

EISA

European Institute for Sciences and Their Applications



Corfu Summer Institute

19th Hellenic School and Workshops on Elementary Particle Physics and Gravity
Corfu, Greece 2019



Workshop on Connecting Insights in Fundamental Physics:

Standard Model and Beyond

August 31 - September 11, 2019

**Highlights from
ATLAS
Louis FAYARD (LAL Orsay)
*on behalf of the ATLAS Collaboration***

- ♪ *Historical introduction , Setting the stage*
- ♪ *Results from (Run-1 and) Run-2*
- ♪ *Future of ATLAS , Run-3 , HL-LHC*
- ♪ *Conclusions*
- ♪ *Backup*

see also recent conferences



The logo features the text "EUROPEAN PHYSICAL SOCIETY CONFERENCE ON HIGH ENERGY PHYSICS" in large, bold, orange letters at the top. Below this is the date "10-17 JULY 2019 — Ghent, Belgium". The background is a blue and yellow gradient. In the center, there is a stylized illustration of the St. Bavo's Cathedral in Ghent, Belgium, with the "HEP2019" logo overlaid. To the right of the cathedral is a globe with the EPS logo. The bottom of the page contains logos for various organizations involved in the conference, including Ghent University, UCLouvain, ULB, VUB, Nikhef, CERN, Springer, and others.

Presentation by Andreas Hoecker

Presentation by Pierre Savard



Large number of results !

I will be selective with only
few details !

For more results : look at backup
and references

*I will insist more on raw results
(less on phenomenological
interpretations)*

Rien n'est cru si fermement que ce
que l'on sait le moins

Nothing is believed more strongly than which we know the least

Montaigne, Essais

- ♪ *Historical introduction , Setting the stage*
- ♪ *Results from (run 1 and) run 2*
- ♪ *Future of ATLAS , run 3 , HL-LHC*
- ♪ *Conclusions*
- ♪ *Backup*

Spontaneous Symmetry breaking

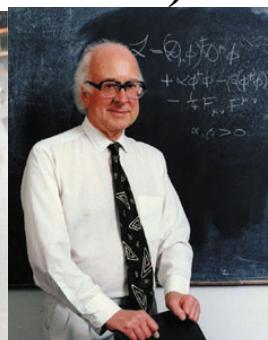
The Brout-Englert-Higgs mechanism

The LHC

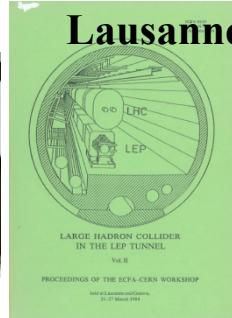


in a

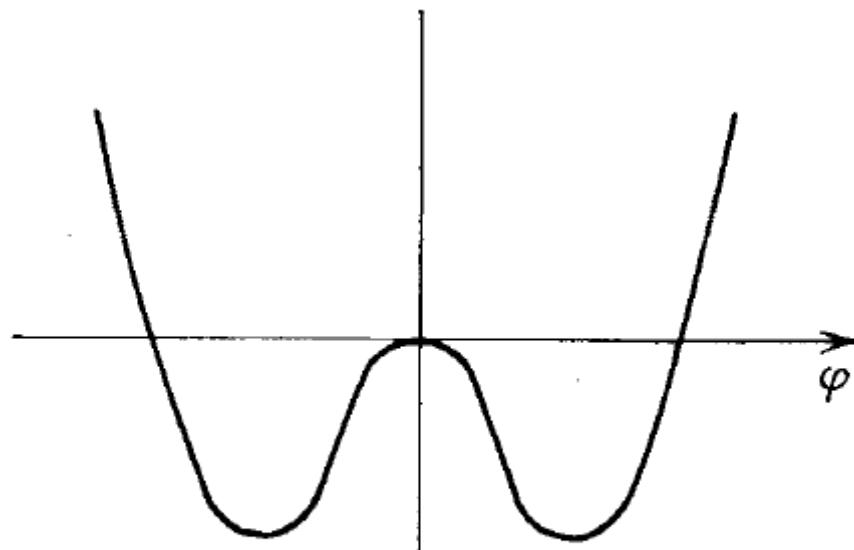
- 1950 Ginzburg-Landau (Meissner-Ochsenfeld effect) → London penetration length $\sim W$ mass
 1959 Nambu
 1960 Goldstone, Gell-Mann Levy, NJL
 1961 Schwinger
 1962 Anderson
 1964 Brout, Englert, Higgs, Guralnik, Hagen, Kibble
 1967 Weinberg, Salam Faddeev, Popov
 1970 Glashow, Iliopoulos,
 Maiani, 't Hooft, Veltman, BRST.....



1983 Rubbia, van der M particles of mass $\sqrt{-2\mu_0^2}$
 1984 Spin⁻ Parallel discovery of W and Z at CERN
 1989 construction beginning
 1992 LOI of 'lai
 1994 TP of ATL
 1995 discovery
 1996 approval o
 1998 approval o
 1999 ATLAS Ph
 2000



2006
 2008 CMS Phys
 2010 ATLAS E
 start-up a
 2012 4th July discovery of boson (m ~ 125 GeV)
 2013 boson like properties
 2014



$$\frac{\mu_0^2}{2} \varphi^2 + \frac{\lambda_0}{24} \varphi^4$$

Nobel prize to Englert and Higgs

Tevatron data ended in Sept 2011

2008

10th september 2008 : first beams around
19th september 2008 : incident

2009

20th november 2009 : first beams around (again)
december 2009 : collisions at 2.36 TeV cms

2010

30th march 2010 : first collisions at 7 TeV cms
august 2010 : luminosity of $10^{31} \text{ cm}^{-2} \text{ s}^{-1}$

2011

may 2011 : luminosity $> 10^{33} \text{ cm}^{-2} \text{ s}^{-1}$
november 2011 : integrated luminosity $\sim 5 \text{ fb}^{-1}$
13th december 2011 : first ‘signal’ around 126 GeV

2012

march 2012 : start again at 8 TeV
(50 ns between bunches)
4th July 2012 : evidence for a new boson
(8 TeV integrated luminosity $\sim 6 \text{ fb}^{-1}$)

2013

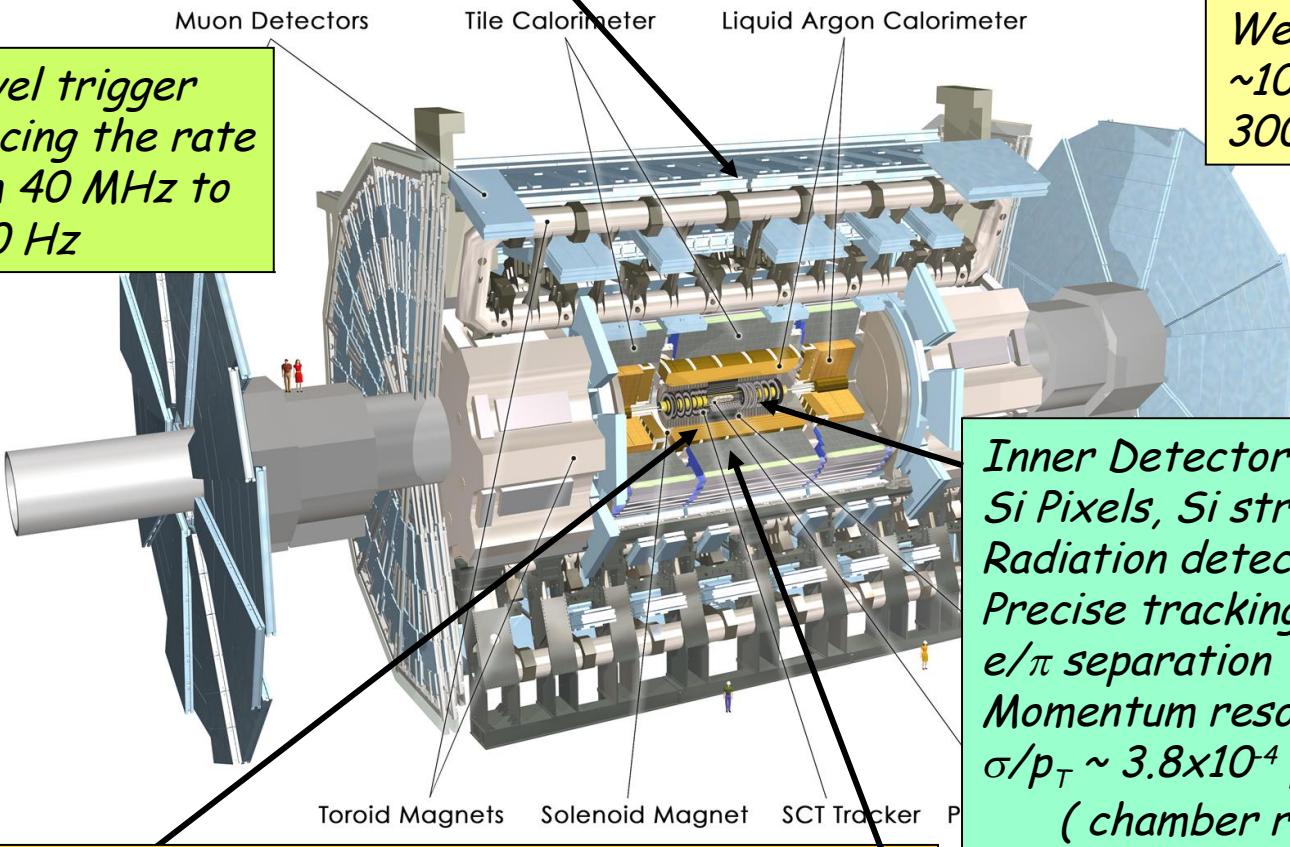
(Standard-Model) boson-like properties
peak luminosity $7 \cdot 10^{33} \text{ cm}^{-2} \text{ s}^{-1}$
integrated luminosity $\sim 5+20 \text{ fb}^{-1}$



end of Run-1

Muon Spectrometer ($|\eta| < 2.7$) : air-core toroids ($B \sim 0.5 / 1T$ in barrel/ end-cap) with gas-based muon chambers Muon trigger and measurement with momentum resolution $< 10\%$ up to $E_\mu \sim 1$ TeV

run 1 ATLAS detector

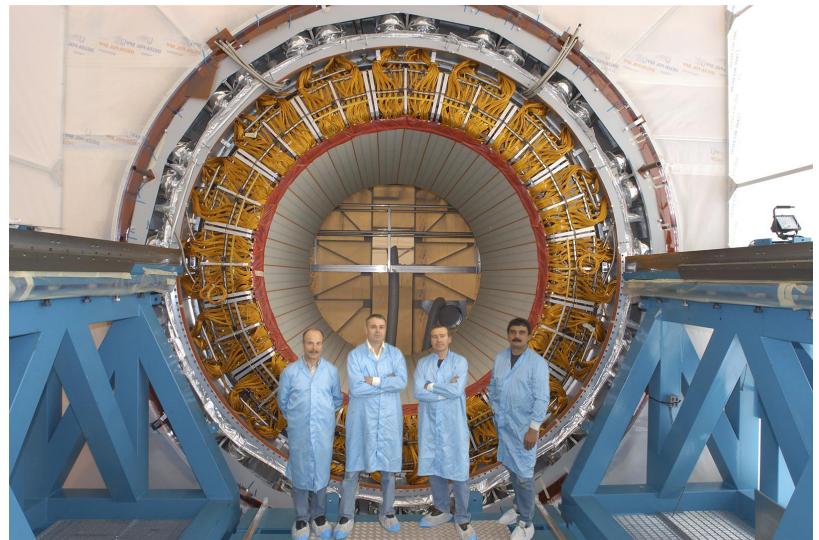
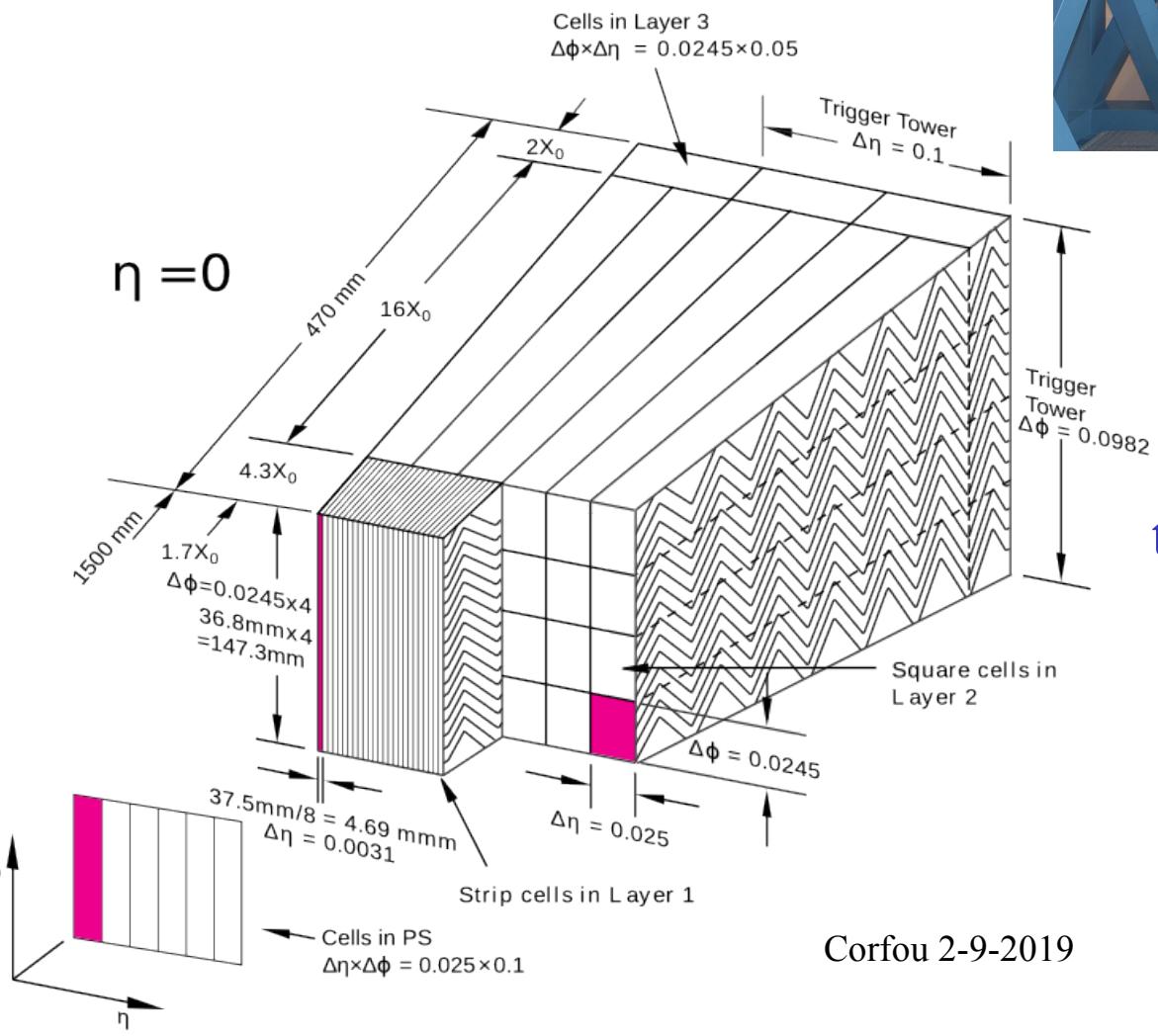


EM calorimeter: Pb-LAr Accordion
e/ γ trigger, identification and measurement
E-resolution: $\sigma/E \sim 10\%/\sqrt{E}$

C

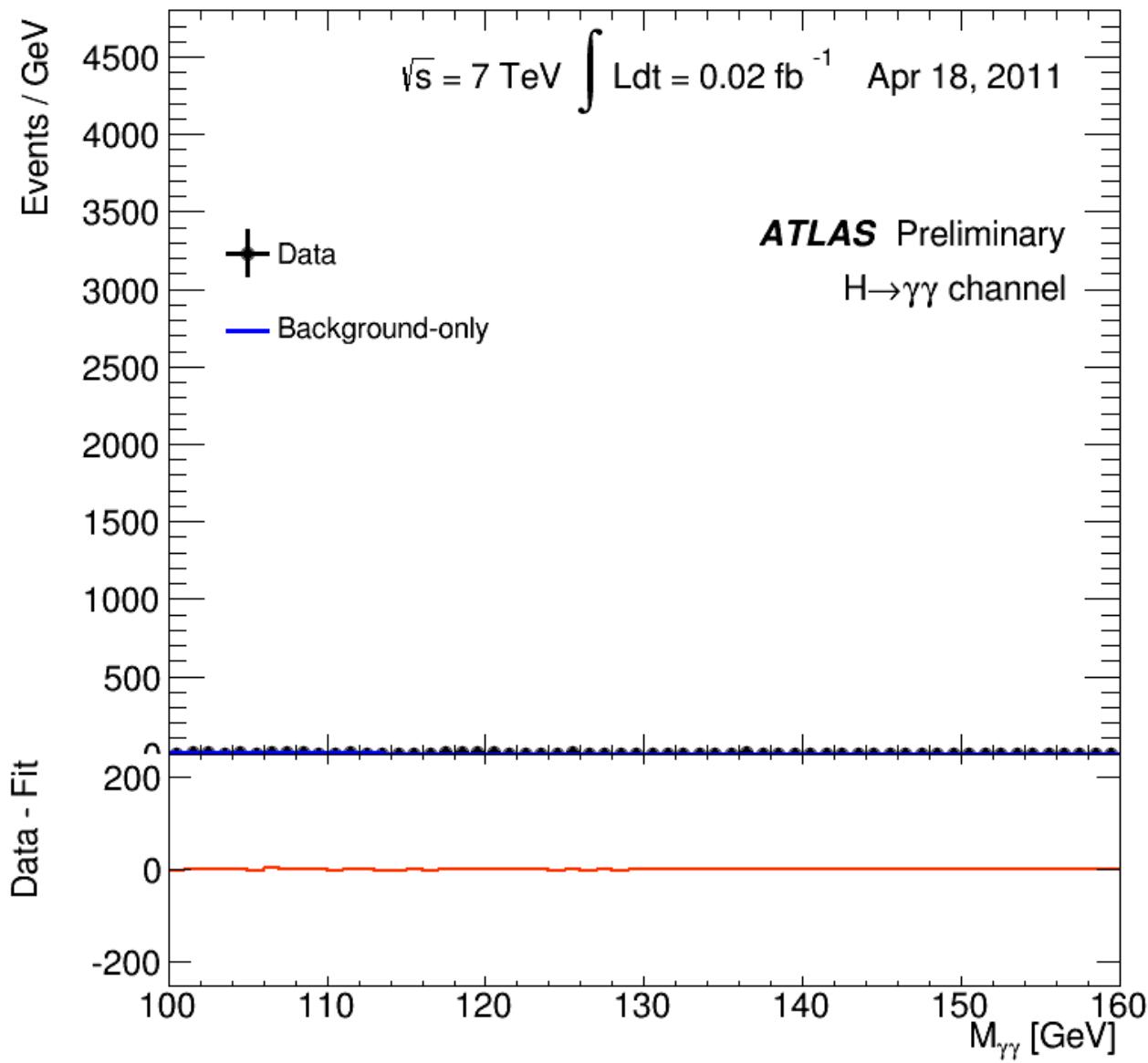
HAD calorimetry ($|\eta| < 5$): segmentation, hermeticity
Fe/scintillator Tiles (central), Cu/W-LAr (fwd)
Trigger and measurement of jets and missing E_T
E-resolution: $\sigma/E \sim 50\%/\sqrt{E} \oplus 0.03$

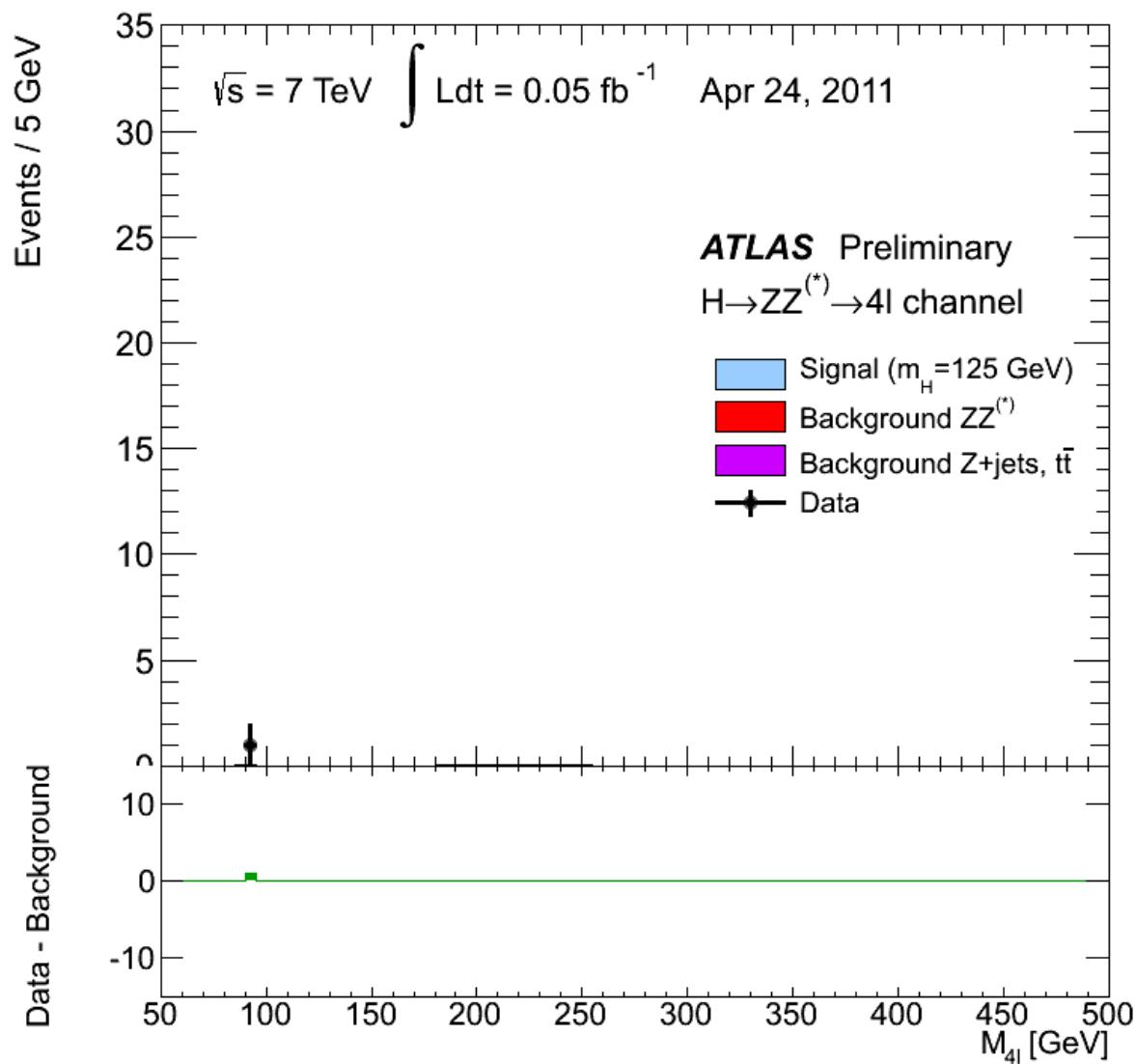
Length : ~ 46 m
Radius : ~ 12 m
Weight : ~ 7000 tons
 $\sim 10^8$ electronic channels
3000 km of cables

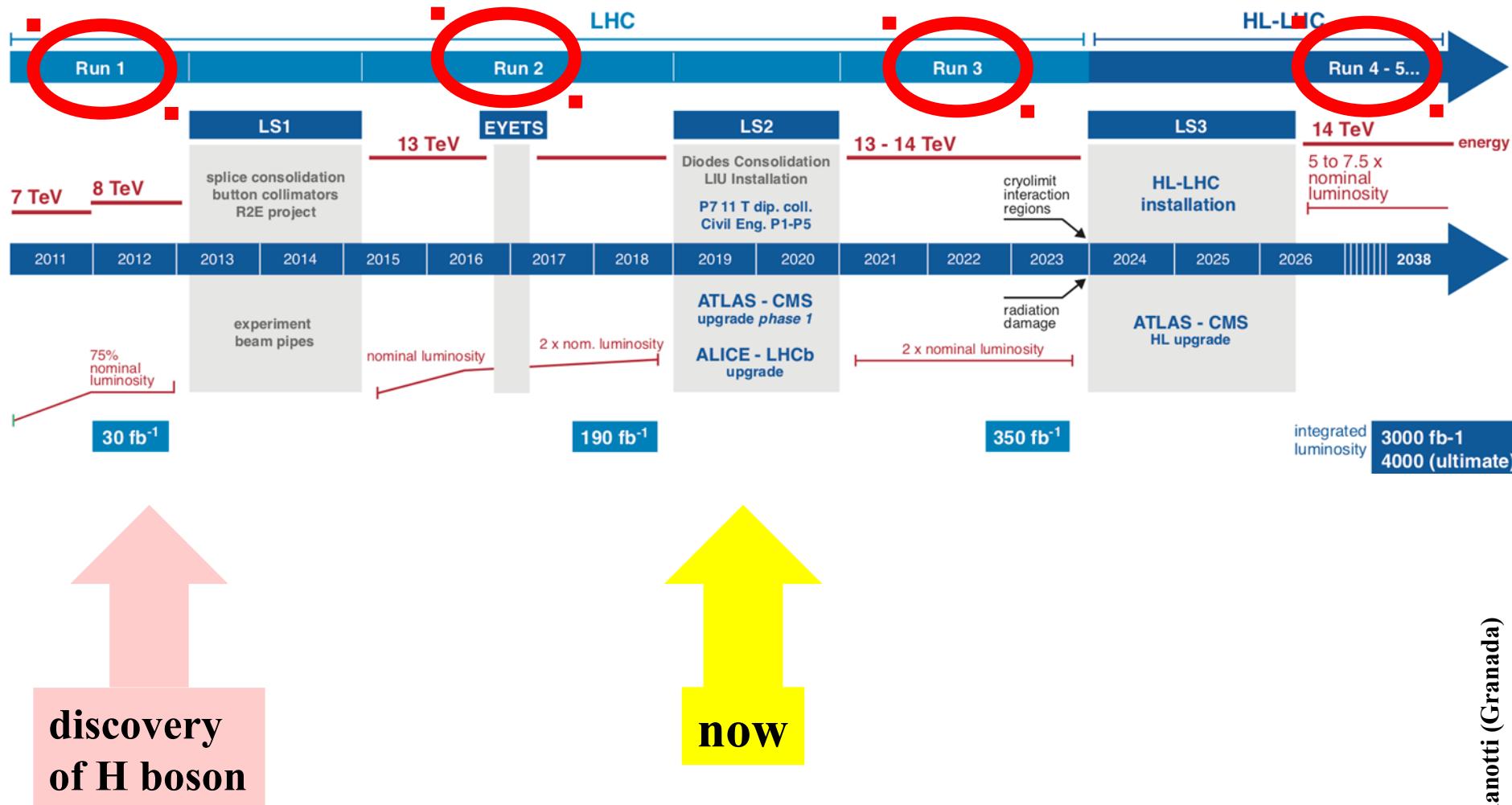


Corfou 2-9-2019

transverse and longitudinal
 segmentation of the EM
ATLAS (Liquid Argon)
 accordion calorimeter
 (very stable - about
 200 000 channels)







- ♪ *Historical introduction , Setting the stage*
- ♪ ***Results from (Run-1 and) Run-2***
- ♪ *Future of ATLAS , Run-3 , HL-LHC*
- ♪ *Conclusions*
- ♪ *Backup*

- *Results from (Run-1 and) Run-2*
 - * *detector*
 - * *SM*
 - * *BSM*
 - * *(B-E)H*
 - * *Vector-boson scattering*

short summary

1 > No new physics (yet) *outside the discovery of the H boson*

2> We are entering precision physics area

Large sample of various particles produced in Run-2

W bosons	$12 \cdot 10^9$
Z bosons	$3 \cdot 10^9$
Top	$300 \cdot 10^6$
B quarks	$40 \cdot 10^{12}$
BEH bosons	$8 \cdot 10^6$

♪ Results from (run 1 and) run 2

* *detector*

* *SM*

* *BSM*

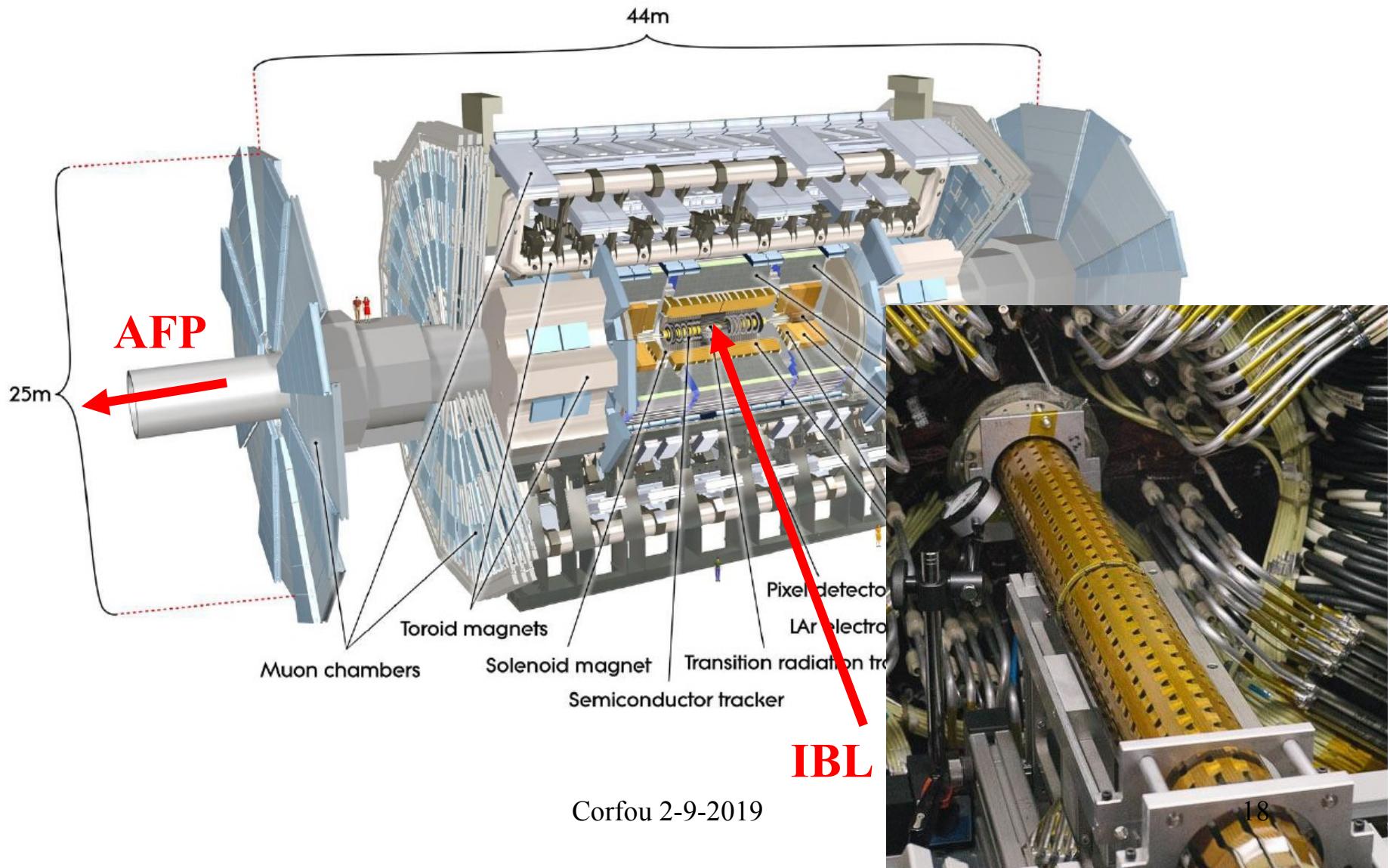
* *(B-E)H*

* *Vector-boson scattering*

ATLAS in Run-2

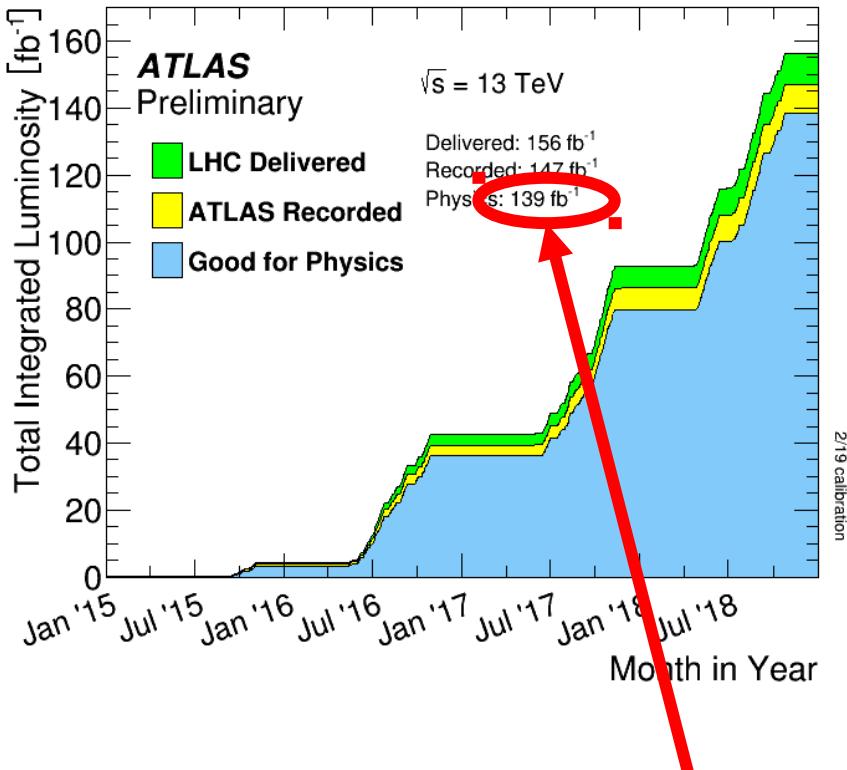
New detectors in Run-2:

- Innermost pixel layer IBL, 3.4cm from interaction point
- Forward proton detectors (one arm in 2016, 210m from IP)



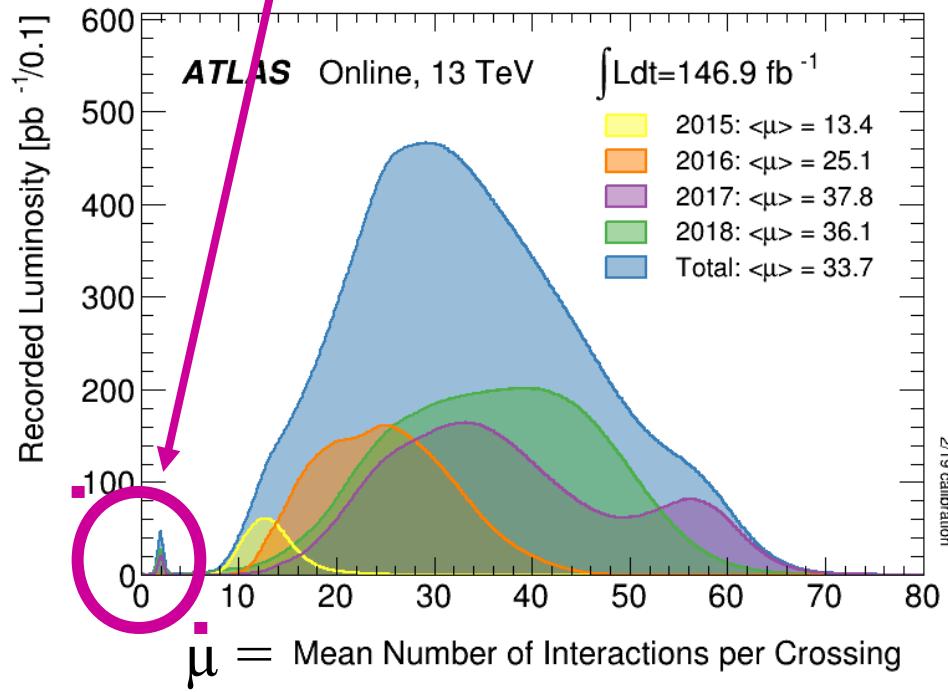
Integrated pp luminosity during Run-2

Also collected Pb-Pb p-Pd Xe-Xe data



measured to 1.7% precision
(ATLAS-CONF-2019-021)

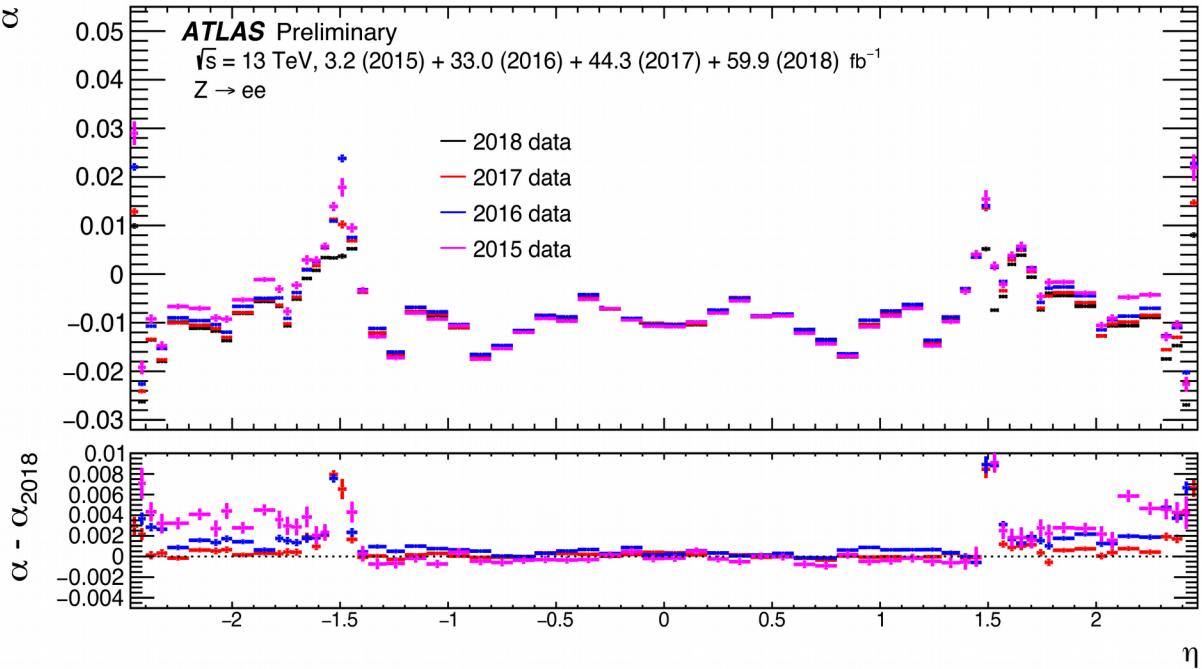
low μ data for high precision W physics



High-luminosity comes with a challenge

All dogmas need to be revisited

Like the fact that the response of the calorimeter is constant w.r.t time
(there are also short time-scale variations due to T change)



energy mis-calibration defined by α : $E^{\text{data}} = E^{\text{MC}}(1 + \alpha_i)$

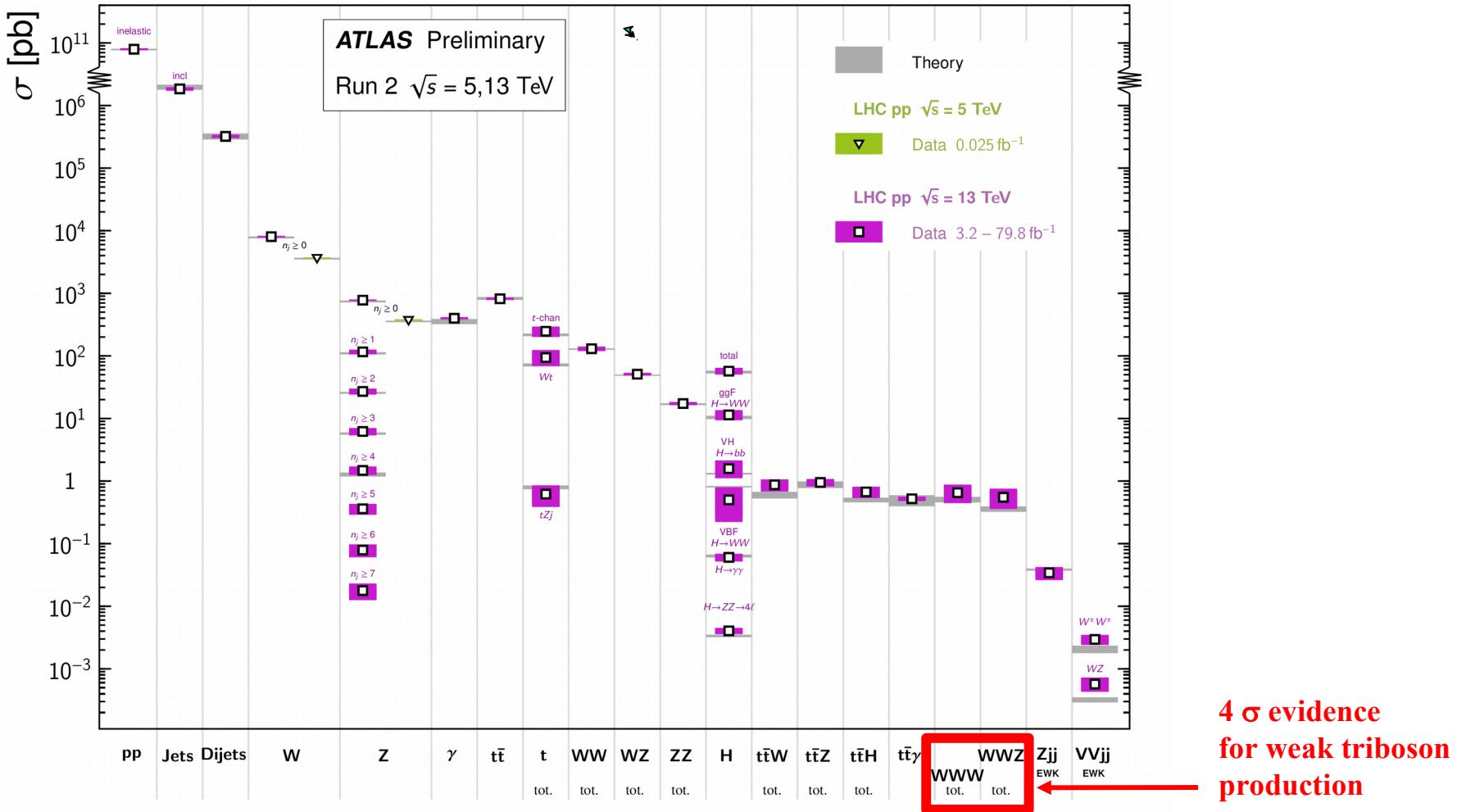
We have a lot of data in order to make precise calibrations

But the needs for precision physics are very important !

- ♪ Results from (run 1 and) run 2
 - * detector
 - * **SM**
 - * **BSM**
 - * **(B-E)H**
 - * **Vector-boson scattering**

Standard Model Production Cross Section Measurements

Status: July 2019

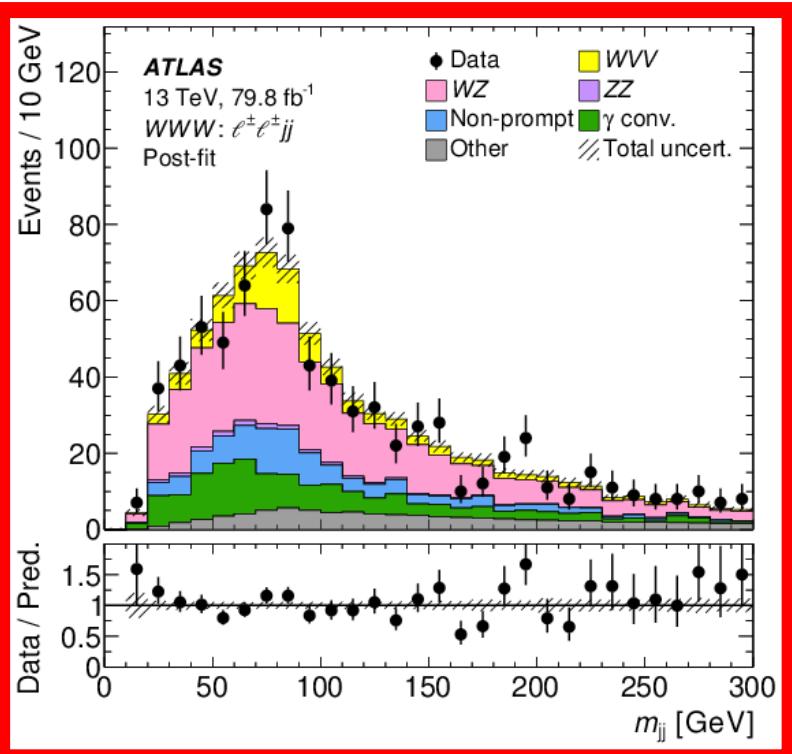
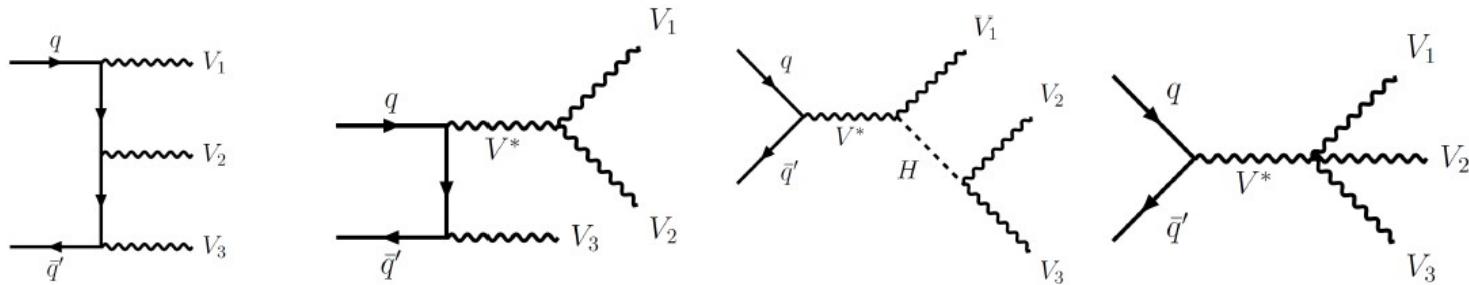


Theory agrees so far with the measured cross sections
on 15 orders of magnitude

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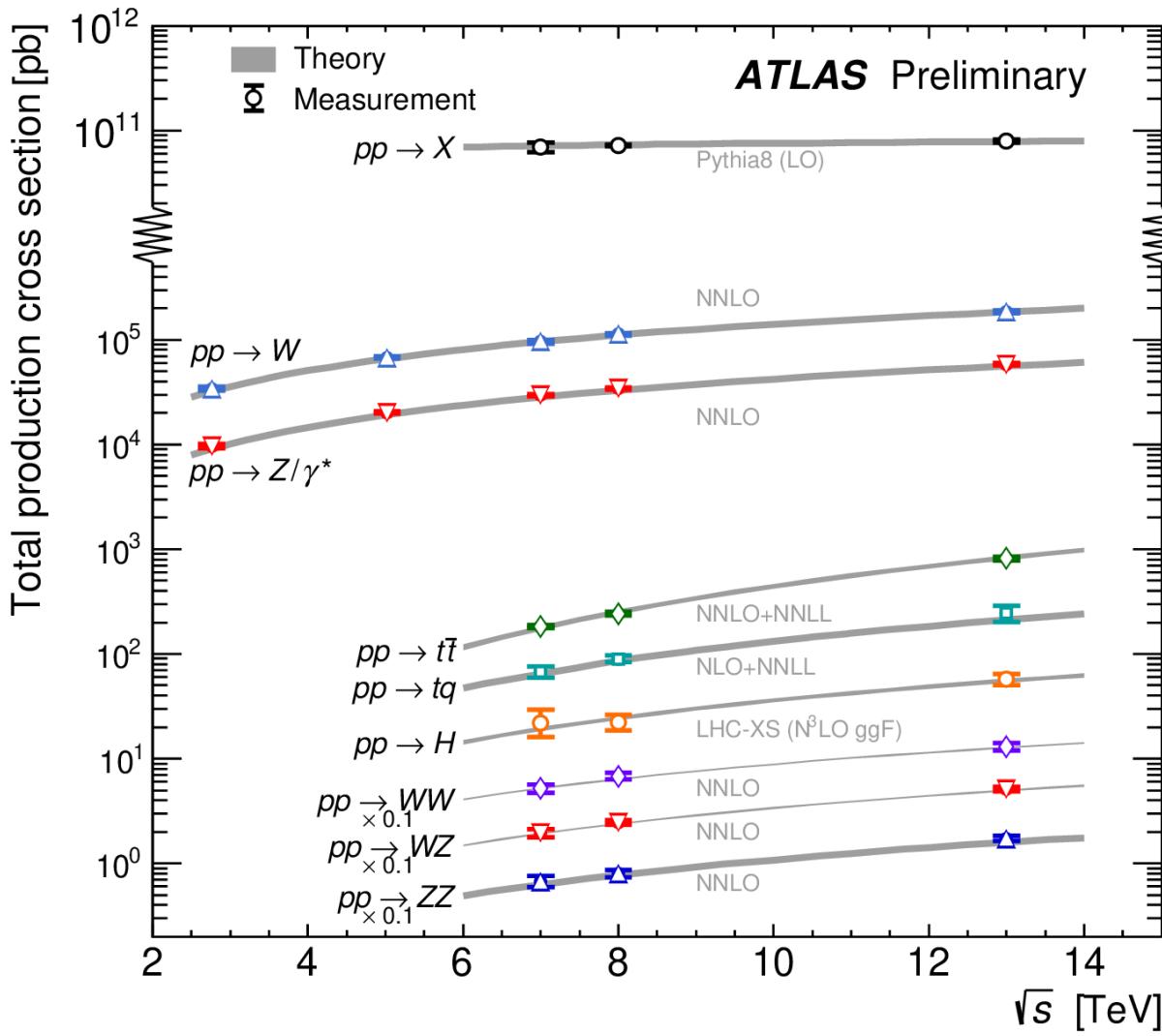
22

4 σ evidence for weak triboson production using 2015-2017 data



Decay channel	Significance	
	Observed	Expected
$WW\bar{W}$ combined	3.2σ	2.4σ
$WW\bar{W} \rightarrow \ell\nu\ell\nu qq$	4.0σ	1.7σ
$WW\bar{W} \rightarrow \ell\nu\ell\nu\ell\nu$	1.0σ	2.0σ
$WW\bar{Z}$ combined	3.2σ	2.0σ
$WW\bar{Z} \rightarrow \ell\nu qq\ell\ell$	0.5σ	1.0σ
$WW\bar{Z} \rightarrow \ell\nu\ell\nu\ell\ell/qq\ell\ell\ell\ell$	3.5σ	1.8σ
$WW\bar{V}$ combined	4.1σ	3.1σ

For different center-of-mass energies

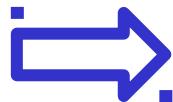


- \square $pp \rightarrow X$
7 TeV, $20\text{ }\mu\text{b}^{-1}$, Nat. Commun. 2, 463 (2011)
- \square $pp \rightarrow W$ \square $pp \rightarrow Z/\gamma^*$
8 TeV, $500\text{ }\mu\text{b}^{-1}$, Phys.Lett. B761 158 (2016)
- \square $pp \rightarrow Z/\gamma^*$
13 TeV, $60\text{ }\mu\text{b}^{-1}$, Phys. Rev. Lett. 117 182002 (2016)
- \triangle $pp \rightarrow W$ \square $pp \rightarrow Z/\gamma^*$
2.76 TeV, 4 pb^{-1} , arXiv:1907.03567 (for Z/W)
- \square $pp \rightarrow Z/\gamma^*$
5 TeV, 25 pb^{-1} , Eur. Phys. J. C79 (2019) 128 (for Z/W)
- \square $pp \rightarrow Z/\gamma^*$
7 TeV, 4.6 fb^{-1} , Eur. Phys. J. C77 (2017) 367 (for Z/W)
- \square $pp \rightarrow Z$
8 TeV, 20.2 fb^{-1} , JHEP 02, 117 (2017) (for Z)
- \square $pp \rightarrow W$
8 TeV, 20.2 fb^{-1} , arXiv:1904.05631 (for W)
- \square $pp \rightarrow W$
13 TeV, 81 pb^{-1} , PLB 759 (2016) 601 (for W)
- \square $pp \rightarrow Z$
13 TeV, 3.2 fb^{-1} , JHEP 02, 117 (2017) (for Z)
- \square $pp \rightarrow t\bar{t}$
7 TeV, 4.6 fb^{-1} , Eur. Phys. J. C 74:3109 (2014)
- \square $pp \rightarrow t\bar{t}$
8 TeV, 20.3 fb^{-1} , Eur. Phys. J. C 74:3109 (2014)
- \square $pp \rightarrow t\bar{t}$
13 TeV, 3.2 fb^{-1} , Phys. Lett. B 761 (2016)
- \square $pp \rightarrow tq$
7 TeV, 4.6 fb^{-1} , PRD 90, 112006 (2014)
- \square $pp \rightarrow tq$
8 TeV, 20.3 fb^{-1} , Eur. Phys. J. C 77 (2017) 531
- \square $pp \rightarrow tq$
13 TeV, 3.2 fb^{-1} , JHEP 1704 (2017) 086
- \square $pp \rightarrow H$
7 TeV, 4.5 fb^{-1} , Eur. Phys. J. C76 (2016) 6
- \square $pp \rightarrow H$
8 TeV, 20.3 fb^{-1} , Eur. Phys. J. C76 (2016) 6
- \square $pp \rightarrow H$
13 TeV, 36.1 fb^{-1} , Phys. Lett. B 786 (2018) 114
- \square $pp \rightarrow WW$
7 TeV, 4.6 fb^{-1} , PRD 87, 112001 (2013)
- \square $pp \rightarrow WW$
8 TeV, 20.3 fb^{-1} , JHEP 09 029 (2016)
- \square $pp \rightarrow WW$
13 TeV, 36.1 fb^{-1} , arXiv:1905.04242
- \square $pp \rightarrow WZ$
7 TeV, 4.6 fb^{-1} , Eur. Phys. J. C (2012) 72:2173
- \square $pp \rightarrow WZ$
8 TeV, 20.3 fb^{-1} , PRD 93, 092004 (2016)
- \square $pp \rightarrow WZ$
13 TeV, 36.1 fb^{-1} , arXiv:1902.05759
- \triangle $pp \rightarrow ZZ$
7 TeV, 4.6 fb^{-1} , JHEP 03, 128 (2013)
- \triangle $pp \rightarrow ZZ$
8 TeV, 20.3 fb^{-1} , JHEP 01, 099 (2017)
- \triangle $pp \rightarrow ZZ$
13 TeV, 36.1 fb^{-1} , Phys. Rev. D 97 (2018) 032005

Measurement of fiducial and differential W+ W- production

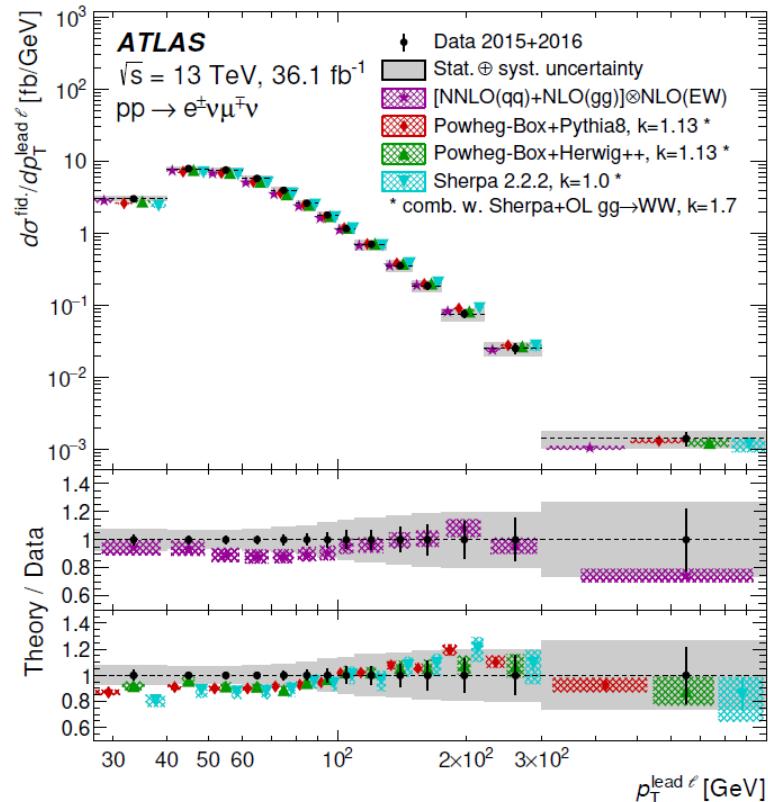
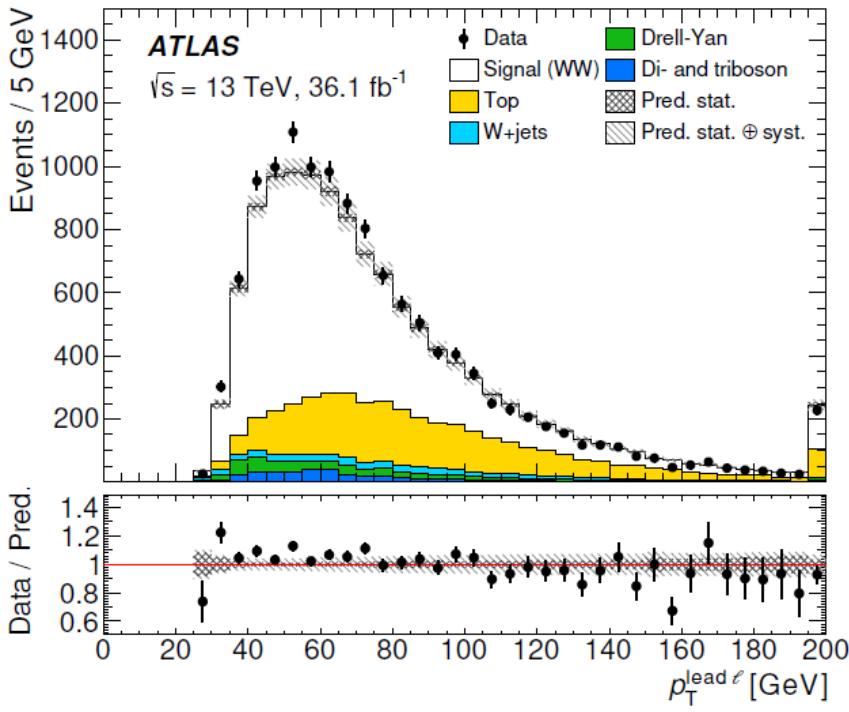
cross-sections at $\sqrt{s}=13$ TeV with the ATLAS detector

Suppress
top-quark background



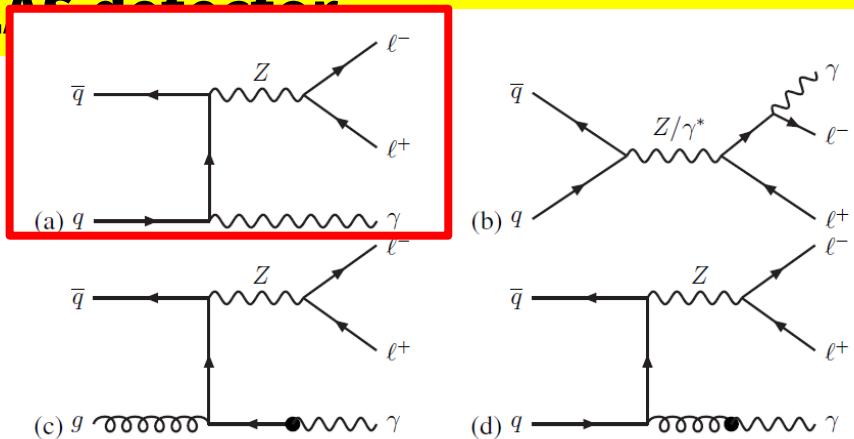
events containing jets with a transverse momentum exceeding 35 GeV are not included in the measurement phase space

$$WW \rightarrow e^\pm \nu \mu^\mp \nu$$

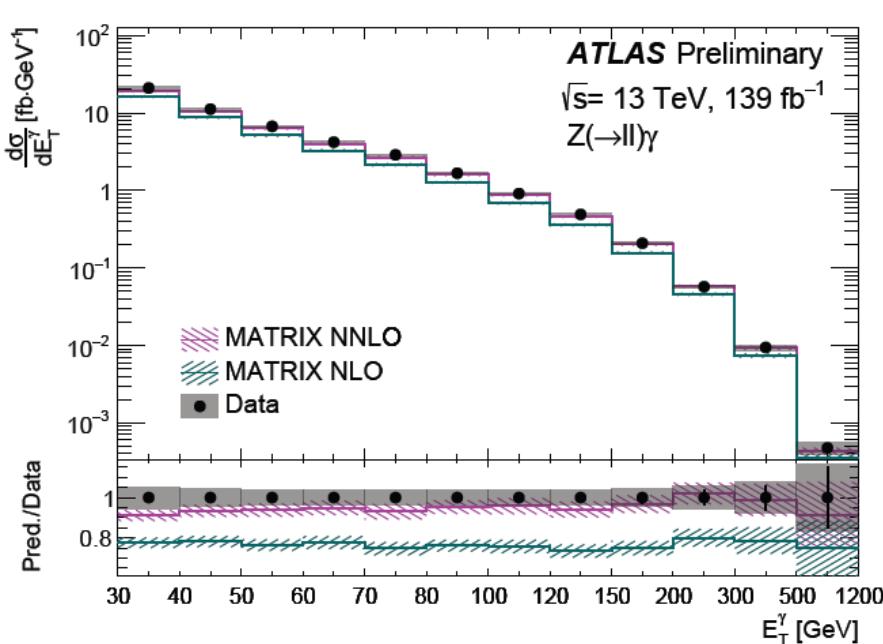
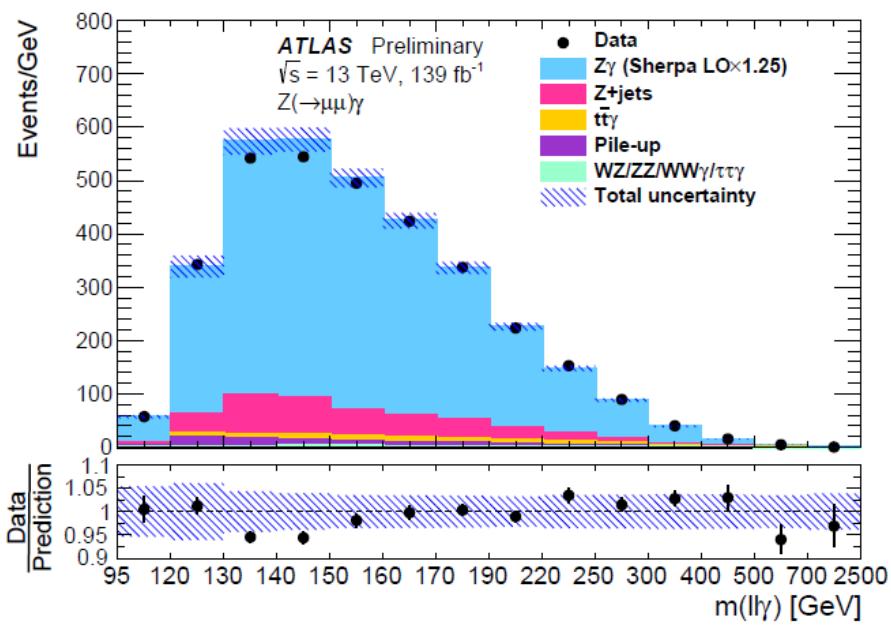


Measurement of $Z(\rightarrow l^+l^-)\gamma$ differential cross-sections in pp collisions at $\sqrt{s} = 13$ TeV with the ATLAS detector

Full Run-2



FSR and
quark/gluon
fragmentation
removed (isolation)



Inclusive and differential measurements of the charge asymmetry in $t\bar{t}$ events at the Tevatron-inspired detector

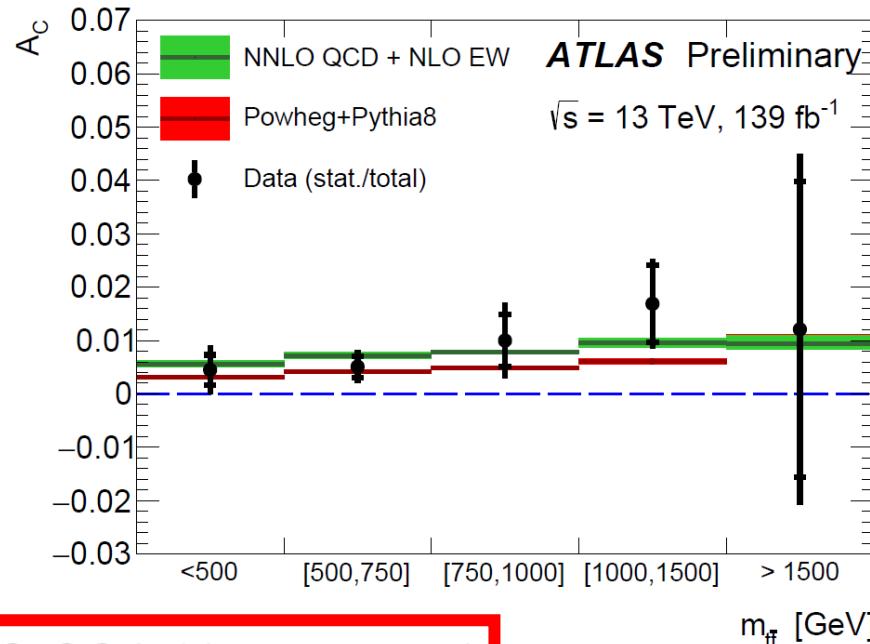
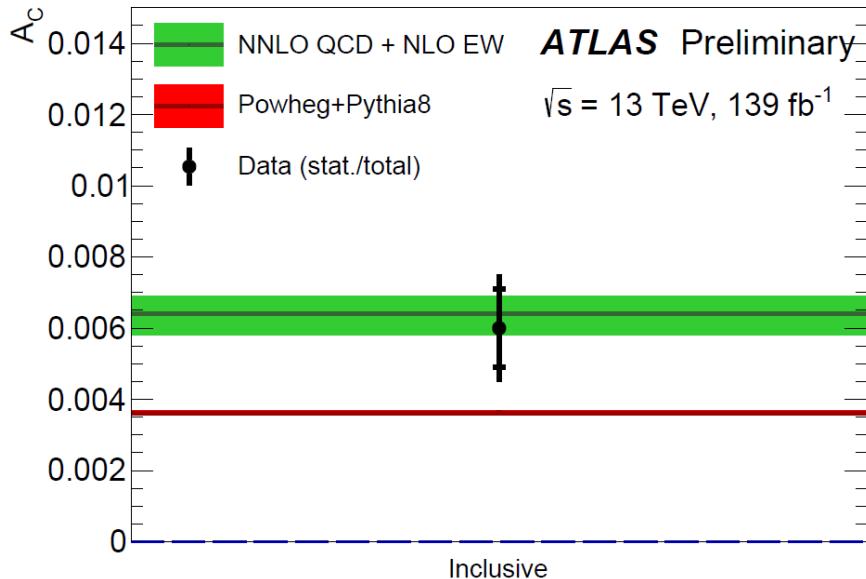
central-forward charged asymmetry is defined

Full Run-2

$$A_C = \frac{N(\Delta|y| > 0) - N(\Delta|y| < 0)}{N(\Delta|y| > 0) + N(\Delta|y| < 0)}$$

$$\Delta|y| = |y_t| - |\bar{y}_t|$$

different from 0 because of interference at NLO

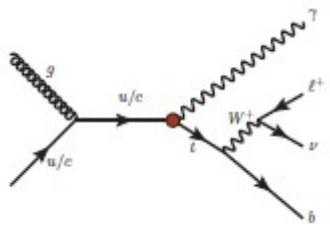


$$A_C = 0.0060 \pm 0.0015(\text{stat+syst.})$$

FCNC (Flavour-Changing Neutral Current)

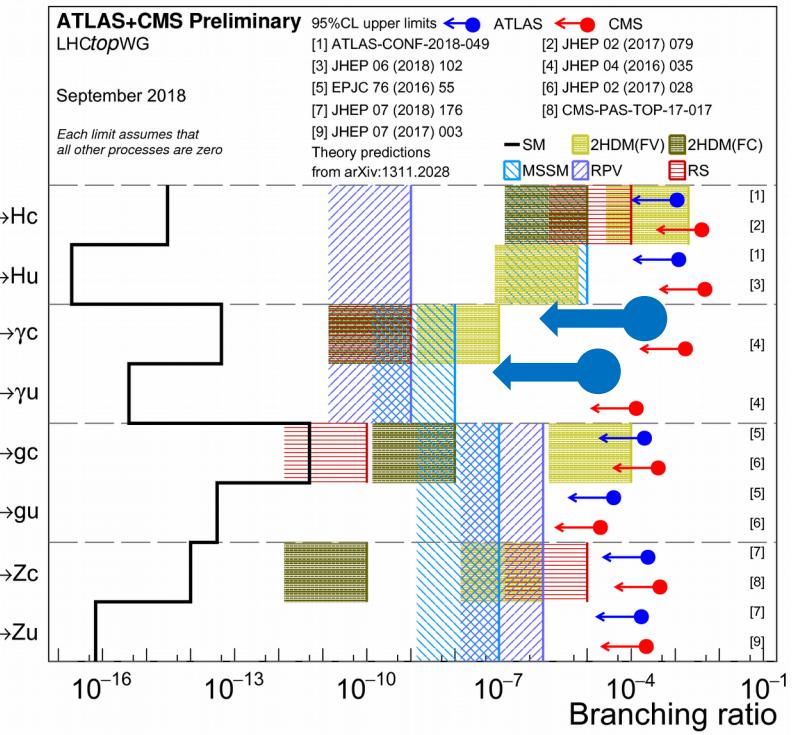
window for new physics

tq γ 80 fb $^{-1}$ 13 TeV



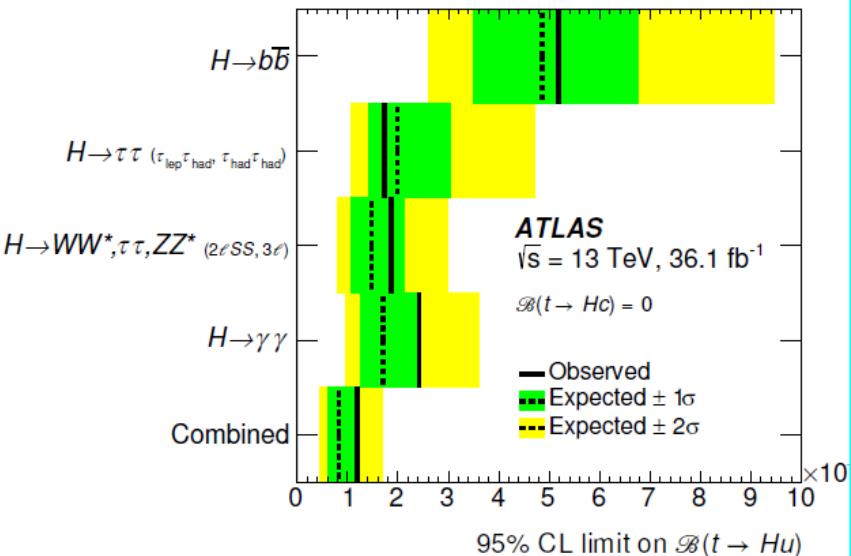
topology
1 e/ μ , 1 γ
1 b-jet, MET

NN to separate S and B



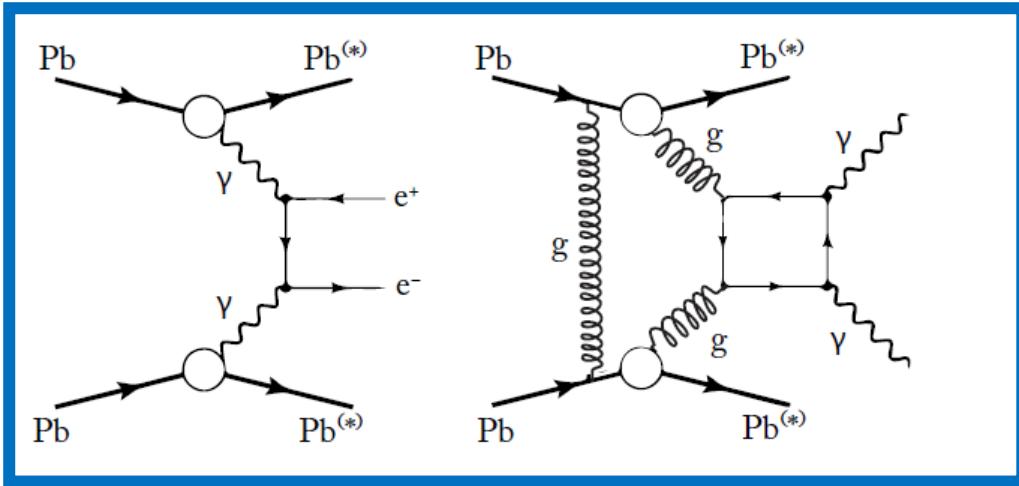
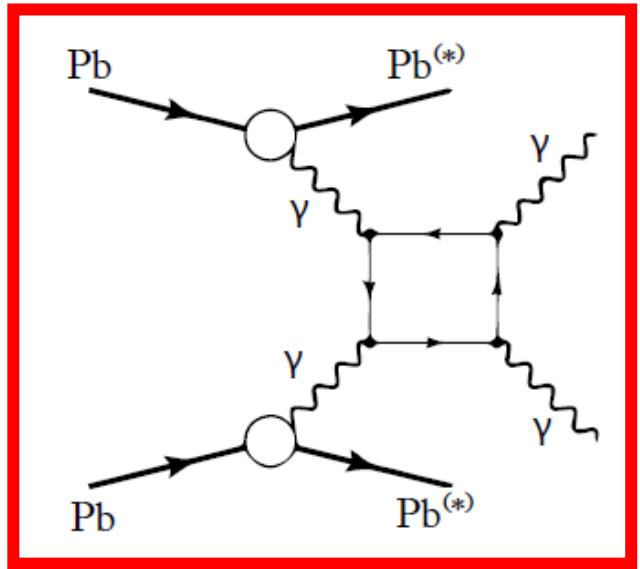
B(t → γc) < $2.2 \cdot 10^{-4}$ ($2.7 \cdot 10^{-4}$ exp)
 B(t → γu) < $2.8 \cdot 10^{-5}$ ($4.0 \cdot 10^{-5}$ exp)

t → Hq 36fb $^{-1}$



B(t → Hu) < $1.2 \cdot 10^{-3}$ ($0.8 \cdot 10^{-4}$ exp)

Observation of light-by-light scattering in ultraperipheral Pb+Pb collisions with the ATLAS detector



Signal : 2 photons with very low $p_T(\gamma\gamma)$

: $< 1 \text{ GeV}/c$ and no further activity in the detector

Field strength up to 10^{25} V/m

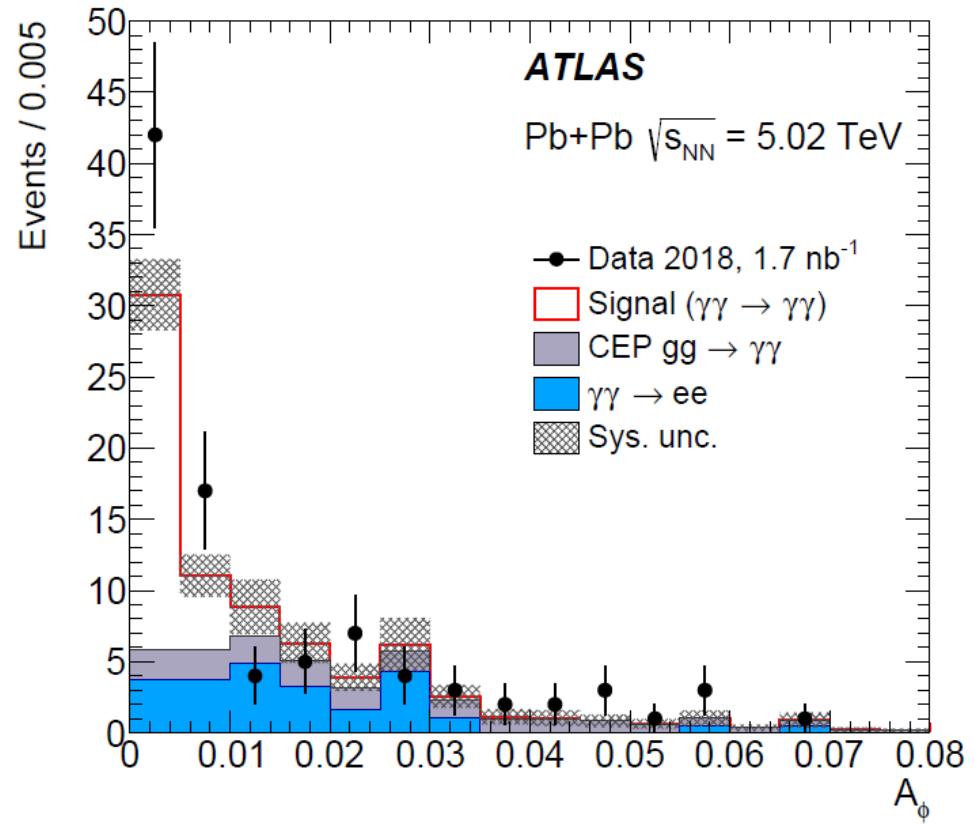
$\gamma\gamma$ luminosity $\sim [Z^4] \sim 5 \cdot 10^7$

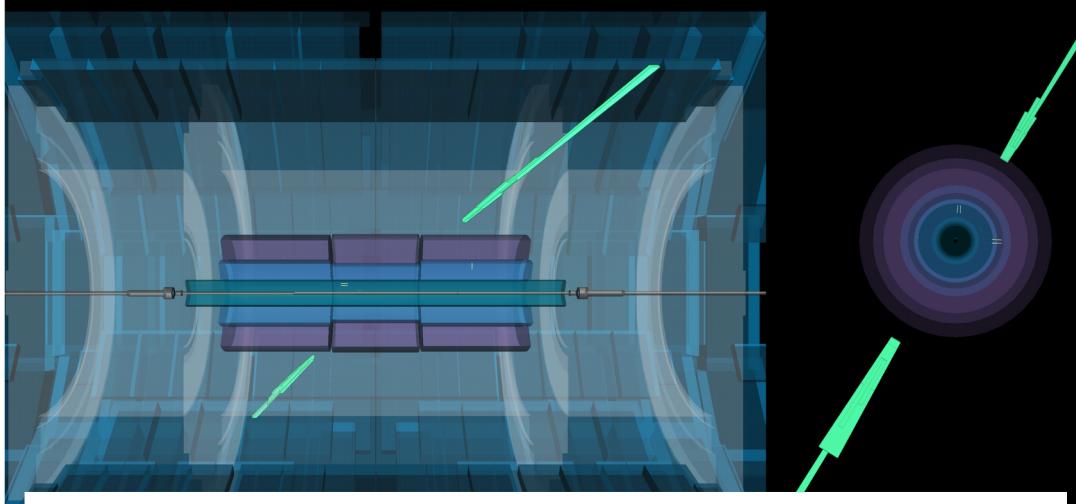
backgrounds

Look at low-energy back-to-back photon pairs with no additional activity in detector

small acoplanarity

$$A_\phi = (1 - |\Delta\phi_{\gamma\gamma}|/\pi) < 0.01$$



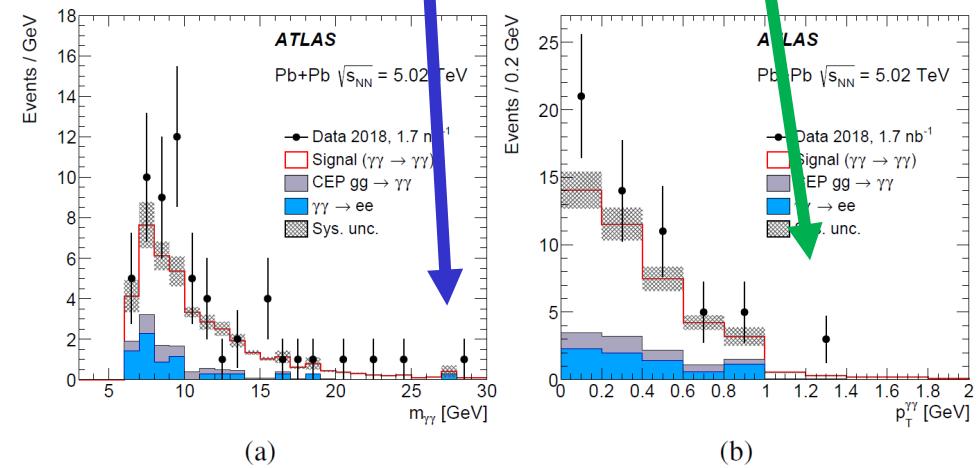


no additional activity in the detector

59 $\gamma\gamma \rightarrow \gamma\gamma$ events observed
 with an expected background
 of 12 ± 3 (8.2σ)

**Event display for an
 exclusive $\gamma\gamma \rightarrow \gamma\gamma$
 candidate.**

**Two back-to-back
 photons ($E_T^{\gamma 1} = 11$ GeV
 and $E_T^{\gamma 2} = 10$ GeV) with
 $m_{\gamma\gamma} = 29$ GeV, $A_\phi = 0.002$
 $P_T^{\pi} = 1.2$ GeV**



EW precision measurements

Weak angle $\sin^2\theta_{\text{eff}}^l$

8 TeV data

W mass m_W

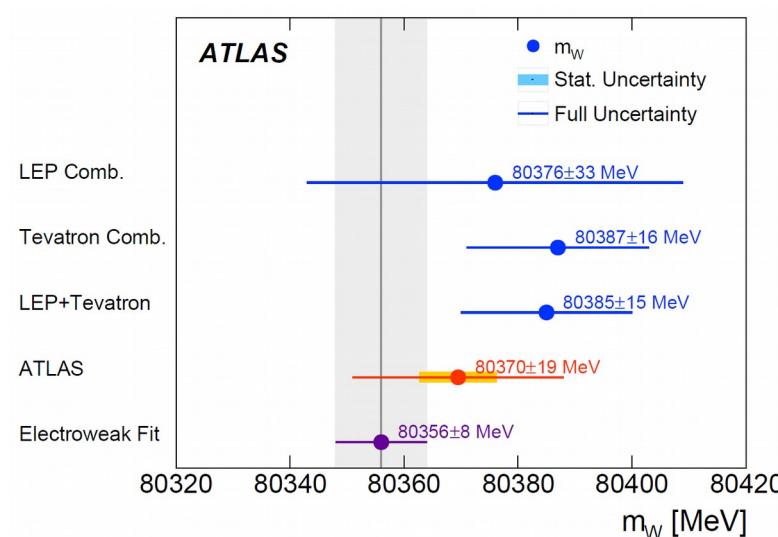
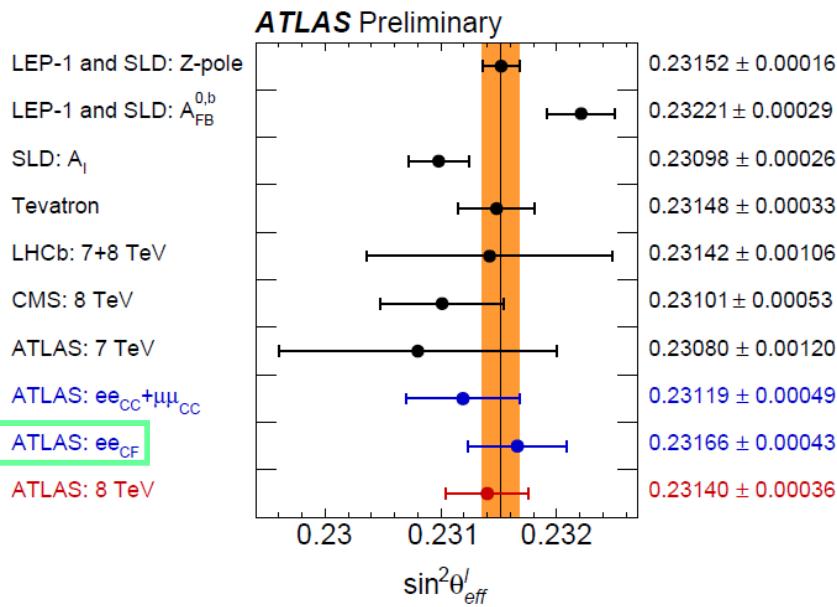
7 TeV data

One wants to have measurements with uncertainties close to the results of the EW fit $\sin^2\theta_{\text{eff}}^l = .23153 \pm .00006$ $m_W = 80354 \pm 7 \text{ MeV}$

arXiv:1803.01853

F-B asymmetry

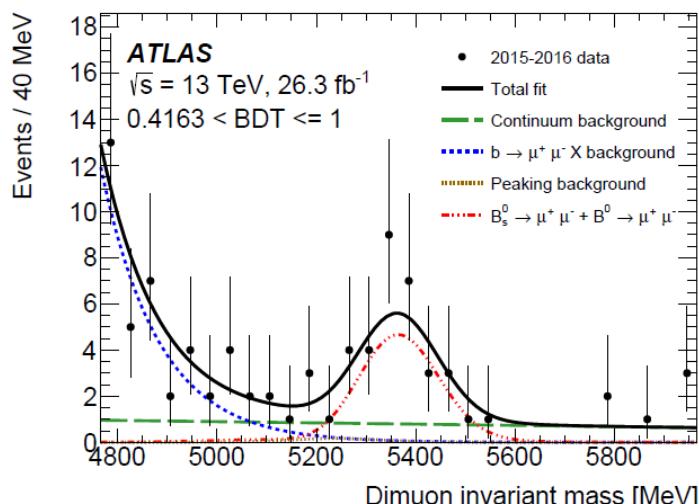
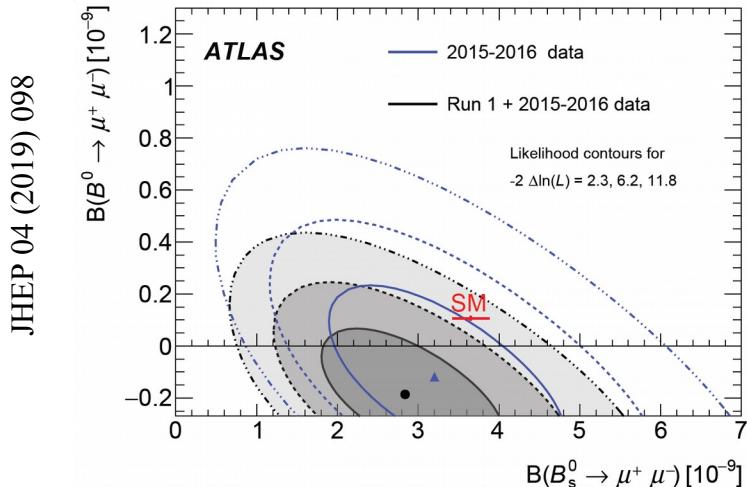
$q\bar{q} \rightarrow Z/\gamma^* \rightarrow \ell^+ \ell^-$



Run-1 + 26fb⁻¹ of Run-2

$$\mathcal{B}(B_s^0 \rightarrow \mu^+ \mu^-) = (2.8^{+0.8}_{-0.7}) \times 10^{-9}$$

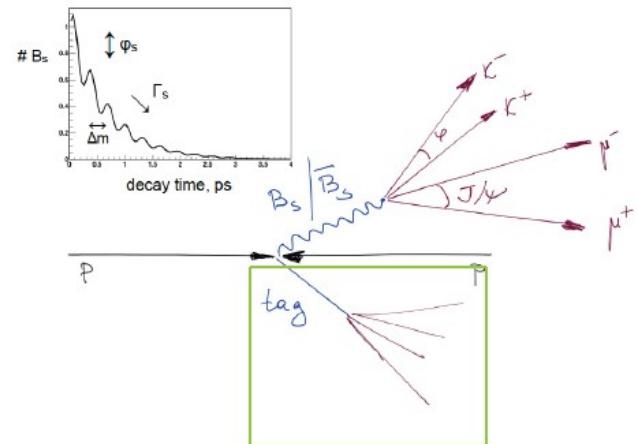
$$\mathcal{B}(B_s^0 \rightarrow \mu^+ \mu^-) < 2.1 \times 10^{-10} \text{ 95% CL}$$



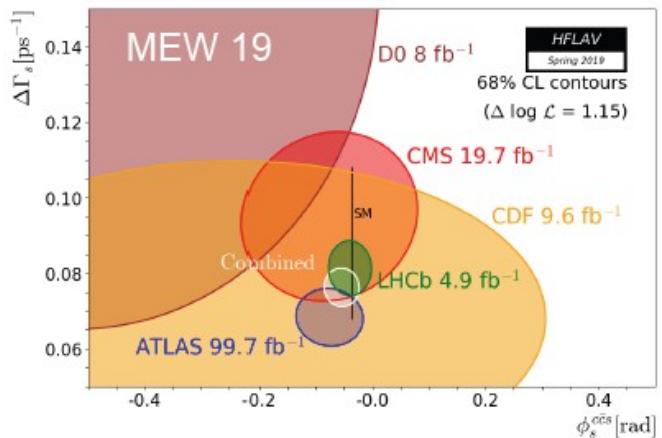
B physics

$$B_s^0 \rightarrow J/\psi \phi$$

Measurement of CP-violating phase ϕ_s

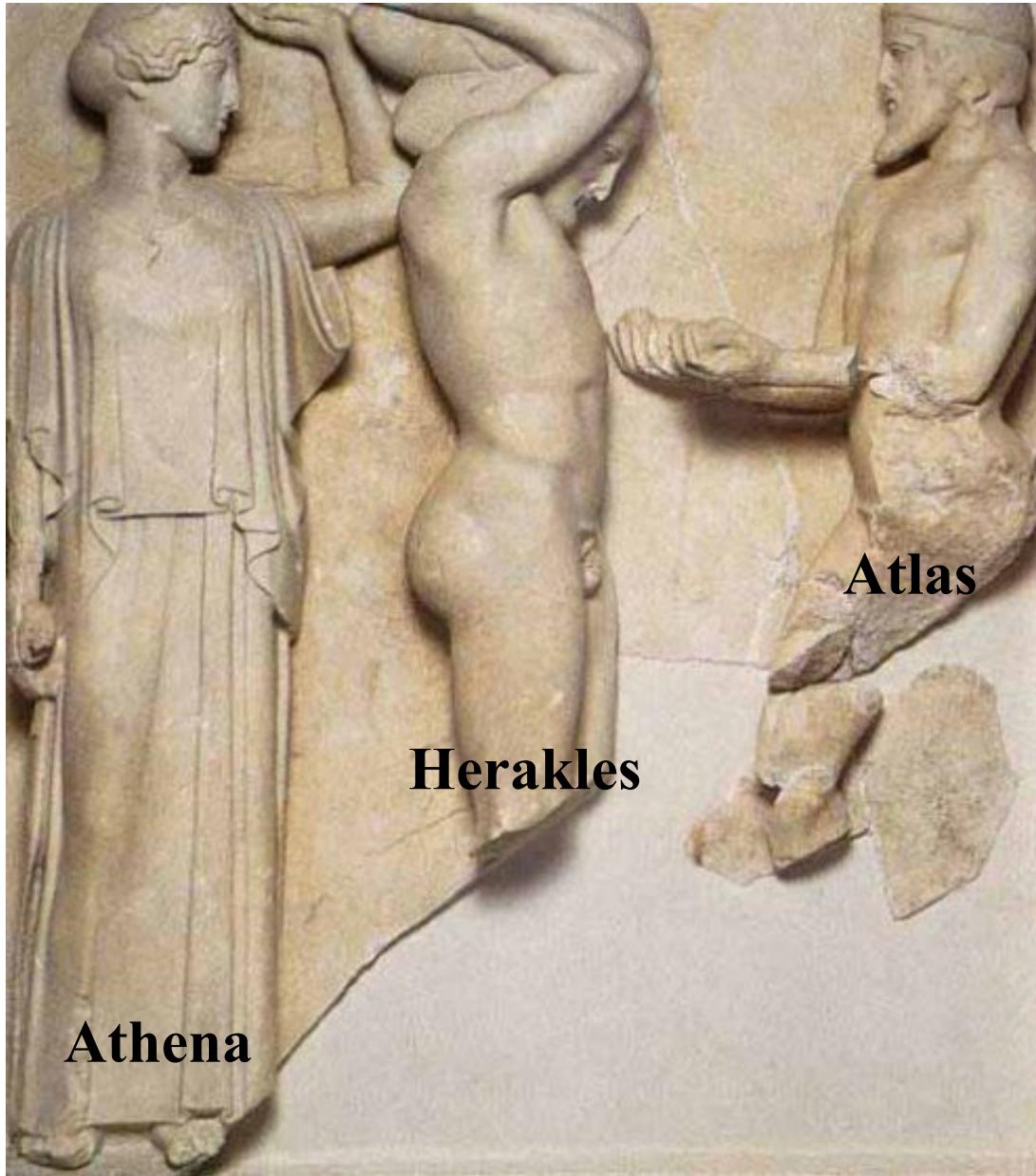


$$\phi_s = -0.076 \pm 0.034 \text{ (stat.)} \pm 0.019 \text{ (syst.) rad}$$



- ♪ *Results from (run 1 and) run 2*
- * *detector*
- * *SM*
- * ***BSM***
- * *(B-E)H*
- * *Vector-boson scattering*

interlude



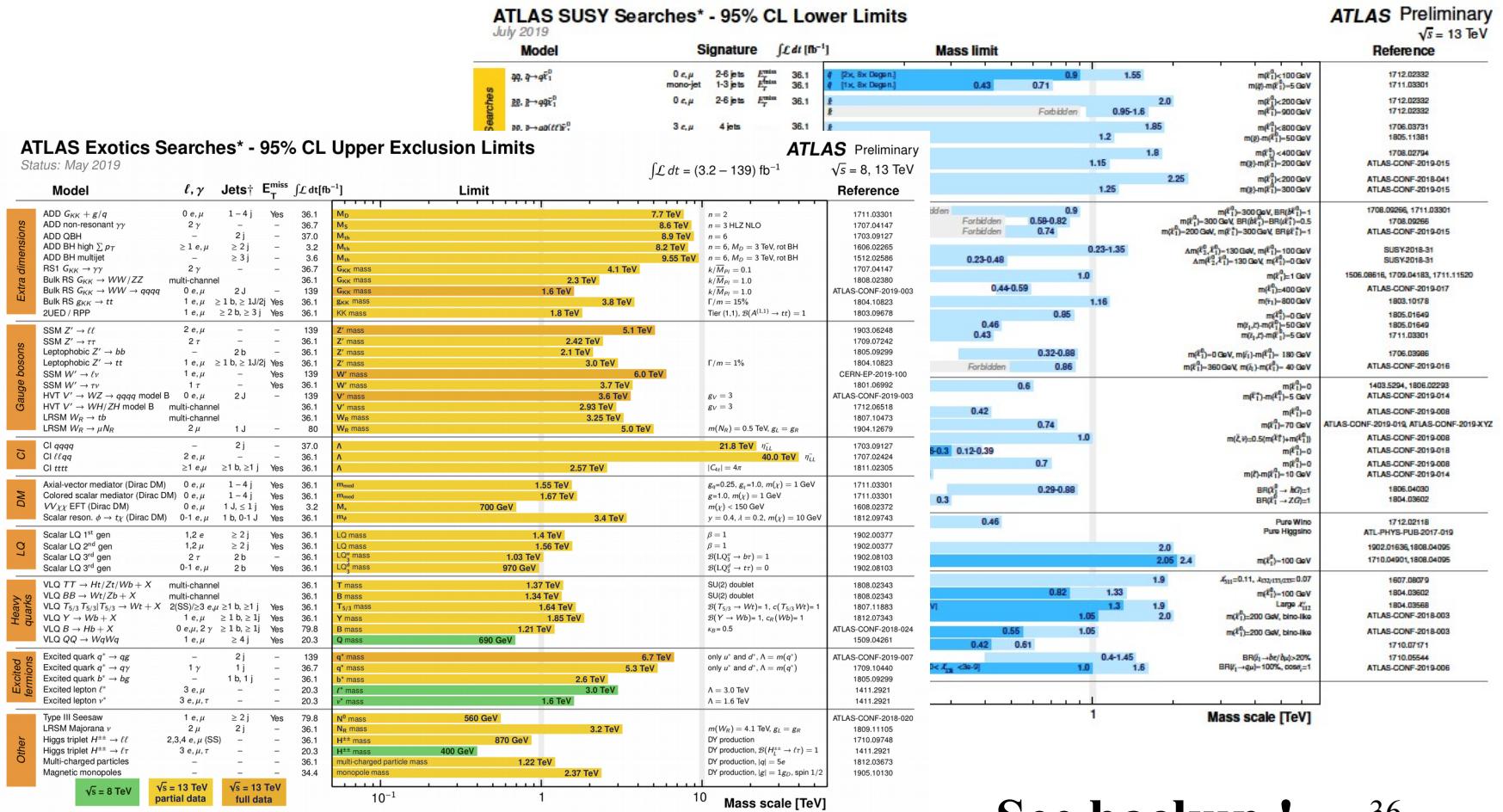
**Search for
Physics BSM
(Beyond the
Standard
Model)**

**Some
(temporary)
rest**

**in a lot of places
Full Run-2 results**

Search for Physics BSM (1)

**A vast programme covering searches for high and low mass particles,
small couplings, long lived particles, forbidden decays, ...**



*Only a selection of the available mass limits on new states or phenomena is shown.

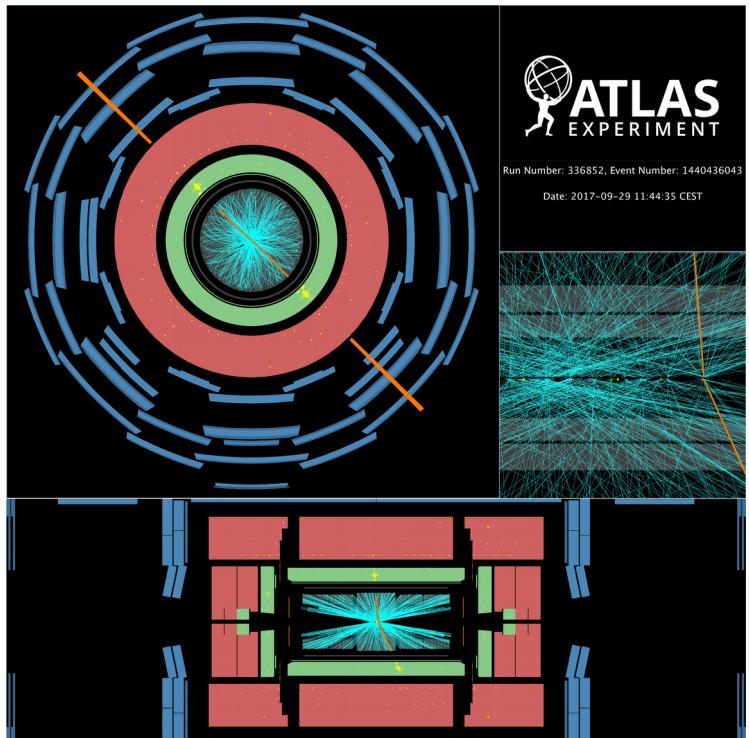
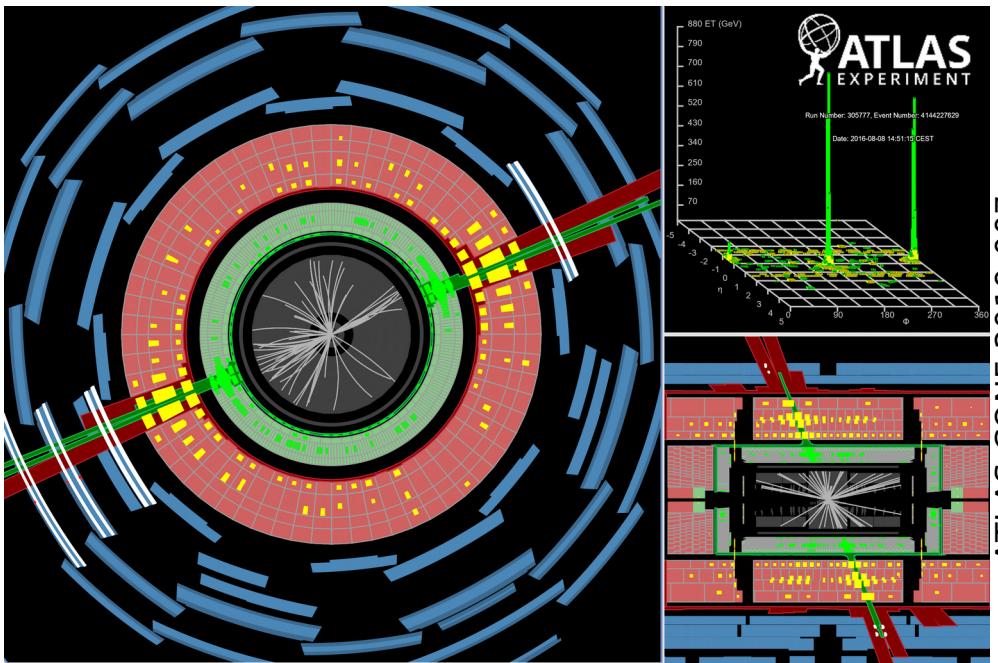
[†]Small-radius (large-radius) jets are denoted by the letter j (J).

See backup !

Search for Physics BSM (2)

Full Run-2

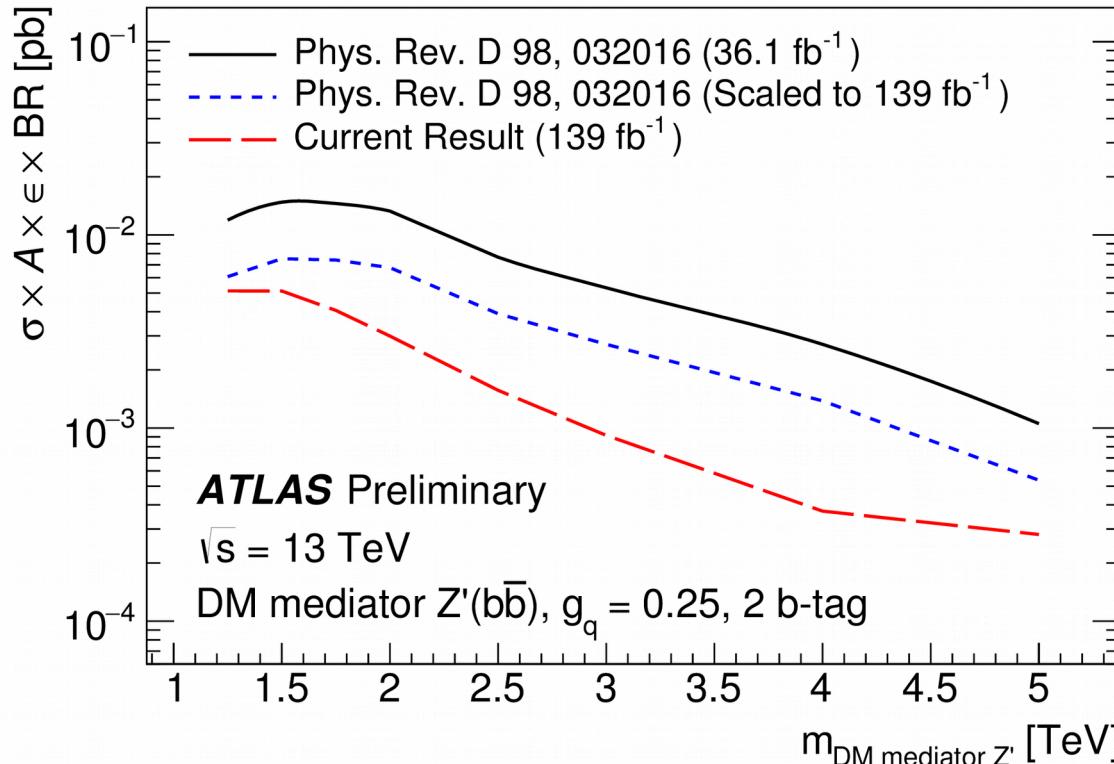
highest-mass **dijet** event
the two central high- p_T jets
each have p_T of 3.74 TeV
their invariant mass is
8.02 TeV.



dielectron candidate with the
highest invariant mass in the
2015-2018 data taking period
with $m_{ee} = 4.06$ TeV
search for Z' and W'

Analysis and detector improvements very important !

2 b-jets



width/mass = 3%

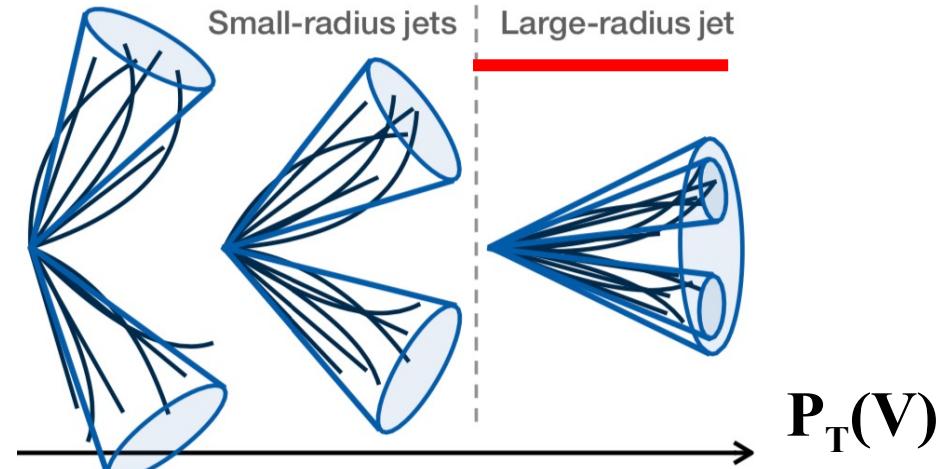
©Nishu priv com

Significant improvement in b-tagging performance at high p_T

Resonances decaying to VV (WW, WZ, ZZ)

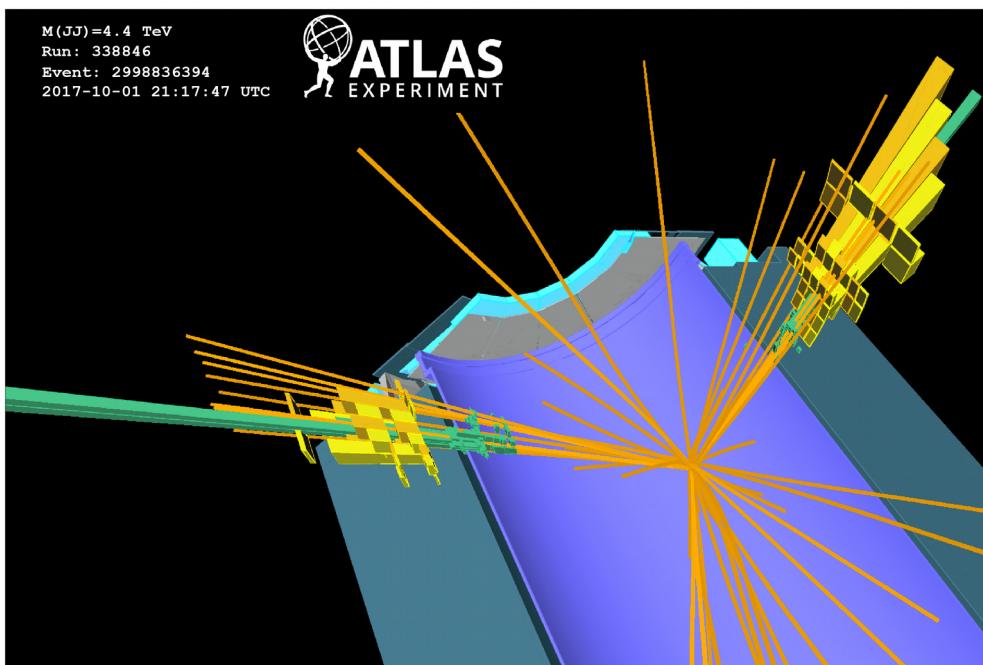
The diboson system is reconstructed using pairs of high transverse momentum, large-radius jets

diboson resonances with masses greater than 1.3 TeV



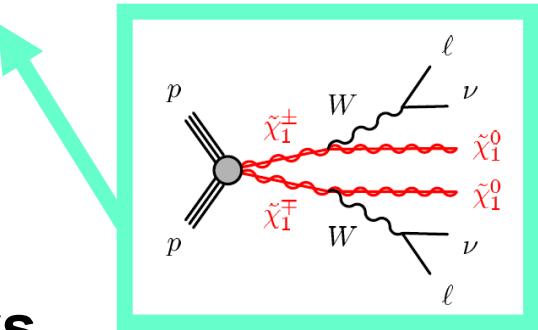
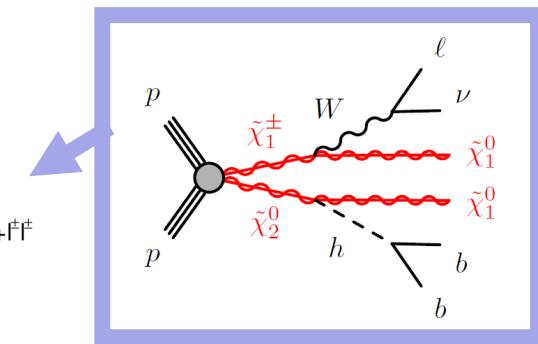
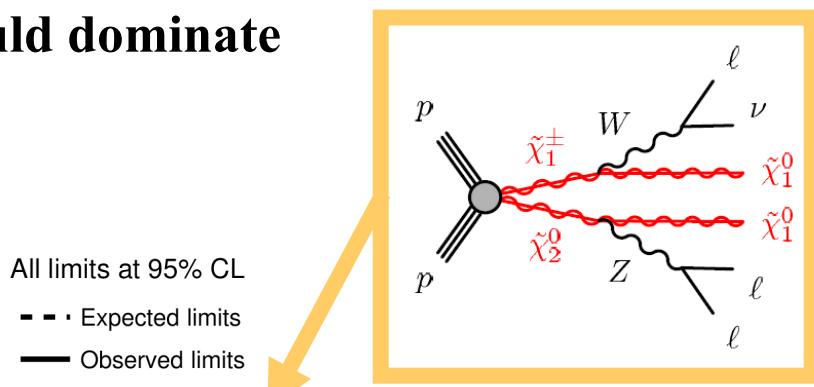
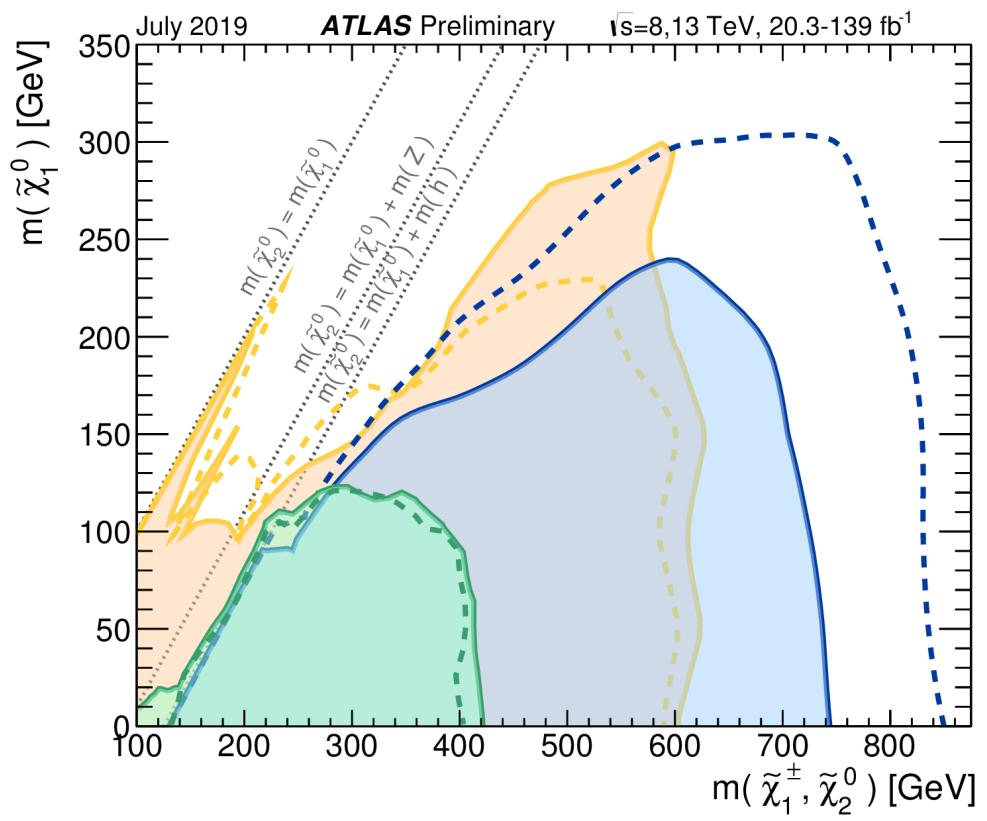
Highest m_{JJ} (=4440 GeV)
diboson candidate

The leading (*subleading*) jet has a p_T of 2136 GeV (2291 GeV), a mass of 89.5 GeV (62.5 GeV)



Search for Physics BSM (5)

SUSY Electroweak production (could dominate if squarks and gluinos heavy)



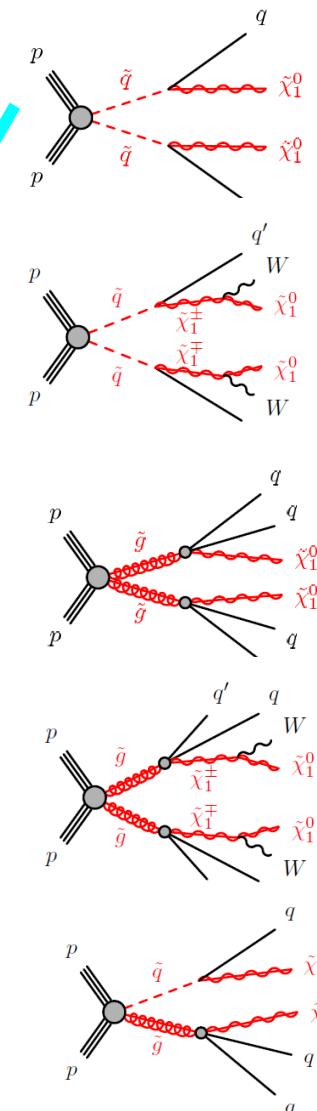
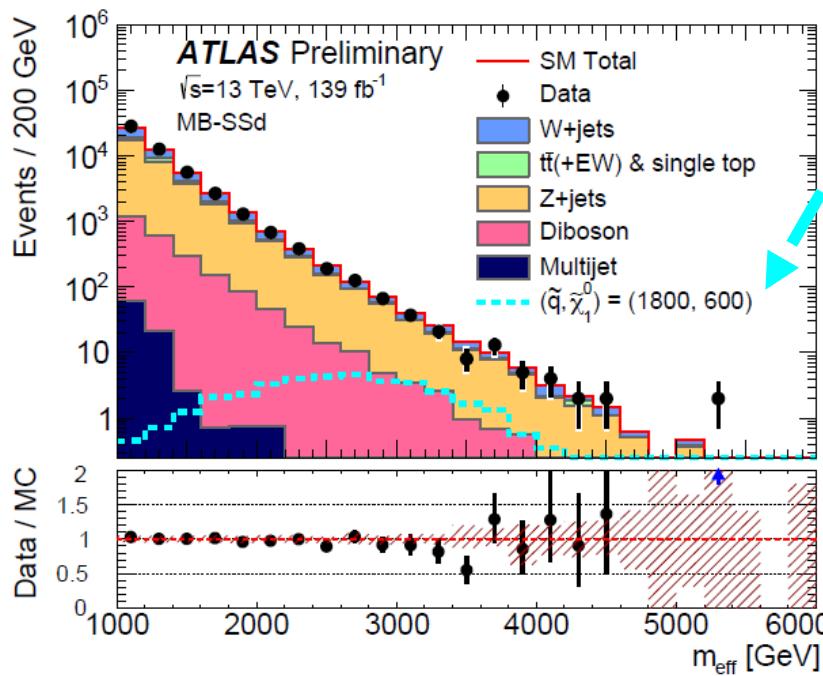
Example : weak-boson-mediated decays

SUSY Strong production of Squarks and Gluinos Golden mode

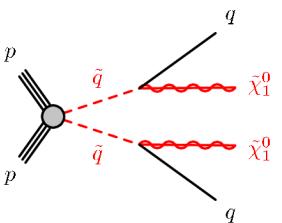
Sensitive searches for squarks and gluinos (in R-parity conserving scenarios) with neutralino as LSP (*no leptons*)

Many different scenarios investigated
with cut-based analyses and BDTs

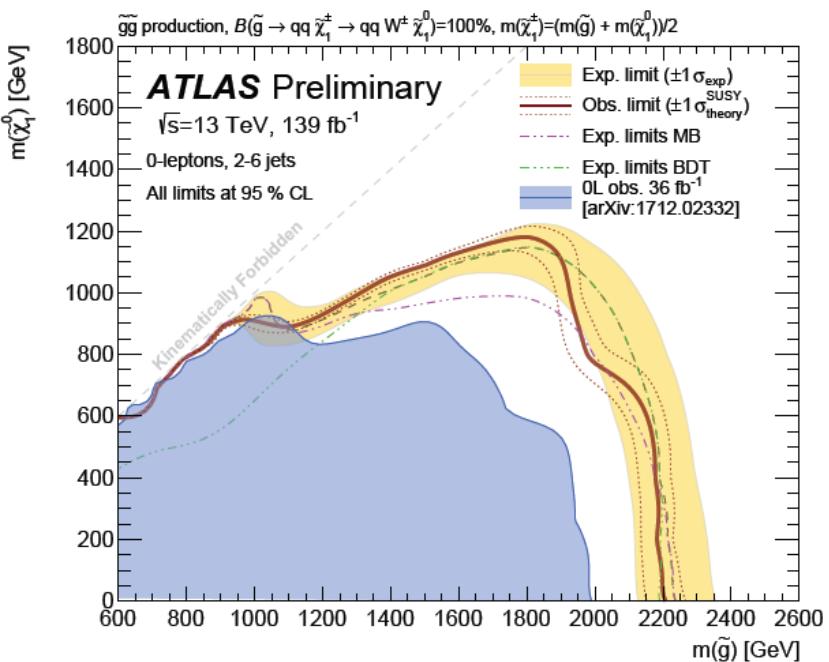
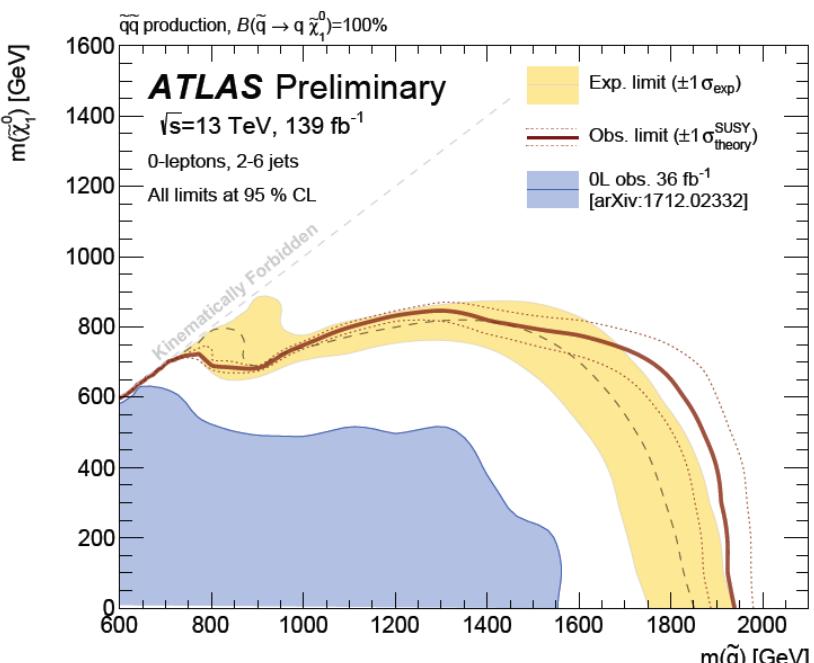
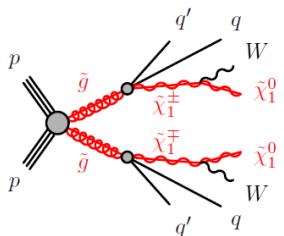
$$M_{\text{eff}} = \text{sum} (\text{p}_T \text{jets} > 50 \text{ GeV} + E_T^{\text{miss}})$$



Search for Physics BSM (7)



Used simplified scenarios



Search for Physics BSM (8)

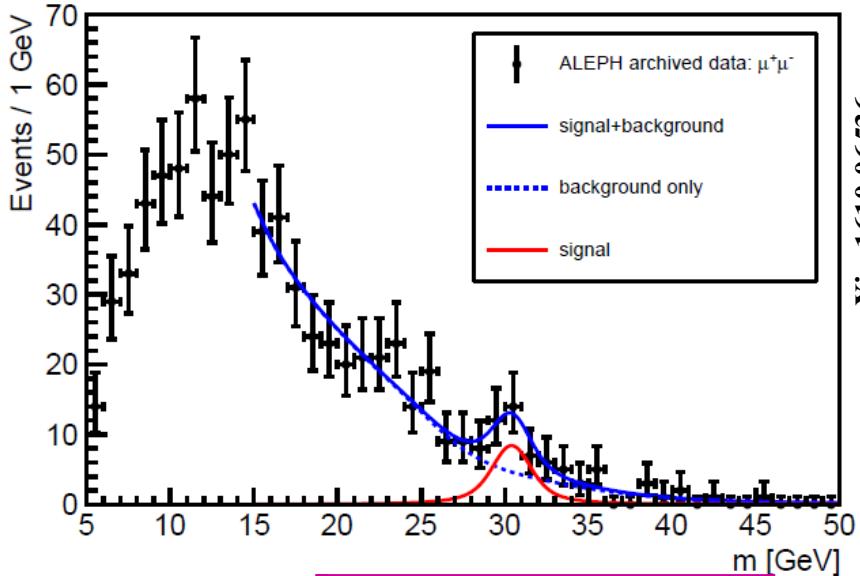
Fluctuation reported in $m_{\mu\mu}$ spectrum

arXiv:1610.06536

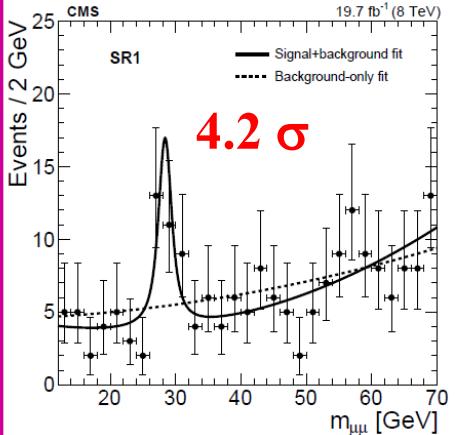
$m_{\mu\mu}$ spectrum obtained from analysis of archived ALEPH data (with b bbar)

CMS

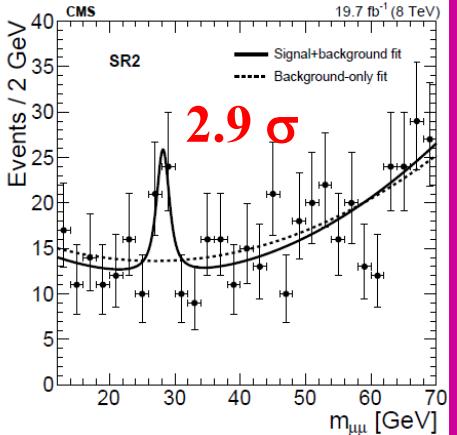
JHEP 1811 (2018) 161



8TeV 19.7 fb^{-1}

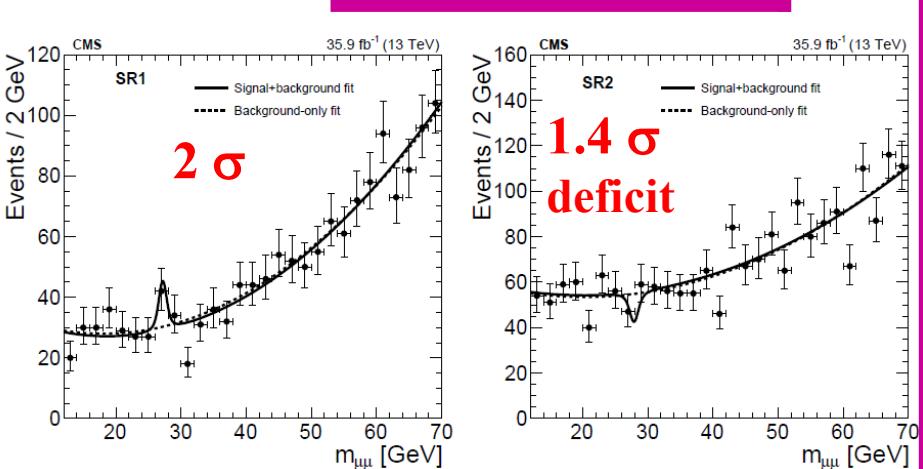


central-forward

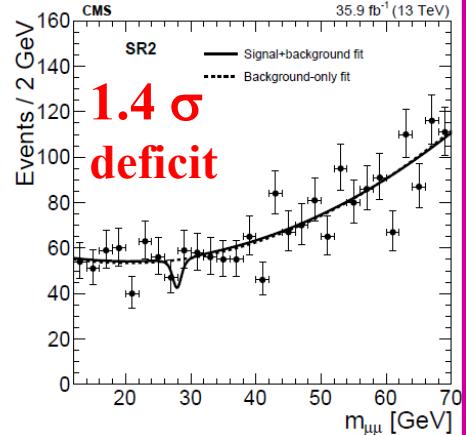


2 central

13TeV 35.9 fb^{-1}



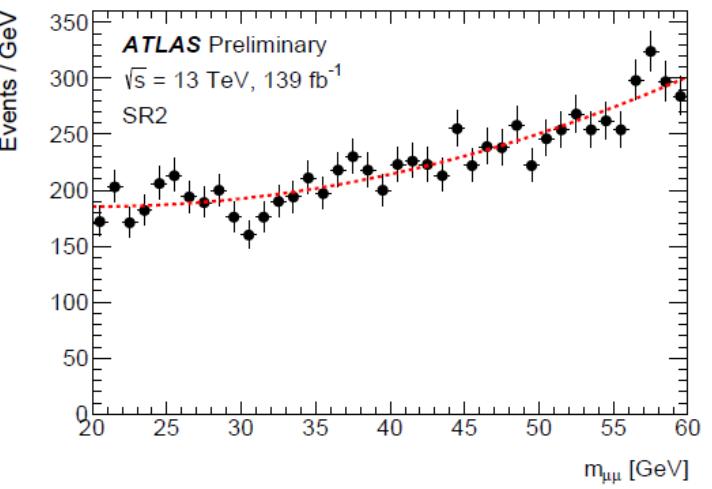
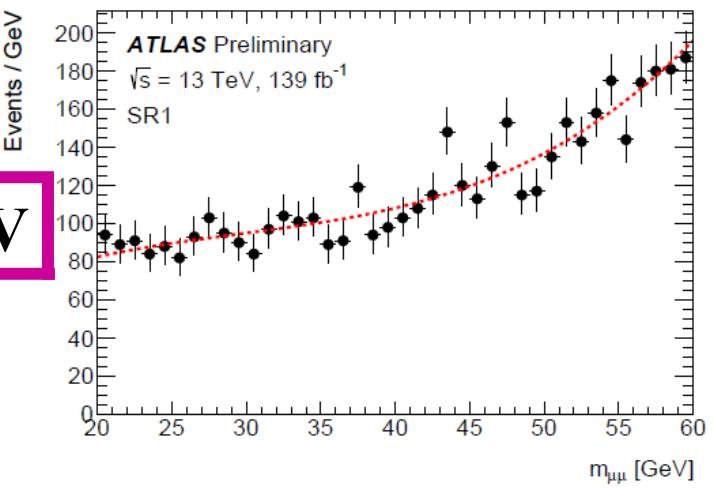
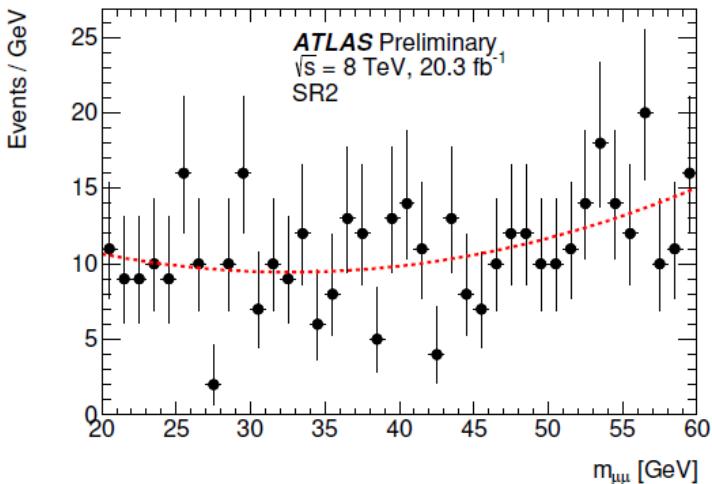
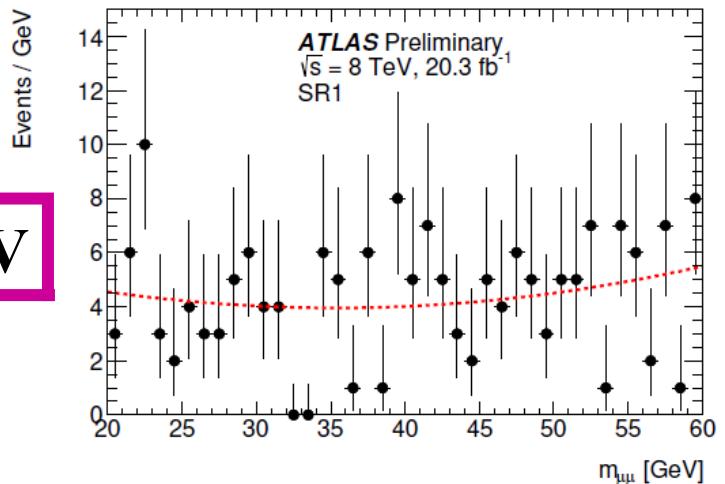
central-forward



2 central

Search for Physics BSM (8)

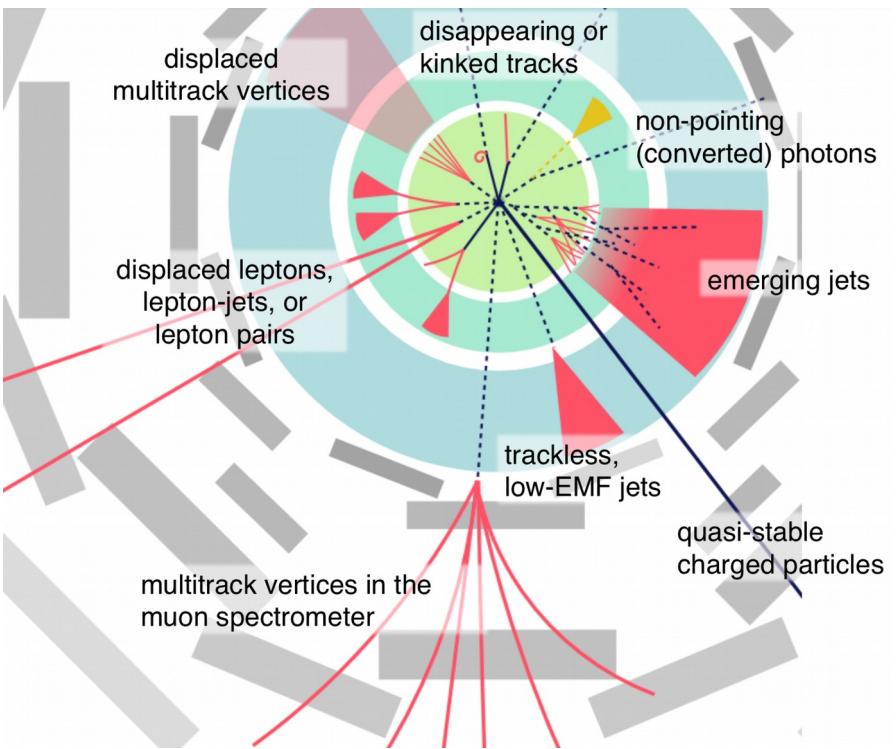
Full Run-2 ~ same cuts than CMS



no significant excess observed

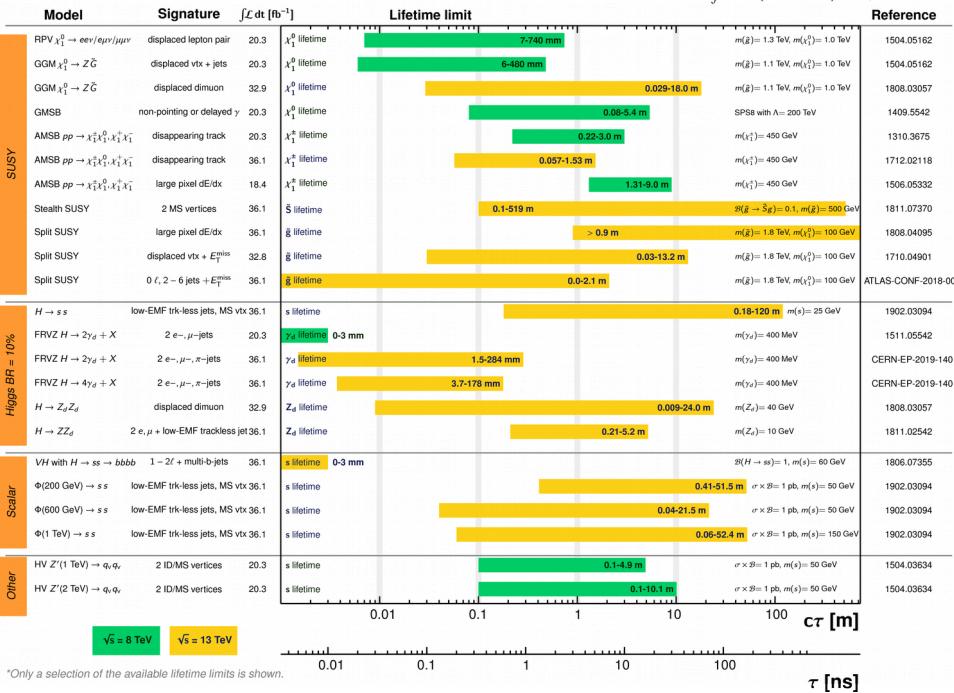
Multiple reasons to be long lived .. small couplings .. intermediate states

many challenging signatures



ATLAS Long-lived Particle Searches* - 95% CL Exclusion

Status: July 2019



- ♪ *Results from (run 1 and) run 2*
- * *detector*
- * *SM*
- * *BSM*
- * ***(B-E)H***
- * *Vector-boson*

The (Brout-Englert-) Higgs = BEH boson(s)

1 Additional BEH bosons

2 The SM BEH boson

3 Search for a pair of BEH bosons

1 Additional BEH bosons (1)

General recipe : SM Higgs Doublet + Additional Field = Additional H bosons

SM + 1 additional H doublet = 2HDM (Two Higgs Doublet

**Model) that corresponds to 5 physical Higgs bosons
 h, H, A, H^+, H^-**

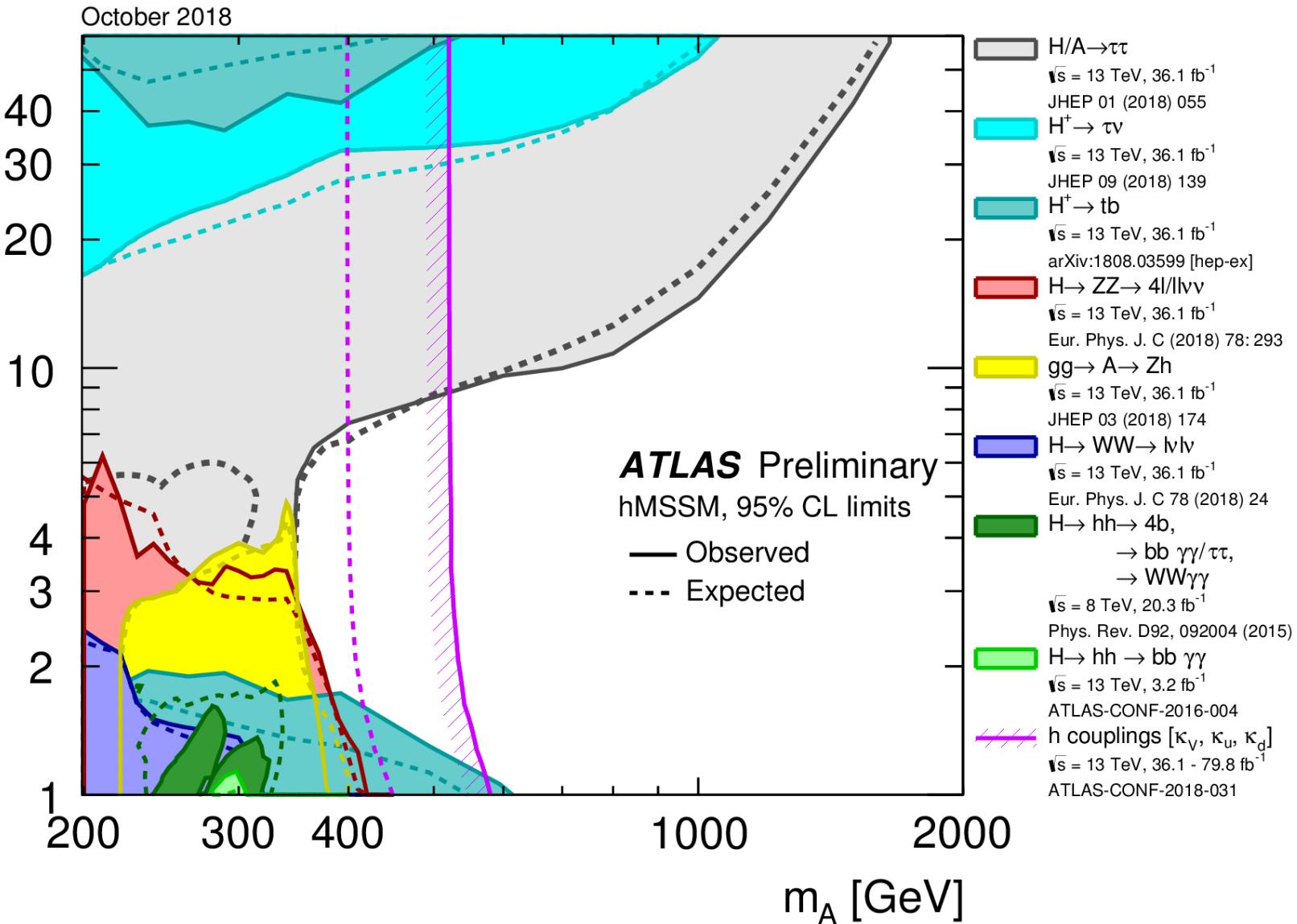
Four v:

Coupling scale factor	Type I	Type II	Lepton-specific	Flipped
κ_V			$\sin(\beta - \alpha)$	
κ_u			$\cos(\alpha) / \sin(\beta)$	
κ_d	$\cos(\alpha) / \sin(\beta)$	$-\sin(\alpha) / \cos(\beta)$	$\cos(\alpha) / \sin(\beta)$	$-\sin(\alpha) / \cos(\beta)$
κ_ℓ	$\cos(\alpha) / \sin(\beta)$	$-\sin(\alpha) / \cos(\beta)$	$-\sin(\alpha) / \cos(\beta)$	$\cos(\alpha) / \sin(\beta)$

MSSM \subset type II HDM .. Numerous benchmark models like hMSSM

1 Additional BEH bosons (2)

$\tan \beta$



1 Additional BEH bosons (3) $\gamma\gamma$ excess at 95 GeV

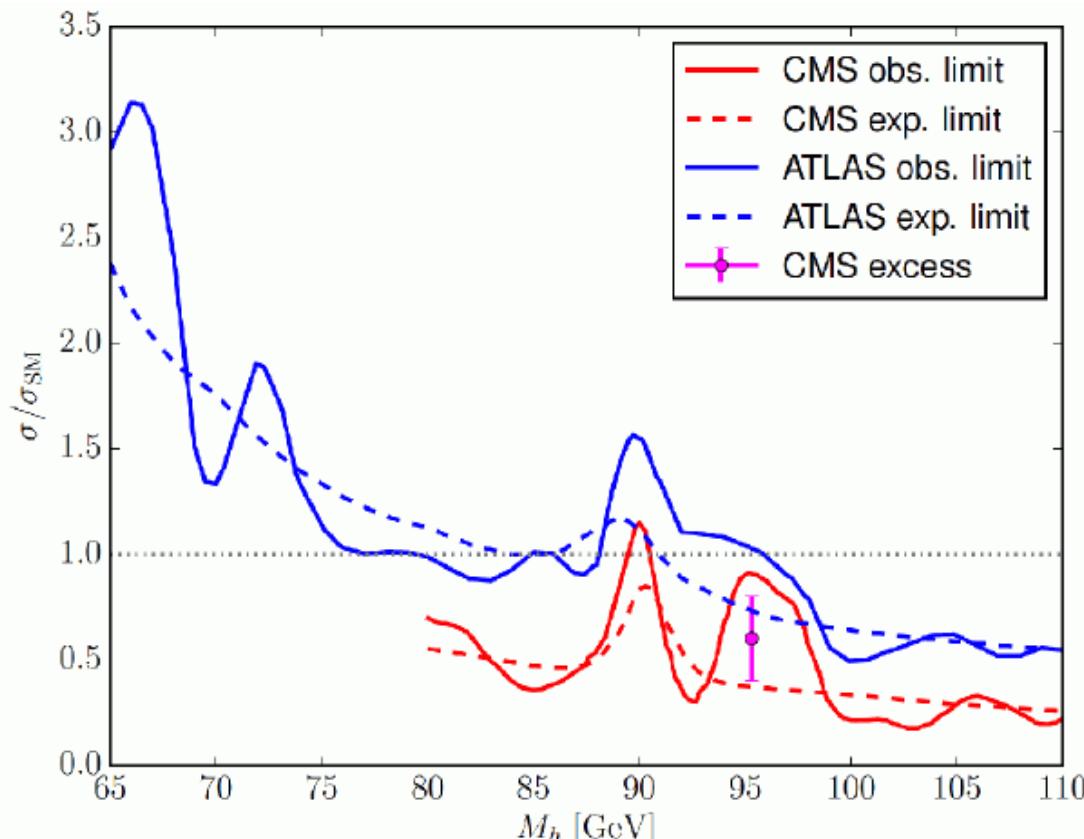
comparison between CMS and ATLAS results (*Sven Heinemeyer*)

CMS PAS HIG-17-013

20 fb⁻¹ (8 TeV) + 36 fb⁻¹ (13 TeV)

ATLAS-CONF-2018-025

80 fb⁻¹ (13 TeV)



2 The SM BEH boson (1) executive summary

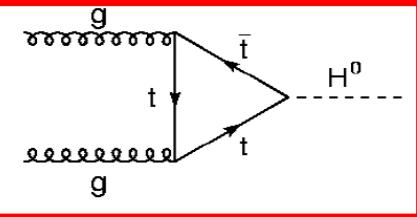
7 years after the discovery we have now a much clearer picture of the BEH boson properties

- ♠ It is **spin 0** and its interactions with bosons are mainly **CP-even**
- ♠ We know its **mass** at **0.2% accuracy**

BEH boson couples to mass → couplings to be measured

Increasing precision in all measurements

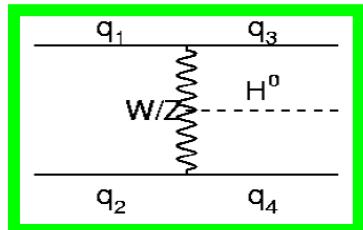
- **bosonic sector** : inclusive measurement at **~10% precision**
differential measurements probing extended phase space
with increasing accuracy
- **fermionic sector** : **3rd generation** (τ, t, b) established
with uncertainties approaching **~20% level**. Most
promising channel for **2nd generation** is $H \rightarrow \mu\mu$



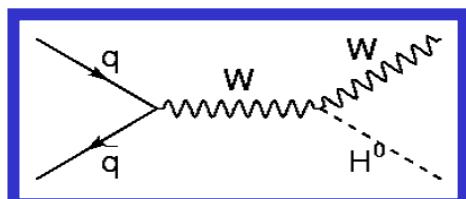
2 The SM BEH boson (2)

A.Djouadi Phys.Rept.457:1-216

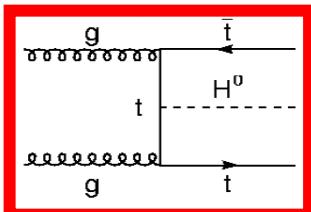
GF $H \rightarrow WW, ZZ, \gamma\gamma, (bb), \tau\tau$



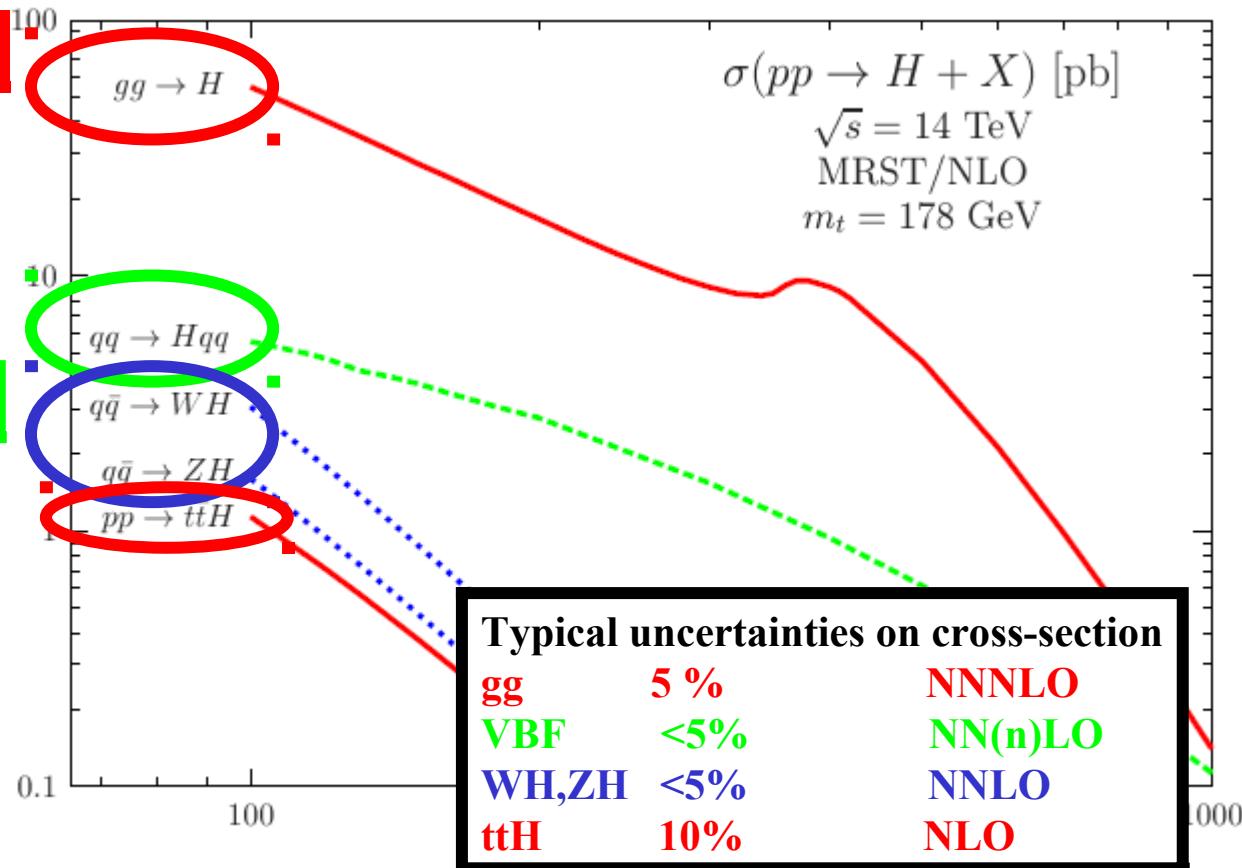
VBF $H \rightarrow WW, ZZ, \gamma\gamma, bb, \tau\tau$



$H \rightarrow WW, \gamma\gamma, bb$



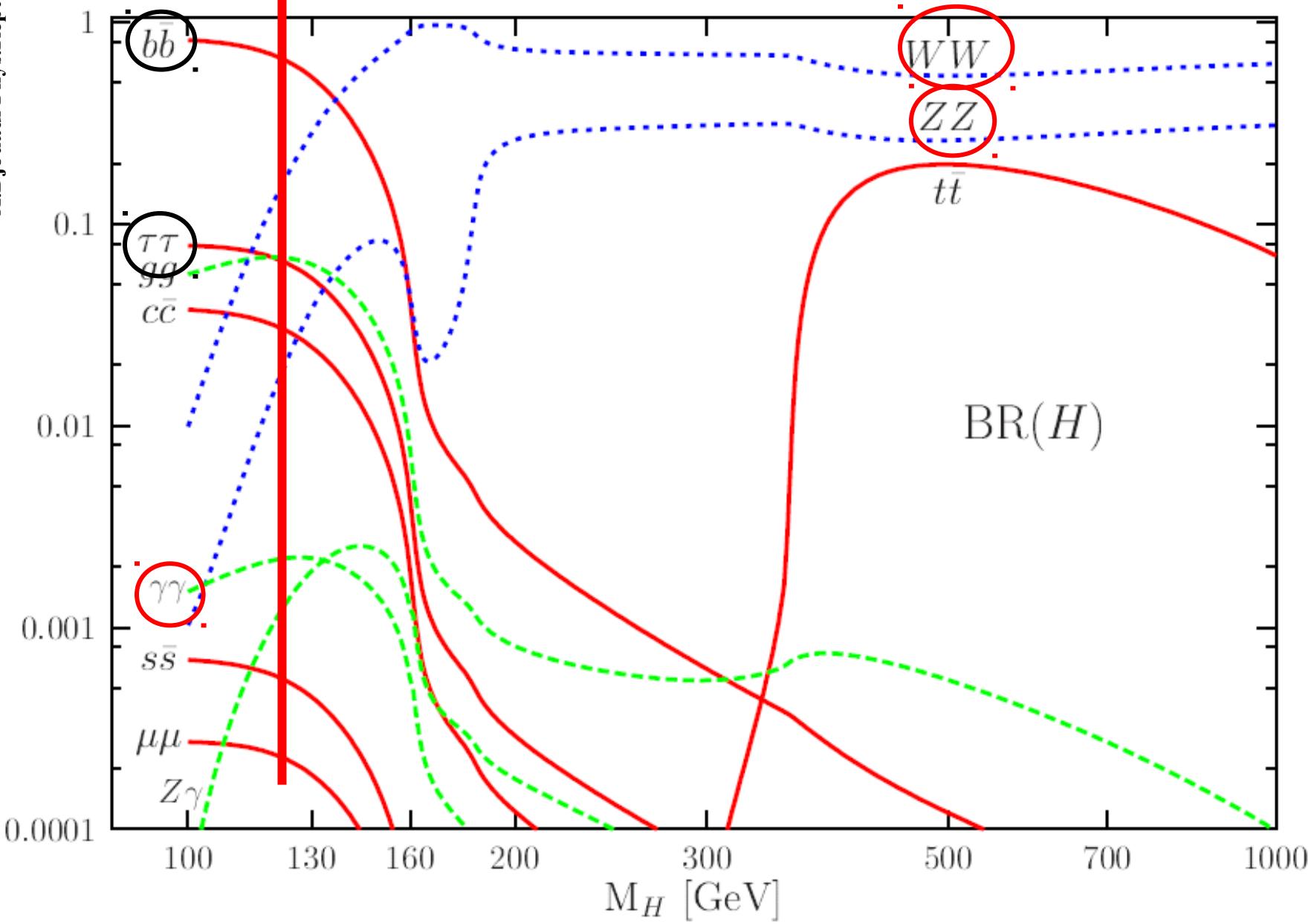
$H \rightarrow WW, \gamma\gamma, \tau\tau, ZZ, bb$



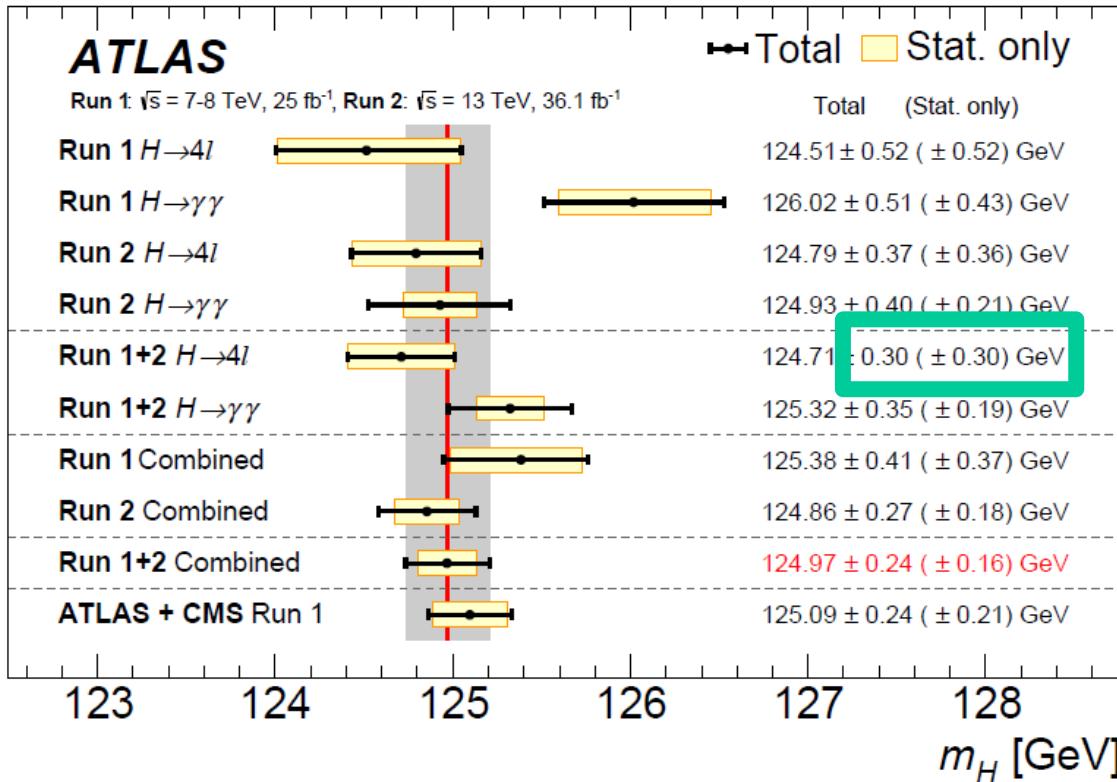
These production cross sections have to be used with the decays **bb**, **$\tau\tau$** , **WW**, **ZZ** , **$\gamma\gamma$**

channels with good mass resolution

2 The SM BEH boson (3)



2 The SM BEH boson (4) The H mass



uncertainty on mass < 0.2 %

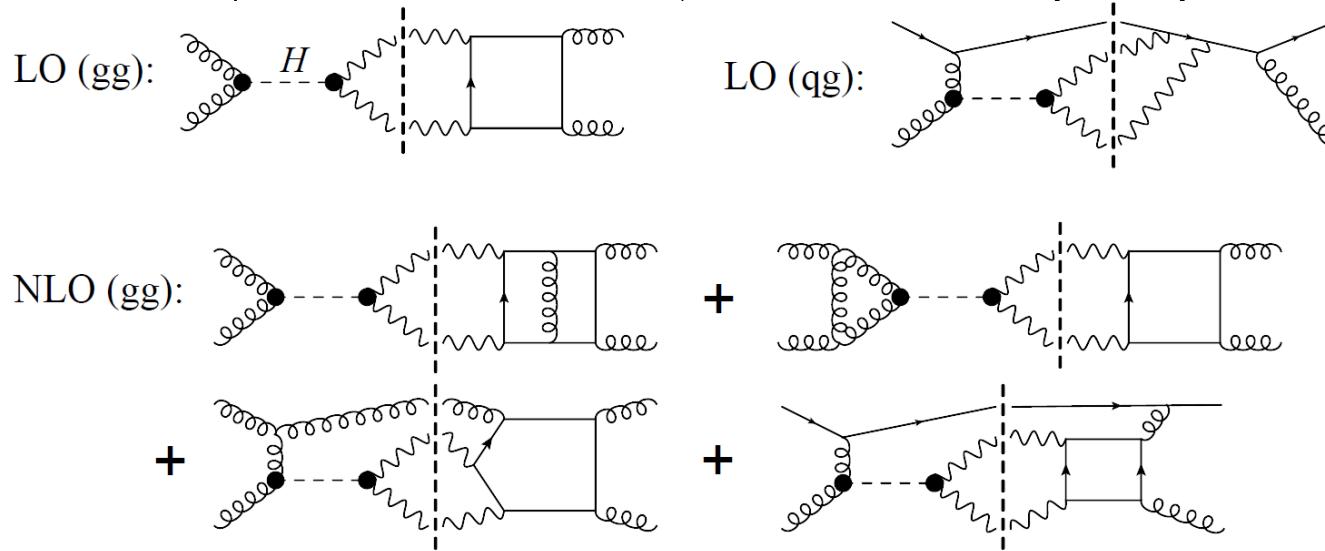
Remember ATLAS has an uncertainty on W mass of 19 MeV Eur.Phys.J. C78 (2018) no.2, 110
 note that $\Delta m_H = 0.1 \text{ GeV} \rightarrow \Delta(\text{BR}(H \rightarrow ZZ)) / \text{BR}(H \rightarrow ZZ) \sim 1\%$

At longer term uncertainty will be dominated by 4l
 (for $H \rightarrow \gamma\gamma$: need to extrapolate from e to γ !)

2 The SM BEH boson (5) Interference in $\gamma\gamma$ (between signal and background)

start to be sensitive to **subtle effects** like interference
(between signal and background) in $\gamma\gamma$

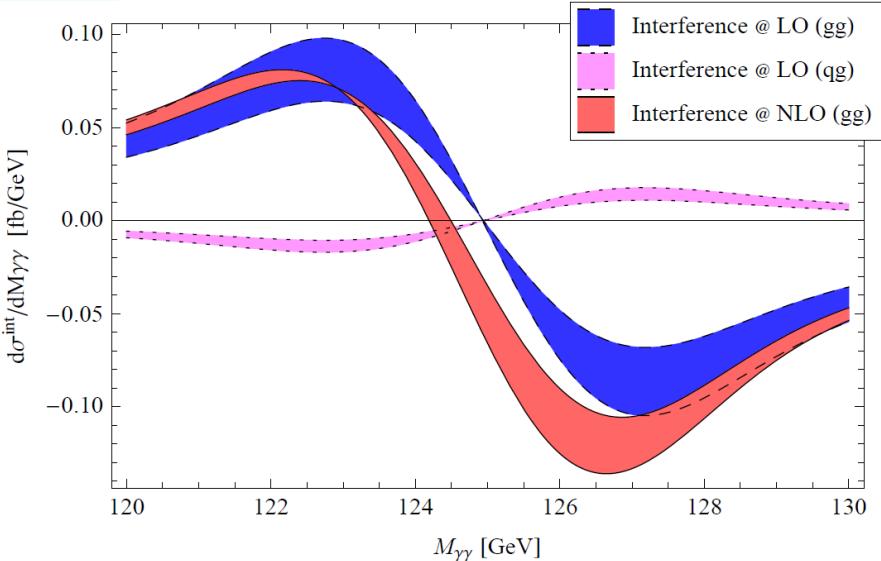
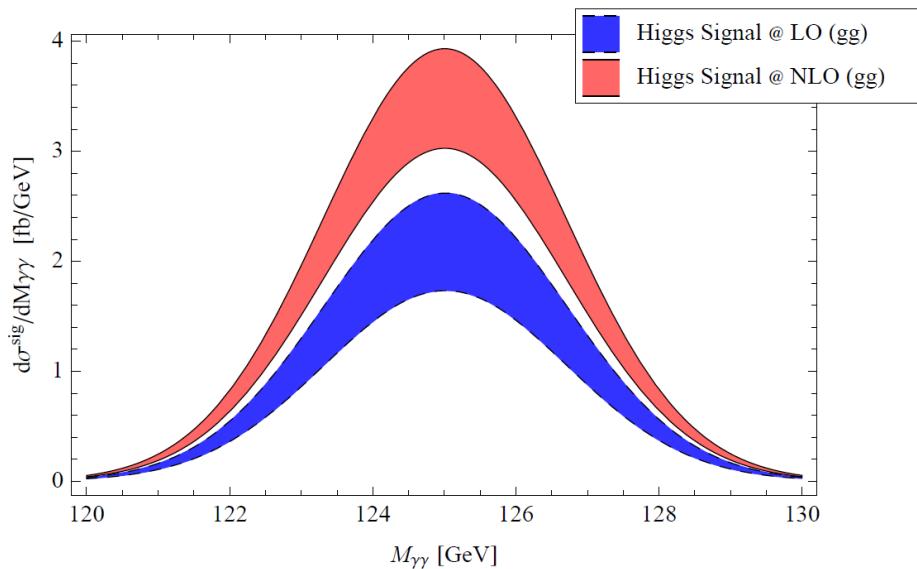
Martin, Dixon and Li Phys.Rev.Lett. 111 (2013) 11180



Interference depends of S/B , therefore is smaller at high $p_T(H)$
where S/B is larger

some work can be done at high pT ($H+2j$) see for instance
Phys.Rev. D92 (2015) no.1, 013004

2 The SM BEH boson (6) Mass shift



The expected median mass shift is found to be $\Delta m_H = -35 \pm 9$ MeV assuming a Standard Model width of 4 MeV for the Higgs boson

However the effect is larger for larger H width
→ could constrain the H width

ATL-PHYS-PUB-2016-009

There is also an effect on the cross-section measurement (few %) to be taken into account

2 The SM BEH boson (7) definition of μ

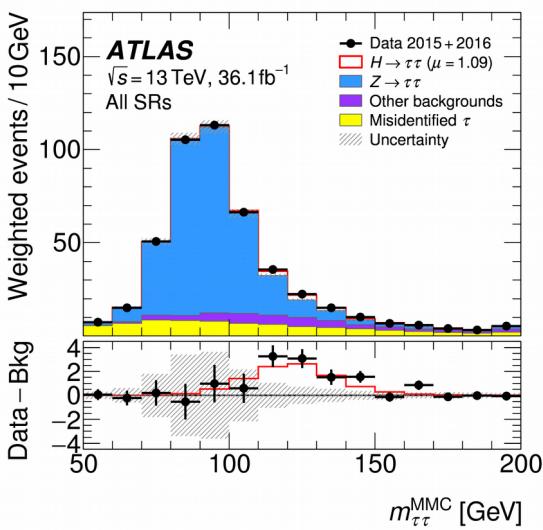
$$\mu = (\sigma \cdot \text{BR}) / (\sigma \cdot \text{BR})_{\text{SM}}$$

2 The SM BEH boson (8) some fermionic results

H $\rightarrow\tau\tau$

Analysis targets VBF
and gg \rightarrow Hj production
Three decay channels

$\tau_{\text{lep}}\tau_{\text{had}}$, $\tau_{\text{lep}}\tau_{\text{lep}}$, $\tau_{\text{had}}\tau_{\text{had}}$



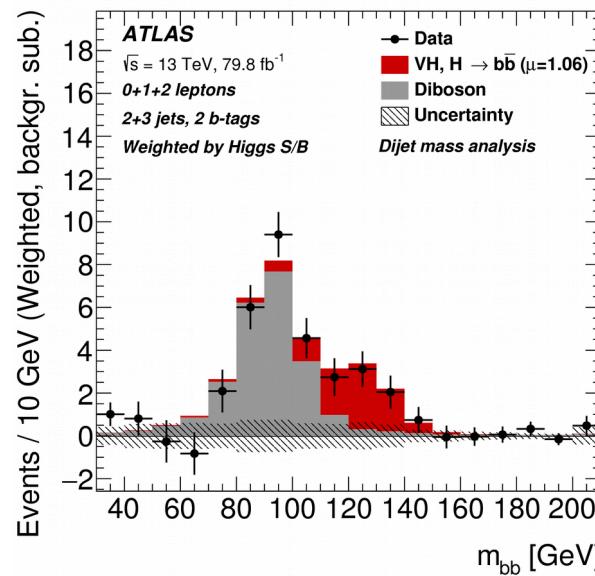
Phys. Rev. D 99, 072001 (2019)

Observation of H $\rightarrow\tau\tau$
(6.4σ obs - 5.4σ exp)
when combined with Run-

1

H $\rightarrow bb$

Main analysis is targetting VH but
also start to look at ggH and VBF modes



H $\rightarrow bb$: 5.4σ obs 5.5σ exp

$$\mu = 1.01 \pm 0.12(\text{stat.})^{+0.16}_{-0.15}(\text{syst.})$$

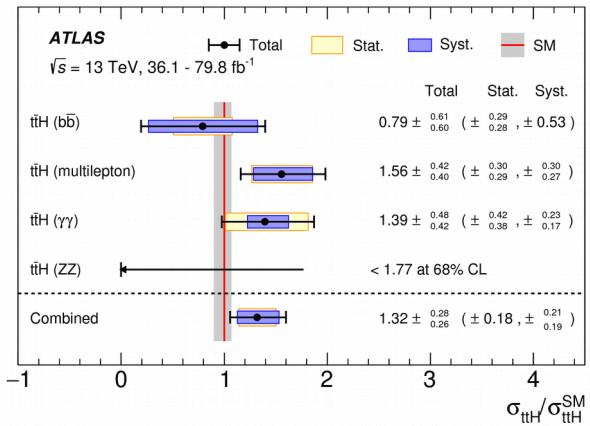
Combination of VH channels gives

significance obs(exp) of 5.3σ (4.8σ)

Phys.Lett. B786 (2018) 59-86

2 The SM BEH boson (9) ttH

Observation of ttH (june 2018)

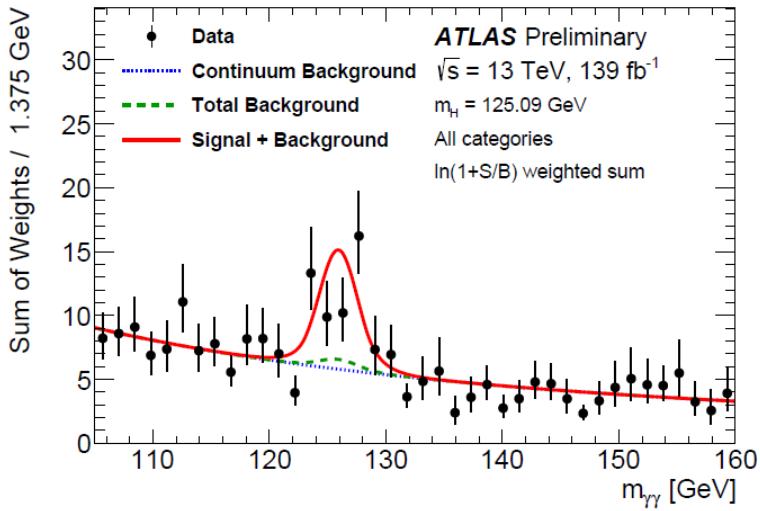


Phys.Lett.B784 (2018) 173-191

Combined with Run-1
obs(exp) significance of
6.3 (5.1) σ

Full Run-2

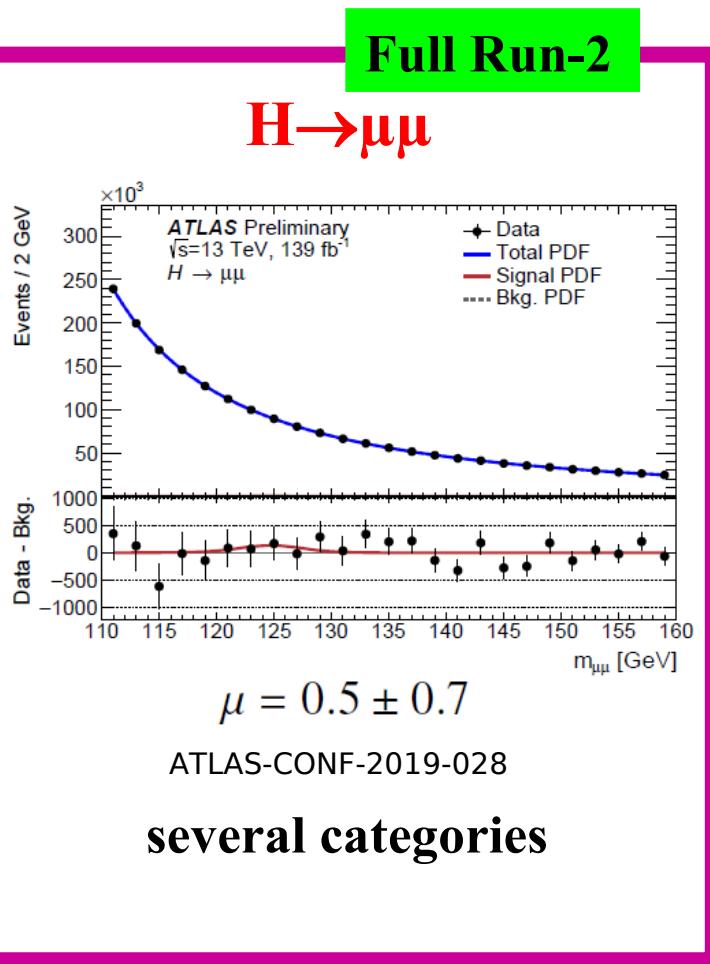
ttH($\rightarrow\gamma\gamma$) full Run-2



Observed (exp)
significance of
4.9(4.2) σ

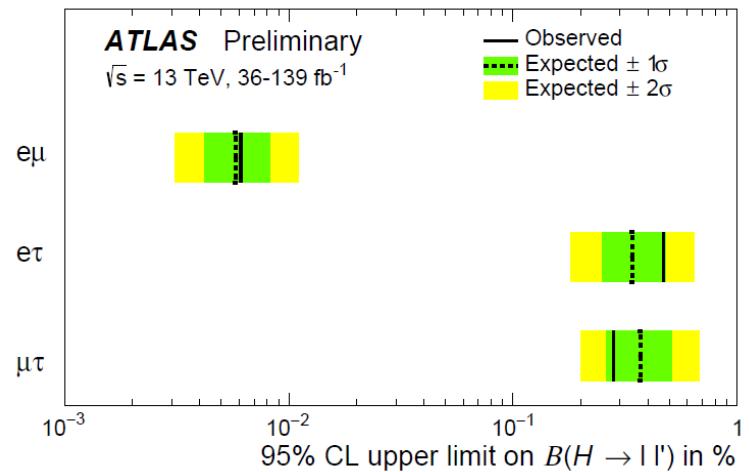
ATLAS-CONF-2019-001

2 The SM BEH boson (10) dileptons



$H \rightarrow ee, e\mu, e\tau, \tau\mu$

no evidence for $H \rightarrow ee$ ☺



$$\text{BR}(H \rightarrow ee) < 3.6 \text{ (3.5)} 10^{-4}$$

obs exp

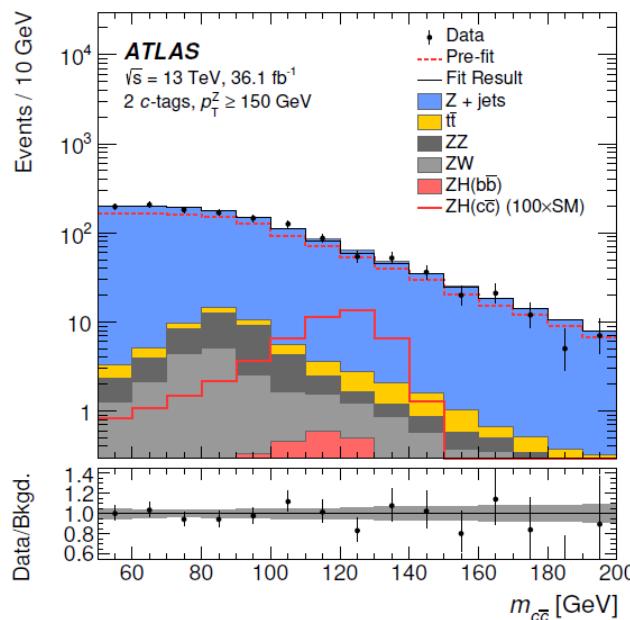
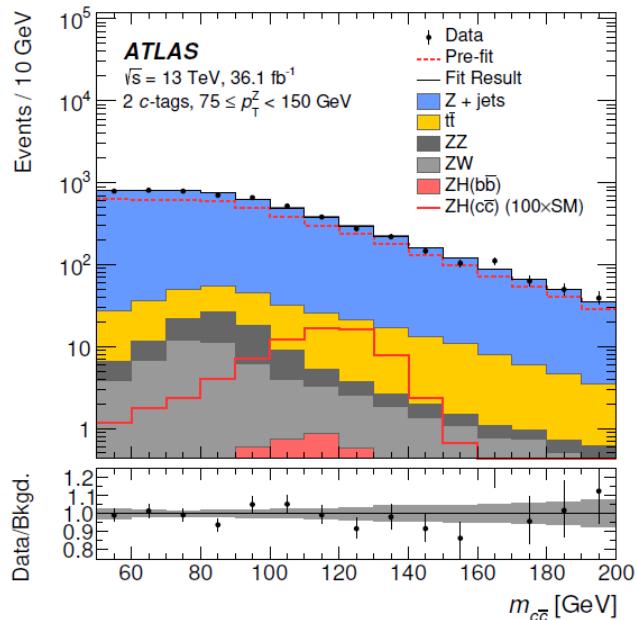
2 The SM BEH boson (11) $H \rightarrow cc$

Phys.Rev.Lett. 120 (2018) no.21, 211802

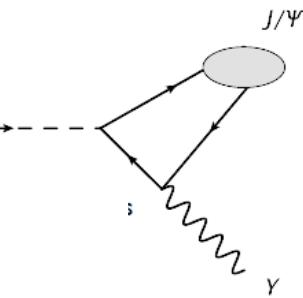
H \rightarrow cc

in mode $Z(\rightarrow l^+l^-)H$

observed (expected) upper limit on μ at the 95% CL of 110 (150^{+80}_{-40})

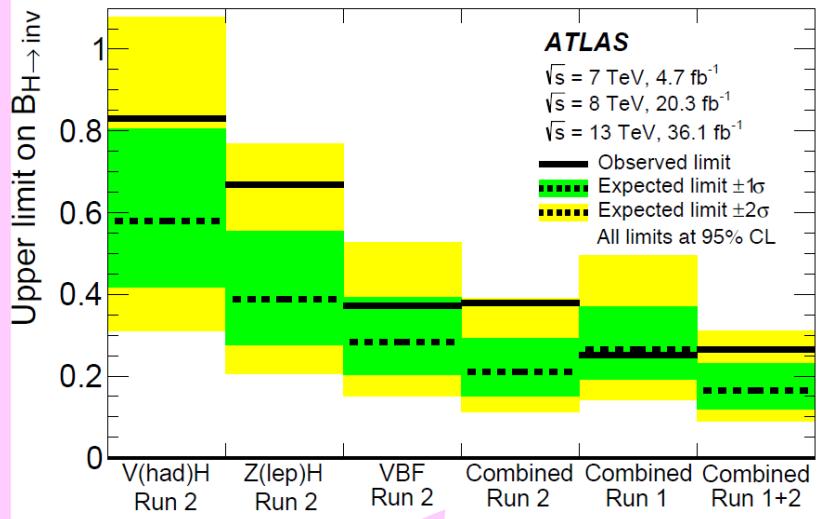


also $H \rightarrow J/\psi \gamma$

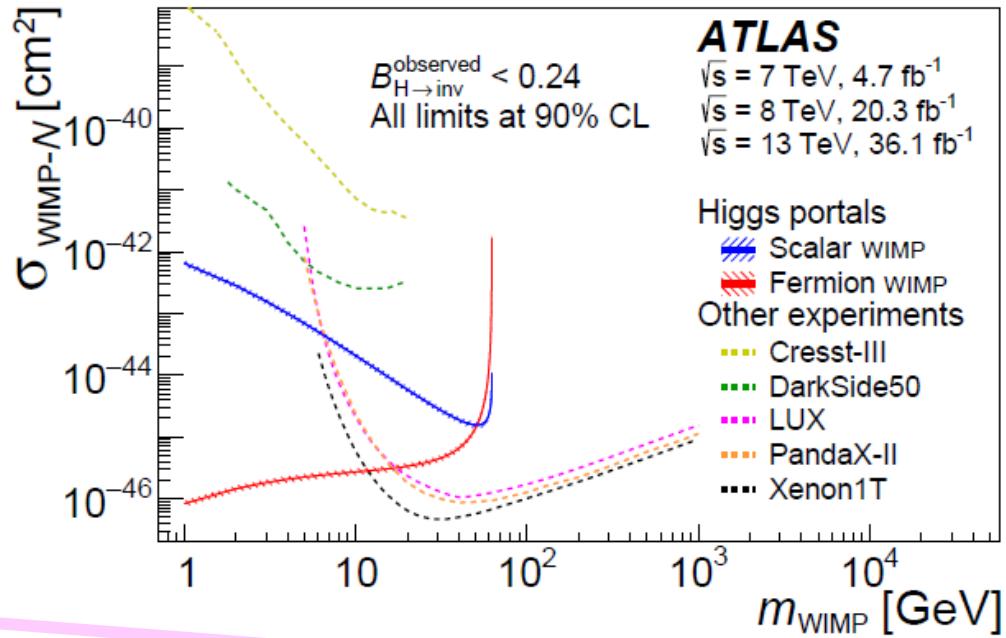


Invisible H

Upper limit at 95% CL on
 $\text{BR}(\text{H} \rightarrow \text{inv})$ obs(exp) 0.26(0.17)



Comparison with direct
detection experiments



SM $\text{BR}(\text{H} \rightarrow 4\nu) \sim 1.2 \cdot 10^{-3}$

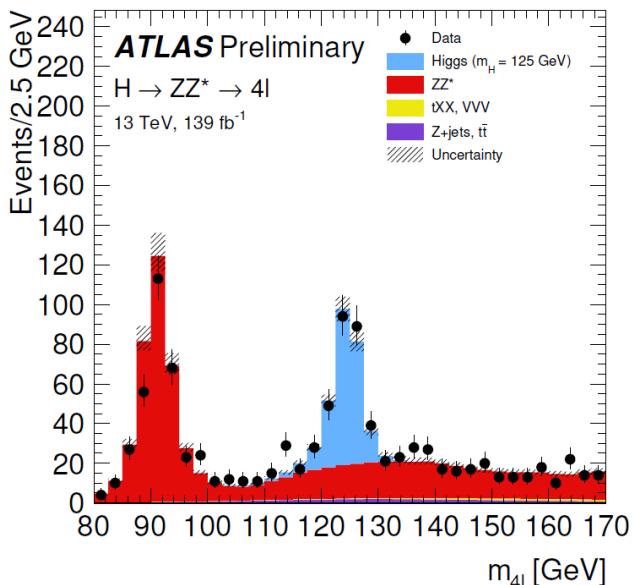
$E_T^{\text{miss}} + 2 \text{ jets}$

2 The SM BEH boson (13) $H \rightarrow 4l$ $H \rightarrow \gamma\gamma$

Full Run-2

Full Run-2

$H \rightarrow 4l$

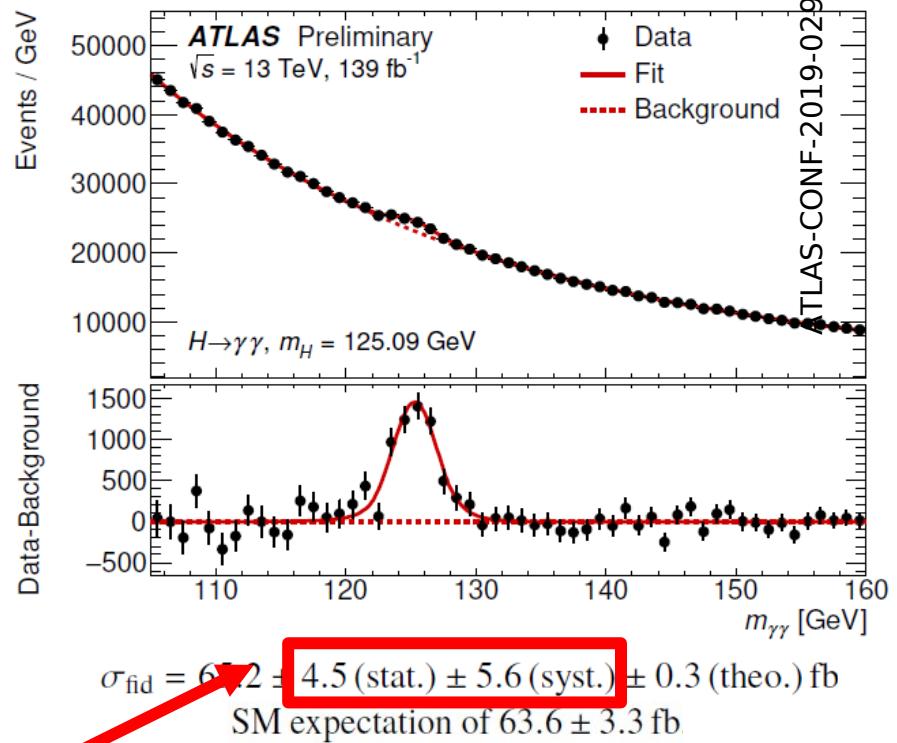


$$\mu = 1.04^{+0.09}_{-0.08}(\text{stat.})^{+0.04}_{-0.03}(\text{exp.})^{+0.06}_{-0.05}(\text{th.}) = 1.04^{+0.12}_{-0.10}$$

ATLAS-CONF-2019-025

syst (background modelling)
 + energy resolution + > stat !

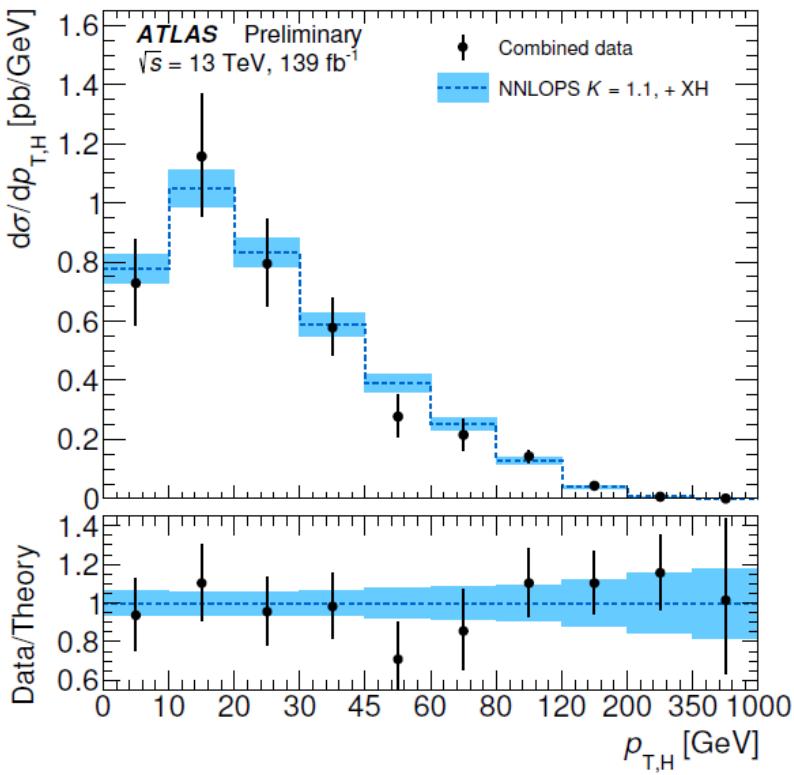
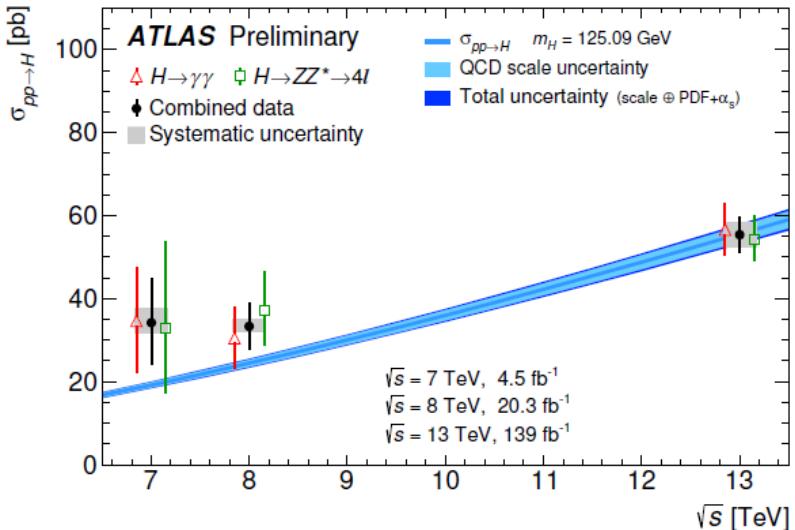
$H \rightarrow \gamma\gamma$



constraint on charm coupling
 through pT distribution (see also next slide)
 ($gg \rightarrow H$ and $c\bar{c} \rightarrow H$)

Coefficient	Observed 95% CL limit	Expected 95% CL limit
κ_c	$[-19, 24]$	$[-15, 19]$

2 The SM BEH boson (14) combined H \rightarrow 4l and H $\rightarrow\gamma\gamma$



Combined inclusive pp \rightarrow H cross section

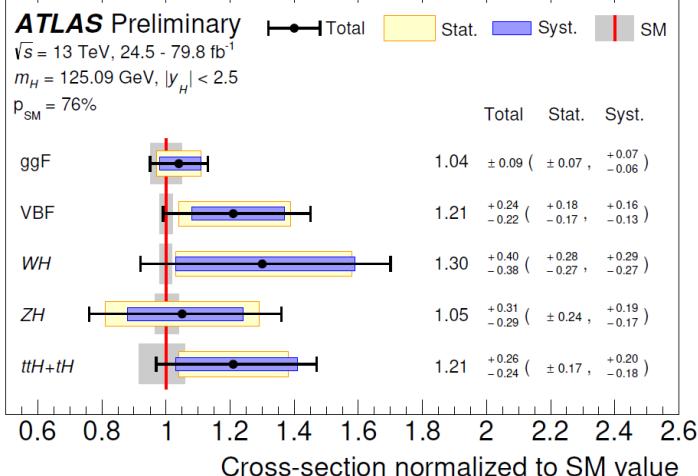
$55.4^{+4.3}_{-4.2}$ pb (± 3.1 (stat.) $^{+3.0}_{-2.8}$ (sys.))

SM = 55.6 ± 2.5 pb

2 The SM BEH boson (15) H combination

ATLAS-CONF-2019-005

$$\mu = 1.11^{+0.09}_{-0.08} = 1.11 \pm 0.05 \text{ (stat.)} {}^{+0.05}_{-0.04} \text{ (exp.)} {}^{+0.05}_{-0.04} \text{ (sig. th.)} \pm 0.03 \text{ (bkg. th.)}$$

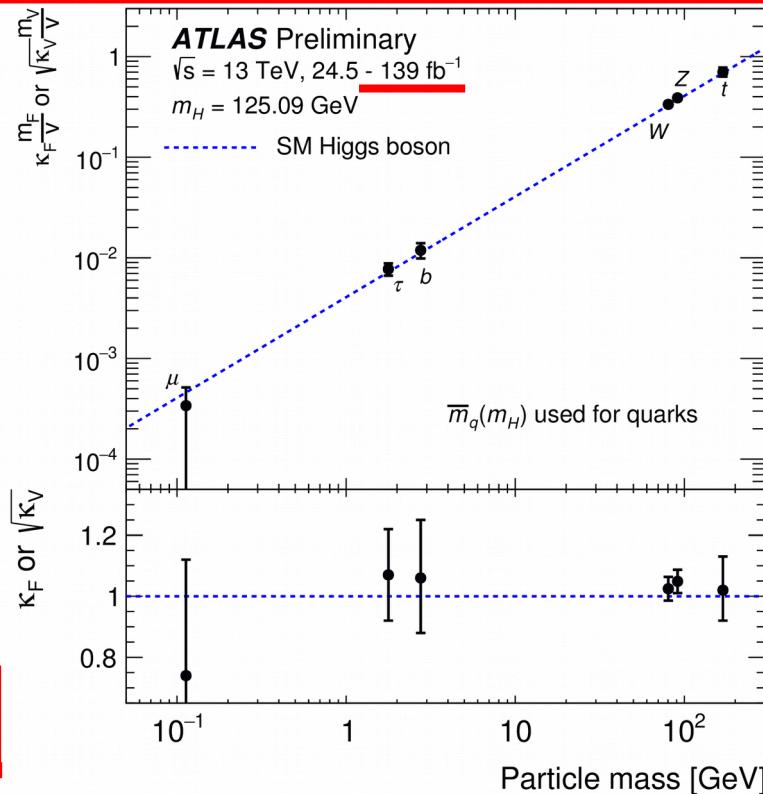


assumption of SM branching fractions

13 TeV
up to 80 fb^{-1}

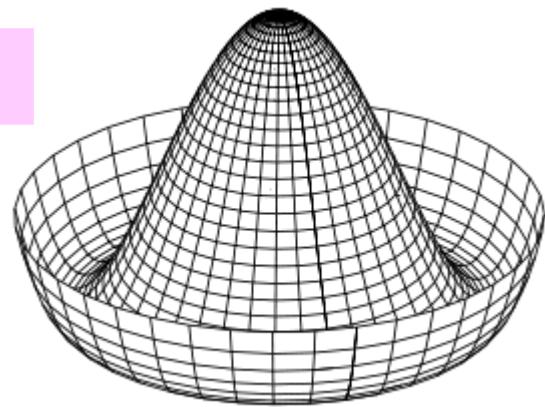
checked

iconic H plot
The SM BEH mechanism predicts relations between couplings and masses



3 Search for a pair of BEH bosons (1)

After discovering the Higgs boson, the ultimate probe of the Standard Model is to fully measure the Higgs potential.

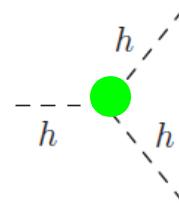


$$\Phi \rightarrow v + h$$

$$V(\phi) = \frac{1}{2} \mu^2 \phi^2 + \frac{1}{4} \lambda \phi^4 = \boxed{\lambda v^2 h^2} + \boxed{\lambda v h^3} + \frac{1}{4} \lambda h^4$$

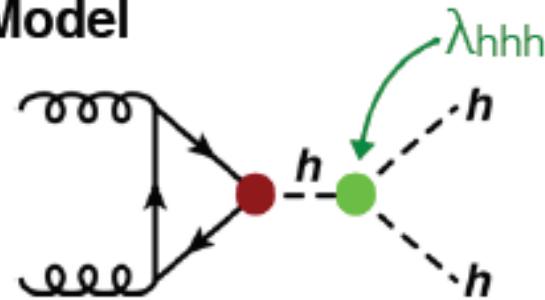
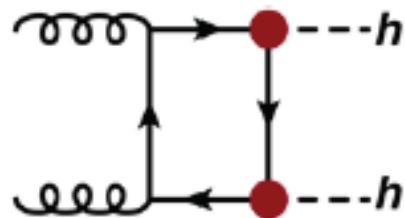
mass term self coupling terms

$$\frac{1}{2} m_h^2 h^2$$



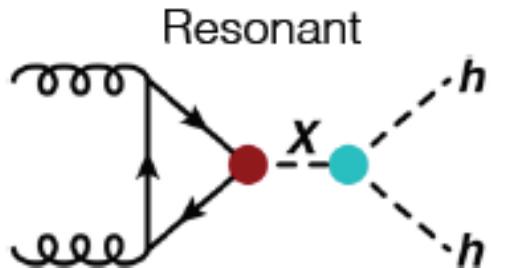
3 Search for a pair of BEH bosons (2)

Standard Model

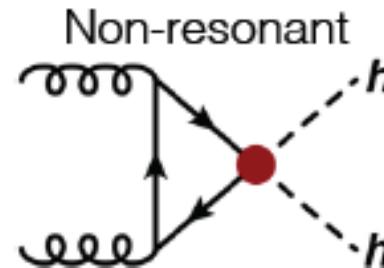


33.4 fb @ 13 TeV (significant destructive interference)

Beyond Standard Model

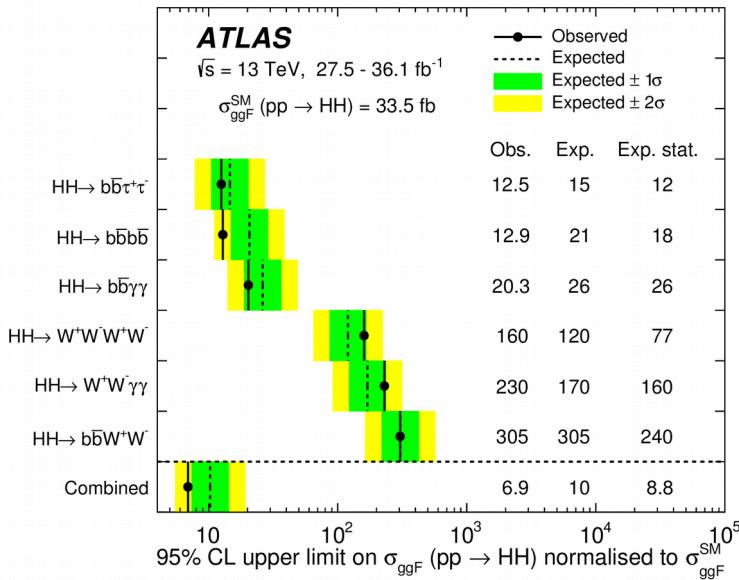


Resonant
KK gravitons, heavy higgs, ...



Non-resonant
tthh vertex, coloured scalars, ...

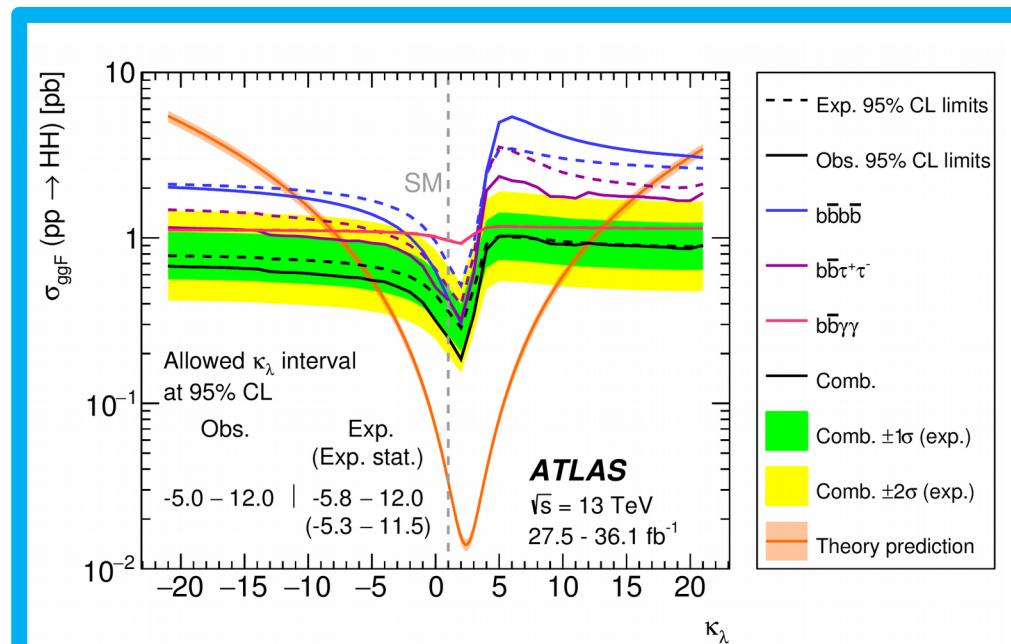
3 Search for a pair of BEH bosons (3)



95% CL limit for $\kappa_\lambda = 1$:

6.9 (10) X SM obs. (exp.)

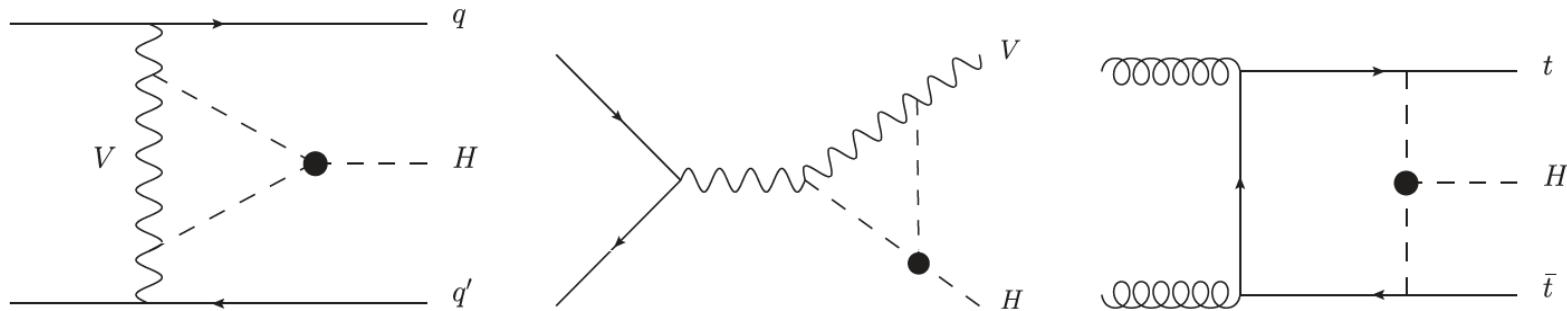
$$\kappa_\lambda = \lambda_{HHH} / \lambda_{HHH}^{\text{SM}}$$



see also new result on 4b channel (VBF) *ATLAS-CONF-2019-030*
and (better) new result on $b\bar{b}l\bar{l}v\bar{v}$ *arXiv:1908.06765*

3 Search for a pair of BEH bosons (4) constraint of the H self-coupling from H differential production and decay mesurements

The Higgs boson cross sections, the branching fractions and the Higgs boson kinematics are affected by the Higgs-boson self coupling contribution through next to leading order electroweak corrections.

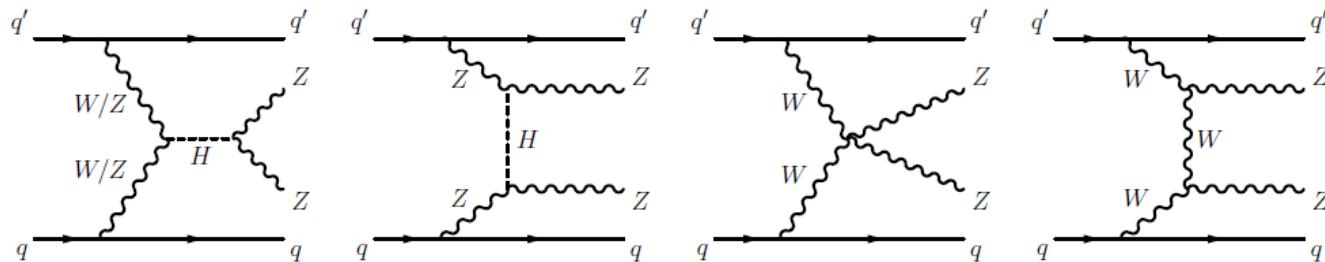


With the assumption that new physics affects only the Higgs boson self-coupling (λ_{HHH}), the ratio $\lambda_{HHH}/\lambda_{HHH}^{\text{SM}}$ is determined to be $\lambda_{HHH}/\lambda_{HHH}^{\text{SM}} = 4.0^{+4.3}_{-4.1}$, excluding values outside the interval $-3.2 < \lambda_{HHH}/\lambda_{HHH}^{\text{SM}} < 11.9$ at the 95% C.L.

Results similar to di-Higgs direct search

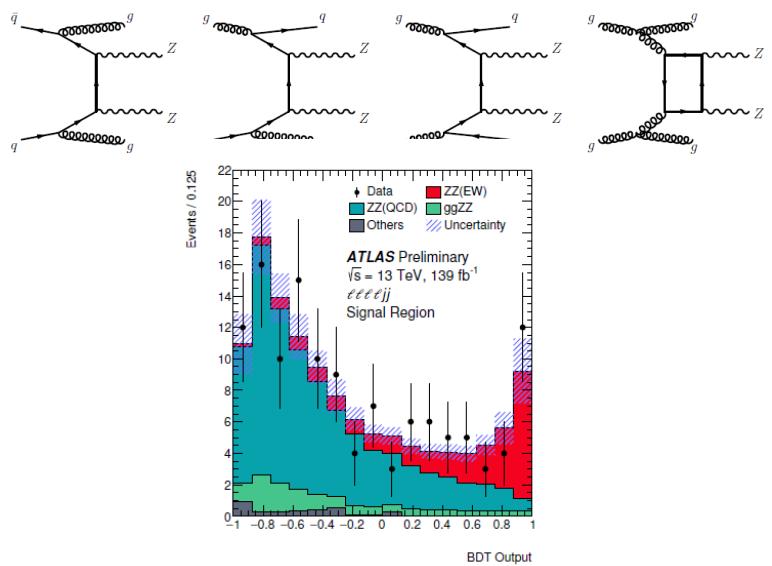
- ♪ Results from (run 1 and) run 2
 - * detector
 - * SM
 - * BSM
 - * (B-E)H
 - * **Vector-boson scattering**

H boson regularizes the EW weak boson scattering at high energies



**ZZ
scattering**

MVA to separate EW from QCD



Observed (expected)
significance for EW
production: 5.5σ (4.3σ)
 $\sigma_{fid}(EW) = 0.82 \pm 0.21$ fb
SM pred. = 0.61 ± 0.03 fb

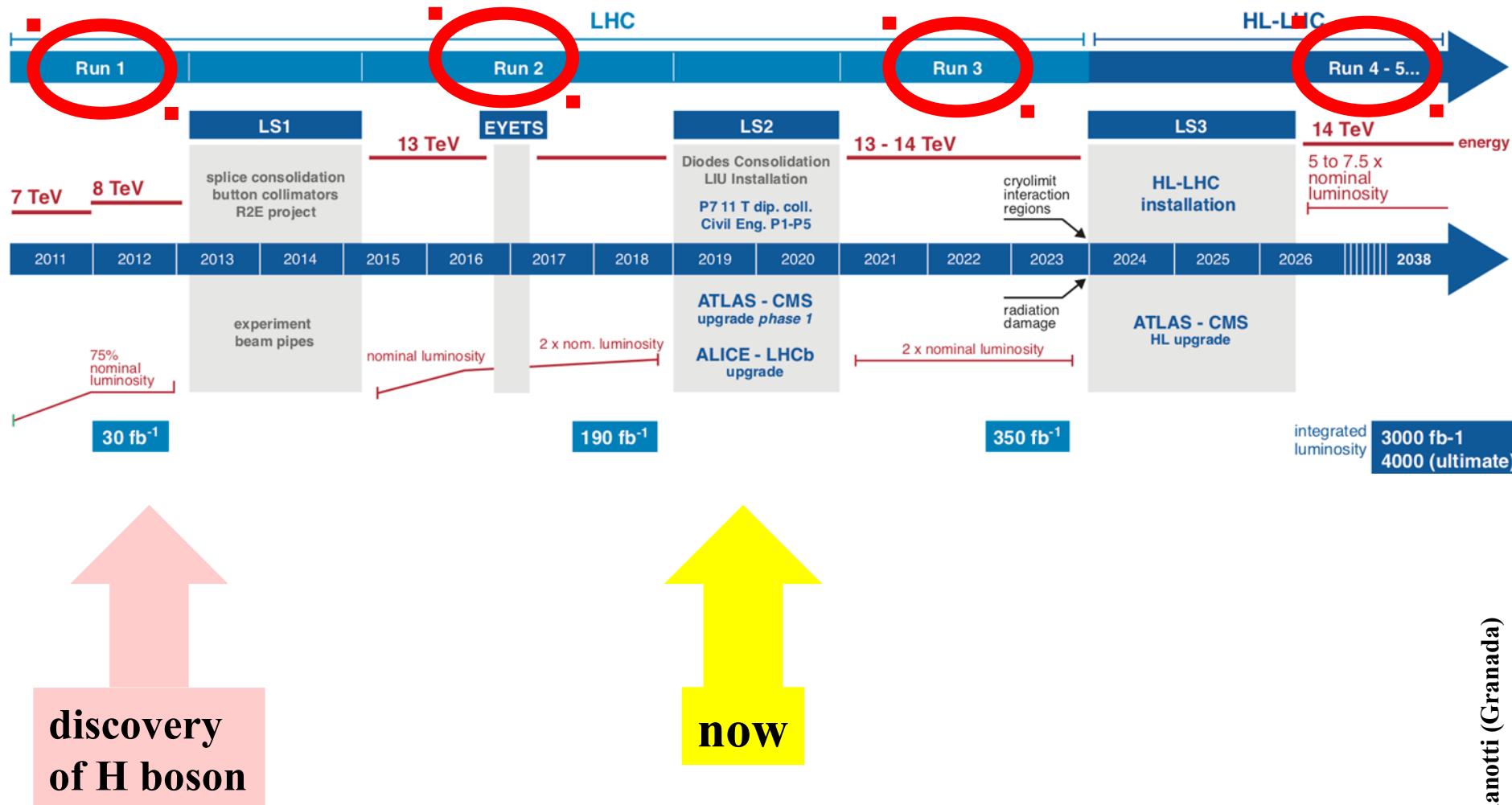
ATLAS observed vector
boson scattering at:

- 6.9σ in WW channel
- 5.3σ in WZ channel

**WW ZZ WZ
observed**

also EW Z γ ATLAS-CONF-2019-039

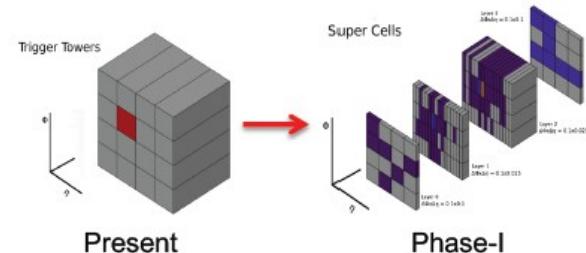
- ♪ *Historical introduction , Setting the stage*
- ♪ *Results from (run 1 and) run 2*
- ♪ ***Future of ATLAS , run 3 , HL-LHC***
- ♪ *Conclusions*
- ♪ *Backup*



ATLAS Phase-I Upgrade

(i) Liquid Argon Calorimeter Electronics

Aim to improve the Level-1 calorimeter decision for Run 3 and beyond
(enhanced jet-rejection and pile-up subtraction)



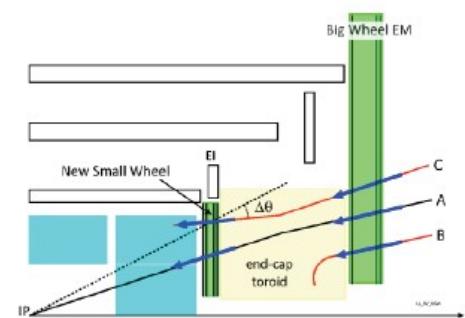
(ii) Trigger / DAQ upgrade

Take full advantage of the finer segmentation available with LAr electronics upgrade, and improved muon trigger information (NSW)

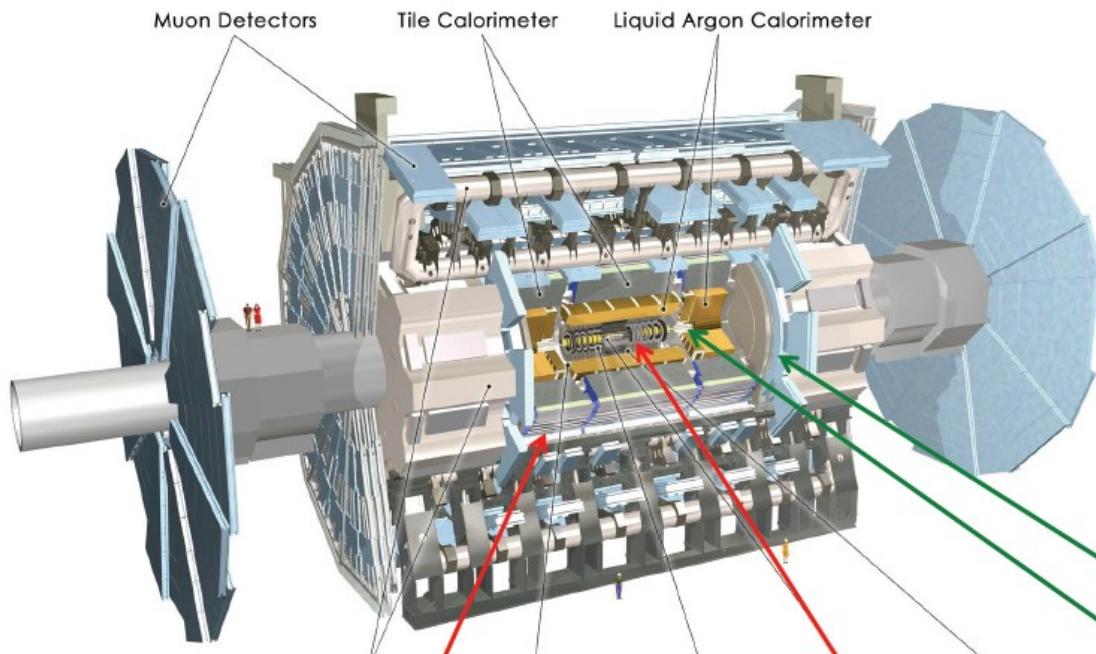


(iii) Muon System: New Small Wheel

Replacement of the inner muon stations in the endcap regions of the detector;
→ reduced muon fake trigger rate, preserve position resolution and efficiency at HL-LHC



ATLAS Phase-II Upgrade



New muon chambers
in the inner barrel region

New Inner Tracking Detector
(all silicon tracker, up to $|\eta| = 4$)

Upgraded Trigger and
Data Acquisition System:

- L0: 1 MHz
- Improved High-Level
Trigger

Electronics Upgrade :

- LAr Calorimeter
- Tile Calorimeter
- Muon system

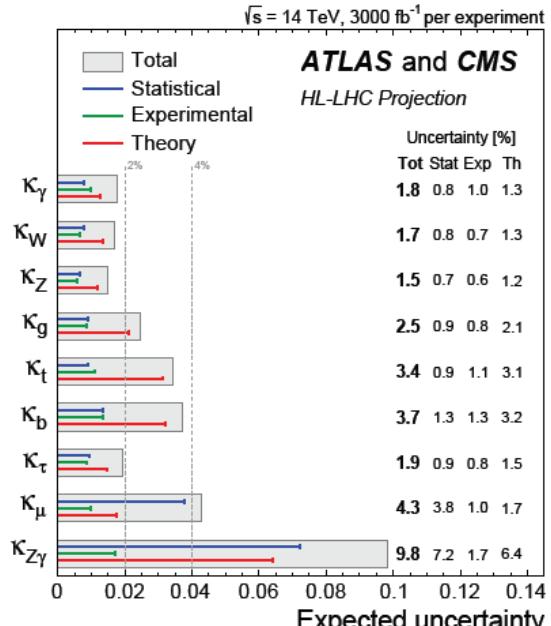
Options:

- High granularity timing detector
(forward region)
- High- η muon tagger

It is very hard to predict, especially the future.

N.Bohr

arXiv:1902.00134



$$\kappa_j^2 = \sigma_j / \sigma_j^{\text{SM}}$$

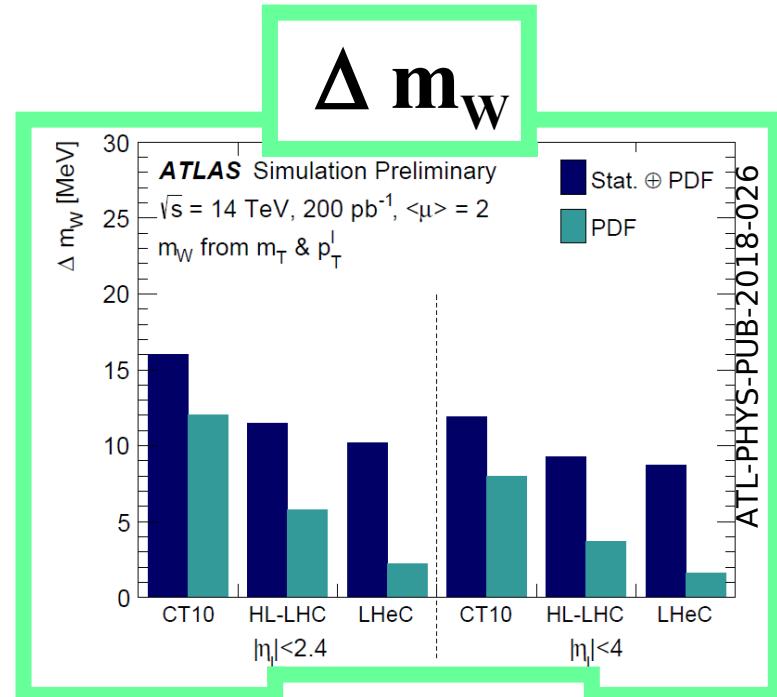
Sensitivity to hh direct search

50% uncertainty on $\kappa_3 \equiv \frac{\lambda_3}{\lambda_3^{\text{SM}}}$

self coupling normalized to SM
 3000 fb^{-1}

arXiv:1905.03764

HH



will require
special runs

- ♪ *Historical introduction , Setting the stage*
- ♪ *Results from (run 1 and) run 2*
- ♪ *Future of ATLAS , run 3 , HL-LHC*
- ♪ ***Conclusions***
- ♪ *Backup*

- ▶ **Fantastic Run-2 dataset , thanks to the outstanding performance of the LHC and ATLAS**
- ▶ **During Run-3 emphasis on precision**
- ▶ **< 5% of the data that will be delivered by HL-LHC
⇒ a lot to do !**

Thanks for your attention

- ♪ *Historical introduction , Setting the stage*
- ♪ *Results from (run 1 and) run 2*
- ♪ *Future of ATLAS , run 3 , HL-LHC*
- ♪ *Conclusions*
- ♪ ***Backup***

- ♪ *Historical introduction , Setting the stage*
- ♪ *Results from (run 1 and) run 2*
- ♪ *Future of ATLAS , run 3 , HL-LHC*
- ♪ *Conclusions*
- ♪ *Backup*

- ♪ *Historical introduction , Setting the stage*
- ♪ ***Results from (Run-1 and) Run-2***
- ♪ *Future of ATLAS , run 3 , HL-LHC*
- ♪ *Conclusions*
- ♪ *Backup*

♪ Results from (run 1 and) run 2

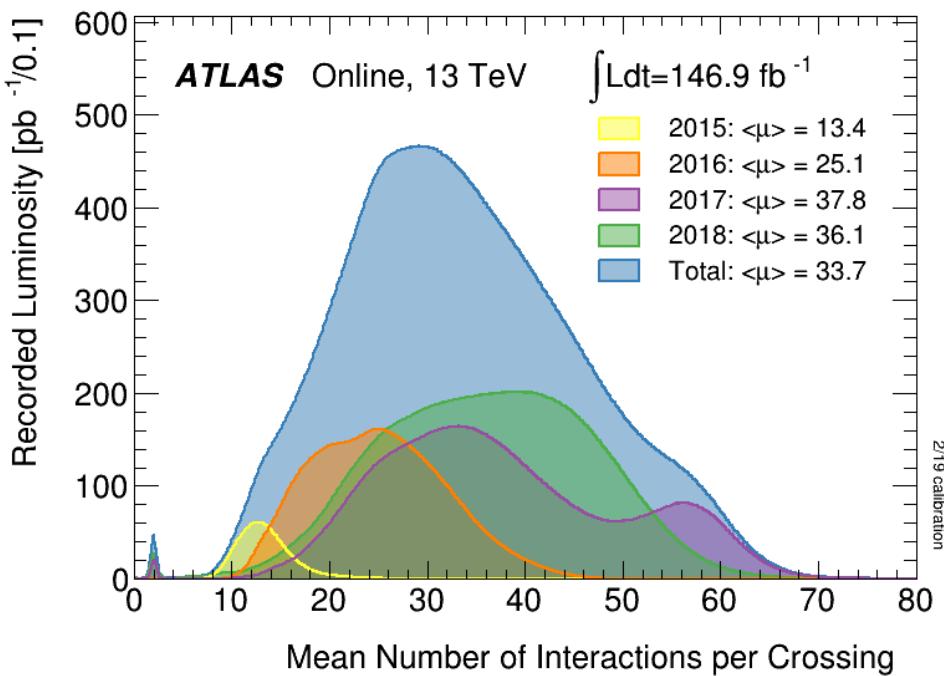
* *detector*

* *SM*

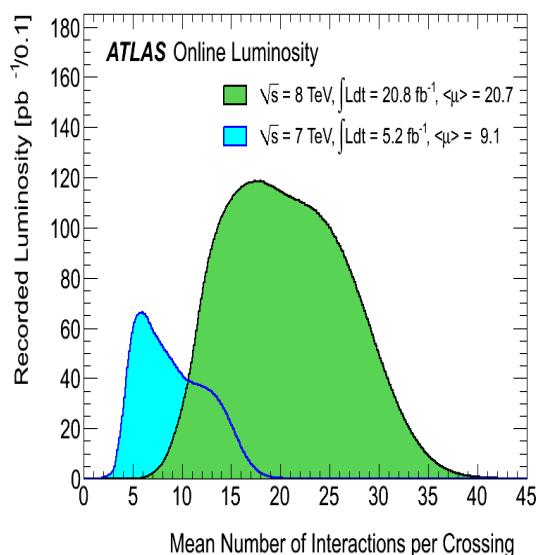
* *BSM*

* *(B-E)H*

* *Vector-boson scattering*

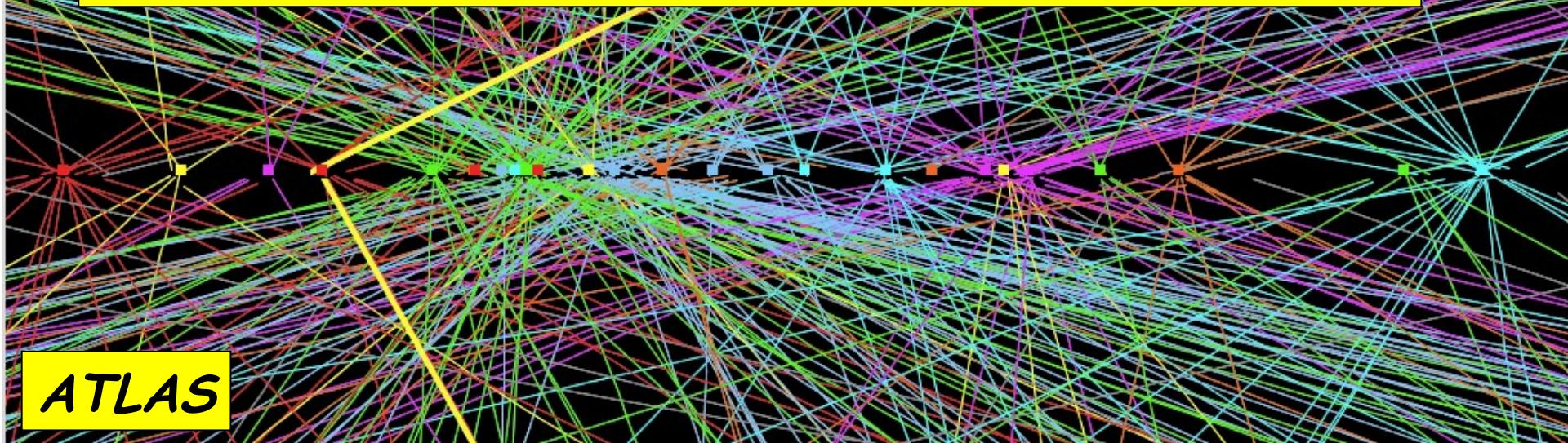


$\Delta t = 25 \text{ ns}$



$\Delta t = 50 \text{ ns}$

$Z \rightarrow \mu\mu$ event from 2012 data with 25 reconstructed vertices

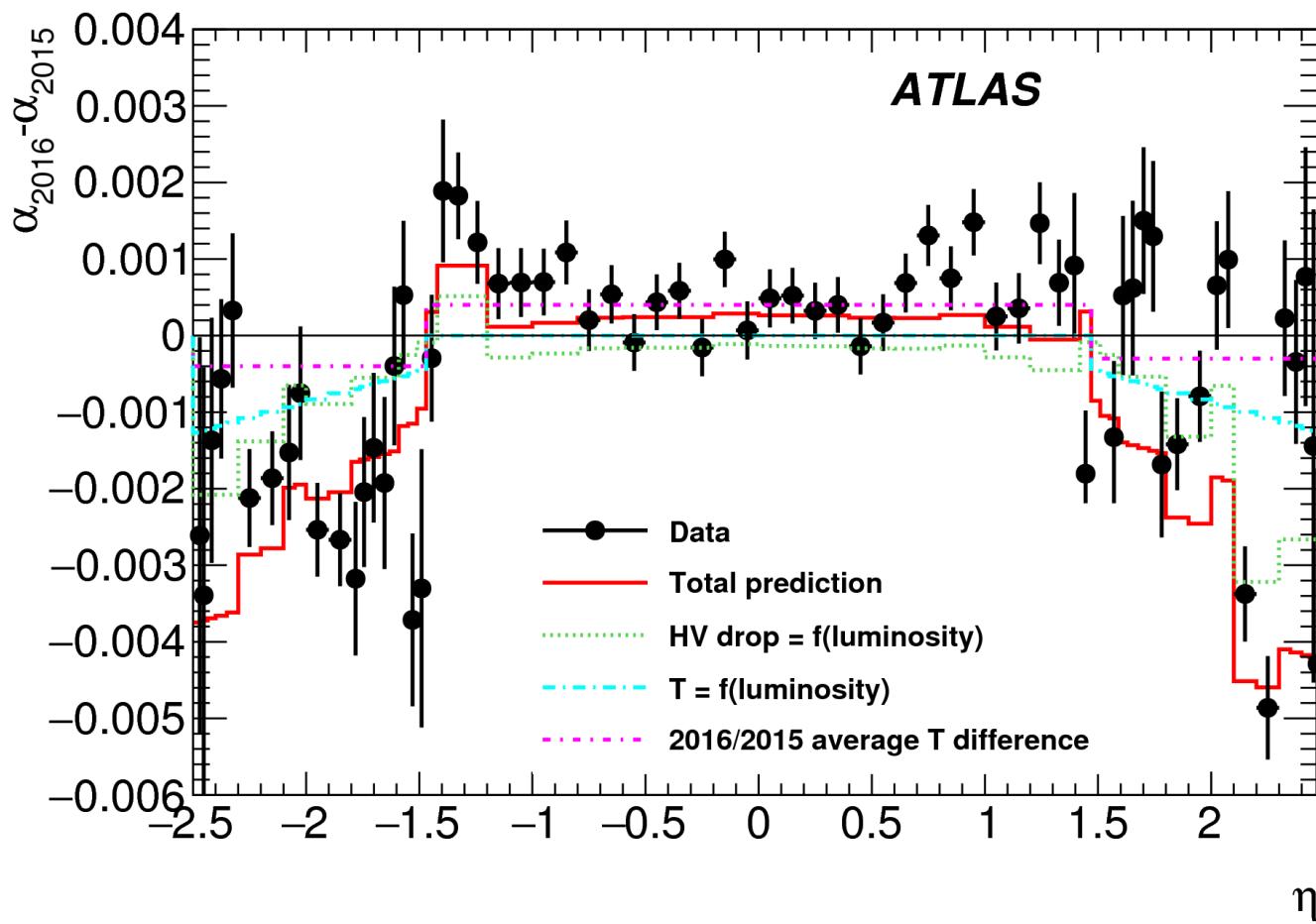


Pile up increases at higher energy (higher luminosity + higher cross sections)

→

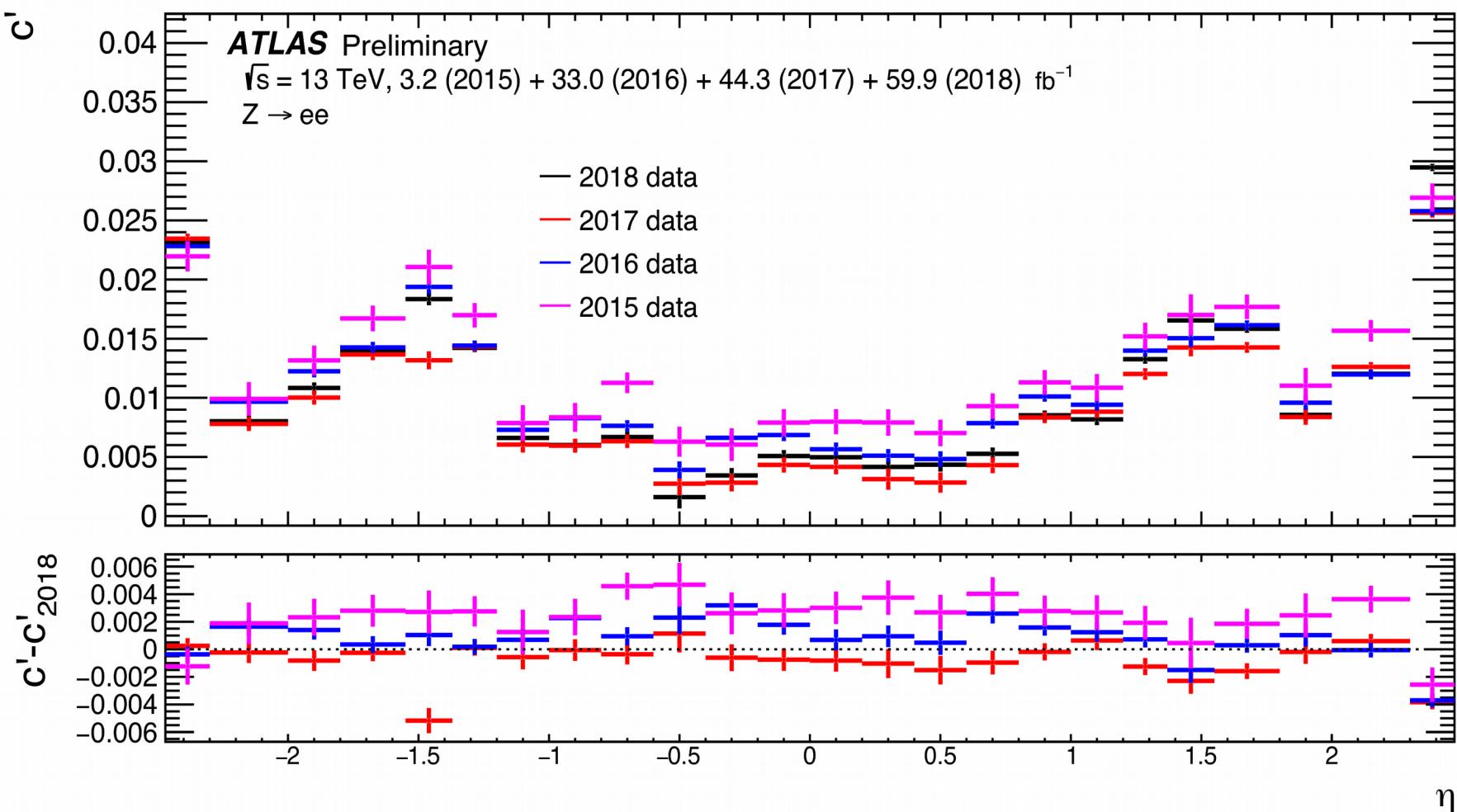
Experiments have requested 25 ns (instead of 50 ns) operation at 13 TeV

But if the time constant is larger than 50 ns (i.e integrating time of the LAr calorimeter) then the pile-up is independent of the bunch spacing (for a given luminosity)

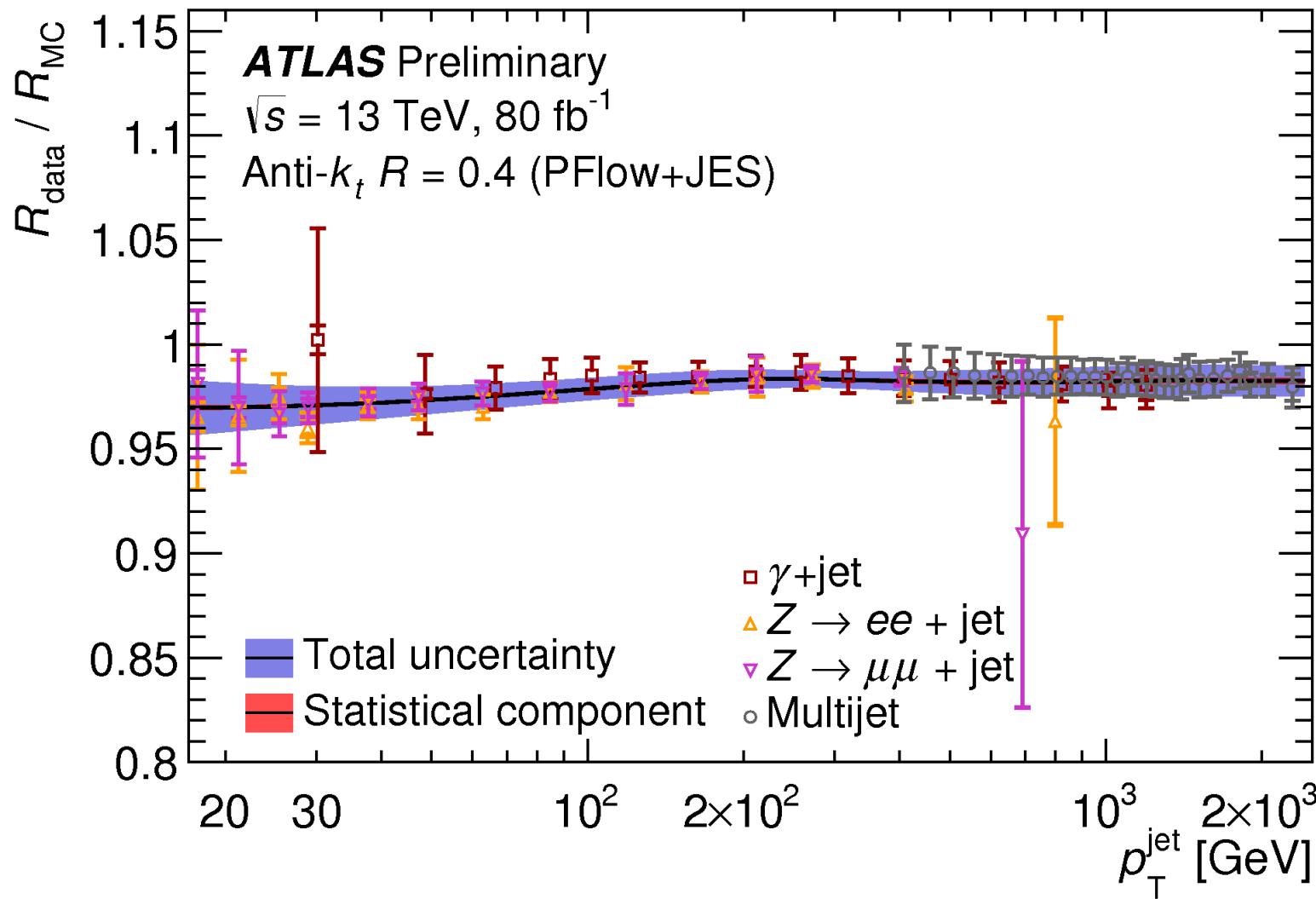


Comparison between the energy scale corrections derived from $Z \rightarrow ee$ events in 2015 and 2016 as a function of η . The difference of the energy scales measured in the data are compared with predictions taking into account the luminosity-induced high-voltage reduction and LAr temperature changes as well as the small overall difference in LAr temperature between 2015 and 2016

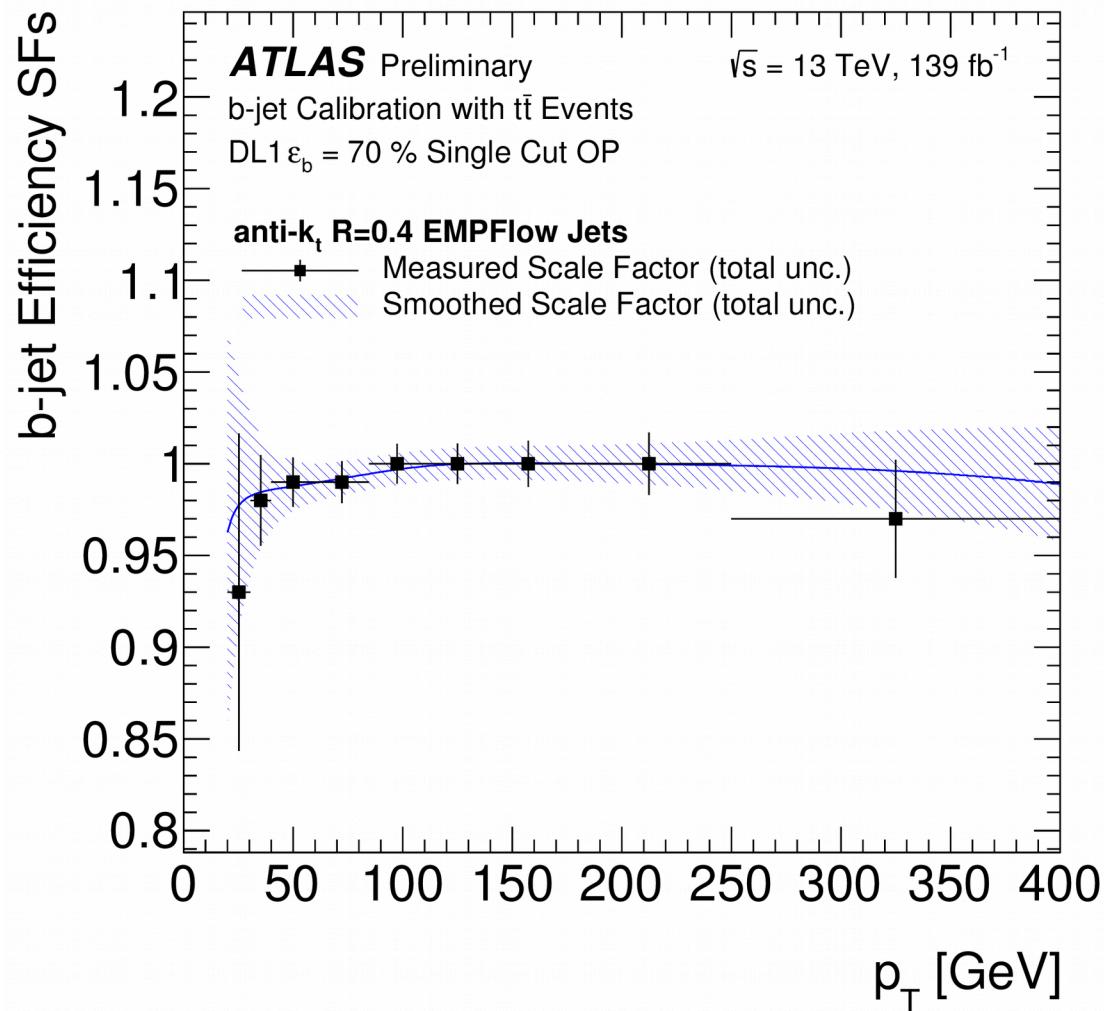
additional constant term c as a function of eta



Data driven energy calibraton of standard particle flow jets w.r.t p_T



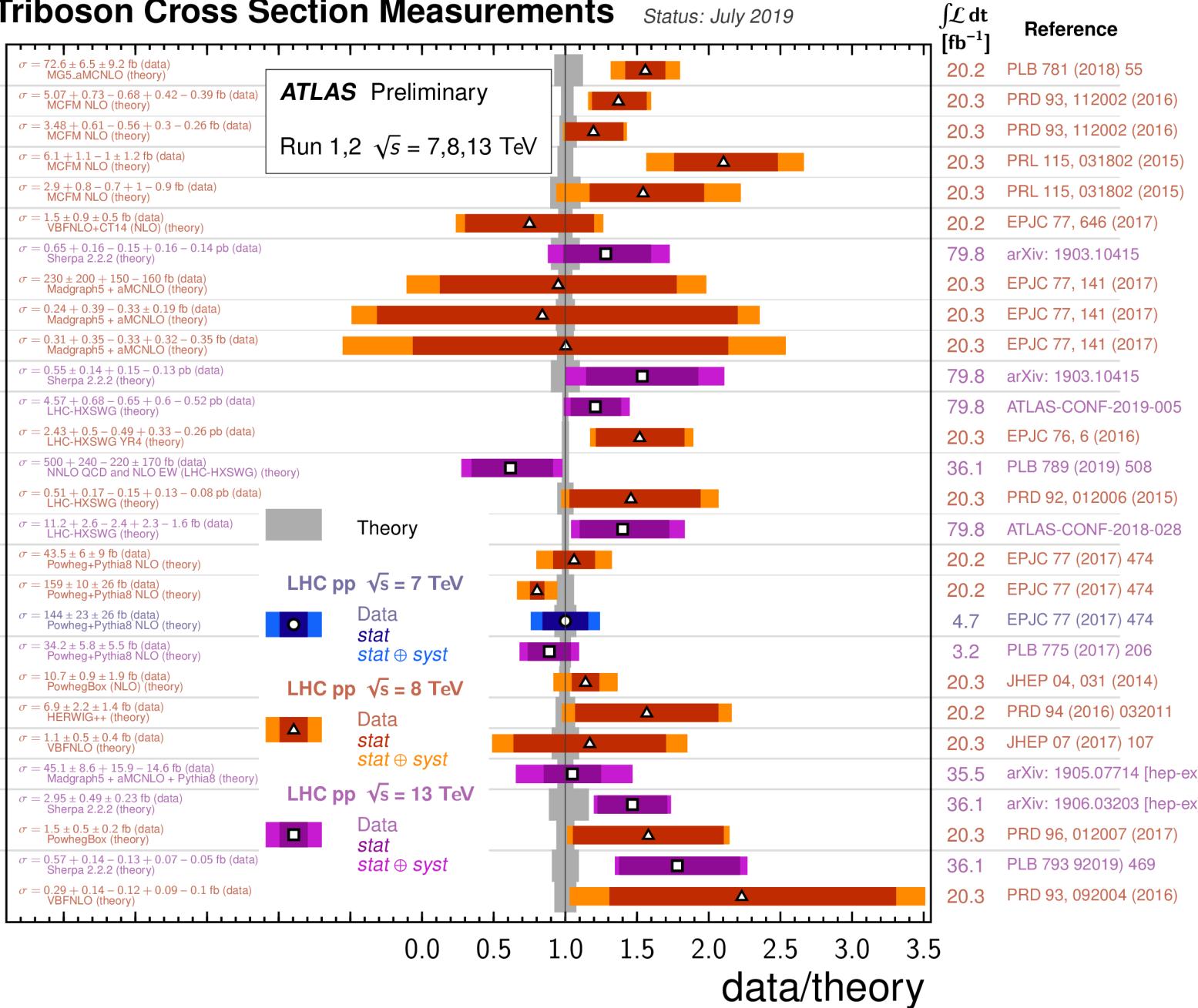
Data driven b-jet tagging efficiency w.r.t p_T

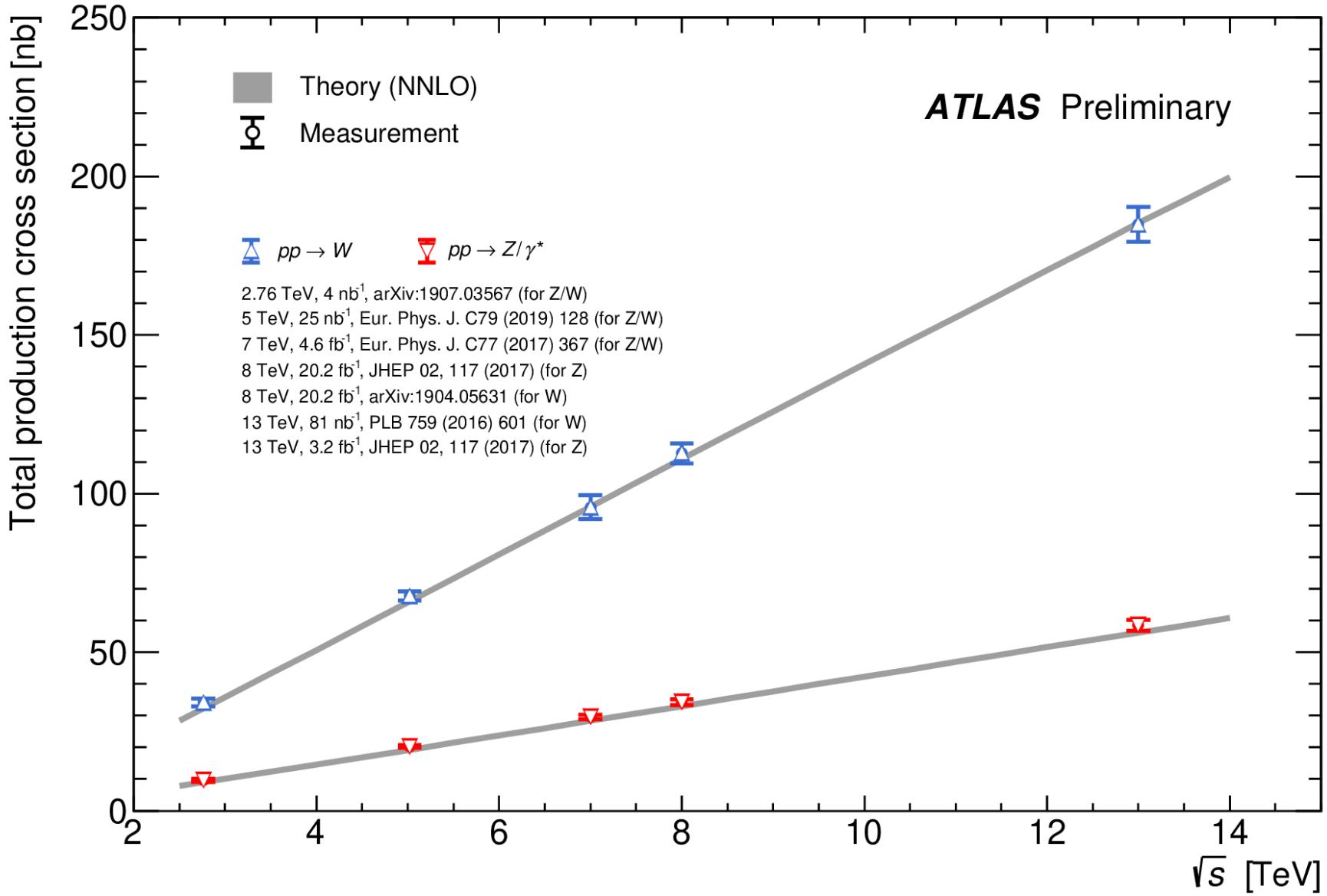


- ♪ Results from (run 1 and) run 2
 - * detector
 - * **SM**
 - * **BSM**
 - * **(B-E)H**
 - * **Vector-boson scattering**

VBF, VBS, and Triboson Cross Section Measurements

Status: July 2019





4 σ evidence for weak triboson production using 2015-2017 data

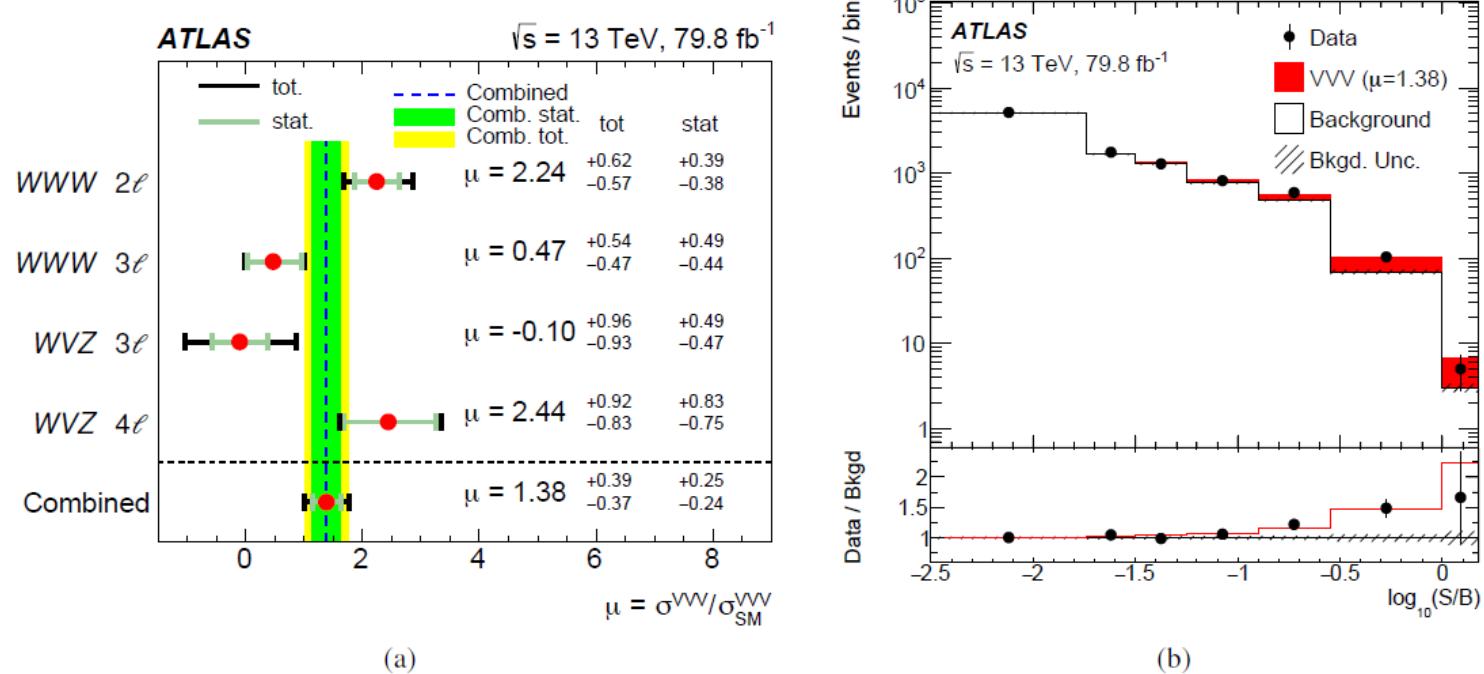


Figure 5: (a) Extracted signal strengths μ for the four analysis regions and for the combination. (b) Event yields as a function of $\log_{10}(S/B)$ for data, background B and the signal S. Events in all eleven signal regions are included. The background and signal yields are shown after the global signal-plus-background fit. The hatched band corresponds to the systematic uncertainties, and the statistical uncertainties are represented by the error bars on the data points. The lower panel shows the ratio of the data to the expected background estimated from the fit, compared to the expected distribution including the signal (red line).

Measurement of fiducial and differential $W^+ W^-$ production cross-section

sector

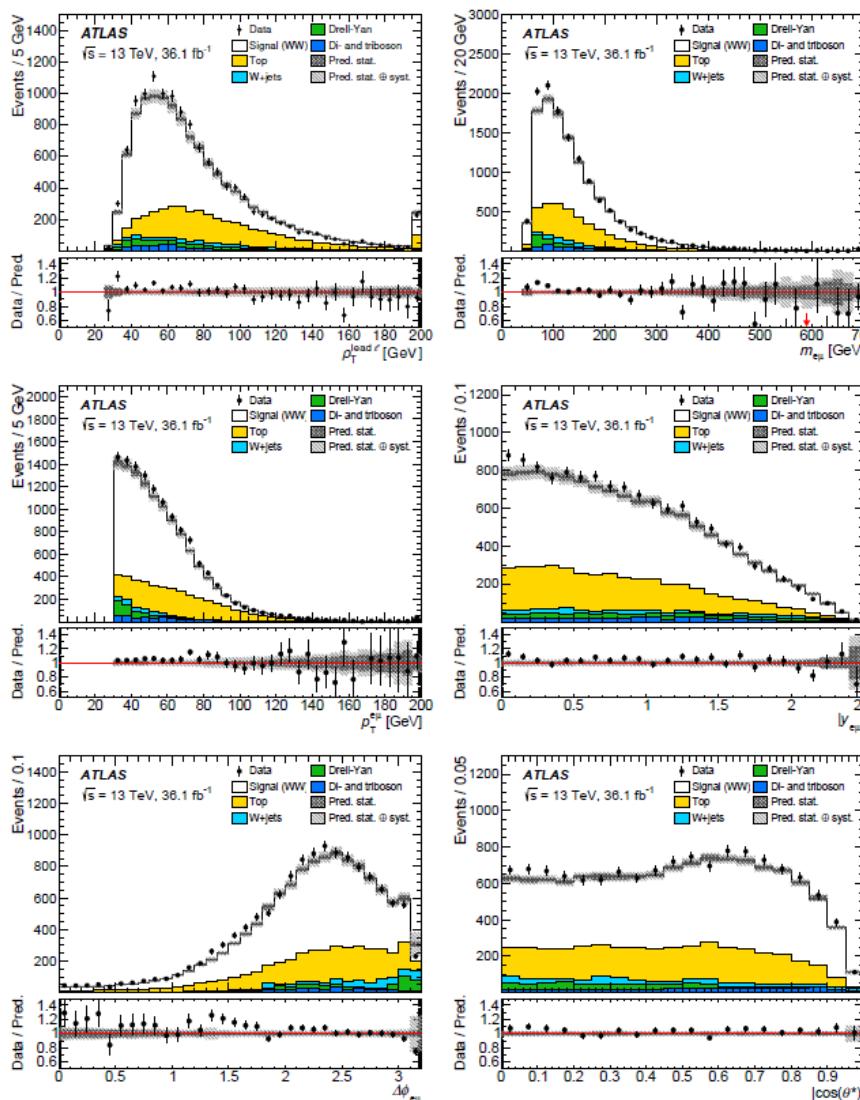


Figure 4: Kinematic distributions of the selected data events after the full event selection (from left to right and top to bottom): $p_T^{\text{lead } \ell}$, $m_{e\mu}$, $p_T^{e\mu}$, $|y_{e\mu}|$, $\Delta\phi_{e\mu}$ and $|\cos\theta^*|$. Data are shown together with the predictions of the signal and background production processes. Statistical and systematic uncertainties in the predictions are shown as hatched bands. The lower panels show the ratio of the data to the total prediction. An arrow indicates that the point is off-scale. The last bin includes the overflow.

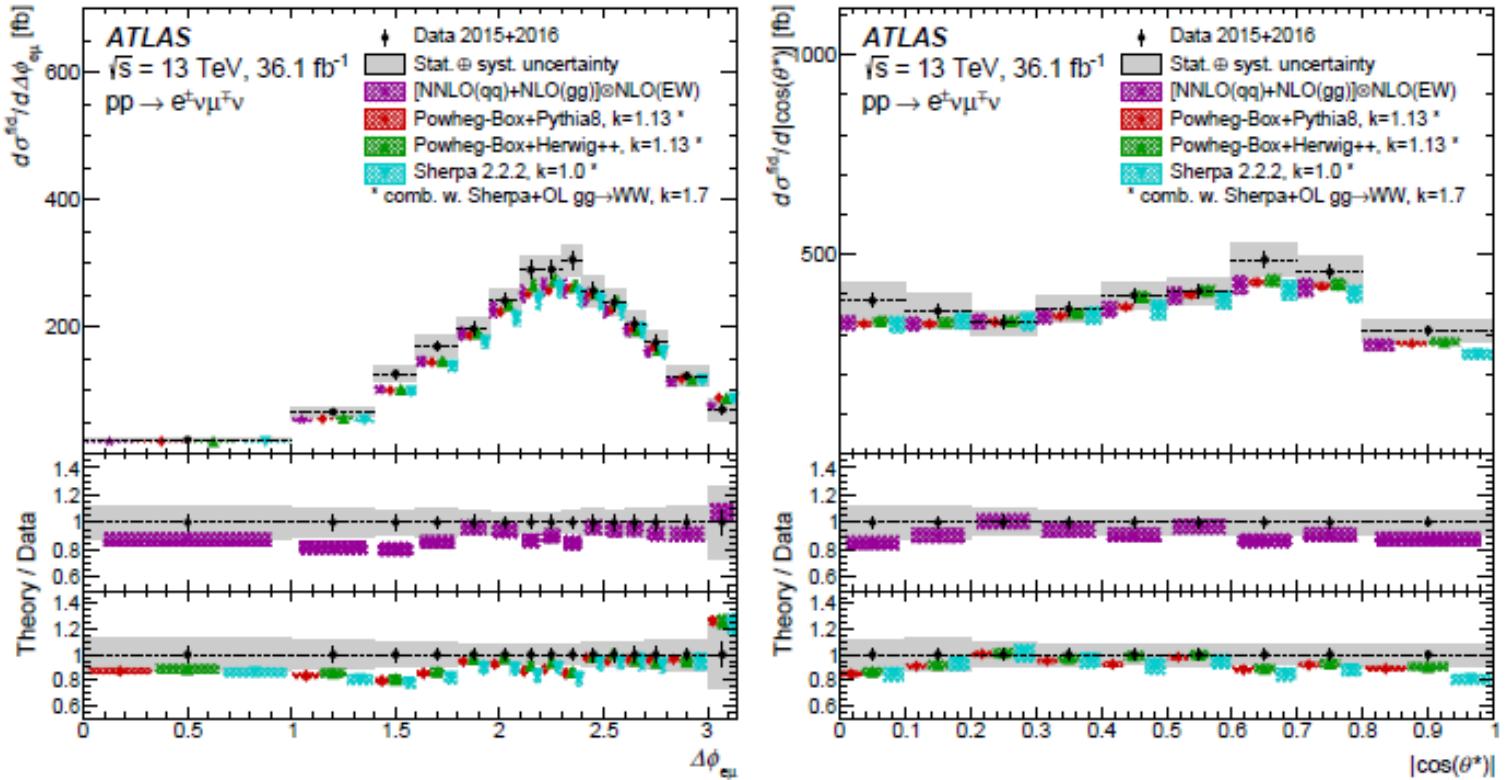


Figure 8: Measured fiducial cross-sections of $WW \rightarrow e\mu$ production for two of the six observables: $\Delta\phi_{e\mu}$ and $|\cos\theta^*|$. The measured cross-section values are shown as points with error bars giving the statistical uncertainty and solid bands indicating the size of the total uncertainty. The results are compared with the NNLO prediction with extra NLO EW corrections and NLO corrections for $gg \rightarrow WW$ production, and with NLO+PS predictions from POWHEG-BOX+PYTHIA 8, POWHEG-BOX+HERWIG++ and SHERPA 2.2.2 for $q\bar{q}$ initial states, combined with SHERPA+OPENLOOPS (LO+PS) for the gg initial states. All three $q\bar{q}$ NLO+PS predictions are normalized to the NNLO theoretical prediction for the total cross-section, with the gg LO+PS contribution normalized to NLO. Theoretical predictions are indicated as markers with hatched bands denoting PDF+scale uncertainties.

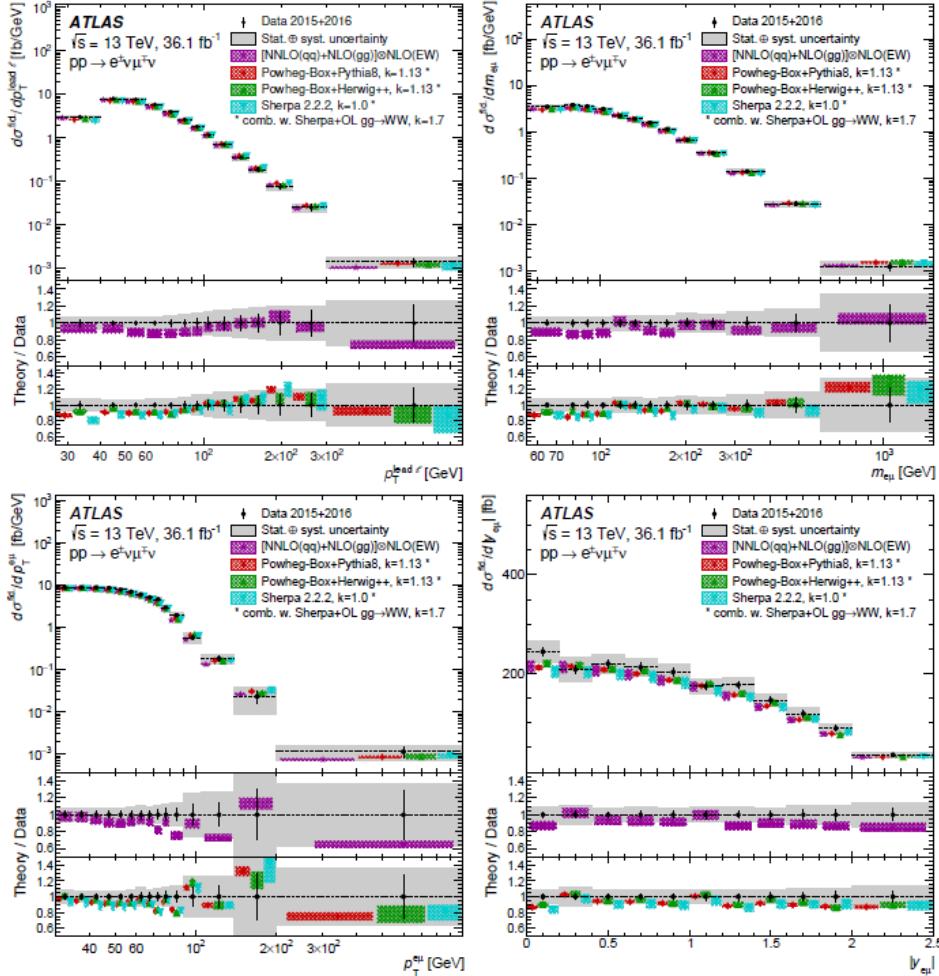
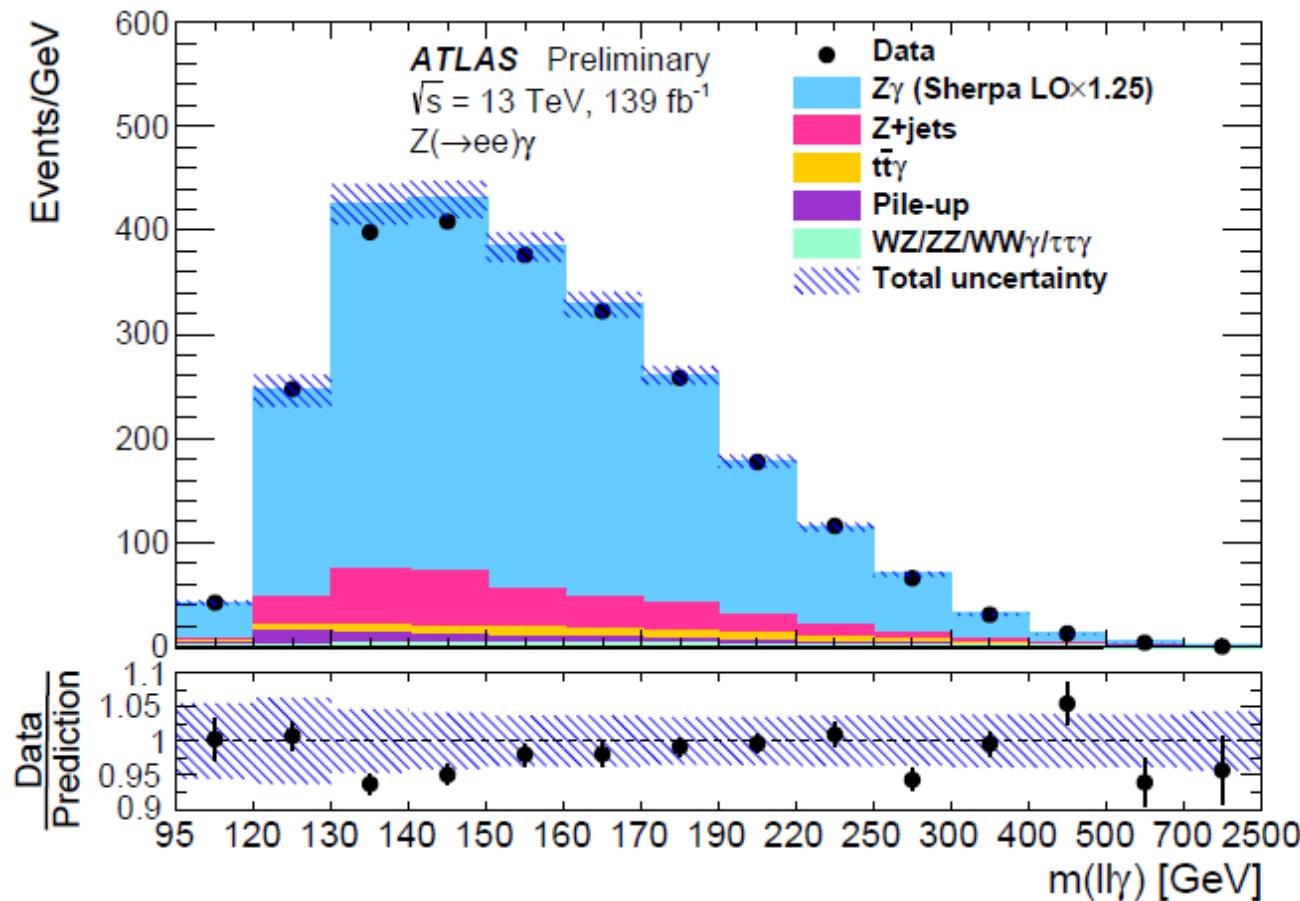


Figure 7: Measured fiducial cross-sections of $WW \rightarrow e\mu$ production for four of the six observables (from left to right and top to bottom): $p_T^{\text{lead } \ell}$, $m_{e\mu}$, $p_T^{e\mu}$, and $|y_{e\mu}|$. The measured cross-section values are shown as points with error bars giving the statistical uncertainty and solid bands indicating the size of the total uncertainty. The results are compared with the NNLO prediction with extra NLO EW corrections and NLO corrections for $gg \rightarrow WW$ production, and with NLO+PS predictions from Powheg-Box+PYTHIA 8, Powheg-Box+HERWIG++ and SHERPA 2.2.2 for $q\bar{q}$ initial states, combined with SHERPA+OPENLOOPS (LO+PS) for the gg initial states. All three $q\bar{q}$ NLO+PS predictions are normalized to the NNLO theoretical prediction for the total cross-section, with the gg LO+PS contribution normalized to NLO. Theoretical predictions are indicated as markers with hatched bands denoting PDF+scale uncertainties.

Measurement of $Z(\rightarrow l^+l^-)\gamma$ differential cross-sections in pp collisions at $\sqrt{s} = 13$ TeV with the ATLAS detector



Inclusive and differential measurements of the charge asymmetry in $t + \bar{t}$ events

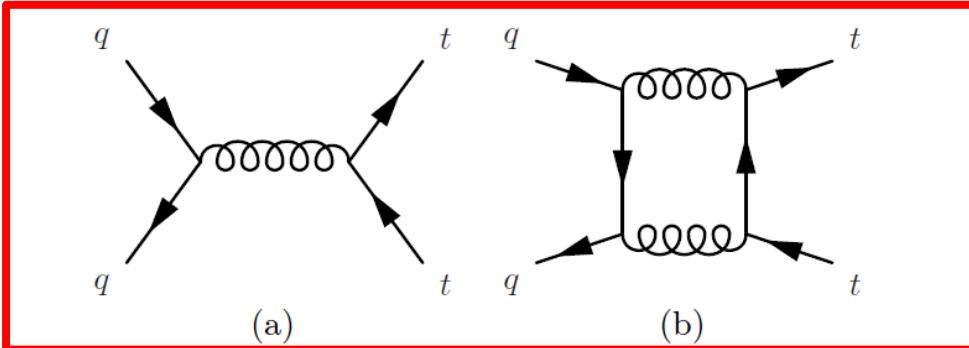
$p\bar{p}$

$$A_{FB} = \frac{N(\Delta y > 0) - N(\Delta y < 0)}{N(\Delta y > 0) + N(\Delta y < 0)}$$

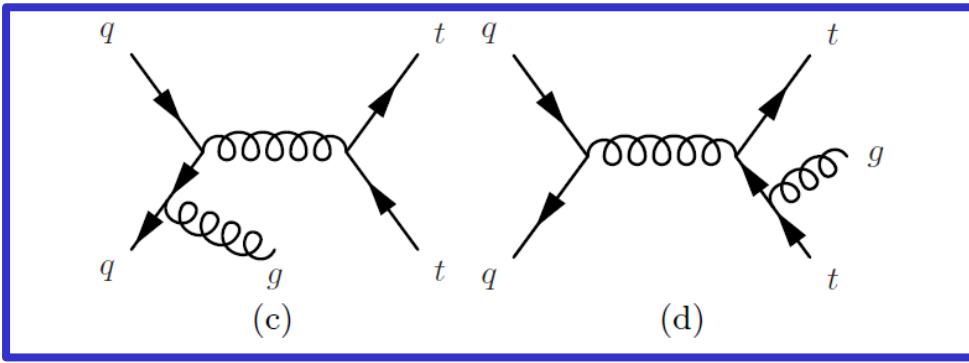
TLAS

$$\Delta y = |y_t| - |\bar{y}_t|$$

different from 0 because of interference

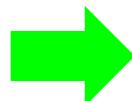


positive asymmetry



negative asymmetry

in $p\bar{p}$ collisions a FB asymmetry with a fixed \hat{z} axis vanishes



$$A_C = \frac{N(\Delta|y| > 0) - N(\Delta|y| < 0)}{N(\Delta|y| > 0) + N(\Delta|y| < 0)}$$

$$\Delta|y| = |y_t| - |\bar{y}_t|$$

Observation of light-by-light scattering in ultraperipheral Pb+Pb collisions with the ATLAS detector

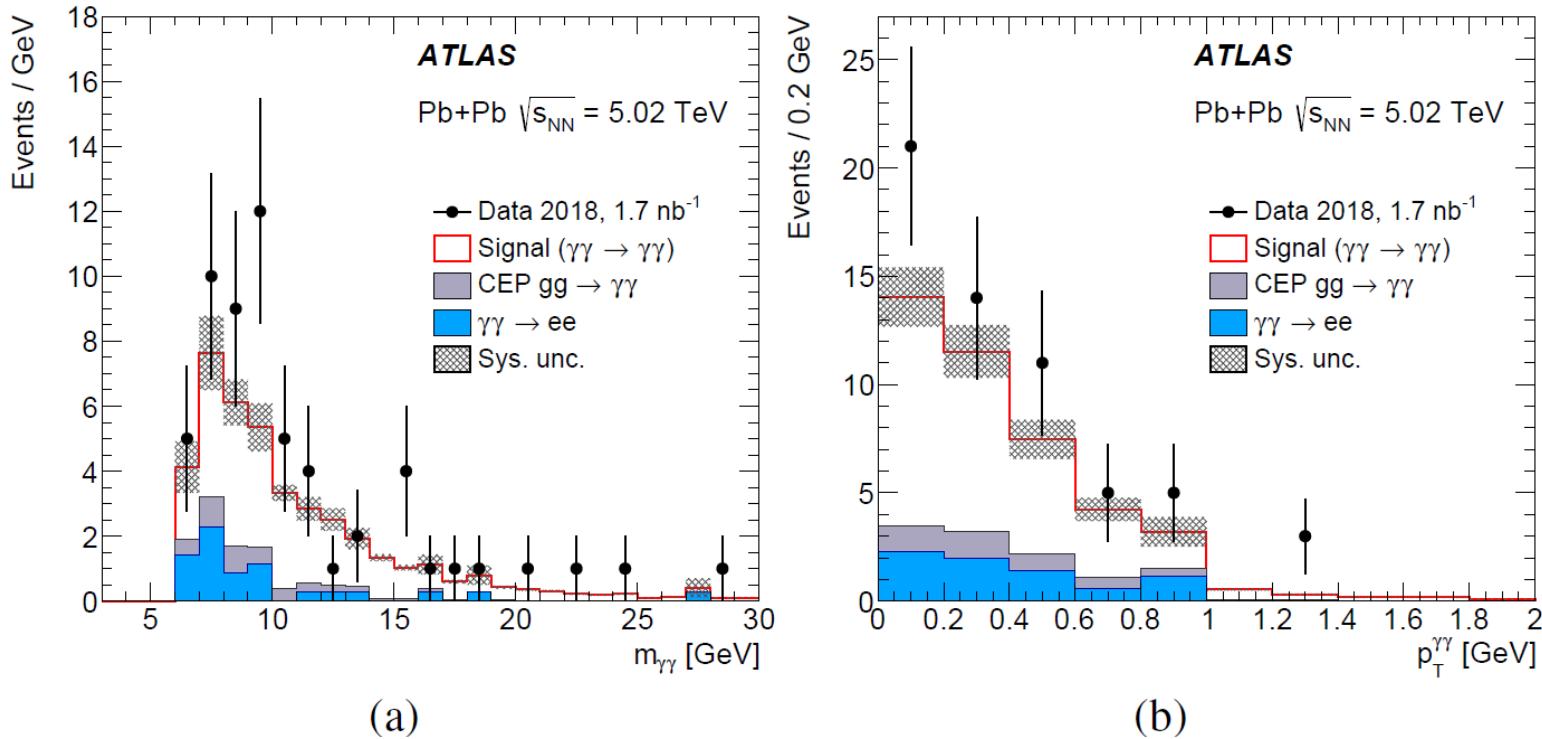


Figure 3: Kinematic distributions for $\gamma\gamma \rightarrow \gamma\gamma$ event candidates: (a) diphoton invariant mass, (b) diphoton transverse momentum. Data (points) are compared with the sum of signal and background expectations (histograms). Systematic uncertainties of the signal and background processes, excluding that of the luminosity, are shown as shaded bands.

$$A_\phi = (1 - |\Delta\phi_{\gamma\gamma}|/\pi) < 0.01$$

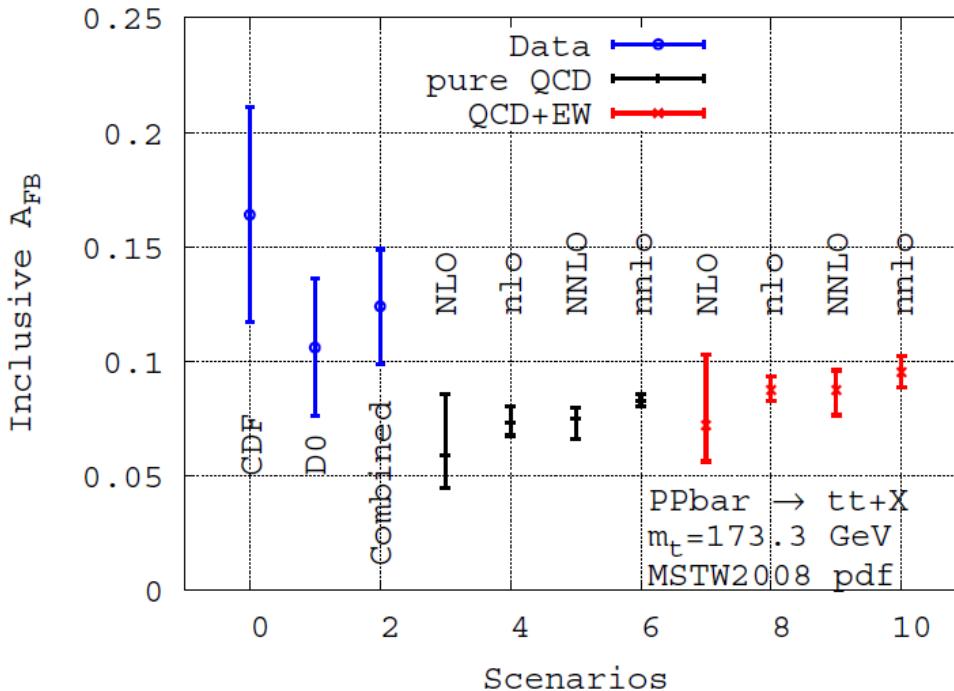


FIG. 1: The inclusive asymmetry in pure QCD (black) and QCD+EW[28] (red). Capital letters (NLO, NNLO) correspond to the unexpanded definition (2), while small letters (nlo, nnlo) to the definition (3). The CDF/DØ (naive) average is from Ref. [29]. Error bands are from scale variation only. Our final prediction corresponds to scenario 10.

FCNC (Flavour-Changing Neutral Current)

arXiv:1908.08461

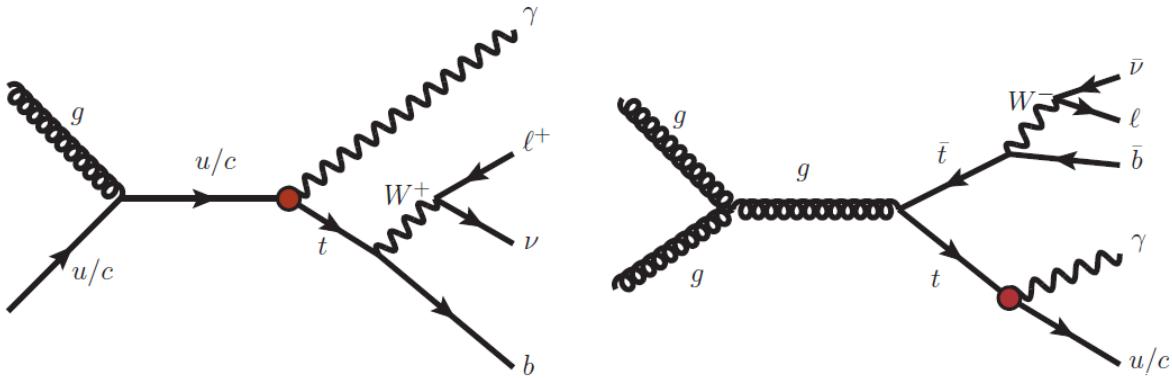
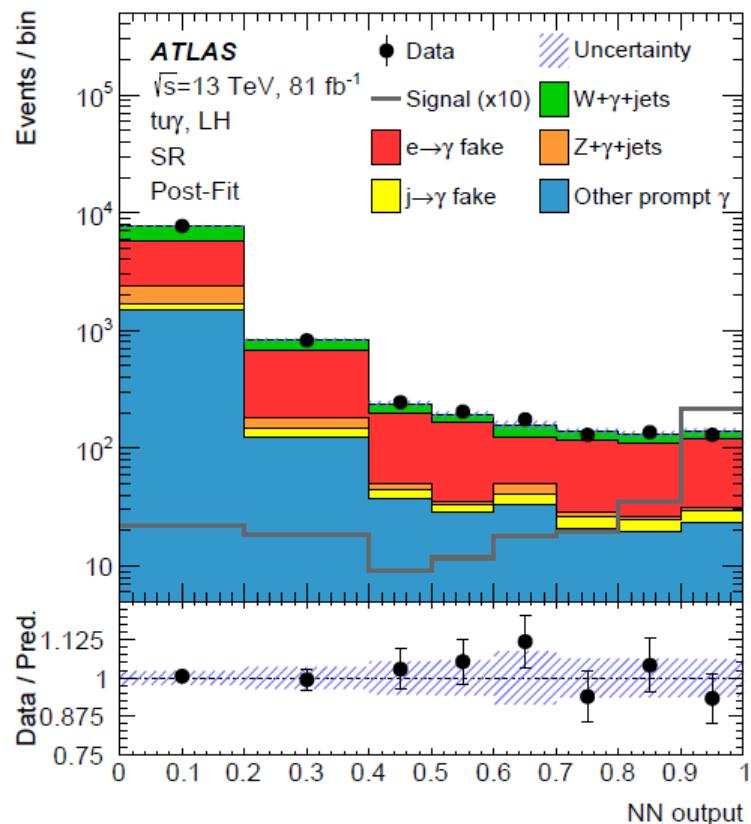


Figure 1: Tree-level Feynman diagrams for top-quark production (left) and decay (right) via FCNCs. The $tq\gamma$ vertex, which is not present in the SM, is highlighted.

Table 1: Observed (expected) 95% CL limits on the effective coupling strengths for different vertices and couplings, the production cross section, and the branching ratio. For the former, the energy scale is assumed to be $\Lambda = 1$ TeV.

Observable	Vertex	Coupling	Obs.	Exp.
$ C_{uW}^{(13)*} + C_{uB}^{(13)*} $	$t u \gamma$	LH	0.19	$0.22^{+0.04}_{-0.03}$
$ C_{uW}^{(31)} + C_{uB}^{(31)} $	$t u \gamma$	RH	0.27	$0.27^{+0.05}_{-0.04}$
$ C_{uW}^{(23)*} + C_{uB}^{(23)*} $	$t c \gamma$	LH	0.52	$0.57^{+0.11}_{-0.09}$
$ C_{uW}^{(32)} + C_{uB}^{(32)} $	$t c \gamma$	RH	0.48	$0.59^{+0.12}_{-0.09}$
$\sigma(pp \rightarrow t\gamma)$ [fb]	$t u \gamma$	LH	36	52^{+21}_{-14}
$\sigma(pp \rightarrow t\gamma)$ [fb]	$t u \gamma$	RH	78	75^{+31}_{-21}
$\sigma(pp \rightarrow t\gamma)$ [fb]	$t c \gamma$	LH	40	49^{+20}_{-14}
$\sigma(pp \rightarrow t\gamma)$ [fb]	$t c \gamma$	RH	33	52^{+22}_{-14}
$\mathcal{B}(t \rightarrow q\gamma)[10^{-5}]$	$t u \gamma$	LH	2.8	$4.0^{+1.6}_{-1.1}$
$\mathcal{B}(t \rightarrow q\gamma)[10^{-5}]$	$t u \gamma$	RH	6.1	$5.9^{+2.4}_{-1.6}$
$\mathcal{B}(t \rightarrow q\gamma)[10^{-5}]$	$t c \gamma$	LH	22	27^{+11}_{-7}
$\mathcal{B}(t \rightarrow q\gamma)[10^{-5}]$	$t c \gamma$	RH	18	28^{+12}_{-8}



FCNC (Flavour-Changing Neutral Current)

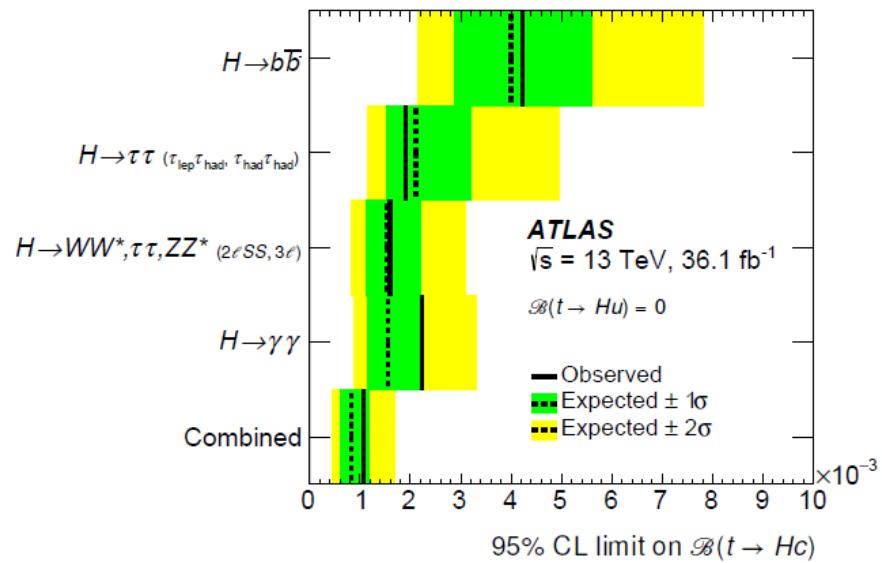
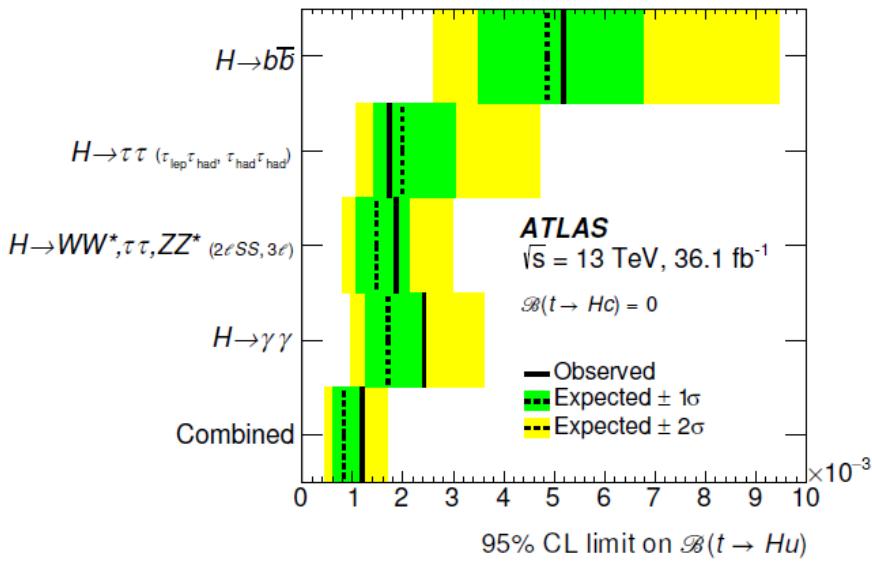


Table 3: Summary of 95% CL upper limits on $\mathcal{B}(t \rightarrow Hc)$ and $\mathcal{B}(t \rightarrow Hu)$, in each case neglecting the other decay mode. Signatures with two same-charge (three) leptons and no τ_{had} candidates are denoted by 2 ℓ SS (3 ℓ).

	95% CL upper limits on $\mathcal{B}(t \rightarrow Hc)$ Observed (Expected)	95% CL upper limits on $\mathcal{B}(t \rightarrow Hu)$ Observed (Expected)
$H \rightarrow b\bar{b}$	4.2×10^{-3} (4.0×10^{-3})	5.2×10^{-3} (4.9×10^{-3})
$H \rightarrow \tau\tau (\tau_{\text{lep}}, \tau_{\text{had}}, \tau_{\text{had}})$	1.9×10^{-3} (2.1×10^{-3})	1.7×10^{-3} (2.0×10^{-3})
$H \rightarrow WW^*, \tau\tau, ZZ^* (2\ell\text{SS}, 3\ell)$ [22]	1.6×10^{-3} (1.5×10^{-3})	1.9×10^{-3} (1.5×10^{-3})
$H \rightarrow \gamma\gamma$ [21]	2.2×10^{-3} (1.6×10^{-3})	2.4×10^{-3} (1.7×10^{-3})
Combination	1.1×10^{-3} (8.3×10^{-4})	1.2×10^{-3} (8.3×10^{-4})

EW precision measurements

Weak angle $\sin^2\theta_{\text{eff}}^{\text{l}}$

$$\frac{d\sigma}{dp_T^{\ell\ell} dy^{\ell\ell} dm^{\ell\ell} d\cos\theta d\phi} = \frac{3}{16\pi} \frac{d\sigma^{U+L}}{dp_T^{\ell\ell} dy^{\ell\ell} dm^{\ell\ell}} \\ \left\{ (1 + \cos^2\theta) + \frac{1}{2} A_0(1 - 3\cos^2\theta) + A_1 \sin 2\theta \cos\phi \right. \\ \left. + \frac{1}{2} A_2 \sin^2\theta \cos 2\phi + A_3 \sin\theta \cos\phi + A_4 \cos\theta \right. \\ \left. + A_5 \sin^2\theta \sin 2\phi + A_6 \sin 2\theta \sin\phi + A_7 \sin\theta \sin\phi \right\}$$

$$A_{\text{FB}} = 3/8 \times A_4$$

EW precision measurements

W mass m_W

arXiv:1701.07240

$$\begin{aligned} m_W &= 80369.5 \pm 6.8 \text{ MeV(stat.)} \pm 10.6 \text{ MeV(exp. syst.)} \pm 13.6 \text{ MeV(mod. syst.)} \\ &= 80369.5 \pm 18.5 \text{ MeV,} \end{aligned}$$

Combined categories	Value [MeV]	Stat. Unc.	Muon Unc.	Elec. Unc.	Recoil Unc.	Bckg. Unc.	QCD Unc.	EW Unc.	PDF Unc.	Total Unc.	χ^2/dof of Comb.
m_T - p_T^ℓ , W^\pm , e - μ	80369.5	6.8	6.6	6.4	2.9	4.5	8.3	5.5	9.2	18.5	29/27

B physics

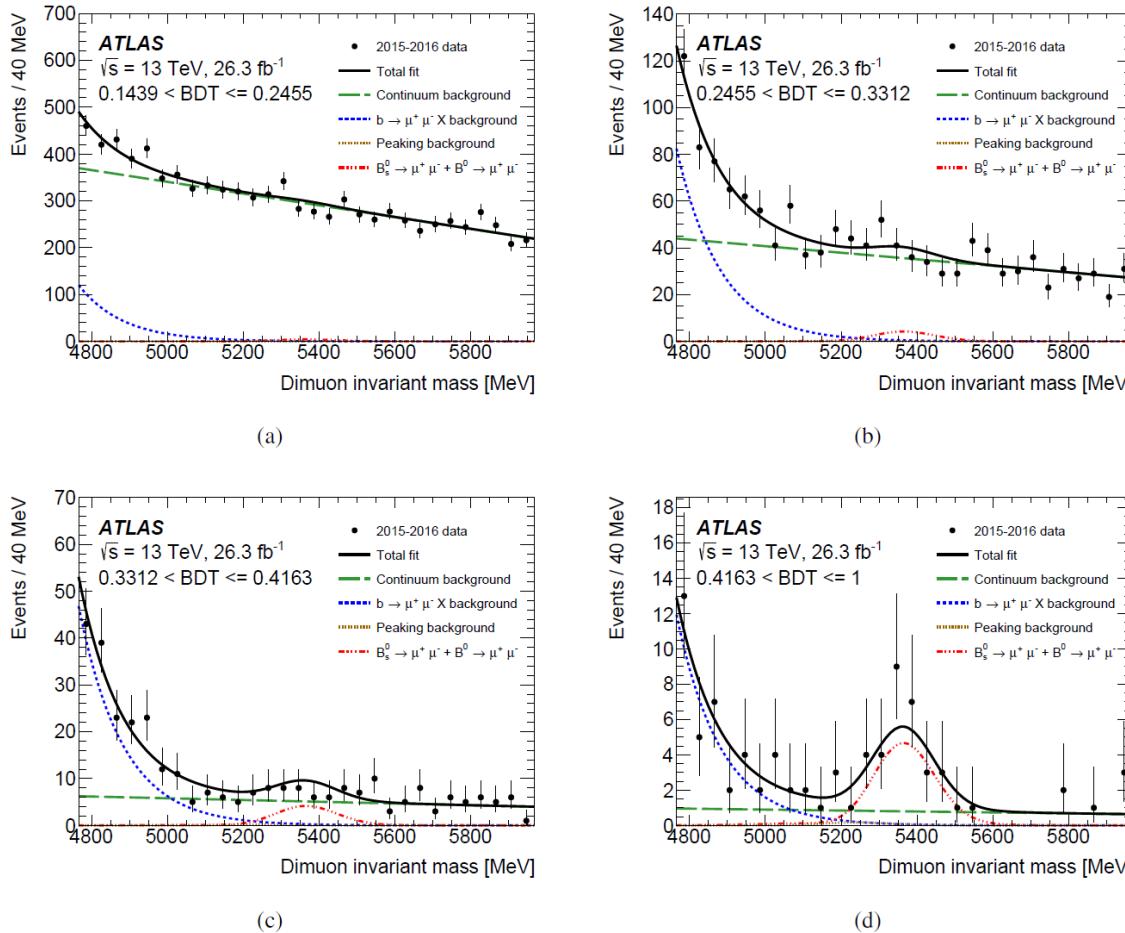
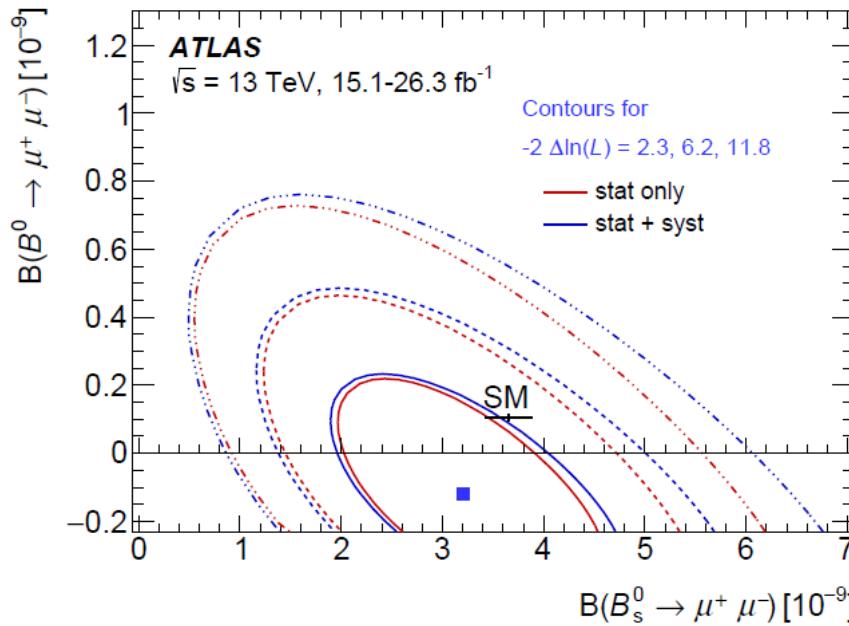
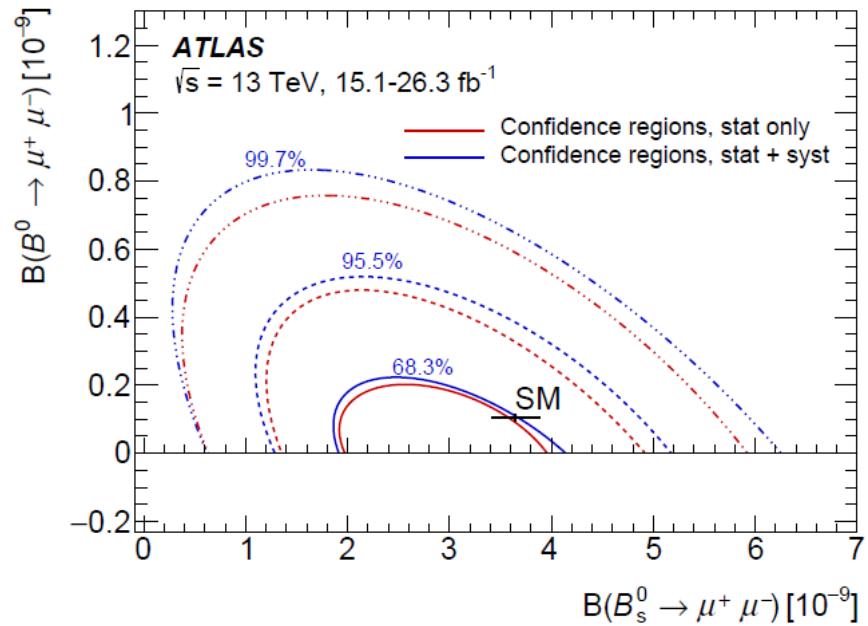


Figure 8: Dimuon invariant mass distributions in the unblinded data, in the four intervals of BDT output. Superimposed is the result of the maximum-likelihood fit. The total fit is shown as a continuous line, with the dashed lines corresponding to the observed signal component, the $b \rightarrow \mu\mu X$ background, and the continuum background. The signal components are grouped in one single curve, including both the $B_s^0 \rightarrow \mu^+\mu^-$ and the (negative) $B^0 \rightarrow \mu^+\mu^-$ component. The curve representing the peaking $B_{(s)}^0 \rightarrow hh'$ background lies very close to the horizontal axis in all BDT bins.

B physics



(a)



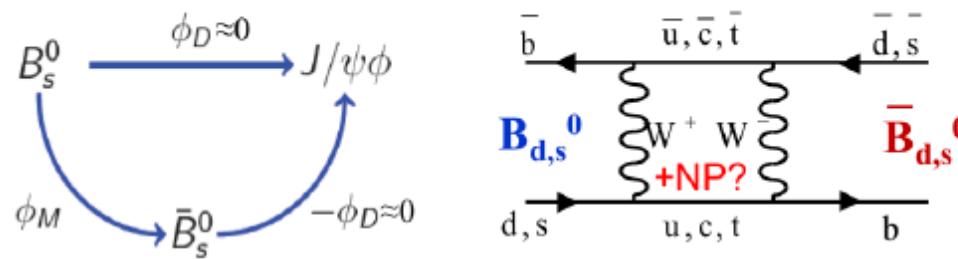
(b)

Figure 9: (a) Likelihood contours for the simultaneous fit to $\mathcal{B}(B_s^0 \rightarrow \mu^+ \mu^-)$ and $\mathcal{B}(B^0 \rightarrow \mu^+ \mu^-)$, for values of $-2\Delta \ln(\mathcal{L})$ equal to 2.3, 6.2 and 11.8. The SM prediction with uncertainties is indicated. (b) Neyman contours in the $\mathcal{B}(B_s^0 \rightarrow \mu^+ \mu^-)$ – $\mathcal{B}(B^0 \rightarrow \mu^+ \mu^-)$ plane for 68.3%, 95.5% and 99.7% coverage. At each $-2\Delta \ln(\mathcal{L})$ or coverage value, the inner contours are statistical uncertainty only, while the outer ones include statistical and systematic uncertainties. The construction of these contours makes use of both the dimuon (26.3 fb^{-1}) and the reference channel (15.1 fb^{-1}) datasets.

$$\begin{aligned}\phi_s &= -0.076 \pm 0.034 \text{ (stat.)} \pm 0.019 \text{ (syst.) rad} \\ \Delta\Gamma_s &= 0.068 \pm 0.004 \text{ (stat.)} \pm 0.003 \text{ (syst.) ps}^{-1} \\ \Gamma_s &= 0.669 \pm 0.001 \text{ (stat.)} \pm 0.001 \text{ (syst.) ps}^{-1}\end{aligned}$$

Measurement of $\Delta\Gamma_s$ and ϕ_s in $B_s \rightarrow J/\psi(\mu\mu) \varphi(KK)$

- CP violation in $B_s \rightarrow J/\psi \varphi$ occurs through the interference in mixing and decay.
- The time evolution of flavour tagged is very sensitive New Physics
- B_s mixing:
 - Mass difference
 - $\Delta m = m_H - m_L$
 - Mixing phase ϕ_s
 - Decay width difference $\Delta\Gamma_s = \Gamma_L - \Gamma_H$
- 9 Physics parameters describe $B_s \rightarrow J/\psi \varphi$ decay:



$\Gamma_s, \Delta\Gamma_s$	decay with and decay width difference
$\phi_s (\approx 2\beta_s)$	CP violating phase
$ A_0 ^2, A_\parallel ^2$	CP state amplitudes
$\delta_\parallel, \delta_\perp$	Strong phases
$ A_\circ ^2, \delta_\circ$	S-wave parameters

- ♪ Results from (run 1 and) run 2
 - * detector
 - * SM
 - * **BSM**
 - * $(B-E)H$
 - * Vector-boson scattering

Search for Physics BSM

ATLAS SUSY Searches* - 95% CL Lower Limits

July 2019

ATLAS Preliminary

$\sqrt{s} = 13 \text{ TeV}$

Model	Signature	$\int \mathcal{L} dt [\text{fb}^{-1}]$	Mass limit	Reference
Inclusive Searches	$q\bar{q}, q\rightarrow q\tilde{\chi}_1^0$ mono-jet	0 e, μ 1-3 jets E_T^{miss}	36.1 36.1	$m(\tilde{\chi}_1^0) < 100 \text{ GeV}$ $m(q\bar{q}-m(\tilde{\chi}_1^0)) < 5 \text{ GeV}$
	$q\bar{q}, q\rightarrow q\tilde{\chi}_1^0$	0 e, μ 2-6 jets E_T^{miss}	36.1	$m(\tilde{\chi}_1^0) < 200 \text{ GeV}$ $m(\tilde{\chi}_1^0) < 900 \text{ GeV}$
	$q\bar{q}, q\rightarrow q\tilde{\chi}_1^0 (\ell\ell)\tilde{\chi}_1^0$	3 e, μ $ee, \mu\mu$ 2 jets E_T^{miss}	36.1 36.1	$m(\tilde{\chi}_1^0) < 800 \text{ GeV}$ $m(\tilde{\chi}_1^0) < 50 \text{ GeV}$
	$q\bar{q}, q\rightarrow q\tilde{W}Z\tilde{\chi}_1^0$	0 e, μ SS e, μ 7-11 jets E_T^{miss}	36.1 139	$m(\tilde{\chi}_1^0) < 400 \text{ GeV}$ $m(\tilde{W})-m(\tilde{\chi}_1^0) < 200 \text{ GeV}$
	$q\bar{q}, q\rightarrow t\tilde{\chi}_1^0$	0-1 e, μ SS e, μ 3 b 6 jets E_T^{miss}	79.8 139	$m(\tilde{\chi}_1^0) < 200 \text{ GeV}$ $m(\tilde{t})-m(\tilde{\chi}_1^0) < 300 \text{ GeV}$
3 rd Gen. direct production	$b_1 b_1, b_1 \rightarrow b\tilde{\chi}_1^0/\tilde{\chi}_1^\pm$	Multiple	36.1	$m(\tilde{\chi}_1^0) < 300 \text{ GeV}, BR(b\tilde{\chi}_1^0) < 1$
	$b_1 b_1, b_1 \rightarrow b\tilde{\chi}_1^0/b\tilde{\chi}_1^\pm$	Multiple	36.1	$m(\tilde{\chi}_1^0) < 300 \text{ GeV}, BR(b\tilde{\chi}_1^0)-BR(b\tilde{\chi}_1^\pm) < 0.5$
	$b_1 b_1, b_1 \rightarrow b\tilde{\chi}_1^0 \rightarrow b\tilde{W}\tilde{\chi}_1^0$	0 e, μ 6 b E_T^{miss}	139	$m(\tilde{\chi}_1^0) < 200 \text{ GeV}, m(\tilde{W})-m(\tilde{\chi}_1^0) < 1$
	$t_1 \bar{t}_1, \bar{t}_1 \rightarrow W\tilde{b}_1^0 \text{ or } \chi_1^0$	0-2 e, μ 0-2 jets/1-2 b E_T^{miss}	36.1	$m(\tilde{\chi}_1^0) < 1 \text{ GeV}$
	$t_1 \bar{t}_1, \bar{t}_1 \rightarrow W\tilde{b}_1^0$	1 e, μ 3 jets/1 b E_T^{miss}	139	$m(\tilde{\chi}_1^0) < 400 \text{ GeV}$
	$t_1 \bar{t}_1, \bar{t}_1 \rightarrow T_1 b\bar{v}_1, T_1 \rightarrow rG$	1 $\tau + 1 \ e, \mu, \tau$ 2 jets/1 b E_T^{miss}	36.1	$m(\tilde{\chi}_1^0) < 800 \text{ GeV}$
	$t_1 \bar{t}_1, \bar{t}_1 \rightarrow c\tilde{\chi}_1^0/2\tilde{e}, \tilde{e} \rightarrow c\tilde{\chi}_1^0$	0 e, μ 2 c E_T^{miss}	36.1	$m(\tilde{\chi}_1^0) < 0 \text{ GeV}$
	$t_1 \bar{t}_1, \bar{t}_1 \rightarrow c\tilde{\chi}_1^0/2\tilde{e}, \tilde{e} \rightarrow c\tilde{\chi}_1^0$	0 e, μ mono-jet E_T^{miss}	36.1	$m(\tilde{\chi}_1^0) < 50 \text{ GeV}$
	$t_1 \bar{t}_2, \bar{t}_2 \rightarrow \bar{t}_1 + h$	1-2 e, μ 4 b E_T^{miss}	36.1	$m(\tilde{\chi}_1^0) < 180 \text{ GeV}, m(\tilde{t}_1)-m(\tilde{\chi}_1^0) < 10 \text{ GeV}$
	$t_1 \bar{t}_2, \bar{t}_2 \rightarrow \bar{t}_1 + Z$	3 e, μ 1 b E_T^{miss}	139	$m(\tilde{\chi}_1^0) < 360 \text{ GeV}, m(\tilde{t}_1)-m(\tilde{\chi}_1^0) < 40 \text{ GeV}$
EW direct	$\tilde{\chi}_1^{\pm 0} \text{ via } WZ$	2-3 e, μ $ee, \mu\mu$ ≥ 1 E_T^{miss}	36.1 139	$m(\tilde{\chi}_1^0)-0$
	$\tilde{\chi}_1^{\pm 0} \text{ via } WW$	2 e, μ E_T^{miss}	139	$m(\tilde{\chi}_1^0)-0$
	$\tilde{\chi}_1^{\pm 0} \text{ via } Wh$	0-1 e, μ 2 $b/2 \gamma$ E_T^{miss}	139	$m(\tilde{\chi}_1^0)-70 \text{ GeV}$
	$\tilde{\chi}_1^{\pm 0} \text{ via } \tilde{t}_1 \tilde{e}_1 \gamma$	2 e, μ 2τ E_T^{miss}	139	$m(\tilde{t}_1, \tilde{e}_1)-0.5(m(\tilde{\chi}_1^0)+m(\tilde{t}_1))$
	$\tilde{t}_1 \bar{t}_1, \tilde{t}_1 \rightarrow \tilde{t}_1 \tilde{\chi}_1^0$	2 e, μ 0 jets E_T^{miss}	139	$m(\tilde{\chi}_1^0)-0$
	$\tilde{t}_1 \bar{t}_1, \tilde{t}_1 \rightarrow \tilde{t}_1 \tilde{\chi}_1^0$	2 e, μ ≥ 1 E_T^{miss}	139	$m(\tilde{t}_1)-m(\tilde{\chi}_1^0)-10 \text{ GeV}$
	$B\bar{B}, B \rightarrow h\tilde{G}/Z\tilde{G}$	0 e, μ 4 e, μ $\geq 3 \ b$ 0 jets E_T^{miss}	36.1 36.1	$BR(\tilde{\chi}_1^0 \rightarrow h\tilde{\chi}_1^0)=1$ $BR(\tilde{\chi}_1^0 \rightarrow Z\tilde{\chi}_1^0)=1$
Long-lived particles	Direct $\tilde{\chi}_1^{\pm 0}$ prod., long-lived $\tilde{\chi}_1^{\pm}$	Disapp. trk 1 jet E_T^{miss}	36.1	Pure Wino Pure Higgsino
	Stable $\tilde{\chi}_1^0$ R-hadron	Multiple	36.1	
	Metastable $\tilde{\chi}_1^0$ R-hadron, $\tilde{\chi}_1^0 \rightarrow q\bar{q}\tilde{\chi}_1^0$	Multiple	36.1	$m(\tilde{\chi}_1^0) > 100 \text{ GeV}$
RPV	LFV $p\bar{p} \rightarrow \tau \nu, \tau \nu \rightarrow e\mu/e\tau/\mu\tau$	$e\mu, e\tau, \mu\tau$	3.2	$\chi_{111}=0.11, \chi_{133}/\chi_{111}=0.07$
	$\tilde{\chi}_1^{\pm 0} \tilde{\chi}_2^{\mp 0} \rightarrow WW/Z\tilde{t}_1 \tilde{t}_2 \nu \nu$	4 e, μ 0 jets E_T^{miss}	36.1	$m(\tilde{\chi}_1^0)-100 \text{ GeV}$
	$B\bar{B}, B \rightarrow q\bar{q}\tilde{\chi}_1^0, \tilde{\chi}_1^0 \rightarrow q\bar{q}$	4-5 large- R jets Multiple	36.1 36.1	Large χ_{111} $m(\tilde{\chi}_1^0)=200 \text{ GeV}, \text{bino-like}$
	$\pi, \pi \rightarrow \tilde{\chi}_1^0, \tilde{\chi}_1^0 \rightarrow b\bar{s}$	Multiple	36.1	$m(\tilde{\chi}_1^0)=200 \text{ GeV}, \text{bino-like}$
	$\tilde{\chi}_1^0, \tilde{\chi}_1^0 \rightarrow b\bar{s}$	2 jets + 2 b	36.7	$BR(\tilde{\chi}_1^0 \rightarrow b\bar{s})=20\%$
$\tilde{\chi}_1^0, \tilde{\chi}_1^0 \rightarrow q\ell$	$\tilde{\chi}_1^0, \tilde{\chi}_1^0 \rightarrow b\bar{s}$	2 e, μ 1 μ DV	36.1	$BR(\tilde{\chi}_1^0 \rightarrow q\ell)=100\%, \cos\theta_l=1$
	$\tilde{\chi}_1^0, \tilde{\chi}_1^0 \rightarrow q\ell$	1 μ DV	136	$BR(\tilde{\chi}_1^0 \rightarrow q\ell)=100\%, \cos\theta_l=1$



*Only a selection of the available mass limits on new states or phenomena is shown. Many of the limits are based on simplified models, c.f. refs. for the assumptions made.

Search for Physics BSM

ATLAS Exotics Searches* - 95% CL Upper Exclusion Limits

Status: May 2019

ATLAS Preliminary

$$\int \mathcal{L} dt = (3.2 - 139) \text{ fb}^{-1}$$

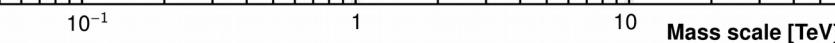
$$\sqrt{s} = 8, 13 \text{ TeV}$$

Model	ℓ, γ	Jets [†]	E_T^{miss}	$\int \mathcal{L} dt [\text{fb}^{-1}]$	Limit	Reference
Extra dimensions	ADD $G_{KK} + g/q$	0 e, μ	1 – 4 j	Yes	36.1	M_D 7.7 TeV
	ADD non-resonant $\gamma\gamma$	2 γ	–	–	36.7	M_S 8.6 TeV
	ADD QBH	–	2 j	–	37.0	M_{th} 8.9 TeV
	ADD BH high $\sum p_T$	≥ 1 e, μ	≥ 2 j	–	3.2	M_{th} 8.2 TeV
	ADD BH multijet	–	≥ 3 j	–	3.6	M_{th} 9.55 TeV
	RS1 $G_{KK} \rightarrow \gamma\gamma$	2 γ	–	–	36.7	G_{KK} mass 4.1 TeV
	Bulk RS $G_{KK} \rightarrow WW/ZZ$	multi-channel		36.1	G_{KK} mass 2.3 TeV	
	Bulk RS $G_{KK} \rightarrow WW \rightarrow qqqq$	0 e, μ	2 J	–	139	G_{KK} mass 1.6 TeV
	Bulk RS $g_{KK} \rightarrow tt$	1 e, μ	≥ 1 b, ≥ 1 J/2j	Yes	36.1	g_{KK} mass 3.8 TeV
	2UED / RPP	1 e, μ	≥ 2 b, ≥ 3 j	Yes	36.1	KK mass 1.8 TeV
Gauge bosons	SSM $Z' \rightarrow \ell\ell$	2 e, μ	–	–	139	Z' mass 5.1 TeV
	SSM $Z' \rightarrow \tau\tau$	2 τ	–	–	36.1	Z' mass 2.42 TeV
	Leptophobic $Z' \rightarrow bb$	–	2 b	–	36.1	Z' mass 2.1 TeV
	Leptophobic $Z' \rightarrow tt$	1 e, μ	≥ 1 b, ≥ 1 J/2j	Yes	36.1	Z' mass 3.0 TeV
	SSM $W' \rightarrow \ell\nu$	1 e, μ	–	Yes	139	W' mass 6.0 TeV
	SSM $W' \rightarrow \tau\nu$	1 τ	–	Yes	36.1	W' mass 3.7 TeV
	HVT $V' \rightarrow WZ \rightarrow qqqq$ model B	0 e, μ	2 J	–	139	V' mass 3.6 TeV
	HVT $V' \rightarrow WH/ZH$ model B	multi-channel		36.1	V' mass 2.93 TeV	
	LRSM $W_R \rightarrow tb$	multi-channel		36.1	W_R mass 3.25 TeV	
	LRSM $W_R \rightarrow \mu N_R$	2 μ	1 J	–	80	W_R mass 5.0 TeV
CI	Cl $qqqq$	–	2 j	–	37.0	Λ 21.8 TeV
	Cl $\ell\ell qq$	2 e, μ	–	–	36.1	Λ 40.0 TeV
	Cl $t\bar{t} t\bar{t}$	≥ 1 e, μ	≥ 1 b, ≥ 1 j	Yes	36.1	Λ 2.57 TeV
DM	Axial-vector mediator (Dirac DM)	0 e, μ	1 – 4 j	Yes	36.1	m_{med} 1.55 TeV
	Colored scalar mediator (Dirac DM)	0 e, μ	1 – 4 j	Yes	36.1	m_{med} 1.67 TeV
	$VV\chi\chi$ EFT (Dirac DM)	0 e, μ	1 J, ≤ 1 j	Yes	3.2	M_\ast 700 GeV
	Scalar reson. $\phi \rightarrow t\bar{t}$ (Dirac DM)	0-1 e, μ	1 b, 0-1 J	Yes	36.1	m_ϕ 3.4 TeV
LQ	Scalar LQ 1 st gen	1,2 e	≥ 2 j	Yes	36.1	LQ mass 1.4 TeV
	Scalar LQ 2 nd gen	1,2 μ	≥ 2 j	Yes	36.1	LQ mass 1.56 TeV
	Scalar LQ 3 rd gen	2 τ	2 b	–	36.1	LQ_3^u mass 1.03 TeV
	Scalar LQ 3 rd gen	0-1 e, μ	2 b	Yes	36.1	LQ_3^d mass 970 GeV
Heavy quarks	VLQ $TT \rightarrow Ht/Zt/Wb + X$	multi-channel		36.1	T mass 1.37 TeV	
	VLQ $BB \rightarrow Wh/Zb + X$	multi-channel		36.1	B mass 1.34 TeV	
	VLQ $T_{5/3} T_{5/3} T_{5/3} \rightarrow Wt + X$	$2(SS) \geq 3$ e, μ	≥ 1 b, ≥ 1 j	Yes	36.1	$T_{5/3}$ mass 1.64 TeV
	VLQ $Y \rightarrow Wh + X$	1 e, μ	≥ 1 b, ≥ 1 j	Yes	36.1	Y mass 1.85 TeV
	VLQ $B \rightarrow Hb + X$	0 e, μ , 2 γ	≥ 1 b, ≥ 1 j	Yes	79.8	B mass 1.21 TeV
	VLQ $QQ \rightarrow WqWq$	1 e, μ	≥ 4 j	Yes	20.3	Q mass 690 GeV
Excited fermions	Excited quark $q^+ \rightarrow qg$	–	2 j	–	139	q^+ mass 6.7 TeV
	Excited quark $q^+ \rightarrow q\gamma$	1 γ	1 j	–	36.7	q^+ mass 5.3 TeV
	Excited quark $b^+ \rightarrow bg$	–	1 b, 1 j	–	36.1	b^+ mass 2.6 TeV
	Excited lepton ℓ^+	3 e, μ	–	–	20.3	ℓ^+ mass 3.0 TeV
	Excited lepton ν^+	3 e, μ , τ	–	–	20.3	ν^+ mass 1.6 TeV
Other	Type III Seesaw	1 e, μ	≥ 2 j	Yes	79.8	N^0 mass 560 GeV
	LRSM Majorana ν	2 μ	2 j	–	36.1	N_R mass 3.2 TeV
	Higgs triplet $H^{\pm\pm} \rightarrow \ell\ell$	2,3,4 e, μ (SS)	–	–	36.1	$H^{\pm\pm}$ mass 870 GeV
	Higgs triplet $H^{\pm\pm} \rightarrow \tau\tau$	3 e, μ , τ	–	–	20.3	$H^{\pm\pm}$ mass 400 GeV
	Multi-charged particles	–	–	–	36.1	multi-charged particle mass 1.22 TeV
	Magnetic monopoles	–	–	–	34.4	monopole mass 2.37 TeV

$\sqrt{s} = 8 \text{ TeV}$

$\sqrt{s} = 13 \text{ TeV}$
partial data

$\sqrt{s} = 13 \text{ TeV}$
full data



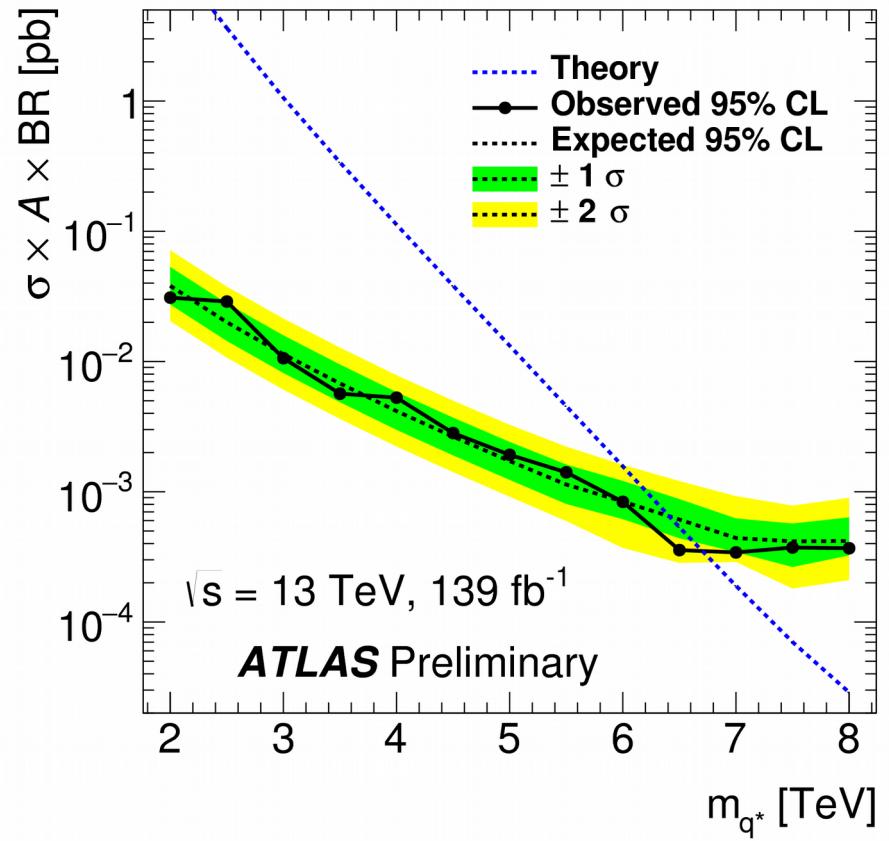
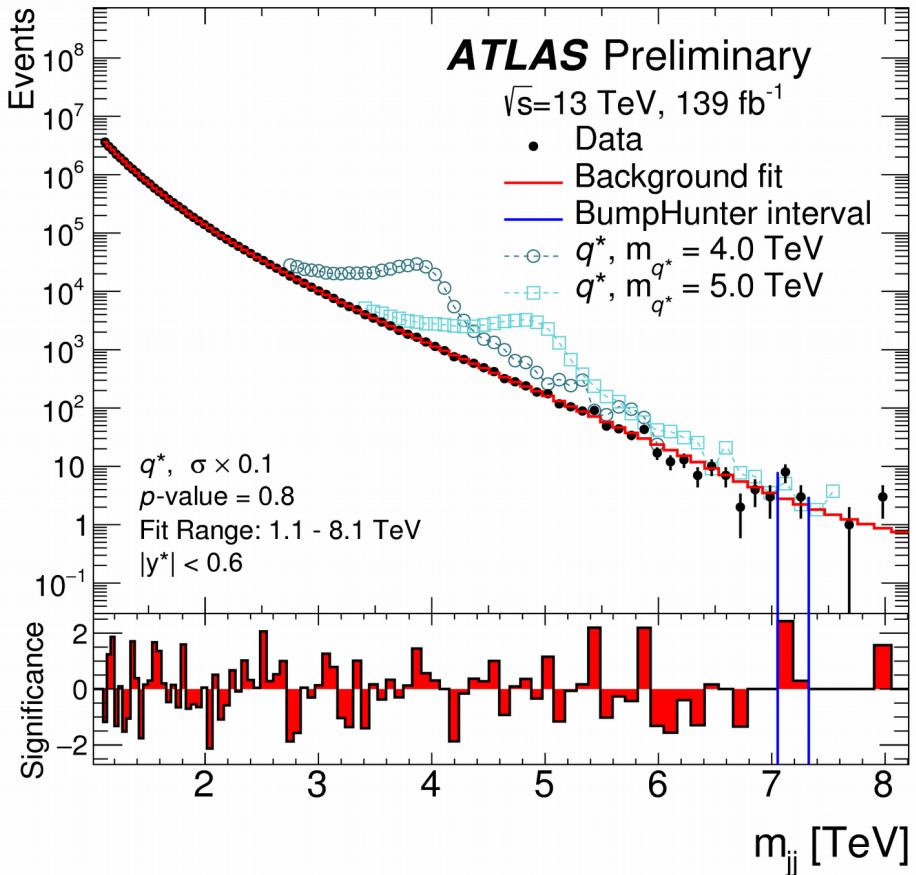
*Only a selection of the available mass limits on new states or phenomena is shown.

[†]Small-radius (large-radius) jets are denoted by the letter j (R).

Search for Physics BSM

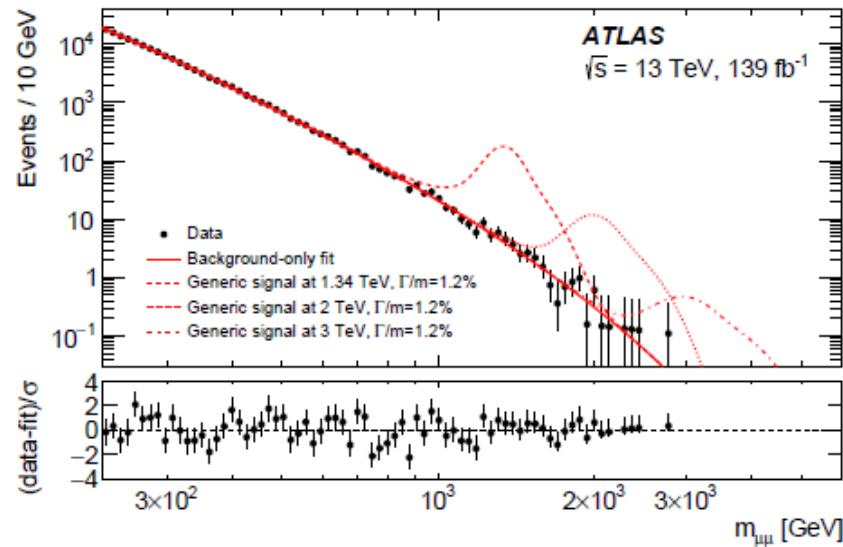
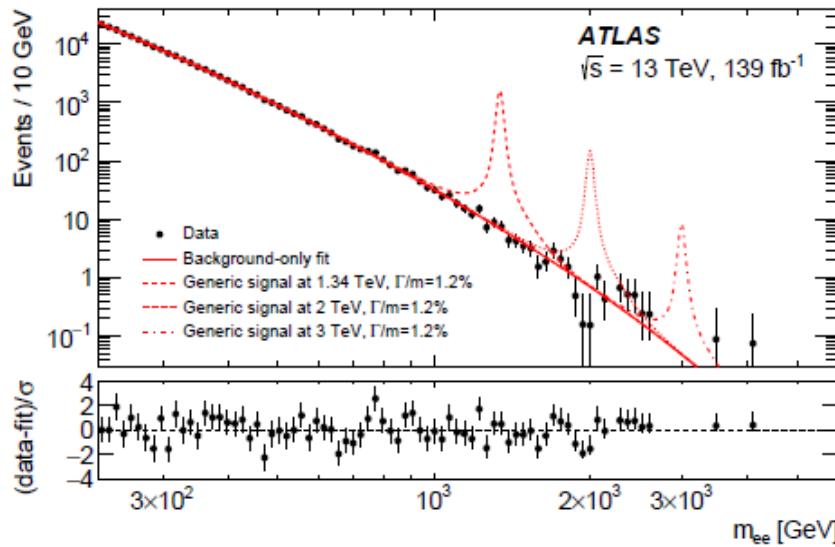
Full Run-2

parton-level generators assuming spin-½ excited quarks with the same coupling constants as SM quarks
the intrinsic width of the q^* signals is comparable to the detector resolution



Search for Physics BSM

Full Run-2

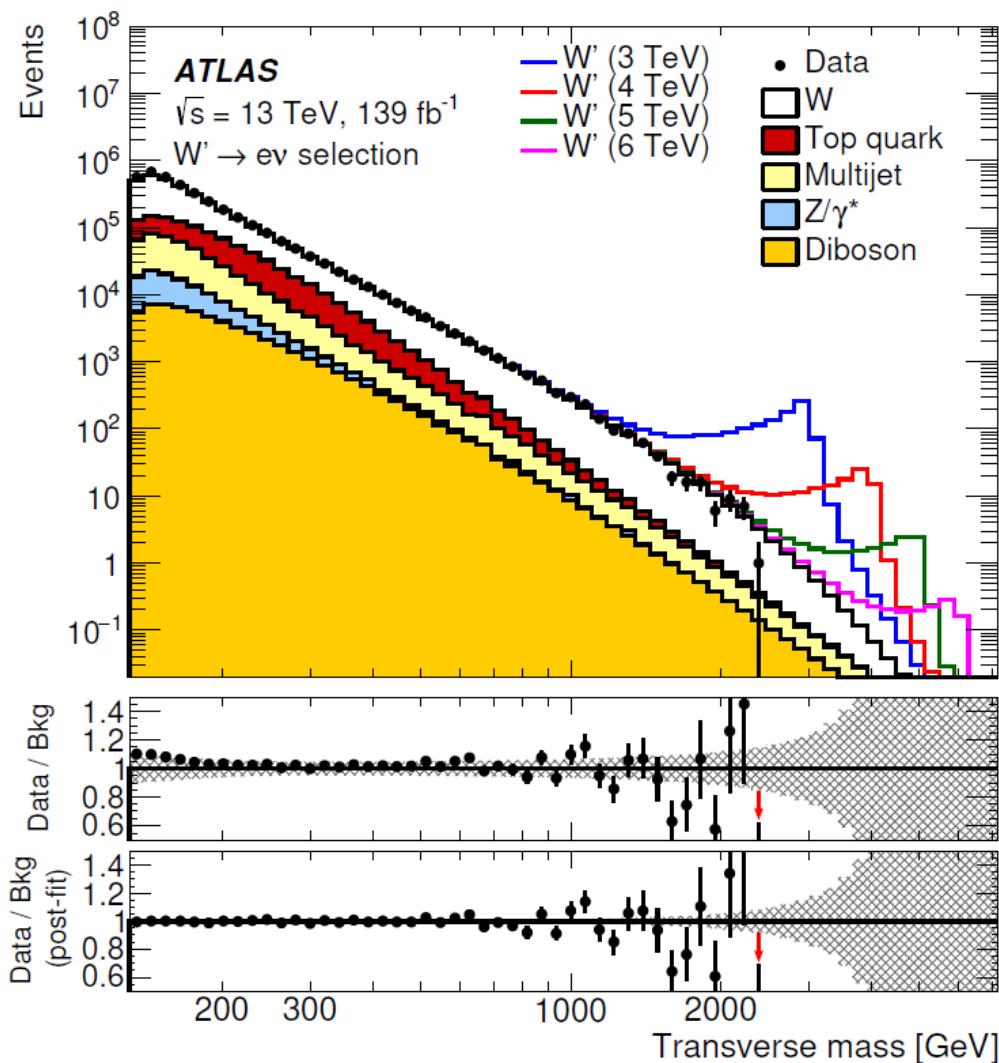


Observed and expected 95% CL lower limits on $m_{Z'}$ for three Z' gauge boson models

Model	Lower limits on $m_{Z'} [\text{TeV}]$					
	ee		$\mu\mu$		$\ell\ell$	
	obs	exp	obs	exp	obs	exp
Z'_ψ	4.3	4.3	4.0	3.8	4.5	4.5
Z'_χ	4.6	4.6	4.2	4.1	4.8	4.7
Z'_{SSM}	4.9	4.9	4.5	4.4	5.1	5.0

Search for Physics BSM

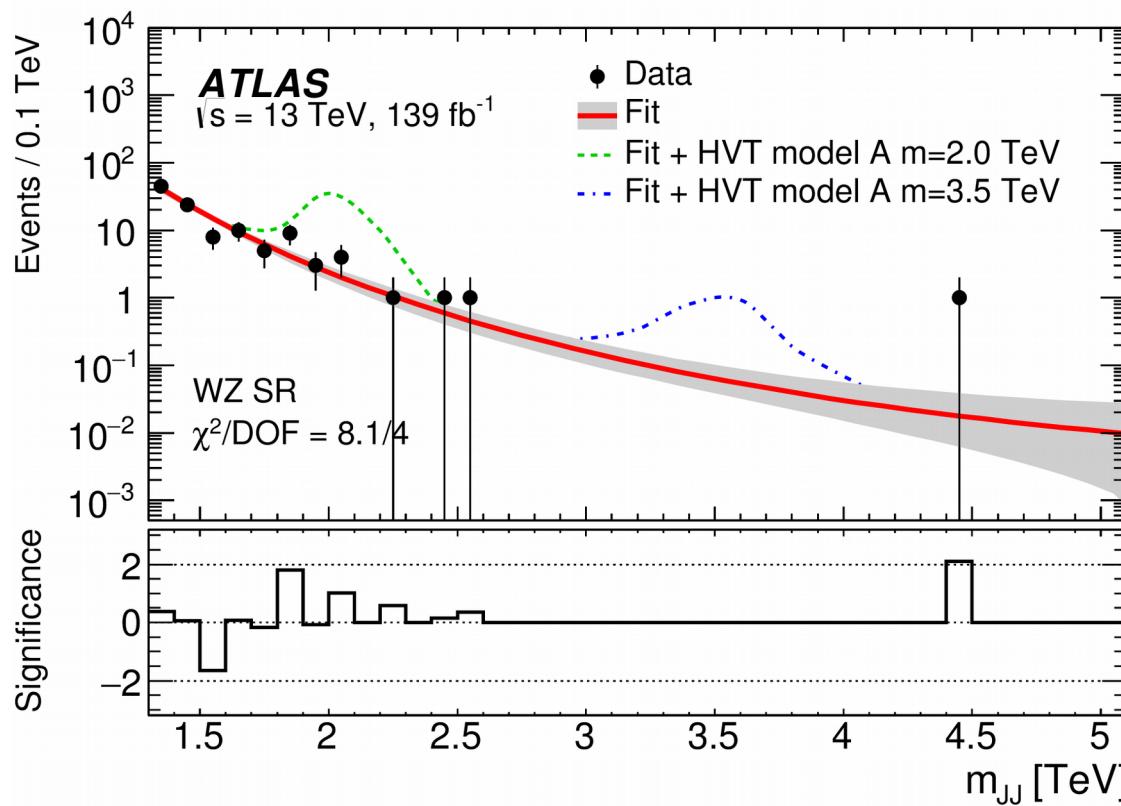
Full Run-2



Observed and expected 95% CL
lower limits on the W' mass

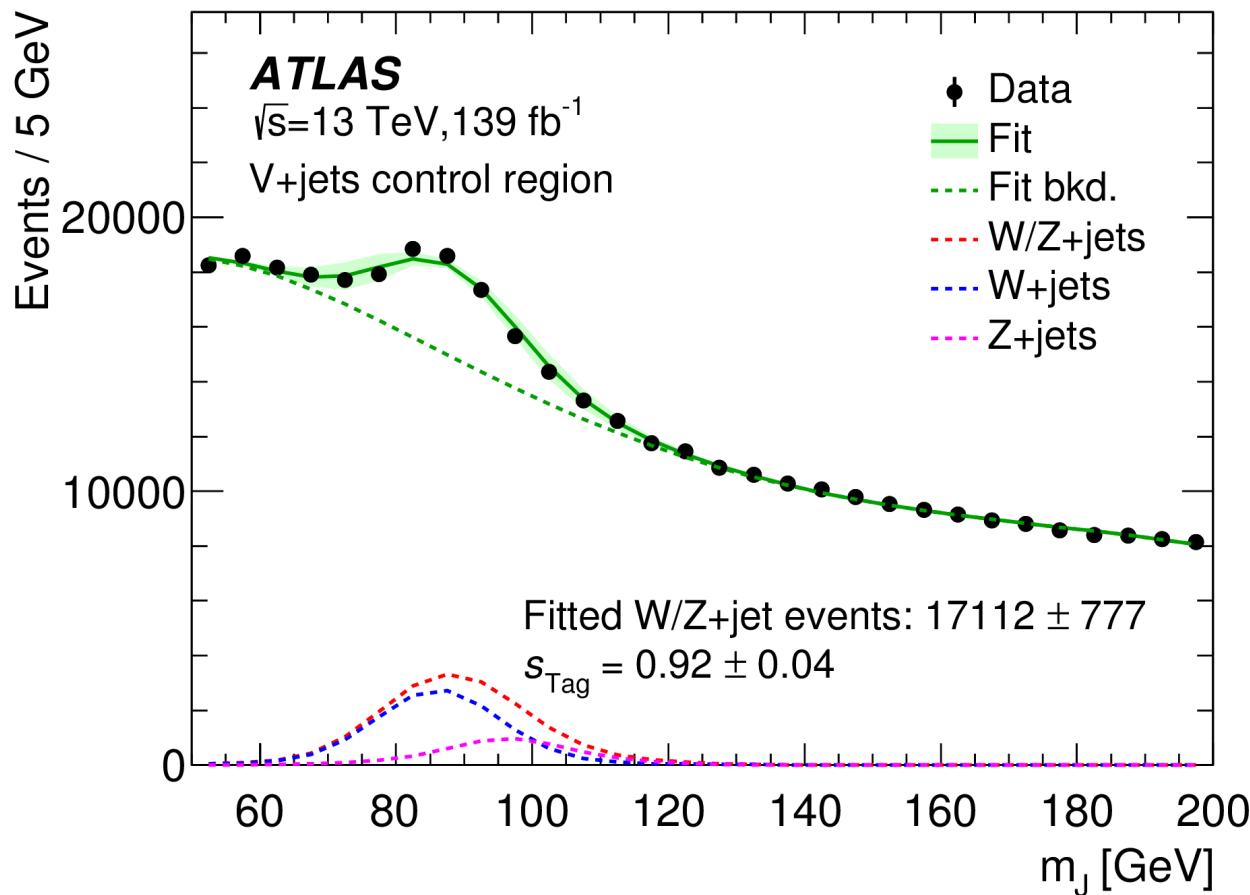
Search for Physics BSM

Full Run-2



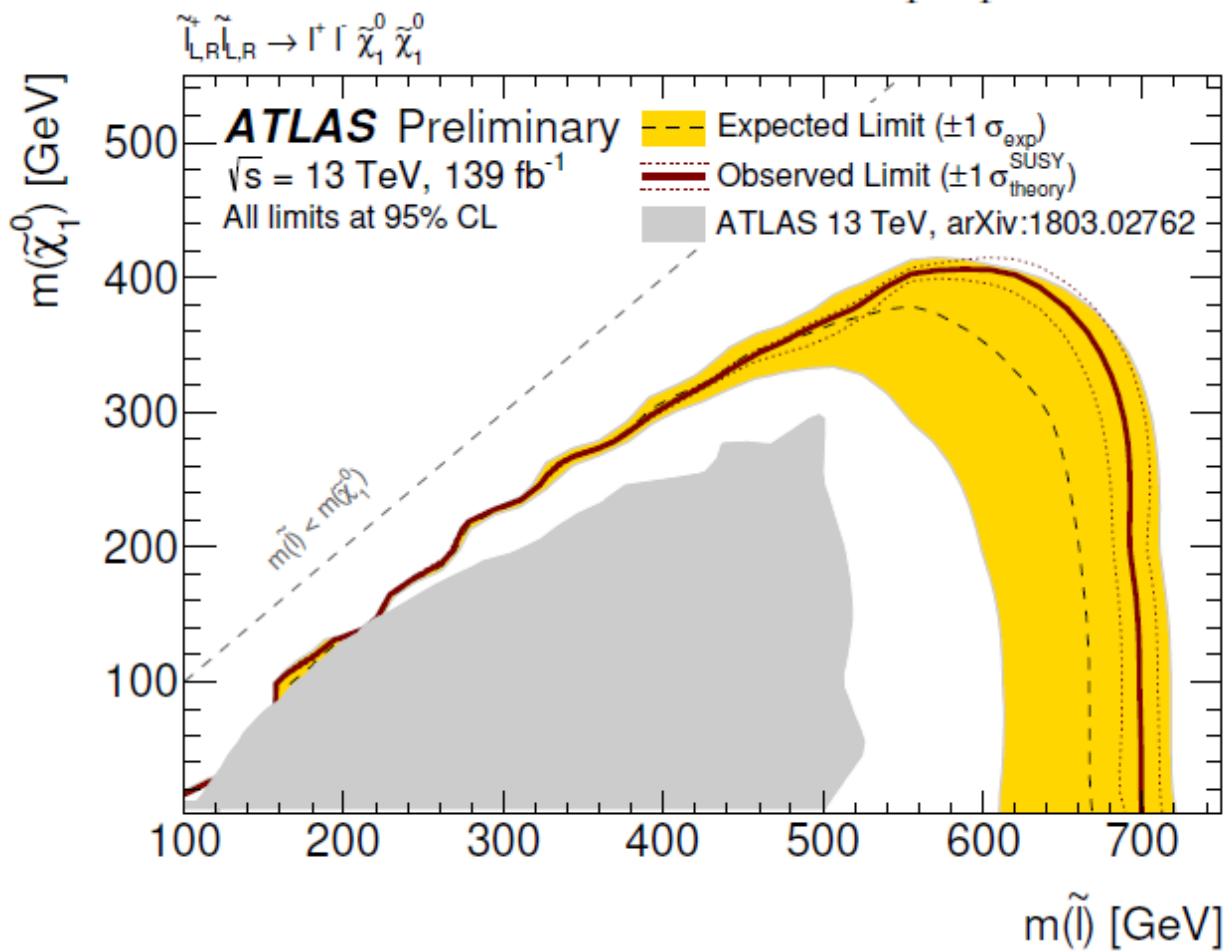
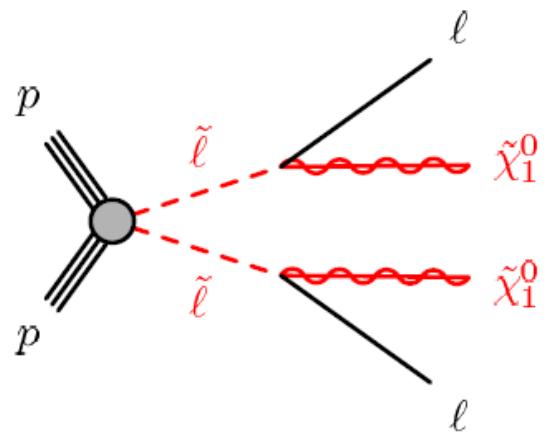
Search for Physics BSM

Full Run-2



**Jet mass distribution for data in
the region enhanced in V+jets
events after boson tagging**

Search for Physics BSM



Search for Physics BSM

	MB-SSd	MB-GGd	MB-C
N_j	≥ 2	≥ 4	≥ 2
$p_T(j_1)$ [GeV]	> 200	> 200	> 600
$p_T(j_{i=2,\dots,N_{j_{\min}}})$ [GeV]	> 100	> 100	> 50
$ \eta(j_{i=1,\dots,N_{j_{\min}}}) $	< 2.0	< 2.0	< 2.8
$\Delta\phi(j_{1,2,(3)}, p_T^{\text{miss}})_{\text{min}}$	> 0.8	> 0.4	> 0.4
$\Delta\phi(j_{i>3}, p_T^{\text{miss}})_{\text{min}}$	> 0.4	> 0.2	> 0.2
Aplanarity	-	> 0.04	-
$E_T^{\text{miss}}/\sqrt{H_T}$ [GeV $^{1/2}$]	> 10	> 10	> 10
m_{eff} [GeV]	> 1000	> 1000	> 1600

Table 3: Summary of preselection criteria used for the multi-bin search.

Search for Physics BSM

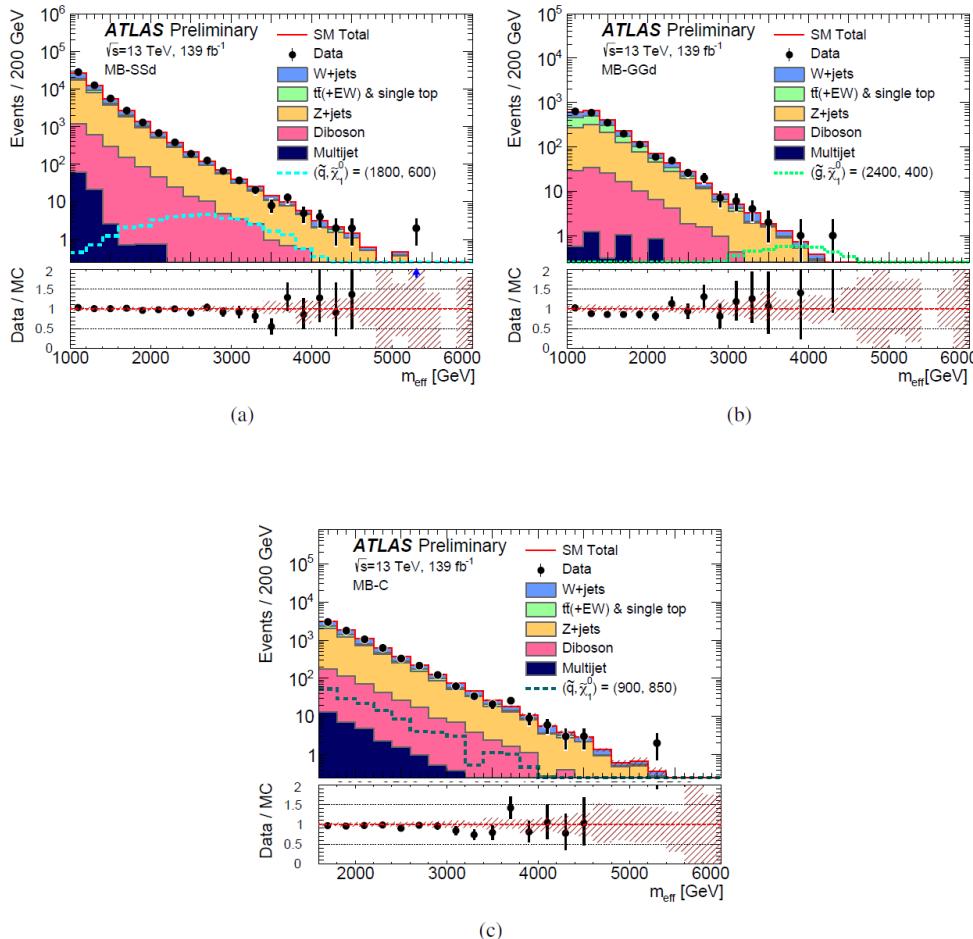


Figure 8: Observed m_{eff} distributions for the (a) MB-SSd, (b) MB-GGd and (c) MB-C regions obtained after applying the selection criteria from Table 3, before the final binning selections on this quantity. The histograms show the MC background predictions prior to the fits described in the text, normalized to the cross-section times integrated luminosity. The hatched (red) error bands indicate the combined experimental and MC statistical uncertainties. Expected distributions for benchmark signal model points, normalized using NLO+NLL cross-section (Section 3) times integrated luminosity, are also shown for comparison (masses in GeV).

Search for Physics BSM

Table 1: Event selection in the two search regions. A dash means that the variable is not used for selection.

Event category	SR1	SR2
	Additional forward jet	Additional central jet
Muons	OS, $p_T > 25 \text{ GeV}, \eta < 2.1$	
$m_{\mu\mu}$		$m_{\mu\mu} > 12 \text{ GeV}$
b-tagged jet		$p_T > 30 \text{ GeV}, \eta \leq 2.4$
Additional jet	$p_T > 30 \text{ GeV}, 2.4 < \eta < 4.7$	$p_T > 30 \text{ GeV}, \eta \leq 2.4$
Jet veto	No other jets $p_T > 30 \text{ GeV}, \eta \leq 2.4$	No jets $p_T > 30 \text{ GeV}, 2.4 < \eta < 4.7$
p_T^{miss}	—	$< 40 \text{ GeV}$
$\Delta\phi(\mu\mu, \text{jj})$	—	$> 2.5 \text{ rad}$

Table 3: The local significances, the measured fiducial signal cross sections with ± 1 s.d. uncertainties, and the upper limits at 95% CL, together with the values of N_S for the two SRs and two collision energies. The reconstruction efficiencies and the integrated luminosities are also listed.

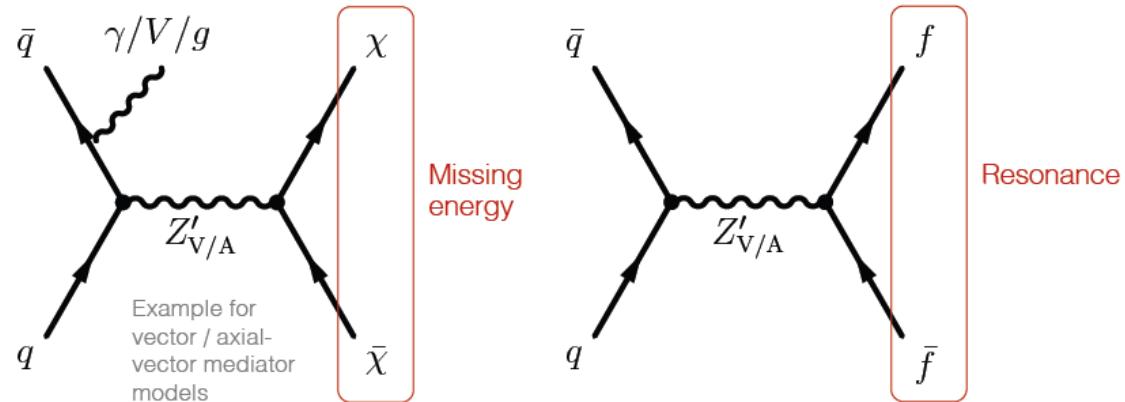
Event category	$\sqrt{s} (\text{TeV})$		8		13	
	SR1	SR2	SR1	SR2	SR1	SR2
Local significance (s.d.)	4.2	2.9	2.0	1.4 deficit		
N_S	22.0 ± 7.6	22.8 ± 9.5	14.5 ± 9.3	-14.9 ± 10.1		
N_S observed upper limit at 95% CL	40.4	44.7	36.9	32.2		
N_S expected upper limit at 95% CL	18.3	27.6	27.6	35.6		
ϵ^{reco}		0.27 ± 0.01		0.28 ± 0.01		
Integrated luminosity, $\mathcal{L} (\text{fb}^{-1})$		19.7 ± 0.5		35.9 ± 0.9		
$\sigma_{\text{fid}} (\text{fb})$	4.1 ± 1.4	4.2 ± 1.7	1.4 ± 0.9	-1.5 ± 1.0		
Observed upper limit at 95% CL (fb)	7.6	8.4	3.7	3.2		
Expected upper limit at 95% CL (fb)	3.4	5.2	2.7	3.5		

	8 TeV		13 TeV	
	SR1	SR2	SR1	SR2
Local significance (28 GeV)	0.5	0.5	0.7	0.2
Max. significance	0.9 (29.5 GeV)	1.1 (29.5 GeV)	0.8 (27.5 GeV)	2.1 (26 GeV)

Table 2: Local significance for a dimuon excess at 28 GeV in each signal region and the maximum observed significance in the dimuon invariant mass probed from 26 to 30 GeV in steps of 0.5 GeV are quoted.

Search for Physics BSM dark matter

If produced at the LHC,
DM interactions will be
mediated by particles
that can also be directly
searched for
— complementarity

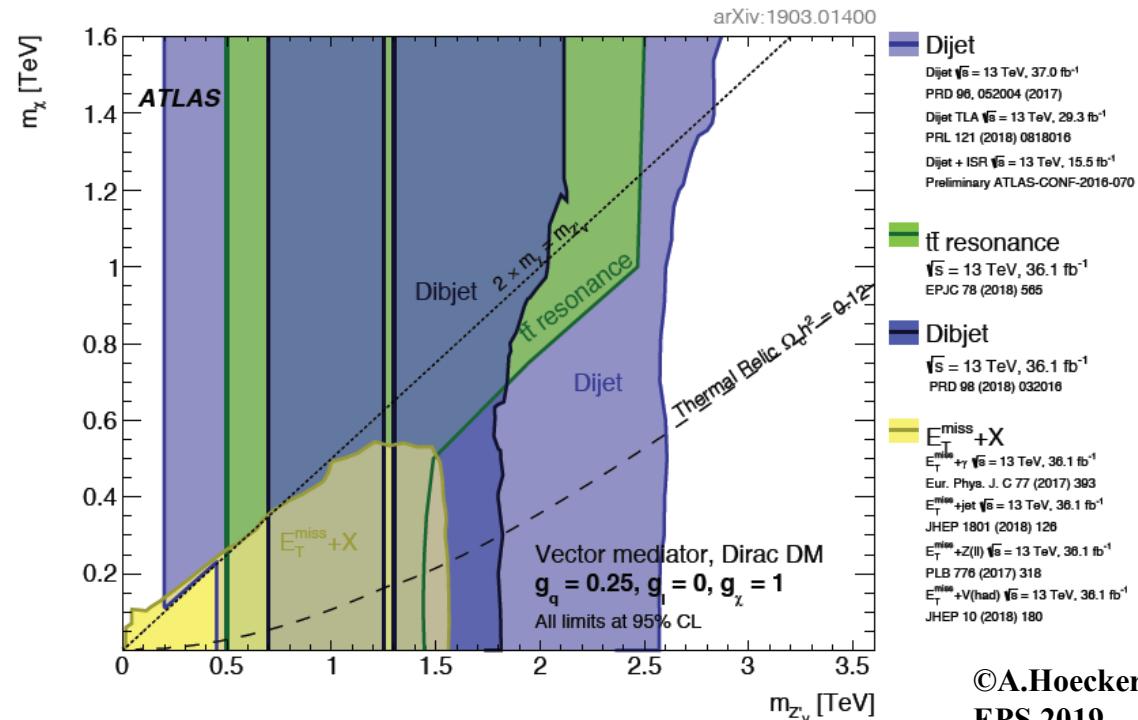


ATLAS released combination
of $E_{T,\text{miss}}$ based DM searches
involving $E_{T,\text{miss}} + X$, $X = \text{jet}, \gamma,$
 $W, Z, H, b(b), t(t)$ using large
number of models

arXiv:1903.01400, up to 37 fb^{-1}

If light enough, Higgs boson can
decay to DM ($H \rightarrow \text{invisible}$)
ATLAS combination:
 $\text{BR}(H \rightarrow \text{invisible}) < 0.26$
(0.17 expected)

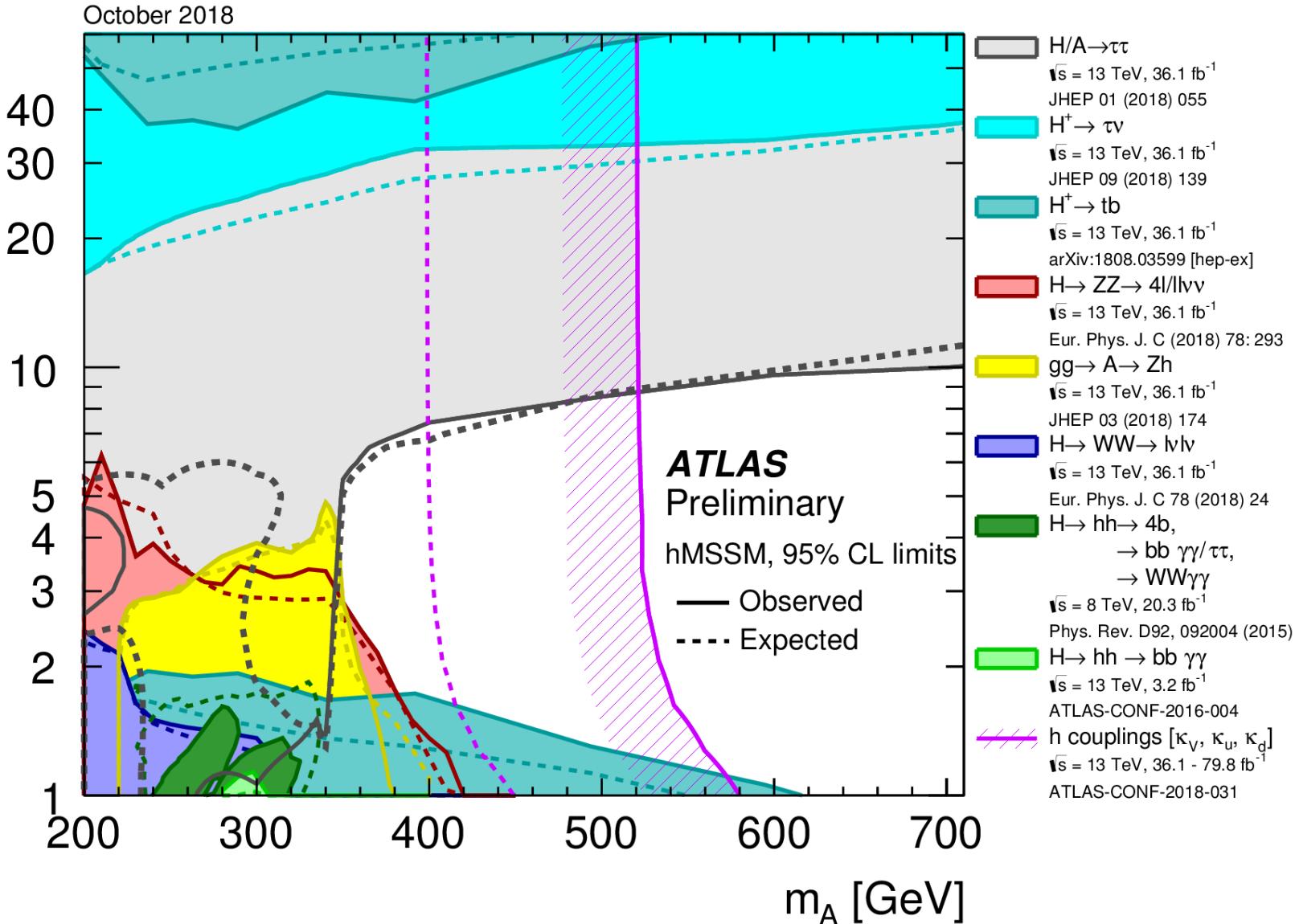
arXiv:1904.05105, 36 fb^{-1}



- ♪ *Results from (run 1 and) run 2*
- * *detector*
- * *SM*
- * *BSM*
- * ***(B-E)H***
- * *Vector-boson*

1 Additional BEH bosons

$\tan \beta$



Corfou 2-9-2019

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1 Additional BEH bosons

$$\kappa_V = \frac{s_d(m_A, \tan \beta) + \tan \beta s_u(m_A, \tan \beta)}{\sqrt{1 + \tan^2 \beta}}$$

$$\kappa_u = s_u(m_A, \tan \beta) \frac{\sqrt{1 + \tan^2 \beta}}{\tan \beta}$$

$$\kappa_d = s_d(m_A, \tan \beta) \sqrt{1 + \tan^2 \beta} \quad ,$$

where the functions s_u and s_d are given by:

$$s_u = \frac{1}{\sqrt{1 + \frac{(m_A^2 + m_Z^2)^2 \tan^2 \beta}{(m_Z^2 + m_A^2 \tan^2 \beta - m_h^2 (1 + \tan^2 \beta))^2}}}$$

$$s_d = \frac{(m_A^2 + m_Z^2) \tan \beta}{m_Z^2 + m_A^2 \tan^2 \beta - m_h^2 (1 + \tan^2 \beta)} s_u$$

1 Additional BEH bosons

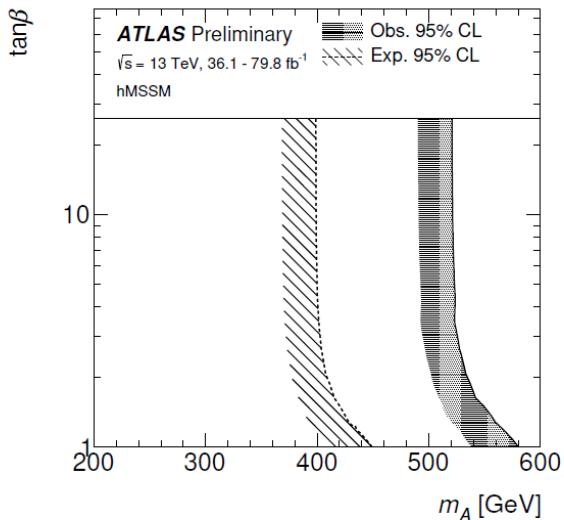


Figure 14: Regions of the $[m_A, \tan\beta]$ plane in the hMSSM excluded by fits to the measured rates of Higgs boson production and decays. Likelihood contours at 95% CL, defined in the asymptotic approximation by $-2 \log \Lambda = 5.99$, are drawn for both the data and the expectation of the SM Higgs sector. The regions to the left of the solid contour are excluded. The decoupling limit, in which all Higgs boson couplings tend to their SM value, corresponds to $m_A \rightarrow \infty$. The hMSSM is a good approximation of the MSSM only for moderate values of $\tan\beta$. For $\tan\beta \gtrsim 10$ the scenario is approximate due to missing supersymmetry corrections in the Higgs boson coupling to b-quarks, and for $\tan\beta$ of $O(1)$ the precision of the approximation depends on m_A [34].

The Higgs boson couplings to vector bosons, up-type fermions and down-type fermions relative to the corresponding SM predictions are expressed as functions of the ratio of the vacuum expectation values of the Higgs doublets, $\tan\beta$, and the masses of the CP-odd scalar (m_A), the Z boson, and of h .

1 Additional BEH bosons

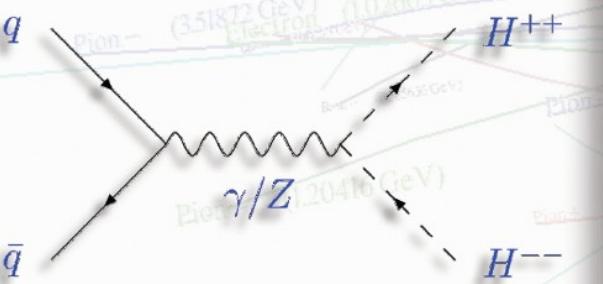
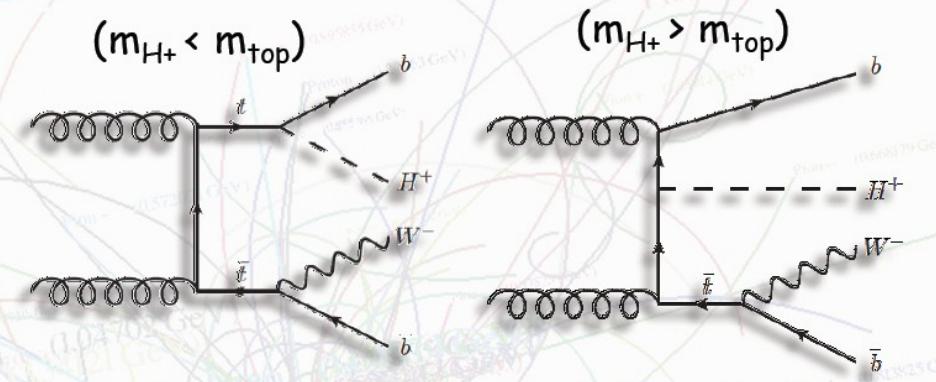
© A.Kaczmarska Higgs Hunting 2019

Charged Higgs searches

H^\pm predicted by 2HDM, Higgs triplets,...

In Type II 2HDM:

- Main production in association with a top quark
- At high mass $H^\pm \rightarrow tb$ is the dominant decay mode
- BR($H^\pm \rightarrow \tau\nu$) significant for a large range of masses for high $\tan\beta$



$H^{\pm\pm}$ predicted by Left-Right Symmetric Models (LRSM), Higgs Triplet Model (HTM), Zee-Babu and Georgi-Machacek models

In LRSM and HTM:

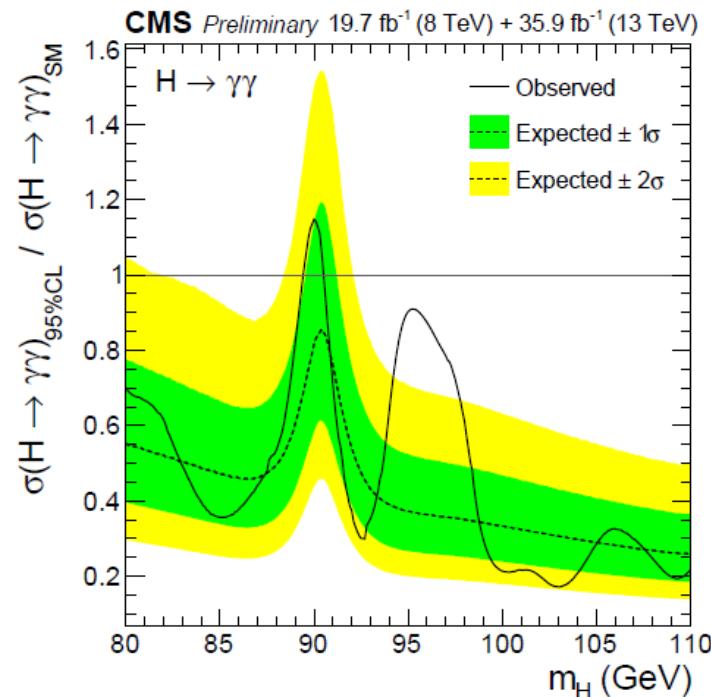
- Dominant production at the LHC: DY pair production
- Decays: $H^{\pm\pm} \rightarrow l^\pm l^\pm$ or $H^{\pm\pm} \rightarrow W^\pm W^\pm$
 - $BR \sim f(m_{H^{\pm\pm}}, vev \text{ of Higgs triplet})$
 - Low $m_{H^{\pm\pm}}$ and low vev: $H^{\pm\pm} \rightarrow l^\pm l^\pm$ dominates

- | | |
|------------------------------------------|------------------------------|
| • $H^\pm \rightarrow \tau\nu$ | JHEP 09 (2018) 139 |
| • $H^\pm \rightarrow tb$ | JHEP 11 (2018) 085 |
| • $H^{\pm\pm} H^{\mp\mp} \rightarrow 4W$ | Eur. Phys. J. C (2019) 79 |
| • $H^{\pm\pm} H^{\mp\mp} \rightarrow 4l$ | Eur. Phys. J. C78 (2018) 199 |

- | | |
|-------------------------------|------------------------------------------|
| • $H^\pm \rightarrow W^\pm Z$ | Phys. Lett. B 787 (2019) 68 |
| • $H^\pm \rightarrow cs$ | Eur. Phys. J. C, 73 6 (2013) 2465, Run1+ |

1 Additional BEH bosons $\gamma\gamma$ excess at 95 GeV

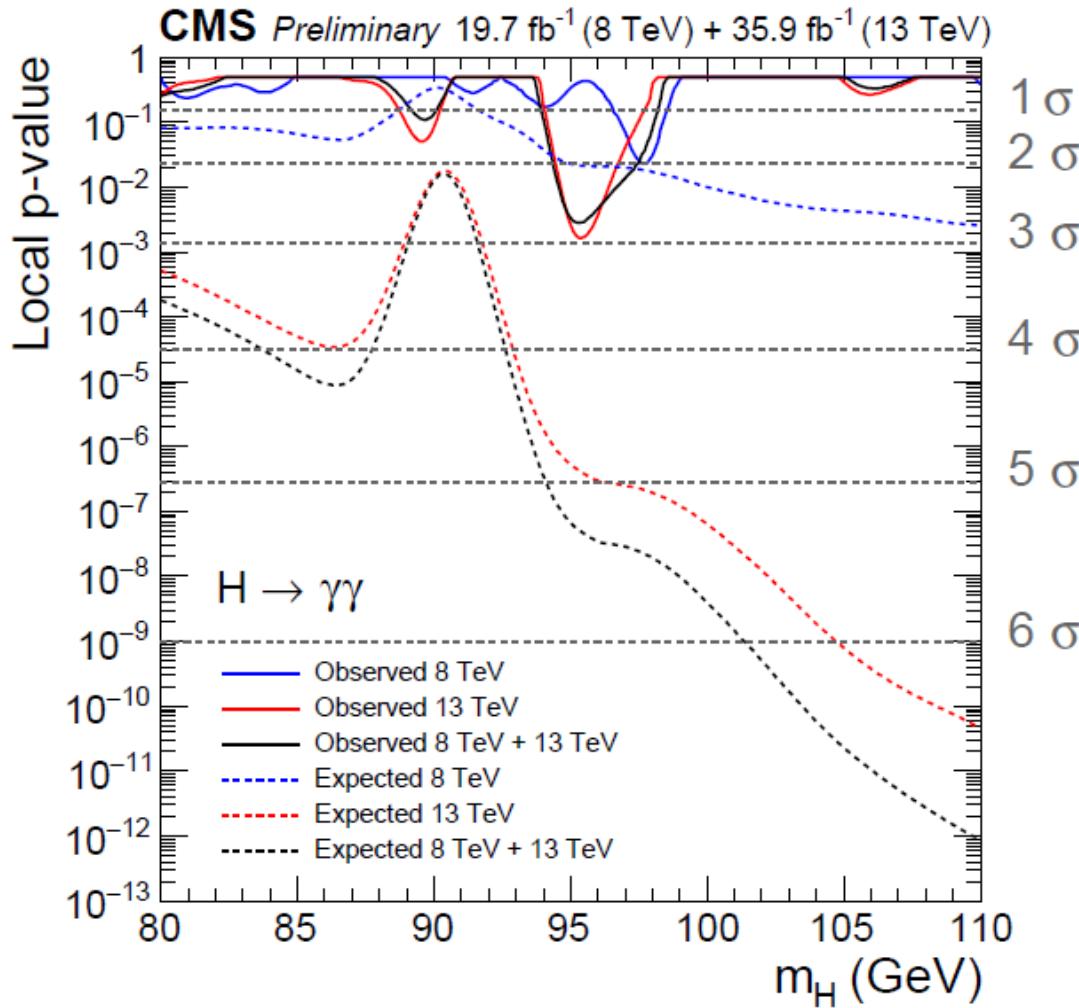
CMS PAS HIG-17-013



these yield an excess with approximately 2.8σ local (1.3σ global) significance for the same hypothesis mass as for the 13 TeV dataset alone, mass of 95.3 GeV.

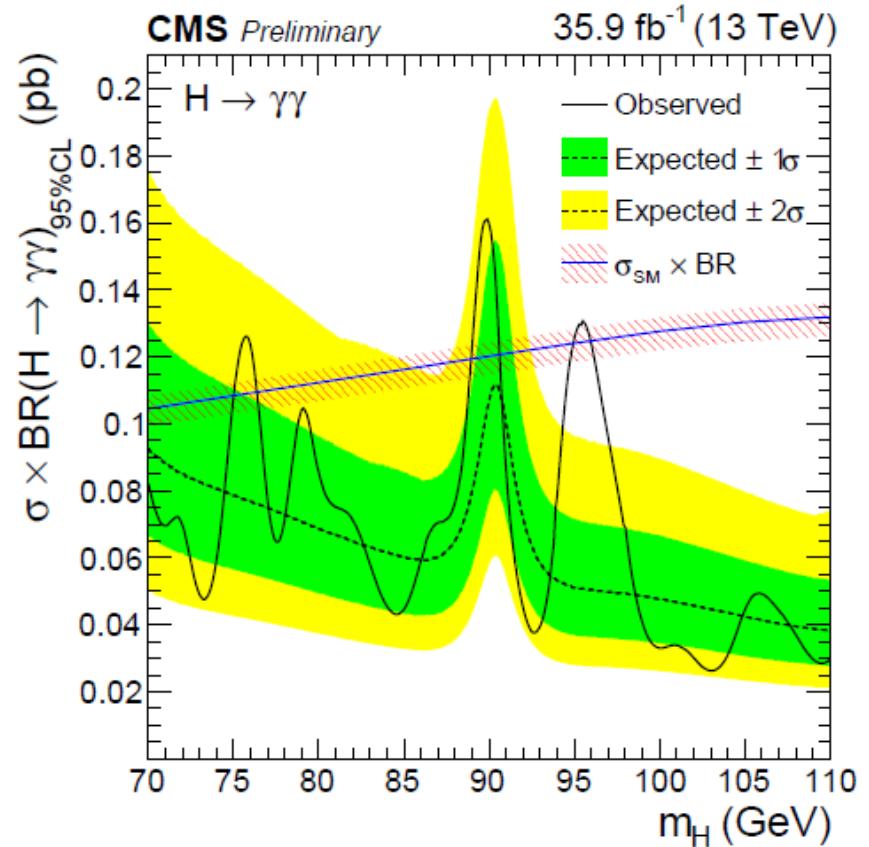
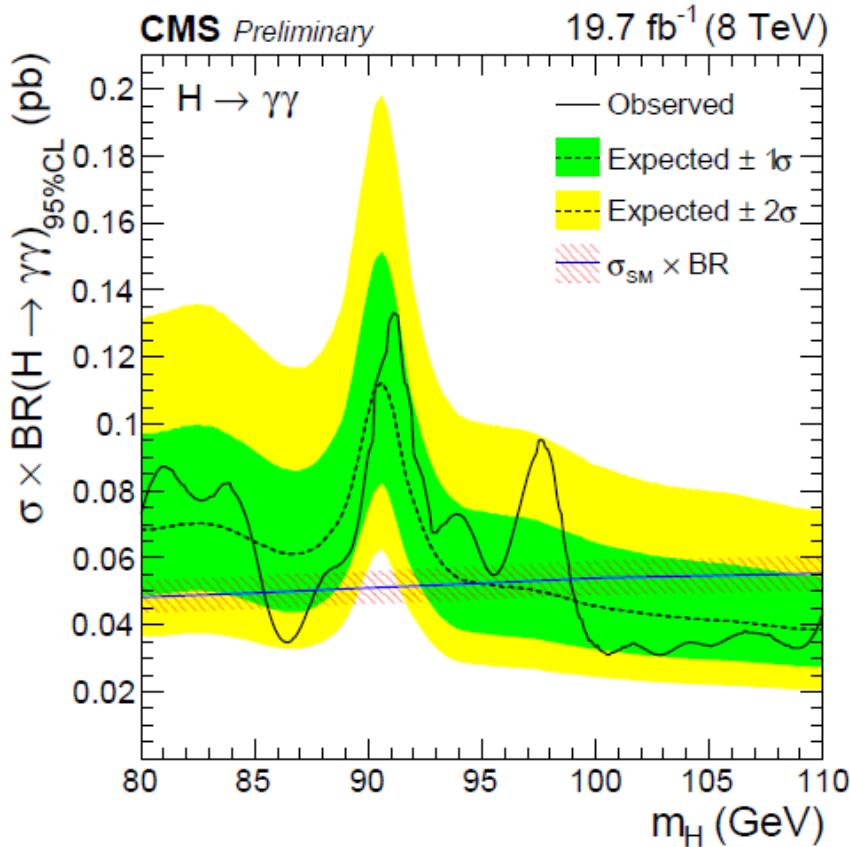
1 Additional BEH bosons $\gamma\gamma$ excess at 95 GeV

CMS PAS HIG-17-013



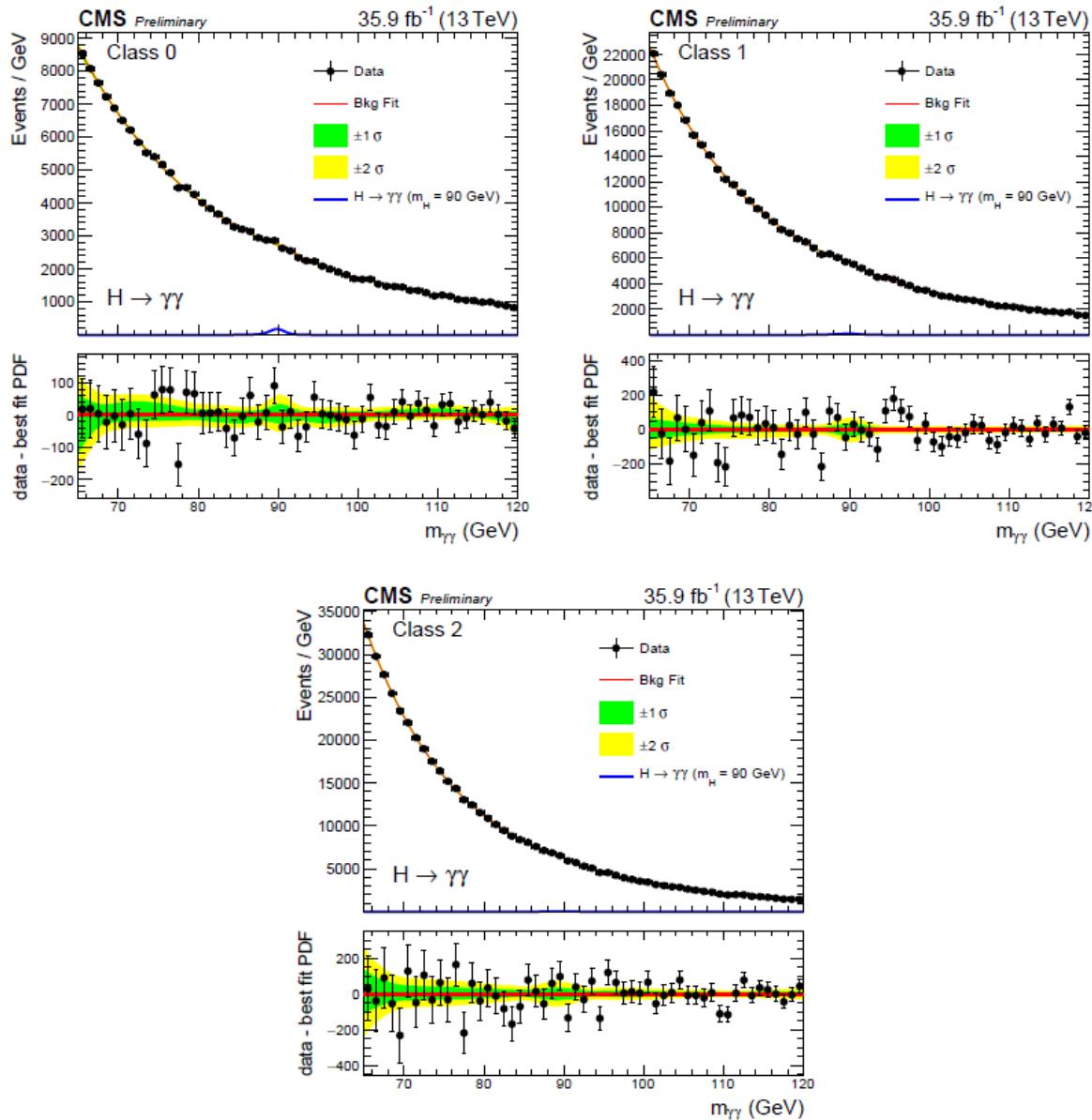
1 Additional BEH bosons $\gamma\gamma$ excess at 95 GeV

CMS PAS HIG-17-013



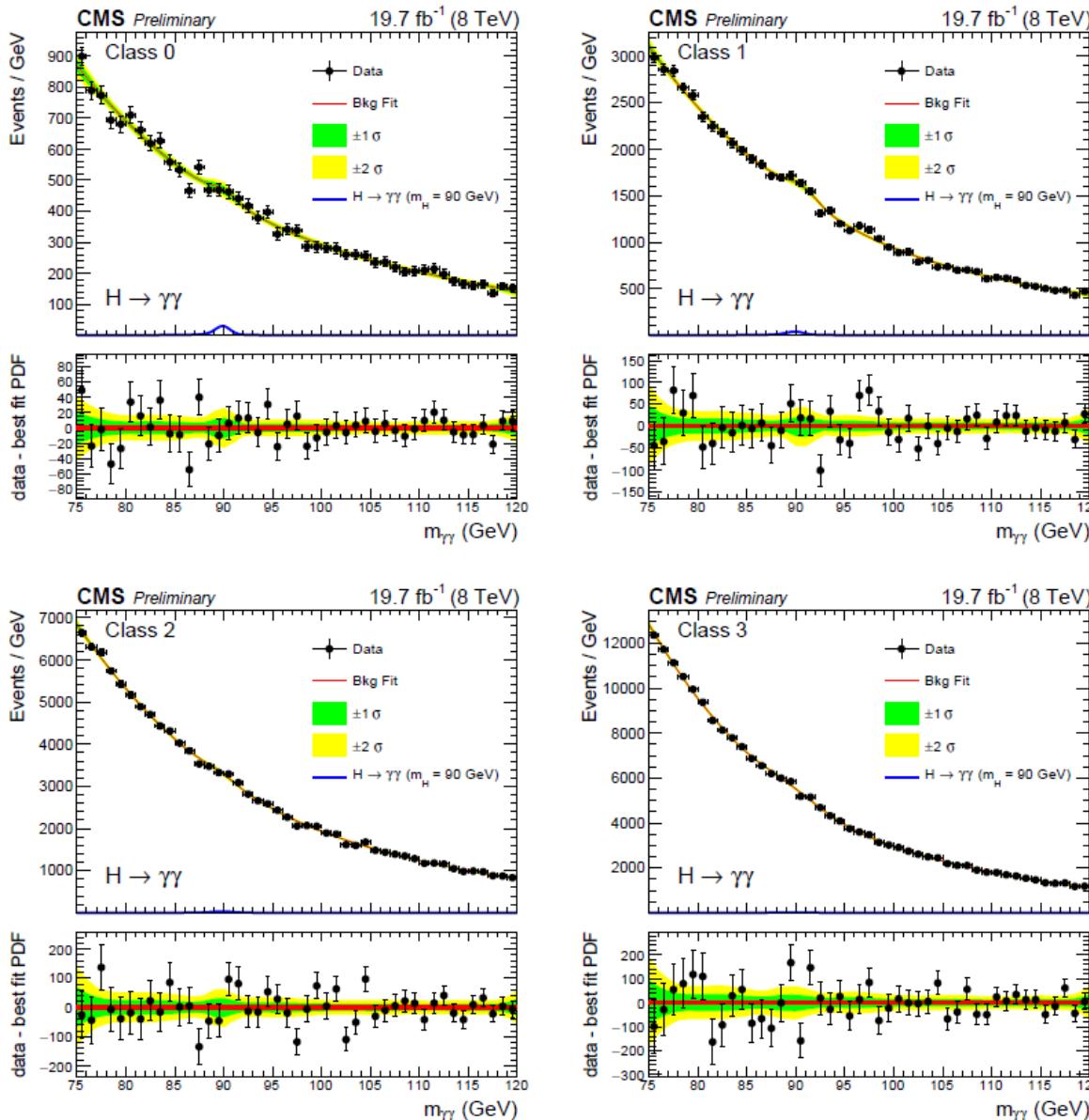
1 Additional BEH bosons $\gamma\gamma$ excess at 95 GeV

CMS PAS HIG-17-013



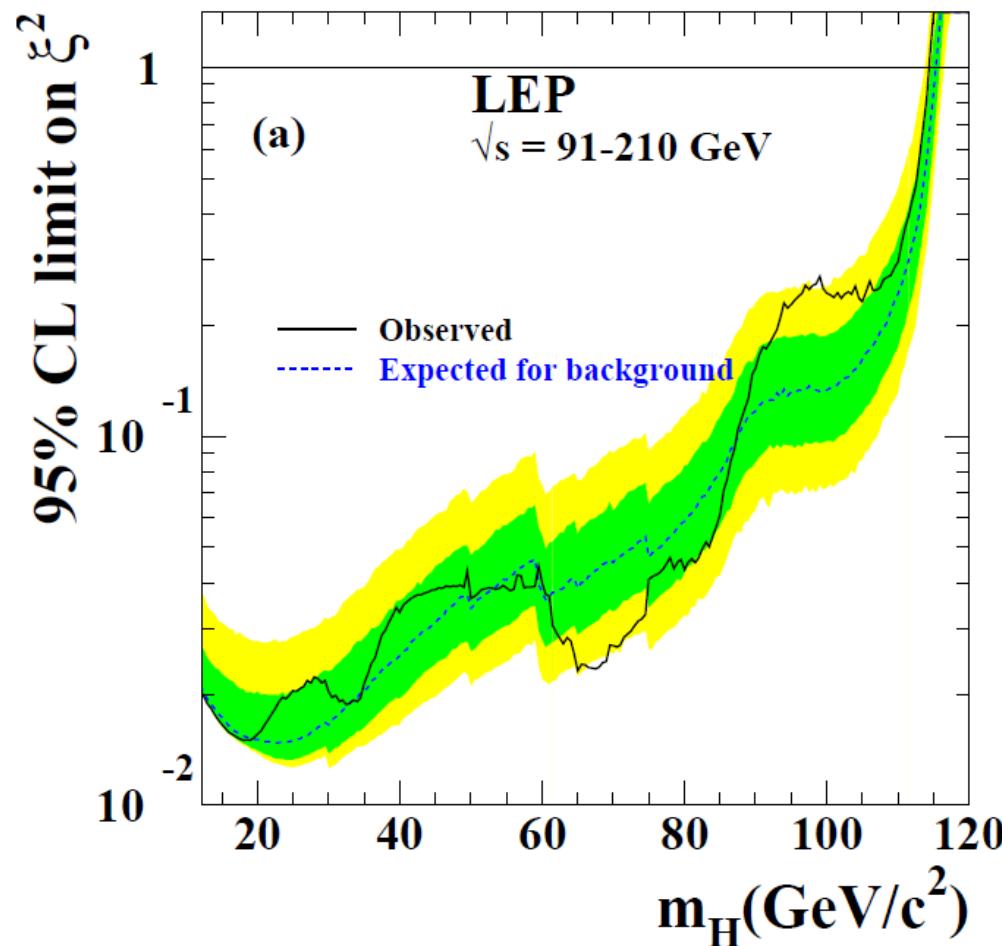
1 Additional BEH bosons $\gamma\gamma$ excess at 95 GeV

CMS PAS HIG-17-013



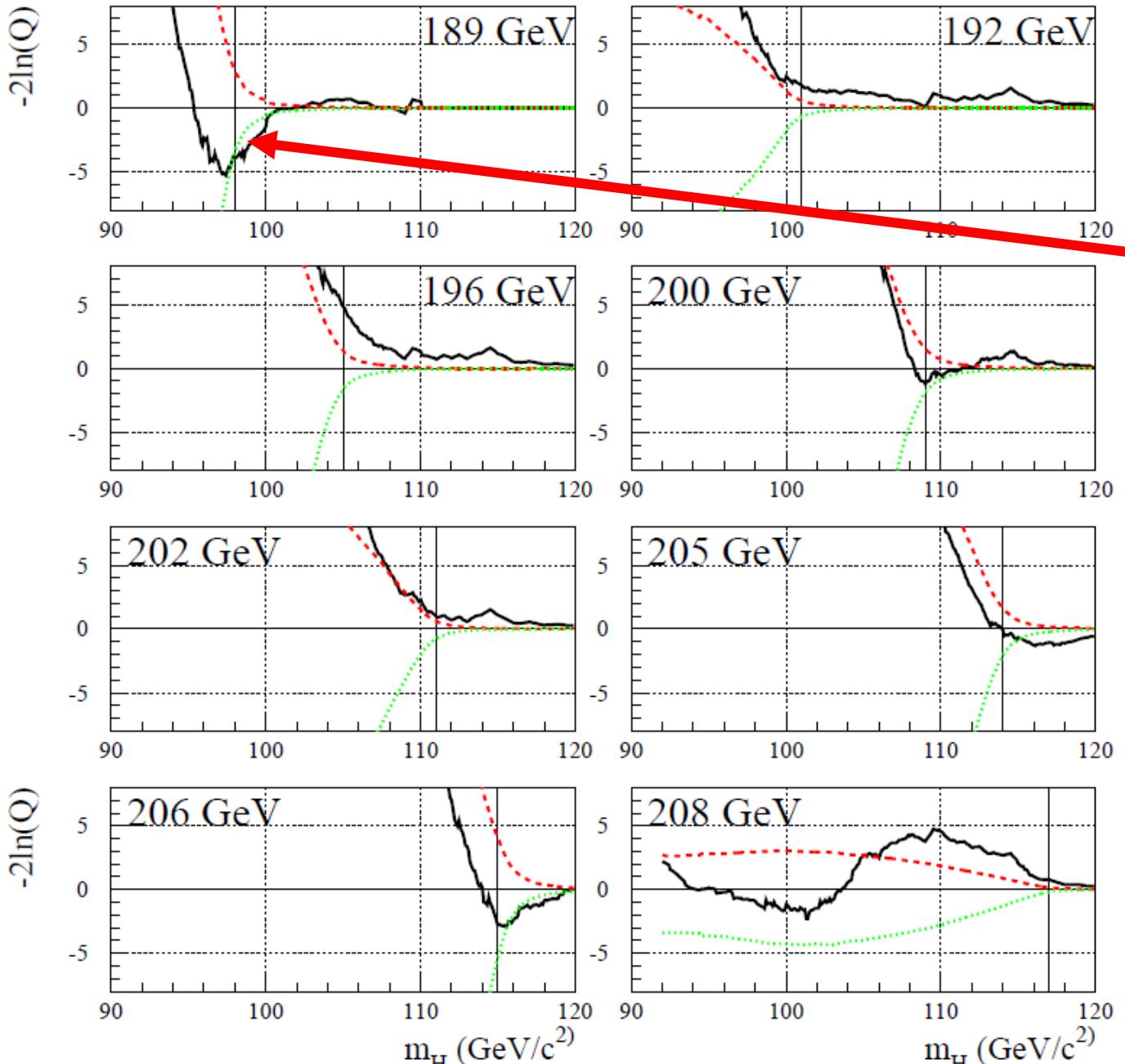
1 Additional BEH bosons $\gamma\gamma$ excess at 95 GeV

Search for the standard model
Higgs boson at LEP



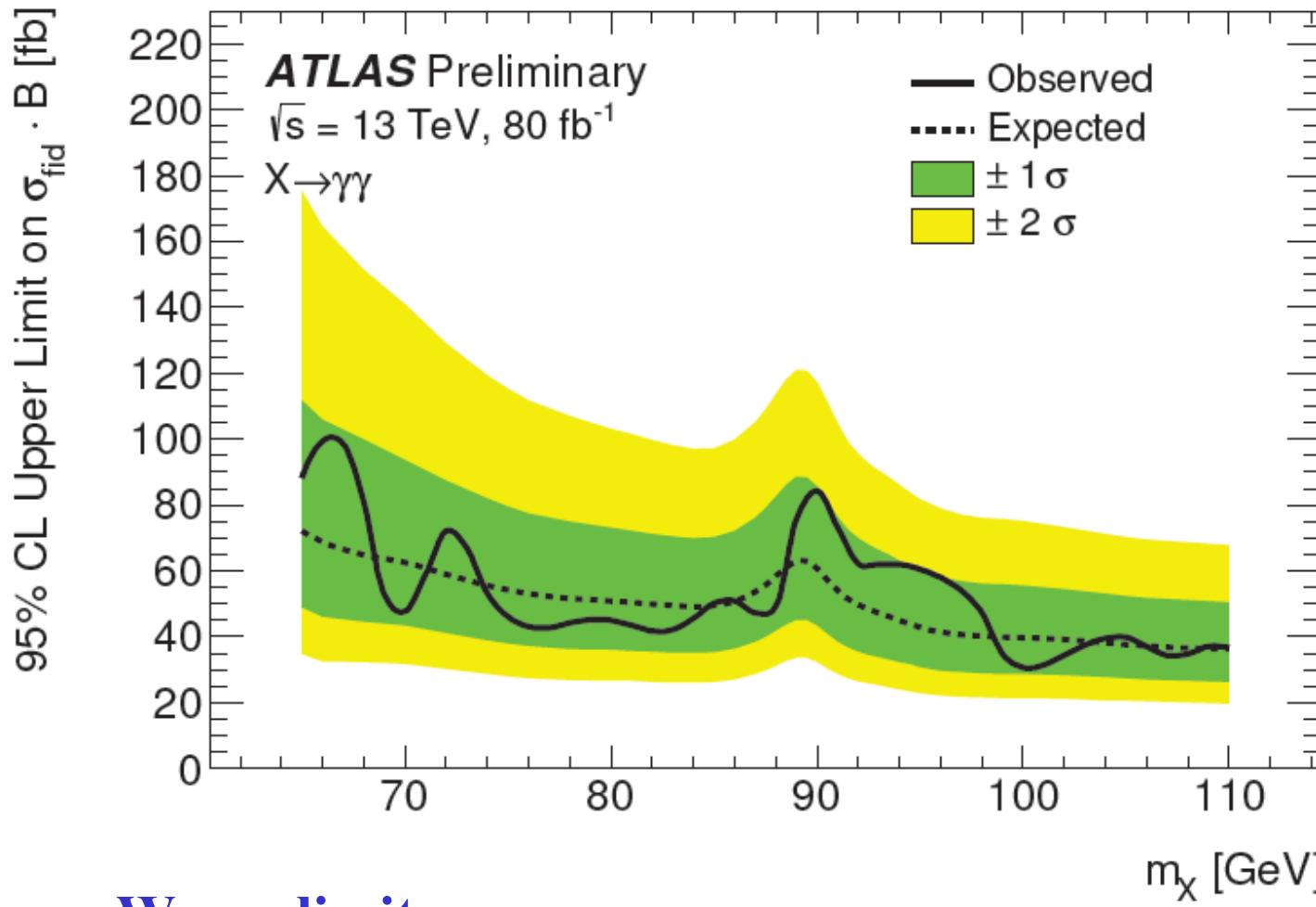
Phys.Lett. B565 (2003) 61-75

1 Additional BEH bosons $\gamma\gamma$ excess at 95 GeV



$$98 = 189 - m_Z$$

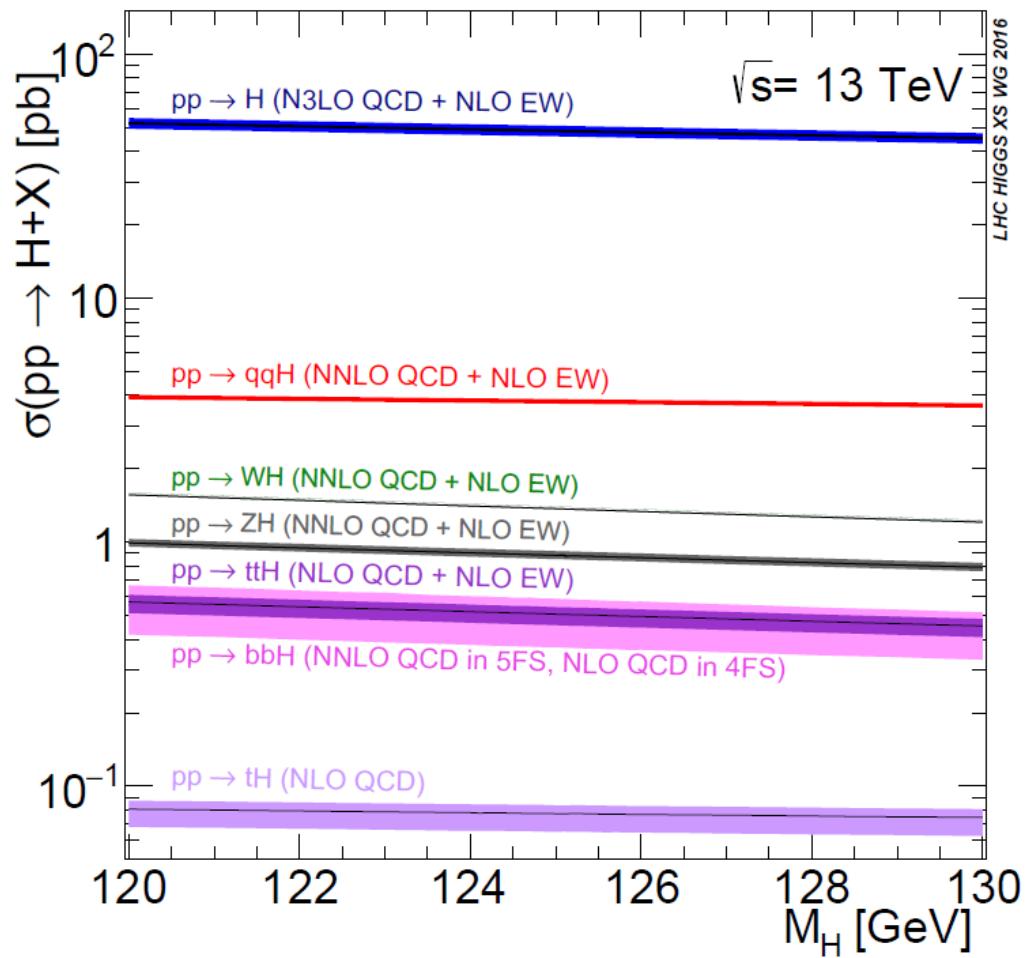
1 Additional BEH bosons $\gamma\gamma$ excess at 95 GeV



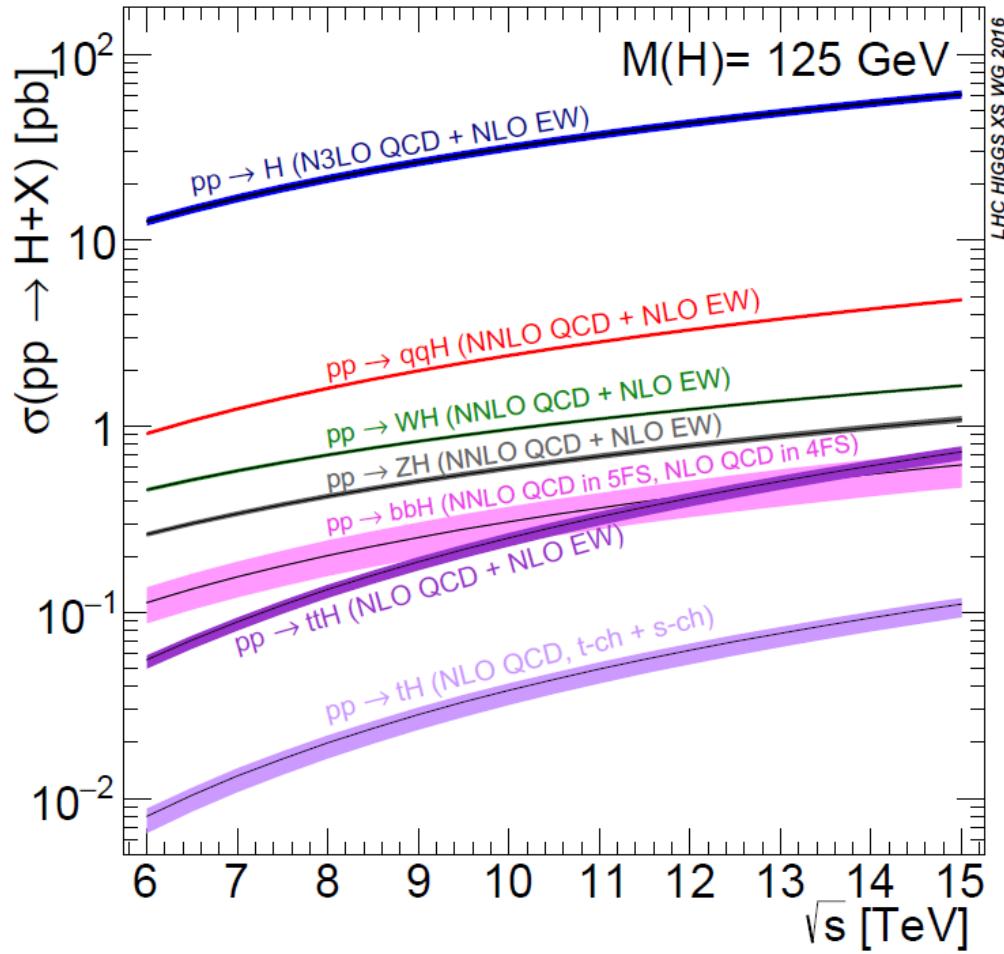
ATLAS-CONF-2018-025

.. Worse limit

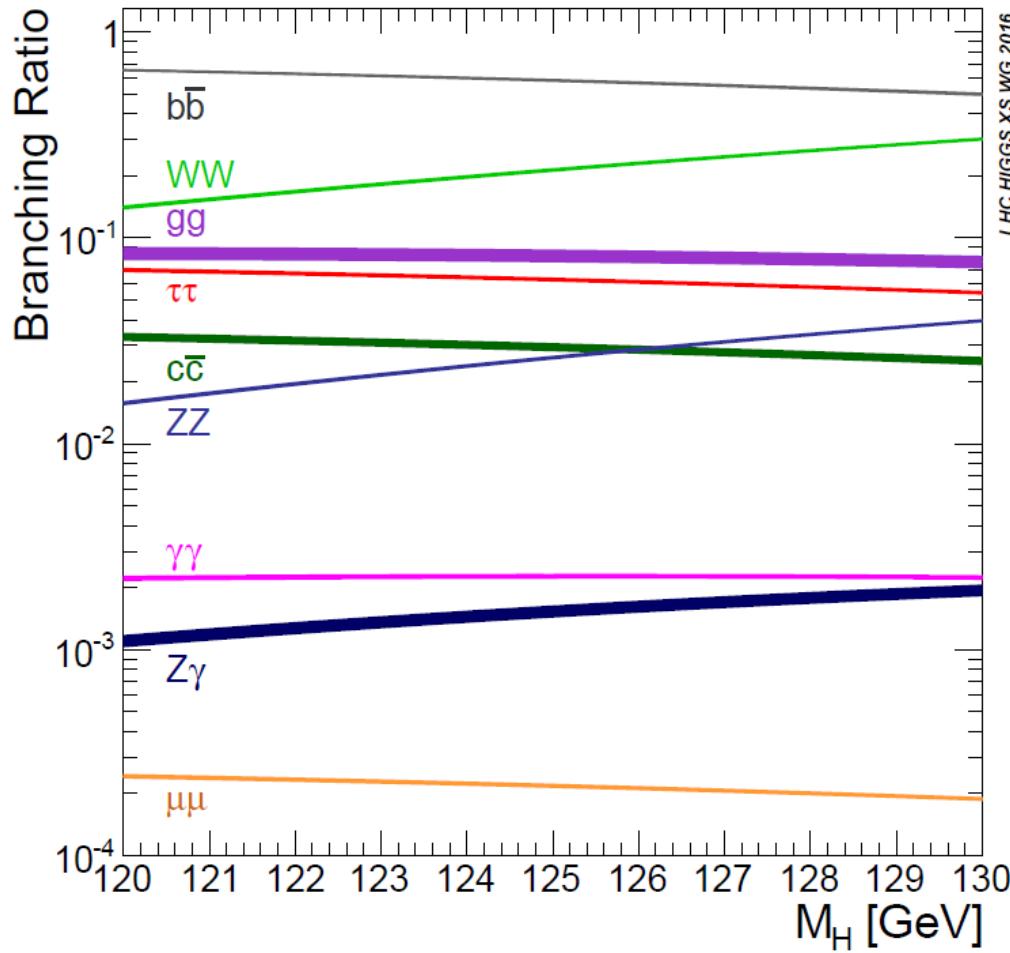
2 The SM BEH boson



2 The SM BEH boson

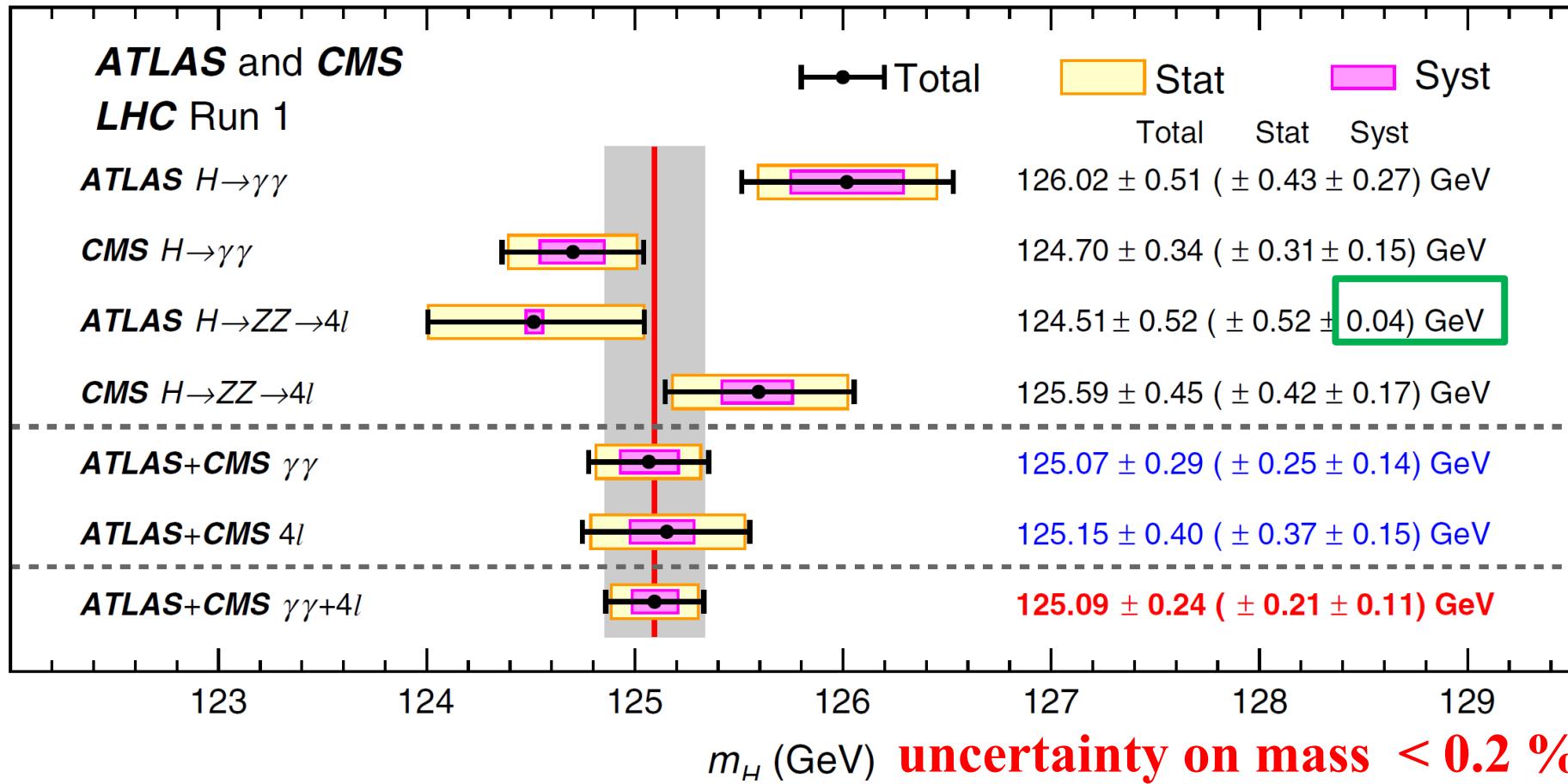


2 The SM BEH boson



2 The SM BEH boson (final Run-1 ATLAS + CMS result)

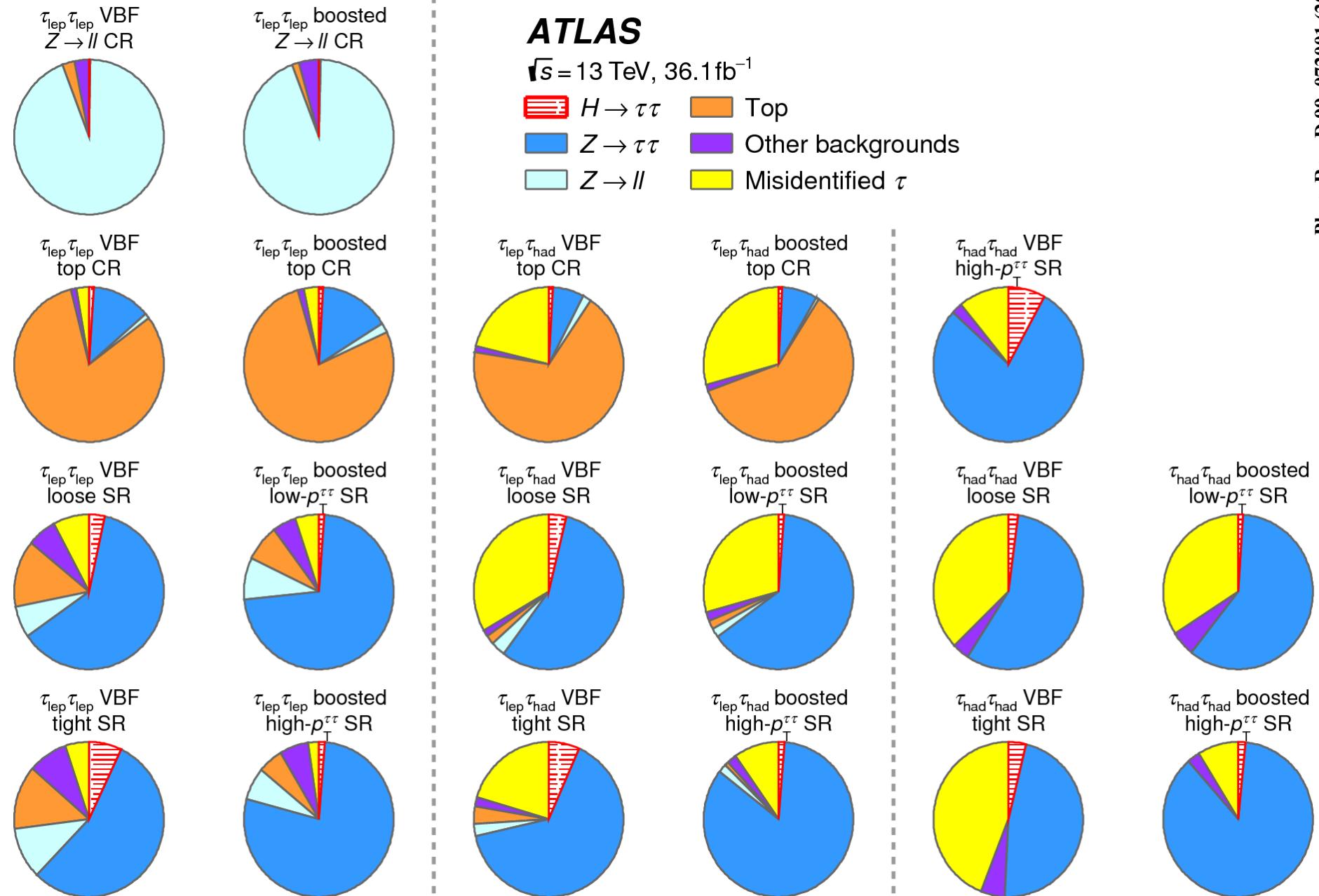
PRL 114, 191803 (2015)



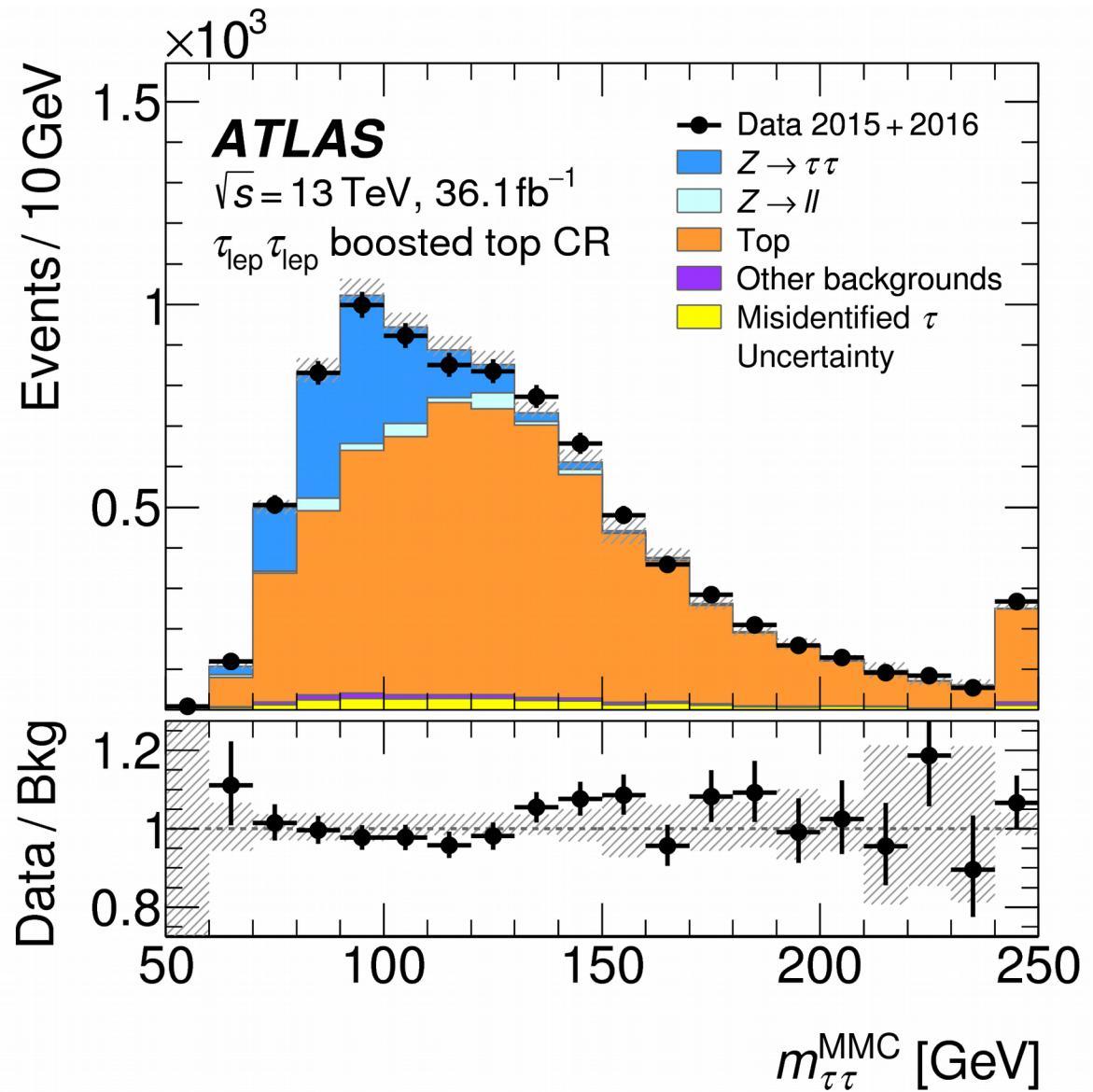
Remember ATLAS has an uncertainty on W mass of 19 MeV Eur.Phys.J. C78 (2018) no.2, 110
 note that $\Delta m_H = 0.1$ GeV $\rightarrow \Delta(\text{BR}(H \rightarrow ZZ)) / \text{BR}(H \rightarrow ZZ) \sim 1\%$

At longer term uncertainty will be dominated by 4l
 (for $H \rightarrow \gamma\gamma$: need to extrapolate from e to γ !)

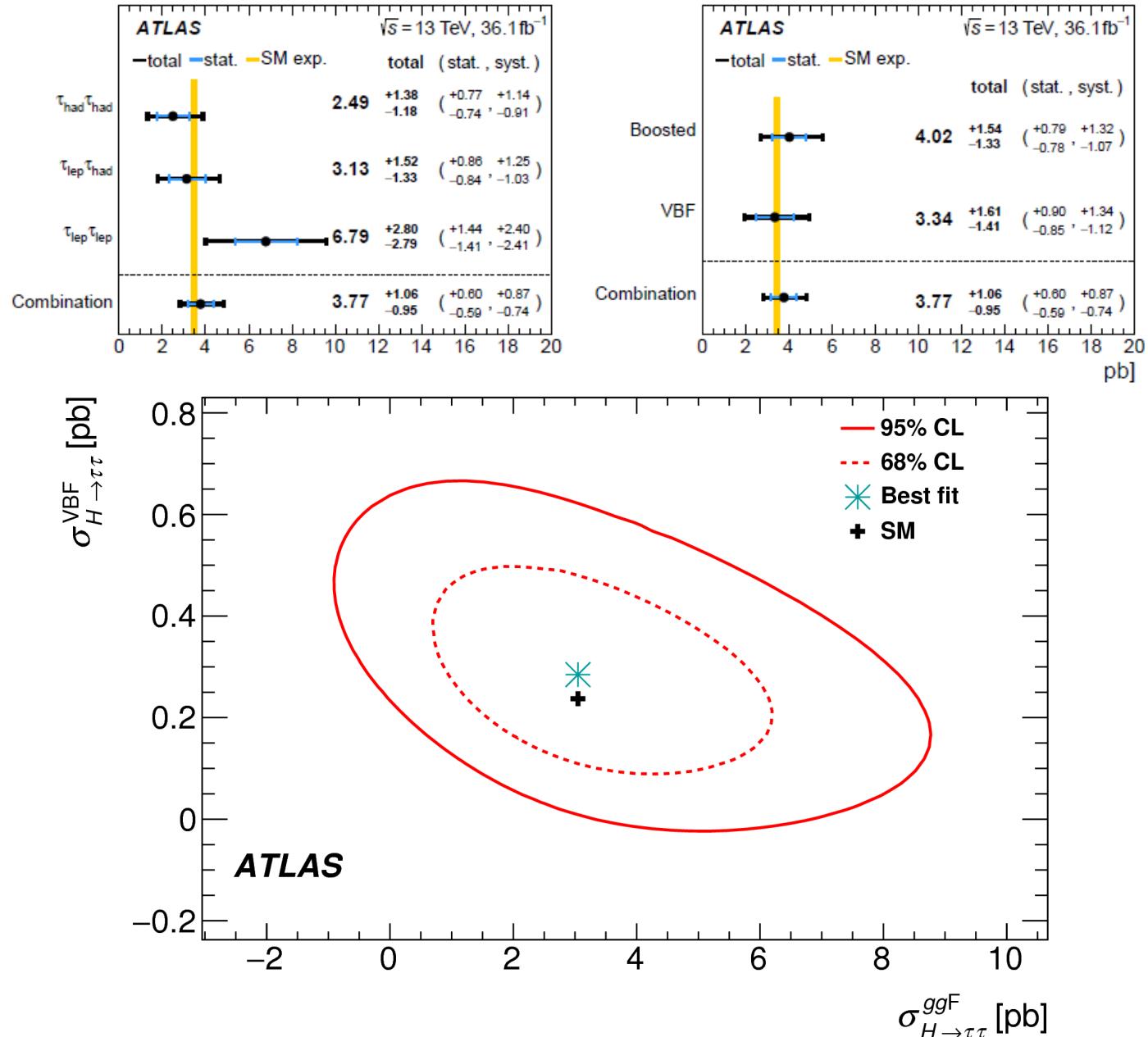
2 The SM BEH boson ($H \rightarrow \tau\tau$)



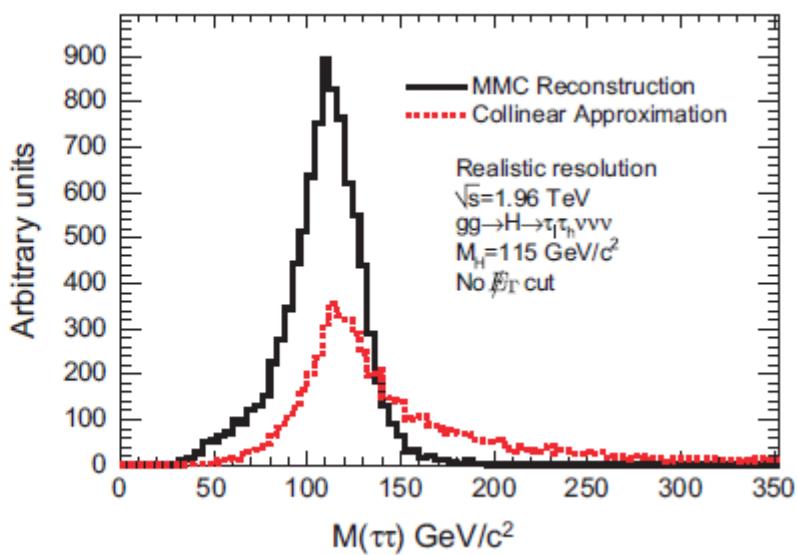
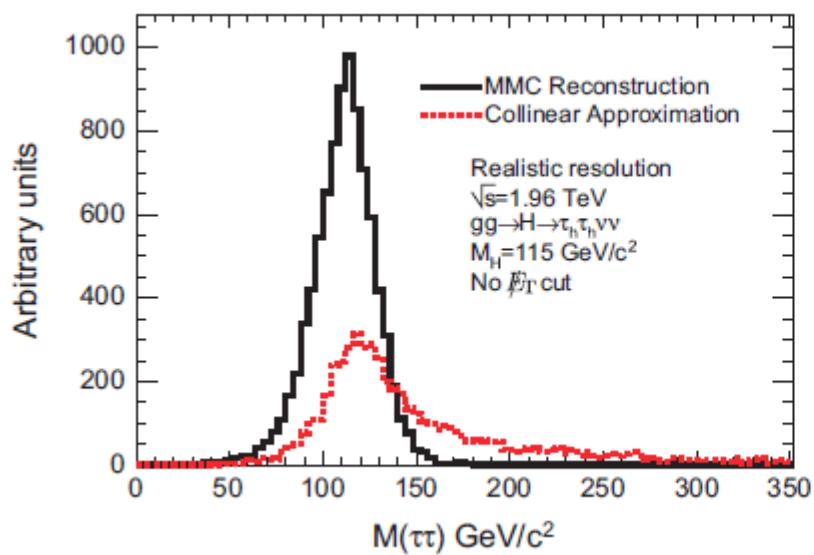
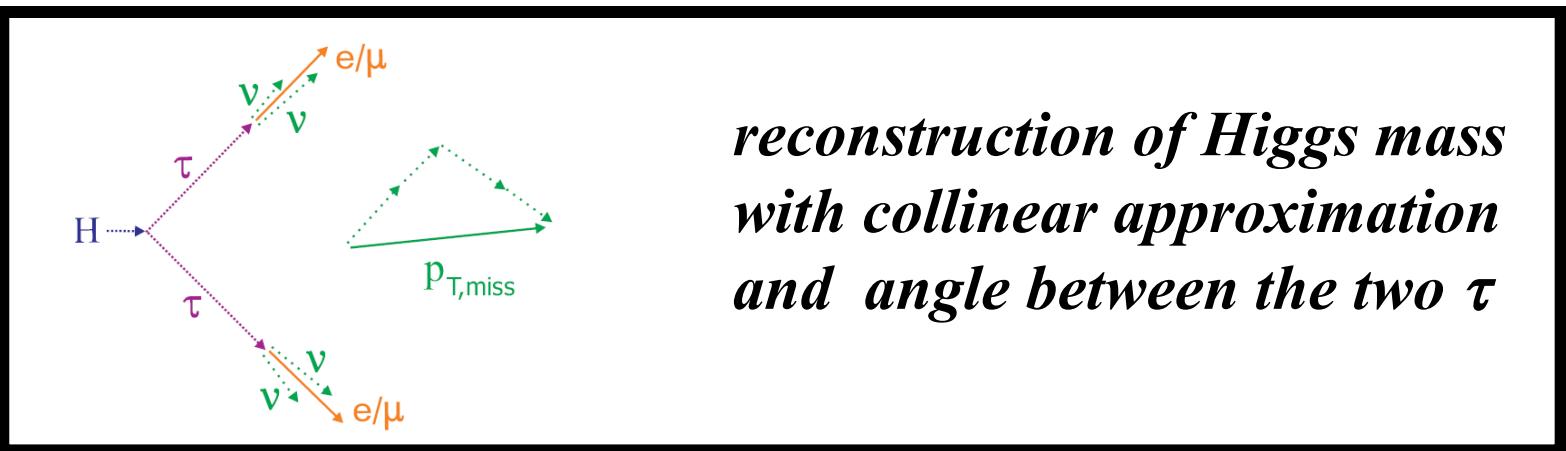
2 The SM BEH boson ($H \rightarrow \tau\tau$)



2 The SM BEH boson ($H \rightarrow \tau\tau$)



2 The SM BEH boson ($H \rightarrow \tau\tau$) τ reconstruction



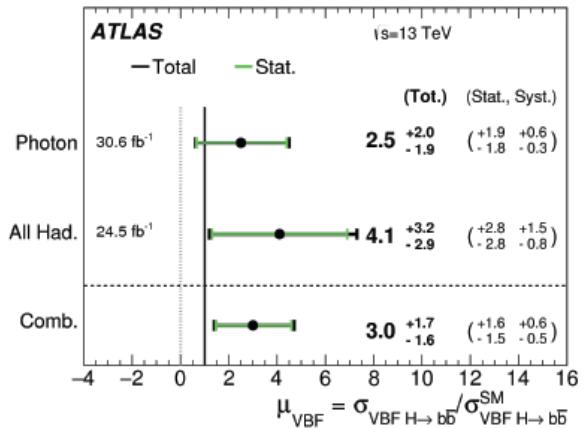
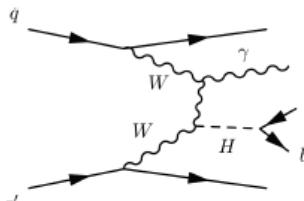
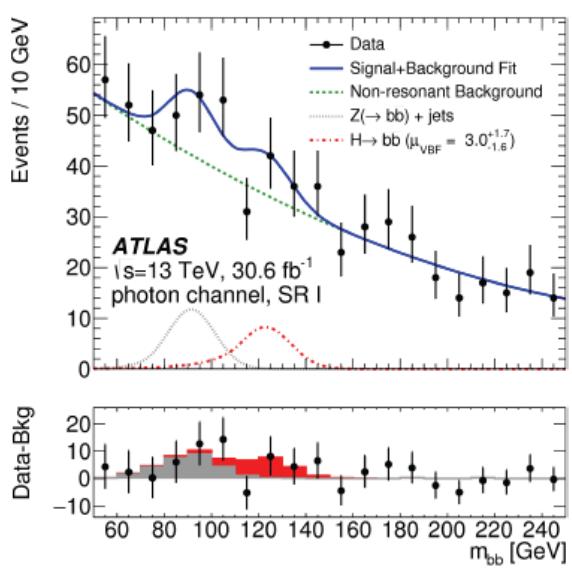
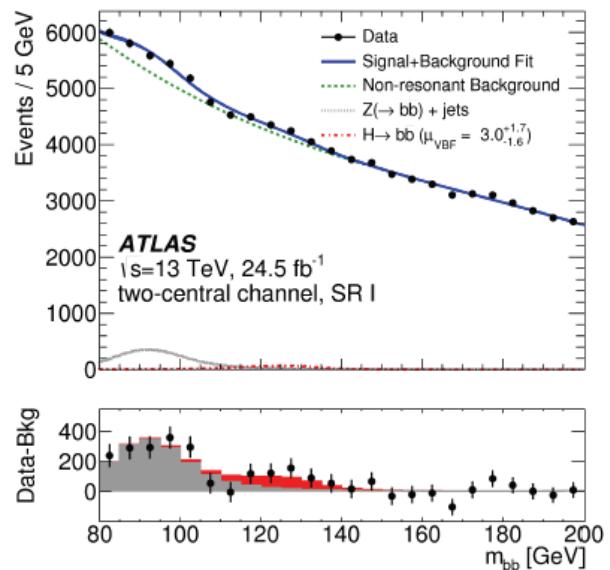
Improvement comes from requiring that the relative orientations of the neutrinos and other decay products are consistent with the mass and kinematics of a τ lepton decay

2 The SM BEH boson ($H \rightarrow bb$) targetting VBF

- Three analysis channels:

- 2 b-tagged jets + forward jet(s)
- 2 b-tagged jets + 2 central jets
- 2 b-tagged jets + 2 forward jets + central high- p_T γ

- Use BDT to select signal regions (Total of 9)



- Fit to the m_{bb} spectrum
- Significance: 1.9σ obs. (0.7σ exp.)
- Observed signal strength for the Higgs boson:
 $\mu_H = 2.5 \pm 1.3(\text{stat.}) \pm 0.6(\text{syst.})$
- Statistically limited, more data needed

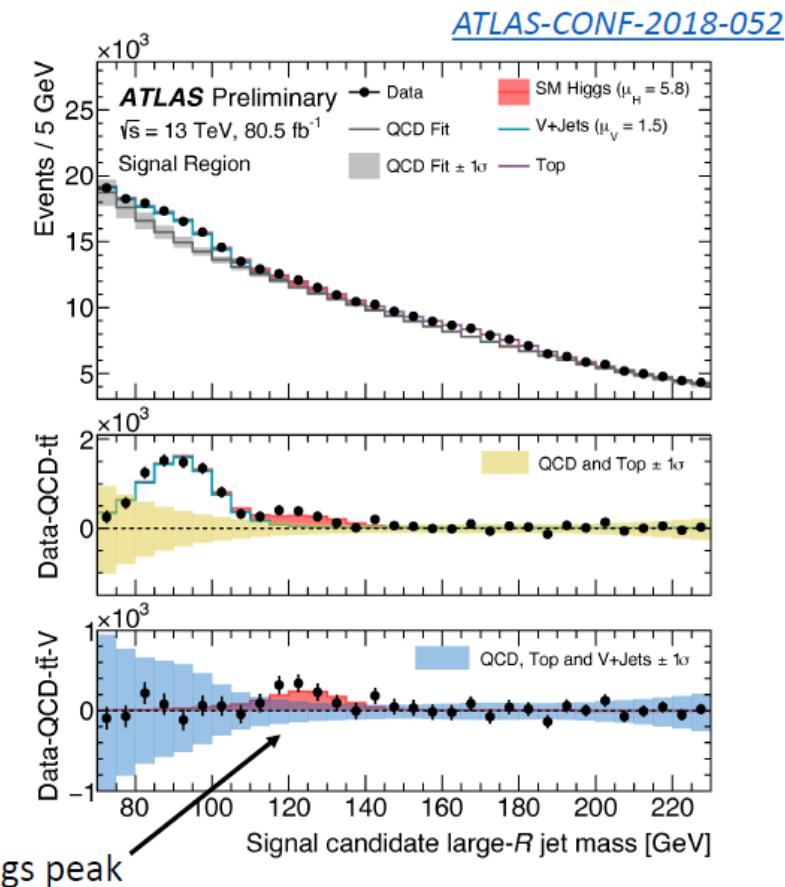
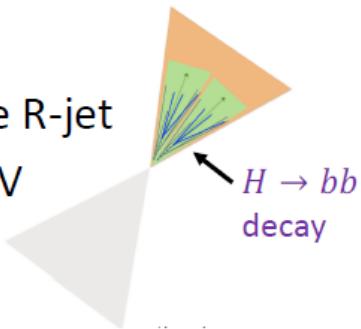
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2 The SM BEH boson ($H \rightarrow bb$) targetting $gg \rightarrow H J$

- Boosted analysis: 2 large-R jets
- 2 b-tagged track jets in one large R-jet
- Higgs-candidate jet: $p_T > 480$ GeV
- Fit to the jet-mass spectrum

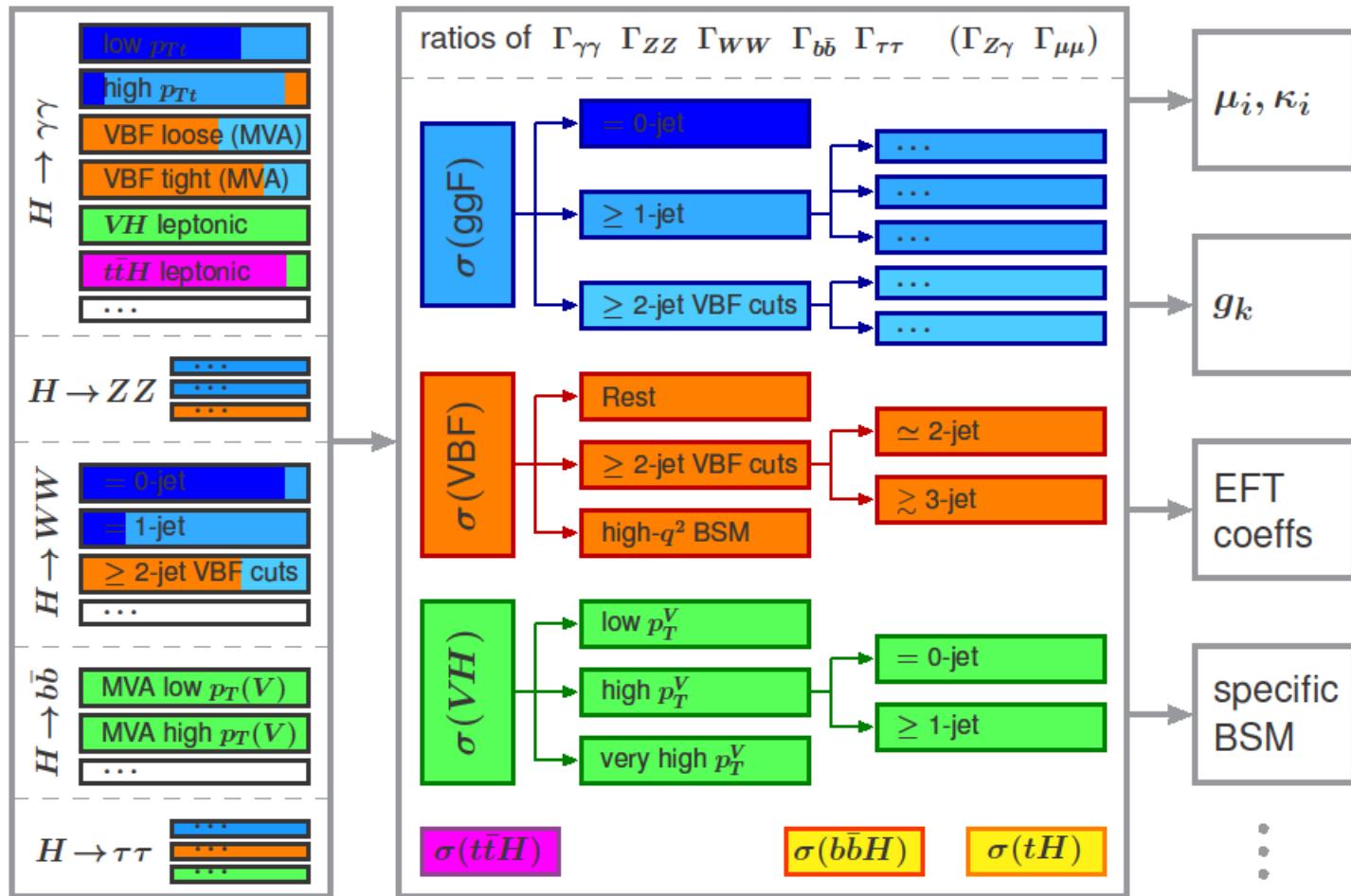
- Significance: 1.6σ obs. (0.28σ exp.)
- Needs more data

- Analysis is sensitive to $p_T^H > 480$ GeV
- Promising analysis to test for deviations from the SM and to include in STXS cross sections!



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2 The SM BEH boson STXS = simplified template cross sections



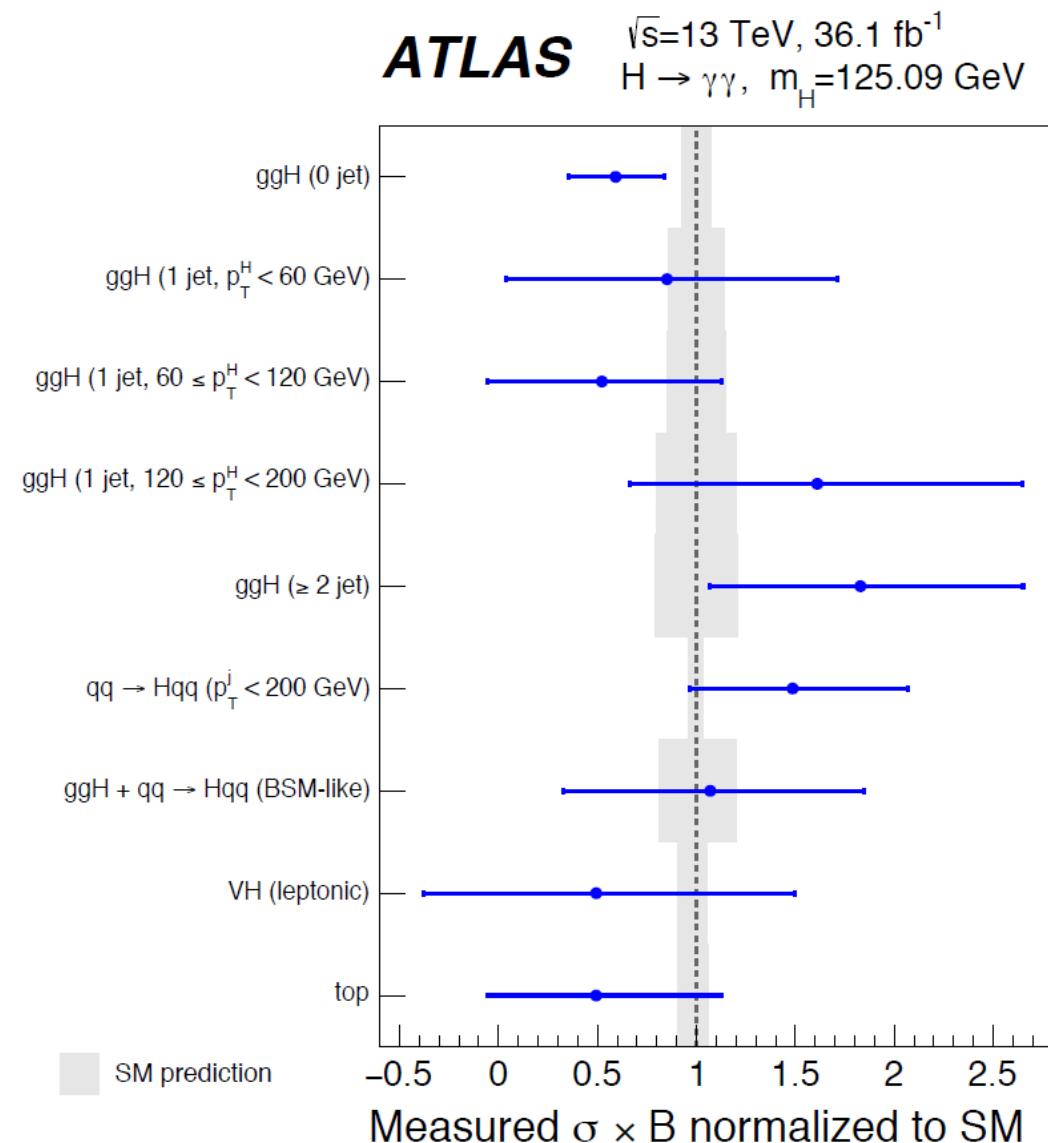
2 The SM BEH boson STXS = simplified template cross sections

designed to measure the different Higgs boson production processes in specific regions of phase space and in a way that can be easily combined with other decay channels

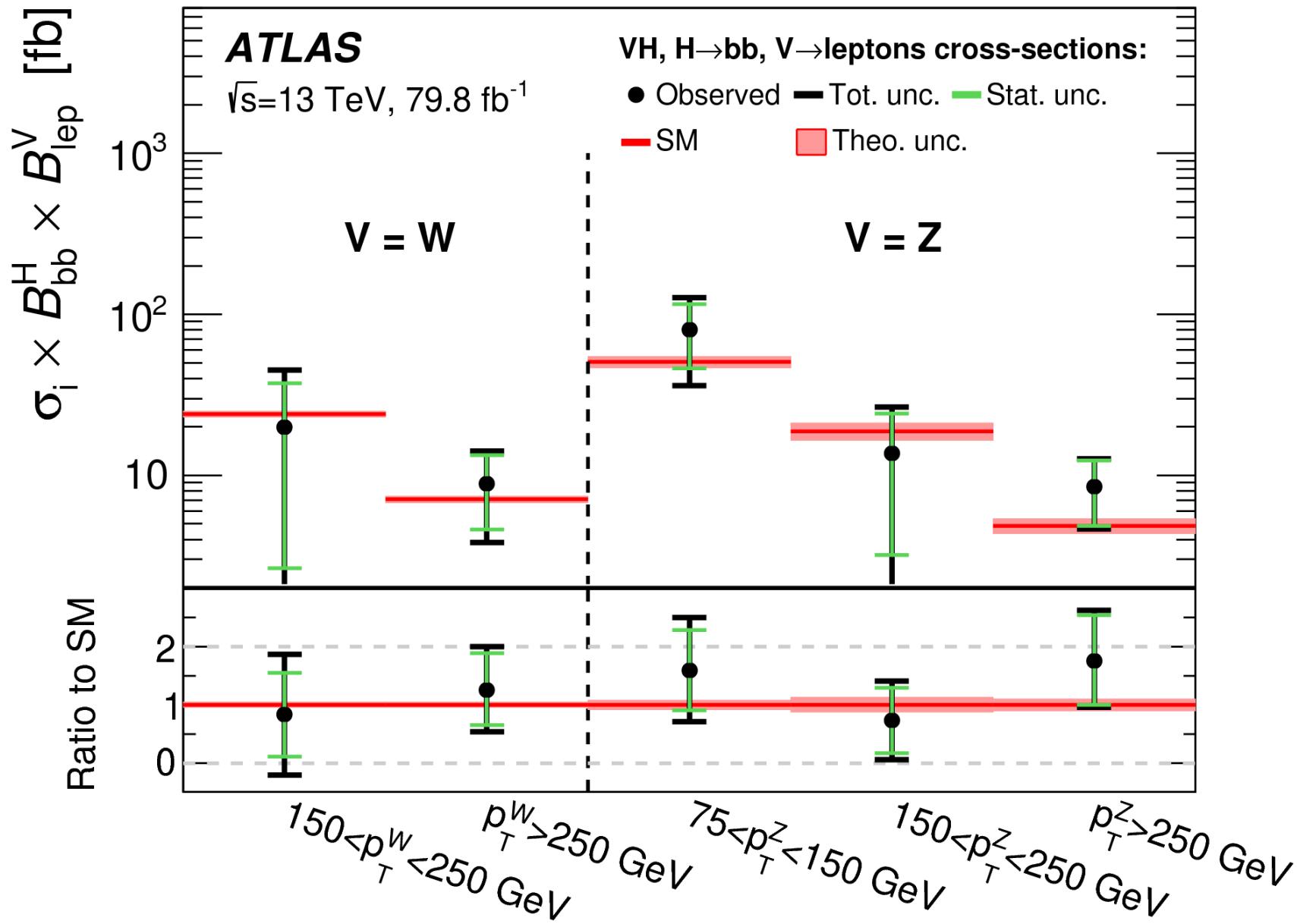
Compared to the signal strength measurements they provide finer granularity

theory uncertainties are smaller

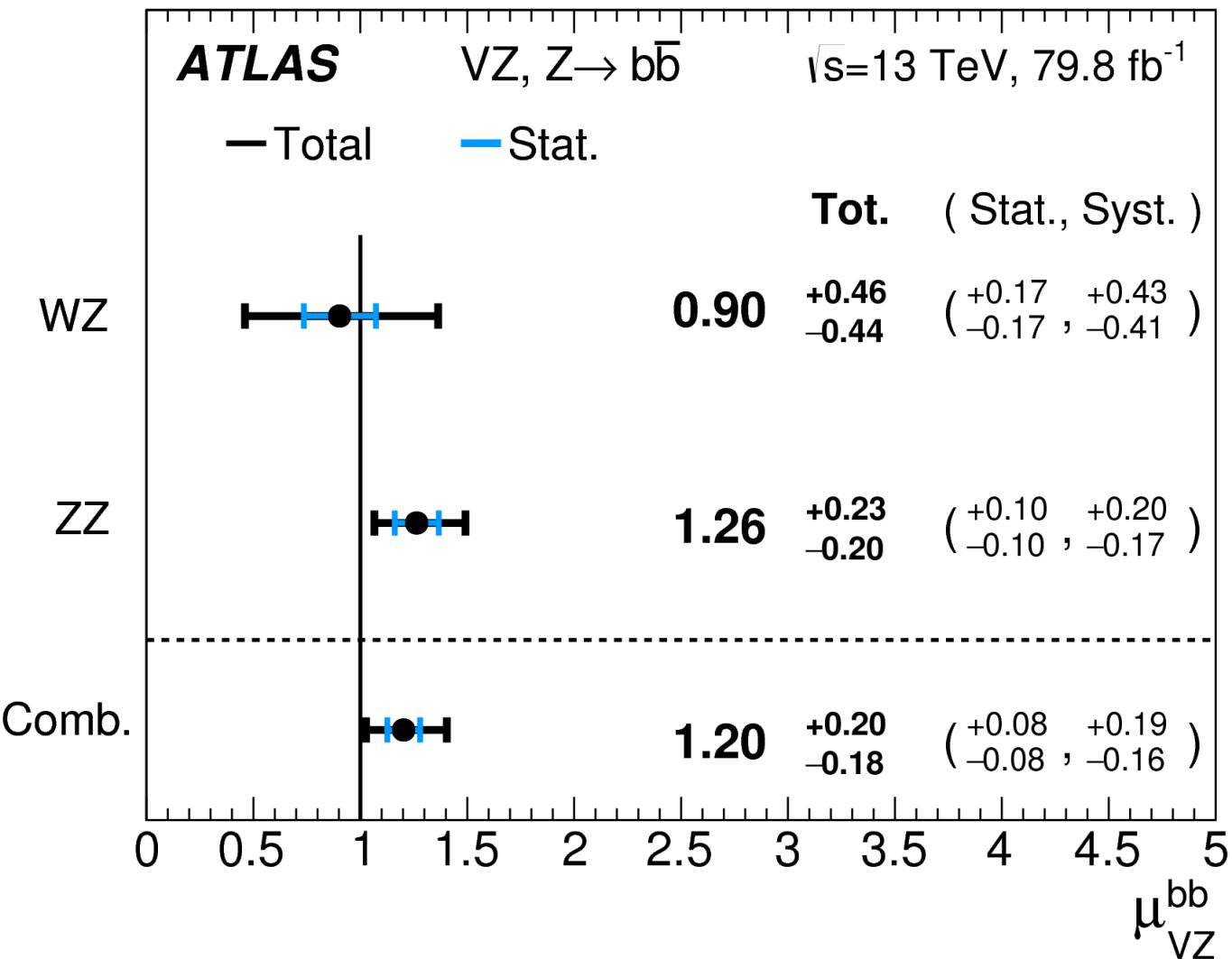
In fact there are 31 STXS , but measure 9 (lack of statistics)



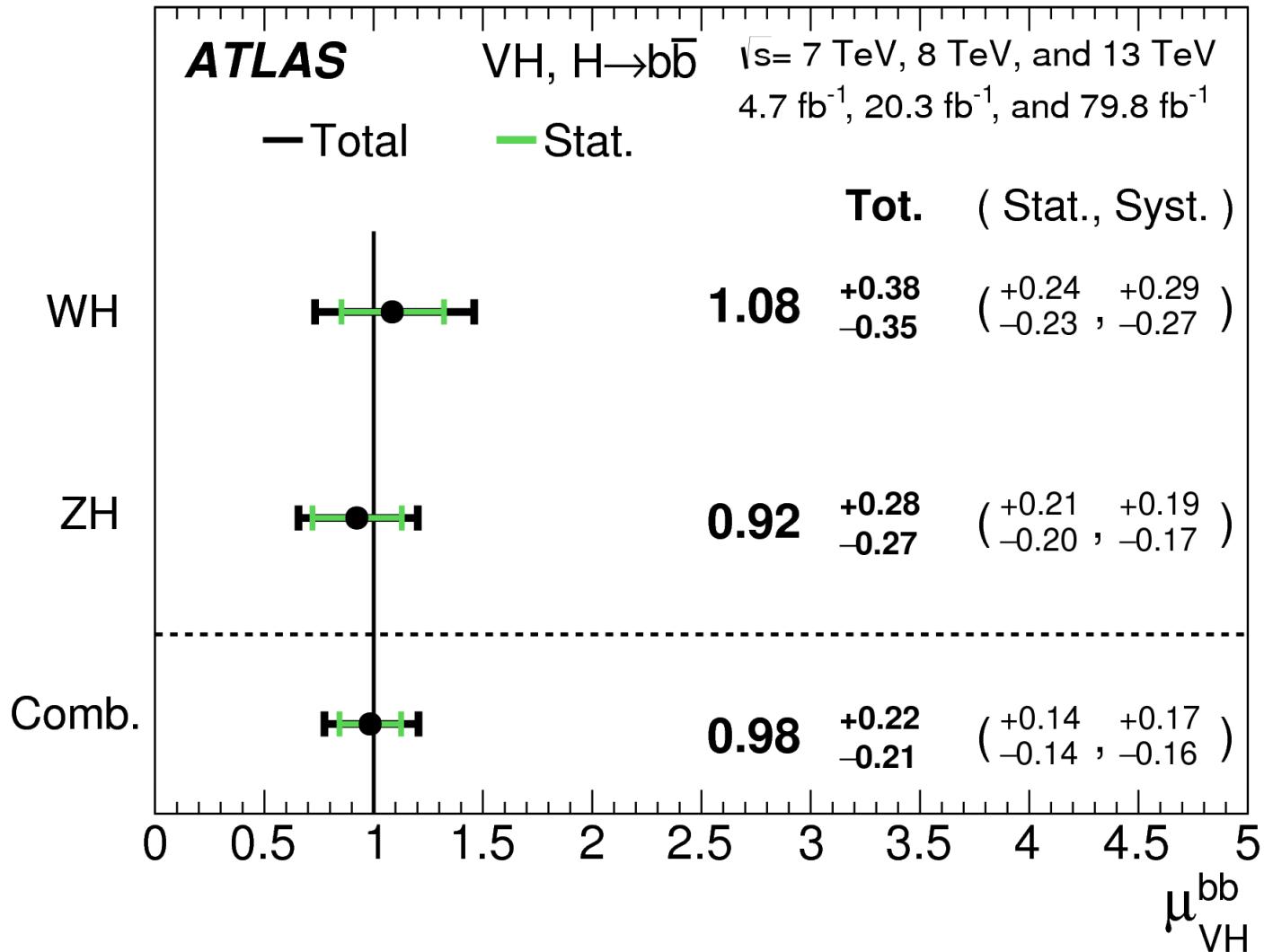
2 The SM BEH boson ($H \rightarrow bb$) STXS



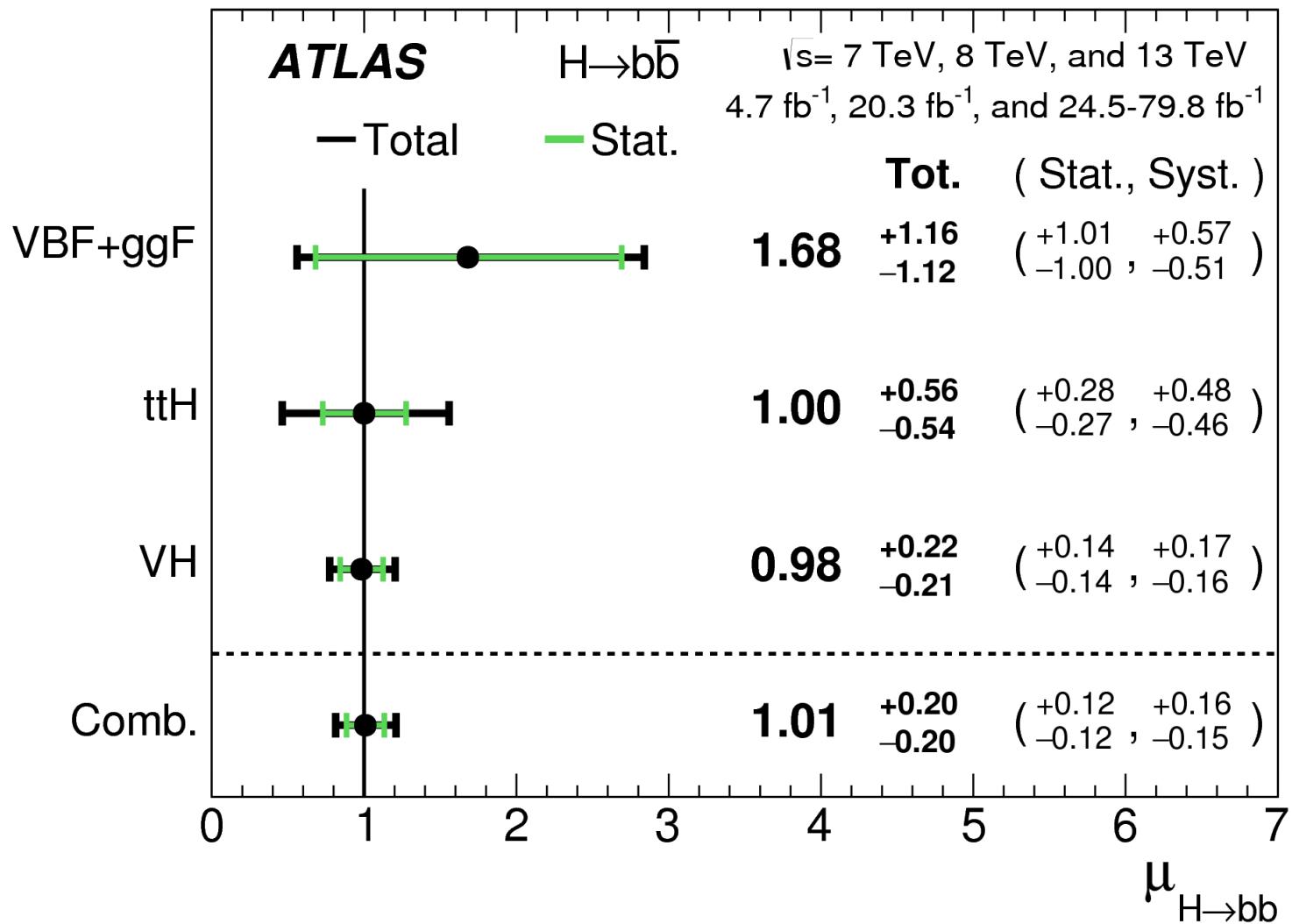
2 The SM BEH boson V ($Z \rightarrow b\bar{b}$)



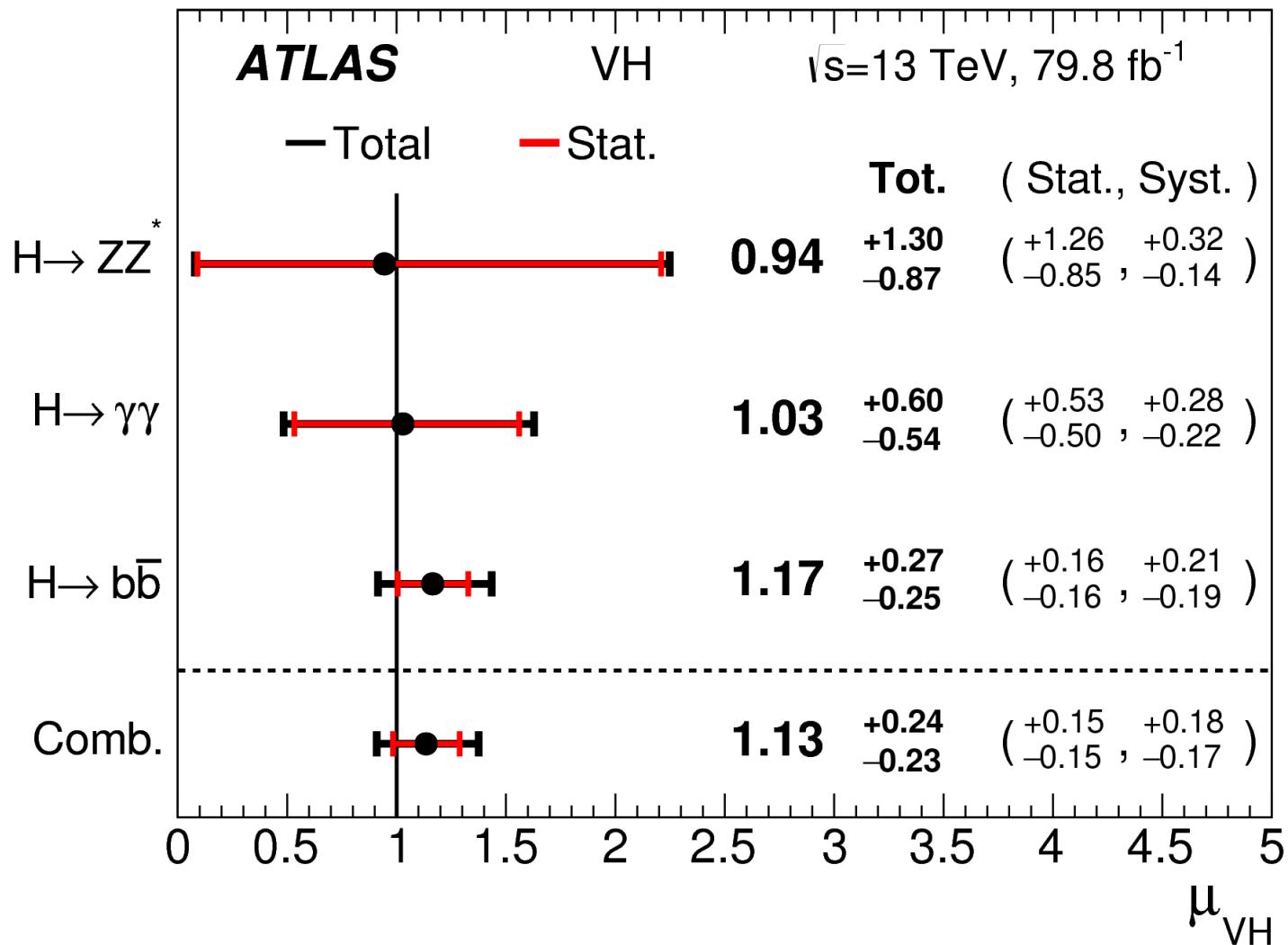
2 The SM BEH boson V ($H \rightarrow b\bar{b}$)



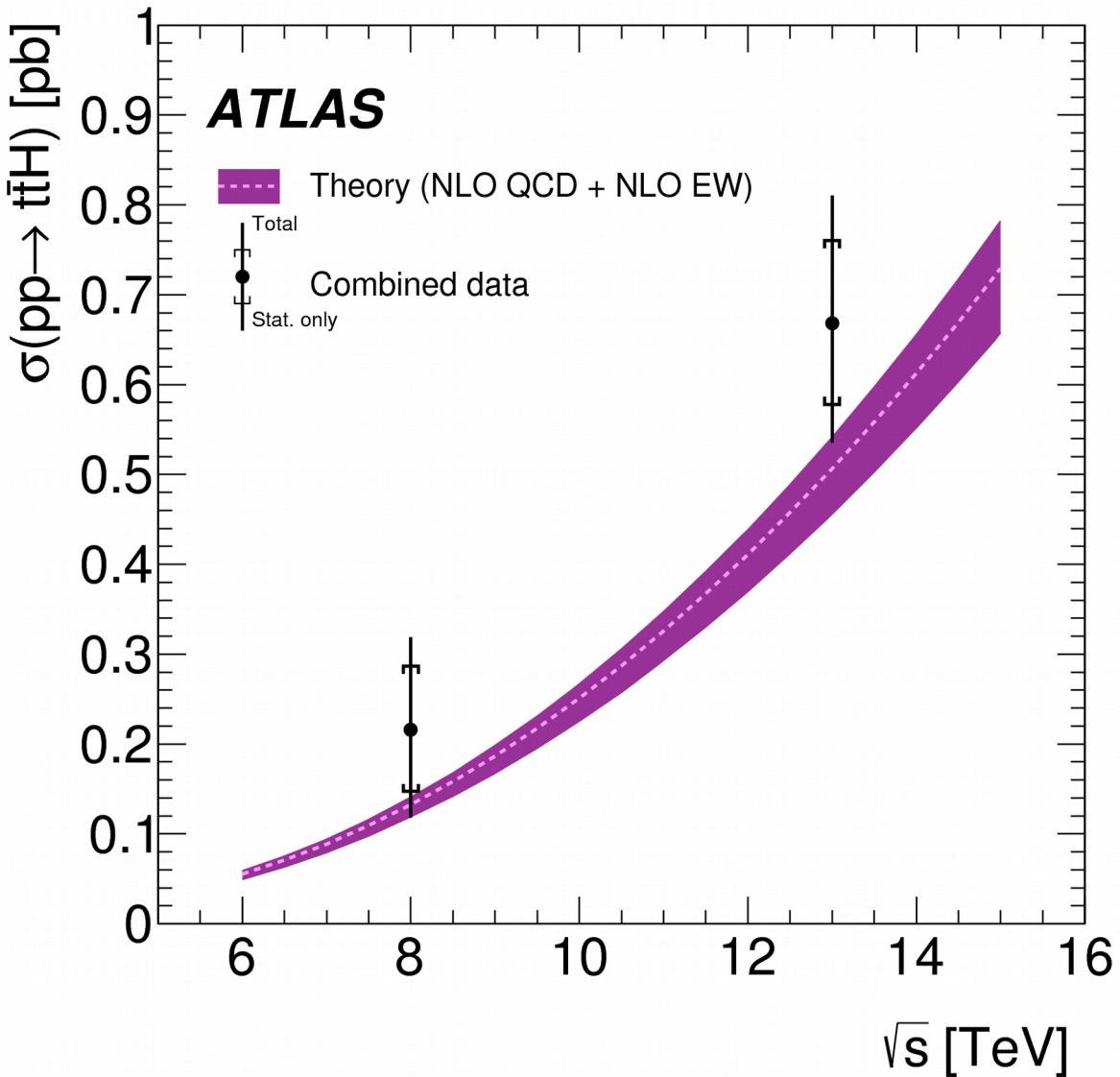
2 The SM BEH boson ($H \rightarrow b\bar{b}$)



2 The SM BEH boson VH

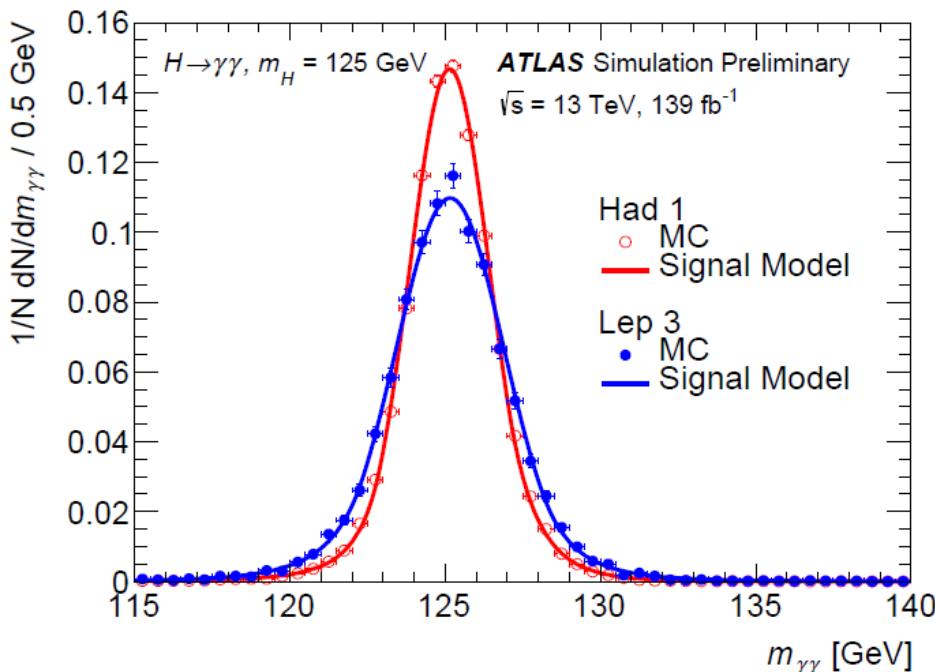
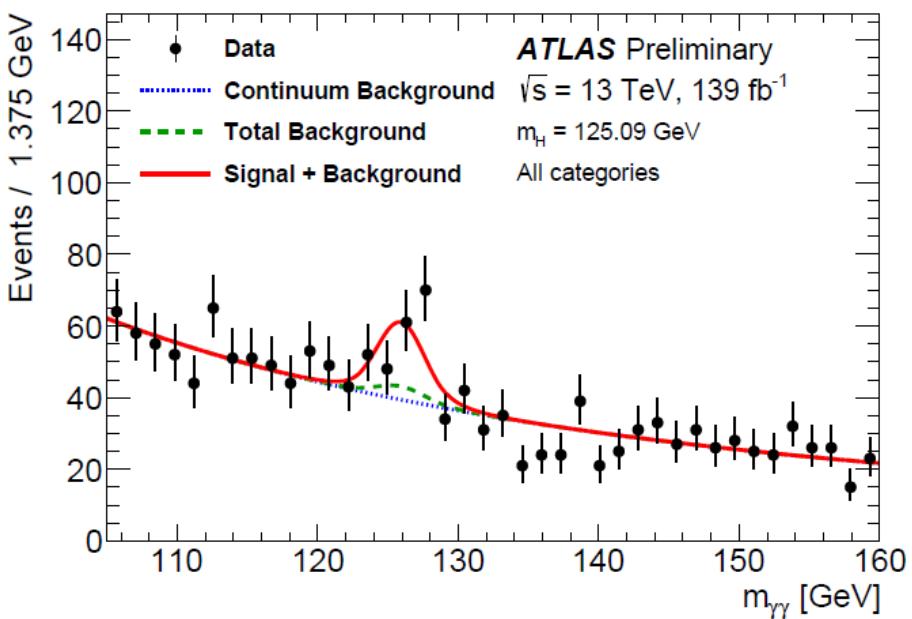


2 The SM BEH boson ttH



2 The SM BEH boson ttH ($\rightarrow \gamma\gamma$)

Category	σ_{68} (GeV)	σ_{90} (GeV)
“Lep” Category 1	1.56	2.80
“Lep” Category 2	1.75	3.13
“Lep” Category 3	1.85	3.30
“Had” Category 1	1.39	2.48
“Had” Category 2	1.58	2.84
“Had” Category 3	1.65	2.96
“Had” Category 4	1.67	3.00



2 The SM BEH boson categories

Several categories are made in order to enhance the sensitivity in order to have different S/B , based on

- number of jets
- different resolutions
- different kinematics giving different S/B

S/B has to be different for various categories

This is needed if we want to gain in statistical significance

if $S_1 / B_1 = S_2 / B_2$

then $S_1 / \sqrt{B_1} \oplus S_2 / \sqrt{B_2} = (S_1 + S_2) / \sqrt{(B_1 + B_2)}$

and one does not gain making categories

(one of) the work of the experimentalist is to find categories with different S/B !

2 The SM BEH boson invisible H decays

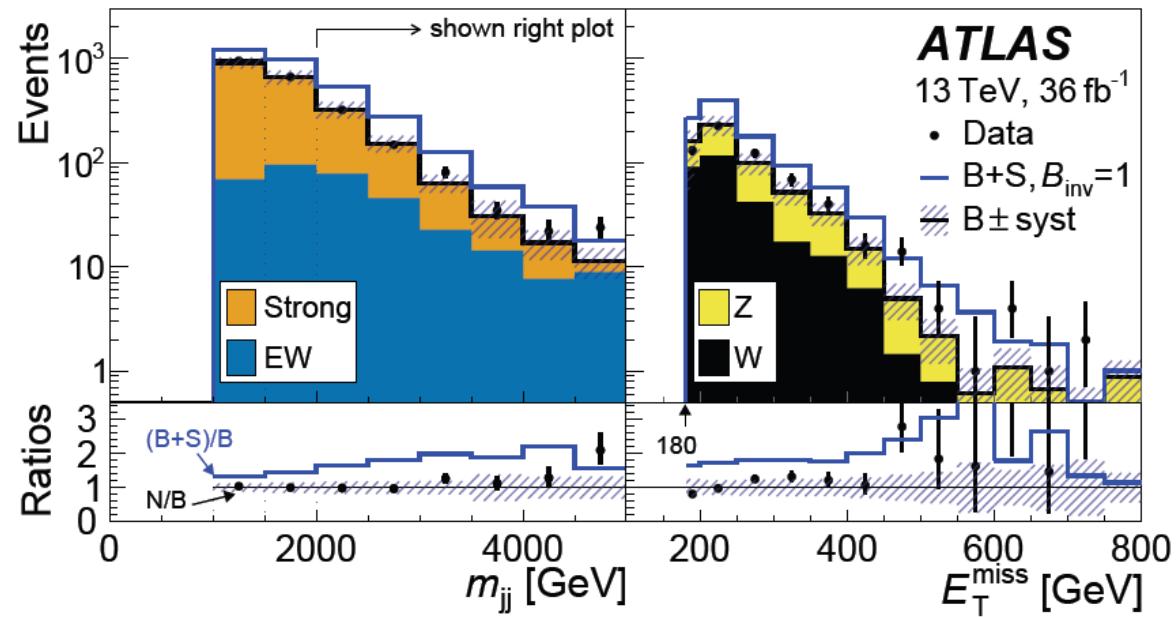
three bins of m_{jj} defined by
boundaries at [1, 1.5, 2, -] TeV

For the SR, an event is required to have

- no isolated electron or muon,
- a leading jet with $p_T > 80$ GeV,
- a subleading jet with $p_T > 50$ GeV,
- no additional jets with $p_T > 25$ GeV,
- $E_T^{\text{miss}} > 180$ GeV,
- $H_T^{\text{miss}} > 150$ GeV.

The two jets are required to have the following properties:

- not be aligned with \vec{E}_T^{miss} , $|\Delta\phi_{j-\text{MET}}| > 1$,
- not be back-to-back, $|\Delta\phi_{jj}| < 1.8$,
- be well separated in η , $|\Delta\eta_{jj}| > 4.8$,
- be in opposite η hemispheres, $\eta_{j_1} \cdot \eta_{j_2} < 0$,
- $m_{jj} > 1$ TeV.

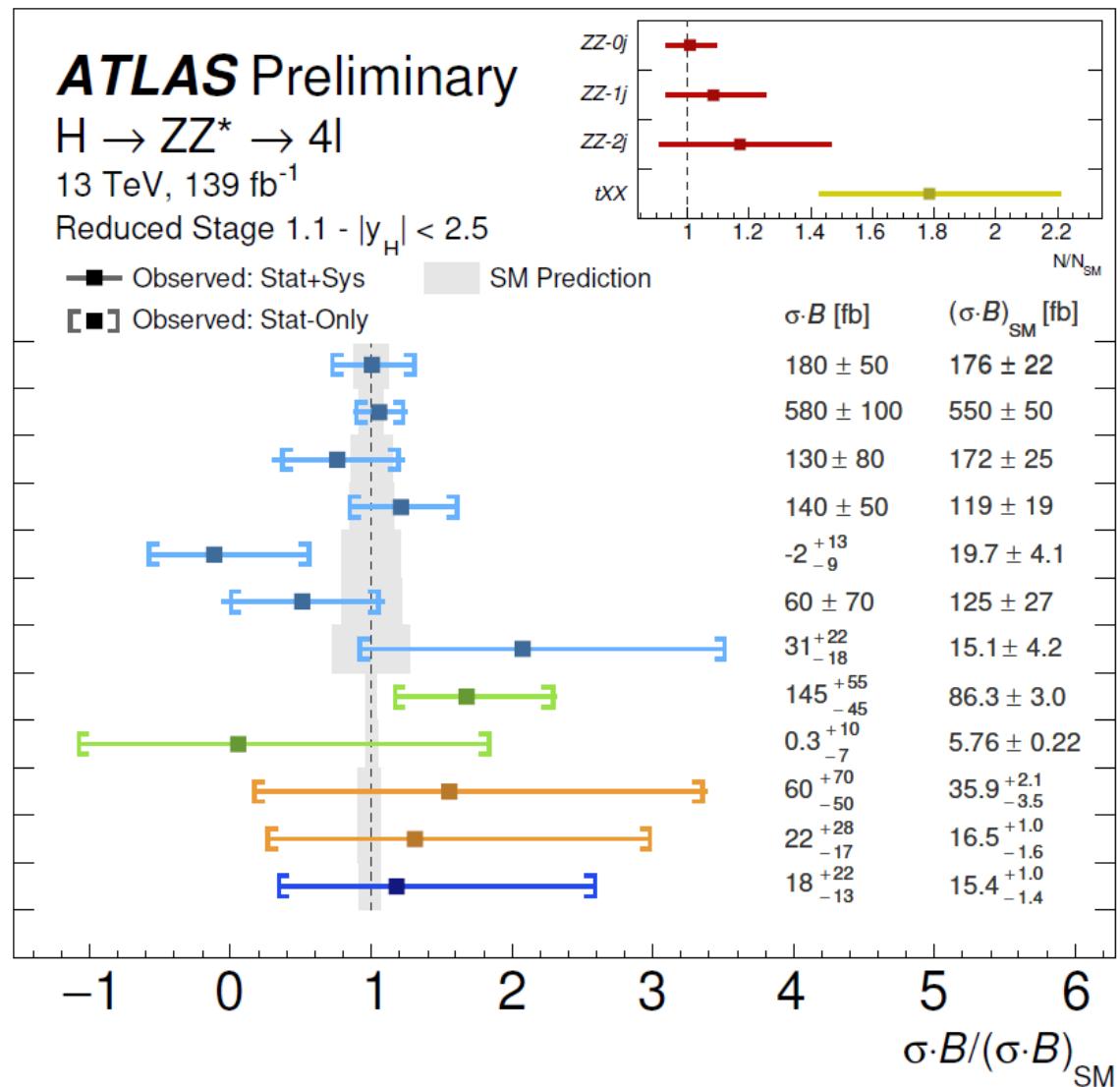


2 The SM BEH boson $H \rightarrow \mu\mu$

Category	Data	S_{SM}	S	B	S/\sqrt{B}	$S/B [\%]$
VBF High	40	4.5	2.3	34	0.39	6.6
VBF Medium	109	5.5	2.8	100	0.28	2.8
VBF Low	450	9.6	4.9	420	0.24	1.2
2-jet High	3400	38	19	3440	0.33	0.6
2-jet Medium	13938	70	35	13910	0.30	0.3
2-jet Low	40747	75	38	40860	0.19	0.1
1-jet High	2885	32	16	2830	0.31	0.6
1-jet Medium	24919	107	54	24890	0.35	0.2
1-jet Low	77482	134	68	77670	0.24	0.1
0-jet High	24777	85	43	24740	0.27	0.2
0-jet Medium	85281	155	79	85000	0.27	0.1
0-jet Low	180478	144	73	180000	0.17	<0.1

ATLAS-CONF-2019-028

2 The SM BEH boson $H \rightarrow 4l$



2 The SM BEH boson $H \rightarrow \gamma\gamma$

Table 3: The breakdown of uncertainties on the inclusive diphoton fiducial cross section measurement. The uncertainties from the statistics of the data and the systematic sources affecting the signal extraction are shown. The remaining uncertainties are associated with the unfolding correction factor and luminosity.

Source	Uncertainty (%)
Statistics	6.9
Signal extraction syst.	7.9
Photon energy scale & resolution	4.6
Background modelling (spurious signal)	6.4
Correction factor	2.6
Pile-up modelling	2.0
Photon identification efficiency	1.2
Photon isolation efficiency	1.1
Trigger efficiency	0.5
Theoretical modelling	0.5
Photon energy scale & resolution	0.1
Luminosity	1.7
Total	11.0

ATLAS-CONF-2019-029

139 fb⁻¹

Source	Uncertainty (%)
Fit (stat.)	10
Fit (syst.)	8.3
Photon energy scale & resolution	4.0
Background modeling (spurious signal)	7.3
Correction factor	5.2
Photon isolation efficiency	4.6
Pileup	1.9
Photon ID efficiency	1.3
Trigger efficiency	0.7
Dalitz Decays	0.4
Theoretical modeling	+0.3 -0.4
Diphoton vertex selection	0.1
Photon energy scale & resolution	0.1
Luminosity	2.0
Total	14

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80 fb⁻¹

Uncertainty in fiducial cross section

Source	Diphoton
Fit (stat.)	17%
Fit (syst.)	6%
Photon energy scale & resolution	4.3%
Background modelling	4.2%
Photon efficiency	1.8%
Jet energy scale/resolution	-
<i>b</i> -jet flavor tagging	-
Lepton selection	-
Pileup	1.1%
Theoretical modeling	0.1%
Signal composition	0.1%
Higgs boson p_T^H & $ y_H $	0.1%
UE/PS	-
Luminosity	3.2%
Total	18%

Uncertainty Group	$\sigma_\mu^{\text{syst.}}$
Theory (QCD)	0.041
Theory ($B(H \rightarrow \gamma\gamma)$)	0.028
Theory (PDF+ α_S)	0.021
Theory (UE/PS)	0.026
Luminosity	0.031
Experimental (yield)	0.017
Experimental (migrations)	0.015
Mass resolution	0.029
Mass scale	0.006
Background shape	0.027

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36 fb⁻¹

2 The SM BEH boson $H \rightarrow \gamma\gamma$

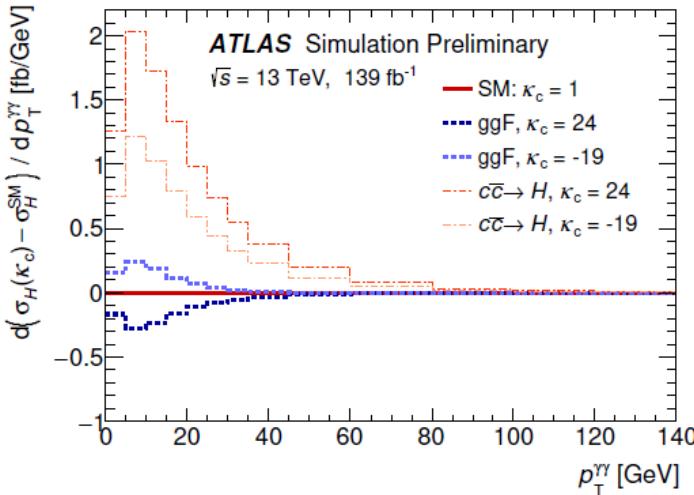


Figure 14: The modification of the $p_T^{\gamma\gamma}$ differential cross section for different values of κ_c , shown separately for the cross sections of the ggF and $c\bar{c} \rightarrow H$ production modes. As expected, for a given value of κ_c , the effect on the $c\bar{c} \rightarrow H$ production cross section is larger than that on the ggF production.

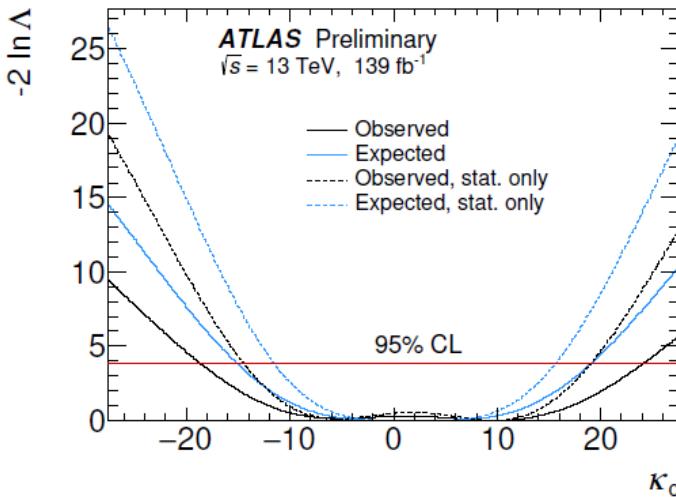


Figure 15: The profile likelihood ratio, λ , shown as a function of κ_c for the fit to the $p_T^{\gamma\gamma}$ distribution. The intersection of the $-2 \ln \Lambda$ curve with the horizontal line provides the 95% confidence intervals.

3 Search for a pair of BEH bosons constraint of the H self-coupling from H differential production and decay mesurements

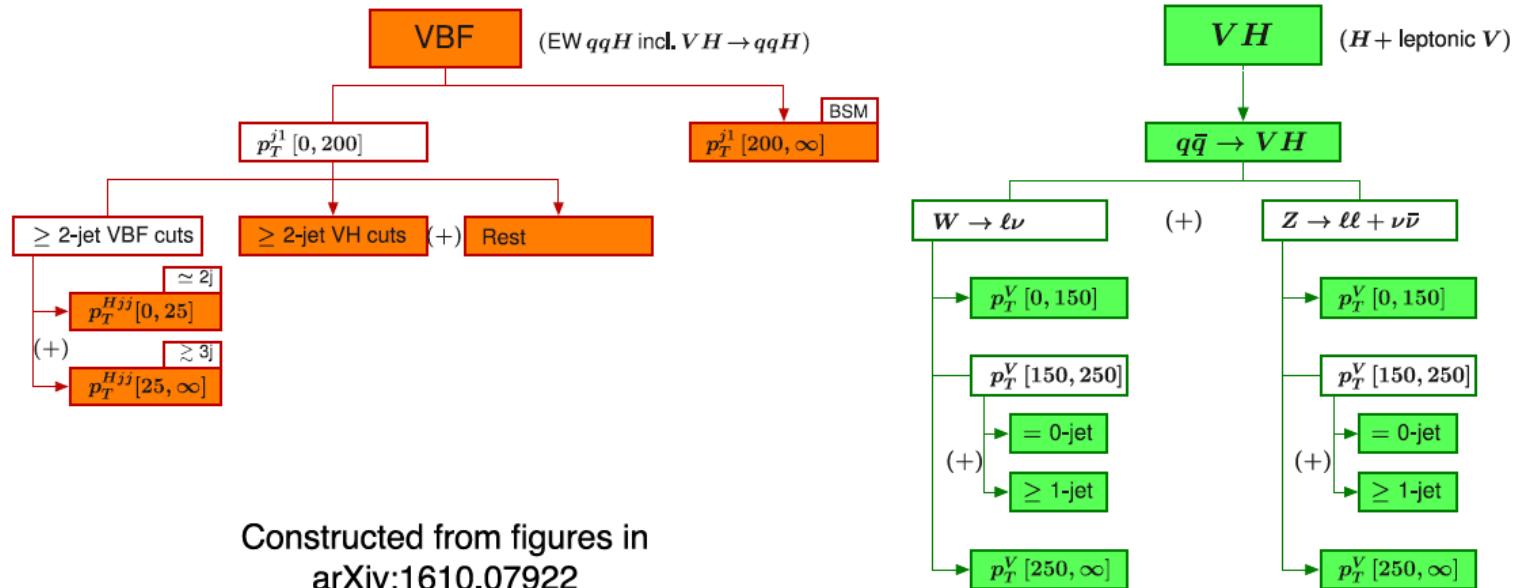


Figure 3: Schematic diagram of the $VBF + V(\text{had})H$ (left) and $V(\text{lep})H$ (right) STXS regions. p_T^{Hjj} is the p_T of the Higgs boson plus two jets system, p_T^V is the p_T of the vector boson V in the VH production mode, p_T^{j1} is the p_T of the jet with the highest p_T . In the $VH, H \rightarrow b\bar{b}$ analysis, the separation in jet number of the $p_T^V [150, 250]$ region in the VH production mode has been ignored, merging the 0 and the ≥ 1 jet regions. The diagrams are obtained from Ref. [14].

- ♪ Results from (run 1 and) run 2
 - * detector
 - * SM
 - * BSM
 - * (B-E)H
 - * **Vector-boson scattering**

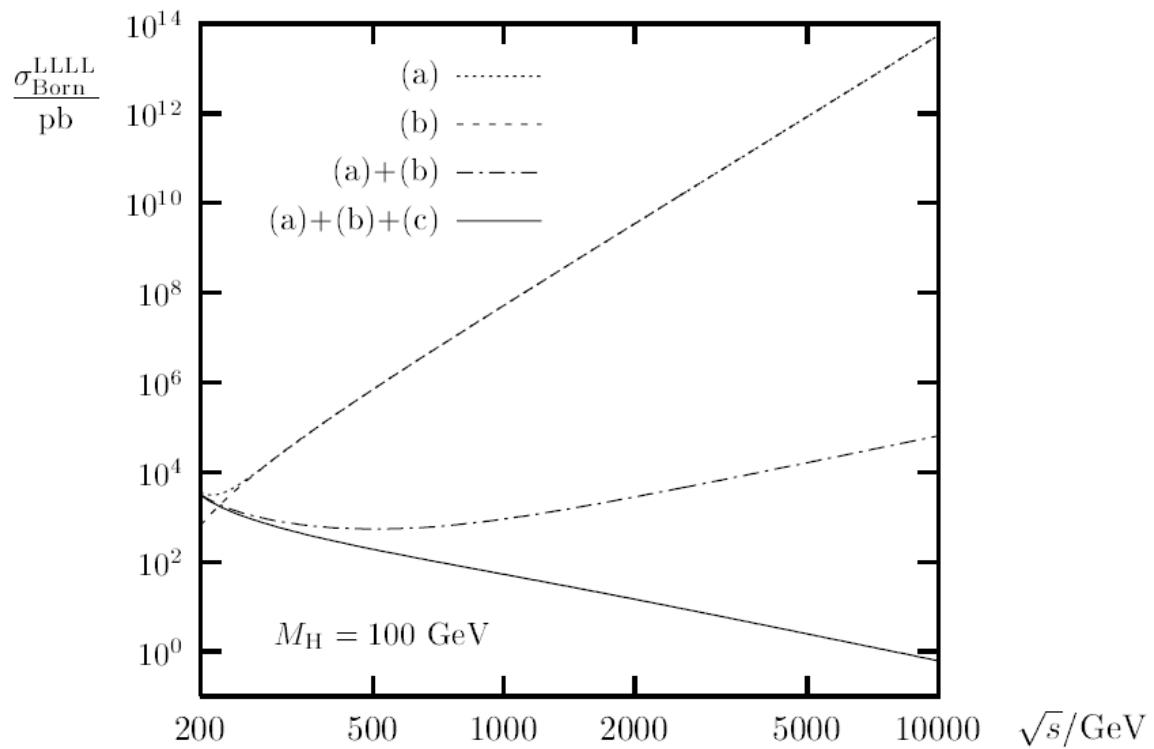


Figure 2: The cross-sections for longitudinal gauge-boson scattering resulting from subsets of the tree-level diagrams: (a) diagrams involving only three-gauge-boson couplings, (b) diagram involving only four-gauge-boson couplings, (c) diagrams involving Higgs bosons.

- ♪ *Historical introduction , Setting the stage*
- ♪ *Results from (run 1 and) run 2*
- ♪ ***Future of ATLAS , run 3 , HL-LHC***
- ♪ *Conclusions*
- ♪ *Backup*

- ♪ *Historical introduction , Setting the stage*
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