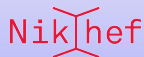


Exploring New Physics in $B \rightarrow \pi K$ Decays

Eleftheria Malami

Nikhef, Theory Group

National Institute for Subatomic Physics



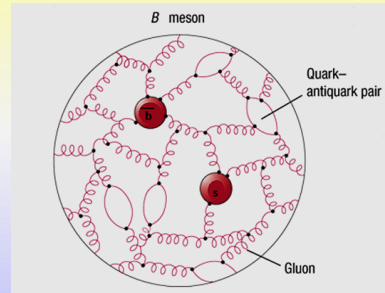
September 2, 2019

Based on:
arXiv:1712.02323 [hep-ph]
arXiv:1806.08783 [hep-ph]

Introduction

B-mesons

- Important system to test the flavour and CP-violating sector of the SM
- After BaBar, Belle, Tevatron, LHCb governs the experimental stage
- In the future:
Belle II will offer new opportunities along with the **LHCb upgrade**

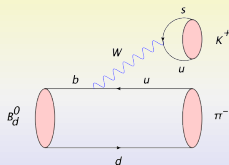


$B \rightarrow \pi K$ channels play a key role in these studies

Non-leptonic $B \rightarrow \pi K$ decays

$B \rightarrow \pi K$ Decays

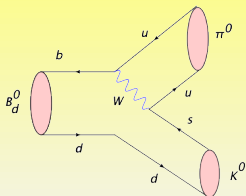
- Naively, we would assume tree contributions playing the leading role **but** they are strongly suppressed by the CKM matrix element $|V_{ub}|$



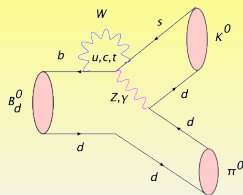
- Dominated by QCD penguin
- In the case of $B^+ \rightarrow \pi^0 K^+$ and $B_d^0 \rightarrow \pi^0 K^0$ colour allowed electro-weak penguins (EWP) enter at the same level as colour allowed trees, contributing $\mathcal{O}(10\%)$ to decay amplitudes

Illustrating contributions to $B_d^0 \rightarrow \pi^0 K_S$ as an example

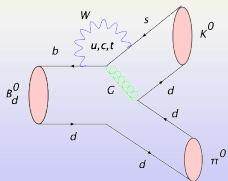
colour
suppressed
tree



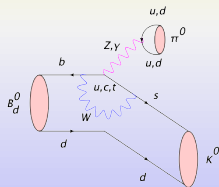
colour
suppressed
EWP



QCD
penguin



colour
allowed
EWP



The decay is:

- dominated by gluonic (QCD) penguins **BUT**
- electroweak penguins: also important

$B_d^0 \rightarrow \pi^0 K_S$ Decays

- New Physics (NP) may enter through EW penguins:
thus, very promising decay
- NP contributions: related to new CP violation sources
(probed through CP violating observables)
- significance of $B_d^0 \rightarrow \pi^0 K_S$:
the only $B \rightarrow \pi K$ mode exhibiting mixing-induced CP violation
outstanding role for testing SM with $B \rightarrow \pi K$ system

Amplitudes

- Non-leptonic decays: challenging due to hadronic matrix elements
- Flavour symmetries of strong interactions show connection between amplitudes of $B \rightarrow \pi K$ and $B \rightarrow \pi\pi$, $B \rightarrow KK$, from which
 - either hadronic amplitudes are eliminated or
 - determined from experimental data of the latter

In our analysis:

- strong interactions theoretical assumptions: as minimal as possible
- use results of QCDF to include SU(3)-breaking corrections

Hadronic Parameters

$$\begin{aligned}
 re^{i\delta} &= \left(\frac{\lambda^2 R_b}{1 - \lambda^2} \right) \left[\frac{T - (P_t - P_u)}{P_t - P_c} \right] \\
 r_c e^{i\delta_c} &= \left(\frac{\lambda^2 R_b}{1 - \lambda^2} \right) \left[\frac{T + C}{P_t - P_c} \right] \\
 \rho_c e^{i\theta_c} &= \left(\frac{\lambda^2 R_b}{1 - \lambda^2} \right) \left[\frac{P_t - P_u}{P_t - P_c} \right] \approx 0 \\
 \rho_n e^{i\theta_n} &= \left(\frac{\lambda^2 R_b}{1 - \lambda^2} \right) \left[\frac{C + (P_t - P_c)}{P_t - P_c} \right] = r_c e^{i\delta_c} - re^{i\delta}
 \end{aligned}$$

$$\begin{aligned}
 \lambda &= |V_{us}| \\
 R_b &= \left(1 - \frac{\lambda^2}{2} \right) \frac{1}{\lambda} \left| \frac{V_{ub}}{V_{cb}} \right|
 \end{aligned}$$

- $re^{i\delta}$ and $r_c e^{i\delta_c}$: non-perturbative (difficult to calculate)
- we calculate them using $B \rightarrow \pi\pi$ and SU(3) flavour symmetry:

$$r_c e^{i\delta_c} = (0.17 \pm 0.06) e^{i(1.9 \pm 23.9)^\circ}$$

$$re^{i\delta} = (0.09 \pm 0.03) e^{i(28.6 \pm 21.4)^\circ}$$

[A. J. Buras, R. Fleischer, S. Recksiegel, F. Schwab (2004)]

Electroweak Penguin Parameters q, ϕ

\Rightarrow Parametrization of EW penguin effects by:

$$qe^{i\phi}e^{i\omega} = \frac{\hat{P}_{EW} + \hat{P}_{EW}^C}{\hat{T} + \hat{C}}$$

ω : strong phase (quite small)

ϕ : CP- violating phase

\hat{T} : color-allowed tree contributions,

\hat{C} : colour-suppressed tree,

\hat{P}_{EW} : color-allowed EWP,

\hat{P}_{EW}^C : color-suppressed EWP

$$qe^{i\phi}e^{i\omega} = -\frac{-3}{2\lambda^2 R_b} \left[\frac{C_9(\mu) + C_{10}(\mu)}{C_1(\mu) + C_2(\mu)} \right] R_q = (0.68 \pm 0.05) R_q$$

$C_i(\mu)$: Wilson coefficients

- The ratio has been calculated for the SM values of $\phi = 0$
- Are there any deviations from the SM values?
- **Do the deviations indicate New Physics?**

CP Asymmetries in $B_d^0 \rightarrow \pi^0 K_S$

- The time-dependent CP asymmetry:

$$\begin{aligned}
 A_{CP}(t) &= \frac{\Gamma(\bar{B}^0(t) \rightarrow \pi^0 K_S) - \Gamma(B^0(t) \rightarrow \pi^0 K_S)}{\Gamma(\bar{B}^0(t) \rightarrow \pi^0 K_S) + \Gamma(B^0(t) \rightarrow \pi^0 K_S)} \\
 &= A_{\pi^0 K_S} \cos(\Delta M t) + S_{\pi^0 K_S} \sin(\Delta M t),
 \end{aligned}$$

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- The direct CP asymmetry:

$$A_{\pi^0 K_S} = \frac{|\bar{A}_{00}|^2 - |A_{00}|^2}{|\bar{A}_{00}|^2 + |A_{00}|^2}$$

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- The mixing-induced CP asymmetry:

(from interference between $B^0 - \bar{B}^0$ mixing and decay processes of B_d^0, \bar{B}_d^0 mesons into $\pi^0 K_S$ final state):

$$S_{\pi^0 K_S} = \frac{2|A_{00}\bar{A}_{00}|}{|\bar{A}_{00}|^2 + |A_{00}|^2} \sin(2\beta - 2\phi_{\pi^0 K_S})$$

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State-of-the-art analysis

- How do we find the angles $2\phi_{\pi^0 K_S} = \arg[\bar{A}_{00} A_{00}^*]$?
 - With the help of isospin relation

$$\begin{aligned}
 3A_{3/2} &= A(B^0 \rightarrow \pi^- K^+) + \sqrt{2}A(B^0 \rightarrow \pi^0 K^0) \\
 &= -(\hat{T} + \hat{C})(e^{i\gamma} - qe^{i\phi} e^{i\omega})
 \end{aligned}$$

we construct triangles in the complex plane

[Y. Nir, H. R. Quinn (1991);
M. Gronau, O. F. Hernandez,
D. London, J. L. Rosner
(1995)]

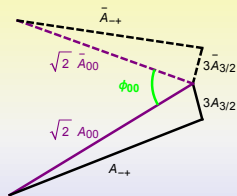
- We get 4 possible combinations between A_{00} and \bar{A}_{00} , thus 4 angles
- So, 4 cases for $S_{\pi^0 K_S} \rightarrow 4$ branches in the $S_{\pi^0 K_S} - A_{\pi^0 K_S}$ plot

Isospin Triangles in the complex plane

$$3A_{3/2} = A(B^0 \rightarrow \pi^- K^+) + \sqrt{2}A(B^0 \rightarrow \pi^0 K^0)$$

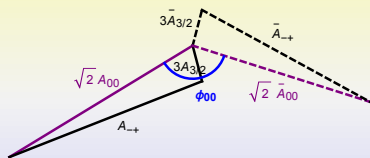
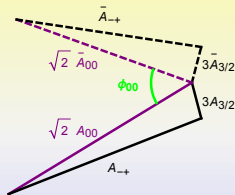
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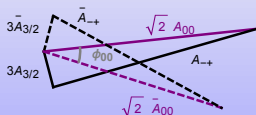
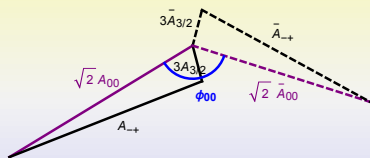
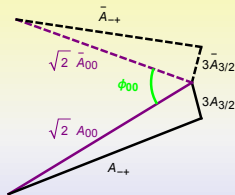
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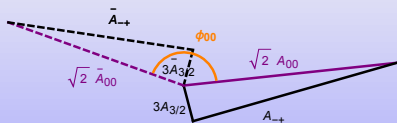
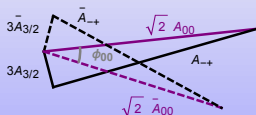
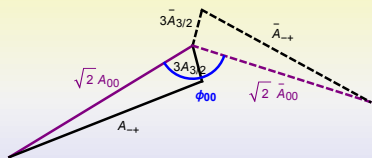
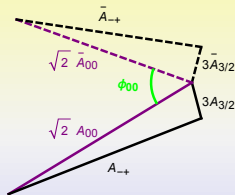
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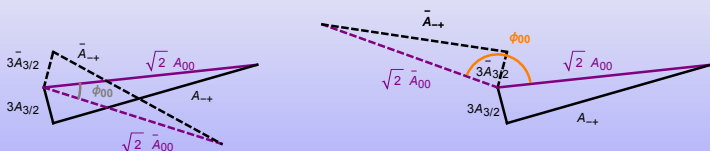
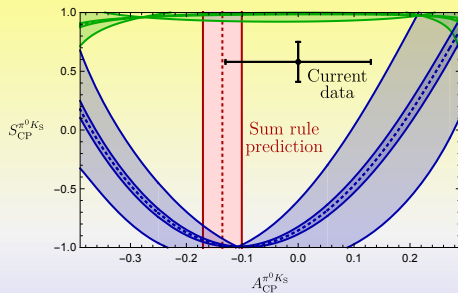


Isospin Triangles in the complex plane

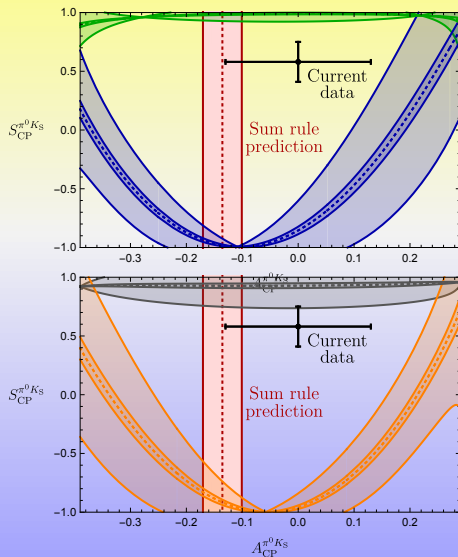
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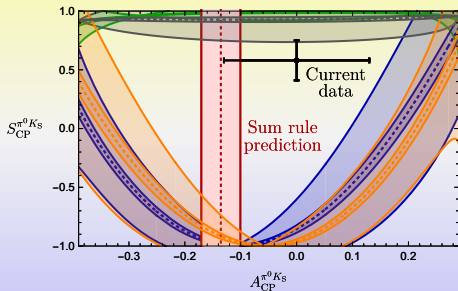
$S_{\pi^0 K_S} - A_{\pi^0 K_S}$ plot



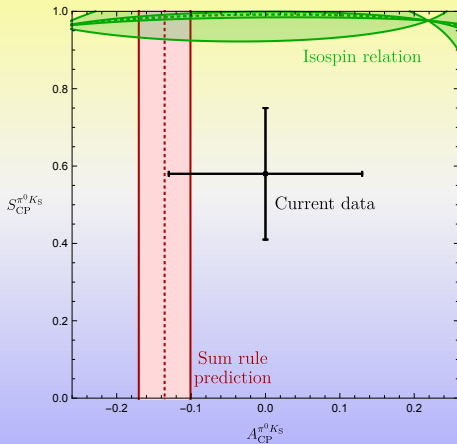
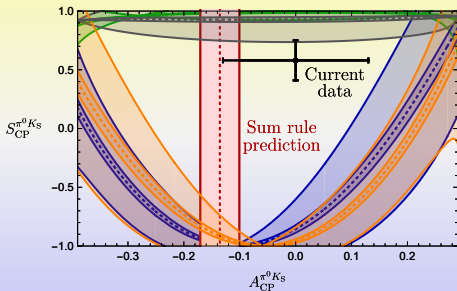
$S_{\pi^0 K_S} - A_{\pi^0 K_S}$ plot



$S_{\pi^0 K_S} - A_{\pi^0 K_S}$ plot



$S_{\pi^0 K_S} - A_{\pi^0 K_S}$ plot



Sum Rule

Using the information encoded in the CP-averaged branching ratios:

$$\begin{aligned} \Delta_{\text{SR}}^{(\text{I})} &= A_{\text{CP}}^{\pi^{\pm}K^{\mp}} + A_{\text{CP}}^{\pi^{\pm}K^0} \frac{\mathcal{B}(B^+ \rightarrow \pi^+ K^0) \tau_{B^0}}{\mathcal{B}(B_d^0 \rightarrow \pi^- K^+) \tau_{B^+}} \\ &\quad - A_{\text{CP}}^{\pi^0 K^{\pm}} \frac{2\mathcal{B}(B^+ \rightarrow \pi^0 K^+) \tau_{B^0}}{\mathcal{B}(B_d^0 \rightarrow \pi^- K^+) \tau_{B^+}} - A_{\text{CP}}^{\pi^0 K^0} \frac{2\mathcal{B}(B_d^0 \rightarrow \pi^0 K^0)}{\mathcal{B}(B_d^0 \rightarrow \pi^- K^+)} \\ &= 0 + \mathcal{O}(r_{(c)}^2, \rho_c^2), \end{aligned}$$

which offers an interesting test of the SM.

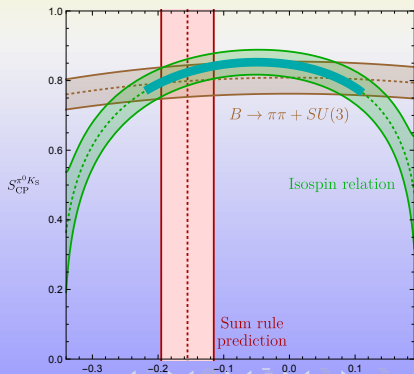
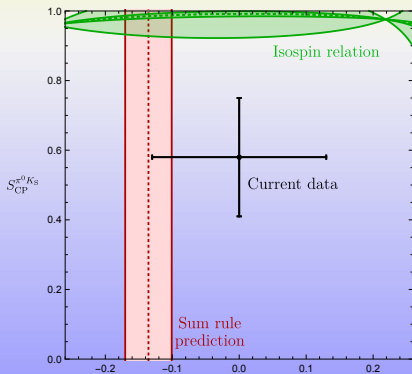
$$\Delta_{\text{SR}}^{(\text{I})}|_{\text{SM}} = -0.009 \pm 0.013$$

$$\Delta_{\text{SR}}^{(\text{I})}|_{\text{exp}} = -0.15 \pm 0.14$$

Can we resolve $B \rightarrow \pi K$ puzzle?

1 change of data?

- prime candidate: branching ratio (due to large experimental uncertainty on $B^0 \rightarrow \pi^0 K^0$)
- lowering branching ratio's central value (2.5σ) gives consistent picture with SM



Can we resolve $B \rightarrow \pi K$ puzzle?

2 effects of NP?

- very promising sector for NP signals: EW penguin sector
- affecting values of q , ϕ
- sensitivity to new CP violation sources
- ★ NP scenarios with extra Z' boson
 - links to anomalies in rare B-decays

Using charged $B \rightarrow \pi K$ decays

To determine q and ϕ , we make use of the **charged** $B \rightarrow \pi K$ decays

- only direct CP asymmetry (not mixing induced)
- following similar isospin analysis for the charged case

$$\begin{aligned} 3A_{3/2} &= A(B^+ \rightarrow \pi^+ K^0) + \sqrt{2}A(B^+ \rightarrow \pi^0 K^+) \\ &= (\hat{T} + \hat{C})(e^{i\gamma} - qe^{i\phi}e^{i\omega}) \end{aligned}$$

Important ratio of branching ratio:

$$R_c = 2 \left[\frac{Br(B \rightarrow \pi^0 K)}{Br(B \rightarrow \pi K^0)} \right] = 1.09 \pm 0.06$$

New strategy to determine q and ϕ

- Defining the angle $\Delta\phi_{3/2}$ and converting $A_{3/2}$ to the quantity N

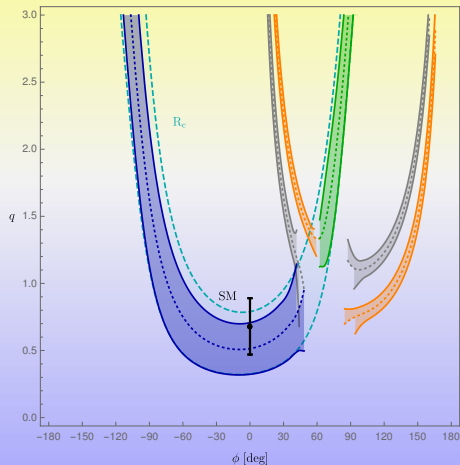
$$\Delta\phi_{3/2} = \phi_{3/2} - \bar{\phi}_{3/2}, \quad \sqrt{N} = 3 \left| \frac{A_{3/2}}{\hat{T} + \hat{C}} \right|$$

$$c = \pm\sqrt{N} \cos\left(\frac{\Delta\phi_{3/2}}{2}\right), \quad s = \pm\sqrt{N} \sin\left(\frac{\Delta\phi_{3/2}}{2}\right)$$

$$q = \pm\sqrt{N + 1 - 2c \cos\gamma - 2s \sin\gamma}$$

$$\tan\phi = \frac{\sin\gamma - s}{\cos\gamma - c}$$

The picture we get for the current data



four contours come
from isospin analysis

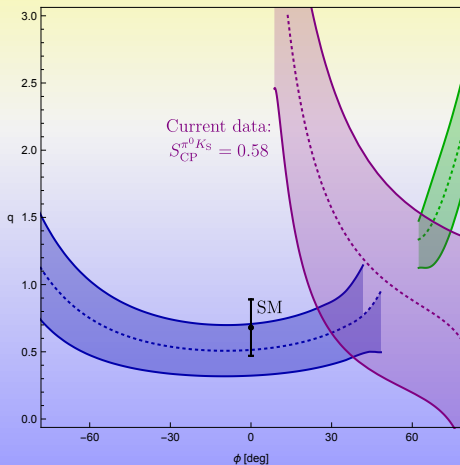
cyan dotted line:
complement analysis with R_c

$$R_c = 1 - 2r_c \cos \delta_c (\cos \gamma - q \cos \phi) + \mathcal{O}(r_c^2)$$

Utilizing the mixing induced CP violation in $B^0 \rightarrow \pi^0 K_S$

How can we get further info for the q, ϕ determination?

Using $S_{CP}^{\pi^0 K_S} \Rightarrow$ extraction of ϕ_{00} phase



$$\phi_{00} = (7.7 \pm 12.1)^\circ$$

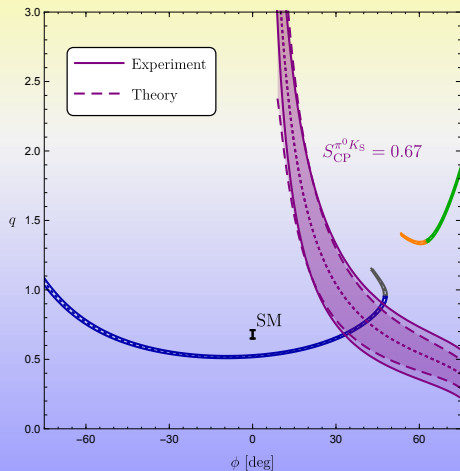
with hadronic par. values:

contributions from CS EW penguins
 $S_{CP}^{\pi^0 K_S}$ converts to purple contour

isospin analysis contours
 agree with R_c

Illustrating a future scenario

Demonstrating a future application of our strategy
considering a scenario for $S_{CP}^{\pi^0 K_S}$ measurements



Constraints from $S_{CP}^{\pi^0 K_S}$ and
isospin determination

Experimental precision and
theory
can be matched

Conclusions

To Sum Up...

Final Remarks

- $B \rightarrow \pi K$ decays: important to test SM and search for NP
- There is tension in the data \Rightarrow Something has to happen
 - either data should move to confirm the SM
 - or maybe there is NP
- We proposed a new strategy to determine q and ϕ
- Data from Belle II and LHCb upgrade will allow exciting new opportunities

Thank you!