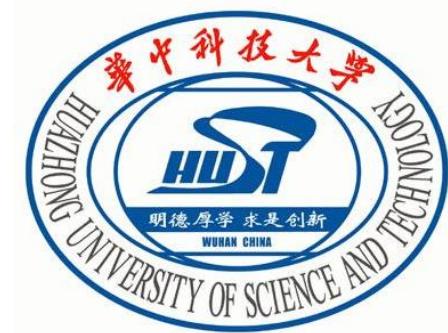


Relativistic effects in atom gravimeters

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2017.09.24 Corfu2017



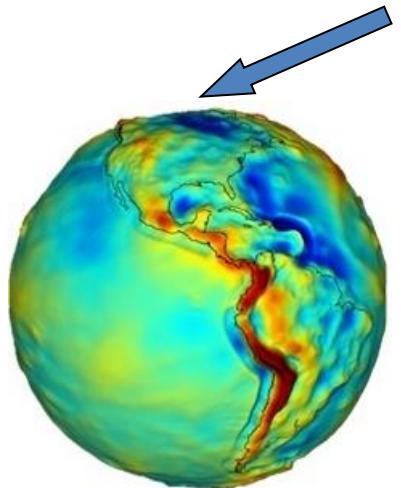
Outline

- Research motivation
- AG's Working principle
- An analytical study method
- Conclusion

1. Research motivation

● Applications of gravimeters

Gravimeter: measure the gravitational acceleration g



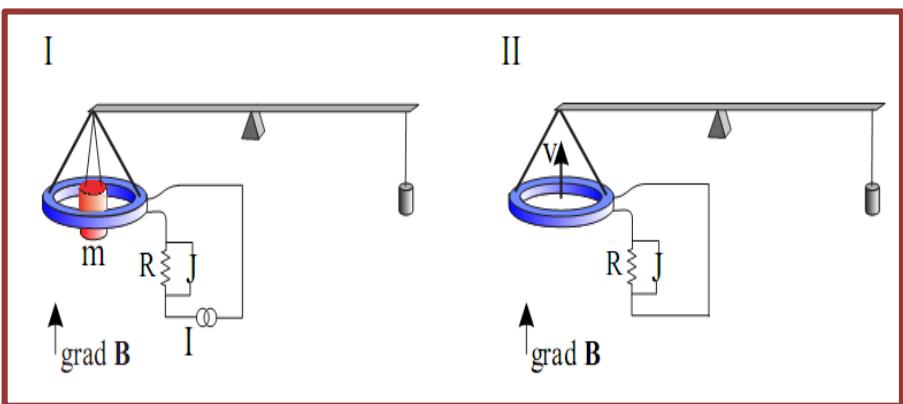
Geophysics



Resources exploration



Gravity navigation



Other scientific researches (such as Watt balance)

1. Research motivation

● Classification of gravimeters

Relative gravimeters



CG-5

Absolute gravimeters

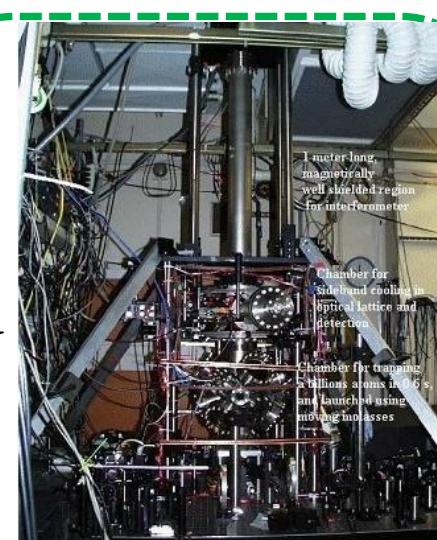


FG-5

Accuracy— μGal
 $1\mu\text{Gal} = 1 \times 10^{-8} \text{ m/s}^2$



GWR



AG

1. Research motivation

● Necessity for establishing a relativistic model of AG

High-precision gravimeter →

High-precision tide model;
High-precision geoid;
Post-Newton gravity theory;
.....

Measured gravity acceleration:

Newtonian effects

$$g_{\text{measured}} = g_0 + \text{temperature, pressure, Earth's rotation, gravity gradient....}$$

$$+ B \frac{v}{c}$$

+

$$\left[C \left(\frac{v}{c} \right)^2 \right]$$

+

$$D \frac{\phi}{c^2} + \dots$$

Finite speed of light (FSL) effect

$10 \mu\text{Gal}$

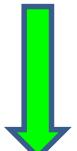
Second-order Doppler effect

$10^{-8} \mu\text{Gal}$

General Relativity

$10^{-7} \mu\text{Gal} (\Delta\phi/c^2)$

Relativistic effect



Model



Establish a more complete relativistic model for AG,
and improve the precision of AG.

1. Research motivation

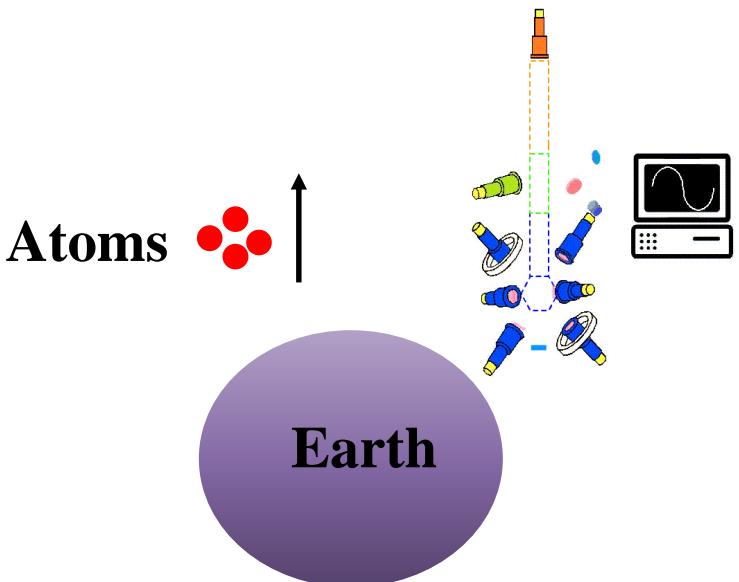
● Research status for the AG's relativistic model

Measured gravity acceleration:

$$g_{\text{measured}} = g_N + \boxed{\Delta g_{\text{PN}} + \Delta g_{\text{AG}}} \quad \text{Relativistic effects}$$

Post-Newton effect of gravitational field^[1]

Relativistic effects related to the instruments Δg_{AG}



[1] S.Wajima *et al.*, Physical Review D, 55 (1997)1964.

1. Research motivation

● Research status for the AG's relativistic model

$$g_{\text{measured}}^{\text{peters}} = g_0 \left[1 + 2 \frac{v(T)}{c} \frac{\alpha_1 + \alpha_2}{\alpha_1 - \alpha_2} \right] \quad [2-3]$$

[2] A. Peters *et al.*, Metrologia 38 (2001) 25-61.

[3] B. Cheng *et al.*, Phys. Rev. A 92 (2015) 063617.

$$g_{\text{measured}}^{\text{Dimopoulos}} = g_0 \left[1 + 3 \frac{v(T)}{c} \right] \quad [4]$$

Disagree!

FSL effect

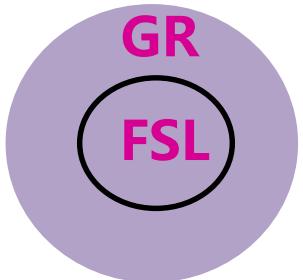
GR effects

[4] S. Dimopoulos *et al.*, PRD 78 (2008) 042003.

Scalar expression!

Need to explore:

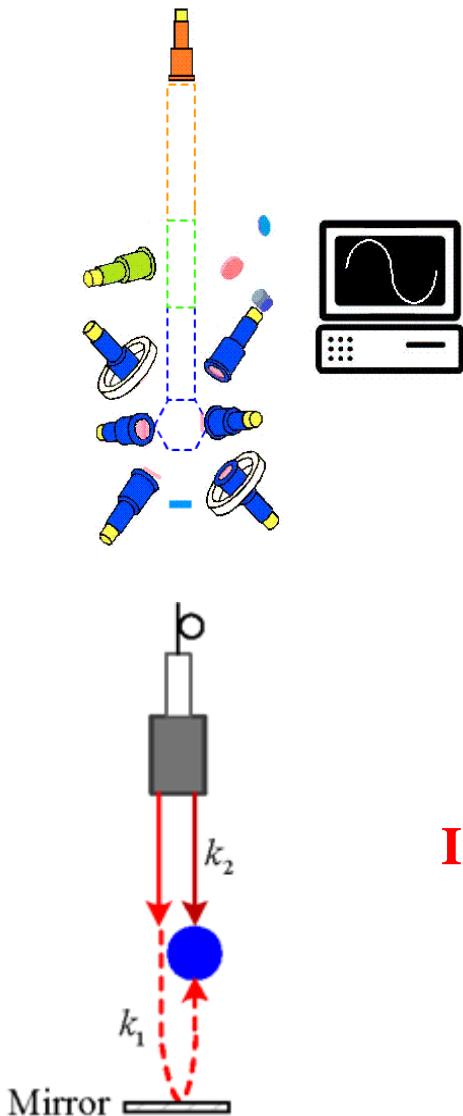
- Disagreement for FSL effect;
- More complete relativistic model



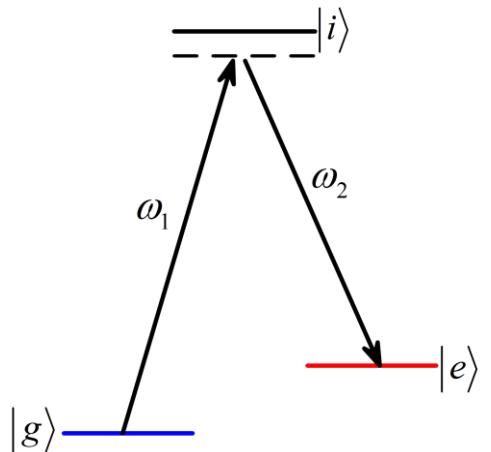
	GR phase shift	Size (rad)	Interpretation
1.	$-k_{\text{eff}} g T^2$	3×10^8	Newtonian gravity
2.	$-k_{\text{eff}} (\partial_r g) v_L T^3$	-2×10^3	1st gradient
3.	$-\frac{7}{12} k_{\text{eff}} (\partial_r g) T^4$	9×10^2	
4.	$-3k_{\text{eff}} g^2 T^3$	-4×10^1	
5.	$-3k_{\text{eff}} g v_L T^2$	4×10^1	finite speed of light and Doppler shift corrections
6.	$-\frac{k_{\text{eff}}^2}{2m} (\partial_r g) T^3$	-7×10^{-1}	1st gradient recoil
7.	$(\omega_{\text{eff}} - \omega_a) g T^2$	-4×10^{-1}	detuning
8.	$(2 - 2\beta - \gamma) k_{\text{eff}} g \phi T^2$	-2×10^{-1}	GR (nonlinearity)
9.	$-\frac{3k_{\text{eff}}^2}{2m} g T^2$	2×10^{-2}	
10.	$-\frac{7}{12} k_{\text{eff}} v_L^2 (\partial_r^2 g) T^4$	8×10^{-3}	2nd gradient
11.	$-\frac{35}{4} k_{\text{eff}} (\partial_r g) g v_L T^4$	6×10^{-4}	
12.	$-4k_{\text{eff}} (\partial_r g) v_L^2 T^3$	-3×10^{-4}	
13.	$2\omega_a g^2 T^3$	2×10^{-4}	
14.	$2\omega_a g v_L T^2$	-2×10^{-4}	
15.	$-\frac{7k_{\text{eff}}^2}{12m} v_L (\partial_r^2 g) T^4$	7×10^{-6}	2nd gradient recoil
16.	$-12k_{\text{eff}} g^2 v_L T^3$	-7×10^{-6}	
17.	$-7k_{\text{eff}} g^3 T^4$	4×10^{-6}	
18.	$-5k_{\text{eff}} g v_L^2 T^2$	3×10^{-6}	
19.	$(2 - 2\beta - \gamma) k_{\text{eff}} \partial_r(g\phi) v_L T^3$	2×10^{-6}	GR (velocity-dependent force)
20.	$\frac{7}{12} (4 - 4\beta - 3\gamma) k_{\text{eff}} \phi (\partial_r g) g T^4$	-2×10^{-6}	GR 1st gradient
			GR

2. AG's Working principle

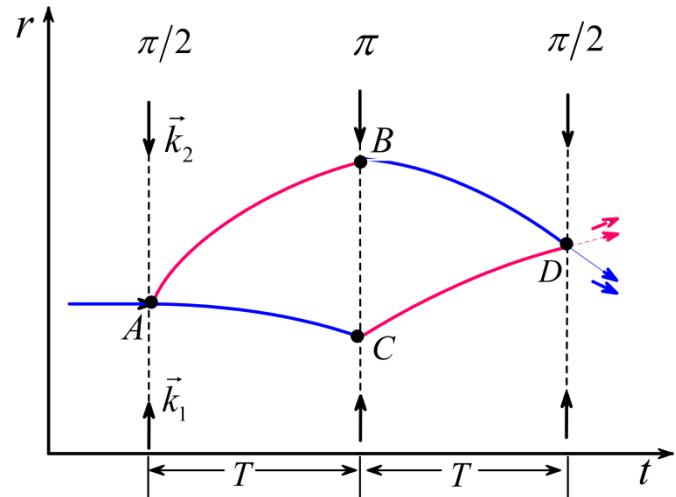
Three-Raman-laser-pulse
atom gravimeter



Rabi oscillation



Split, reflect, recombine
atomic wave packets



Interference phase



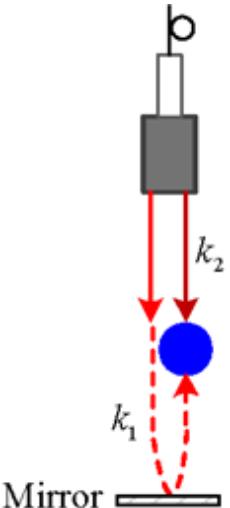
gravitational acceleration

3. An analytical study method

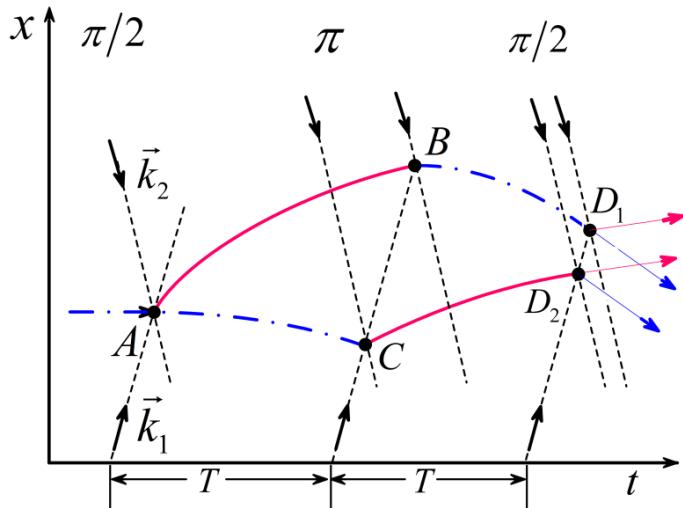
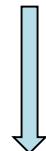
◆ Conventional calculating method^[4]

Solve geodesic equation

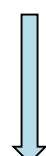
$$\frac{d^2x^\mu}{d\tau^2} + \Gamma_{\alpha\beta}^\mu \frac{dx^\alpha}{d\tau} \frac{dx^\beta}{d\tau} = 0$$



Trajectories of atoms and photons

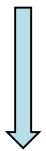


5 intersections



Path integral

Phase shift $\Delta\phi_{\text{tot}} = \boxed{\Delta\phi_{\text{propagation}}} + \Delta\phi_{\text{laser}} + \Delta\phi_{\text{separation}}$

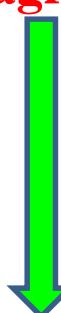


Gravitational acceleration

[4] S. Dimopoulos *et al.*, PRD 78 (2008) 042003.

Refer the space-time diagram

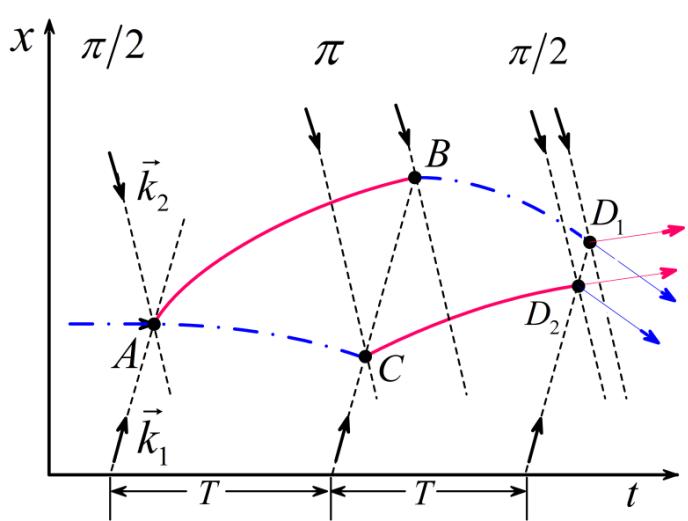
$$\int_{t_A}^{t_B} L_{\text{upper}} dt - \int_{t_A}^{t_C} L_{\text{lower}} dt$$



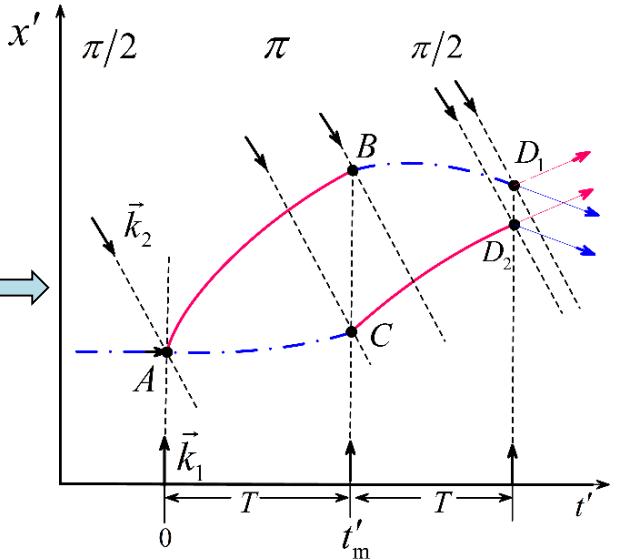
3. An analytical study method

◆ An analytical study method

Idea: transfer FSL effect into atom's Lagrangian



Control light
Coordinate transformation



$$\int_{t_A}^{t_B} L_{\text{upper}} dt - \int_{t_A}^{t_C} L_{\text{lower}} dt \Rightarrow \int_0^{t_m'} (L'_{\text{upper}} - L'_{\text{lower}}) dt' \quad \text{simplify the calculation}$$

After transformation:

- Motion equation of light \vec{k}_1 :

$$dx'/dt' = \infty$$

- Lagrangian of atoms: $L' = L'_0 + L'_1$

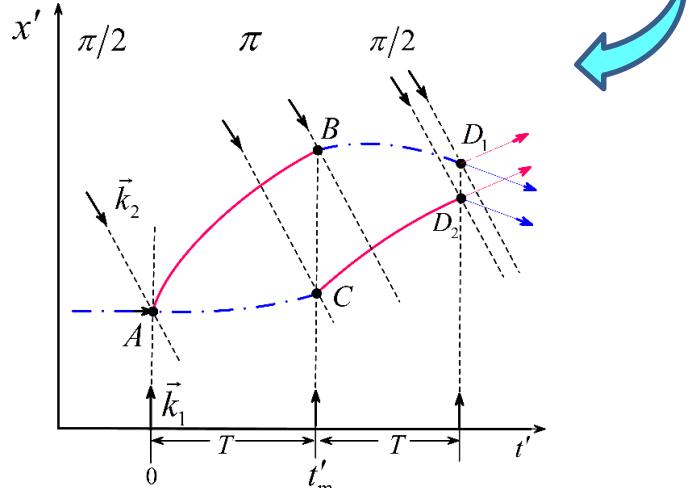
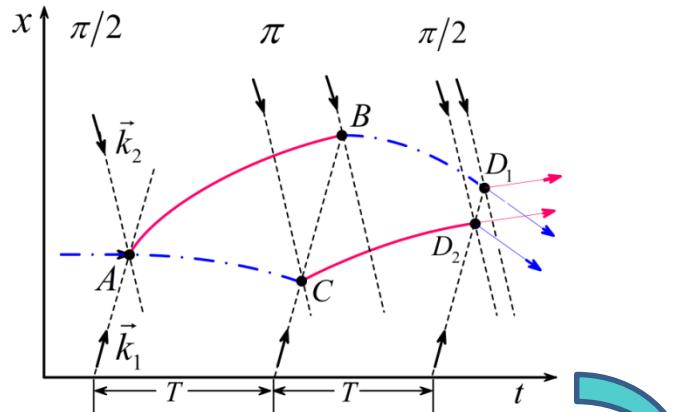
FSL disturbance
Original form

3. An analytical study method

◆ Calculating the FSL effect

Calculation in Special Relativity frame

1. Coordinate transformation:



① Solve the motion equations of AG

photon: $t = \frac{\vec{n} \cdot \vec{x}}{c}$

atom: $L = \frac{1}{2} m \vec{v}^2 + m \vec{g} \cdot \vec{r}$

② Coordinate transformation for \vec{k}_1

$$\begin{cases} \vec{x}' = \vec{x} \\ t' = t - \frac{\vec{n} \cdot \vec{x}}{c} \end{cases}$$

③ Motion equations of AG after transformation

Photon (\vec{k}_1): $dx' / dt' = \infty$

atom: $L' = L'_0 + L'_1$

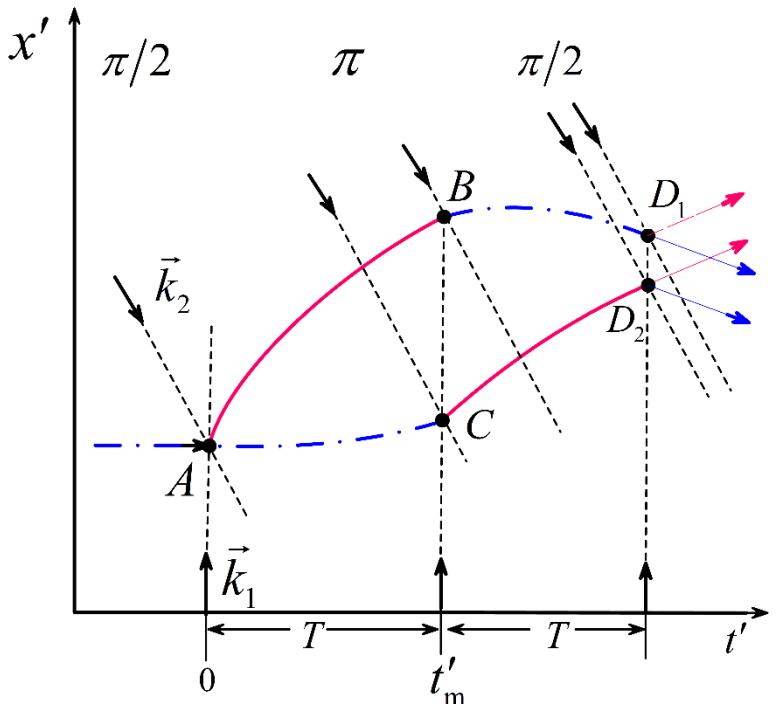
Original form

FSL disturbance

3. An analytical study method

◆ Calculating the FSL effect

2. Total phase shift :



ABCD Matrix, “Perturbation” approach, Laser phase

$$\Delta\phi_{\text{tot}} = \Delta\phi_{\text{propagation}} + \Delta\phi_{\text{separation}} + \Delta\phi_{\text{laser}} = \boxed{\Delta\phi_{L'_0+\text{separation}} + \Delta\phi_{L'_1} + \Delta\phi_{\text{laser}}}$$

3. An analytical study method

◆ Calculating the FSL effect

3. FSL effect in AG :

Measured gravitational acceleration^[5]:

$$g = g_0 \left[1 + 3 \frac{\vec{v}(T) \cdot \vec{e}_k}{c} - 2 \frac{\vec{v}(T) \cdot \vec{e}_k}{c} \frac{\alpha_1 - \alpha_2}{\vec{k}_{\text{eff}} \cdot \vec{g}_0} + \frac{2\vec{v}(T)}{c} \cdot \frac{\alpha_1 \vec{n}_1 - \alpha_2 \vec{n}_2}{\vec{k}_{\text{eff}} \cdot \vec{g}_0} \right]$$

$$g_{\text{measured}}^{\text{Dimopoulos}} = g_0 \left[1 + 3 \frac{v(T)}{c} \right]$$

Time delay

missed
Coupling of time delay
and frequency chirp

$$g_{\text{measured}}^{\text{peters}} = g_0 \left[1 + 2 \frac{v(T)}{c} \frac{\alpha_1 + \alpha_2}{\alpha_1 - \alpha_2} \right]$$

frequency chirp

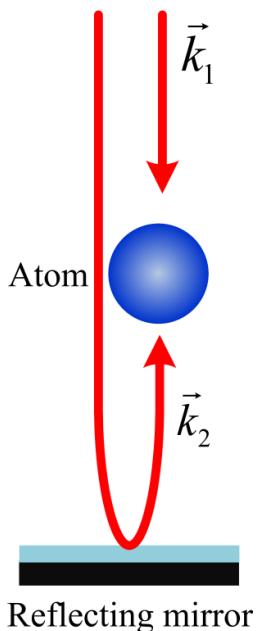
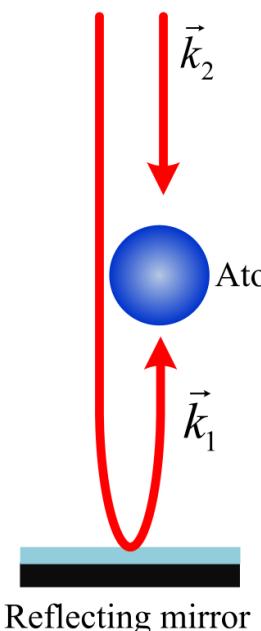
3. An analytical study method

◆ Calculating the FSL effect

$$g = g_0 \left[1 + \frac{\vec{v}(T) \cdot \vec{e}_k}{c} + \frac{2\vec{v}(T)}{c} \cdot \frac{\alpha_1 \vec{n}_1 - \alpha_2 \vec{n}_2}{(\vec{k}_{\text{eff}} \cdot \vec{g}_0)} \right] = g_0 (1 + A + B)$$

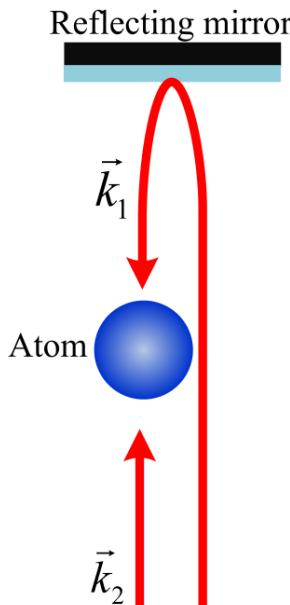
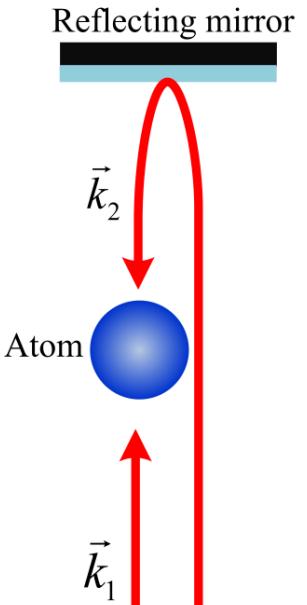
\vec{e}_k : direction of the control light (the light is reflected by the mirror)

\vec{k}_{eff} : $\vec{k}_1 - \vec{k}_2$ effective wave vector



①② or ③④ B
 ①③ or ②④ A
 ①④ or ②③ A+B

Experimental design



①
 $\vec{e}_k \uparrow \vec{k}_{\text{eff}} \uparrow$

②
 $\vec{e}_k \uparrow \vec{k}_{\text{eff}} \downarrow$

③
 $\vec{e}_k \downarrow \vec{k}_{\text{eff}} \uparrow$

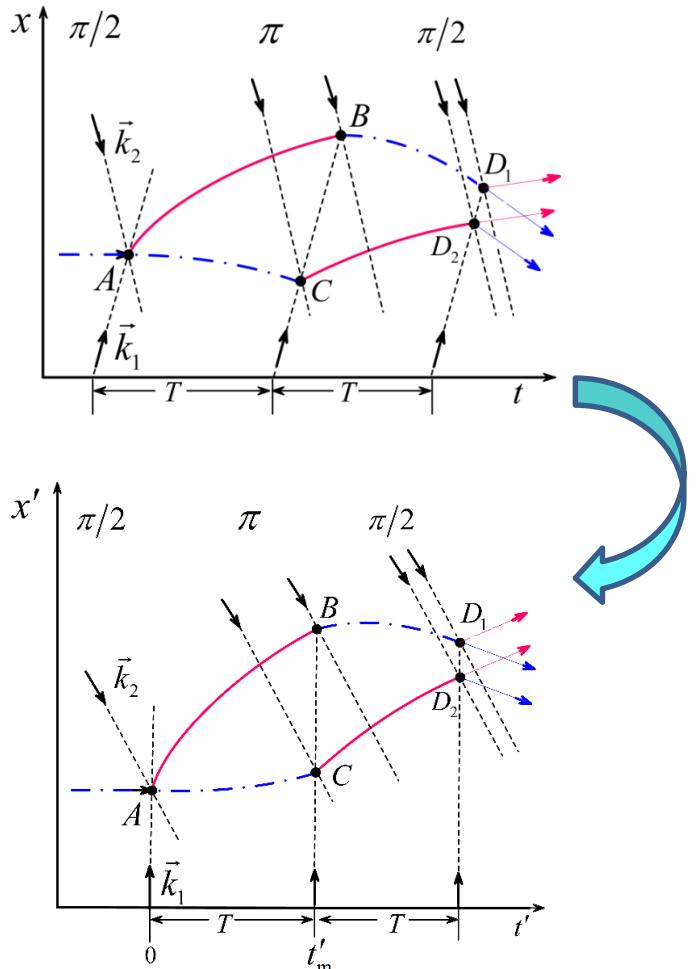
④
 $\vec{e}_k \downarrow \vec{k}_{\text{eff}} \downarrow$

3. An analytical study method

◆ calculating the relativistic effects

Calculation in General Relativity frame

1. Coordinate transformation:



① Solve the motion equations of AG

photon: $t = f(\vec{x})$

atom: L

$$ds^2 = - \left[1 + \frac{2\phi}{c^2} + \frac{2\beta\phi^2}{c^4} - \left(1 - \frac{2\gamma\phi}{c^2} \right) \frac{\Omega^2 r^2 \sin^2 \theta}{c^2} \right] (cdt)^2$$

$$+ \left(1 - \frac{2\gamma\phi}{c^2} \right) 2\Omega r^2 \sin^2 \theta dt d\varphi_r$$

$$+ \left(1 - \frac{2\gamma\phi}{c^2} \right) (dr^2 + r^2 d\theta^2 + r^2 \sin^2 \theta d\varphi_r^2)$$

② Coordinate transformation for \vec{k}_1

$$\begin{cases} t' = t - f(\vec{x}) \\ \vec{x}' = \vec{x} \end{cases}$$

③ Motion equations of AG after Coordinate transformation

Photon (\vec{k}_1): $dx'/dt' = \infty$

atom: $L' = L'_0 + L'_1$

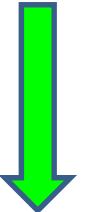
3. An analytical study method

◆ calculating the relativistic effects

2. Relativistic model for AG :

Relativistic phase shift^[6] :

$$\Delta\phi_{\text{tot}} = \Delta\phi_{\text{Newtonian}} + \Delta\phi_{\text{relativistic}}^{\Omega-\text{independent}} + \Delta\phi_{\text{relativistic}}^{\Omega-\text{dependent}}$$



Keep the sensitivity of gravitational acceleration within 10^{-15} g

3. An analytical study method

TABLE I: A list of the non-relativistic phase shifts of the three-pulse interferometer, $\Delta\phi_{\text{classical}}$. Column 2, vector expression of each term; column 3, a rough order estimate for the terms in column 2, and the relative sizes of each term to $k_{\text{eff}}a_0T^2$ are given, with the mentioned design of sensitive to acceleration $\sim 10^{-15}a_0$, and the 2nd gradient $\partial_r\gamma_0=1.4\times 10^{-12}(\text{ms}^2)^{-1}$; column 4, the interpretation for the terms.

Non-relativistic phase shift	Relative size	Interpretation
1 $-T^2\vec{k}_{\text{eff}}\cdot\vec{a}_0$	1	Acceleration due to gravity
2 $-2T^2\vec{k}_{\text{eff}}\cdot[(\vec{v}_0+\vec{a}_0T)\times\vec{\Omega}]$	$v_0\Omega/a_0$	$\sim 1\times 10^{-4}$ Coriolis acceleration
3 $-T^3\vec{k}_{\text{eff}}\vec{\gamma}_0(\vec{v}_0+\frac{7T^4}{12}\vec{a}_0T)$	γ_0T^2	$\sim 5\times 10^{-6}$ 1st gradient
4 $-3T^3\{(\vec{k}_{\text{eff}}\cdot\vec{\Omega})[\vec{\Omega}\cdot(\vec{v}_0+\frac{7T^4}{12}\vec{a}_0T)]\}$ + $3T^3\{\Omega^2[\vec{k}_{\text{eff}}\cdot(\vec{v}_0+\frac{7T^4}{12}\vec{a}_0T)]\}$	$(\Omega T)^2$	$\sim 9\times 10^{-9}$
5 $-\frac{1}{6}T^4[(\vec{\Omega}\times\vec{k}_{\text{eff}})\vec{\gamma}_0(7\vec{v}_0+3\vec{a}_0T)]$ - $\frac{1}{6}T^4\{\vec{k}_{\text{eff}}\vec{\gamma}_0[(7\vec{v}_0+3\vec{a}_0T)\times\vec{\Omega}]\}$	$\gamma_0T^2\Omega T$	$\sim 5\times 10^{-10}$ Couple (1st gradient and earth's rotation)
6 $\frac{7}{12}T^4\nabla_i\nabla_j\nabla_k\phi_0(\vec{k}_{\text{eff}})_i(\vec{v}_0)_j(\vec{v}_0)_k$ + $\frac{3}{4}T^5\nabla_i\nabla_j\nabla_k\phi_0(\vec{a}_0)_i(\vec{v}_0)_j(\vec{k}_{\text{eff}})_k$ + $\frac{31}{120}T^6\nabla_i\nabla_j\nabla_k\phi_0(\vec{a}_0)_i(\vec{a}_0)_j(\vec{k}_{\text{eff}})_k$	$(\partial_r\gamma_0)v_0T^3$	$\sim 4\times 10^{-11}$ 2nd gradient
7 $-\frac{1}{360}T^5\vec{k}_{\text{eff}}\vec{\gamma}_0^2(90\vec{v}_0+31\vec{a}_0T)$	$(\gamma_0T^2)^2$	$\sim 3\times 10^{-11}$ Nonlinearity of the 1st gradient
8 $\frac{1}{3}T^4\Omega^2\vec{k}_{\text{eff}}\cdot[(7\vec{v}_0+3\vec{a}_0T)\times\vec{\Omega}]$	$(\Omega T)^3$	$\sim 9\times 10^{-13}$
9 $-\frac{1}{120}T^5[(\vec{k}_{\text{eff}}\cdot\vec{\Omega})\vec{\Omega}\vec{\gamma}_0(90\vec{v}_0+31\vec{a}_0T)]$ + $\frac{1}{60}T^5[\Omega^2\vec{k}_{\text{eff}}\vec{\gamma}_0(90\vec{v}_0+31\vec{a}_0T)]$ - $\frac{1}{90}T^5\{(\vec{\Omega}\times\vec{k}_{\text{eff}})\vec{\gamma}_0[(90\vec{v}_0+31\vec{a}_0T)\times\vec{\Omega}]\}$ - $\frac{1}{120}T^5\{(\vec{k}_{\text{eff}}\vec{\gamma}_0\vec{\Omega})[\vec{\Omega}\cdot(90\vec{v}_0+31\vec{a}_0T)]\}$	$\gamma_0T^2(\Omega T)^2$	$\sim 5\times 10^{-14}$

Have been given by others!

3. An analytical study method

TABLE II: A list of the earth rotation-independence relativistic phase shifts $\Delta\phi_{\text{relativistic}}^{\Omega-\text{independent}}$ of the three-pulse interferometer (with $\omega_1-\omega_2 = 6.8 \text{ GHz}$).

	Relativistic phase shift	Relative size	Interpretation
1	$-[3\vec{k}_{\text{eff}} \cdot \vec{a}_0 + 4\alpha_2](v_n + a_n T)T^2/c$	v_0/c	$\sim 4 \times 10^{-8}$ Finite speed of light
2	$2(1-\beta)\vec{k}_{\text{eff}} \cdot \vec{a}_0 \phi_0 T^2/c^2$	ϕ_0/c^2	$\sim 7 \times 10^{-10}$ Vanish for GR ($\beta=1$)
3	$(-2\vec{k}_{\text{eff}} \vec{\gamma}_0 \vec{v}_0 v_n + k_{\text{eff}} \vec{v}_0 \vec{\gamma}_0 \vec{v}_0)T^3/c$	$\frac{\gamma_0 v_0 T}{a_0} \frac{v_0}{c}$	$\sim 2 \times 10^{-13}$ SR 1st gradient
4	$-\frac{7}{12}[(\vec{k}_{\text{eff}} \times \vec{n}_0 \times \vec{a}_\perp) \vec{\gamma}_0 \vec{v}_0 + (\vec{k}_{\text{eff}} \vec{\gamma}_0 \times \vec{n}_0 \times \vec{a}_\perp) \vec{v}_0]T^4/c$	$\gamma_0 T^2 \frac{v_0}{c}$	$\sim 2 \times 10^{-13}$ SR 1st gradient
	$-\frac{7}{12}[\vec{k}_{\text{eff}} \vec{\gamma}_0 \vec{a}_0 v_n + 2\vec{k}_{\text{eff}} \vec{\gamma}_0 \vec{v}_0 a_n - 3k_{\text{eff}} \vec{a}_0 \vec{\gamma}_0 \vec{v}_0]T^4/c$		
5	$-\frac{1}{4}[(\vec{k}_{\text{eff}} \vec{\gamma}_0 \times \vec{n}_0 \times \vec{a}_\perp) \cdot \vec{a}_0 + (\vec{k}_{\text{eff}} \times \vec{n}_0 \times \vec{a}_\perp) \vec{\gamma}_0 \vec{a}_0]T^5/c$	$\gamma_0 T^2 \frac{a_0 T}{c}$	$\sim 2 \times 10^{-13}$ GR 1st gradient
	$+\frac{3}{4}[k_{\text{eff}} \vec{a}_0 \vec{\gamma}_0 \vec{a}_0 - \vec{k}_{\text{eff}} \vec{\gamma}_0 \vec{a}_0 a_n]T^5/c$		
6	$(\omega_1 - \omega_2)[\vec{a}_\perp \cdot \vec{v}_\perp + 2a_n v_n]T^2/c^2$	$\frac{\omega_1 - \omega_2}{\omega_1 + \omega_2} \frac{v_0}{c}$	$\sim 6 \times 10^{-14}$ SR
7	$(\omega_1 - \omega_2)(a_\perp^2 + 2a_n^2)T^3/c^2$	$\frac{\omega_1 - \omega_2}{\omega_1 + \omega_2} \frac{a_0 T}{c}$	$\sim 6 \times 10^{-14}$ GR
8	$\frac{1}{3}(1-\beta)\vec{k}_{\text{eff}} \vec{\gamma}_0 (6\vec{v}_0 + 7\vec{a}_0 T)\phi_0 T^3/c^2$	$\gamma_0 T^2 \frac{\phi_0}{c^2}$	$\sim 4 \times 10^{-14}$ Vanish for GR ($\beta=1$)
9	$\left\{(1+\gamma)\vec{k}_{\text{eff}} \cdot \vec{a}_n (v_n)^2 + (\vec{k}_{\text{eff}} \cdot \vec{a}_0)[(1-\gamma)(v_0)^2 + \frac{1}{2}(v_n)^2]\right\}T^2/c^2$	$(v/c)^2$	$\sim 2 \times 10^{-15}$ GR
	$+(\vec{k}_{\text{eff}} \cdot \vec{v}_0) \{[3 + (2\gamma + 1)]\vec{a}_0 \cdot \vec{v}_0 + 2a_n v_n\} T^2/c^2$		
10	$-[\gamma(\vec{k}_{\text{eff}} \cdot \vec{a}_n)(\vec{a}_\perp - \gamma\vec{a}_n) \cdot \vec{v}_0 + 3(1+\gamma)(\vec{k}_{\text{eff}} \cdot \vec{a}_n)a_n v_n]T^3/c^2$	$a_0 v_0 T/c^2$	$\sim 2 \times 10^{-15}$ GR
	$+(\vec{k}_{\text{eff}} \cdot \vec{a}_0)[2(\gamma + \beta + 3)\vec{a}_0 \cdot \vec{v}_0 + 4a_n v_n]T^3/c^2$		
	$+(\vec{k}_{\text{eff}} \cdot \vec{v}_0)[(2\gamma + 4)a_0^2 + 2a_n^2]T^3/c^2$		
11	$\frac{7}{12}[(2\beta + 4\gamma + 10)(\vec{k}_{\text{eff}} \cdot \vec{a}_0)a_0^2 + 6(\vec{k}_{\text{eff}} \cdot \vec{a}_0)a_n^2]T^4/c^2$	$(a_0 T/c)^2$	$\sim 2 \times 10^{-15}$ GR
	$-\frac{7}{12}\vec{k}_{\text{eff}} \cdot \vec{a}_n [\gamma(a_\perp^2 - \gamma a_n^2) + 3(1+\gamma)a_n^2]T^4/c^2$		

More general expression !

3. An analytical study method

TABLE III: A list of the earth rotation-dependence relativistic phase shifts $\Delta\phi_{\text{relativistic}}^{\Omega-\text{dependent}}$ of the three-pulse interferometer (with $\omega_1 - \omega_2 = 6.8$ GHz).

Relativistic phase shift	Relative size	Interpretation
1 $2\vec{k}_{\text{eff}} \cdot (\Omega \times \vec{v}_0) v_n T^2 / c$	$\frac{\Omega v_0}{a_0} \frac{v_0}{c} \sim 4 \times 10^{-12}$	SR coriolis acceleration
2 $\vec{k}_{\text{eff}} \cdot \left\{ (\vec{a}_\perp - \gamma \vec{a}_n) (\vec{n}_0 \times \vec{\Omega}) + (\vec{n}_0 \times \vec{\Omega}) (\vec{a}_\perp - \gamma \vec{a}_n) \right\} \cdot \vec{v}_0 T^3 / c$ + $[\vec{k}_{\text{eff}} \times \vec{n}_0 \times \vec{a}_\perp - k_{\text{eff}} (\vec{a}_\perp - \gamma \vec{a}_n)] \cdot (\vec{\Omega} \times \vec{v}_0) T^3 / c$ + $2[\vec{k}_{\text{eff}} \cdot \vec{\Omega} \times (\vec{a}_\perp - \gamma \vec{a}_n)] v_n T^3 / c$	$\Omega T \frac{v_0}{c} \sim 4 \times 10^{-12}$	SR earth's rotation
3 $-\frac{7}{12} [\gamma (\vec{k}_{\text{eff}} \cdot \vec{a}_n) (\vec{n}_0 \times \vec{\Omega}) \cdot \vec{a}_0 - 2\vec{k}_{\text{eff}} \cdot \vec{\Omega} \times (\vec{a}_\perp - \gamma \vec{a}_n) a_n] T^4 / c$ + $\frac{7}{12} [\vec{k}_{\text{eff}} \times \vec{n}_0 \times \vec{a}_\perp - k_{\text{eff}} (\vec{a}_\perp - \gamma \vec{a}_n)] \cdot (\vec{\Omega} \times \vec{a}_0) T^4 / c$	$\Omega T \frac{a_0 T}{c} \sim 4 \times 10^{-12}$	GR earth's rotation
4 $[-(\gamma + 1) \vec{k}_{\text{eff}} \cdot \vec{a}_0 (\vec{\Omega} \times \vec{R}_e)^2 + (\vec{k}_{\text{eff}} \cdot \vec{\Omega} \times \vec{R}_e) (\vec{\Omega} \times \vec{R}_e) \cdot \vec{a}_0] T^2 / c^2$	$(\frac{\Omega R_e}{c})^2 \sim 2 \times 10^{-12}$	SR earth's rotation
5 $-2\vec{k}_{\text{eff}} \cdot (\Omega \times \vec{v}_0) \phi_0 T^2 / c^2$	$\frac{\Omega v_0}{a_0} \frac{\phi_0}{c^2} \sim 7 \times 10^{-14}$	GR coriolis acceleration
6 $-4(1 - \beta) \vec{k}_{\text{eff}} \cdot (\vec{\Omega} \times \vec{a}_0) \phi_0 T^3 / c^2$	$\Omega T \frac{\phi_0}{c^2} \sim 7 \times 10^{-14}$	Vanish for GR ($\beta = 1$)
7 $-(2\gamma + 1) \vec{k}_{\text{eff}} \times (\vec{\Omega} \times \vec{R}_e) \times \vec{a}_0 \cdot \vec{v}_0 T^2 / c^2$ + $(\vec{k}_{\text{eff}} \times \vec{a}_0 \times \vec{n}_0) \cdot \vec{v}_0 - 3(\vec{k}_{\text{eff}} \cdot \vec{a}_0) v_n (\vec{\Omega} \times \vec{R}_e \cdot \vec{n}_0) T^2 / c^2$ + $\left\{ [\vec{k}_{\text{eff}} \times (\vec{\Omega} \times \vec{R}_e) \times \vec{a}_0] \cdot \vec{v}_0 + 3(\vec{k}_{\text{eff}} \cdot \vec{a}_0) [\vec{v}_0 \cdot (\vec{\Omega} \times \vec{R}_e)] \right\} T^2 / c$	$\frac{\Omega R_e v_0}{c^2} \sim 7 \times 10^{-14}$	SR earth's rotation
8 $-(\vec{\Omega} \times \vec{R}_e \cdot \vec{n}_0) (\vec{k}_{\text{eff}} \times \vec{n}_0 \times \vec{a}_\perp) \cdot \vec{a}_0 T^3 / c^2$ + $\left\{ [\vec{k}_{\text{eff}} \times (\vec{\Omega} \times \vec{R}_e) \times \vec{a}_0] \cdot \vec{a}_0 + 3(\vec{k}_{\text{eff}} \cdot \vec{a}_0) (\vec{a}_0 \cdot \vec{\Omega} \times \vec{R}_e) \right\} T^3 / c^2$ + $[(\vec{k}_{\text{eff}} \times \vec{a}_0 \times \vec{n}_0) \cdot \vec{a}_0 - 3(\vec{k}_{\text{eff}} \cdot \vec{a}_0) a_n] (\vec{\Omega} \times \vec{R}_e \cdot \vec{n}_0) T^3 / c^2$ - $\left\{ 2(\gamma + 1) [\vec{k}_{\text{eff}} \times (\vec{\Omega} \times \vec{R}_e) \times \vec{a}_0] \cdot \vec{a}_0 - 3(\vec{k}_{\text{eff}} \cdot \vec{a}_0) [\vec{a}_0 \cdot (\vec{\Omega} \times \vec{R}_e)] \right\} T^3 / c^2$	$\frac{\Omega R_e a_0 T}{c^2} \sim 7 \times 10^{-14}$	GR earth's rotation

New results !

4. Conclusion

- Developed an analytical study method, in which a coordinate transformation is performed to simplify the calculation!
- Based on the method, we studied FSL effect in atom gravimeters, and further proposed a preliminary experimental design to test this effect.

$$g = g_0 \left[1 + 3 \frac{\vec{v}(T) \cdot \vec{e}_k}{c} - 2 \frac{\vec{v}(T) \cdot \vec{e}_k}{c} \frac{\alpha_1 - \alpha_2}{\vec{k}_{\text{eff}} \cdot \vec{g}_0} + \frac{2\vec{v}(T)}{c} \cdot \frac{\alpha_1 \vec{n}_1 - \alpha_2 \vec{n}_2}{\vec{k}_{\text{eff}} \cdot \vec{g}_0} \right] + \dots$$

- Extending the analysis on FSL effect, we established a more complete relativistic model for atom gravimeters!



Thanks for your attention!