

Thoughts on a Hamiltonian formulation of sigma models, doubling and possible susy

U. Lindström¹

¹Department of Physics and Astronomy
Division of Theoretical Physics
University of Uppsala, SE,

Kerkyra, September , 2019



$2d$ (1,1) superspace

$$\int d^2\xi d\theta^+ d\theta^- \mathcal{L} = \int d^2\xi d\theta^+ d\theta^- (D_+ \phi^i E_{ij}(\phi) D_- \phi^j)$$

$$E_{ij} = (E_{(ij)} + E_{[ij]}) := G_{ij} + B_{ij}$$

$$\nabla_+^{(+)} D_- \phi^i = 0 ,$$

$$S_{\pm i} := \frac{\partial \mathcal{L}}{\partial D_{\mp} \phi^i} .$$

$$S_{+i} = -D_+ \phi^j E_{ji}$$

$$S_{-i} = E_{ij} D_- \phi^j .$$

$$\mathcal{H} = [S_{(+i} D_-) \phi^i + \mathcal{L}]_|$$

$$\mathcal{H} = S_{+i} E^{ij} S_{-j}$$

$$E^{ij} E_{jk} = \delta_k^i$$

Canonical Equations

$$D_+ \phi^i = \frac{\partial \mathcal{H}}{\partial S_-}$$

$$D_- \phi^i = \frac{\partial \mathcal{H}}{\partial S_+}$$

$$D_{(+S_-)i} = \frac{\partial \mathcal{H}}{\partial \phi^i} .$$

This system is an equivalent formulation of the motion.

A Phase Space Action

First order Lagrangian

$$\mathcal{L}_1 = S_{(+i} D_-) \phi^i - S_{+i} E^{ij}(\phi) S_{-j} ,$$

Can be obtained by gauge fixing a model on $\mathbb{T} \oplus \mathbb{T}^*$

$$\mathcal{D}_{(+\tilde{\phi}_i} D_-) \phi^i - \mathcal{D}_{+\tilde{\phi}_i} E^{ij}(\phi) \mathcal{D}_{-\tilde{\phi}_j} ,$$

$$\mathcal{D}_{\pm} \tilde{\phi}_i := D_{\pm} \tilde{\phi}_i + S_{\pm i} ,$$

$$\delta S_{\pm i} = D_{\pm} \Lambda_i .$$

The topological term is necessary for equivalence.

The Doubled Model

A $(1, 1)$ Lagrangian on $\mathbb{T} \oplus \mathbb{T}^*$ subject to constraints. In a polarisation:

$$D_{(+}\tilde{\phi}_i D_{-})\phi^i - D_{+}\tilde{\phi}_i E^{ij}(\phi) D_{-}\tilde{\phi}_j ,$$

$$D_{+}\tilde{\phi}_i = -D_{+}\phi^j E_{ji}$$

$$D_{-}\tilde{\phi}_i = E_{ij} D_{-}\phi^j .$$

Inserting the constraints returns the original action on $\mathbb{T}$.
Equivalently this action is obtained via the gauging procedure.

$$D_{(+}\tilde{\phi}_i D_{-)}\phi^i - D_{+}\tilde{\phi}_i E^{ij}(\phi)D_{-}\tilde{\phi}_j + \mu D_{+}\phi^i E_{ij}(\phi)D_{-}\phi^j .$$

- Can we find additional susy for the Doubled Lagrangian

$$\mathcal{L}(\tilde{\phi}, \phi)?$$

c.f. S.Driezen, A. Sevrin, and D.C. Thompson 2016

- Can we find additional susy for the first order Lagrangian

$$\mathcal{L}(\phi, S_{\pm})?$$

U.L. 2004, U.L., R.Minasian, A.Tomasiello and M.Zabzine 2005

- Can we find additional susy for the gauged Lagrangian

$$\mathcal{L}(\tilde{\phi}, \phi, S_{\pm})?$$

A general (1, 1) model

$$\frac{1}{2} D_+ \phi^A E_{AB} D_- \phi^B$$

has additional susy if

$$\delta \phi^A = \epsilon^+ J_{(+B}^A D_+ \phi^B + \epsilon^- J_{(-B}^A D_- \phi^B$$

$$J_{(\pm)}^2 = -\mathbb{1} , \quad J^t G J = G$$

$$\nabla^{(\pm)} J_{(\pm)} = 0 .$$

But the constraints?

One (complex) chiral superfield, Kähler potential $K(\phi, \bar{\phi})$

$$\mathcal{L} = D_+ \bar{\phi} K_{,\phi \bar{\phi}} D_- \phi + c.c.$$

Second susy $\phi := (\phi, \bar{\phi})^t$

$$\delta \phi := \delta \begin{pmatrix} \phi \\ \bar{\phi} \end{pmatrix} = \epsilon^+ \begin{pmatrix} i & 0 \\ 0 & -i \end{pmatrix} D_+ \begin{pmatrix} \phi \\ \bar{\phi} \end{pmatrix} =: \epsilon^+ J D_+ \phi$$

$$\tilde{\phi} = (\tilde{\phi}, \bar{\tilde{\phi}})^t$$

$$2D_{(+}\tilde{\phi}^t D_{-})\phi - D_{+}\tilde{\phi}^t K^{\bar{\phi}\phi} D_{-}\tilde{\phi} + 2\mu D_{+}\phi K_{,\phi\bar{\phi}} D_{-}\bar{\phi} \text{ ,}$$

$$D_{+}\bar{\tilde{\phi}} = -D_{+}\phi K_{,\phi\bar{\phi}}$$

$$D_{-}\tilde{\phi} = K_{,\phi\bar{\phi}} D_{-}\bar{\phi}$$

Not invariant under the unique second susy:

$$\delta\phi = \epsilon^+ JD_{+}\phi$$

$$\delta\tilde{\phi} = \pm\epsilon^+ JD_{+}\tilde{\phi}$$

If in addition

$$\delta \bar{S}_+ = -i\epsilon^+ D_+ \bar{S}_+ + 2i\epsilon^+ K_{,\phi\bar{\phi}\bar{\phi}} D_+ \phi D_+ \bar{\phi}$$

$$\delta S_- = -i\epsilon^+ D_+ S_- - 2i\epsilon^+ (\ln K_{,\phi\bar{\phi}})_{\phi} S_- D_+ \phi$$

then

$$2\mathcal{D}_{(+}\tilde{\phi}^t D_-)\phi - \mathcal{D}_+ \tilde{\phi}^t K^{\bar{\phi}\phi} \mathcal{D}_- \tilde{\phi} + 2\mu D_+ \phi K_{,\phi\bar{\phi}} D_- \bar{\phi}.$$

is (off-shell) supersymmetric. So implementation via gauging is supersymmetric.

A nontrivial b transformation:

$$\begin{aligned}\mathcal{G} &= (e^{-b})^t \mathbb{G} e^{-b} \\ &= (e^{-b})^t \begin{pmatrix} \mu g & 0 \\ 0 & -g^{-1} \end{pmatrix} e^{-b} \\ &= \begin{pmatrix} \mu g + bg^{-1}b & -bg^{-1} \\ g^{-1}b & -g^{-1} \end{pmatrix} .\end{aligned}$$

The b transform of the whole Lagrangian is then

$$\begin{aligned}& (D_+\phi, D_+\tilde{\phi})^t \mathbb{G} (D_-\phi, D_-\tilde{\phi}) \\ & \rightarrow (D_+\phi, -bD_+\phi + D_+\tilde{\phi})^t \mathcal{G} (D_-\phi, -bD_-\phi + D_-\tilde{\phi})\end{aligned}$$

On the constraints

$$D_+\tilde{\phi} = -D_+\phi g$$

$$D_-\tilde{\phi} = gD_-\phi ,$$

and the Lagrangian becomes

$$\begin{aligned} & (D_+\phi, D_+\tilde{\phi})^t \mathbb{G} (D_-\phi, D_-\tilde{\phi}) \\ & \rightarrow (D_+\phi, -eD_+\phi)^t \mathcal{G} (D_-\phi, +e^t D_-\phi) , \end{aligned}$$

i.e., the metric $\mathbb{G}$, gets replaced by $\mathcal{G}$, and g gets replaced by e and e^t in the constraints.

Evaluating the corresponding action for $\mu = -1$ yields

$$D_+\phi(g+b)D_-\phi = D_+\phi e D_-\phi$$

$$\mathcal{G} = \begin{pmatrix} g - bg^{-1}b & bg^{-1} \\ -g^{-1}b & g^{-1} \end{pmatrix}.$$

$$\eta = \begin{pmatrix} 0 & \mathbb{1} \\ \mathbb{1} & 0 \end{pmatrix}$$

$$\mathcal{S} = \eta^{-1}\mathcal{G} = \begin{pmatrix} -g^{-1}b & g^{-1} \\ g - bg^{-1}b & bg^{-1} \end{pmatrix}$$

$$\mathcal{S}^2 = \mathbb{1}$$

$$\mathcal{P}_{(\pm)} := \frac{1}{2}(1 \pm \mathcal{S})$$

$$\Omega = \begin{pmatrix} 0 & -\mathbb{1} \\ \mathbb{1} & 0 \end{pmatrix}$$

$$\mathcal{I} := -\Omega^{-1}\mathcal{G}, \quad \mathcal{K} := \Omega^{-1}\eta$$

$$-\mathcal{I}^2 = \mathcal{S}^2 = \mathcal{K}^2 = \mathbb{1}$$

$$\mathcal{ISK} = -\mathbb{1}, \quad \{\mathcal{I}, \mathcal{S}\} = \{\mathcal{S}, \mathcal{K}\} = \{\mathcal{K}, \mathcal{I}\} = 0$$

$$\Phi^M = \begin{pmatrix} \phi^\mu \\ \tilde{\phi}_{\tilde{\mu}} \end{pmatrix}$$

$$S_2 = \frac{1}{4} \int d^2x d^2\theta (D_+ \Phi)^t \mathcal{K} \mathcal{G} D_- \Phi ,$$

$$\mathcal{P}_{(\pm)} D_\pm \Phi = 0 , \quad \Rightarrow D_+ \tilde{\phi}_{\tilde{\mu}} = -D_+ \phi^\mu e_{\mu\tilde{\mu}} , \quad D_- \tilde{\phi}_{\tilde{\mu}} = e_{\tilde{\mu}\mu} D_- \phi^\mu$$

$$S_2 \rightarrow \frac{1}{2} \int d^2x d^2\theta D_+ \phi e D_- \phi$$

or

$$S_1 \rightarrow \frac{1}{2} \int d^2x d^2\theta D_+ \phi e D_- \phi - D_{(+}\tilde{\phi} D_{-})\phi$$

$$\mathcal{J}^2 = -\mathbb{1} \ , \quad \mathcal{J}^t \mathcal{G} \mathcal{J} = \mathcal{G} \ , \quad \mathcal{J}^t \eta \mathcal{J} = \eta$$

$$[\mathcal{J}, \mathcal{S}] = 0 \ , \quad [\mathcal{J}, \Omega] = 0 \ ,$$

$$\Rightarrow \mathcal{J} = e^b \mathcal{J}_0 e^{-b}$$

$$\mathcal{J}_0 = \begin{pmatrix} J & 0 \\ 0 & -J^t \end{pmatrix} .$$

$$\delta\Phi = \epsilon \mathcal{J} D\Phi$$