Generalised contact geometry as reduced generalised complex geometry

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Given a vector bundle $E \cong TM \oplus \wedge^{i_1} T^*M \oplus \cdots \oplus \wedge^{i_n} T^*M$, there is a natural Leibniz algebroid structure with a prescribed set of symmetries and twists, interpreted as splitting the bundle (Baraglia [arXiv:1101.0856]).

- The Leibniz algebroid (and associated L_{∞}) structure and symmetries can be described explicitly.
- Additional structure is sometimes present (e.g. O(d, d) and $E_{d(d)}$)
- The additional structure should be compatible with the 'twisting symmetry' enlarging the notion of equivalence.
- Generalised contact structures were studied in the literature but the full set of twists were not used.

Generalised contact geometry is the odd-dimensional counterpart to Hitchin's generalised contact geometry. Based on a Courant algebroid structure on

$$E\cong TM\oplus\mathbb{R}\oplus\mathbb{R}\oplus T^*M.$$

Background on Generalised Contact structures

- Iglesias and Wade [arXiv:math/0404519]
- Poon and Wade [arXiv:0912.5314]
- Sekiya [arXiv:1212.6064]
- Aldi and Grandini [arXiv:1312.7471]

• Extra symmetries and twists

• K.W [arXiv:1708.09550]

Taking $E \cong TM \oplus \mathbb{R} \oplus \mathbb{R} \oplus T^*M$:

• Gauge transformations are given by

 $\mathsf{Diff}(M) \ltimes (\Omega^2_{\mathsf{cl}}(M), \Omega^1_{\mathsf{cl}}(M), \Omega^1_{\mathsf{cl}}(M)) \qquad \Big(\mathsf{vs}.\mathsf{Diff}(M) \ltimes \Omega^2_{\mathsf{cl}}(M) \Big).$ $(B, b, a) \in (\Omega^2_{\mathsf{cl}}(M), \Omega^1_{\mathsf{cl}}(M), \Omega^1_{\mathsf{cl}}(M)).$

• The twists are given by

$$(H_3, H_2, F) \in (\Omega^3(M), \Omega^2(M), \Omega^2(M)).$$

 The twists describe non-trivial bundles E ≅ TM ⊕ ℝ ⊕ ℝ ⊕ T*M and allow a description of non-coorientable structures. Generalised Contact structures can be viewed as an S^1 -invariant reduction of generalised complex structures.

Generalised complex geometry	Generalised contact geometry
$TM \oplus T^*M$	$TM \oplus \mathbb{R} \oplus \mathbb{R} \oplus T^*M$
Pure spinor $ ho$	Mixed pair $arphi,\psi$
Integrability $d_H \rho$	Integrability $d_{\mathcal{H}_3,\mathcal{F},\mathcal{H}_2}(arphi,\psi)$
Generalised complex structure	Generalised contact structure
Generalised metric	Generalised metric
Generalised Kähler	Generalised coKähler
T-duality acts Gen. Comp.	T-duality acts Gen. Con.

Generalised Contact algebroid

The Contact Courant algebroid is described on the vector bundle $E \cong TM \oplus \mathbb{R} \oplus \mathbb{R} \oplus T^*M$. Given sections $V = (v, f, g, \xi) \in (\Gamma(TM), C^{\infty}(M), C^{\infty}(M), \Gamma(T^*M))$: $V_1 \circ_{H_3, H_2, F} V_2 = \left([v_1, v_2], v_1(f_2) - v_2(f_1) - \iota_{v_1} \iota_{v_2} F, v_1(g_2) - v_2(g_1) - \iota_{v_1} \iota_{v_2} H_2, \right)$ $\mathcal{L}_{v_1}\xi_2 - \iota_{v_2}d\xi_1 - \iota_{v_1}\iota_{v_2}H_3 + g_2df_1 + f_2dg_1$ + $f_1\iota_{\nu_2}H_2 - f_2\iota_{\nu_1}H_2 + g_1\iota_{\nu_2}F - g_2\iota_{\nu_1}F$; $\langle V_1, V_2 \rangle = \frac{1}{2} (\iota_{v_1} \xi_2 + \iota_{v_2} \xi_1 + f_1 g_2 + g_1 f_2);$ $\rho(V) = \rho((v, f, g, \xi)) = v.$

The Leibniz identity for $\circ_{H_3,H_2,F}$ gives

$$dH_3+H_2\wedge F=0,\quad dH_2=0,\quad dF=0.$$

Sekiya quadruple

Definition

A Sekiya quadruple on an odd-dimensional manifold M is given by the quadruple $(\Phi, e_1, e_2, \lambda) \in (\text{End}(\mathbb{T}M), \Gamma(\mathbb{T}M), \Gamma(\mathbb{T}M), C^{\infty}(M))$, satisfying the following conditions (where $\mathbb{T}M = TM \oplus T^*M$):

$$\begin{split} \langle \mathbf{e}_1, \mathbf{e}_1 \rangle &= 0 = \langle \mathbf{e}_2, \mathbf{e}_2 \rangle, \ \langle \mathbf{e}_1, \mathbf{e}_2 \rangle = \frac{1}{2}; \\ \Phi^* &= -\Phi; \\ \Phi(\mathbf{e}_1) &= \lambda \mathbf{e}_1, \ \Phi(\mathbf{e}_2) = -\lambda \mathbf{e}_2; \\ \Phi^2(\mathbf{e}) &= -\mathbf{e} + 2(1+\lambda^2)(\langle \mathbf{e}, \mathbf{e}_2 \rangle \mathbf{e}_1 + \langle \mathbf{e}, \mathbf{e}_1 \rangle \mathbf{e}_2), \text{ for } \mathbf{e} \in \Gamma(\mathbb{T}M). \end{split}$$

Sekiya quadruples can be associated to generalised almost complex structures on $M' = M \times \mathbb{R}$ (or $M' = M \times S^1$):

$$\mathbb{J}_{\mathsf{inv}} = egin{pmatrix} \Phi & \mu e_1 & \mu e_2 \ -2\mu\langle e_2,\cdot
angle & -\lambda & 0 \ -2\mu\langle e_1,\cdot
angle & 0 & \lambda \end{pmatrix}, \quad (\mathbb{J}^* = -\mathbb{J}, \ \mathbb{J}^2 = -1).$$

Definition

Let *M* be an odd-dimensional manifold of dimension *m*. A generalised almost contact structure is a quadruple (L, e_1, e_2, λ) , where $L \subset \mathbb{T}M \otimes \mathbb{C}$ is a maximal isotropic subspace dim_{\mathbb{R}}(L) = m - 1, and $e_1, e_2 \in \Gamma(\mathbb{T}M)$, satisfy $\langle e_1, e_1 \rangle = 0$, $\langle e_2, e_2 \rangle = 0$, $\langle e_1, e_2 \rangle = \frac{1}{2}$.

- Identifying +i-eigenbundle of Φ ∈ End(𝔅𝑘𝔥) with L we have a correspondence between Sekiya quadruples and generalised contact structures.
- A (H₃, H₂, F)-generalised contact structure is an almost generalised contact structure where the +*i*-eigenbundle of J_{inv} is involutive with respect to ◦<sub>H₃,H₂,F.
 </sub>

Definition

A mixed pair $(\varphi, \psi, e_1, e_2)$ consists of two differential forms $\varphi, \psi \in \Gamma(\Omega^{\text{ev/od}}(M) \otimes \mathbb{C})$ and a choice of two sections $e_1, e_2 \in \Gamma(\mathbb{T}M)$ satisfying

$$\begin{split} (\varphi,\bar{\varphi})_{m-1} &\neq 0, \quad (\psi,\bar{\psi})_{m-1} \neq 0, \quad ((\varphi,\psi),(\bar{\varphi},\bar{\psi}))_{m-1} \neq 0, \\ e_1 \cdot \psi &= 0, \quad \mu e_1 \cdot \varphi = (1+i\lambda)\psi, \quad e_2 \cdot \varphi = 0, \quad \mu e_2 \cdot \psi = (1-i\lambda)\varphi, \\ \text{where } \mu &= \sqrt{1+\lambda^2}, \ \lambda \in C^{\infty}(M). \end{split}$$

$$((\varphi_1,\psi_1),(\varphi_2,\psi_2))_{m-1} := (-1)^{|\varphi_1|} (\alpha(\varphi_1) \wedge \psi_2)_{m-1} + (-1)^{|\psi_1|} (\alpha(\psi_1) \wedge \varphi_2)_{m-1},$$

where $\alpha(dx^1 \otimes dx^2 \otimes \cdots \otimes dx^k) = dx^k \otimes dx^{k-1} \otimes \cdots \otimes dx^1$, $m = \dim(M)$, $|\varphi| = k$ for $\varphi \in \Omega^k(M)$, and $(\cdot)_{m-1}$ is the projection to $\Omega^{m-1}(M)$. The integrability of a generalised contact structure is encoded in the mixed pair description by the requirement:

$$d_{H_3,H_2,F}(\varphi,\psi)=V\cdot(\varphi,\psi),$$

where

$$\begin{aligned} d_{H_3,H_2,F}(\varphi,\psi) &:= (d\varphi + H_3 \land \varphi + F \land \psi, H_2 \land \varphi - d\psi - H_3 \land \psi), \\ V \cdot (\varphi,\psi) &= (v,f,g,\xi) \cdot (\varphi,\psi) = (\iota_v \varphi + \xi \land \varphi + f\psi, g\varphi - \iota_v \psi - \xi \land \psi), \end{aligned}$$

and (H_3, H_2, F) satisfy

$$dH_3 + H_2 \wedge F = 0, \quad dH_2, \quad dF = 0.$$

Example: Symplectic structure

Let $\omega \in \Gamma(\wedge^2 T^*P)$ be a symplectic form, where $P = M \times S^1$ where S^1 has coordinate t, and dim(M) = m = 2n + 1. The symplectic form is S^1 -invariant if it admits the decomposition

$$\omega = \theta + dt \wedge \eta \in \Omega^2(P).$$

$$0 \neq \omega^{n+1} = (\theta + dt \wedge \eta)^{n+1} = (n+1)\theta^n \wedge dt \wedge \eta \Rightarrow \eta \wedge \theta^n \neq 0.$$

There exists a Reeb vector field $R \in \Gamma(TM)$ such that $\iota_R \eta = 1$ and $\iota_R \theta = 0$. θ is non-degenerate on ker (η) , and $\phi(v) := \iota_v \theta$ ($v \in \text{ker}(\eta)$). The generalised complex structure \mathbb{J}_{ω} is reduced to a generalised contact structure $(\Phi, e_1, e_2, \lambda)$:

$$\begin{split} \mathbb{J}_{\omega} &= \begin{pmatrix} 0 & -\omega^{-1} \\ \omega & 0 \end{pmatrix} \quad \Rightarrow \quad \Phi = \begin{pmatrix} 0 & -\theta^{-1} \\ \theta & 0 \end{pmatrix}, \quad e_1 = \eta, \quad e_2 = R, \quad \lambda = 0, \\ \rho_{\mathbb{J}_{\omega}} &= e^{i\omega} \quad \Rightarrow \quad \varphi = e^{i\theta}, \quad \psi = \eta \wedge e^{i\theta}, \end{split}$$

where $\mathbb{J}_{\omega} \in \operatorname{End}(TP \oplus T^*P)$ and $\Phi \in \operatorname{End}(\ker(\eta) \oplus \operatorname{Ann}(R))$.

Example: Complex structure

Consider an almost complex structure $J \in \text{End}(TP)$, on $P = M \times S^1$, and dim M = 2n + 1. Given local coordinates, x, for M, and t for S^1 , the almost complex structure is S^1 -invariant if there exists the decomposition

$$J_{K}^{l}\partial_{l}\otimes dx^{K}=(\phi_{k}^{i})\partial_{i}\otimes dx^{k}+R^{i}\partial_{i}\otimes dt+lpha_{k}\partial_{t}\otimes dx^{k}.$$

The conditions $J^* = -J$ and $J^2 = -1$ give

$$\iota_R \alpha = 1, \quad \Phi(R) = 0 = \phi^*(\alpha), \quad \phi^2(v) = -v + (\iota_v \alpha)R.$$

The generalised almost complex structure \mathbb{J}_J reduces to a generalised almost contact structure $(\Phi, e_1, e_2, \lambda)$:

$$\mathbb{J}_{J} = \begin{pmatrix} J & 0 \\ 0 & -J^{*} \end{pmatrix} \quad \Rightarrow \quad \Phi = \begin{pmatrix} \phi & 0 \\ 0 & -\phi^{*} \end{pmatrix}, \quad e_{1} = \alpha, \quad e_{2} = R, \quad \lambda = 0,$$
$$\rho_{J} = \Omega_{J} \quad \Rightarrow \quad \varphi = \Omega_{\phi}, \quad \psi = \alpha \wedge \Omega_{\phi},$$

where $\Omega_J \in \Omega^{2n+2,0}$ is the decomposable top form giving the pure spinor describing \mathbb{J}_J , and $\alpha \wedge \Omega_{\phi} := \Omega_J$.

$$e^{(B,b,a)}V = \left(v, f + 2\langle v, a \rangle, g + 2\langle v, b \rangle, \xi + \iota_v B - fb - ga - \langle v, a \rangle b - \langle v, b \rangle a\right).$$

Compatibility with symmetries

$$\langle e^{(B,b,a)} V_1, e^{(B,b,a)} V_2 \rangle = \langle V_1, V_2 \rangle, \\ e^{(B,b,a)} V_1 \circ_{(0,0,0)} e^{(B,b,a)} V_2 = e^{(B,b,a)} (V_1 \circ_{H'_3,H'_2,F'} V_2),$$

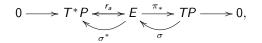
where

$$H_3' = dB + rac{1}{2}(da \wedge b + a \wedge db), \quad H_2' = db, \quad F' = da.$$

Composition of symmetries:

$$(B_2, b_2, a_2) \cdot (B_1, b_1, a_1) = (B_1 + B_2 - \frac{1}{2}(b_1 \wedge a_2 + a_1 \wedge b_2), b_1 + b_2, a_1 + a_2),$$

Splitting bundles



Exact Courant algebroid with $H_3(X_1, X_2) = \sigma^*(\sigma(X_1) \circ_{\mathsf{Dorf}} \sigma(X_2))$, $X \in \Gamma(P)$.



$$0 \longrightarrow P \times \mathbb{R} \xrightarrow{TP/U(1)} TM \longrightarrow 0,$$

Atiyah algebroid with $F(v_1, v_2) = s^*([s(v_1), s(v_2)])$, where $v \in \Gamma(TM)$.

- Create examples based on non-coorientable structures (those not defined by a globally defined one-form) by considering *F* ≠ 0.
 *H*₃, *H*₂, *F* are globally defined forms but *B*, *b*, *a* are not in general global forms: Gerbe structure
- Use reduction to define generalised coKähler and generalised coKähler–Einstein structures.
- T-duality maps a generalised coKähler(-Einstein) structure to another generalised coKähler(-Einstein) structure.