

# Forward-backward asymmetry in the gauge-Higgs unification at the International Linear Collider

Shuichiro Funatsu

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# Gauge-Higgs Unification

$$A_M = (A_\mu, A_y)$$


Higgs boson

The Higgs mass is protected by a gauge symmetry

The Higgs boson obtains **finite** mass at 1-loop level

Effective potential is written by Wilson line phase

$$\theta_H \equiv g \int_C dy \langle A_y \rangle$$

Hosotani mechanism

Y. Hosotani (1983)

H. Hatanaka, T. Inami, C.S. Lim (1998)

# $SO(5) \times U(1)$ GHU

On the Randall-Sundrum warped (AdS) spacetime

$$ds^2 = e^{-2ky} \eta_{\mu\nu} dx^\mu dx^\nu + dy^2$$



Agashe, Contino, Pomarol (2005)

Contino, Da Rold, Pomarol (2007)

Medina, Shah, Wagner (2007)

Sakamura, Hosotani (2007)

Hosotani, Oda, Ohnuma, Sakamura (2008)

SF, Hatanaka, Hosotani, Orikasa, Shimotani (2013)

# $SO(5) \times U(1)$ GHU

$$SO(5) \times U(1)_X$$

$$\xrightarrow{\text{B.C.}} SO(4) \times U(1)_X$$

$$W_R, Z_R$$

$$\simeq SU(2)_L \times SU(2)_R \times U(1)_X$$

$$\xrightarrow{\text{Brane int.}} SU(2)_L \times U(1)_Y$$

$$W, Z, \gamma$$

$$\xrightarrow{\text{Hosotani mechanism}} U(1)_{\text{EM}}$$

# Higgs couplings

$$\begin{array}{c} HWW \\ HZZ \end{array} \simeq \begin{array}{c} \text{SM} \\ \text{value} \end{array} \times \cos \theta_H$$

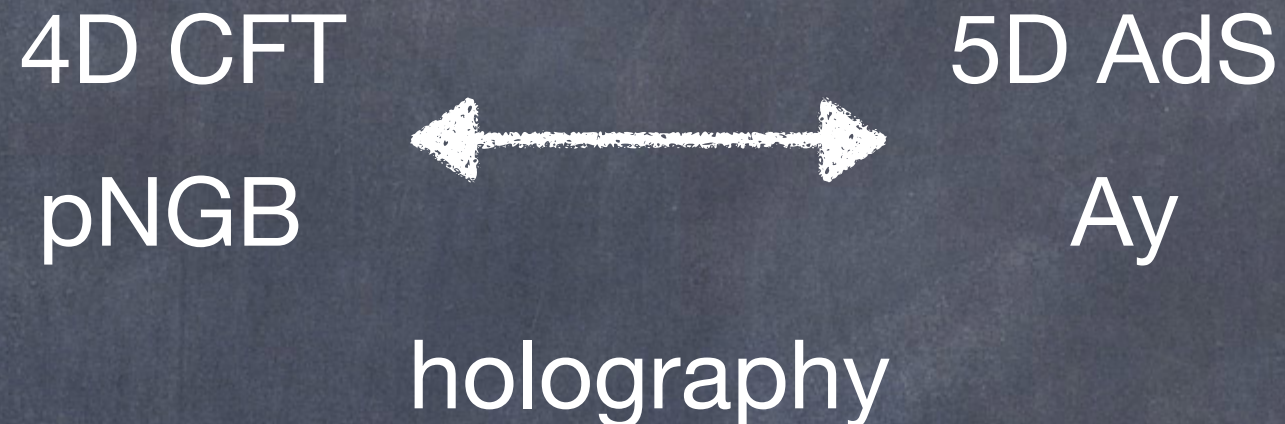
$$\frac{\Gamma(H \rightarrow WW)}{\Gamma(H \rightarrow ZZ)} \simeq \begin{array}{c} \text{SM} \\ \text{value} \end{array} \times \cos^2 \theta_H$$

$\sin \theta_H$  corresponds to  $\xi$  in MCHMs

$\Gamma(H \rightarrow \gamma\gamma)$  and  $\Gamma(H \rightarrow Z\gamma)$  are evaluated in  
SF, Hatanaka, Hosotani, Orikasa, Shimotani (2013)  
SF, Hatanaka, Hosotani (2015)

# Comments on MCHM

Agashe et al, argue

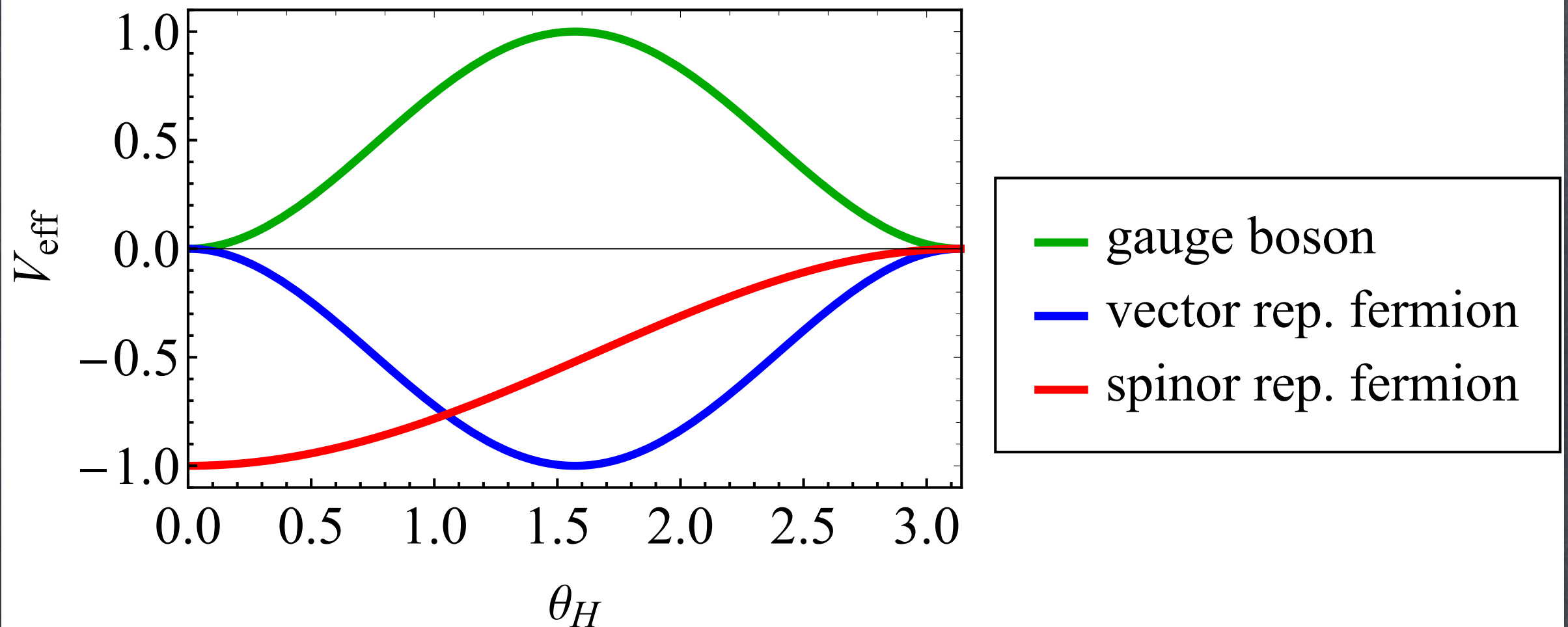


GHU on RS = holographic composite Higgs model

GHU has

- calculable effective potential
- restricted Lagrangian

# Effective Potential



Both vector-rep. and spinor-rep. are necessary to obtain  $\theta_H \neq 0, \pi/2$

# matter content

Two types of the  $SO(5) \times U(1)$  GHU models

- quarks in **5** and non-SM fermions in **4**  
left-handed quarks in **(2, 2)**  
right-handed quarks in **(1, 1)**

SF, Hatanaka, Hosotani, Orikasa, Shimotani (2013)

## Today's talk

- quarks in **4** and non-SM fermions in **5**  
left-handed quarks in **(2, 1)**  
right-handed quarks in **(1, 2)**

SF, Hatanaka, Hosotani, Orikasa, Yamatsu (2019)



# Parameters

- metric:  $k, e^{kL}$   $m_Z$
- gauge:  $g_W^{(5D)}, g_X^{(5D)}$   $\alpha_{\text{EM}}, \sin^2 \theta_W$
- fermion:  $c_t, r_t$   $m_t, m_b$   
 $c_F$   $m_H$   
 $(c \equiv m^{(5D)} / k)$



One parameter  $e^{kL} \rightarrow \theta_H, m_{\text{KK}}$  is determined

# Parameters

Upper bound by LHC



| $\theta_H$ | $e^{kL}$           | $ c_t $  | $m_{KK}$ (GeV) |
|------------|--------------------|----------|----------------|
| 0.10       | $2.90 \times 10^4$ | 0.16116  | 8063           |
| 0.09       | $1.70 \times 10^4$ | 0.11646  | 8721           |
| 0.08       | $1.01 \times 10^4$ | 0.008914 | 9544           |



Lower bound from top-mass realisation

small  $\theta_H \Leftrightarrow$  large KK scale

# Fermion Localisation

$$\Psi_{L,R}(x, y) = \frac{e^{\frac{3}{2}ky}}{\sqrt{L}} \sum_{n=0}^{\infty} \psi_{L,R}^{(n)}(x) \frac{f_{L,R}^{(n)}(y)}{\sqrt{N^{(n)}}},$$

$$f_{L,R}^{(0)}(y) = e^{(\frac{1}{2} \mp c)ky}$$

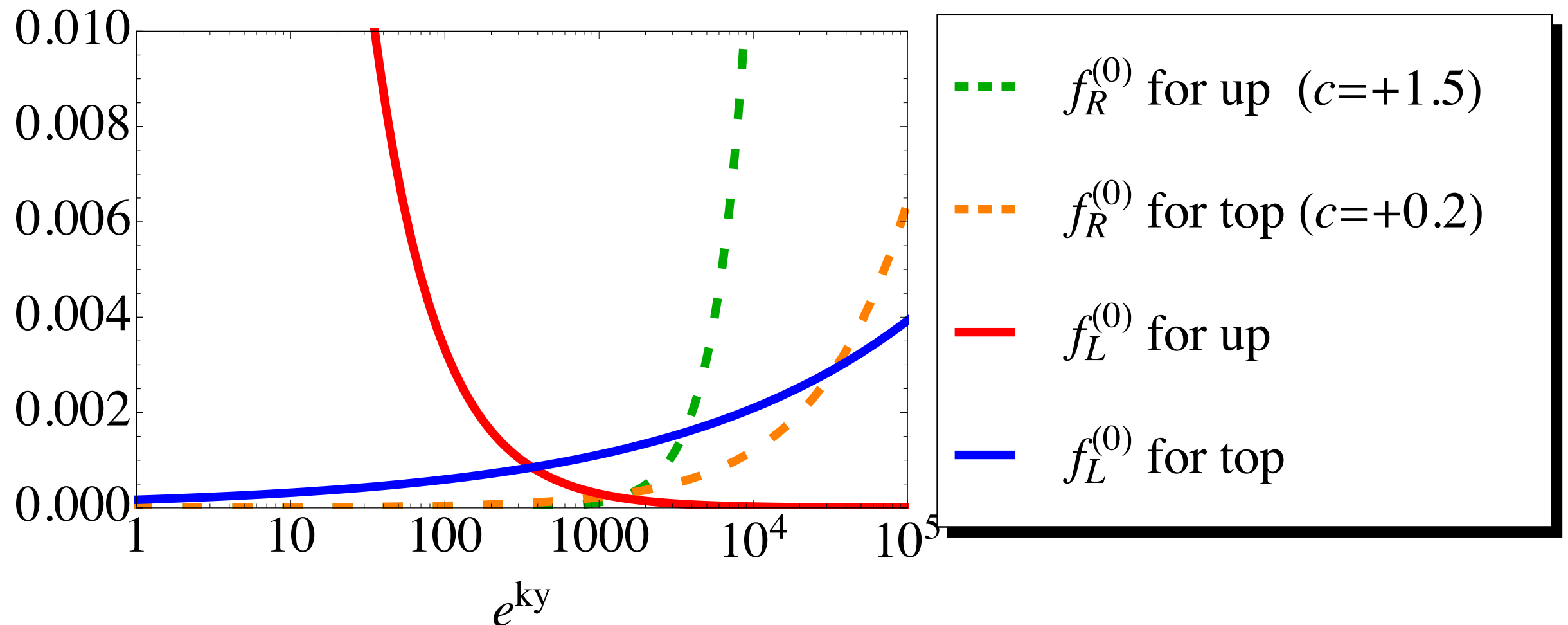
Left-handed : localised toward the UV brane ( $c > +1/2$ )  
the IR brane ( $c < +1/2$ )

Right-handed : localised toward the IR brane ( $c > -1/2$ )  
the UV brane ( $c < -1/2$ )

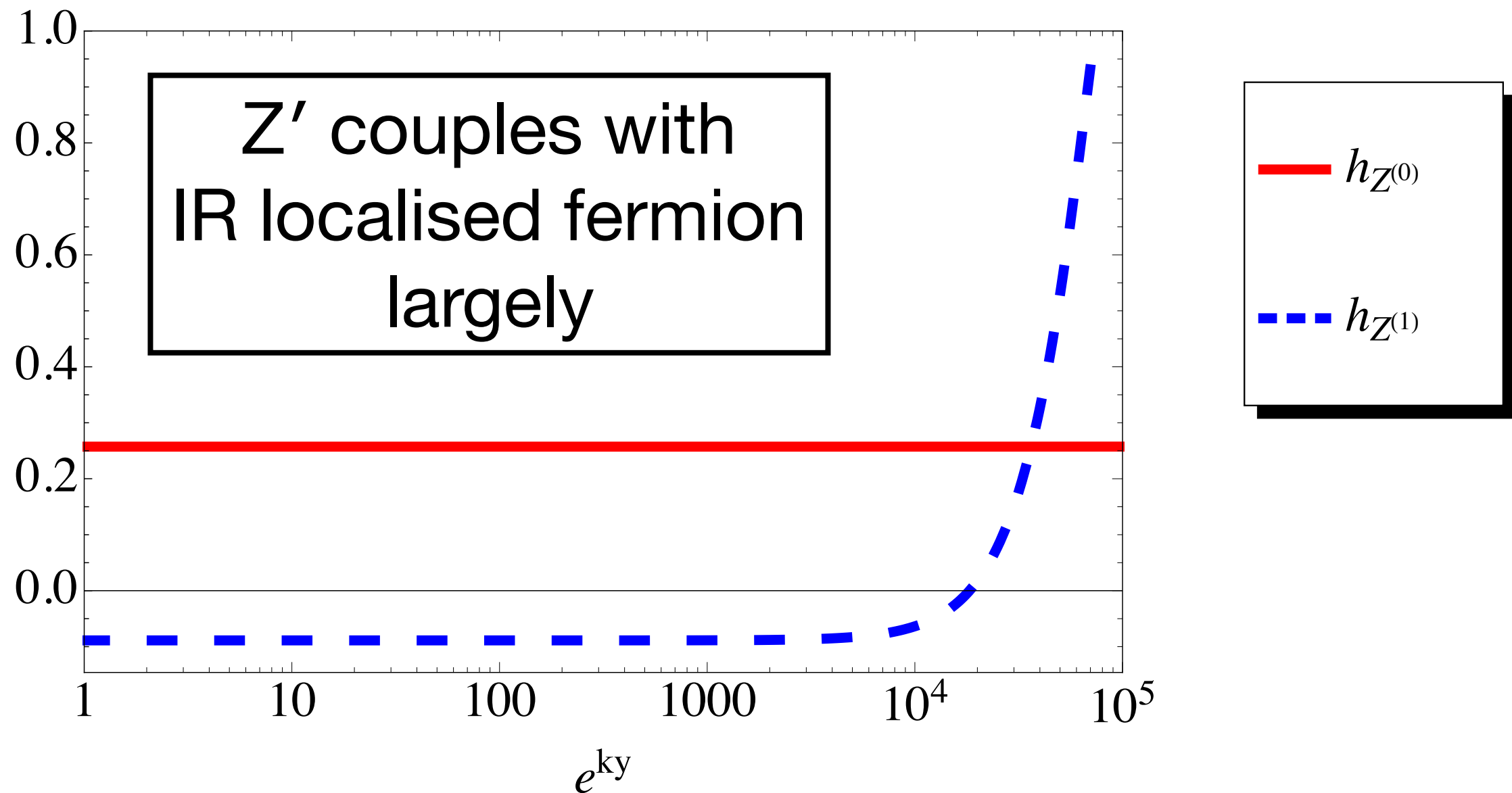
$c \rightarrow -c$ , the left- and right-handed are reversed

mass is invariant

# Fermion Localisation



# Gauge Boson Localisation



$Z^{(0)}$ : localised toward the IR brane very slightly

$Z^{(1)}$ : localised toward the IR brane

# Z-couplings

Z-boson couplings in  $g_W / \cos\theta_W$  unit  
for  $\sin^2\theta_W = 0.2312$  and  $\theta_H = 0.10$ ,

|           | $g_{Ze_L}$ | $g_{Ze_R}$ |
|-----------|------------|------------|
| SM        | $-0.2688$  | $+0.2312$  |
| $c_e > 0$ | $-0.2688$  | $+0.2313$  |
| $c_e < 0$ | $-0.2664$  | $+0.2338$  |

For  $c < 0$  and  $\theta_H = 0.10$ ,  $\sin^2\theta_W = 0.2290$  is required,  
so  $c > 0$  is set.

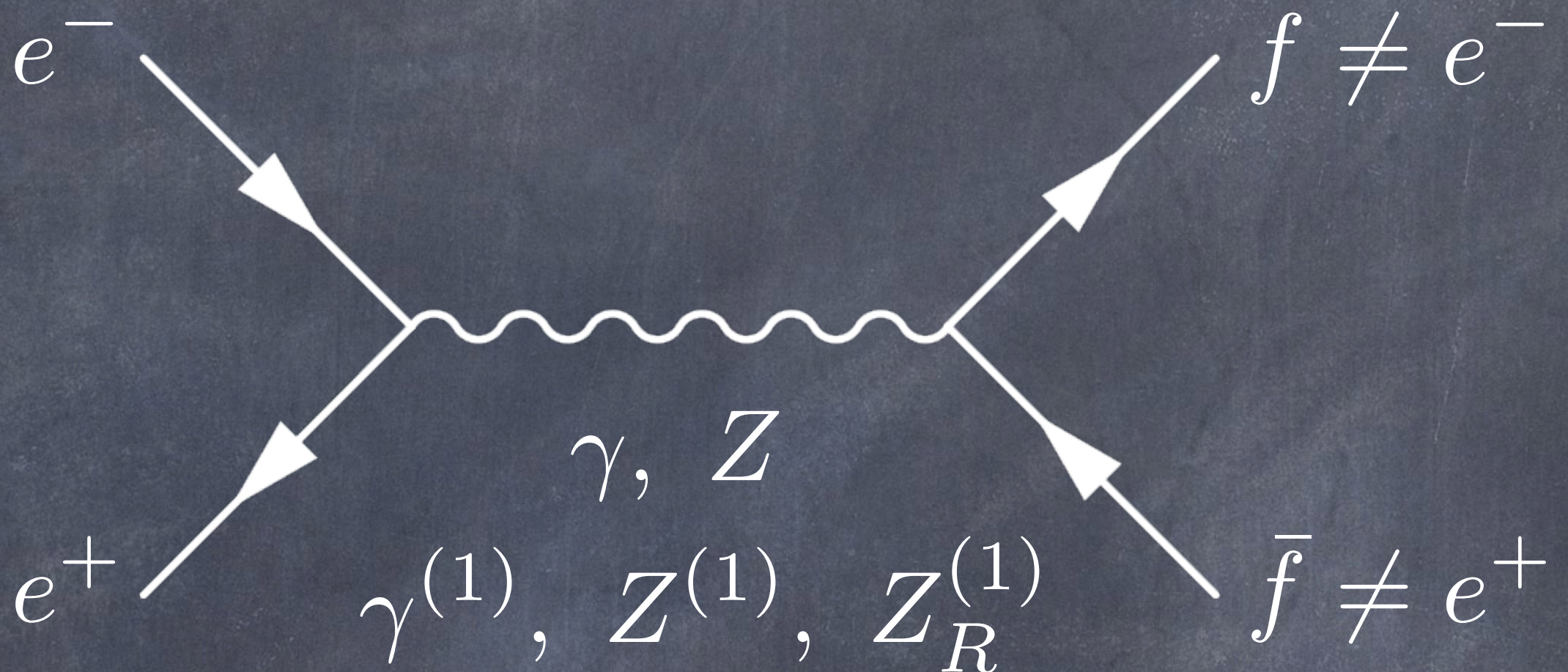
# Z<sup>(1)</sup>-couplings

Z<sup>(1)</sup>-boson couplings in  $g_W / \cos\theta_W$  unit  
for  $\sin^2\theta_W = 0.2312$  and  $\theta_H = 0.10$ ,

|           | $g_{Z^{(1)}e_L}$ | $g_{Z^{(1)}e_R}$ |
|-----------|------------------|------------------|
| $c_e > 0$ | +0.0987          | +0.9148          |
| $c_e < 0$ | −1.0535          | −0.0858          |

| $\theta_H$ | $m_{Z^{(1)}}$ | sign of $(c_l, c_q)$     |          |          |          |
|------------|---------------|--------------------------|----------|----------|----------|
|            |               | $\Gamma_{Z^{(1)}}(+, +)$ | $(+, -)$ | $(-, +)$ | $(-, -)$ |
| 0.10       | 6585          | 429                      | 1632     | 959      | 2162     |
| 0.09       | 7149          | 463                      | 1674     | 1014     | 2225     |
| 0.08       | 7855          | 534                      | 1705     | 1112     | 2283     |

# Processes



Tree level only

Bhabha process: not yet

# Polarisation

Ignoring the Higgs exchange,

$$\frac{d\sigma}{d\cos\theta} = \frac{1}{4} \left[ (1 - P_{e^-})(1 + P_{e^+}) \frac{d\sigma_{LR}}{d\cos\theta} + (1 + P_{e^-})(1 - P_{e^+}) \frac{d\sigma_{RL}}{d\cos\theta} \right]$$

$P = +1$  : right-handed

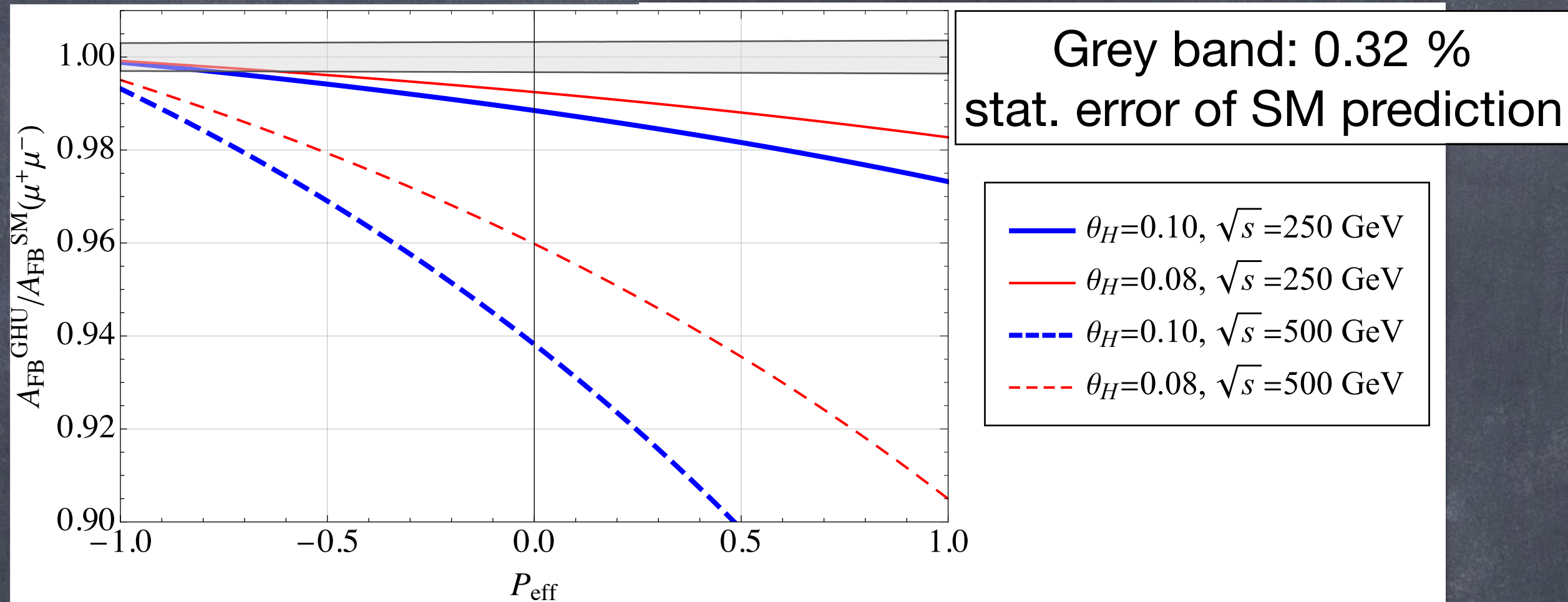
$P = -1$  : left-handed

It is rewritten by  $P_{\text{eff}} = \frac{P_{e^-} - P_{e^+}}{1 - P_{e^-}P_{e^+}}$  as

$$\frac{d\sigma}{d\cos\theta} = \frac{(1 - P_{e^+}P_{e^-})}{4} \left[ (1 - P_{\text{eff}}) \frac{d\sigma_{LR}}{d\cos\theta} + (1 + P_{\text{eff}}) \frac{d\sigma_{RL}}{d\cos\theta} \right]$$

$$e^+e^- \rightarrow \mu^+\mu^-$$

Deviation of AFB from SM ( $c_l > 0$ )



−1.1 % for  $P_{\text{eff}} = 0$ ,  $\theta_H = 0.10, \sqrt{s} = 250 \text{ GeV}$

−2.5 % for  $P_{\text{eff}} = 0.887$ , //

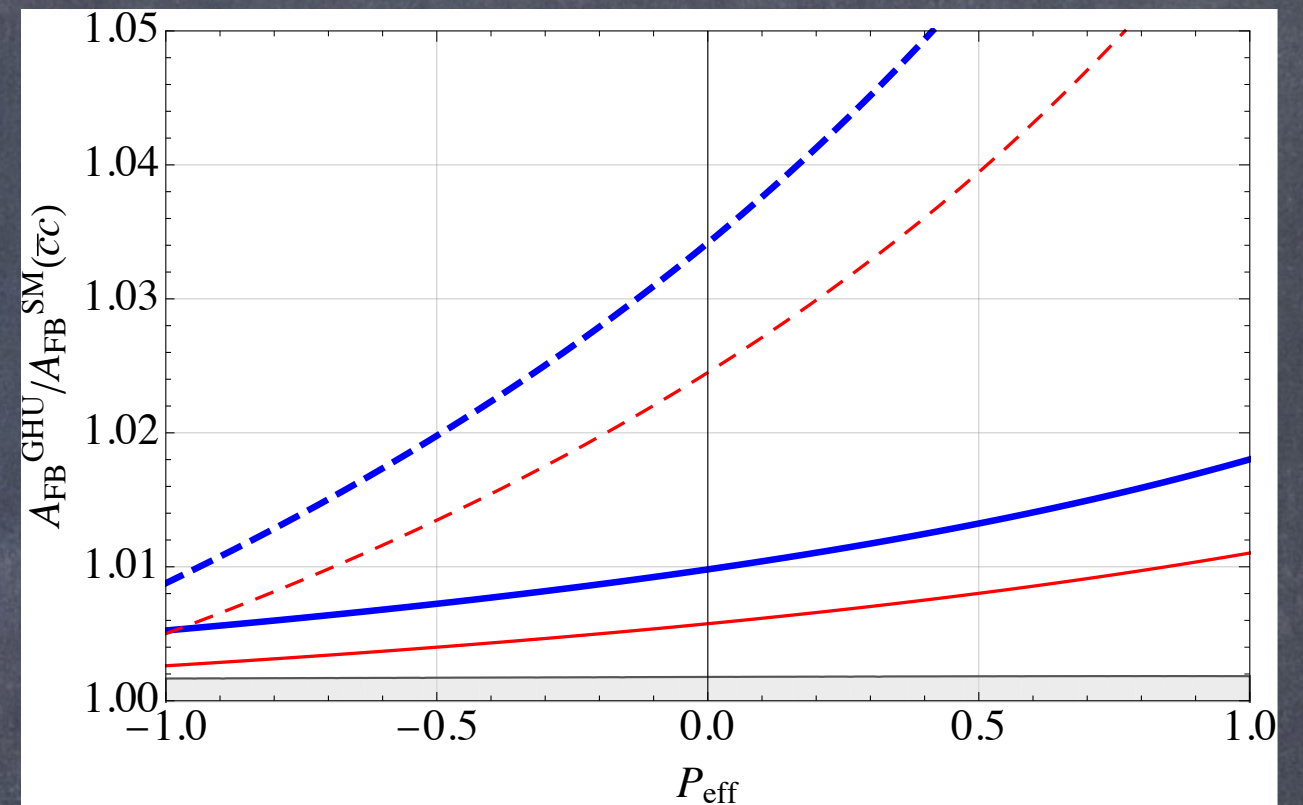
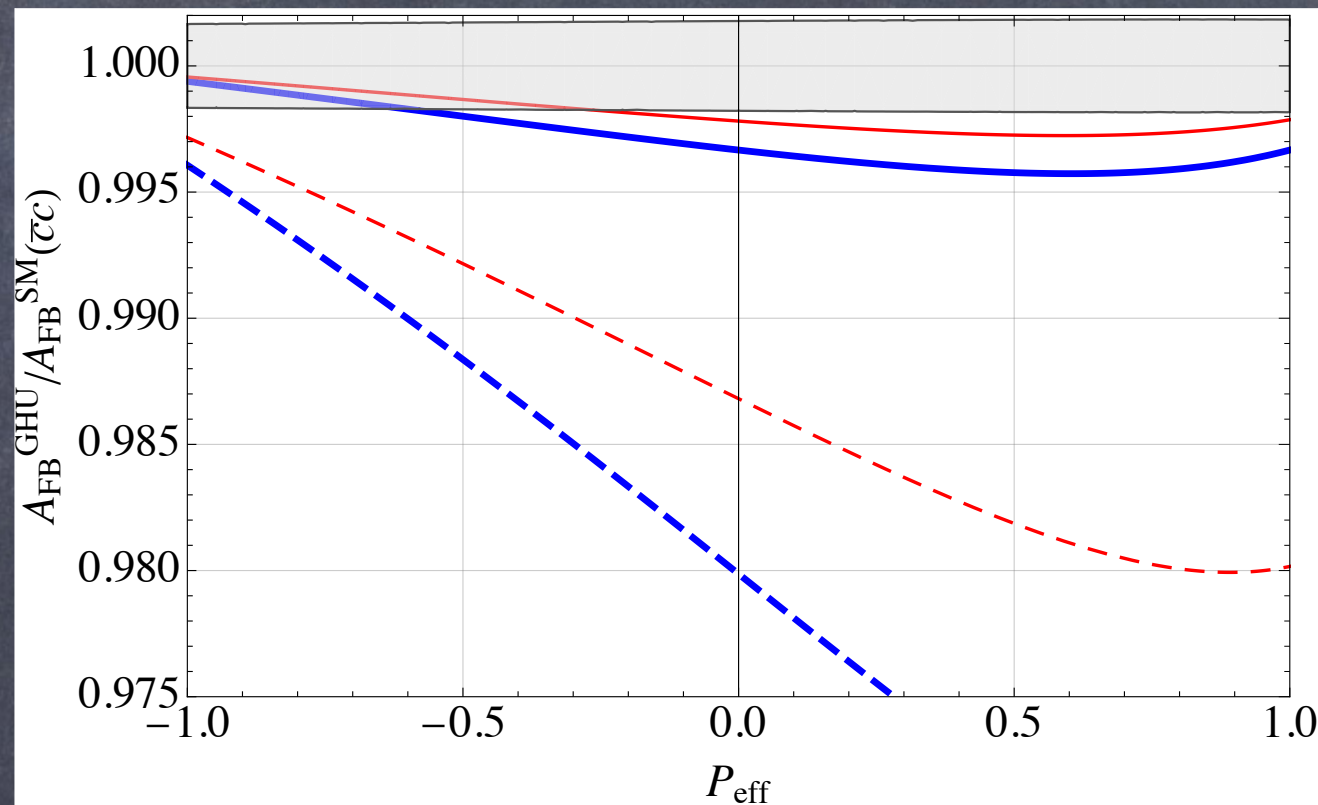
$$e^+e^- \rightarrow \bar{c}c$$

Deviation of AFB from SM

Grey band: 0.17 %  
stat. error of SM prediction

$C_c > 0$

$C_c < 0$



-0.33 % for  $P_{eff} = 0$

+0.98 % for  $P_{eff} = 0$

-0.38 % for  $P_{eff} = 0.887$

+1.68 % for  $P_{eff} = 0.887$

$(\theta_H = 0.10, \sqrt{s} = 250 \text{ GeV})$

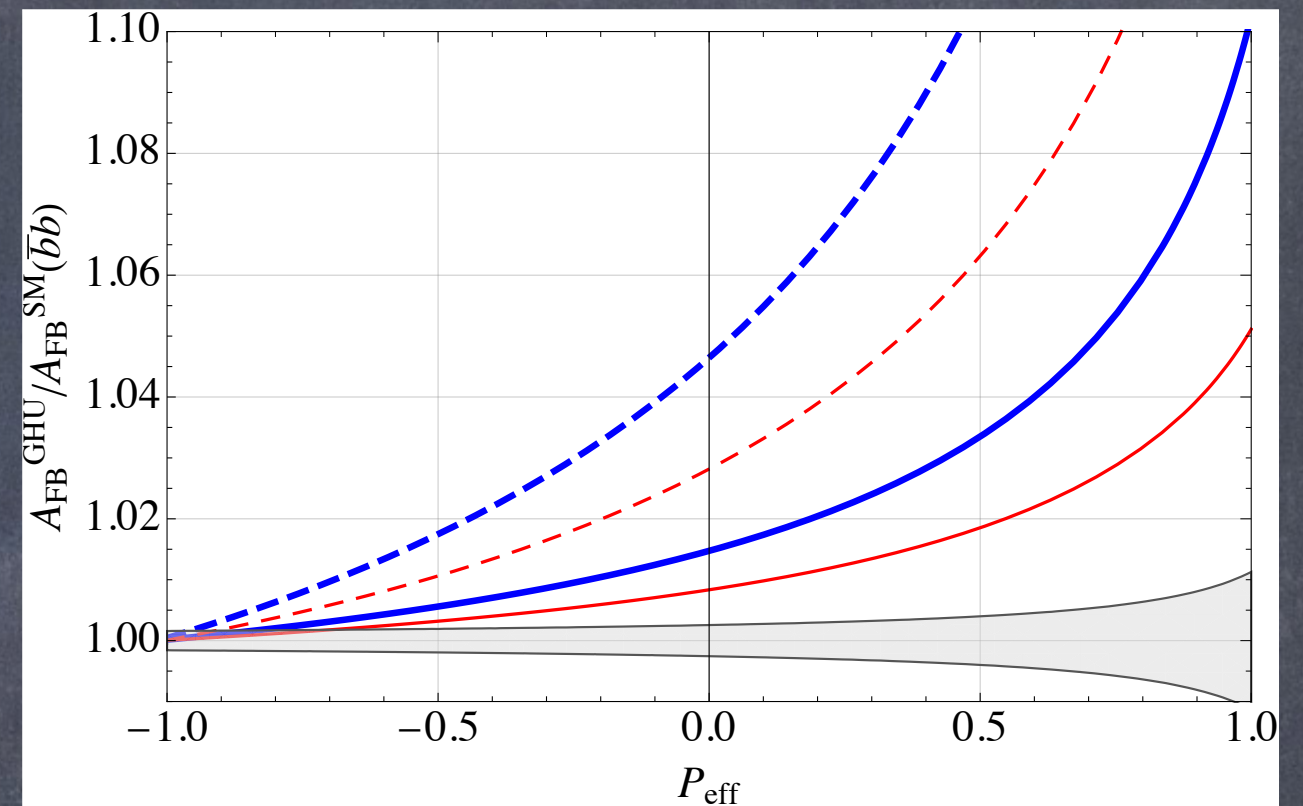
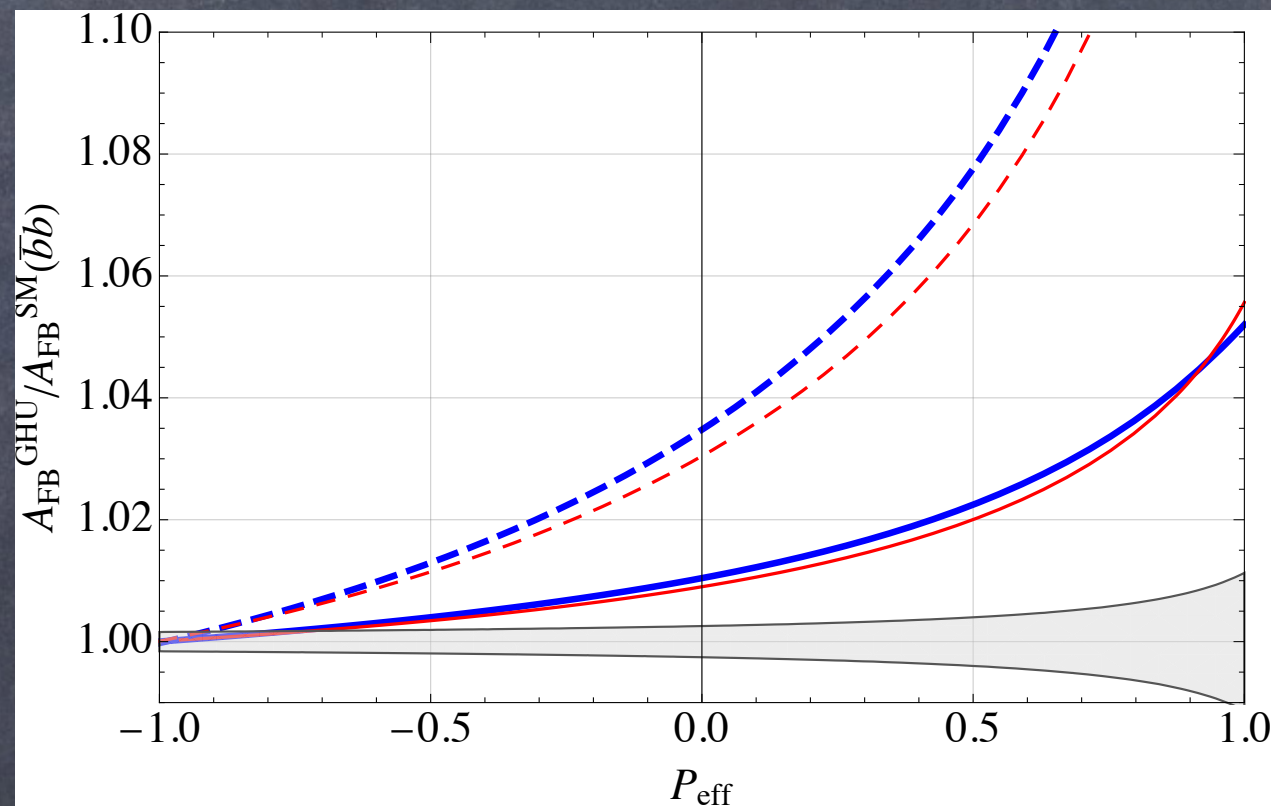
$$e^+e^- \rightarrow \bar{b}b$$

Deviation of AFB from SM

Grey band: 0.11 – 0.80 %  
stat. error of SM prediction

$C_b > 0$

$C_b < 0$



+1.0 % for  $P_{eff} = 0$

+4.2 % for  $P_{eff} = 0.887$

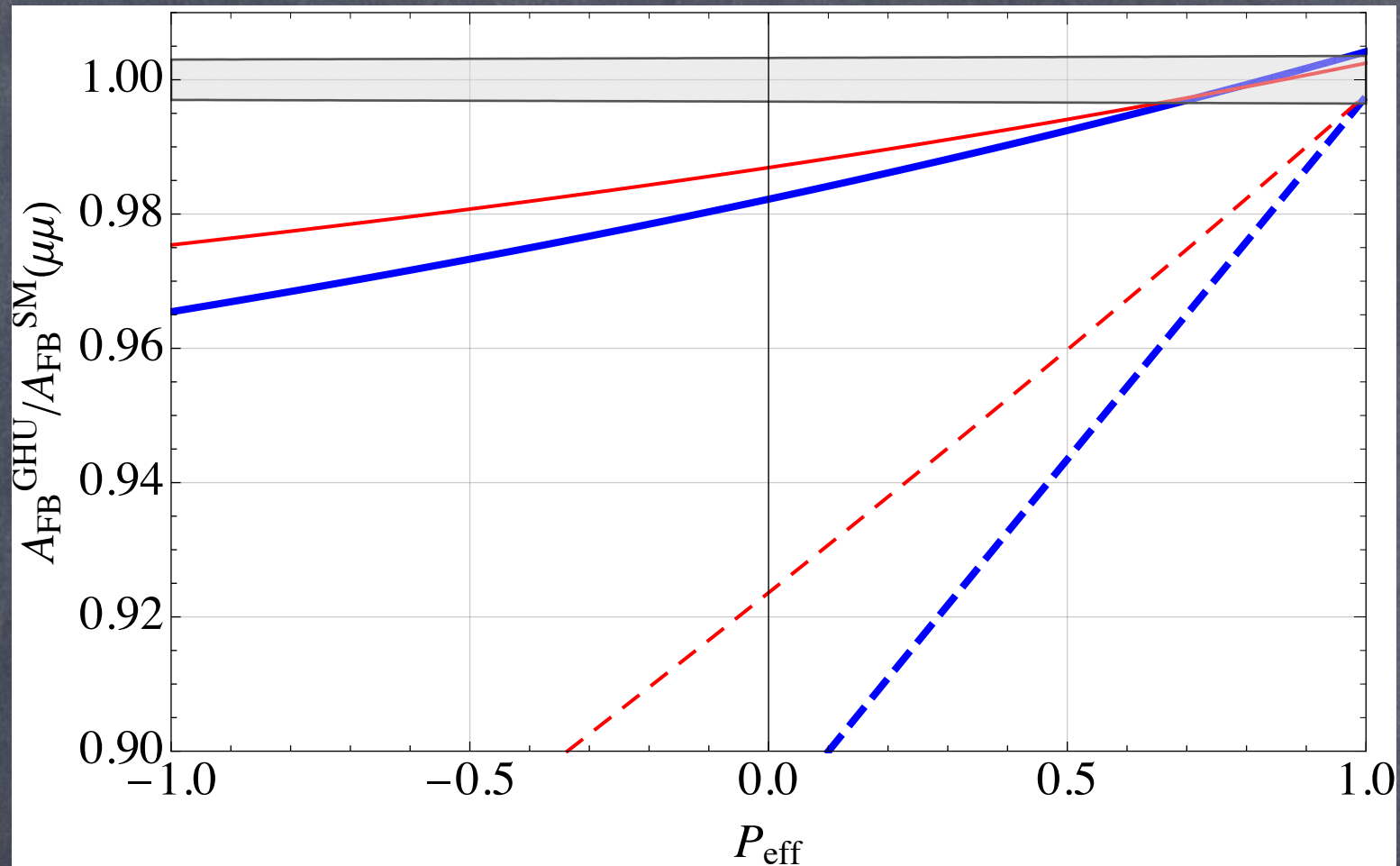
+1.5 % for  $P_{eff} = 0$

+7.3 % for  $P_{eff} = 0.887$

$(\theta_H = 0.10, \sqrt{s} = 250 \text{ GeV})$

# $e^+e^- \rightarrow \mu^+\mu^-$ (preliminary)

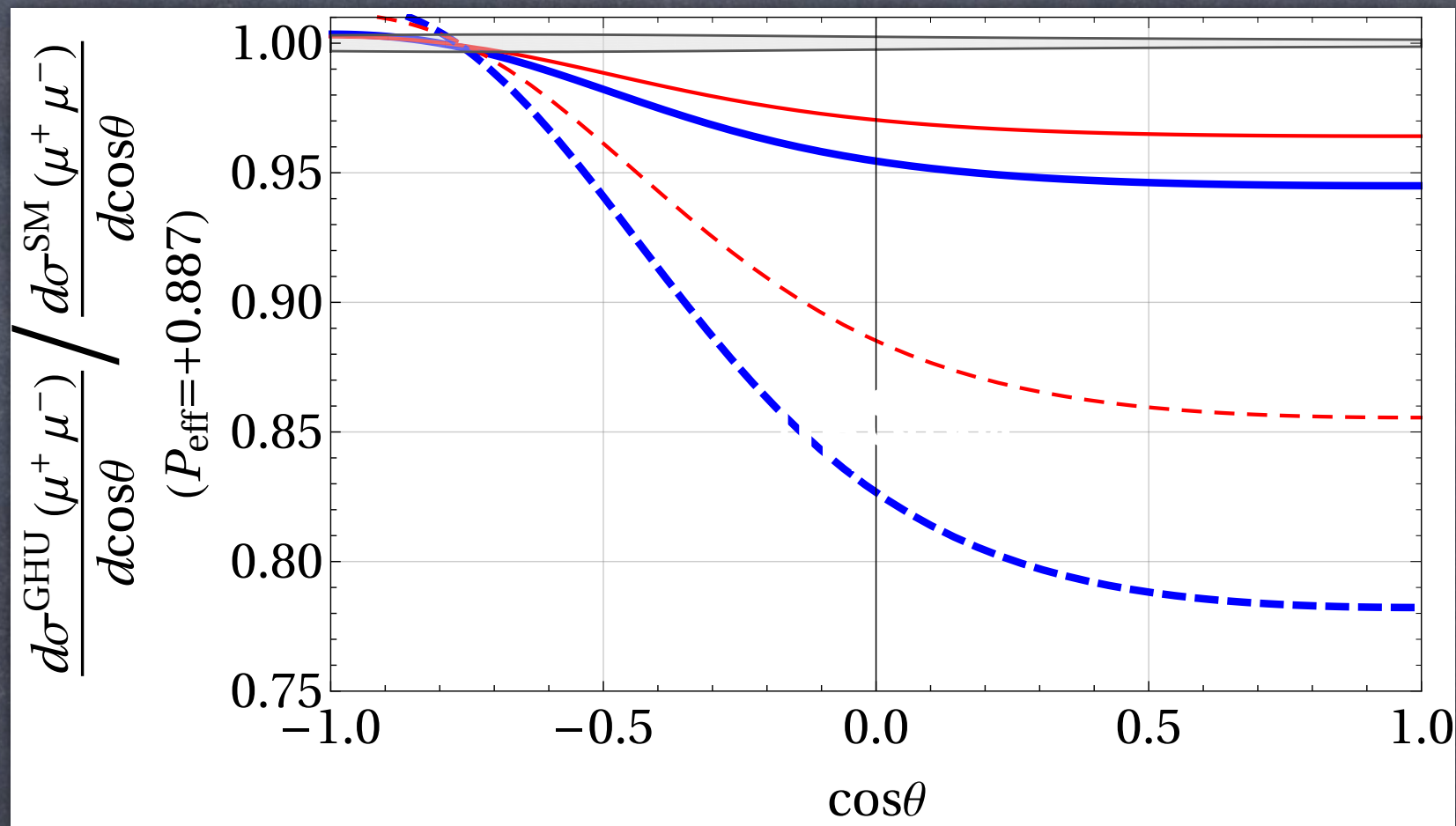
Deviation of AFB from SM ( $c_l < 0$ )



- 1.1 % for  $P_{\text{eff}} = 0$ ,  $\theta_H = 0.10$ ,  $\sqrt{s} = 250$  GeV
- 3.9 % for  $P_{\text{eff}} = -0.887$ , //

# $e^+e^- \rightarrow \mu^+\mu^-$ (preliminary)

Deviation of  $d\sigma/d\cos\theta$  from SM ( $c_l > 0$ )



$5.8 \times 10^3$  events (0.4 % stat. error) in SM  
for  $\cos\theta = [0.8, 0.9]$ ,  $P_{\text{eff}} = 0.887$

# Summary

- gauge-Higgs unification is a solution to the fine-tuning problem of the Higgs mass
- KK photon,  $Z$ ,  $Z_R$  are  $Z'$ 's  
They have large coupling asymmetries
- 7 TeV  $Z'$  effects are seen by  $\sqrt{s} = 250$  GeV