

How Higgs physics can uncover the
nature of flavour?

Family dependent Higgs Physics

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- Kinds of models
- Prospects of detection at colliders
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SM: 3 families of fermions, 1 of Higgs?

Three Generations of Matter (Fermions)				
	I	II	III	
mass→	3 MeV	1.24 GeV	172.5 GeV	0
charge→	$\frac{2}{3}$	$\frac{2}{3}$	$\frac{2}{3}$	0
spin→	$\frac{1}{2}$	$\frac{1}{2}$	$\frac{1}{2}$	1
name→	u up	c charm	t top	γ photon
Quarks	6 MeV	95 MeV	4.2 GeV	0
	$-\frac{1}{3}$	$-\frac{1}{3}$	$-\frac{1}{3}$	0
	$\frac{1}{2}$	$\frac{1}{2}$	$\frac{1}{2}$	1
	d down	s strange	b bottom	g gluon
Leptons	<2 eV	<0.19 MeV	<18.2 MeV	90.2 GeV
	0	0	0	0
	$\frac{1}{2}$	$\frac{1}{2}$	$\frac{1}{2}$	1
	ν_e electron neutrino	ν_μ muon neutrino	ν_τ tau neutrino	Z^0 weak force
	0.511 MeV	106 MeV	1.78 GeV	80.4 GeV
	-1	-1	-1	± 1
	$\frac{1}{2}$	$\frac{1}{2}$	$\frac{1}{2}$	1
	e electron	μ muon	τ tau	W^\pm weak force

H
1 Higgs Boson

- Should not be considering this as well?

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H_1 H_2 H_3

Traditionally

- The Higgs is special:

Just one or family independent extensions

GUT+ family symmetries can solve the problem of hierarchies and mixing for fermions

- Lots of complications if we add more Higgs bosons:

Fine tuning to obtain EW breaking, non-perturbatively, etc.

Kinds of models

- The Higgs is special:

Flavour comes from other sector, e.g. **Flavons**

$$-\mathcal{L}_{\text{Flavon}}(\varphi_n^\psi) = \sum_n \bar{\psi}_{Li} \psi_{Rj} H c_{\varphi_n^\psi ij}^\psi \left(\frac{\varphi_n^\psi}{\Lambda_{\varphi_n^\psi}} \right)^{p_{nij}^\psi} + \text{H.c.},$$

$H \rightarrow$ Higgs boson

$\psi \rightarrow$ SM fermions

$\varphi_n \rightarrow$ flavons

- There are really **family dependent Higgs bosons** : 3 for each kind of fermion

$$-\mathcal{L}_H = \sum_{k=1,3,5} \bar{\psi}_{Li} D_{Rj} H_k + \sum_{k=2,4,6} \bar{\psi}_{Li} U_{Rj} H_k + \text{H.c.},$$

Flavons at Colliders

- The flavon acquires vacuum expectation value

$$\varphi = \frac{(v_\varphi + H_f + iA_f)}{\sqrt{2}}$$

$$-\mathcal{L}_{\text{Flavon}}(\varphi_n^\psi) = c_{ij}^\ell \bar{L}_i e_{Rj} H \left(\frac{\varphi}{\Lambda}\right)^{-\mathcal{Q}_F(q_{L_i}^\ell + q_{R_j}^e)} + c_{ij}^d \bar{Q}_i d_{Rj} H \left(\frac{\varphi}{\Lambda}\right)^{-\mathcal{Q}_F(q_{L_i}^d + q_{R_j}^d)} \\ + c_{ij}^u \bar{Q}_i u_{Rj} \tilde{H} \left(\frac{\varphi}{\Lambda}\right)^{-\mathcal{Q}_F(q_{L_i}^d + q_{R_j}^u)} + \text{H.c.},$$

0911.2149 K. Tsumura, L. V-S, 1603.06950 Bauer et al., L. V-S
1709.xxxx

- Full cancellation of anomalies possible with two flavons.

$U(1)_F$ charges						
Field	\bar{Q}_{Li}	d_{Ri}	u_{Ri}	\bar{L}_{Li}	e_{Ri}	ν_{Ri}
Charge	$q_{L_i}^{\bar{Q}}$	$q_{R_i}^d$	$q_{R_i}^u$	$q_{L_i}^{\bar{L}}$	$q_{R_i}^e$	$q_{R_i}^\nu$

$$Y_{ij}^f = c_{ij}^f \left(\frac{v_\varphi}{\sqrt{2}\Lambda}\right)^{(q_{L_i}^f + q_{R_j}^f)},$$

- Effective Lagrangian for phenomenology

$$-\mathcal{L}_\varphi^{\text{eff}} = m_i^f \overline{f_{L_i}} f_{R_i} \left(1 + \frac{h}{v} \right) + \kappa_{ij}^f \overline{f_{L_i}} f_{R_j} \left(\frac{\varphi}{v_\varphi/\sqrt{2}} \right) + \text{H.c.},$$

- Collider phenomenology controlled by

$$\kappa_{ij}^f = \left[m_j^f \sum_k q_{Lk}^f \left(U_L^f \right)_{ik} \left(U_L^f \right)_{jk}^* + m_i^f \sum_k q_{Rk}^f \left(U_R^f \right)_{ik} \left(U_R^f \right)_{jk}^* \right].$$

- The interaction with extra Z is controlled by

$$\mathcal{L}(\varphi_i, Z_k) = (D^\mu H)^\dagger (D_\mu H) + \sum_i (D^\mu \varphi_i)^\dagger (D_\mu \varphi_i) - V(\varphi_i), \quad k = 1, 2, \quad i = 1, 2$$

$$D_\mu \varphi_i = \frac{1}{\sqrt{2}} \left(\partial_\mu - \frac{i}{2} g_{F_i} \mathcal{Q}_{F_i} (Z_\mu)_k \right) \varphi_i$$

- The most stringent bounds may come from the FCNC decays

$$\Gamma_{\ell_i \rightarrow \ell_j \gamma} = \frac{m_i^3}{4\pi} \left(1 - \frac{m_j^2}{m_i^2}\right)^3 \left(\left|A_{Lij}^\gamma\right|^2 + \left|A_{Rij}^\gamma\right|^2 \right).$$

$$A_{Lij}^\gamma = \frac{1}{(4\pi)^2} \frac{Q_f e}{2v_\varphi^2} \left[-\kappa_{jk}^f \kappa_{ki}^f m_k c_{11} + \kappa_{jk}^f \kappa_{ik}^{*f} [m_j (c_{11} - c_{12} + c_{21} - c_{23}) + m_i (c_{12} + c_{23})] \right],$$

$$A_{Rij}^\gamma = \frac{1}{(4\pi)^2} \frac{Q_f e}{2v_\varphi^2} \left[-\kappa_{kj}^{*f} \kappa_{ik}^{*f} m_k c_{11} + \kappa_{kj}^{*f} \kappa_{ki}^f [m_j (c_{11} - c_{12} + c_{21} - c_{23}) + m_i (c_{12} + c_{23})] \right],$$

- The leading decay channels are given by

$$\Gamma_{A_f \rightarrow f_i \bar{f}_j} = N_C^f \frac{G_F m_\varphi}{8\sqrt{2}\pi} \left(\left| \kappa_{ij}^f \right|^2 + \left| \kappa_{ji}^f \right|^2 \right) \left(\frac{v}{v_\varphi} \right)^2 \beta_{ij} \left[1 - \frac{(m_i - m_j)^2}{m_{A_f}^2} \right],$$

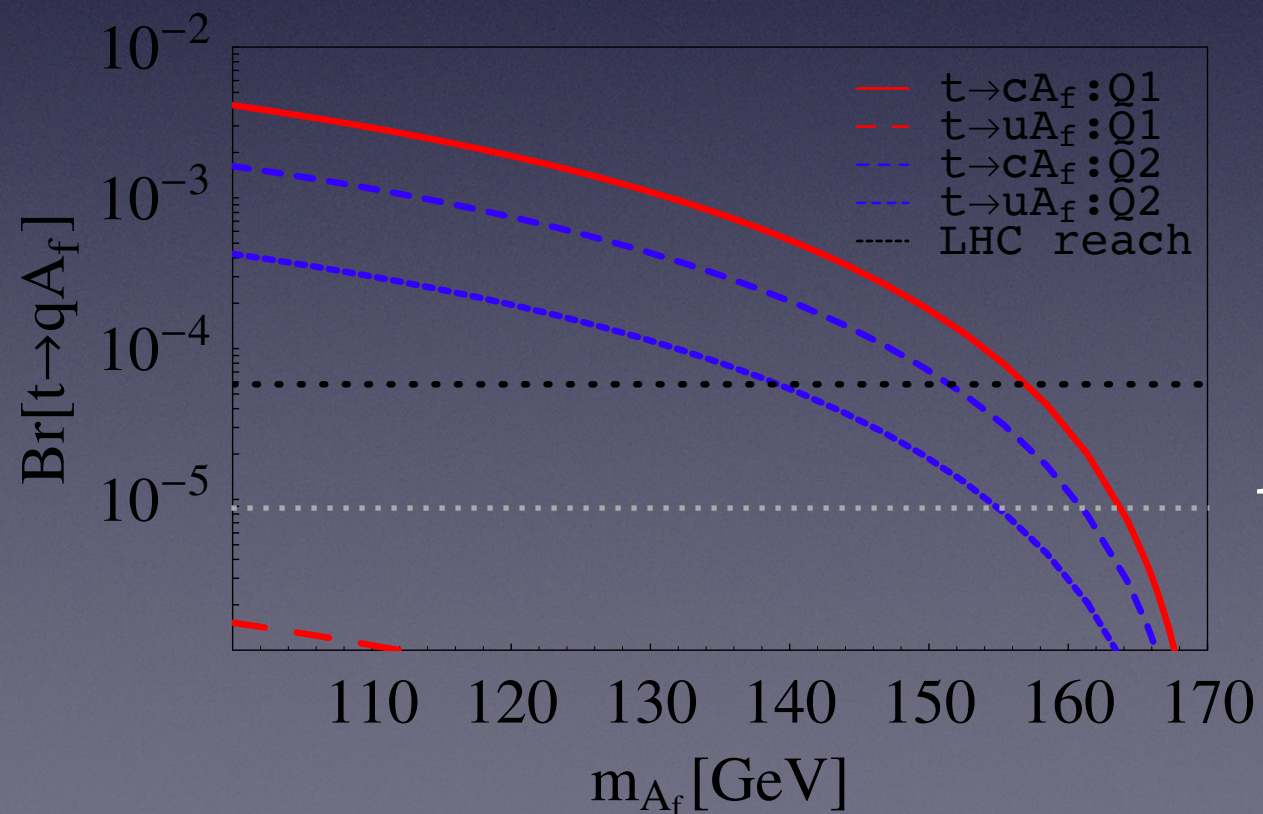
$$\Gamma_{A_f \rightarrow gg} = \frac{G_F \alpha_S^2 m_\varphi^3}{64\sqrt{2}\pi} \left(\frac{v}{v_\varphi} \right)^2 \left| N_C \sum_q \kappa_{qq}^q \frac{I(m_q)}{m_q} \right|^2,$$

Most important constraints

- LEP bounds

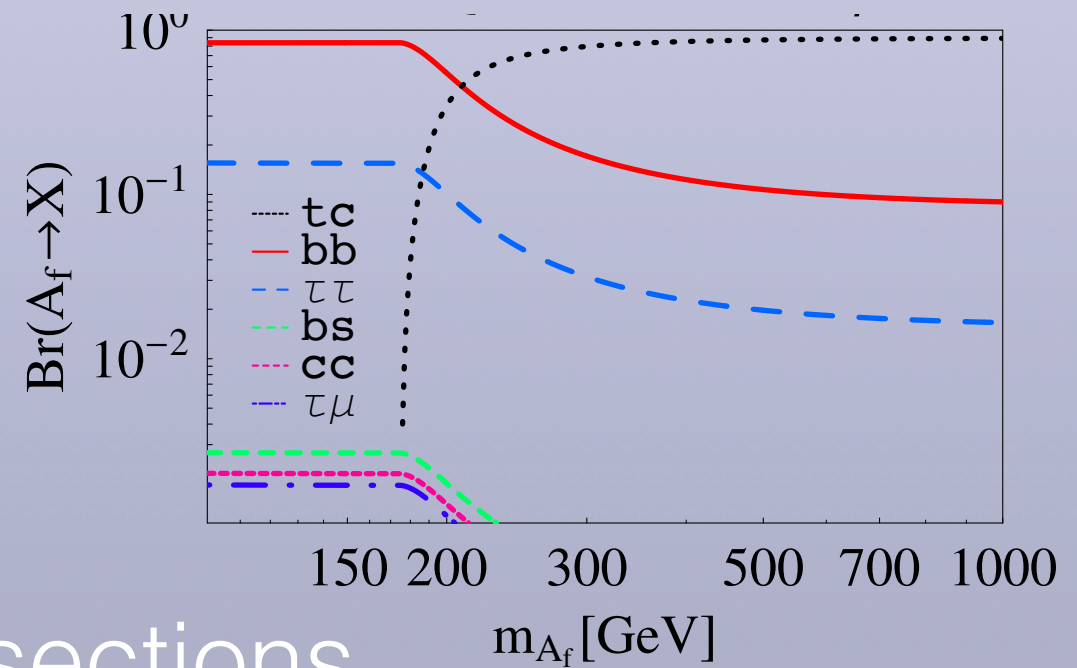
$$e^-e^+ \rightarrow e^-e^+ \quad \frac{M_{Z_1}}{g_{F_1}} \geq 26 \text{ TeV}, \quad \frac{M_{Z_2}}{g_{F_2}} \geq 116 \text{ TeV}$$

- Production through top decays

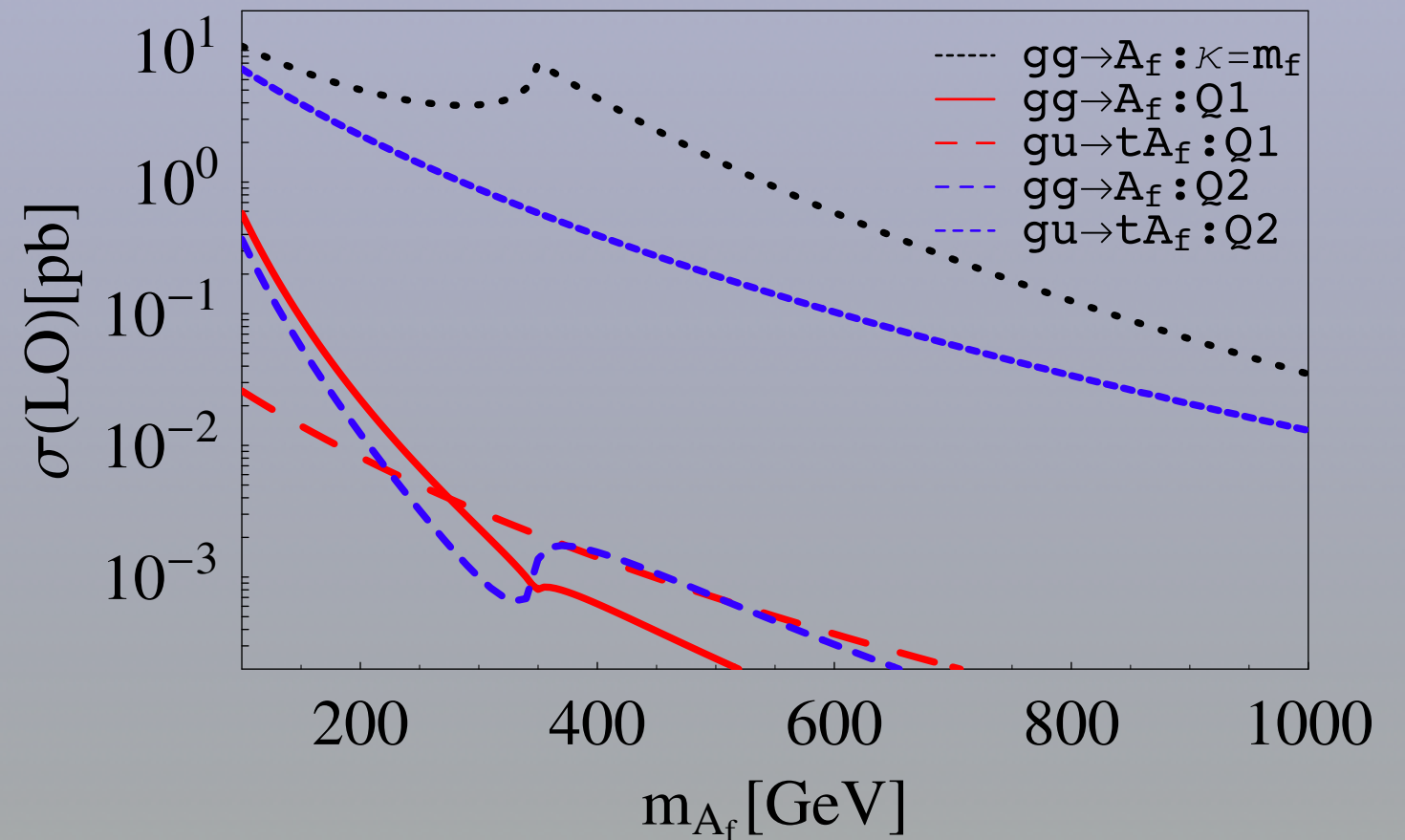


Signatures

- Flavon decays



- Typical production cross sections



Family Dependent Higgs bosons at Colliders

$$-\mathcal{L}_H = \sum_{k=1,3,5} \bar{\psi}_{Li} D_{Rj} H_k + \sum_{k=2,4,6} \bar{\psi}_{Li} U_{Rj} H_k + \text{H.c.},$$

- Non-supersymmetric: Need Z_2 symmetries to forbid arbitrary couplings to D and U quarks
- Supersymmetric: Holomorphicity in the superpotential determines the way they couple: only 3 Higgs doublets couple to D, and 3 different ones to U

$$W(H_k) = \sum_{k=1,3,5} \bar{\psi}_{Li} D_{Rj} H_k + \sum_{k=2,4,6} \bar{\psi}_{Li} U_{Rj} H_k \text{H.c.},$$

- The soft potential is given by

$$V_{\text{soft}} = \sum_{i,j=1,3,5} (m_d^2)_{ij} H_i^\dagger H_j + \sum_{k,l=2,4,6} (m_u^2)_{kl} H_k^\dagger H_l - \sum_{\substack{i=1,3,5 \\ l=2,4,6}} [(B\mu)_{il} H_i H_l + \text{H.c.}] .$$

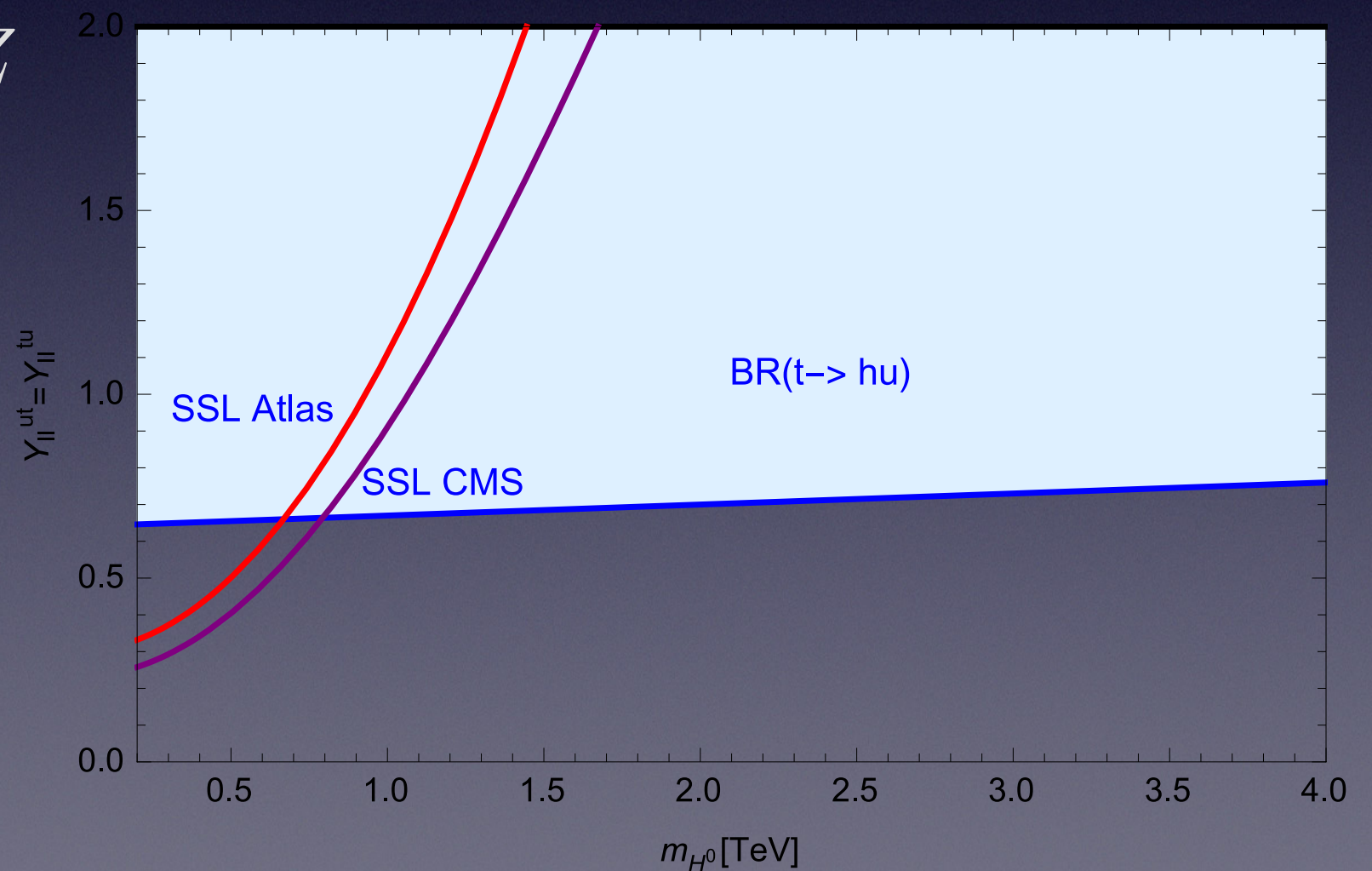
- Symmetries force the following alignment

$$Y_I = \begin{pmatrix} 0 & 0 & 0 \\ 0 & O(1) & O(1) \\ 0 & O(1) & O(1) \end{pmatrix}, \quad Y_{II} = \begin{pmatrix} O(1) & O(1) & O(1) \\ O(1) & O(1) & O(1) \\ O(1) & O(1) & O(1) \end{pmatrix}, \quad Y_{III} = \begin{pmatrix} O(1) & O(1) & O(1) \\ O(1) & 0 & 0 \\ O(1) & 0 & 0 \end{pmatrix}$$

Most important constraints

- Top decay BR ($t \rightarrow hu$)
- Same-sign di-lepton (SSL) + b jet searches

$$H^0 \rightarrow WW, ZZ$$



Summary

- Family dependent Higgs Bosons and/or flavons can give an explanation to CKM matrix and values of fermions masses
- Interesting phenomenology
- Good prospects of test/exclusion at Colliders

Field	Decay	Production	FCNC: Most constraining?	Precision to be discriminated
φ_k Flavon	✓	✓	✓ (From top decays)	O(1%)
H_K Heavy extra Higgs	✓	×	✓ (From top decays)	O(1%)

TABLE I: Processes which can be tested at Colliders.