

# Dark Matter and the Higgs

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**(LAPTh Annecy)**

**The Higgs portal to dark matter**  
**Invisible Higgs at the LHC**  
**Invisible Higgs in the future**  
**Comparison with astroparticle physics**  
**Heavy DM states**  
**Conclusion**

with Giorgio Arcadi and Martti Raidal, arXiv:1903.03616 (to appear in Phys. Rept.)

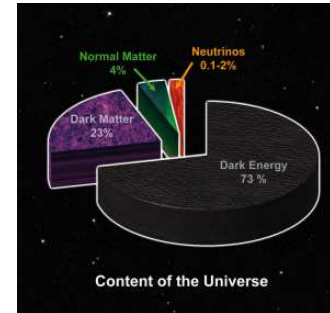
# The Higgs portal to dark matter

- DM indeed exists:
- Rotation curve of (clusters of) galaxies;
  - CMB and large scale structure formation.
- It is a new/unknown particle: no SM particle (eg  $\nu$ ) can do the job.

**A crucial hint for physics beyond the SM of particle physics!**

Makes: 25% of energy budget of Universe  
80% of total matter of Universe

**PLANCK:  $\Omega_{\text{DM}}h^2 = 0.119 \pm 0.001$**



- Neutral particle and hence dark (a charged particle will shine...).
- Cold or not too warm and hence non-relativistic ( $\nu$  do not work...).
- Very weakly interacting (reason why we did not observe it yet...).
- Stable or at least very long lived particle:  $\tau_{\text{DM}} \gg 10^{17}$  s.
- Possibly a relic from the early Universe: a window to the origin.

**In short, a WIMP. It worths the effort searching for it!**

# The Higgs portal to dark matter

WIMPS appear in models for naturalness/hierarchy problems:

- LSP: lightest neutralino of SUSY with  $R_p$  conserved;
- examples: • LKP: lightest Kaluza-Klein state in extra-dimensions;
- LTP: lightest T-odd state in Little Higgs/composite.

All particles not been observed (also their companions) at the LHC.

We choose a different way, simply Agnosticism and Occam razor,

postulate the existence of a weakly interacting massive particle:

- a singlet particle but of any spin i.e. a scalar, vector or fermion;
- QED neutral + isosinglet, no  $SU(2) \times U(1)$  charge: no Z couplings;
- $Z_2$  parity for stability: no couplings or mixing with fermions.

Hence, only couplings with the Higgs bosons  $\Rightarrow$  Higgs portal DM:

- annihilates into SM particles through s-channel Higgs exchange;
- interacts with fermionic matter only through Higgs exchange;
- can be produced in pairs via Higgs boson exchange or decays.

Again Occam razor: assume only the SM-like Higgs boson.

# The Higgs portal to dark matter

## Effective Lagrangians:

$$\Delta\mathcal{L}_s = -\frac{1}{2}M_s^2 S^2 - \frac{1}{4}\lambda_s S^4 - \frac{1}{4}\lambda_{Hss} \Phi^\dagger \Phi S^2$$

$$\Delta\mathcal{L}_v = \frac{1}{2}M_v^2 v_\mu v^\mu + \frac{1}{4}\lambda_v (v_\mu v^\mu)^2 + \frac{1}{4}\lambda_{Hvv} \Phi^\dagger \Phi v_\mu v^\mu$$

$$\Delta\mathcal{L}_\chi = -\frac{1}{2}M_\chi \bar{\chi} \chi - \frac{1}{4} \frac{\lambda_{H\chi\chi}}{\Lambda} \Phi^\dagger \Phi \bar{\chi} \chi$$

Mc Donald

Kanemura, ....

Lebedev, AD, ..

EWSB:  $\Phi \rightarrow \frac{1}{\sqrt{2}}(v + H)$  with  $v=246$  GeV and  $m_x^2 = M_x^2 + \frac{1}{4}\lambda_{Hxx} v^2 \dots$

Two free parameters: DM mass  $m_x$  and DM-Higgs coupling  $\lambda_{Hxx}$

**Of course effective and even non-renormalisable models. EFT for:**

- scalars: inert Higgs doublet models (comes with others states);
- vectors: may be related to a spontaneously broken U(1) symmetry;
- fermions: fourth family, singlet-doublet or vector-like leptons..

**All these are renormalisable, unitary, perturbative (up to cut-off) ...**

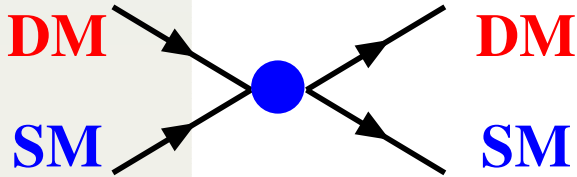
but where are all the new states gone? should have been seen?

- $Z_2$  symmetry: all new particles decay at end into the DM state;
- eventually some mass degeneracy with the invisible DM particle.

**$\Rightarrow$  missing energy and soft SM particle signature: very difficult!**

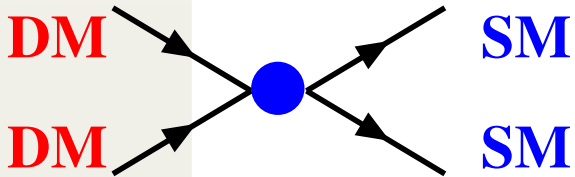
# The Higgs portal to dark matter

## Direct Detection:



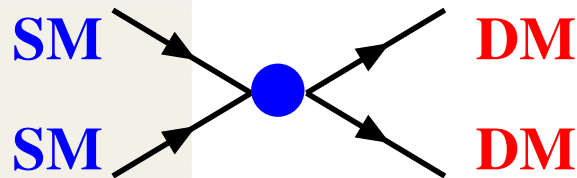
scattering on nucl. target:  
XENON  $\Rightarrow$  LZ, DARWIN

## Indirect Detection:

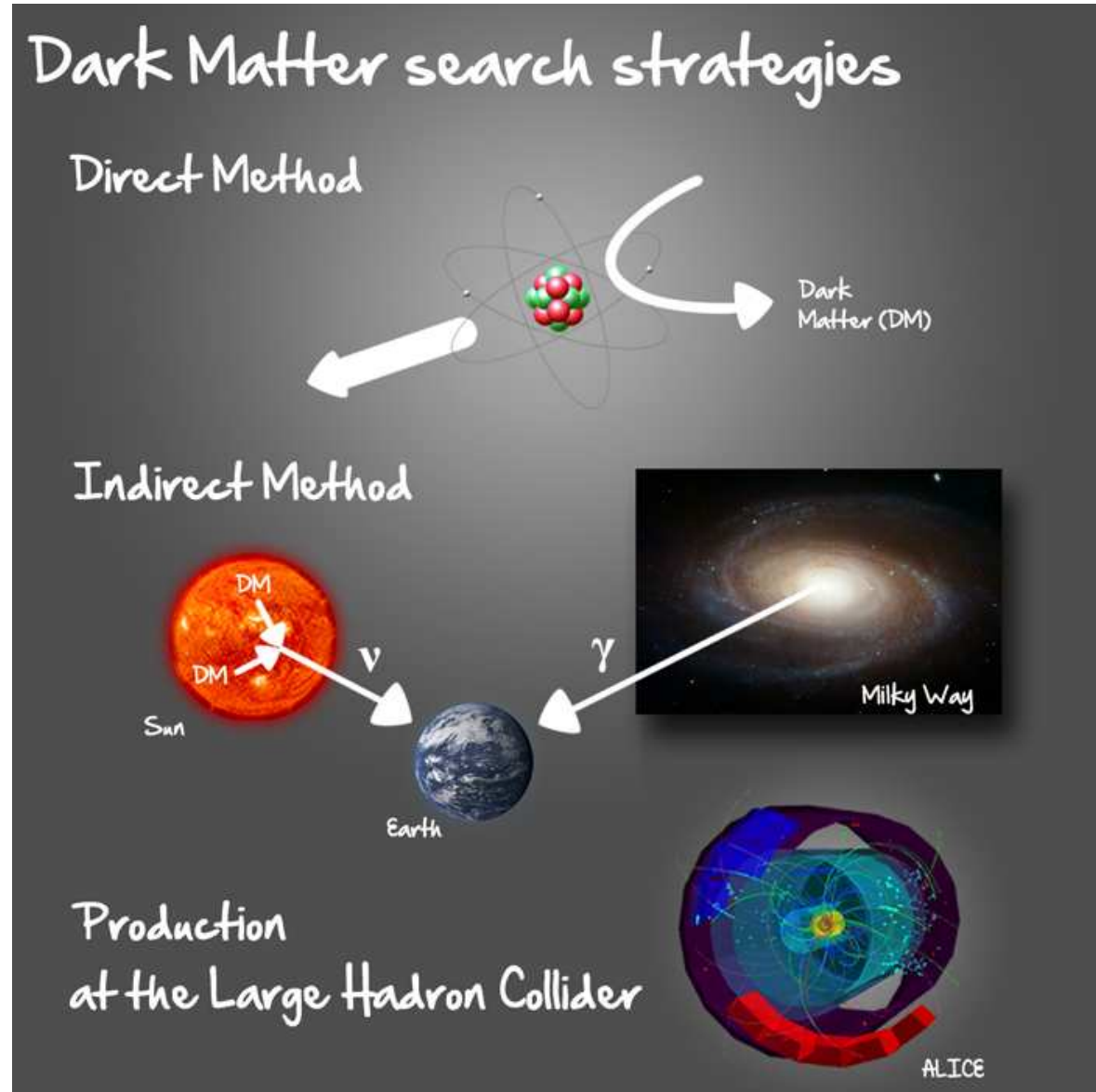


annihilation products:  $\gamma, \nu$   
HESS, Fermi  $\Rightarrow$  CTA, ...

## Detection at colliders:



missing energy signature  
LHC  $\Rightarrow$  HL-LHC,  $e^+e^-$ , pp



# Invisible Higgs at the LHC

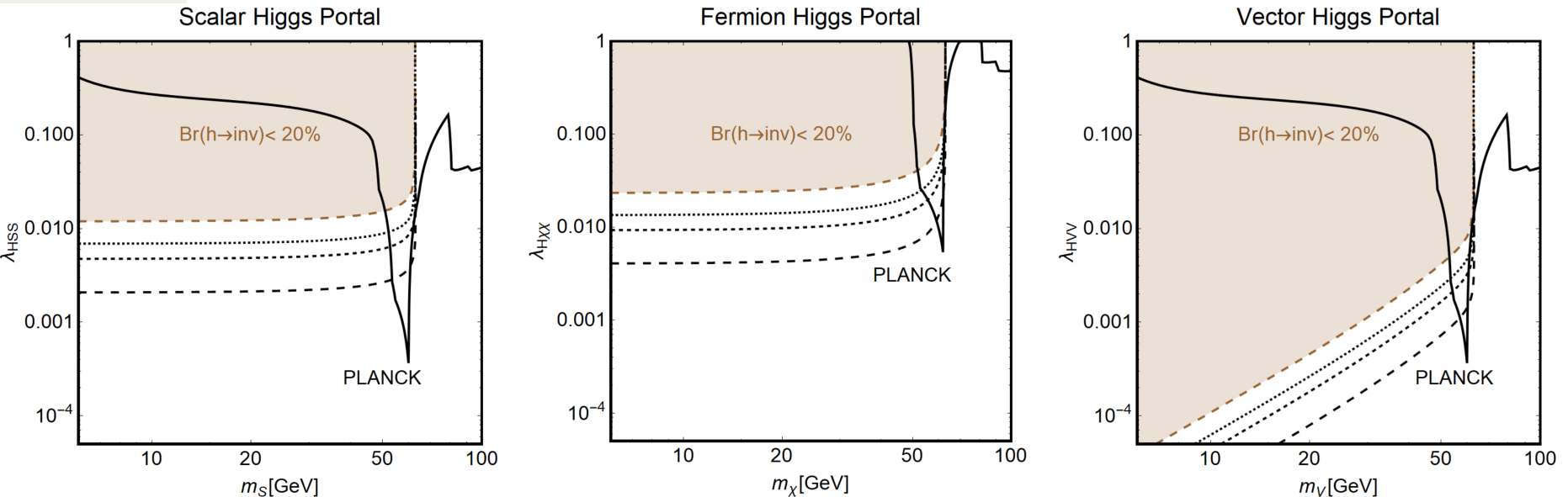
For light DM states, only possible handle at colliders is Higgs decays:

$$\Gamma_{\text{inv}}(\text{H} \rightarrow \text{SS}) = \frac{\lambda_{\text{Hss}}^2 v^2 \beta_s}{64\pi M_{\text{H}}}$$

$$\Gamma_{\text{inv}}(\text{H} \rightarrow \text{VV}) = \frac{\lambda_{\text{Hvv}}^2 v^2 M_{\text{H}}^3 \beta_v}{256\pi M_{\text{V}}^4} \left( 1 - 4 \frac{M_{\text{V}}^2}{M_{\text{H}}^2} + 12 \frac{M_{\text{V}}^4}{M_{\text{H}}^4} \right)$$

$$\Gamma_{\text{inv}}(\text{H} \rightarrow \text{ff}) = \frac{\lambda_{\text{Hff}}^2 v^2 M_{\text{H}} \beta_f^3}{32\pi \Lambda^2}$$

Possible only for  $m_{\text{X}} < \frac{1}{2} M_{\text{H}} \approx 62 \text{ GeV}$ ; depends on  $m_{\text{X}}, \lambda_{\text{HXX}}$ :



One has to check also the relic density/Planck: only one input?  
 maybe no, X not all DM and/or  $\Omega h^2$  obtained via other means...

# Invisible Higgs at the LHC

**Direct: measurement of total Higgs decay width via interference:**

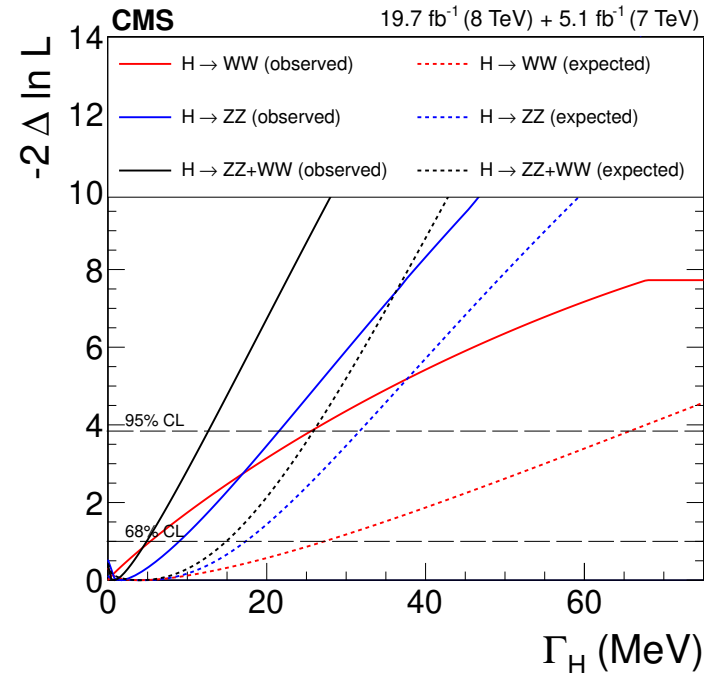
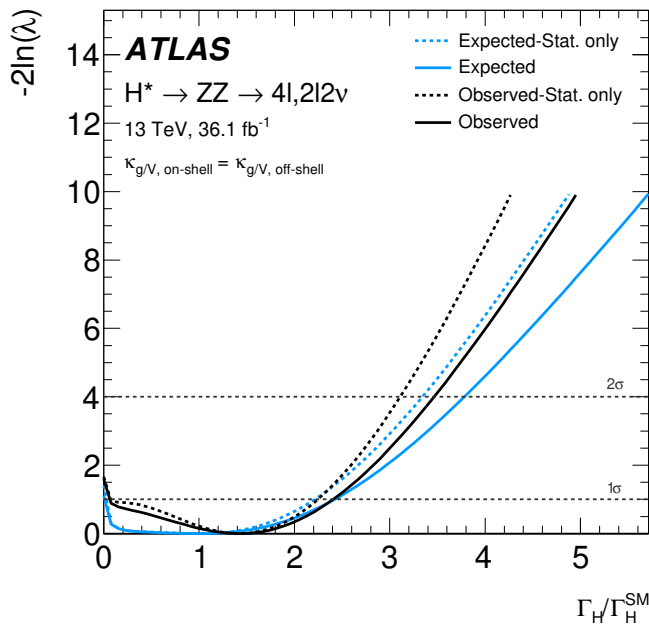
$\Gamma_H^{\text{SM}} = 4.07 \text{ MeV} \Rightarrow$  too small to be resolved experimentally.

If  $M_H \gtrsim 200 \text{ GeV}$ ,  $\Gamma_H > 1 \text{ GeV} \Rightarrow$  possible in  $H \rightarrow ZZ \rightarrow 4l$ .

But in  $pp \rightarrow H \rightarrow ZZ \rightarrow 4l$ , about 20% are for  $M_{4l} \gtrsim 2M_Z$ .

**In fact:**  $\sigma_{\text{gg} \rightarrow \text{H} \rightarrow 4l}^{\text{on-shell}} \propto g_{\text{ggH}}^2$ ,  $\sigma_{\text{gg} \rightarrow \text{H} \rightarrow 4l}^{\text{off-shell}} \propto g_{\text{ggH}}^2 \Gamma_H \Rightarrow \text{interf} \propto g_{\text{ggH}} \sqrt{\Gamma_H}$

**Direct measurement of  $\Gamma_H$  via interference with  $pp \rightarrow ZZ$  continuum:**



**The constraints are starting to be serious:  $\Delta\Gamma_H/\Gamma_H^{\text{SM}} \lesssim 1.25 @ 95\% \text{ CL!}$**

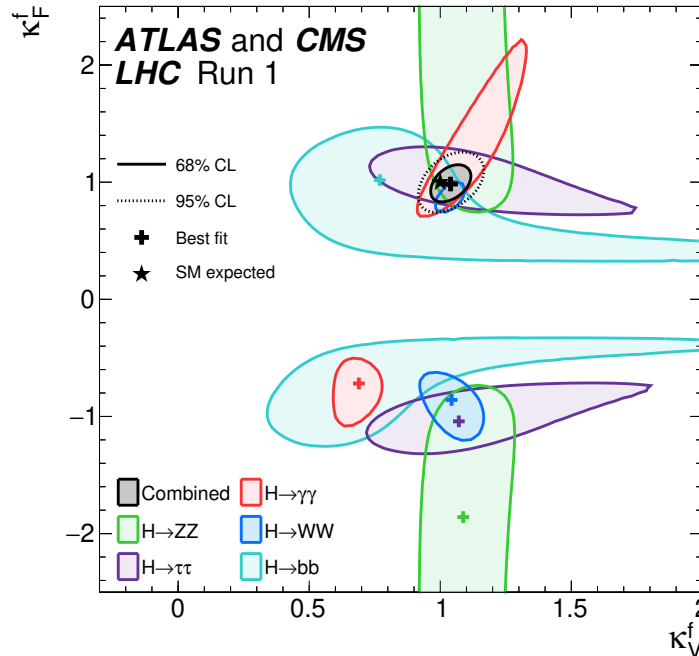
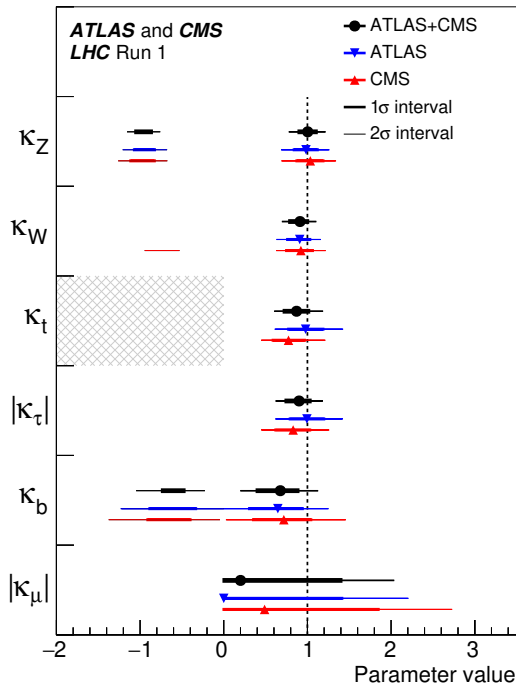
# Invisible Higgs at the LHC

**B) Indirectly: measurements of the Higgs decay branching ratios:**

$$\kappa_{\mathbf{x}}^2 = \sigma(\mathbf{x})/\sigma(\mathbf{x})|_{\text{SM}} = \Gamma(\mathbf{xx})/\Gamma(\mathbf{xx})|_{\text{SM}} = \mathbf{g}_{\text{Hxx}}^2/\mathbf{g}_{\text{Hxx}}^2|_{\text{SM}}$$

$$\Gamma(\mathbf{vv}) \rightarrow \kappa_{\mathbf{v}}^2, \quad \Gamma(\mathbf{ff}) \rightarrow \kappa_{\mathbf{f}}^2, \quad \sigma(\mathbf{vH}) \rightarrow \kappa_{\mathbf{v}}^2, \quad \sigma(\mathbf{VBF}) \rightarrow 0.74\kappa_{\mathbf{W}}^2 + 0.26\kappa_{\mathbf{Z}}^2$$

$$\sigma(\mathbf{\gamma\gamma}) \rightarrow \kappa_{\mathbf{\gamma}}^2 = 1.5\kappa_{\mathbf{W}}^2 + 0.1\kappa_{\mathbf{t}}^2 - 0.7\kappa_{\mathbf{t}}\kappa_{\mathbf{W}}, \quad \sigma(\mathbf{ggH}) \rightarrow \kappa_{\mathbf{g}}^2 = 1.06\kappa_{\mathbf{t}}^2 - 0.07\kappa_{\mathbf{t}}\kappa_{\mathbf{b}}$$



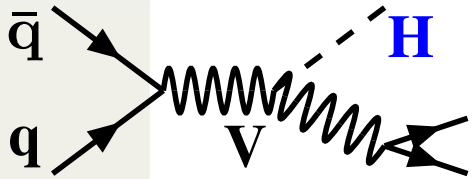
$$\kappa_{\mathbf{H}}^2 = 0.57\kappa_{\mathbf{b}}^2 + 0.22\kappa_{\mathbf{W}}^2 + 0.06\kappa_{\mathbf{\tau}}^2 + 0.03\kappa_{\mathbf{Z}}^2 + 0.03\kappa_{\mathbf{c}}^2 + 0.0023(\kappa_{\mathbf{\gamma}}^2 + \kappa_{\mathbf{Z}\gamma}^2)$$

global ATLAS+CMS fit gives  $\text{BR}(\text{H} \rightarrow \text{invisible}) \lesssim 20\% @ 95\% \text{CL}$



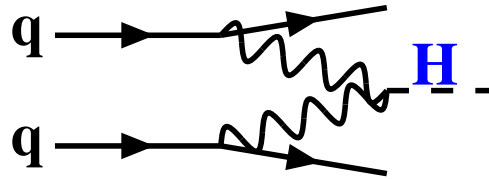
# Invisible Higgs at the LHC

**C) Even more direct: search for Higgs decaying invisibly and  $E_{\cancel{T}}$ :**



$$\begin{aligned} q\bar{q} &\rightarrow WH \rightarrow l\nu + E_{\cancel{T}} \\ q\bar{q} &\rightarrow ZH \rightarrow ll + E_{\cancel{T}} \end{aligned}$$

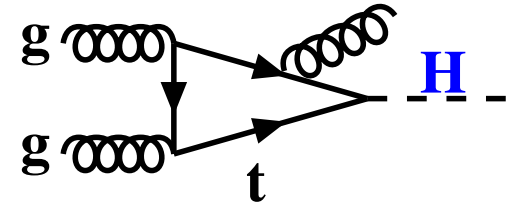
Choudhury+Roy, ...



$$qq \rightarrow qqH \rightarrow jj + E_{\cancel{T}}$$

high-mass,  $p_T$ ,  $\eta$  jets

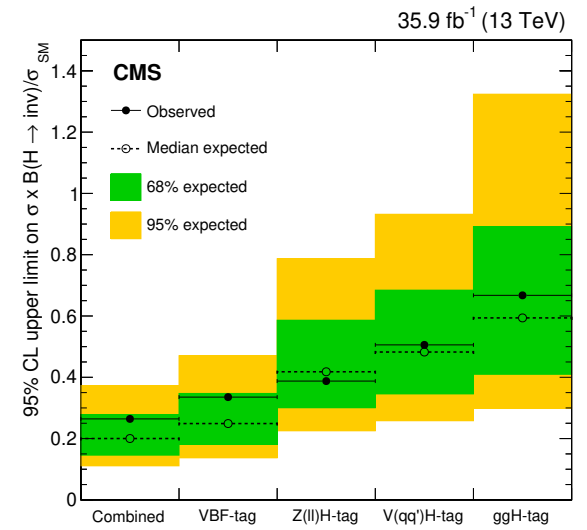
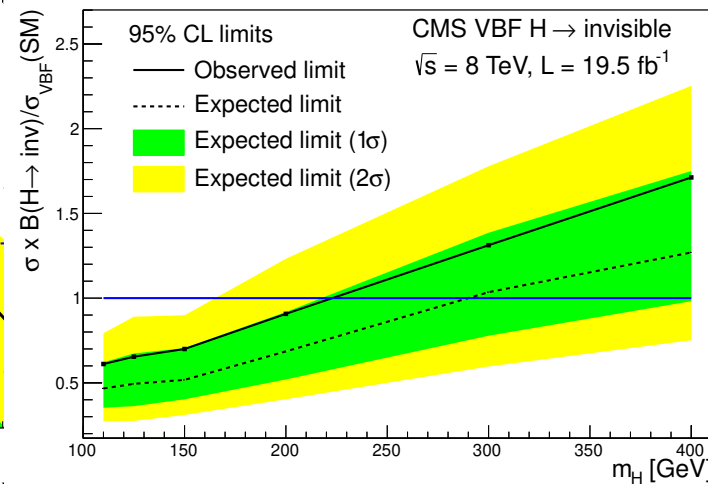
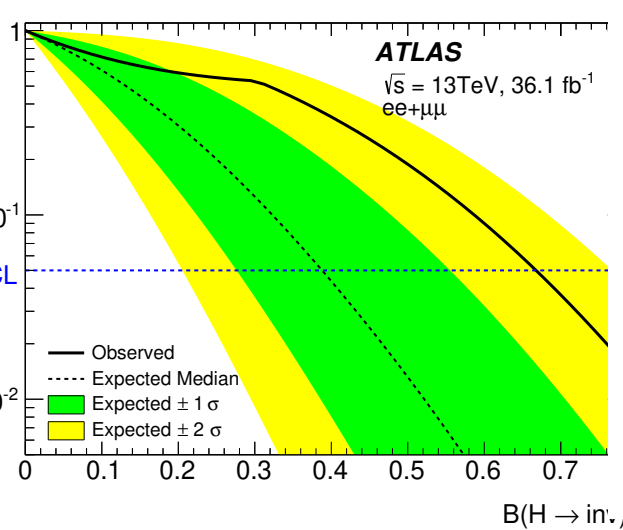
Eboli+Zeppenfeld



$$gg \rightarrow Hg \rightarrow j + E_{\cancel{T}}$$

also 2j, high rate.

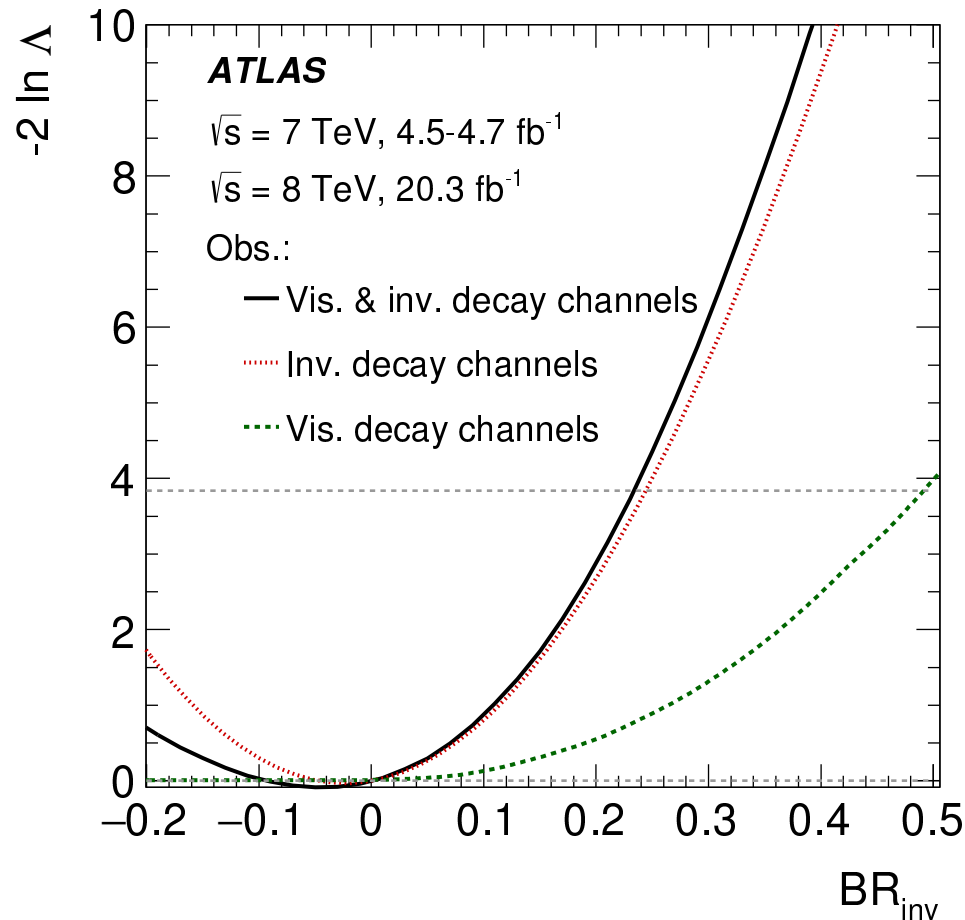
AD,Falkowski,Mambrini...



are also ATLAS+CMS combo gives  $\text{BR}(H \rightarrow \text{invisible}) \lesssim 20\% @ 95\% \text{ CL}$

# Invisible Higgs at the LHC

**D) A combination of all these channels/possibilities:**



**All in all at RunI+RunII LHC:  $BR(H \rightarrow \text{inv}) < 20\%$**

**This not very constraining after all. Can all this can be improved?**

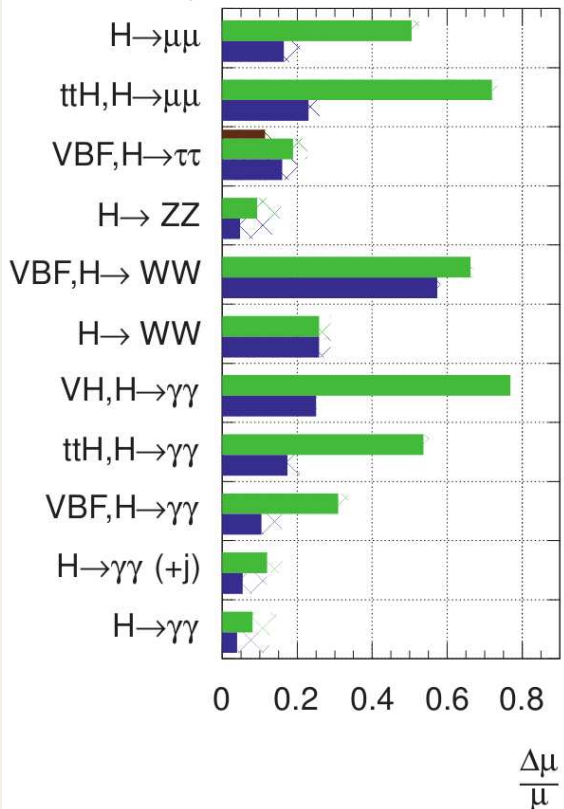
# Invisible Higgs in the future

**High-Luminosity LHC:  $\sqrt{s} = 14$  TeV and  $\mathcal{L} = 3 \text{ ab}^{-1}$ .**

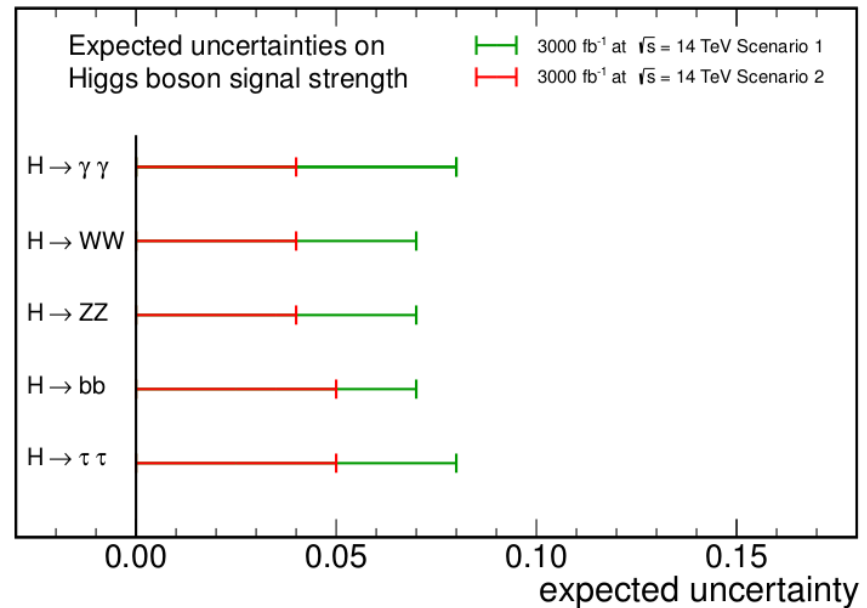
ATLAS Simulation

$\sqrt{s} = 14$  TeV:  $\int L dt = 300 \text{ fb}^{-1}$ ;  $\int L dt = 3000 \text{ fb}^{-1}$

$\int L dt = 300 \text{ fb}^{-1}$  extrapolated from 7+8 TeV



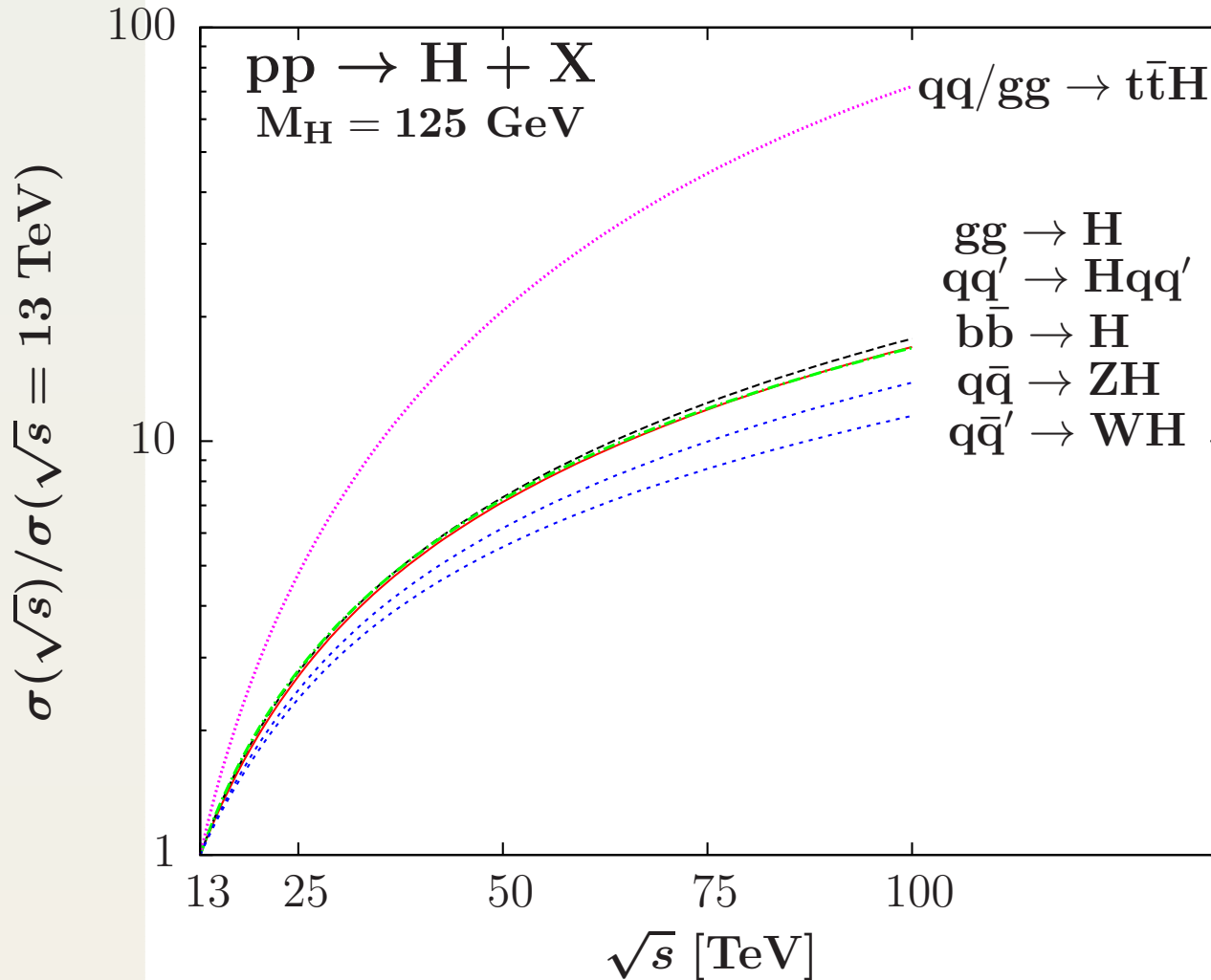
CMS Projection



**Improve sensitivity by factor two from BRs:  $\text{BR}(H \rightarrow \text{invisible}) \lesssim 10\%$   
 Large improvement in direct searches not expected; tougher conditions?)**

# Invisible Higgs in the future

higher energy pp colliders  $\Rightarrow \sqrt{s} = 28 - 100$  TeV and  $\mathcal{L} = \text{few ab}^{-1}$

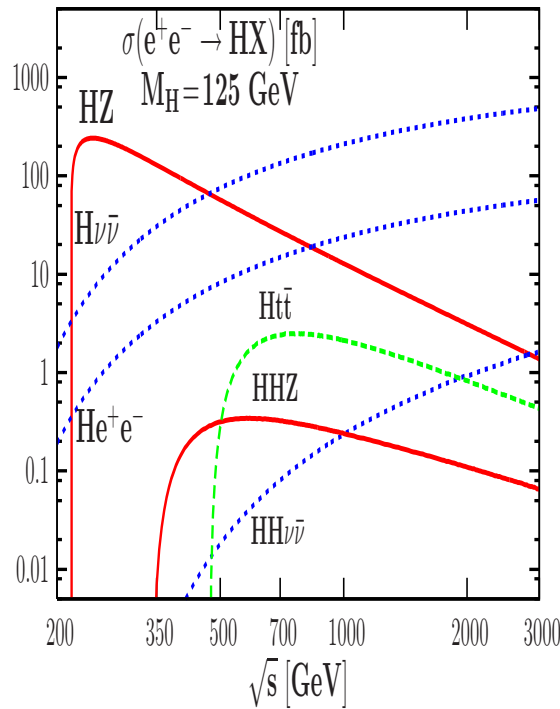
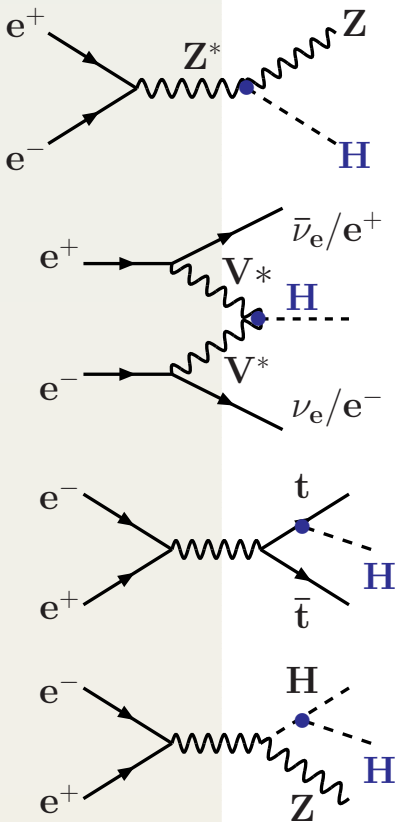


Baglio, ...

prove the sensitivity by another factor of two:  $BR(H \rightarrow \text{invisible}) \lesssim 5\%$

# Invisible Higgs in the future

**Best deal: an  $e^+e^-$  (non ILC?) machine:  $\sqrt{s} = 240\text{GeV}$  and  $\mathcal{L} = 1\text{ab}^{-1}$**



**Very precise measurements  
mostly at  $\sqrt{s} \approx 250\text{ GeV}$   
and mainly in  $e^+e^- \rightarrow ZH$**

$g_{HZZ}$	$\pm 0.01$
$g_{HWW}$	$\pm 0.01$
$g_{Hbb}$	$\pm 0.01$
$g_{Hcc}$	$\pm 0.03$
$g_{H\tau\tau}$	$\pm 0.03$
$g_{Htt}$	$\pm 0.05$
$\lambda_{HHH}$	$\pm 0.2$
$M_H$	$\pm 0.0004$
CP	$\pm 0.03$
$\Gamma_H$	$\pm 0.05$

**Improve the sensitivity very seriously:  $\text{BR}(H \rightarrow \text{invisible}) \lesssim 1\%$**

**In fact: a  $\text{BR}(H \rightarrow \text{inv})$  of 1% can be observed at the  $5\sigma$  level,  
and a  $\text{BR}(H \rightarrow \text{inv})$  of 5% can be measured with a 10% accuracy...**

# Comparison with astroparticle experiments

- Relic density  $\propto 1/\langle\sigma(\mathbf{x}\mathbf{x} \rightarrow \mathbf{H} \rightarrow \mathbf{f}\bar{\mathbf{f}})\mathbf{v}_r\rangle$  annihilation rate.

$$\langle\sigma_{\text{ferm}}^{\mathbf{x}} \mathbf{v}_r\rangle = \frac{\lambda_{\mathbf{H}\mathbf{x}\mathbf{x}}^2 m_{\text{ferm}}^2}{16\pi} \frac{1}{(4m_{\mathbf{x}}^2 - M_{\mathbf{H}}^2)^2} \delta_{\mathbf{x}}, \quad \delta_{\mathbf{s}} = 1, \delta_{\mathbf{v}} = \frac{1}{3}, \delta_{\mathbf{f}} = \frac{1}{2} \frac{v_r^2}{\Lambda^2}$$

In principle needs to fit the Planck value:  $\Omega_{\text{DM}} h^2 = 0.119 \pm 0.001$

- Spin-independent direct detection, simple for s, v, f DM states:

$$\sigma_{\mathbf{x}-\mathbf{N}}^{\text{SI}} = \frac{\lambda_{\mathbf{H}\mathbf{x}\mathbf{x}}^2}{16\pi M_{\mathbf{H}}^4} \frac{m_{\mathbf{N}}^4 f_{\mathbf{N}}^2}{(m_{\mathbf{x}} + m_{\mathbf{N}})^2} \delta'_{\mathbf{x}}, \quad \delta_{\mathbf{s}} = \delta_{\mathbf{v}} = 1, \delta_{\mathbf{f}} = \frac{4}{\Lambda^2}$$

$$\text{BR}_{\mathbf{x}}^{\text{inv}} \equiv \text{BR}(\mathbf{H} \rightarrow \text{inv}) = \frac{\Gamma(\mathbf{H} \rightarrow \mathbf{x}\mathbf{x})}{\Gamma_{\mathbf{H}}^{\text{SM}} + \Gamma(\mathbf{H} \rightarrow \mathbf{x}\mathbf{x})} = \frac{\sigma_{\mathbf{x}\mathbf{p}}^{\text{SI}}}{\Gamma_{\mathbf{H}}^{\text{tot}}/r_{\mathbf{x}} + \sigma_{\mathbf{x}\mathbf{p}}^{\text{SI}}}$$

**Direct detection:**

$$\text{BR}_{\mathbf{x}}^{\text{inv}} \simeq (\sigma_{\mathbf{x}\mathbf{p}}^{\text{SI}}/10^{-9} \text{pb}) [\delta_{\mathbf{x}} + (\sigma_{\mathbf{x}\mathbf{p}}^{\text{SI}}/10^{-9} \text{pb})]^{-1}$$

$$\delta_{\mathbf{s}} = 400 (10 \text{ GeV}/m_{\mathbf{s}})^2, \quad \delta_{\mathbf{v}} = 4 \times 10^{-2} (m_{\mathbf{v}}/10 \text{ GeV})^2, \quad \delta_{\mathbf{f}} = 3.5$$

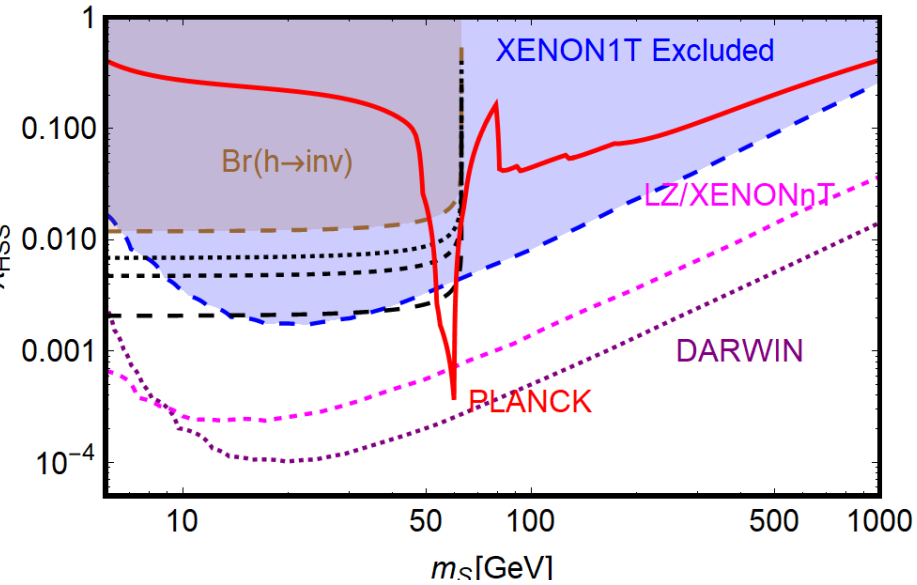
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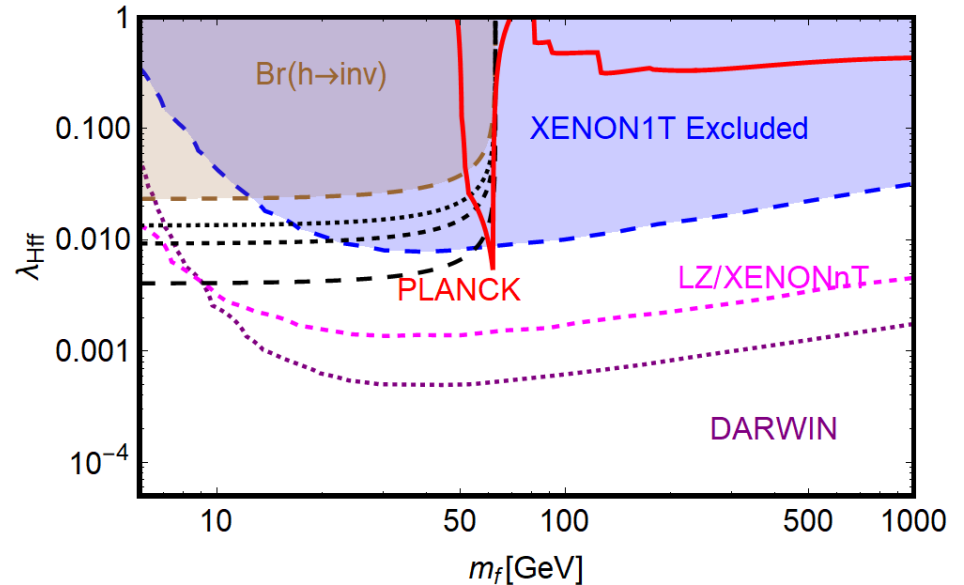
$$\delta_{\mathbf{s}} = 2.4 \times 10^{-2}, \quad \delta_{\mathbf{v}} = 1.3 \times 10^{-10} (m_{\mathbf{v}}/\text{GeV})^4, \quad \delta_{\mathbf{f}} = 3.9 \times 10^{-13} (m_{\mathbf{f}}/\text{GeV})^2$$

# Comparison with astroparticle experiments

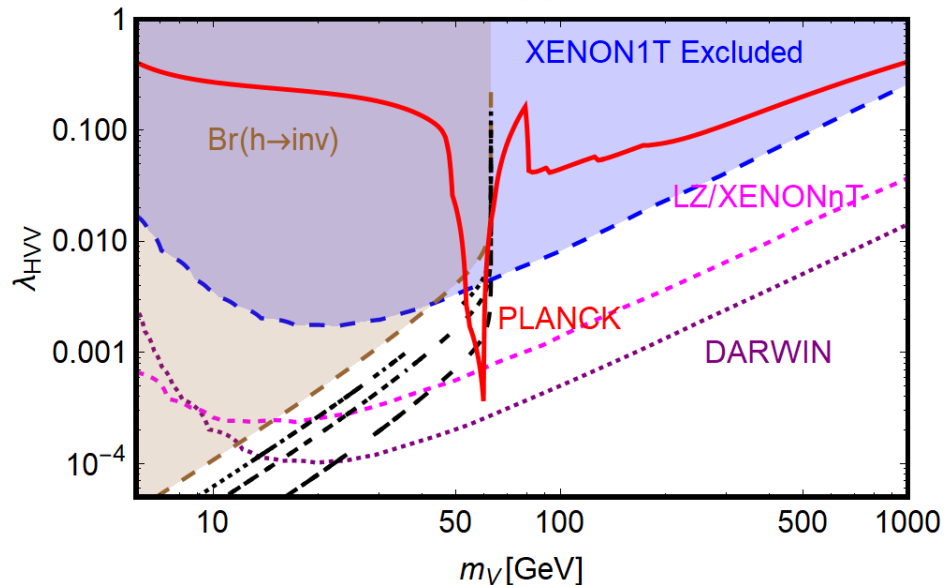
Scalar Higgs Portal



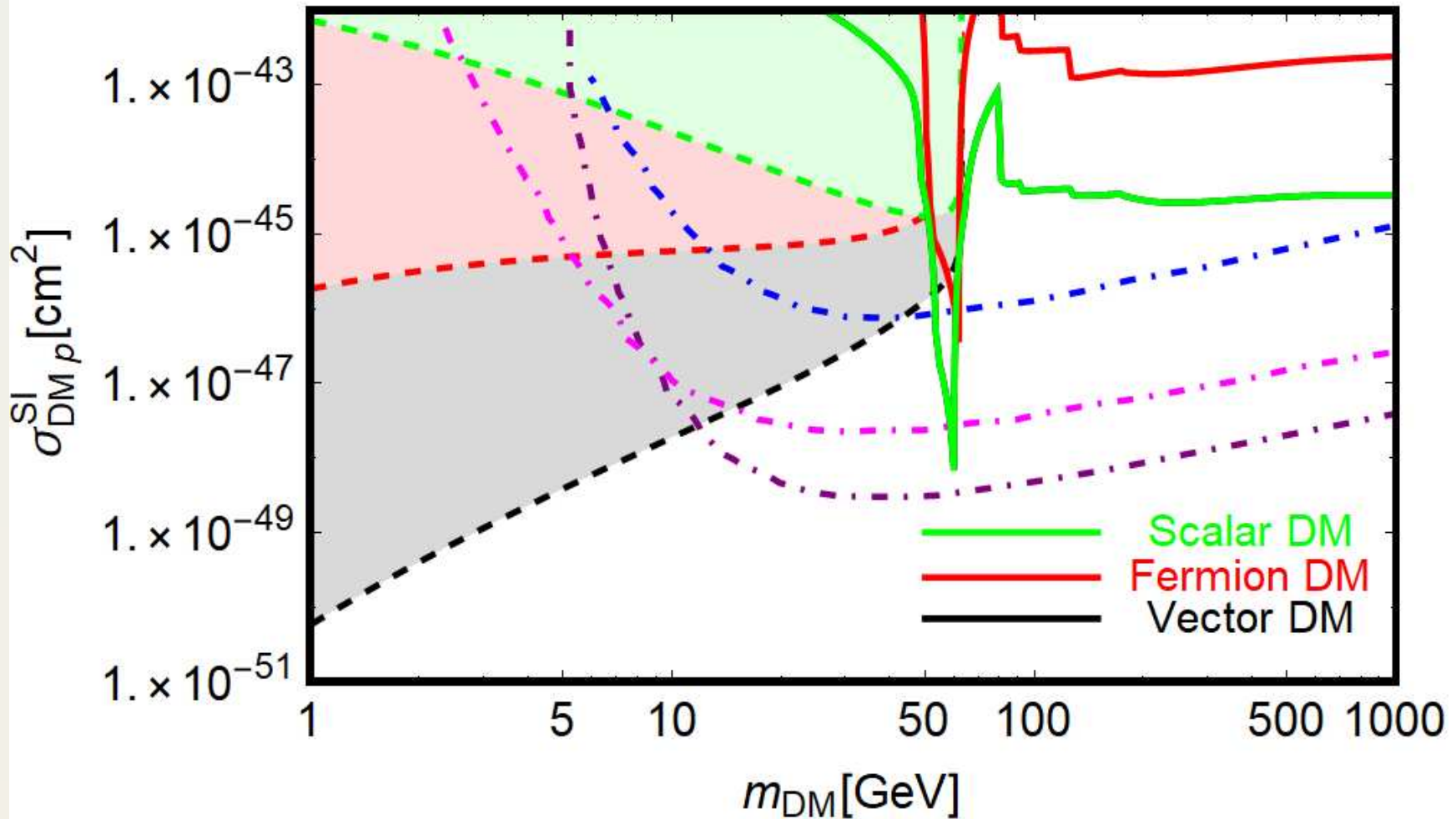
Fermion Higgs Portal



Vector Higgs Portal



# Comparison with astroparticle experiments



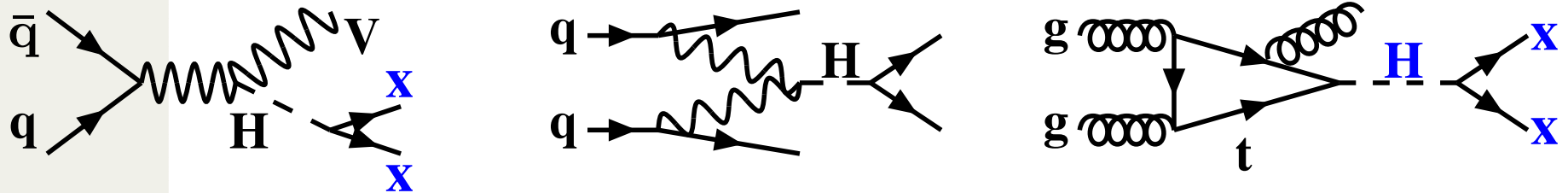
Arcadi, AD  
Raidal



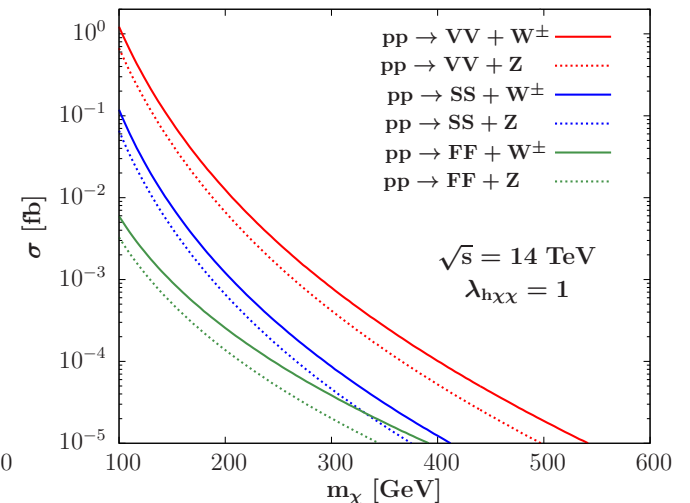
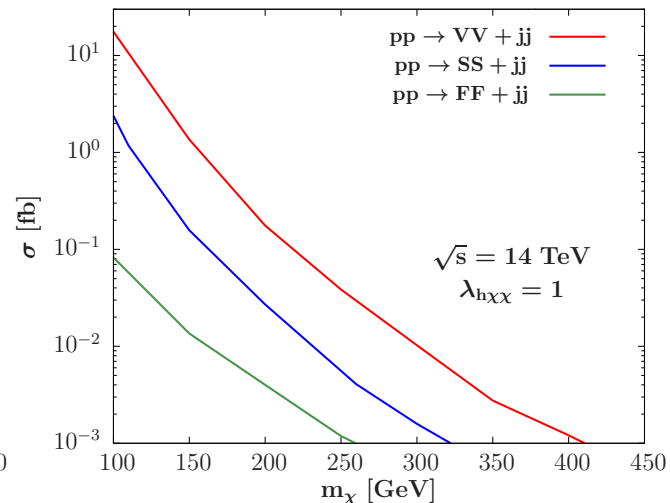
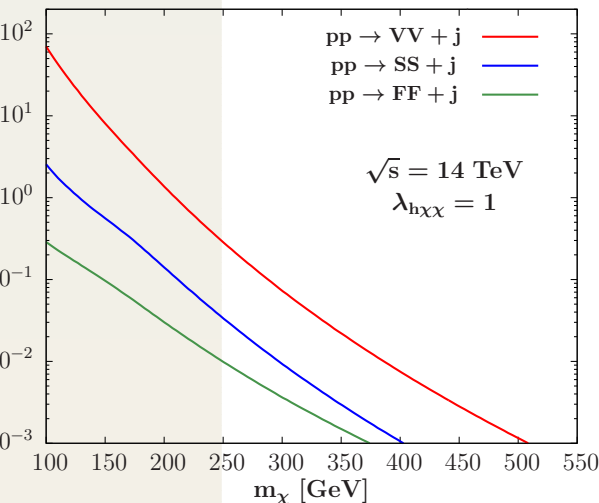
# Heavy DM states

Invisible decays work only for light DM; what if  $m_{\text{DM}} \gtrsim \frac{1}{2}M_{\text{H}}$ ?

Only way: produce them in pairs in continuum. At pp colliders:



- Exactly same channels as before but with an off-shell Higgs boson.
- Suppressed by the Higgs virtuality and the small couplings to DM.
- Needs high energies, very high luminosities and some real efforts...

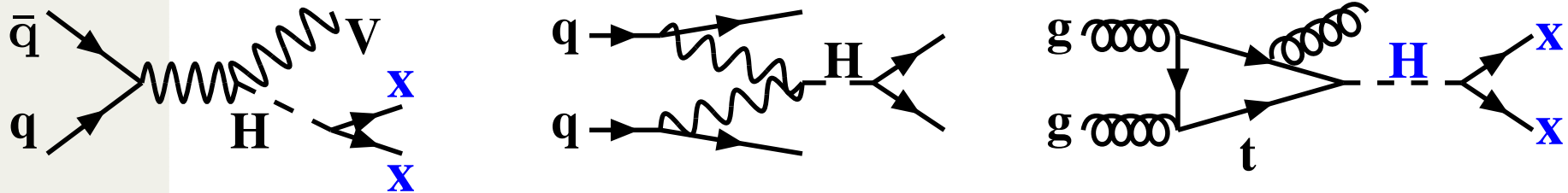


Quevillon, ...

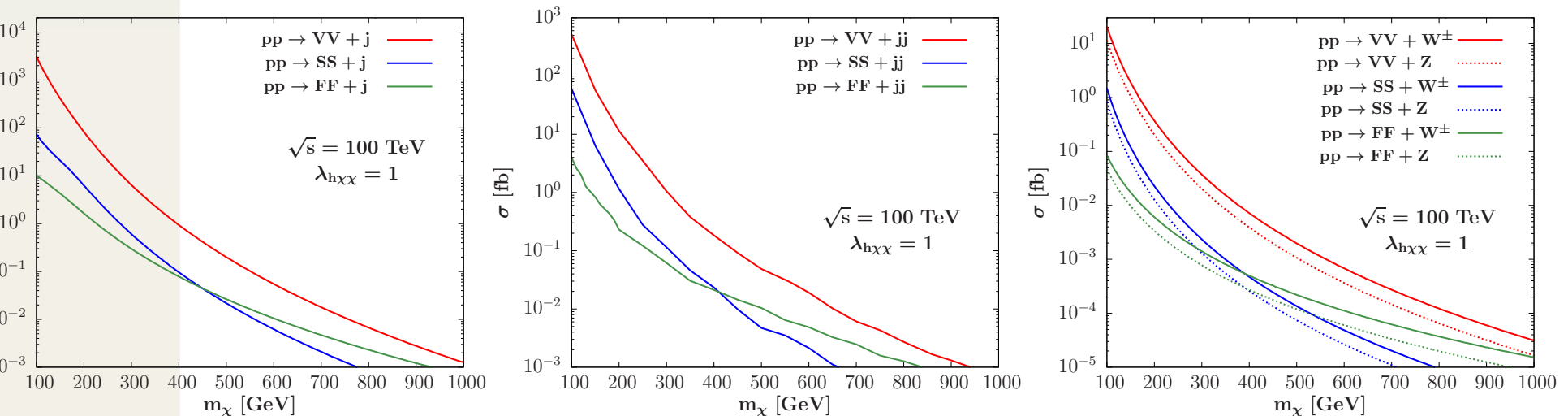
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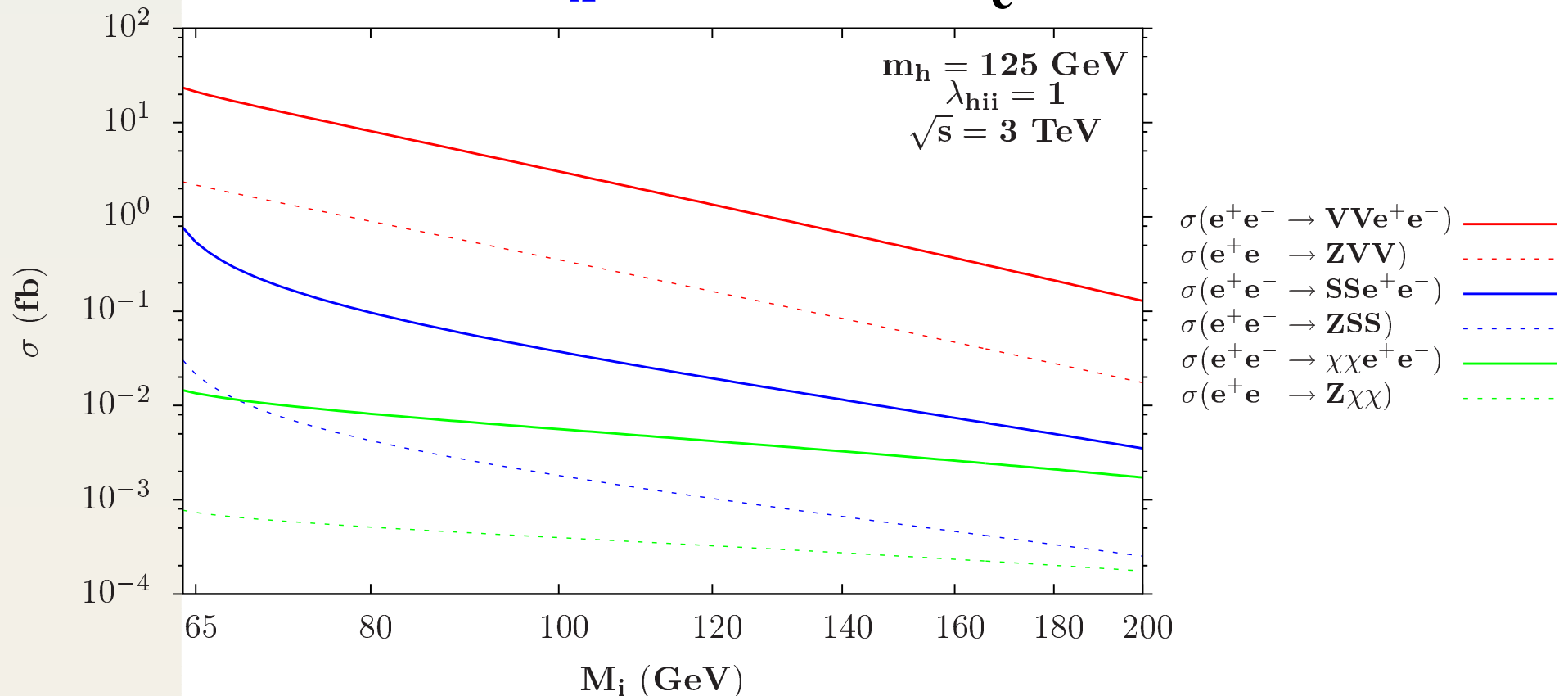
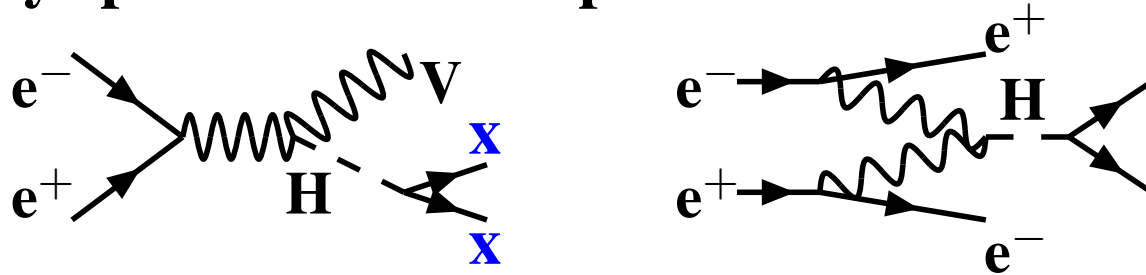
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Only way: produce them in pairs in continuum. At  $e^+e^-$  colliders:



# Conclusions

- Dark matter exists, maybe only reachable sign of new physics?
  - The Higgs portal to DM is one of the most minimal realizations.
  - Even more minimal if only the SM-like Higgs state is considered.
  - Scenario being tested at LHC in Higgs searches/measurements:
    - measurements of total Higgs width and various visible BRs;
    - direct searches for missing energy signature in VH, VBF, ggF;
    - possibility of going off-shell for  $m_{\text{DM}} \gg \frac{1}{2}M_{\text{H}}$  not that promising.
  - Limits from LHC challenged by future astrophysics sensitivity but only with some assumptions on the DM relic density, profile, etc...
  - But I didn't tell you the really interesting part of the story:
    - concrete realizations of Higgs-DM and search for DM companions;
    - extending the Higgs sector makes it even more interesting...
- ⇒ arXiv:1903.03616.**

**Needs further investigations at LHC and also future  $e^+e^-$  and pp colliders besides all experiments in astroparticle physics.**