angelantonj, Antoniadis, D'Appollonio, Dudas, Sagnotti '00]

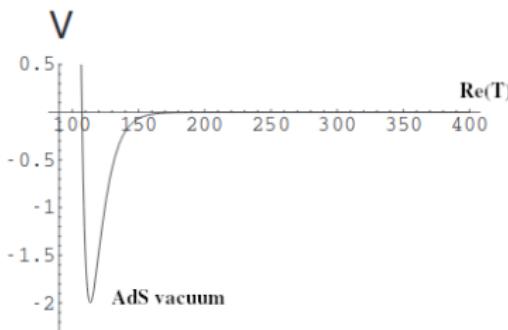
- ▶ The first proposal for a dS construction in string theory with massless fields and tunable  $\Lambda = \mathcal{V}_{min}$  has been given by [Kachru, Kallosh, Linde, Trivedi '03]

# KKLT - moduli stabilization

[Kachru, Kallosh, Linde, Trivedi '03;  
Kachru, Kallosh, Linde, Maldacena, McAllister, Trivedi '03]

**Setup:** IIB compactified on CY with orientifold

1. Stabilize CS moduli and axio-dilaton using fluxes  
[Giddings, Kachru, Polchinski '01]
2. Stabilize Kähler moduli using non-perturbative effects



1.  $K = -3\log(T + \bar{T})$
2.  $W = W_0 + Ae^{-aT}$

$$\mathcal{V}_{KKLT} = e^K (|\nabla_T W|^2 - 3|W|^2)$$

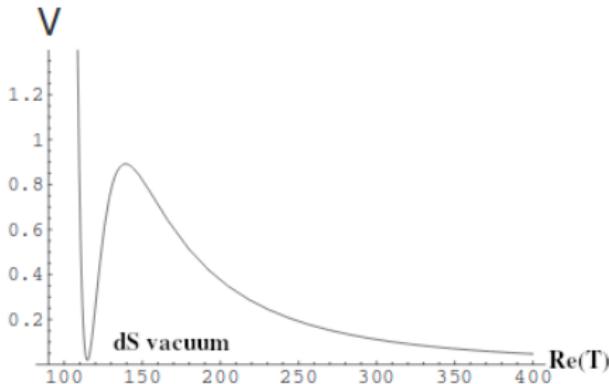
**Stable and supersymmetric AdS vacuum**

# KKLT - supersymmetry breaking and uplift

[Kachru, Kallosh, Linde, Trivedi '03;  
Kachru, Kallosh, Linde, Maldacena, McAllister, Trivedi '03]

- ▶ Add  $\overline{D3}$  brane to break supersymmetry and uplift the vacuum energy

$$\mathcal{V} = \mathcal{V}_{KKLT} + \frac{\mu^4}{(T + \bar{T})^2} \quad \text{metastable dS}$$



## Anti-D3-brane in KKLT

- ▶ Most delicate step in the procedure.
- ▶ It breaks supersymmetry **spontaneously** → **goldstino**
- ▶ On the world volume fields susy is non-linear  
[Kallosh, Wrase '14; Bergshoeff, Dasgupta, Kallosh, Van Proeyen, Wrase '15; Kallosh, Quevedo, Uranga '15; Bandos, Heller, Kuzenko, Martucci, Sorokin '15; García-Etxebarria, Quevedo, Valandro '15; Vercnocke, Wrase '16; Kallosh, Vercnocke, Wrase '16]
- ▶ Can we still use the standard SUSY language?
- ▶ It would help in understanding better (this step of) KKLT.

## The low energy description

# What is the effective description?

**Setup:** 4d  $\mathcal{N} = 1$  supergravity + **goldstino**

**First step: scalar potential** [Ferrara, Kallosh, Linde '14]

- ▶ Add a **goldstino** chiral multiplet  $X$ , such that  $X^2 = 0$

$$K = -3\log(T + \bar{T} - X\bar{X}) \quad W = W_0 + Ae^{-aT} + \mu^2 X$$

$$\mathcal{V} = \mathcal{V}_{KKLT} + \frac{\mu^4}{(T + \bar{T})^2}$$

**However**

- ▶ Full action? [Garcia del Moral, Parameswaran, Quiroz, Zavala '17]

# Supersymmetry on the brane

## ► D3-brane: linear SUSY

$$\begin{array}{ll} \text{3 scalars } H^a \longleftrightarrow \text{3 fermions } \chi^a & \text{chiral multiplet } \{H^a, \chi^a\} \\ \text{1 fermion } \lambda \longleftrightarrow \text{1 vector } v_\mu & \text{vector multiplet } \{\lambda, v_\mu\} \end{array}$$

## ► anti-D3-brane: non-linear SUSY      $\lambda \sim \text{goldstino}$

$$\begin{array}{ll} H^a \longleftrightarrow \partial H^a \lambda + \dots & \chi^a \longleftrightarrow \chi^a \lambda + \lambda^2 + \dots \\ \lambda \longleftrightarrow \lambda^2 + \dots & v_\mu \longleftrightarrow \lambda \not{v} + \dots \end{array}$$

- More complicated transformations than linear SUSY
- All of them involve the goldstino
- How to deal with this?

# Non-linear supersymmetry

[Volkov, Akulov '73;

Rocek '78; Lindstrom, Rocek '79; Casalbuoni, De Curtis, Dominici, Feruglio, Gatto '89]

- ▶ Recent progress in non-linear supersymmetry.  
[Antoniadis, Bergshoeff, Dall'Agata, Derendinger, Dudas, Farakos, Ferrara, Freedman, Kallosh, Kehagias, Komargodski, Kuzenko, Markou, Poratti, Sagnotti, Seiberg, Sorokin, Tartaglino-Mazzucchelli, Tournoy, Van Proeyen, Wrane, Yamada, Zwirner,...]
- ▶ General recipe to describe ANY low energy spectrum.  
[Dall'Agata, Dudas, Farakos '16; NC, Dall'Agata, Farakos '17]
- ▶ Apply it to de Sitter and KKLT

# The strategy

[NC, Roupec, Wräse, Yamada '19]

- ▶ Embed all worldvolume fields in **constrained superfields**

worldvolume field	linear SUSY	non-linear SUSY
1 goldstino	$X$	$X^2 = 0$
3 massive fermions	$Y^i$	$XY^i = 0$
3 complex scalars	$H^a$	$X\bar{H}^a = \text{chiral}$
1 vector	$W_\alpha$	$XW_\alpha = 0$

- ▶ Couple them to closed string fields in **4d  $\mathcal{N} = 1$  supergravity**
  1. Kähler potential  $K$
  2. superpotential  $W$
  3. gauge kin. function  $f(\tau)$
- ▶ Match with anti-D3-brane action

[Martucci, Rosseel, Van den Bleeken, Van Proeyen '05; Bergshoeff, Kallosh, Kashani-Poor, Sorokin, Tomasiello '05]
- ▶ Non-trivial restriction: modular invariance

# A particular problem

[NC, Roupec, Wrášek, Yamada '19]

D3 action       $\xleftarrow{\text{RR sign flip}}$        $\overline{D3}$  action

- ▶ Therefore  $f(\tau)$        $\xleftarrow{\text{RR sign flip}}$        $f(\bar{\tau})$       **antiholom.**  
**NOT compatible with susy**
- ▶ Solution: **non-linear supersymmetry**

$$f_{\overline{D3}} = \bar{D}^2 \left( \frac{\bar{X} f(\bar{\tau})}{\bar{D}^2 \bar{X}} \right) = f(\bar{\tau}) + \text{goldstino} \quad \text{chiral!}$$

- ▶ Inspired by new D-term in supergravity  
[NC, Farakos, Tournay, Van Proeyen '17]
- ▶ A similar problem happens for the fermionic mass

# The supersymmetric anti-D3-brane action in KKLT

[NC, Roupec, Wräse, Yamada '19]

$$\begin{aligned} K &= -\log(-i(\tau - \bar{\tau})) - 3 \log \left[ (T + \bar{T}) \left( -i \int \Omega \wedge \bar{\Omega} \right)^{\frac{1}{3}} + k(H^a, \bar{H}^a) \right] \\ &\quad - 3 \log \left( 1 - C_1 X \bar{X} - C_2 \delta_{i\bar{j}} Y^i \bar{Y}^{\bar{j}} \right) \\ W &= \int G_3 \wedge \Omega + A e^{-aT} + \mu^2 X + \frac{1}{2} M_{ij} Y^i Y^j \\ f_{\overline{D3}} &= \bar{D}^2 \left( \frac{\bar{X} f(\bar{\tau})}{\bar{D}^2 \bar{X}} \right), \quad M_{ij} = \bar{D}^2 \left( \frac{\bar{X} e^{-4A-2u} m_{ij}}{\bar{D}^2 \bar{X} (-i(\tau - \bar{\tau}))^{\frac{1}{2}} (-i \int \Omega \wedge \bar{\Omega})^{-\frac{1}{6}}} \right) \\ \mathcal{V}_{\overline{D3}} &= 2 e^{4A(H^a, \bar{H}^a) - 8u}, \end{aligned}$$

where

$$\begin{aligned} e^{4u} &= (T + \bar{T}) + k(H^a, \bar{H}^a) \left( -i \int \Omega \wedge \bar{\Omega} \right)^{-\frac{1}{3}} = (V_6)^{\frac{2}{3}} \\ C_1 &= \frac{e^{-4A(H^a, \bar{H}^a)}}{3(-i(\tau - \bar{\tau}))(T + \bar{T} + k(H^a, \bar{H}^a) (-i \int \Omega \wedge \bar{\Omega})^{-\frac{1}{3}}) (-i \int \Omega \wedge \bar{\Omega})} \\ C_2 &= \frac{e^{-4A(H^a, \bar{H}^a)}}{3(-i(\tau - \bar{\tau}))(T + \bar{T} + k(H^a, \bar{H}^a) (-i \int \Omega \wedge \bar{\Omega})^{-\frac{1}{3}})^2 (-i \int \Omega \wedge \bar{\Omega})^{\frac{1}{3}}} \end{aligned}$$

# Conclusion

- ▶ A possible explanation for Dark Energy is de Sitter vacuum
- ▶ Microscopic origin: string theory and brane susy breaking
- ▶ Particular setup:  $\overline{D3}$  brane used in KKLT
- ▶ We derived the complete low-energy effective theory for the  $\overline{D3}$  brane, using 4d  $\mathcal{N} = 1$  language

Thank you!