

# Cosmic microwave background as a probe of dark relics

*Corfu Summer Institute - Workshop on the SM and beyond, 08/09/2018*



# RATIONALE AND OUTLINE

**What if new physics is very weakly coupled?**

E.g. what if DM does little more than its job of “gravitating”?

**Is this type of BSM undetectable?**

Not quite, there’s still hope!

Notably from **cosmology**, where indirect techniques offer interesting probes even of “almost invisible” relics

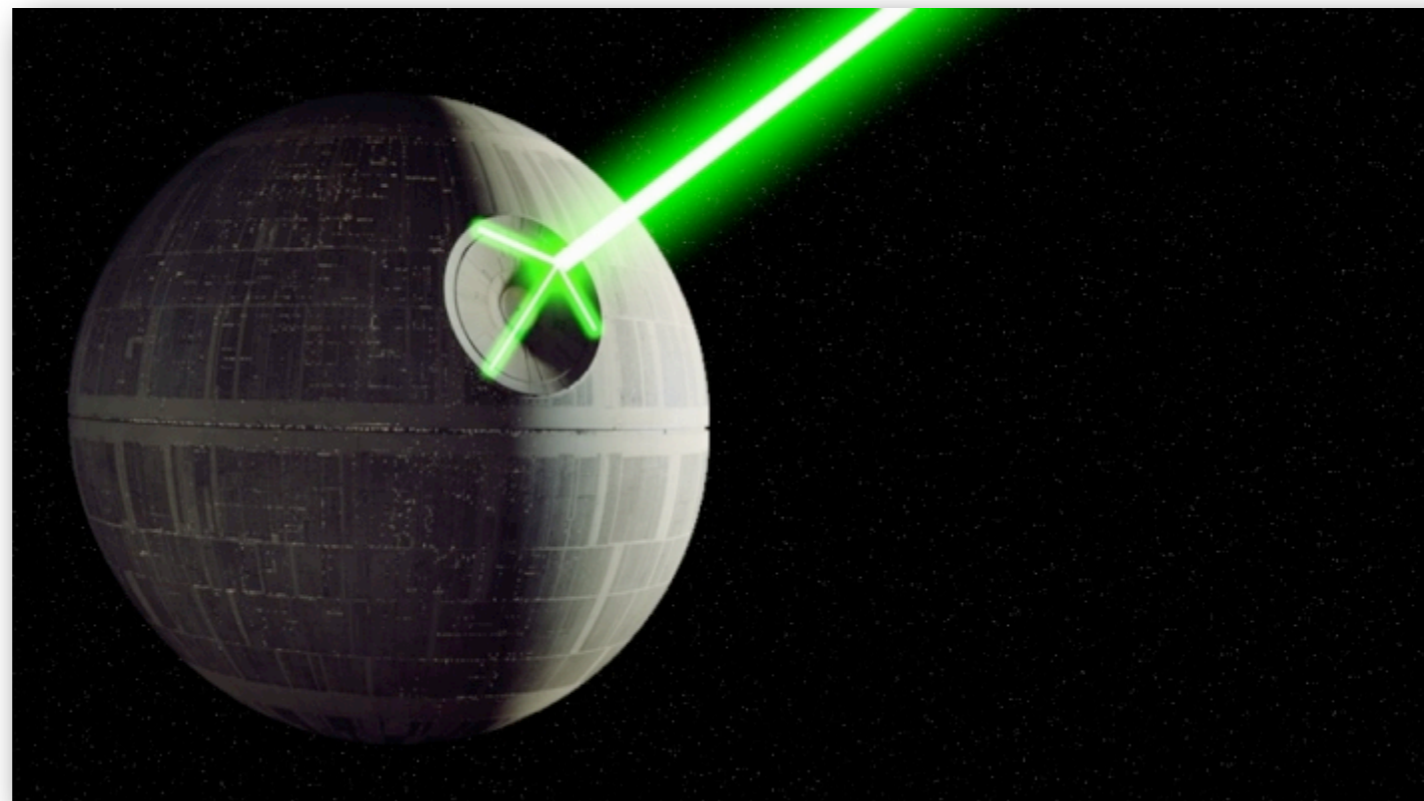
- ▶ **“Dark matter” (part or all) conversion into *dark radiation***
  - Purely gravitational effects! Rationale of the argument & applications
- ▶ **Massive relics injecting some electromagnetic energy**
  - Principles & expected sensitivity
  - Application to annihilating WIMPs, decaying relics
  - Bonus (time permitting), applications to primordial BHs
- ▶ **Sensitivity to dark radiation & conclusions**

# Episode I

possibly the next-to-closest thing to  
an "undiscoverable" DM candidate...



“Dark Matter” conversion into “Dark Radiation”:  
Gravitational effects



# On a decaying DM fraction

Assume a stable component in DM, plus an unstable relic, whose fraction of the initial total is  $f$ , decaying into “dark” relativistic species (DR).

$$\begin{aligned}\Omega_{\text{dm}} &= \Omega_{\text{sdm}} + \Omega_{\text{dcdm}} \\ &= (1 - f_{\text{dcdm}})\Omega_{\text{dm}}^{\text{ini}} + f_{\text{dcdm}} \exp(-\Gamma_{\text{dcdm}}t)\Omega_{\text{dm}}^{\text{ini}}\end{aligned}$$

To some extent also describes DM'  $\rightarrow$  “lighter” DM, which has however additional constraints

The smooth background equations can be easily derived, e.g. from  $\nabla_{\mu}T^{\mu\nu} = 0$

homogeneous equations given by  
(prime=derivative with respect to conformal time)

$$\rho'_{\text{dcdm}} = -3\frac{a'}{a}\rho_{\text{dcdm}} - a\Gamma_{\text{dcdm}}\rho_{\text{dcdm}}$$

$$\rho'_{\text{dr}} = -4\frac{a'}{a}\rho_{\text{dr}} + a\Gamma_{\text{dcdm}}\rho_{\text{dcdm}}$$

dilution factors for energy density of  
matter  
radiation

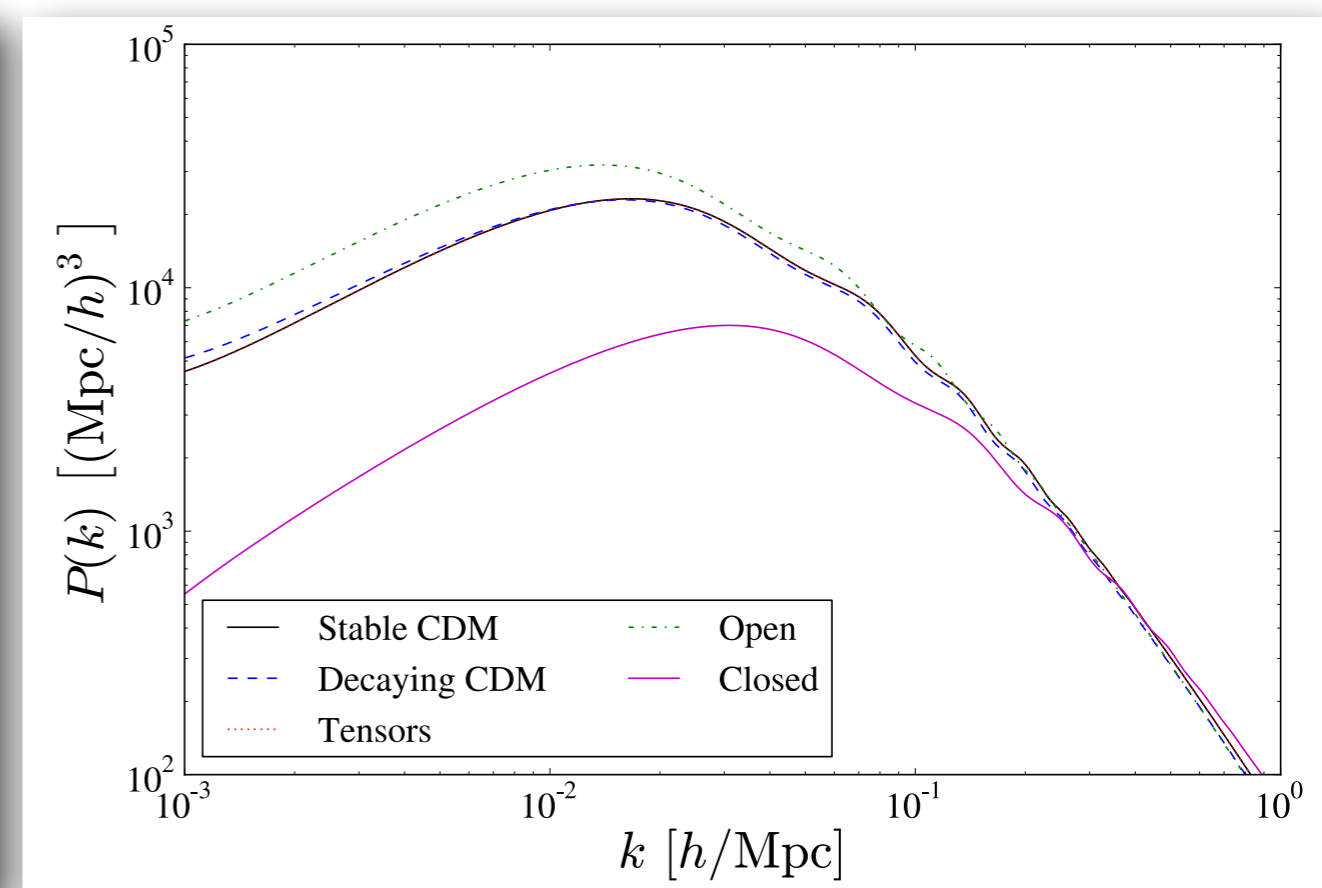
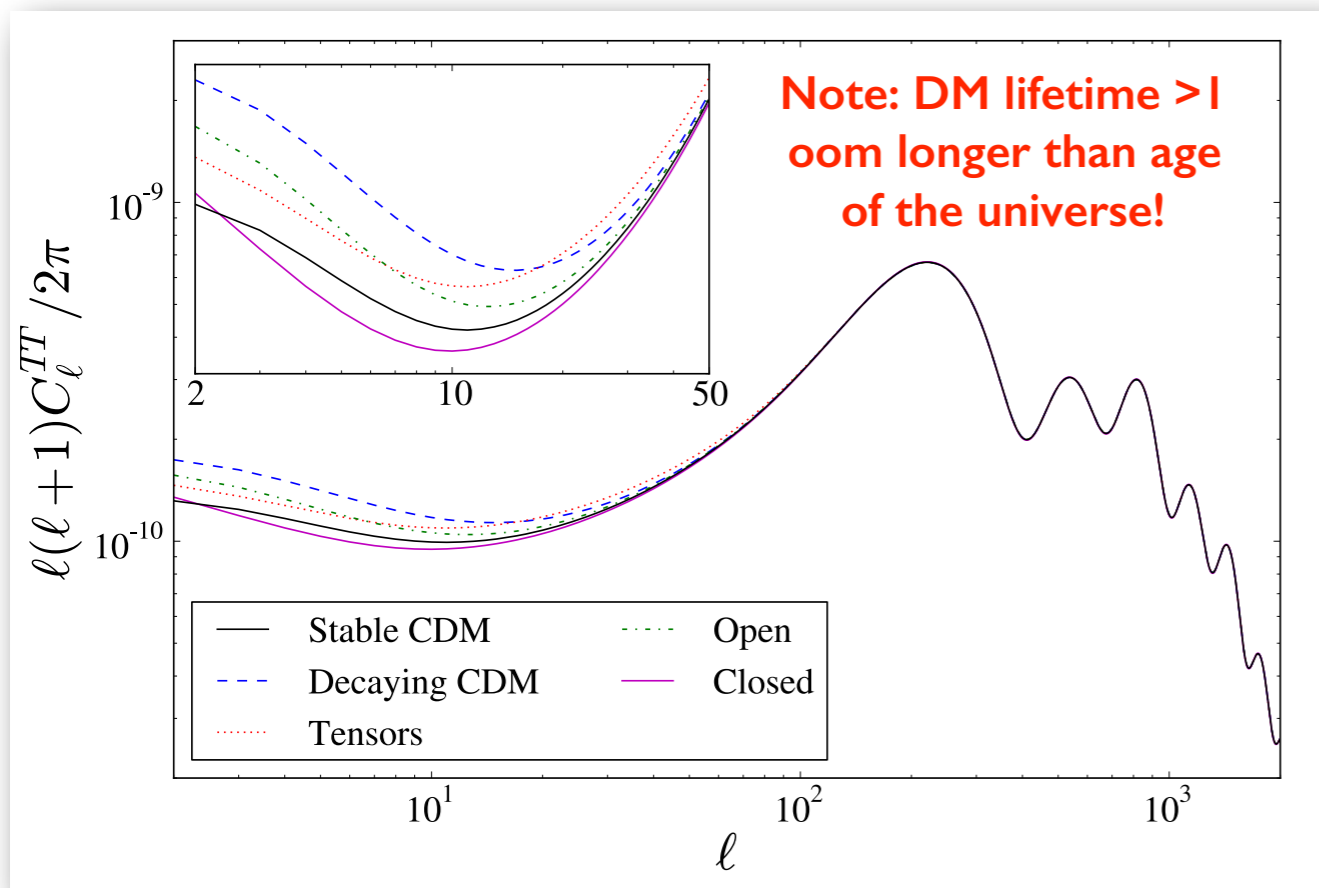
$$\frac{a'}{a} = aH = \mathcal{H}$$

For perturbations, must be careful about gauge choice/fixing... I won't enter in details, if interested see [V. Poulin, P.D.S. and J. Lesgourgues, JCAP 1608, 036 \(2016\) \[1606.02073\]](#)

# Effects of decaying DM ( $f_{\text{dcdm}}=1$ , first)

CMB affected (mostly) by late integrated Sachs-Wolfe effect (modification of homogeneous & perturbed DM density at late times affects evolution of metric fluctuation) LSS helps in breaking partial degeneracy with curvature & tensor modes

Model implemented in CLASS, <http://class-code.net/>



Case for  $f_{\text{dcdm}}=1$ , from

*B. Audren et al. JCAP 1412, 028 (2014) [1407.2418]*

Current bounds:  $\tau \gtrsim 160 \text{ Gyr}$  (CMB only)

$\tau \gtrsim 170 \text{ Gyr}$  (with other consistent data)

*V. Poulin, P.D.S. and J. Lesgourgues, JCAP 1608, 036 (2016) [1606.02073]*

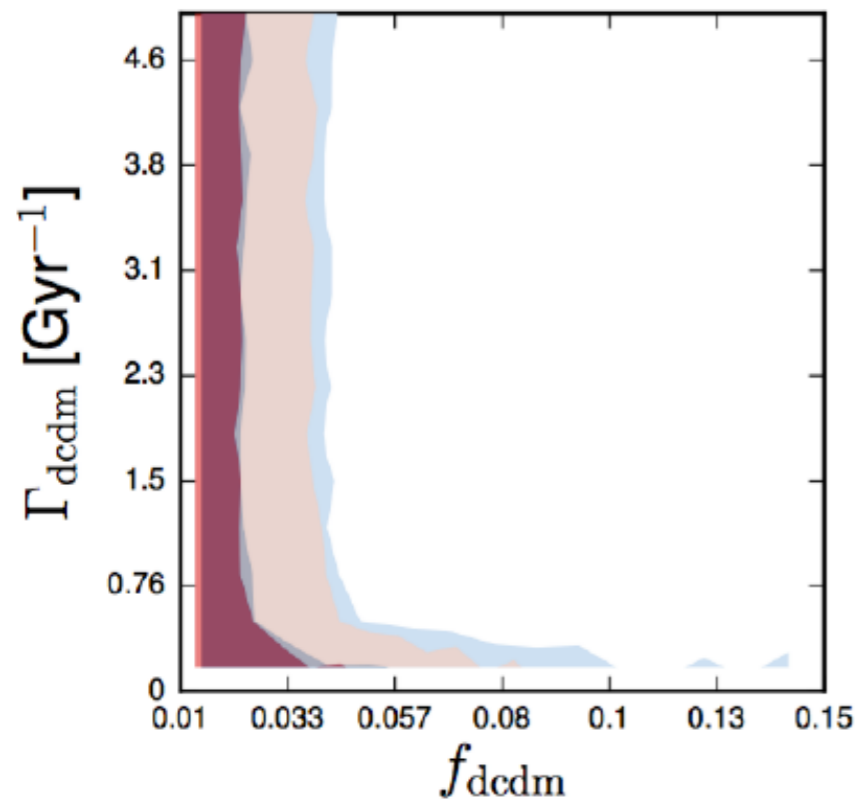
# Bounds: 3 timescale regimes

If the lifetime is very **long**, to first order data are only sensitive to the product  $\Gamma f$

$\Gamma f < 0.0063$  (**0.0059**)  $\text{Gyr}^{-1}$  CMB only (**+consistent data**)

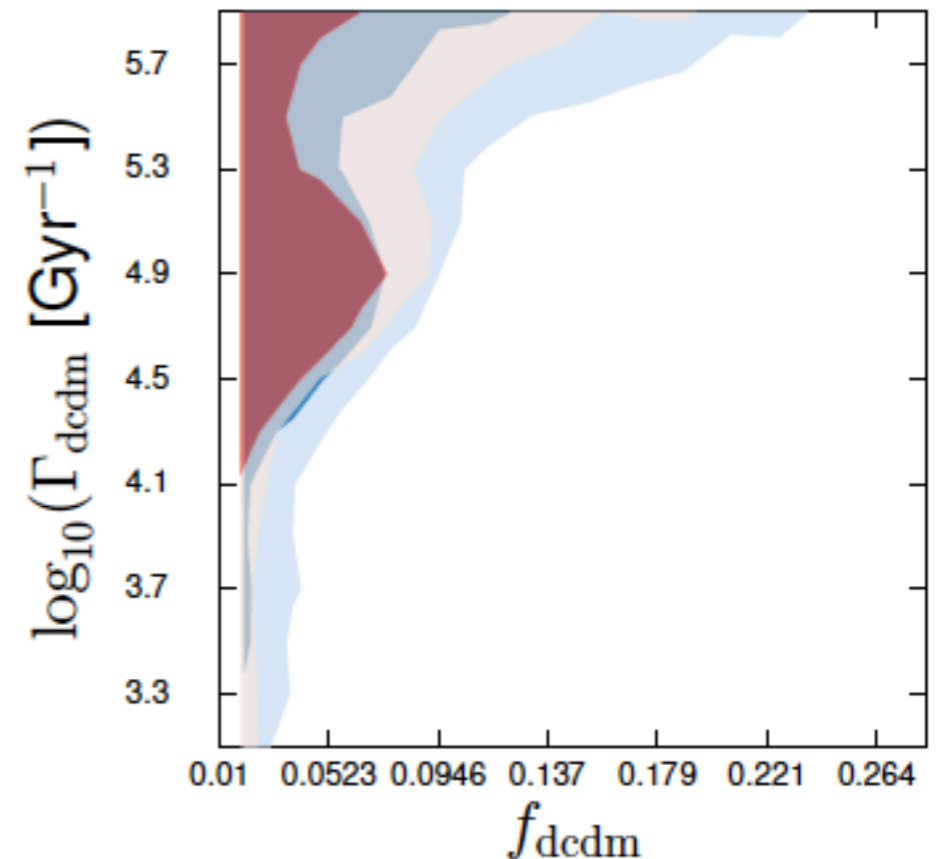
$$\begin{aligned}\Omega_{\text{dm}} &= \Omega_{\text{sdm}} + \Omega_{\text{dcdm}} \\ &= (1 - f_{\text{dcdm}})\Omega_{\text{dm}}^{\text{ini}} + f_{\text{dcdm}} \exp(-\Gamma_{\text{dcdm}} t)\Omega_{\text{dm}}^{\text{ini}} \\ &= (1 - f_{\text{dcdm}})\Omega_{\text{dm}}^{\text{ini}} + f_{\text{dcdm}}[1 - \Gamma_{\text{dcdm}} t + \mathcal{O}((\Gamma_{\text{dcdm}} t)^2)]\Omega_{\text{dm}}^{\text{ini}} \\ &= [1 - f_{\text{dcdm}}\Gamma_{\text{dcdm}} t + \mathcal{O}((\Gamma_{\text{dcdm}} t)^2)]\Omega_{\text{dm}}^{\text{ini}}.\end{aligned}$$

bounds  $\sim$  independent of lifetime **between** recombination and recent times (bounds apply also to complicated, non-decaying DM)



Less than 3.8% of DM has converted into any invisible radiation from recombination to now!

bounds on  $f_{\text{dcdm}}$  relax for very **short** lifetimes, accompanied by an increase in the value of  $\Omega^{\text{ini}}$



...but for a “sweet spot” where “the produced DR” alters  $N_{\text{eff}}$  appreciably

# Numerous applications

## Examples in the literature:

**within SUSY**, if the LSP and NLSP are gravitinos, axions/axinos, RH sneutrinos...

*for a recent ex. see e.g. R. Allahverdi et al. "Dark Matter from Late Invisible Decays to/of Gravitinos," Phys. Rev D 91, 055033 (2015)*

**BSM models** (including string-inspired) accompanied by dark sectors; generically the lightest particle expected in the dark sector and the **lightest "visible" SUSY partner is metastable**

*B. S. Acharya, S. Ellis, G. Kane, B. Nelson & M. Perry, "The lightest visible-sector supersymmetric particle is likely to be unstable," Phys. Rev. Lett. 117, 181802 (2016)*

**non SUSY examples:** keV-scale majoron, decaying into neutrinos

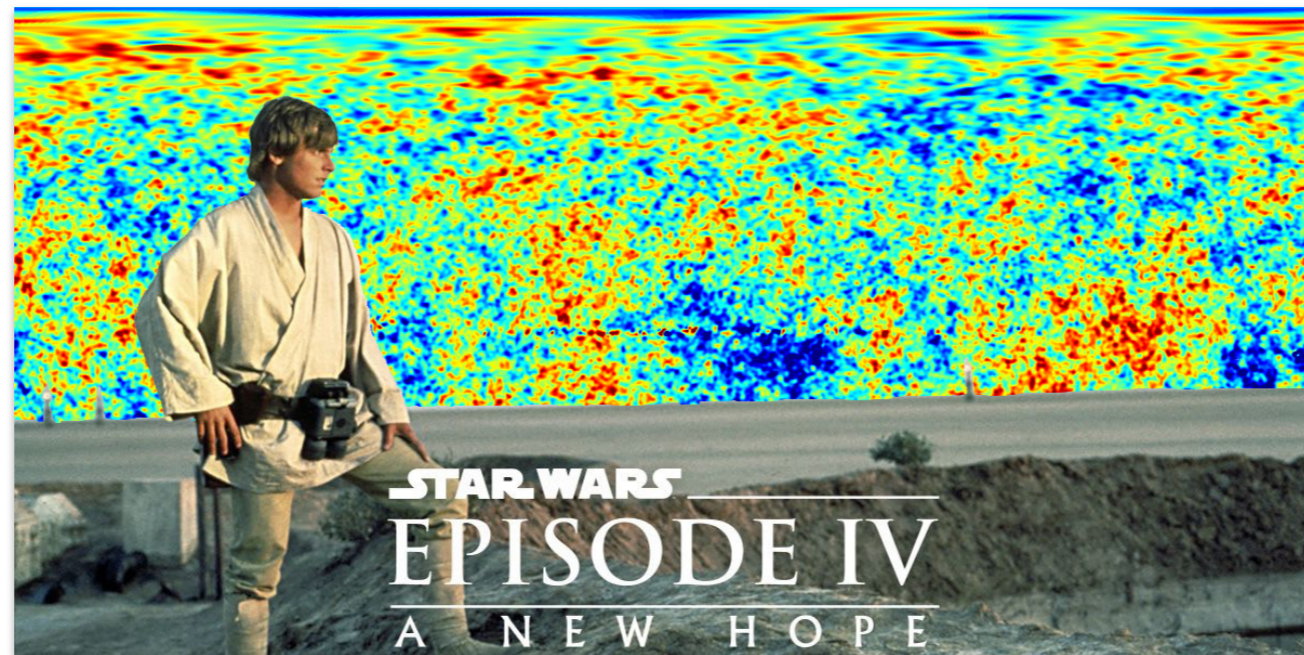
*e.g. M. Lattanzi and J.W. F. Valle, "Decaying warm dark matter and neutrino masses," Phys. Rev. Lett., vol. 99, p. 121301, 2007*

**"non-particle" example:** Primordial Black Holes (DR = GW due to merging)

➔ **either PBH do not make a sizable fraction of the DM or their mass function evolution should be negligible**

# Episode II (or IV?)

Slightly less invisible relics:  
rays of hope





# What if a relic injects interacting SM particles?

associated to a number of processes, like

- Annihilating relics (like WIMP DM)
- Decaying relics such as sterile  $\nu$ 's, Super-WIMP progenitors
- Evaporating (hence “light”) primordial black holes
- Accreting (hence “stellar mass or heavier”) primordial black holes

What happens e.g. to CMB observables?

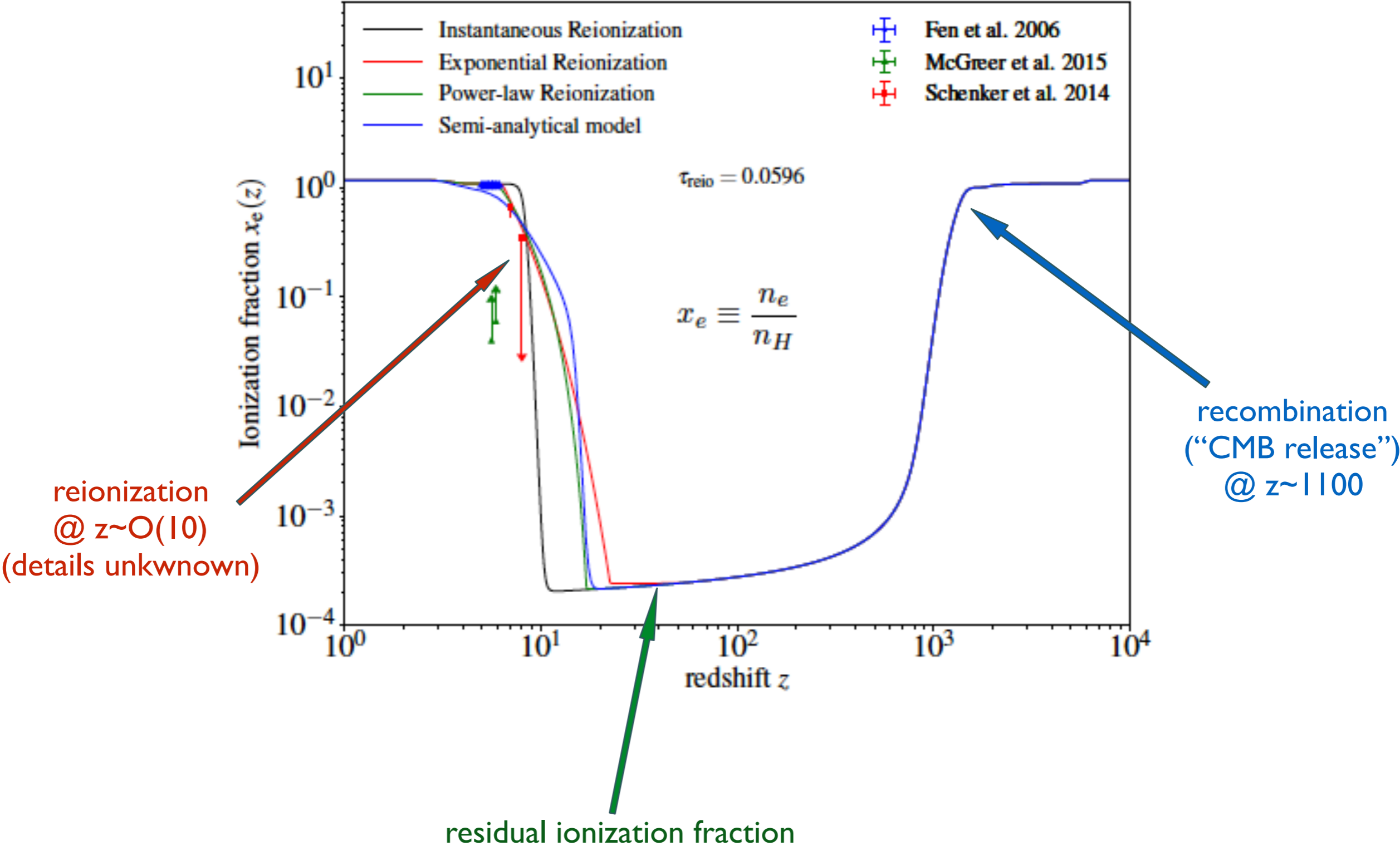
the energy of the injected non-thermal particles is **not negligible wrt the kinetic energy of the baryonic gas.**

The e.m. interacting part of the injection can eventually **heat up** (alter  $T_M$ ) and especially **ionize the gas** (alter  $x_e$ )! (hence alteration in the optical depth experienced by the CMB photons)

*CMB is very sensitive to that!*

# Basic estimates

Have a look at the standard ionization and gas temperature evolution



# Basic estimates

Have a look at the standard ionization  
and gas temperature evolution

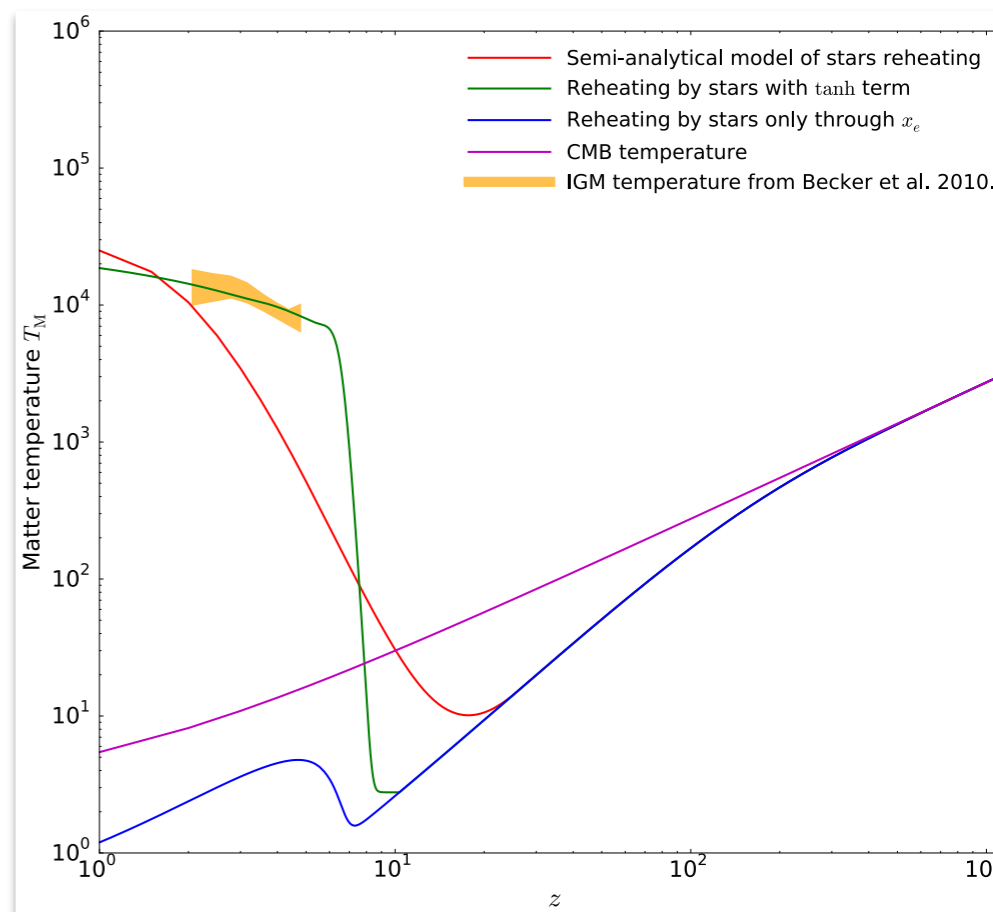
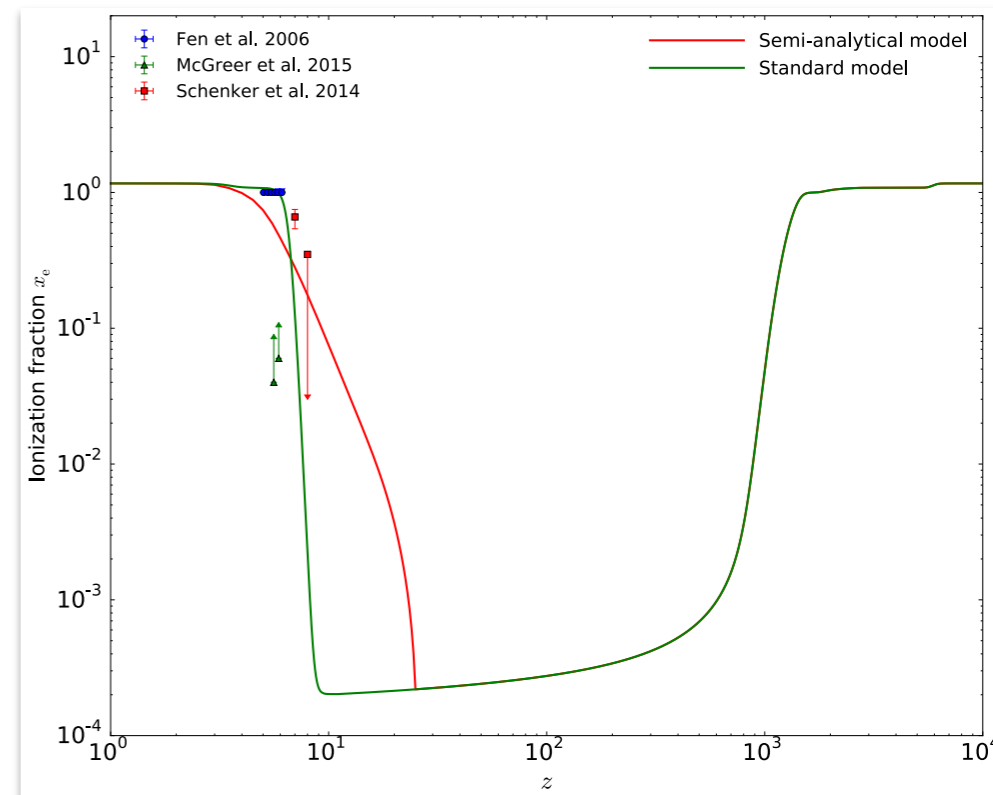
## Note:

$\mathcal{O}(100)$  eV/baryons more than enough to ionize  
all atoms!

In the DM, in principle  $\sim 5$  GeV/baryon “stored”

The reionization fraction in the standard  
expectation drops to  $\sim 5 \cdot 10^{-4}$

a “visible” b.r. of  $\mathcal{O}(10^{-11})$  may be sufficient to  
induce major alterations in  $x_e$  or  $T_M$ !



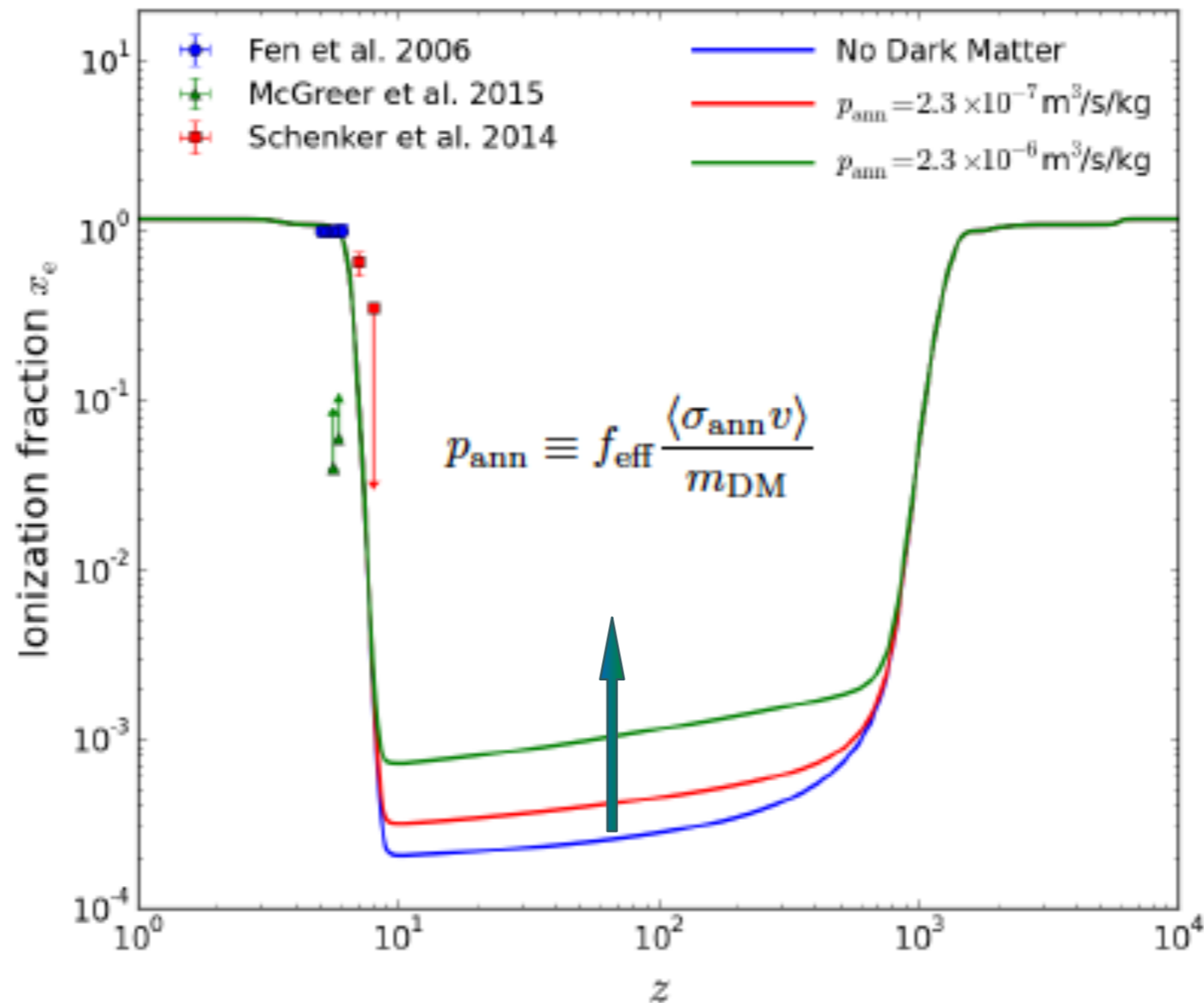
# For instance, what do WIMPs do on CMB?

Via annihilation byproducts, they inject energy in the medium

*key-parameter  
linked to particle physics*

$$\frac{dE}{dV dt} = \rho_c^2 (1+z)^6 \Omega_{\text{DM}}^2 p_{\text{ann}}$$

$$p_{\text{ann}} = \frac{\langle \sigma v \rangle}{8\pi m_X^2} [4\pi] [2m_X] f(z) = f(z) \frac{\langle \sigma v \rangle}{m_X}$$



# For instance, what do WIMPs do on CMB?

Via annihilation byproducts, they inject energy in the medium

$$\frac{dE}{dV dt} = \rho_c^2 (1+z)^6 \Omega_{\text{DM}}^2 p_{\text{ann}}$$

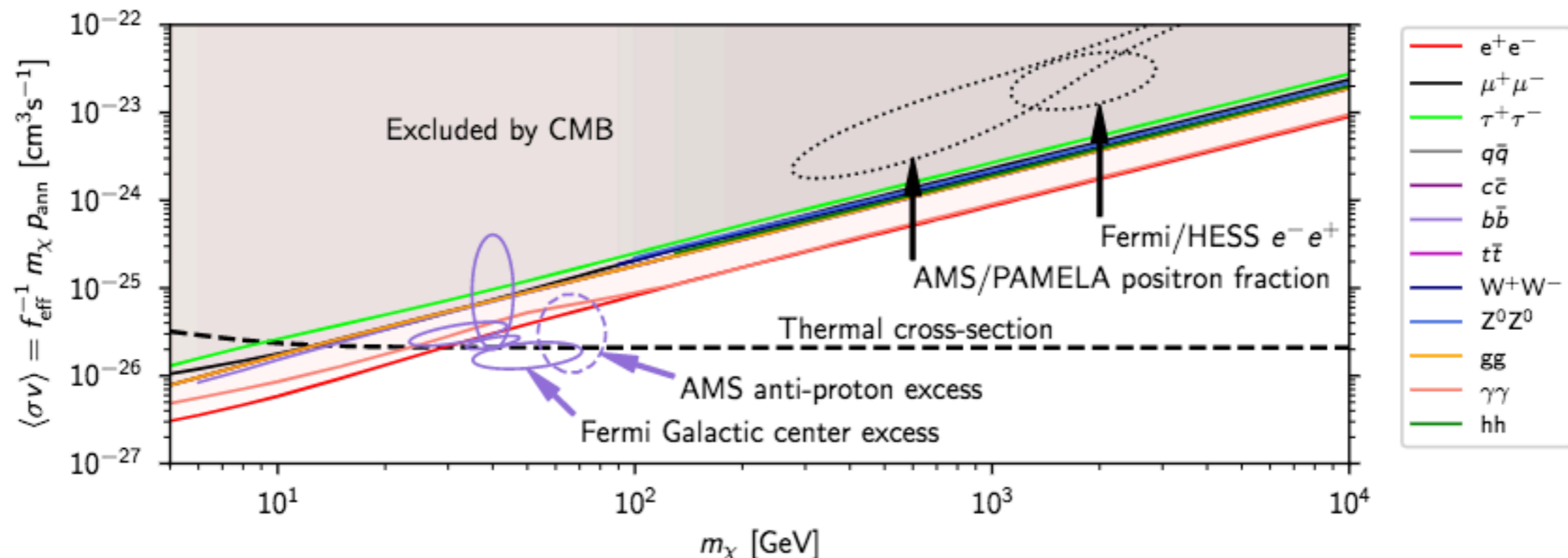
key-parameter  
linked to particle physics

$$p_{\text{ann}} = \frac{\langle \sigma v \rangle}{8\pi m_X^2} [4\pi] [2m_X] f(z) = f(z) \frac{\langle \sigma v \rangle}{m_X}$$

$$p_{\text{ann}} < 3.2 \times 10^{-28} \frac{\text{cm}^3}{\text{s GeV}}$$

(Planck VI, 2018, 1807.06209)

computed with ExoCLASS module, P. Stöcker, M. Krämer, J. Lesgourgues and V. Poulin, JCAP1803 (2018)



Same ballpark of “low-z”, astrophysical constraints **but calorimetric and independent of astro details**

# Example of application: relic decay (superWIMP, sterile $\nu$ ...)

$$\left. \frac{dE}{dV dt} \right|_{\text{inj}} = (1+z)^3 \Xi \Omega_{\text{DM}} \rho_c c^2 \Gamma e^{-\Gamma t}$$

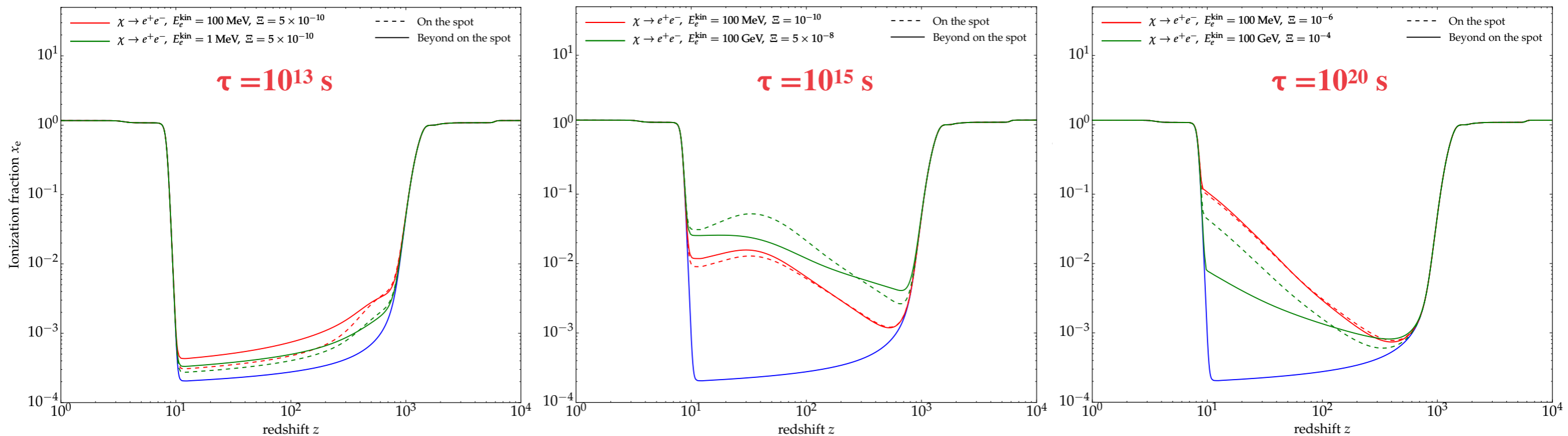
$\Xi$  is the relative amount of energy released into e.m. for a single decay. For instance, a species constituting 1% of the total DM abundance decaying into  $\nu \gamma$  corresponds to  $\Xi = 1/200$ .

We can define the efficiency  $f$ -functions, and compute the corresponding evolution of  $x_e$  and  $T_M$  which show a certain variety, notably due to the large range of  $\Gamma$  allowed

similar to WIMP annihilation

peculiarly bumpy!

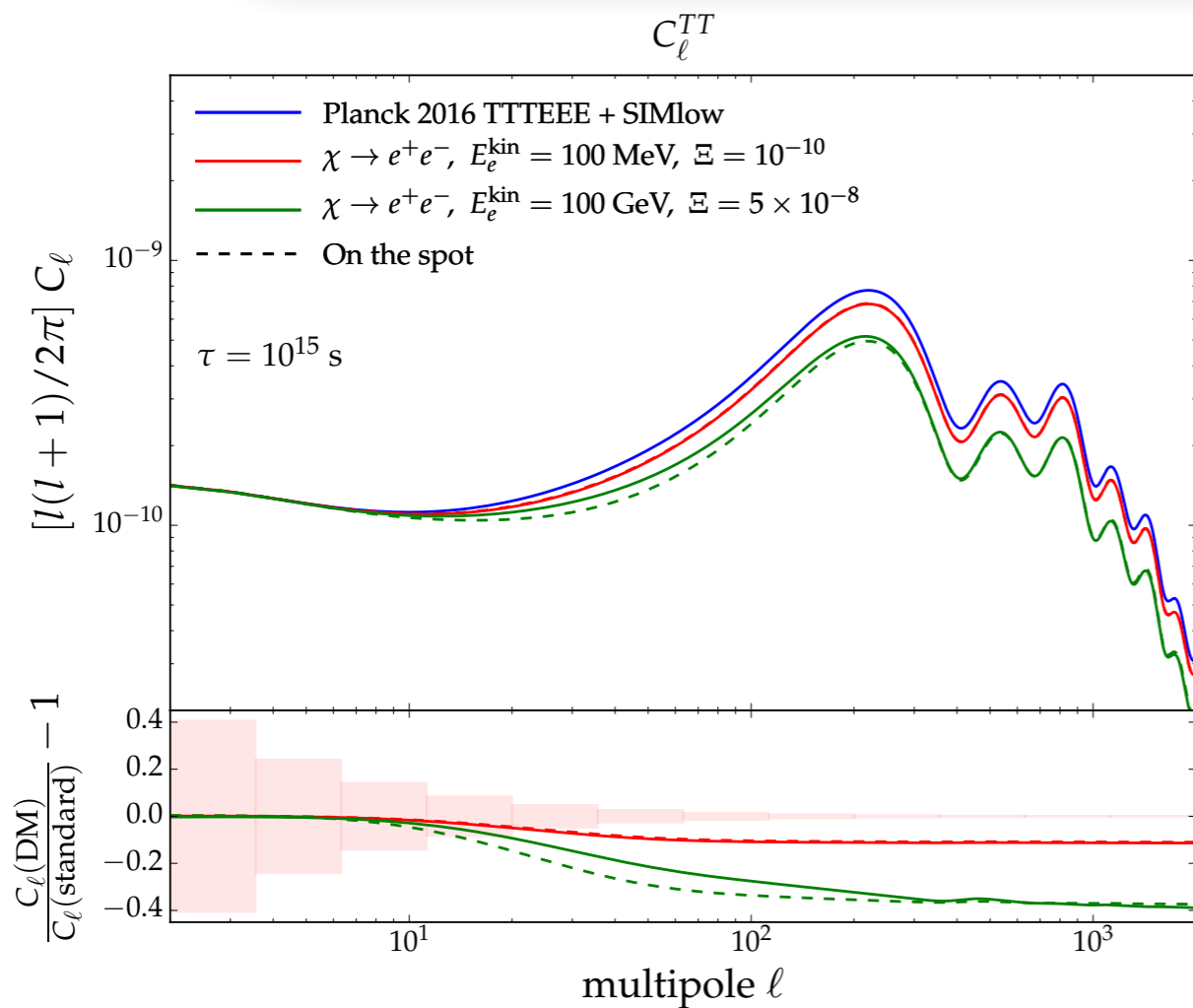
similar to early star formation



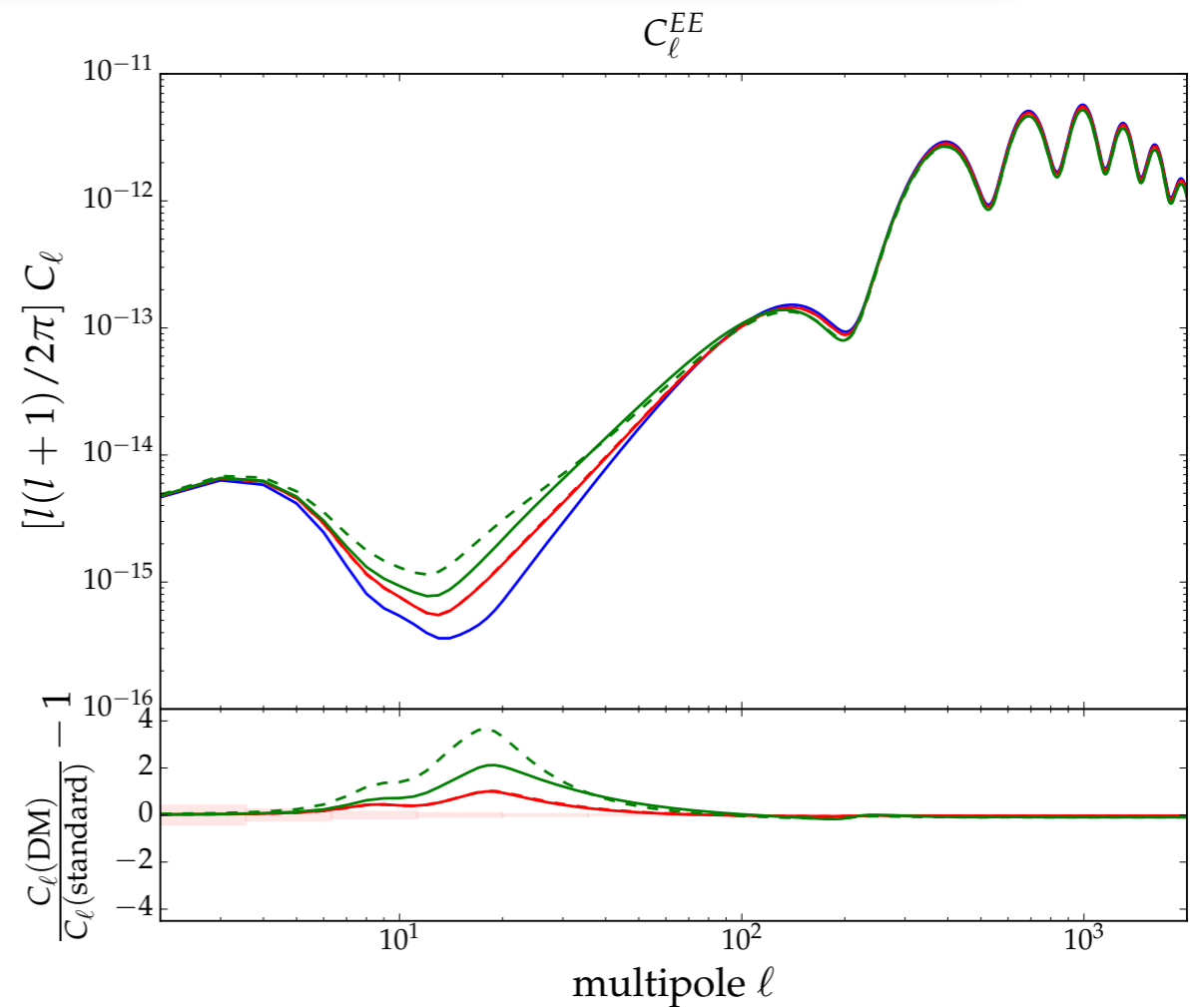
# relic decay, $C_l$ 's

- ➔ Recombination delay: shifts of the peak, more diffusion damping
- ➔ Higher freeze-out plateau: reionization bump higher, higher optical depth

Major physical change is to reionization, since we keep  $\exp(-2\tau_{\text{reio}})A_s$  fixed (i.e. same suppression at large scales)



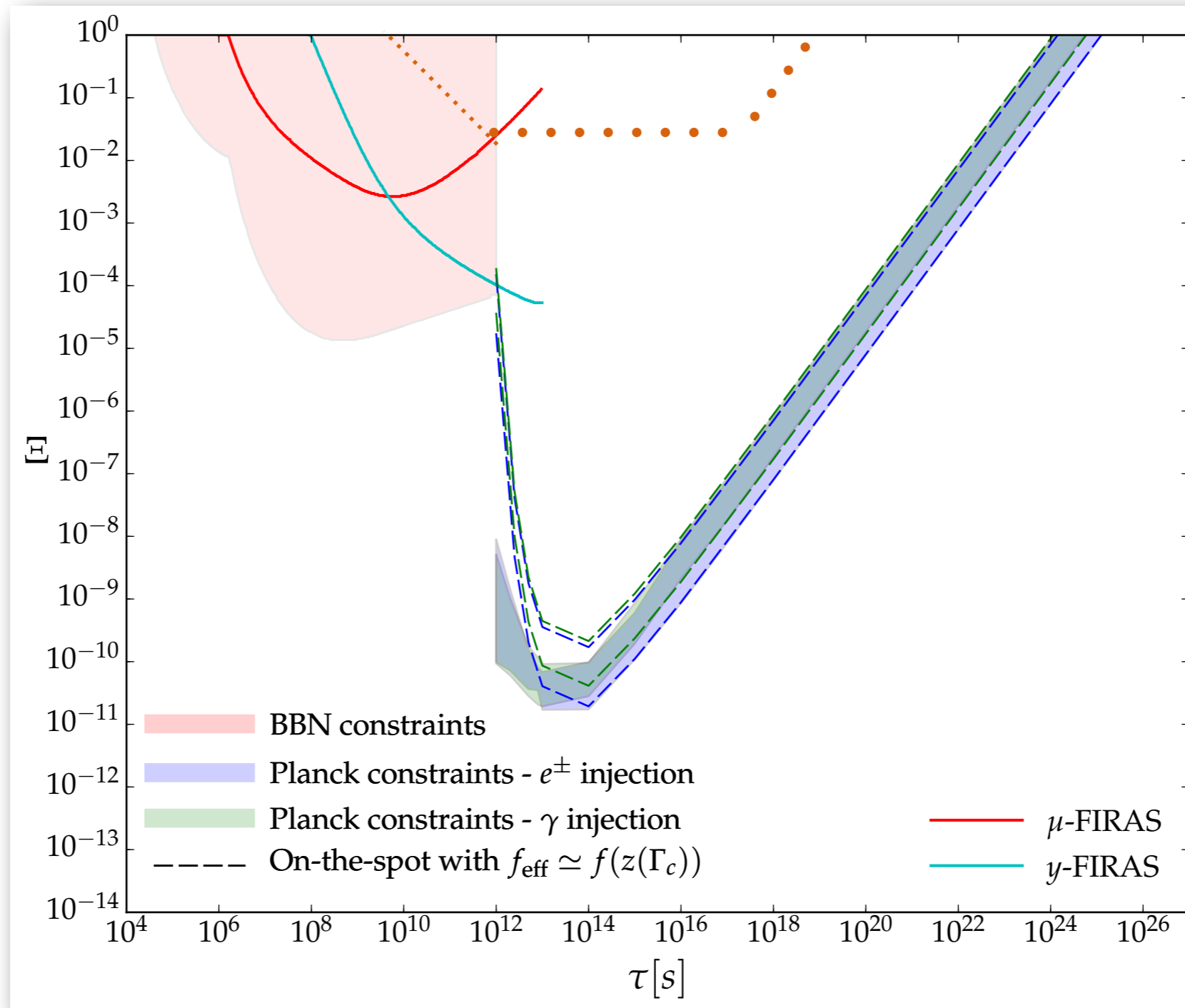
On TT the main effect is visible on the small-scale normalization.



As for EE, effect at relatively large scales (known to be good probe of reionization)

# Results

## & Complementarity of different probes



### Notes:

- 1) we do reach the  $10^{-11}$  level maximal sensitivity estimated at the beginning, for stuff decaying around recombination time
- 2) Much better than purely **gravitational!**
- 3) Complementarity (timescales and actually energies, too!) with other probes



# Bonus Episode

Application to primordial black holes



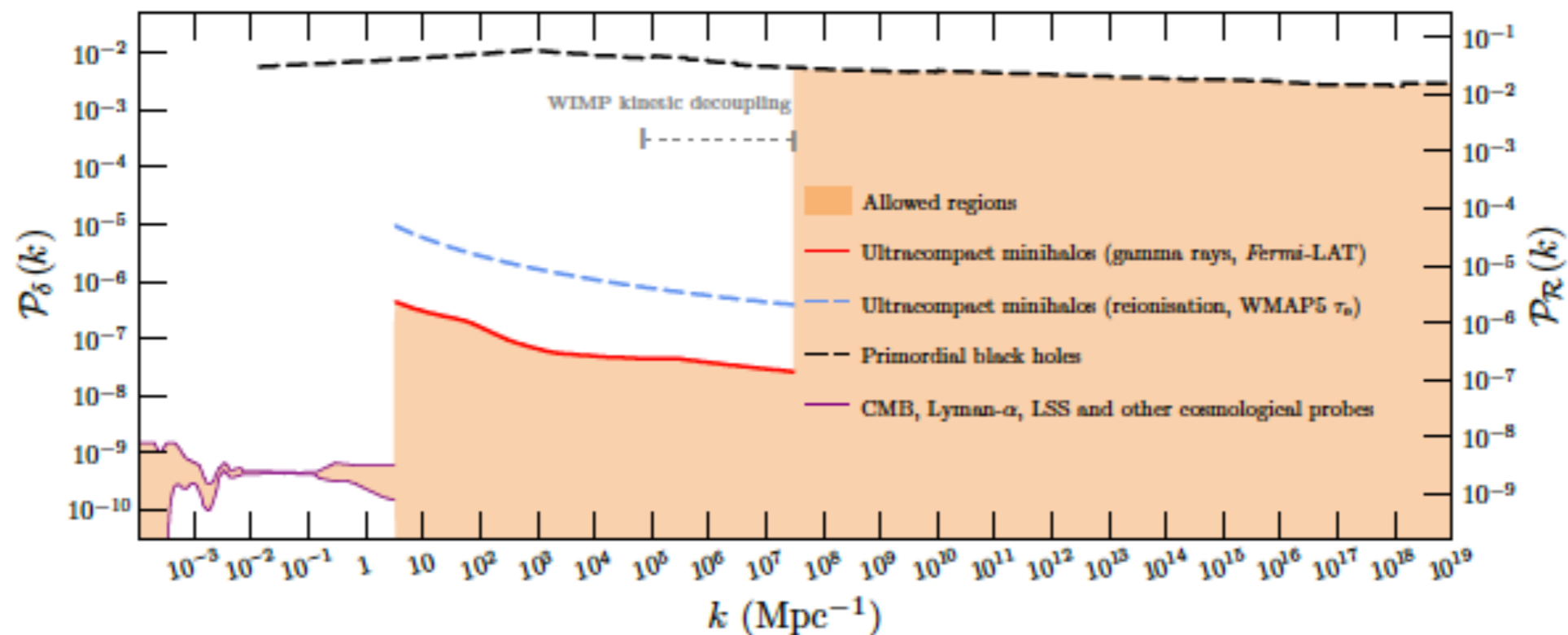
# Since I mentioned Primordial Black Holes (PBH):

When thinking of “early universe relics” (like CMB photons) we usually think of *particles*

Yet, **PBH** (*Zeldovich, Hawking...*) are possibly macroscopic relics which can originate **from gravitational collapse of sufficiently large density fluctuations, at scales much smaller** ( $k \gg \text{Mpc}^{-1}$ ) **than the CMB ones**, typically associated to *non-trivial inflationary dynamics or phase transitions* (the kind of ingredients seen in the previous discussion!)

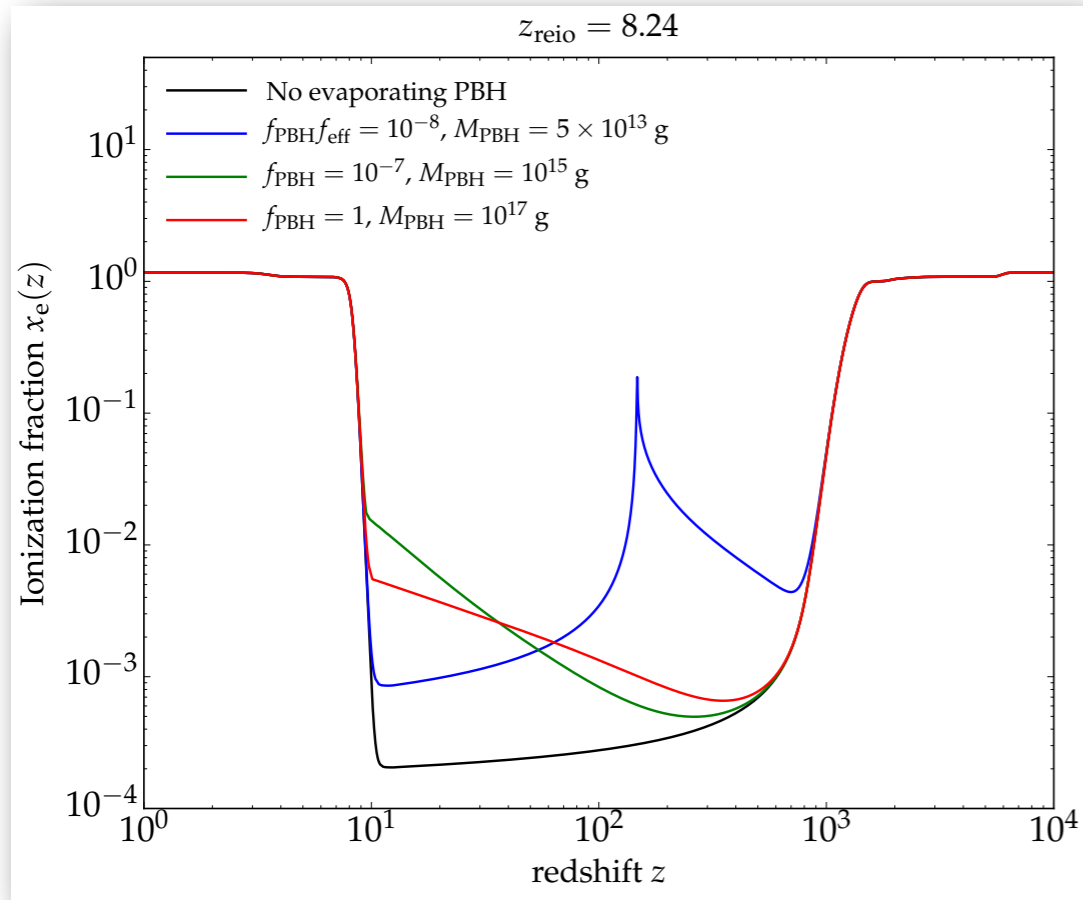
**Such scales are almost unconstrained**

(avoiding PBH overproduction which would over close the Universe is one of the few bounds)



T. Bringmann, P. Scott and Y. Akrami, *Phys. Rev. D* 85, 125027 (2012) [1110.2484]

# Evaporating PBH effects on $x_e$ & CMB bounds



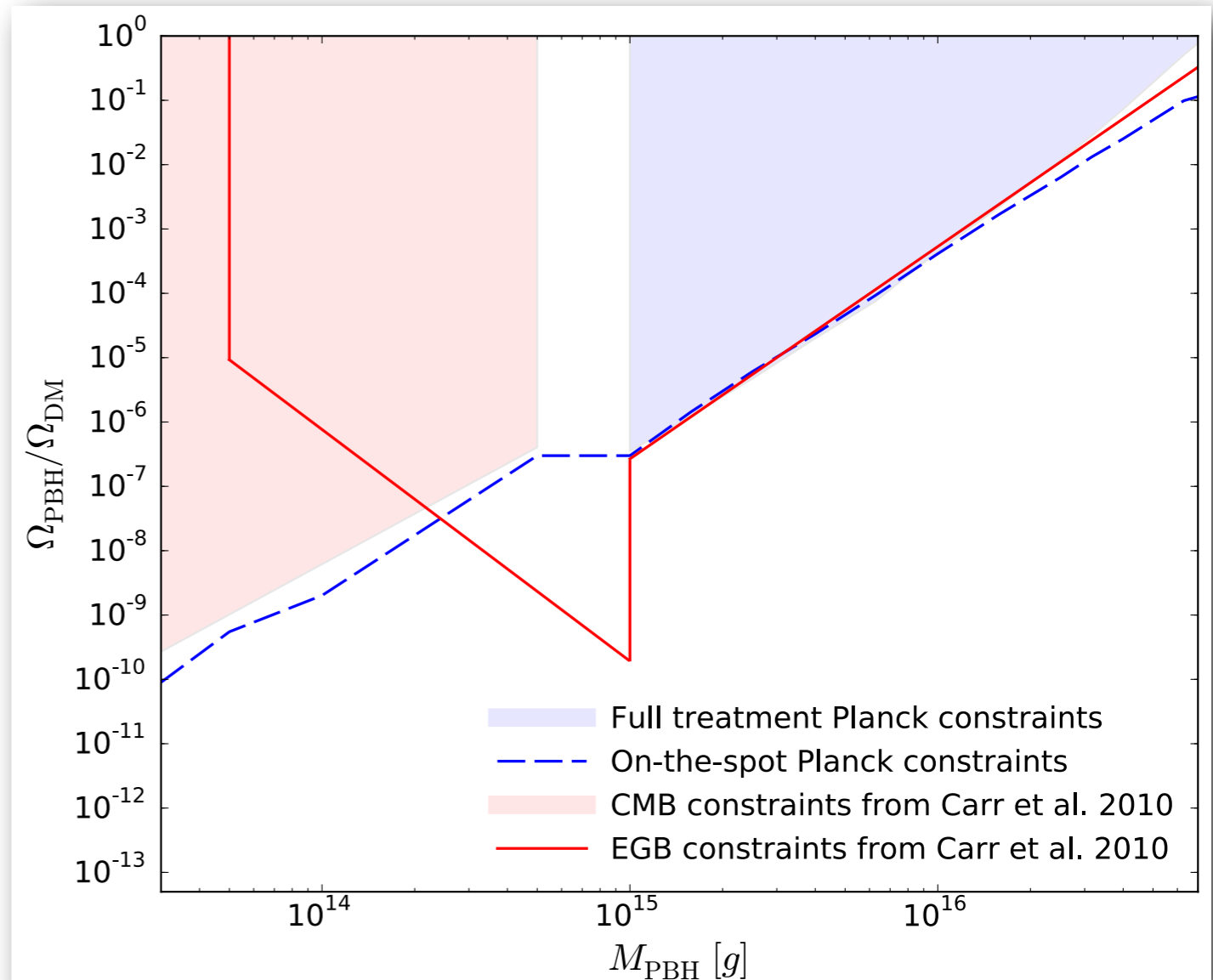
In particular, "light" PBH evaporation injects  $e^+e^-, \gamma \dots$  at a rate

$$\frac{dE}{dV dt} = \frac{\Omega_{\text{DM}} \rho_c c^2 (1+z)^3 f_{\text{PBH}}}{M_{\text{PBH}}^{\text{ini}}} \left. \frac{dM}{dt} \right|_{\text{e.m.}}$$

Peculiar modification of  $x_e$  possible

Bounds comparable or better than existing ones from diffuse gamma-ray background, for a certain range of masses

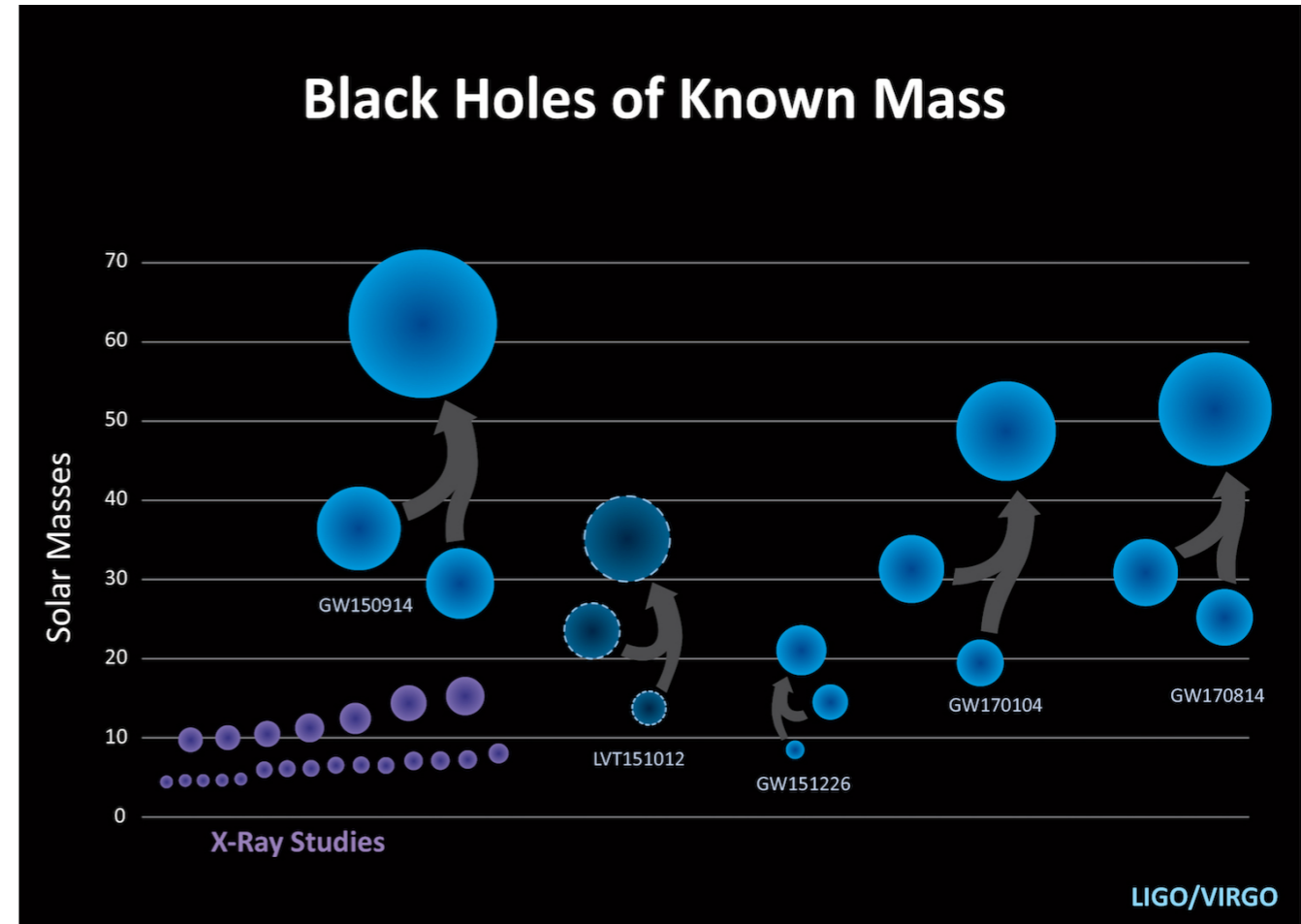
*V. Poulin, J. Lesgourgues and P.D.S., "Cosmological constraints on exotic injection of electromagnetic energy," JCAP 1703, 043 (2017) [1610.10051]*



# Did LIGO detect PBH dark matter?

LIGO/Virgo has detected relatively massive BH mergers, starting from the seminal (Nobel-prize worth!)

*B.P. Abbott et al. [LIGO & Virgo],  
PRL 116, 061102 (2016) [1602.03837]*



Hypothesis that they are primordial & explain (part of) the DM considered in several papers:

*S. Bird et al. "Did LIGO detect dark matter?," PRL 116, 201301 (2016) [1603.00464]*

*S. Clesse and J. García-Bellido, Phys. Dark Univ. 10, 002 (2016) [1603.05234]*

*M. Sasaki, T. Suyama, T. Tanaka, S. Yokoyama, PRL 117, 061101 (2016) [1603.08338]*

*K. Inomata, M. Kawasaki, K. Mukaida, Y. Tada, T.T. Yanagida, PRD 96, 043504 (2017) [1701.02544]*

...

**If true, this has consequences for the Cosmic Microwave Background!**

# Accreting PBH & CMB

Mass falling from “infinity to the BH” converts a sizable part of its potential energy into radiative emission, mostly X-rays: Up to 6%-40% of its mass energy, depending on BH spin, most efficient mechanism known in astrophysics! Invoked for powering Quasars, UHECRs, etc.

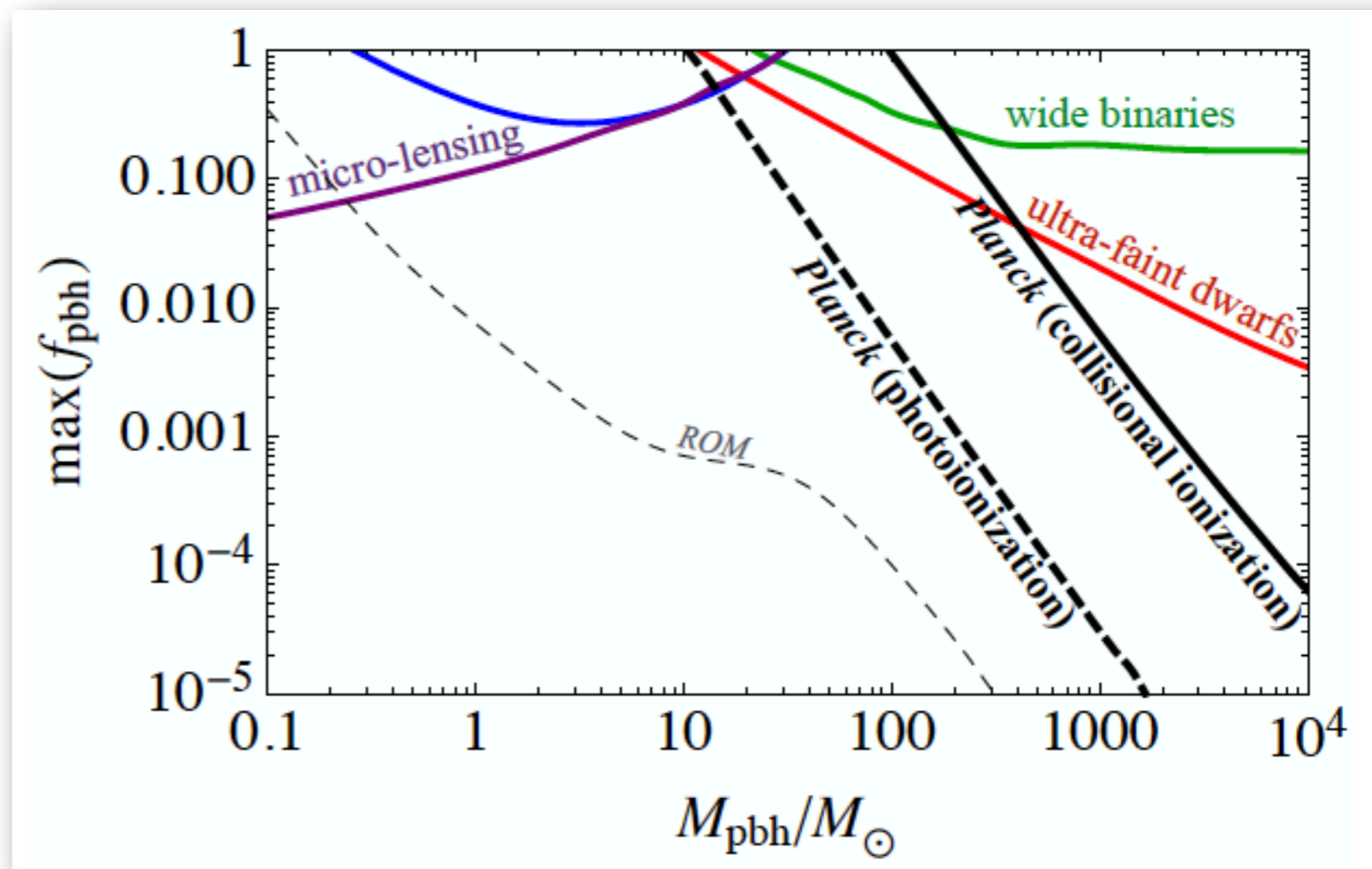
*For stellar mass PBH, cosmological gas accreting onto PBH radiates, affecting CMB*

Pioneering & stringent bounds obtained a decade ago (*Ricotti et al. 2008*) have been shown to be incorrect & inconsistent.

*Y. Ali-Haïmoud and M. Kamionkowski, “Cosmic microwave background limits on accreting primordial black holes,” PRD 95, 043534 (2017) [1612.05644]*

Conservative bounds for a *spherical accretion* flow yield

$f_{\text{PBH}} < 1$  for  $M > 10\text{-}100 M_{\odot}$



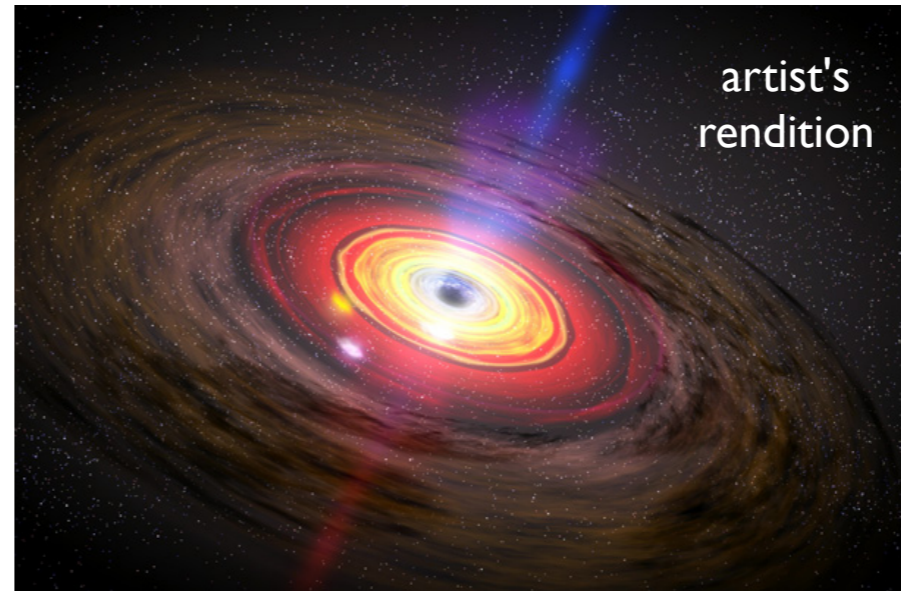
# Our Contribution

**Little problem:**

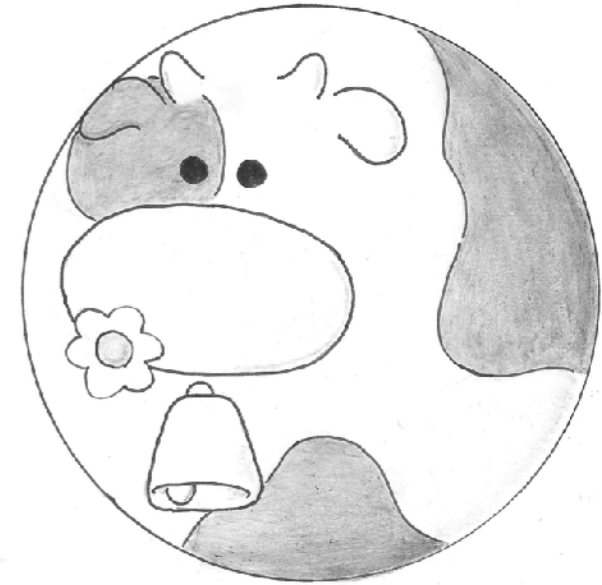
Nobody has ever seen a BH emission associated to a **spherical** accretion!

**Only disks!**

**Why? Is cosmology different?**  
Or it's a *spherical cow* approximation?



NASA/Dana Berry, SkyWorks Digital



by Christopher Berry

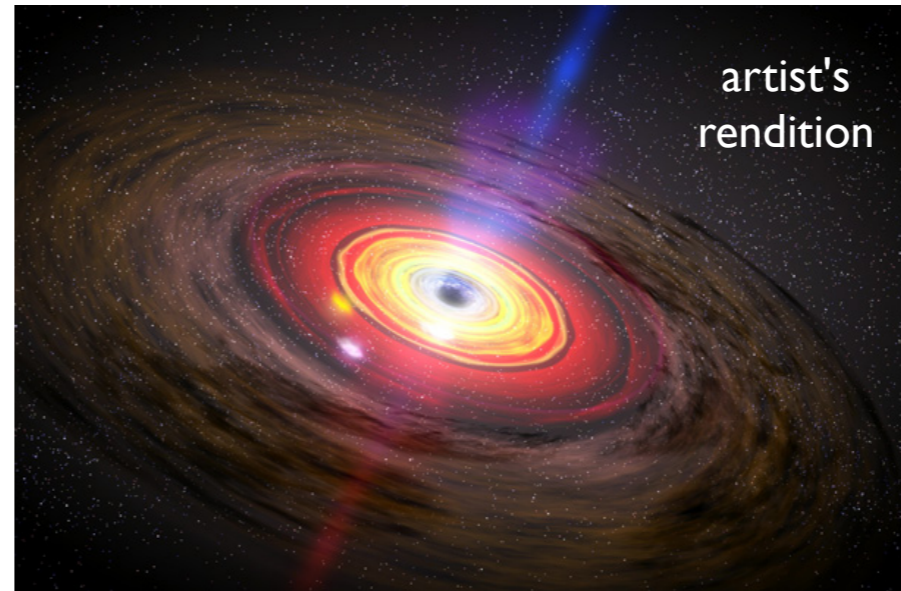
# Our Contribution

## Little problem:

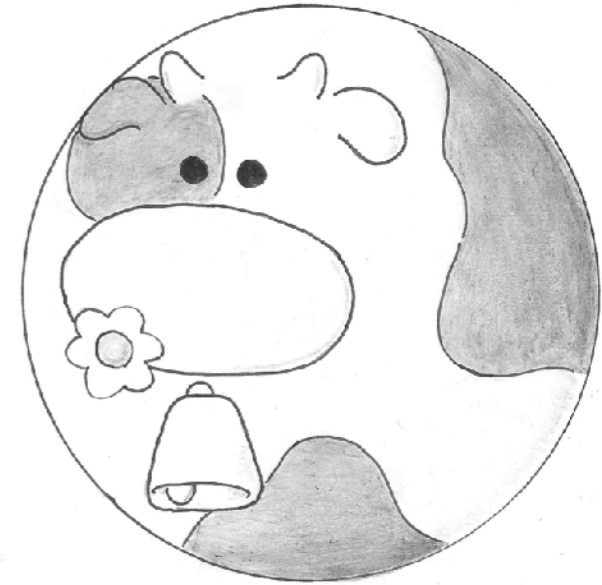
Nobody has ever seen a BH emission associated to a spherical accretion!

Only disks!

Why? Is cosmology different?  
Or it's a *spherical cow* approximation?



NASA/Dana Berry, SkyWorks Digital



by Christopher Berry

## Disk criterion:

If the accreted material has sufficient angular momentum wrt the BH to settle in Keplerian orbit at distance  $\gg 3 R_{\text{Schw}}$  then emission is disk-dominated (inner radii dominate the flux)

Criterion going back to *Shapiro & Lightman, ApJ 1976, Agol & Kamionkowski MNRAS 2002...*

The gas-PBH angular momentum cannot be computed exactly, since it depends upon non-linear physics. But several independent arguments (e.g. BH in binaries, supersonic motions) suggest that it is unavoidable to pass this threshold.

Also consistent with BH accretion *disks* being the only kind we've ever seen!

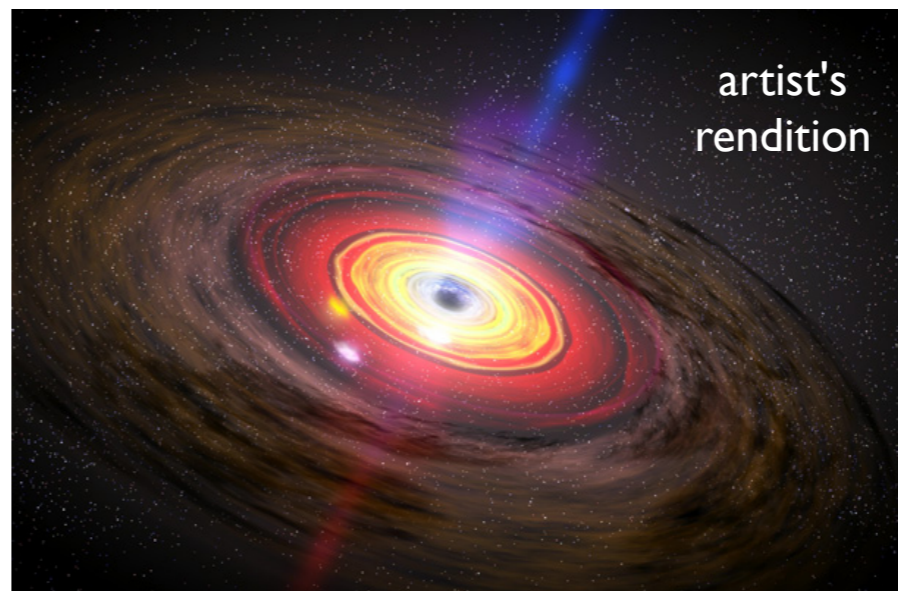
# Our Contribution

## Little problem:

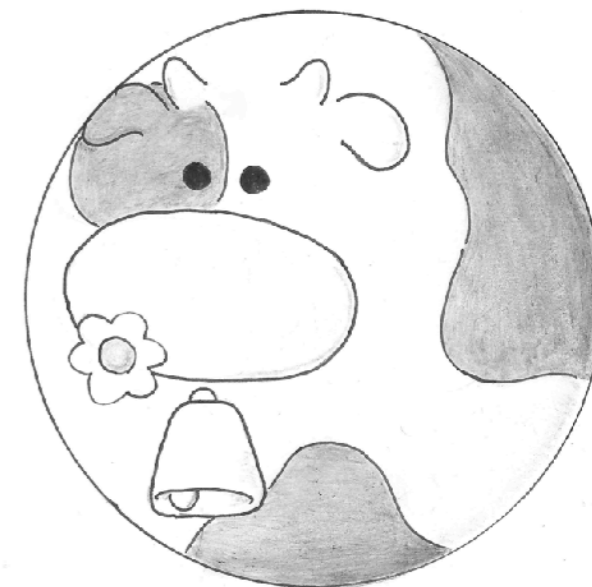
Nobody has ever seen a BH emission associated to a spherical accretion!

Only disks!

Why? Is cosmology different?  
Or it's a *spherical cow* approximation?



NASA/Dana Berry, SkyWorks Digital



by Christopher Berry

## Disk criterion:

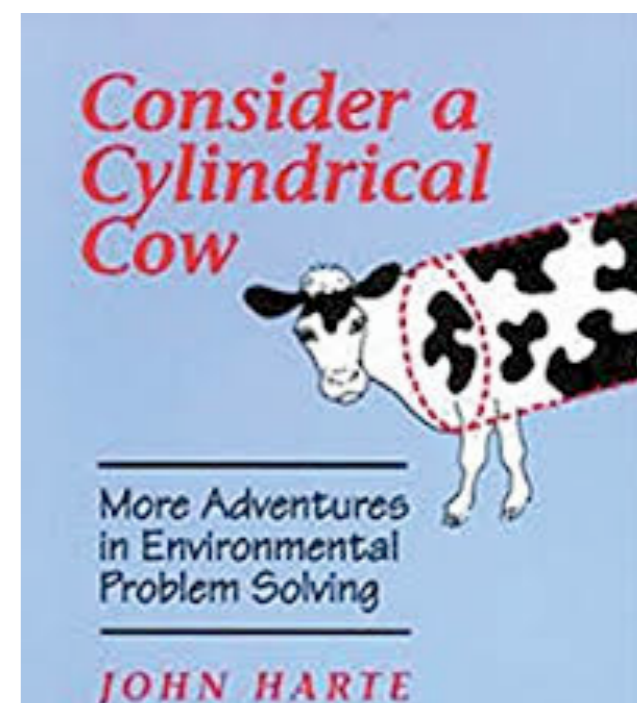
If the accreted material has sufficient angular momentum wrt the BH to settle in Keplerian orbit at distance  $\gg 3 R_{\text{Schw}}$  then emission is disk-dominated (inner radii dominate the flux)

Criterion going back to *Shapiro & Lightman, ApJ 1976, Agol & Kamionkowski MNRAS 2002...*

The gas-PBH angular momentum cannot be computed exactly, since it depends upon non-linear physics. But several independent arguments (e.g. BH in binaries, supersonic motions) suggest that it is unavoidable to pass this threshold.

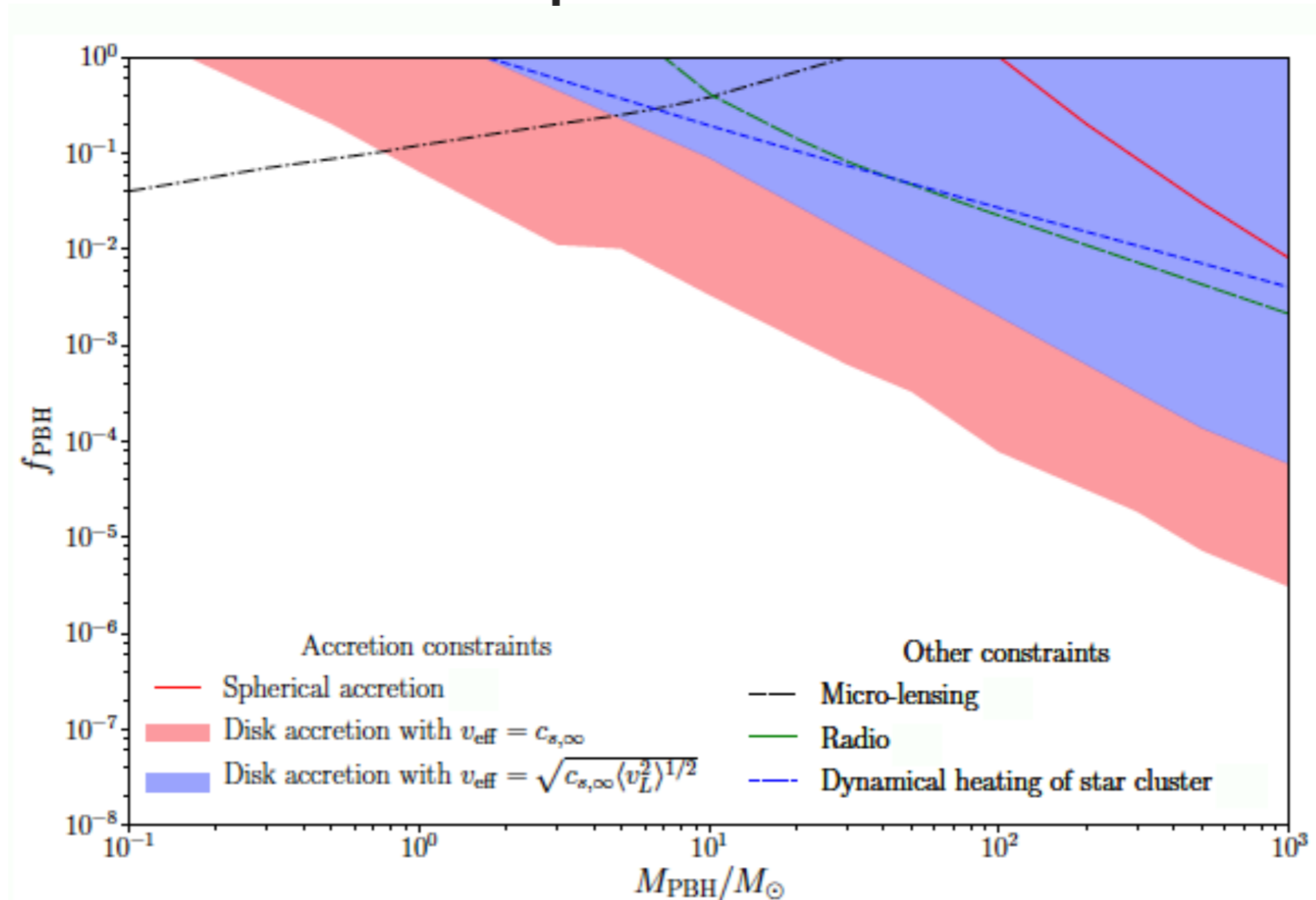
Also consistent with BH accretion *disks* being the only kind we've ever seen!

*V. Poulin, P. D. Serpico, F. Calore, S. Clesse and K. Kohri, "CMB bounds on disk-accreting massive primordial black holes," Phys. Rev. D 96, 083524 (2017) [1707.04206]*





# Results: disk vs. spherical accretion results



**Much stronger constraints**

CMB excludes PBH as totality of DM down to solar masses, when eventually lensing constraints take over.

V. Poulin, P. D. Serpico, F. Calore, S. Clesse and K. Kohri, “CMB bounds on disk-accreting massive primordial black holes,” *Phys. Rev. D* 96, 083524 (2017) [1707.04206]

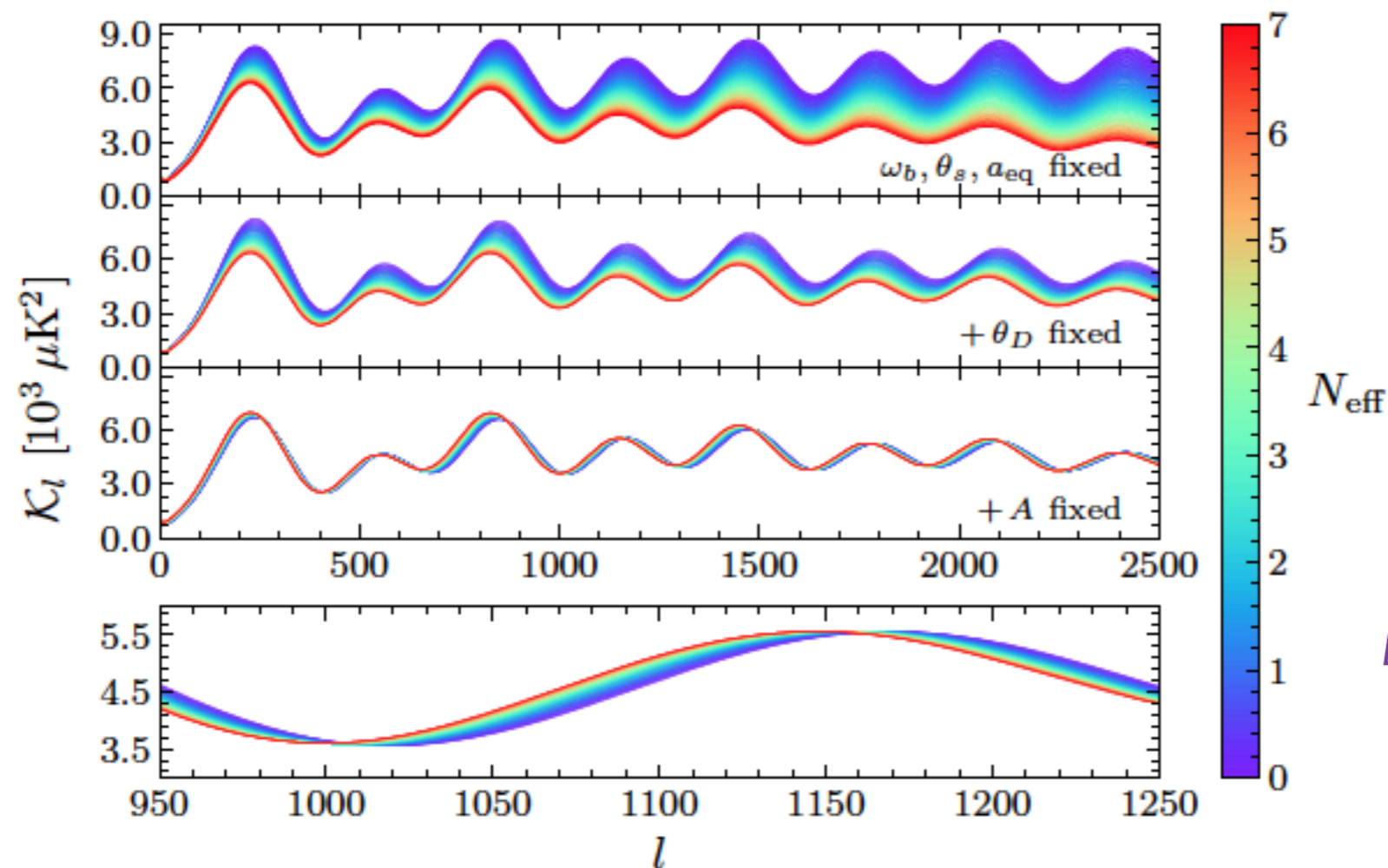
**Key message:** analytical toy models fail. State of the art “recipes” suggest strong bounds. To check and improve over them, BH accretion in a cosmo context requires dedicated simulations

# Episode III

on “dark radiation”: sensitivity and perspectives



# Effects of “dark radiation” (neutrino-like) on CMB



*D. Baumann, D. Green and B. Wallisch,  
JCAP1808, 08, 029 (2018)  
[1712.08067]*

If holding the matter-radiation equality fixed, more dark radiation  $\rightarrow$  increased damping (*1st panel*)

*Partially degenerate with  $Y_p$ , adjusting it one can fix the damping scale. Still, residual effect!*

Dark radiation free streaming causes anisotropic stress, leading in turn to

- amplitude shift at small scales (visible in the *2nd panel*)
- Phase shift of acoustic peaks at small scales (visible in *3rd panel*, renormalized at 4th peak; zoomed in *4th panel*)

# Standard model expectation and present sensitivity

$$\rho_{\text{DR}} = \frac{7}{8} \left( \frac{4}{11} \right)^{4/3} N_{\text{eff}} \rho_{\gamma}$$

$$N_{\text{eff}}^{\text{SM}} = 3.045 - 3.046 \text{ (th. \& num. error } \sim \mathcal{O}(0.001))$$

Non-instantaneous, momentum-dependent decoupling, accounting for finite temperature QED corrections (effective electron and photon masses, that in turn modify the equation of state of the plasma) + neutrino oscillations (act on last digit at most)

*G. Mangano, G. Miele, S. Pastor, T. Pinto, O. Pisanti and PDS, “Relic neutrino decoupling including flavor oscillations,” Nucl. Phys. B 729, 221 (2005) [hep-ph/0506164]*

*P. F. de Salas and S. Pastor, “Relic neutrino decoupling with flavour oscillations revisited,” JCAP 1607, 051 (2016) [1606.06986]*

*(new code, plus improved treatment of off-diagonal dumping terms in density matrix evolution)*

Planck+BAO only allow at 95% CL  $\Delta N_{\text{eff}} \equiv N_{\text{eff}} - N_{\text{eff}}^{\text{SM}} \lesssim 0.3$

*(Planck I, 2018, 1807.06205)*

disfavours at 95%CL any light, thermal relics that froze out after the QCD phase transition

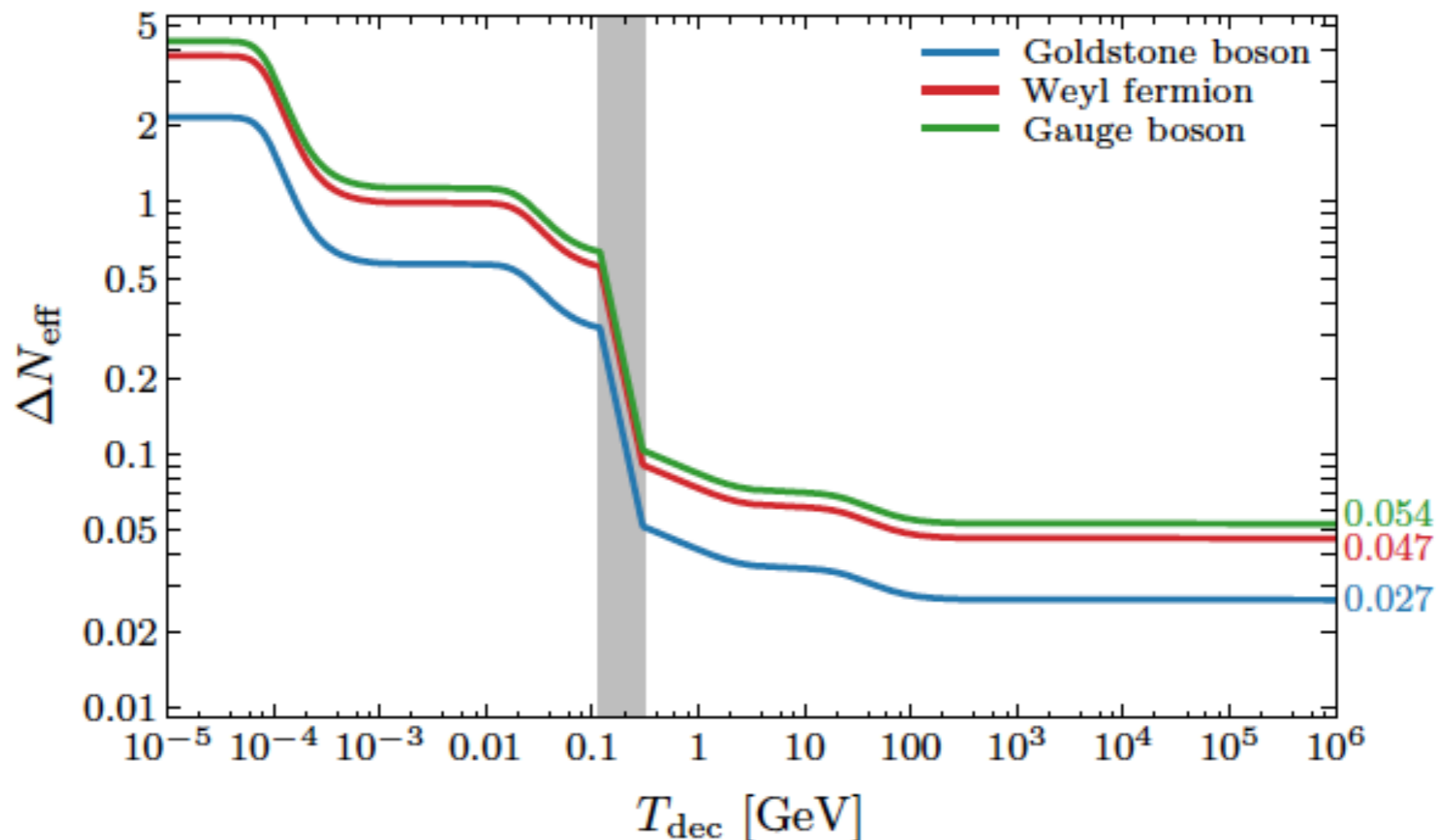
# Future CMB sensitivity to “dark radiation”

The estimated sensitivity of future ground-based CMB-S4 surveys is about five times better: Should be marginally sensitive to the non-instantaneous  $\nu$  decoupling & to the presence of any BSM relativistic thermal relic (no matter  $T_{\text{dec}}$  in the limit where the SM dofs are all there is...)

$$\sigma(N_{\text{eff}}) \simeq 0.03$$



*D. Baumann, D. Green and B. Wallisch, JCAP1808, 08, 029 (2018) [1712.08067]*



# summary and conclusions

- ▶ Cosmology (and CMB in particular) is sensitive to even extremely suppressed interaction rates of (meta)stable species present in the cosmic soup.
- ▶ The example of an **invisible** decay mode of (a fraction of) DM is noteworthy: For instance, it limits to  $<3.8\%$  the conversion of DM mass into “dark” radiation (like GW, low-E  $\nu$ 's...)
- ▶ If even a tiny fraction of the energy stored in the DM mass is released into “visible” (e.m.) form, CMB constraints can be quite tight (due to gas ionization and heating phenomena). DM annihilation, DM decay, evaporating PBH, accreting PBH are examples to which this can be applied
- ▶ CMB is also sensitive directly to relic relativistic species (dark radiation): currently excluding any relic decoupling after QCD phase transition, in the future testing for any relativistic relic decoupling before EW phase transition is within reach!

# summary and conclusions

- ▶ Cosmology (and CMB in particular) is sensitive to even extremely suppressed interaction rates of (meta)stable species present in the cosmic soup.
- ▶ The example of an **invisible** decay mode of (a fraction of) DM is noteworthy: For instance, it limits to  $<3.8\%$  the conversion of DM mass into “dark” radiation (like GW, low-E  $\nu$ 's...)
- ▶ If even a tiny fraction of the energy stored in the DM mass is released into “visible” (e.m.) form, CMB constraints can be quite tight (due to gas ionization and heating phenomena). DM annihilation, DM decay, evaporating PBH, accreting PBH are examples to which this can be applied
- ▶ CMB is also sensitive directly to relic relativistic species (dark radiation): currently excluding any relic decoupling after QCD phase transition, in the future testing for any relativistic relic decoupling before EW phase transition is within reach!

ευχαριστώ πολύ