

Neutrino physics and implications

José W F Valle

Lecture II



<https://www.facebook.com/ific.ahep/>

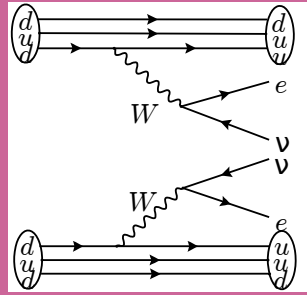
Summer School and Workshop on Standard Model and Beyond – Corfu 2016

Outline

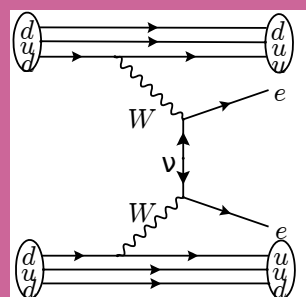
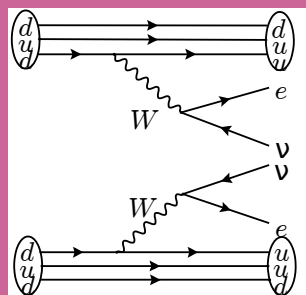
Lecture II

- Neutrinoless double beta decay
- Flavor predictions for neutrinoless double beta decay
- Long versus short range mechanisms
- Black-box theorem
- Neutrinos and electroweak vacuum
- Stability & perturbativity
- Invisible Higgs decays
- Gravity and the Standard Model
- Extra dimensions & unification
- Neutrino predictions from warped flavor
- String completion of $SU(3)_c \otimes SU(3)_L \otimes U(1)$ model
- Neutrinos in cosmology

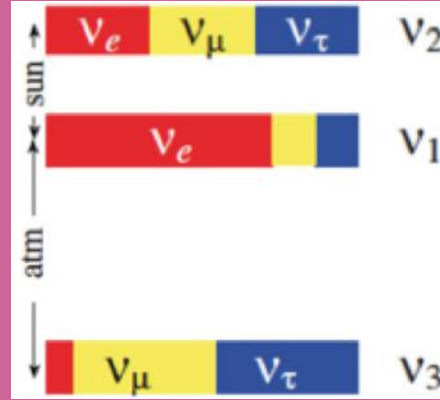
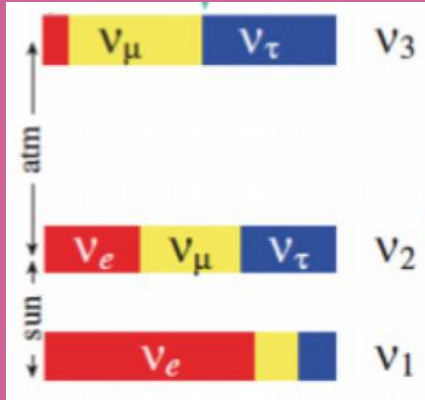
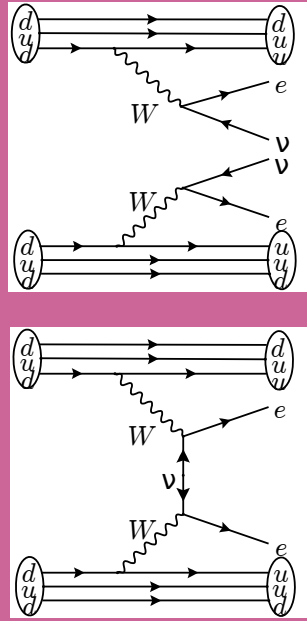
Neutrinoless double beta decay



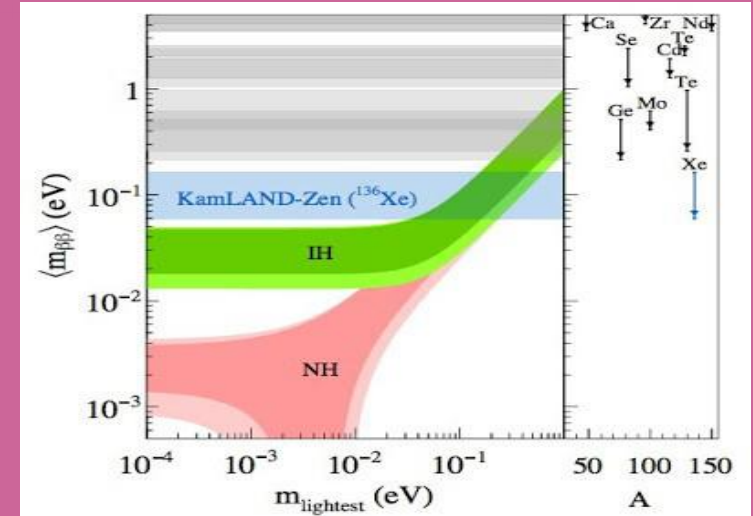
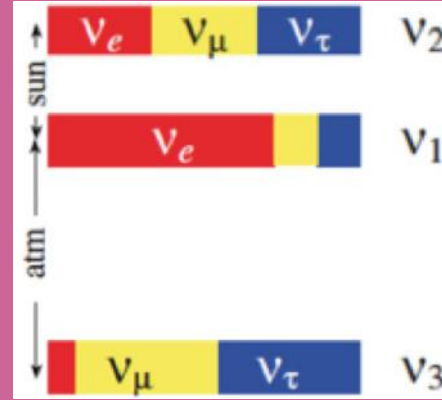
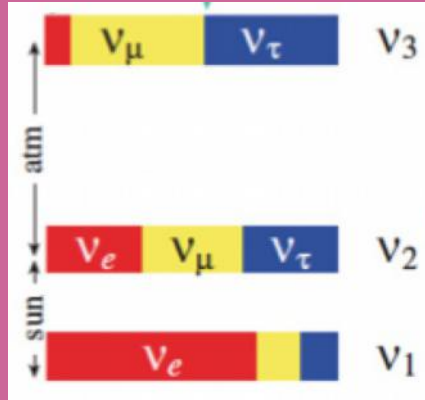
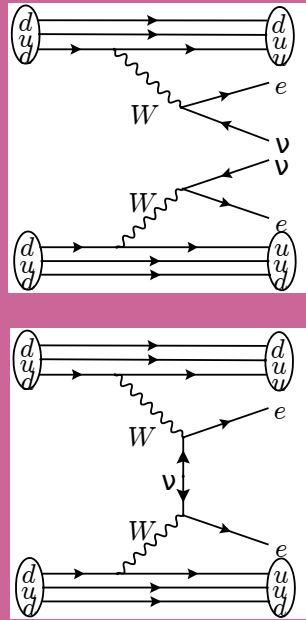
Neutrinoless double beta decay



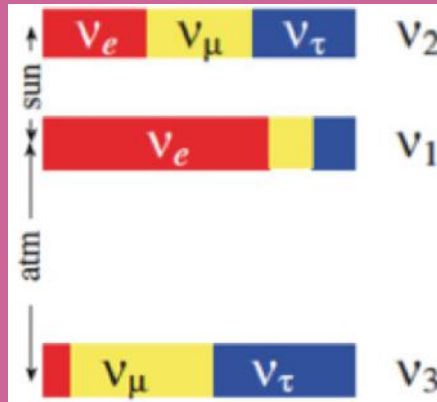
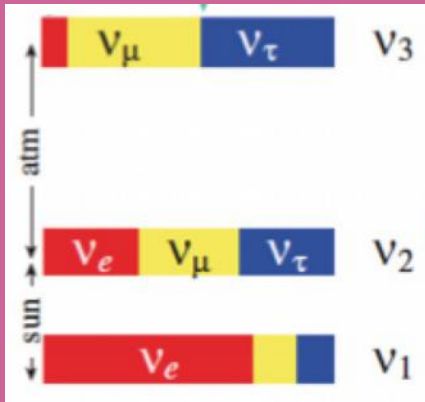
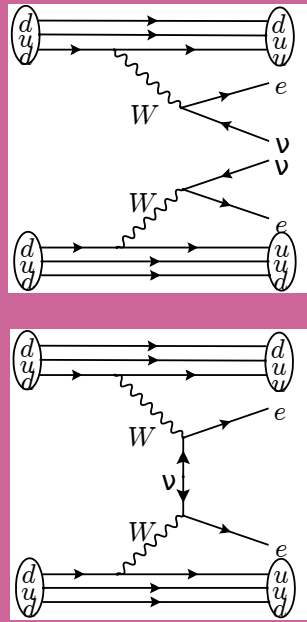
Neutrinoless double beta decay



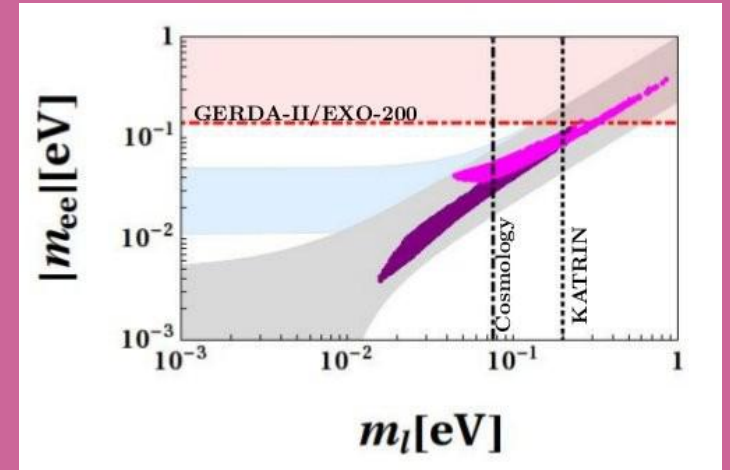
Neutrinoless double beta decay



Neutrinoless double beta decay

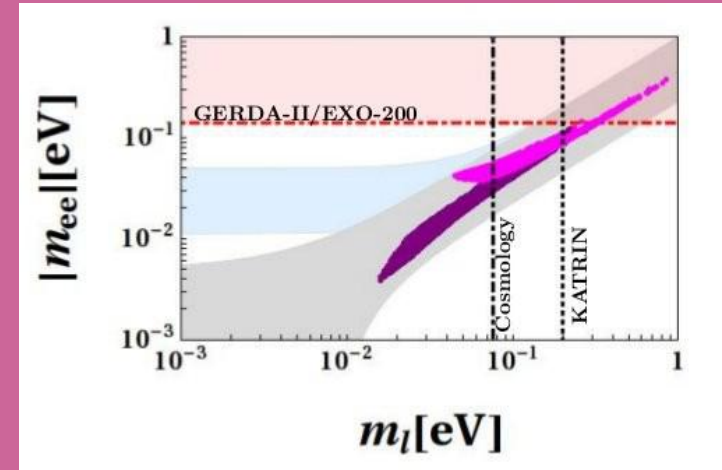
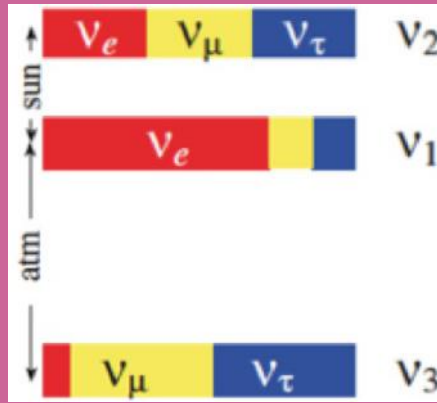
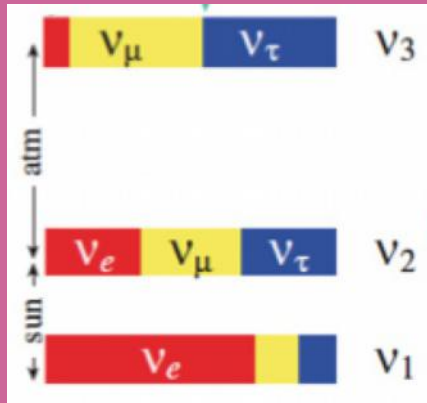
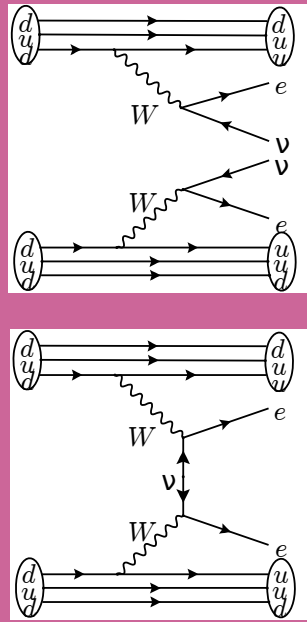


Flavor Sensitivity



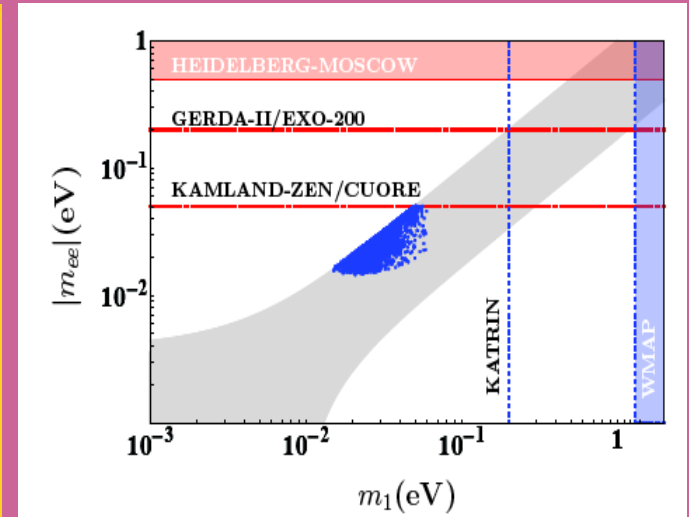
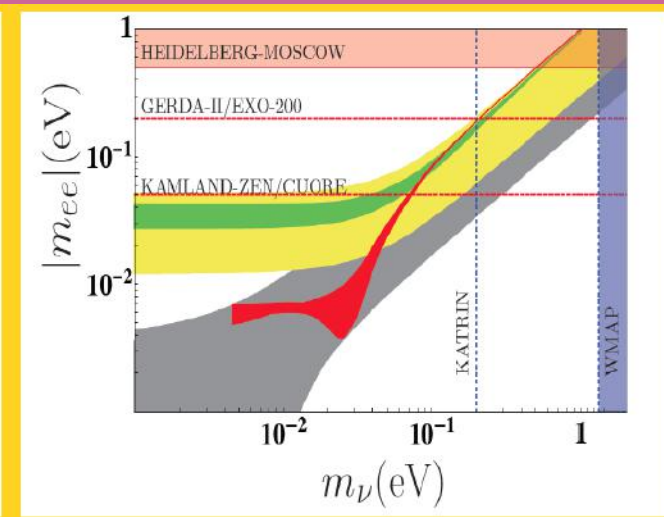
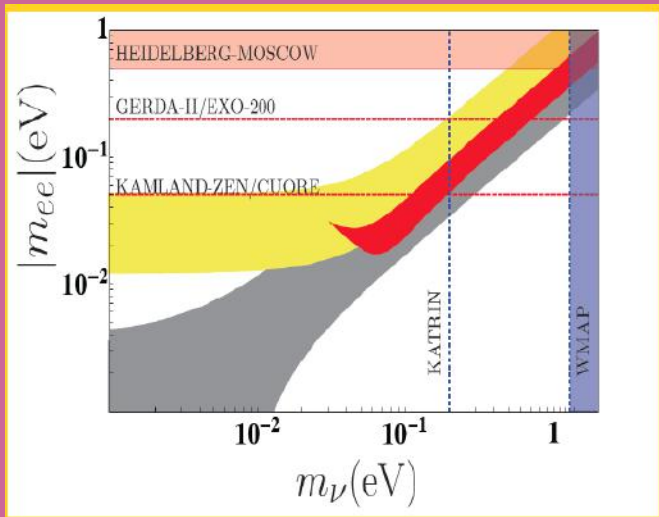
Bonilla et al arXiv:1411.4883

Neutrinoless double beta decay



Flavor Sensitivity

Bonilla et al arXiv:1411.4883



Dorame et al

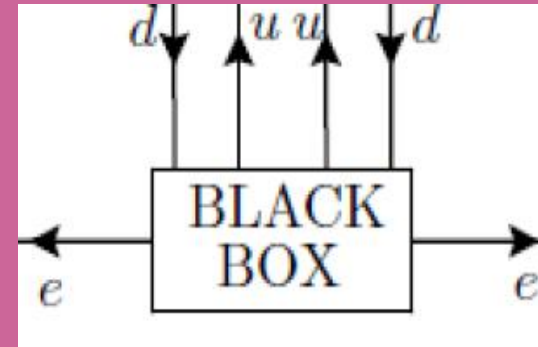
NPB861 (2012) 259-270

PhysRevD.86.056001

King et al Phys. Lett. B 724 (2013) 68



The Majorana connection



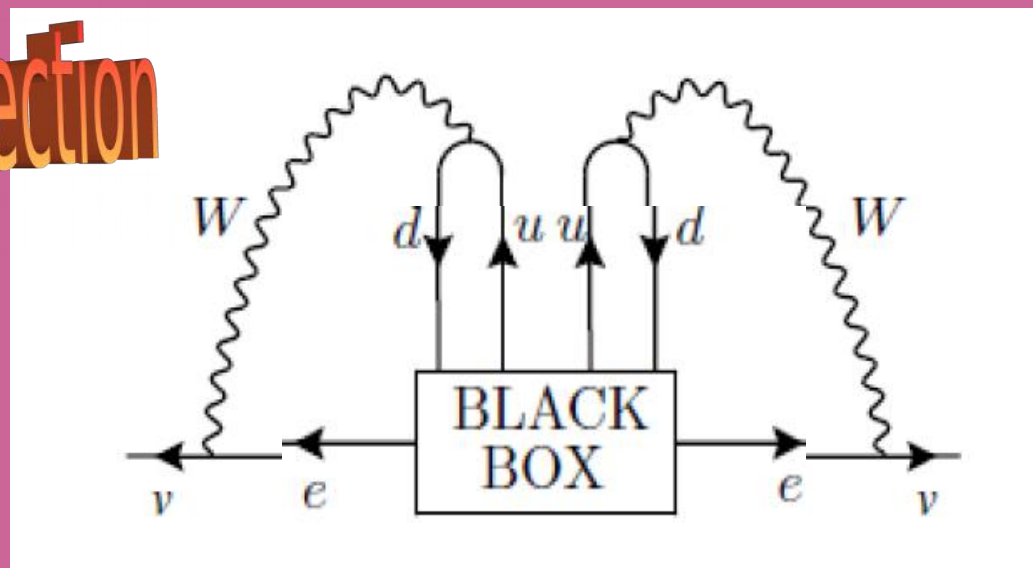
Schechter, JWFV 82

Lindner et al JHEP 1106 (2011) 091



The Majorana connection

*Even if mediated by
short-range mechanism ...*

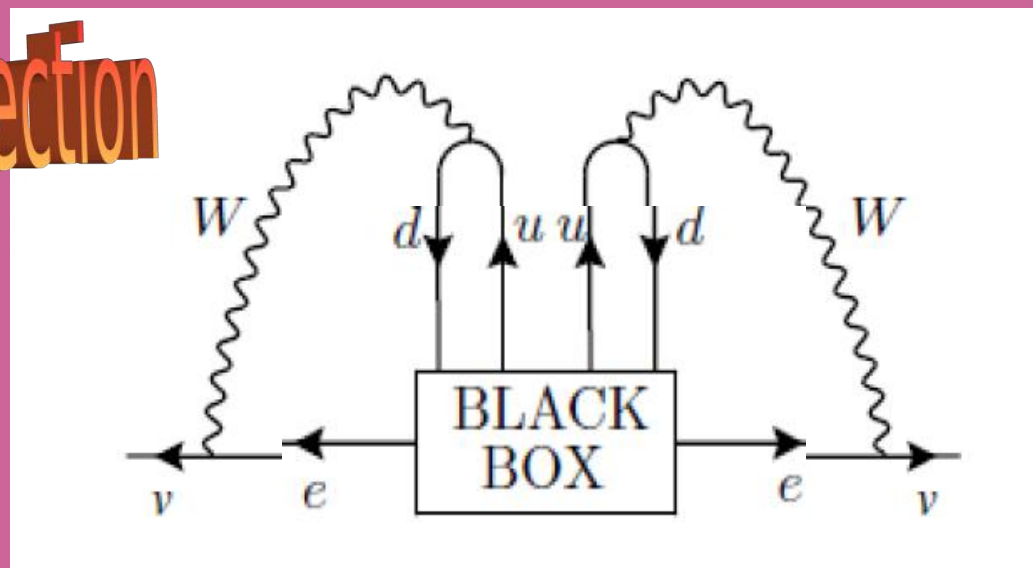


Schechter, JWFV 82

Lindner et al JHEP 1106 (2011) 091



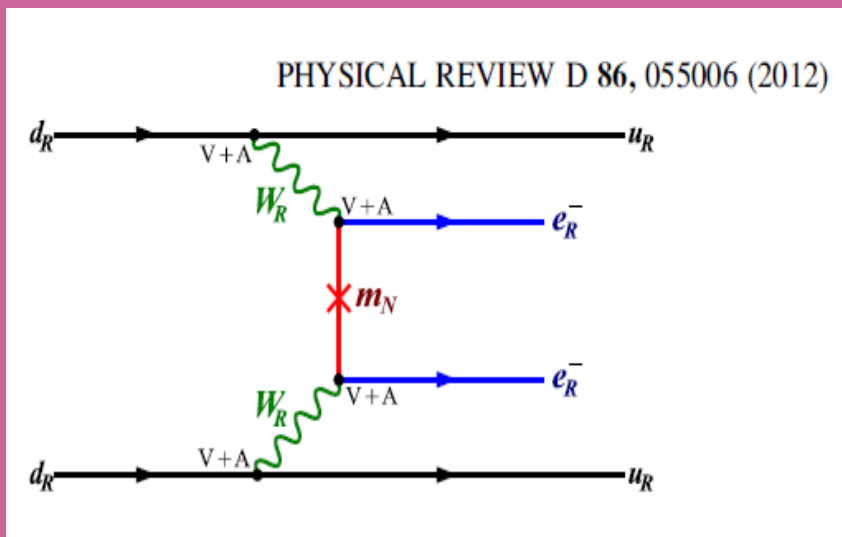
The Majorana connection



Even if mediated by short-range mechanism ...

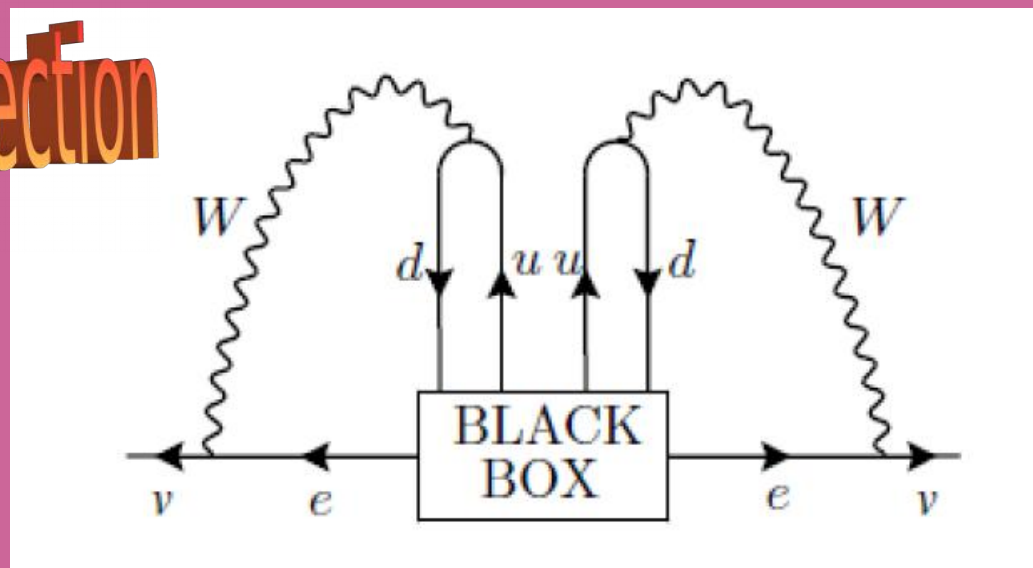
Schechter, JWFV 82
Lindner et al JHEP 1106 (2011) 091

Heavy mediators





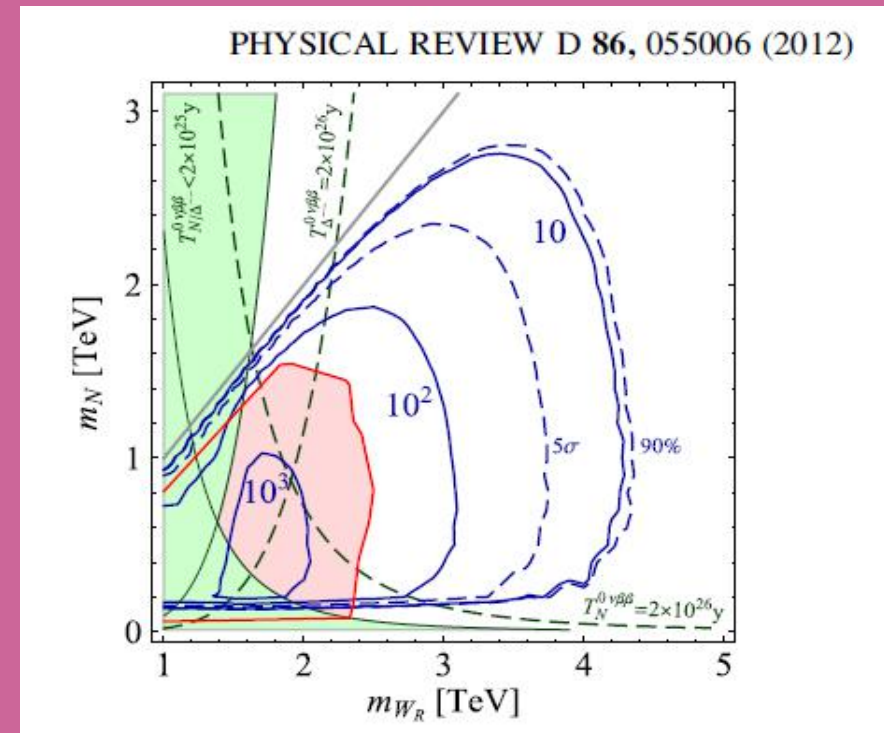
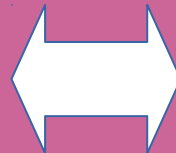
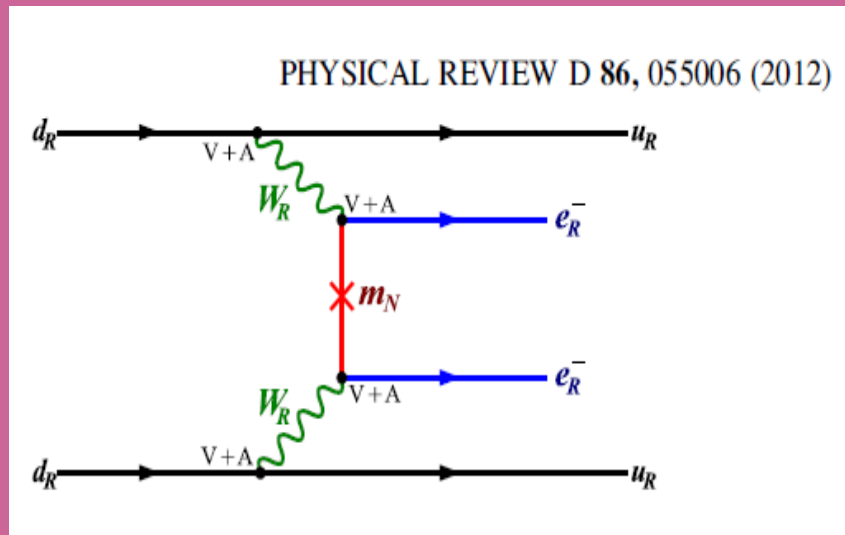
The Majorana connection



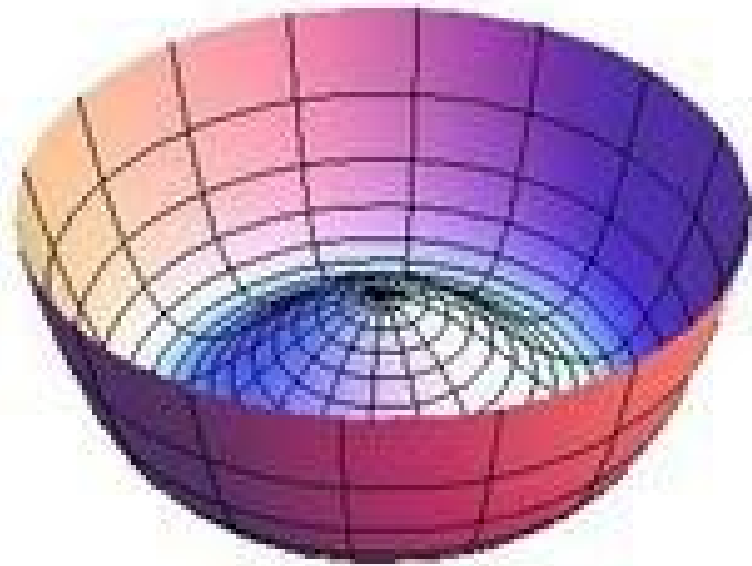
Even if mediated by short-range mechanism ...

Schechter, JWFV 82
Lindner et al JHEP 1106 (2011) 091

Heavy mediators



Spontaneous electroweak symmetry breaking
The role of neutrino mass generation



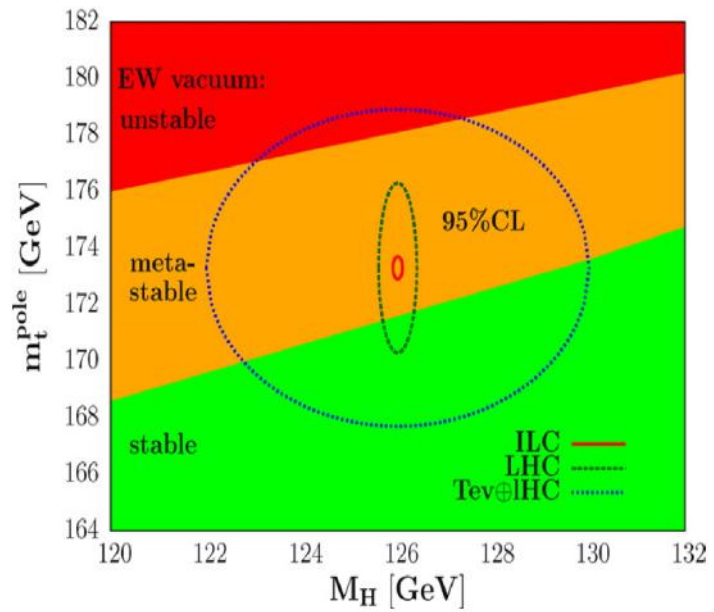


SM vacuum

SM vacuum

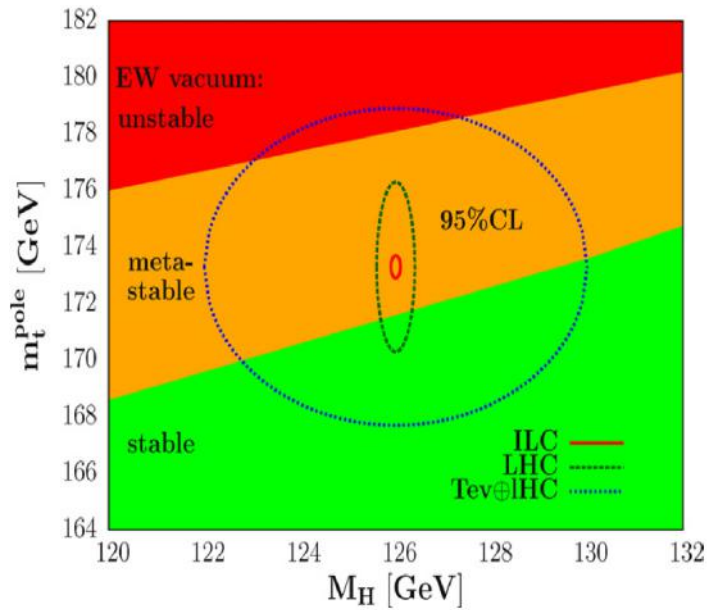


Physics Letters B 716 (2012) 214–219



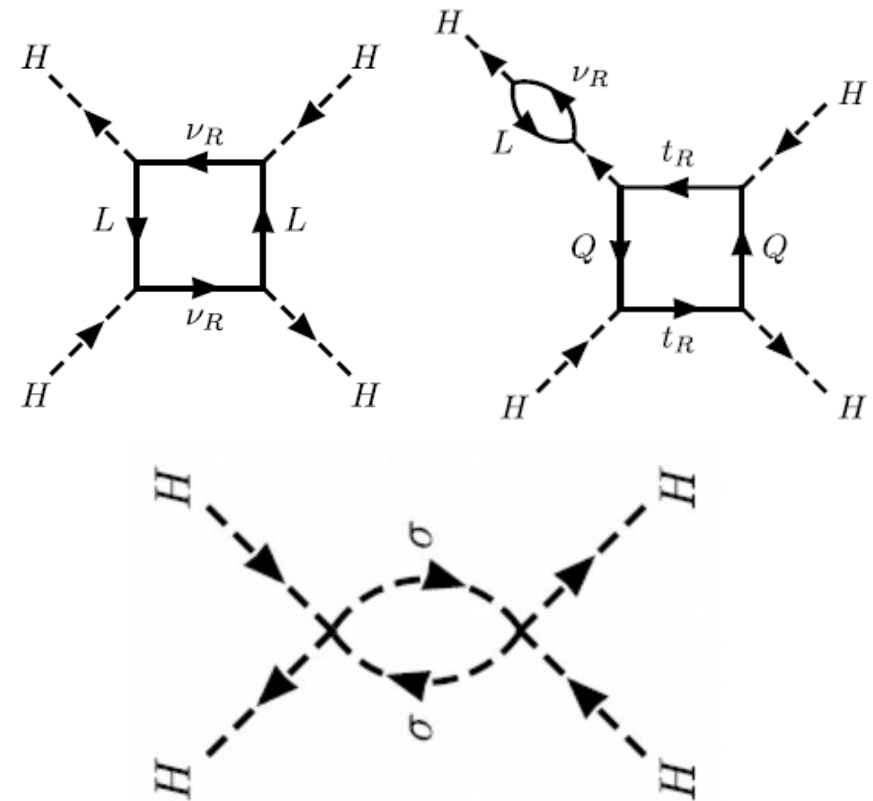


Physics Letters B 716 (2012) 214–219



SM vacuum and neutrinos

Physics Letters B 756 (2016) 345–349



Stability from neutrinos

Physics Letters B 756 (2016) 345–349

In addition to SM gauge invariance must break lepton number to give masses to neutrinos ..

$$(4\pi)^2 \frac{dg_i}{dt} = b_i g_i^3 \text{ with } b_i = \left(\frac{41}{10}, -\frac{19}{6}, -7 \right),$$

$$(4\pi)^2 \frac{dY_t}{dt} = \left(-\frac{17}{20}g_1^2 - \frac{9}{4}g_2^2 - 8g_3^2 + \frac{9}{2}Y_t^2 + Y_\nu^2 \right) Y_t,$$

$$(4\pi)^2 \frac{dY_\nu}{dt} = \left(-\frac{9}{20}g_1^2 - \frac{9}{4}g_2^2 + 3Y_t^2 + \frac{5}{2}Y_\nu^2 \right) Y_\nu,$$

$$(4\pi)^2 \frac{dY_S}{dt} = 6Y_S^3,$$

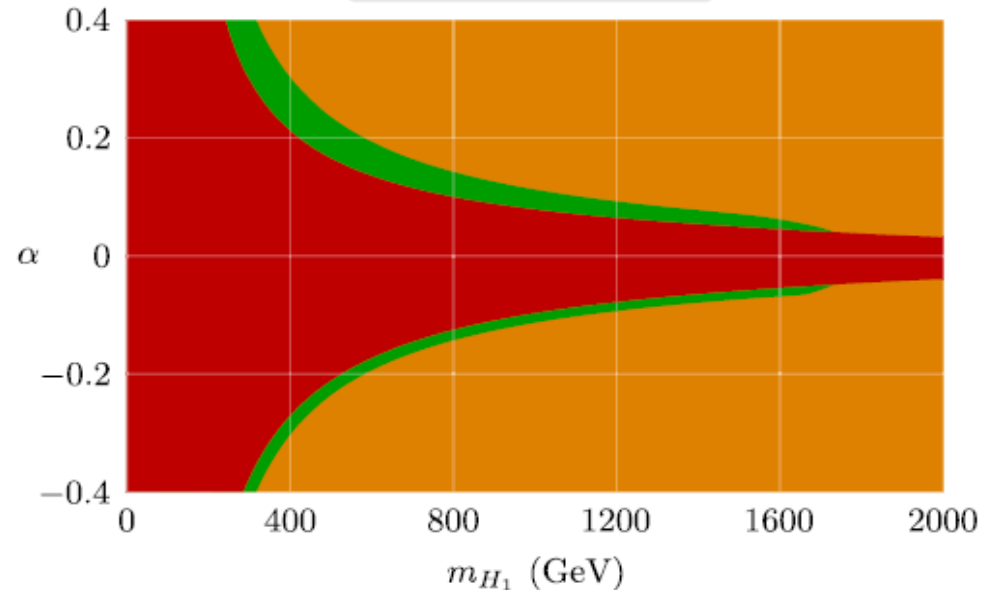
$$(4\pi)^2 \frac{d\lambda_1}{dt} = 20\lambda_1^2 + 2\lambda_{12}^2 + 8\lambda_1 Y_S^2 - 16Y_S^4,$$

$$(4\pi)^2 \frac{d\lambda_2}{dt} = \frac{27}{200}g_1^4 + \frac{9}{20}g_1^2 g_2^2 + \frac{9}{8}g_2^4 - \left(\frac{9}{5}g_1^2 + 9g_2^2 \right) \lambda_2 + 24\lambda_2^2 + \lambda_{12}^2 + \lambda_2 \left(12Y_t^2 + 4Y_\nu^2 \right) - \left(6Y_t^4 + 2Y_\nu^4 \right),$$

$$(4\pi)^2 \frac{d\lambda_{12}}{dt} = \left[- \left(\frac{9}{10}g_1^2 + \frac{9}{2}g_2^2 \right) + 6Y_t^2 + 2Y_\nu^2 + 4Y_S^2 + 8\lambda_1 + 12\lambda_2 + 4\lambda_{12} \right] \lambda_{12}.$$

$$V(\sigma, H) = \mu_1^2 |\sigma|^2 + \mu_2^2 H^\dagger H + \lambda_1 |\sigma|^4 + \lambda_2 (H^\dagger H)^2 + \lambda_{12} (H^\dagger H) |\sigma|^2.$$

$Y_\nu = 0.5, v_\sigma = 3 \text{ TeV}$



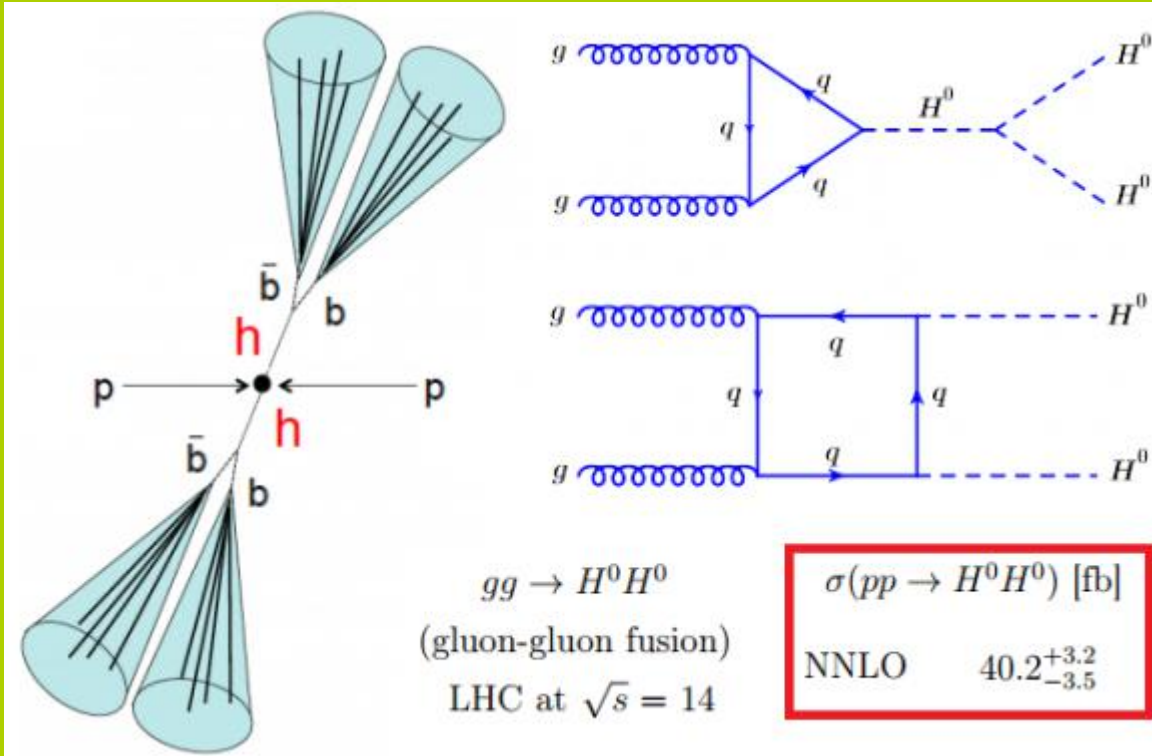
Neutrino as higgs benchmark

Where is the New physics?

Large Hadron Collider



Higgs production & decay



channel	ATLAS	CMS
$\mu_{\gamma\gamma}$	1.17 ± 0.27	$1.14^{+0.26}_{-0.23}$
μ_{WW}	$1.00^{+0.32}_{-0.29}$	0.83 ± 0.21
μ_{ZZ}	$1.44^{+0.40}_{-0.35}$	1.00 ± 0.29
$\mu_{\tau^+\tau^-}$	$1.4^{+0.5}_{-0.4}$	0.91 ± 0.27
$\mu_{b\bar{b}}$	$0.2^{+0.7}_{-0.6}$	0.93 ± 0.49

Neutrino mass and invisible Higgs decays at the LHCCesar Bonilla,^{1,*} Jorge C. Romão,^{2,†} and José W. F. Valle^{1,‡} $H_i \rightarrow JJ$ and $H_2 \rightarrow 2H_1 \rightarrow 4J$ $\left(\text{when } m_{H_1} < \frac{m_{H_2}}{2}\right).$

arXiv:1502.01649

channel	ATLAS	CMS
$\mu_{\gamma\gamma}$	1.17 ± 0.27	$1.14^{+0.26}_{-0.23}$
μ_{WW}	$1.00^{+0.32}_{-0.29}$	0.83 ± 0.21
μ_{ZZ}	$1.44^{+0.40}_{-0.35}$	1.00 ± 0.29
$\mu_{\tau^+\tau^-}$	$1.4^{+0.5}_{-0.4}$	0.91 ± 0.27
$\mu_{b\bar{b}}$	$0.2^{+0.7}_{-0.6}$	0.93 ± 0.49

Neutrino mass and invisible Higgs decays at the LHC

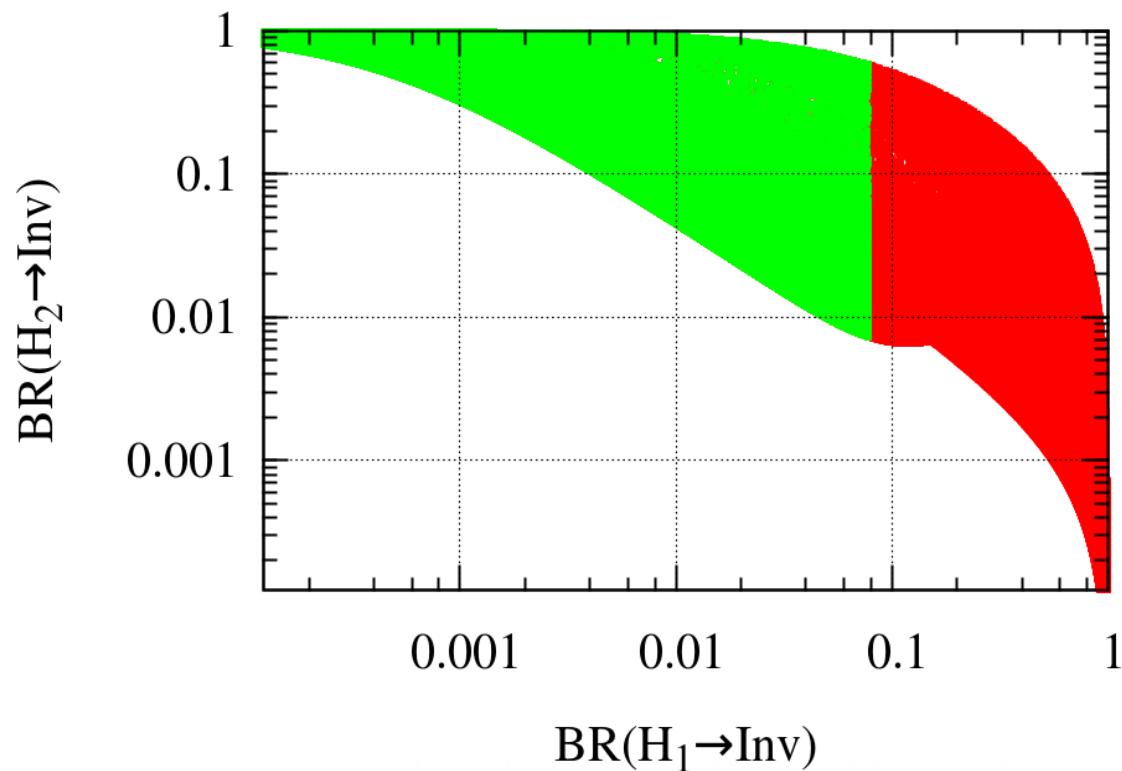
Cesar Bonilla,^{1,*} Jorge C. Romão,^{2,†} and José W.F. Valle^{1,‡}

$v_\sigma=3$ TeV

$$H_i \rightarrow JJ \quad \text{and} \quad H_2 \rightarrow 2H_1 \rightarrow 4J$$

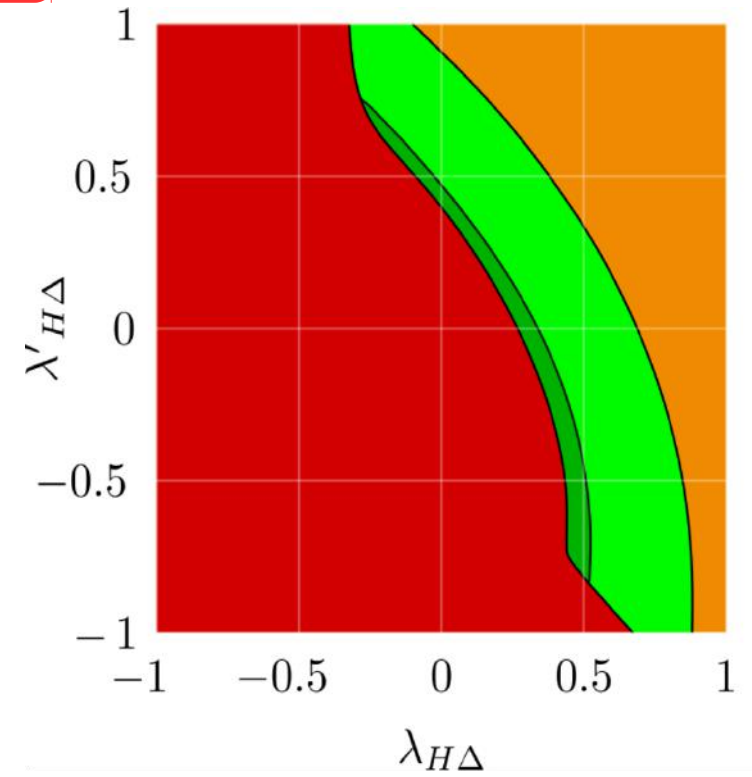
$$\left(\text{when } m_{H_1} < \frac{m_{H_2}}{2} \right).$$

channel	ATLAS	CMS
$\mu_{\gamma\gamma}$	1.17 ± 0.27	$1.14^{+0.26}_{-0.23}$
μ_{WW}	$1.00^{+0.32}_{-0.29}$	0.83 ± 0.21
μ_{ZZ}	$1.44^{+0.40}_{-0.35}$	1.00 ± 0.29
$\mu_{\tau^+\tau^-}$	$1.4^{+0.5}_{-0.4}$	0.91 ± 0.27
$\mu_{b\bar{b}}$	$0.2^{+0.7}_{-0.6}$	0.93 ± 0.49

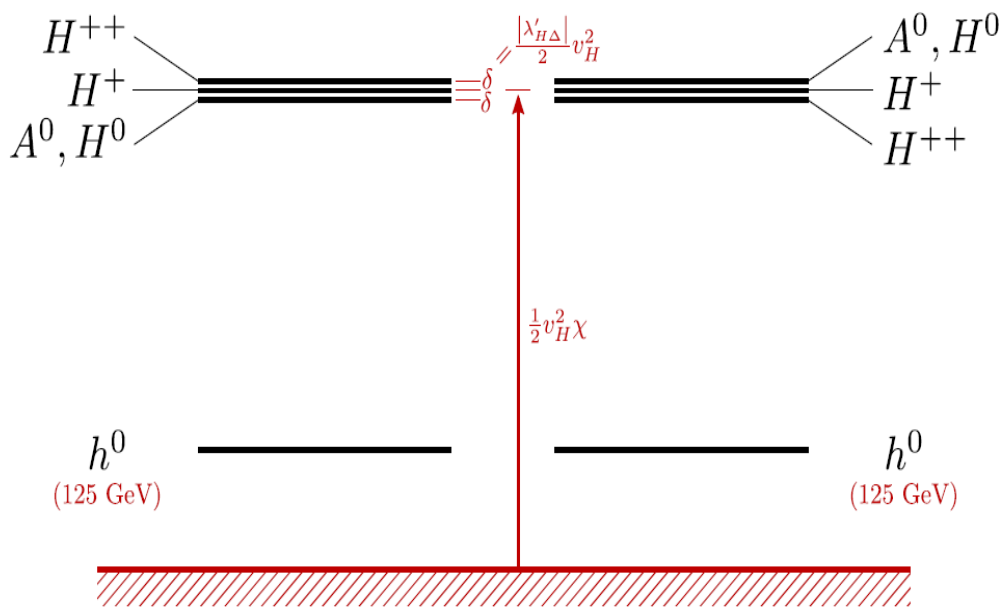
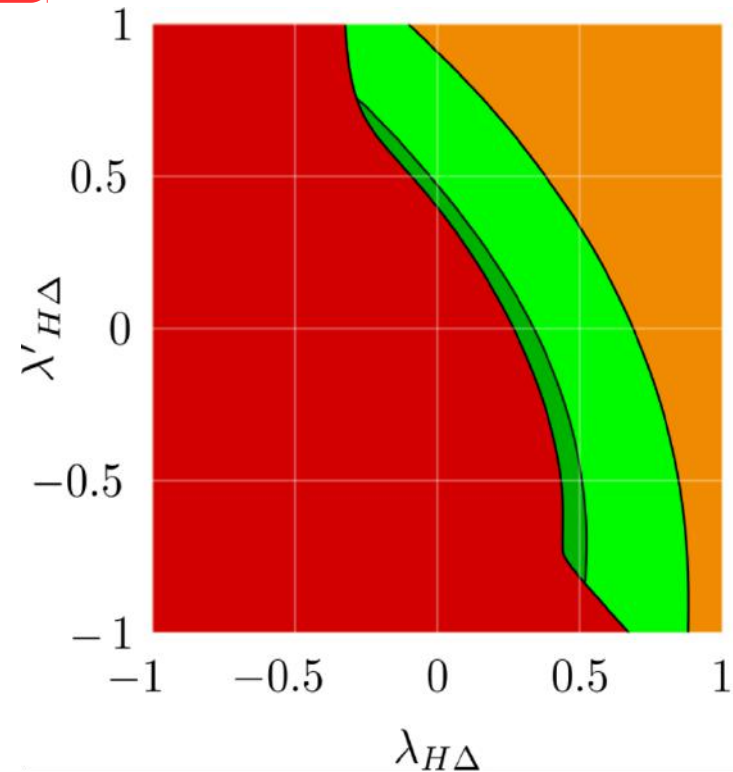
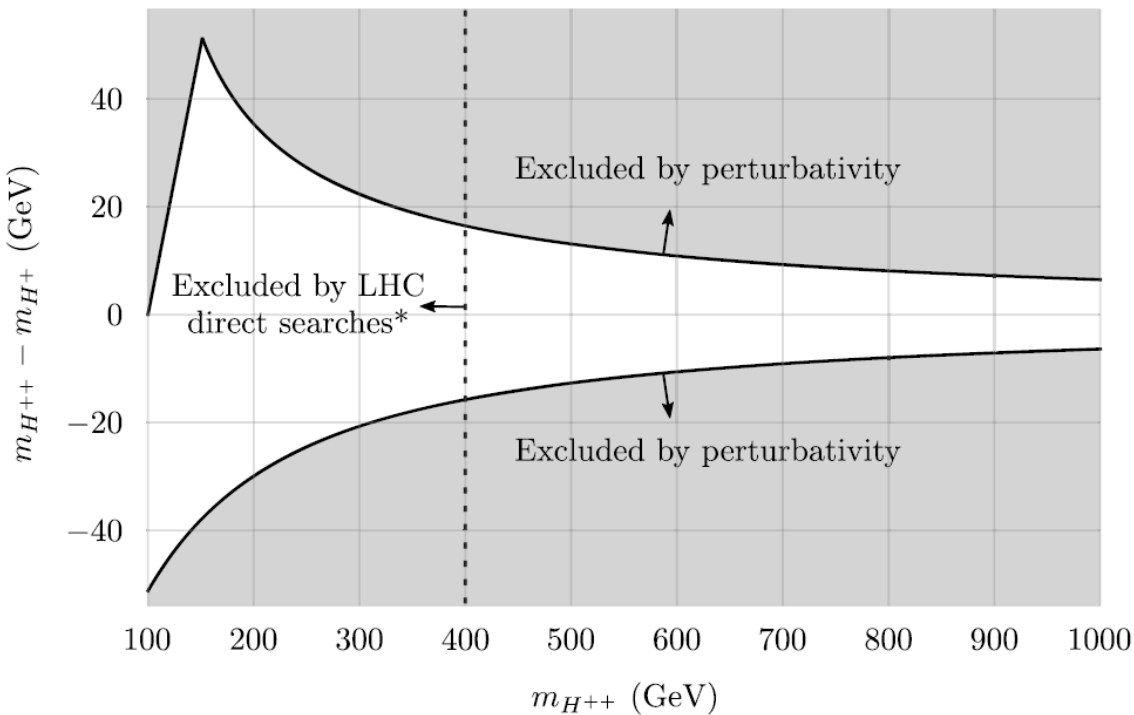


**Neutrino as
higgs benchmark**

Consistency of the triplet seesaw model revisited



Consistency of the triplet seesaw model revisited



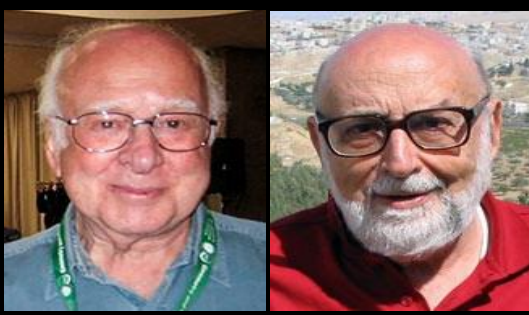
IOP Publishing

New J. Phys. 18 (2016) 033033

Neutrino as higgs benchmark

The background is a deep space scene. It features a dense field of stars of various colors and sizes. In the upper right quadrant, there is a prominent, bright yellow star with a four-pointed diffraction pattern. In the lower center, there is a colorful nebula with a core of bright blue and purple, surrounded by wisps of red and orange gas. The overall color palette is dominated by dark blues and blacks, punctuated by the colors of the stars and nebula.

How about gravity?



THE STANDARD MODEL

FERMIONS (matter)

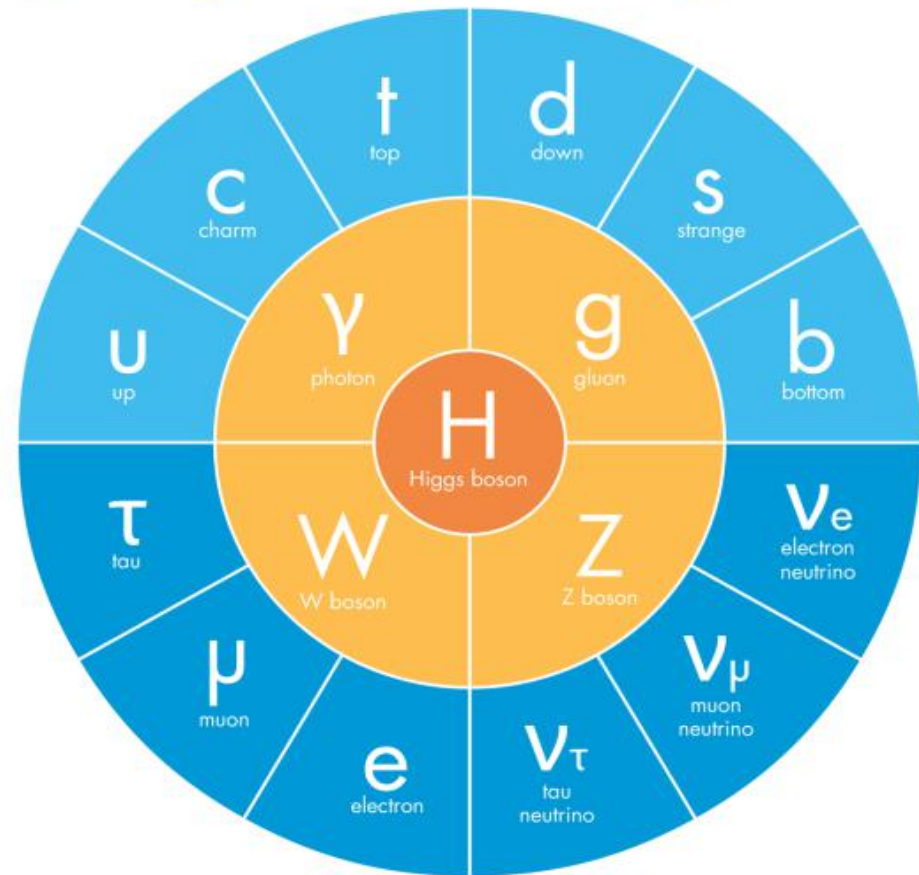
● Quarks

● Leptons

BOSONS (force carriers)

● Gauge bosons

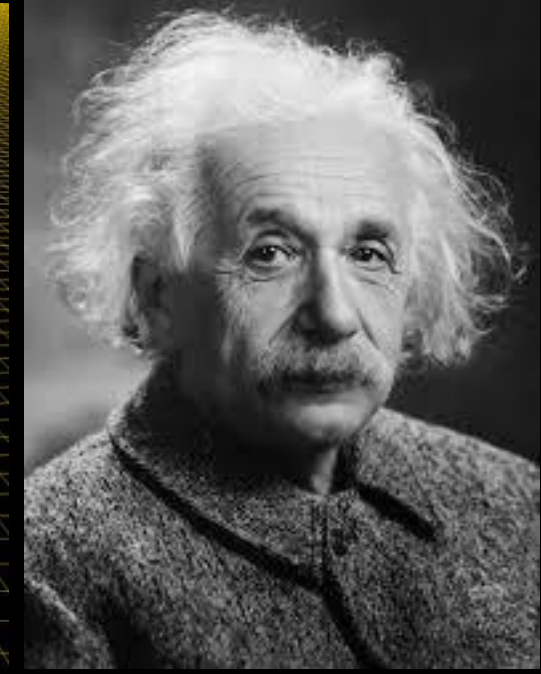
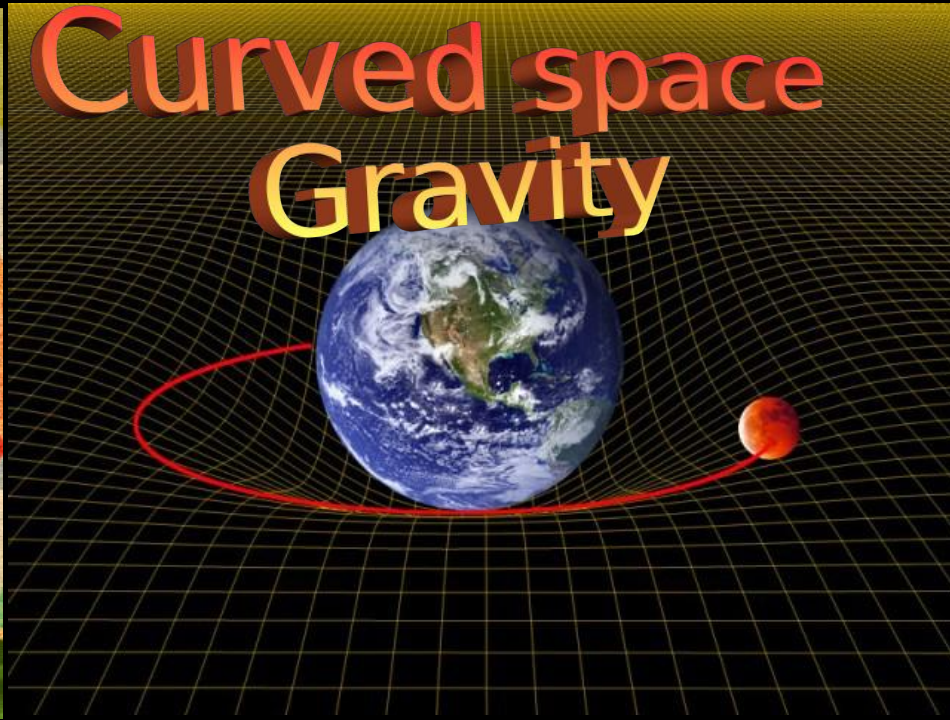
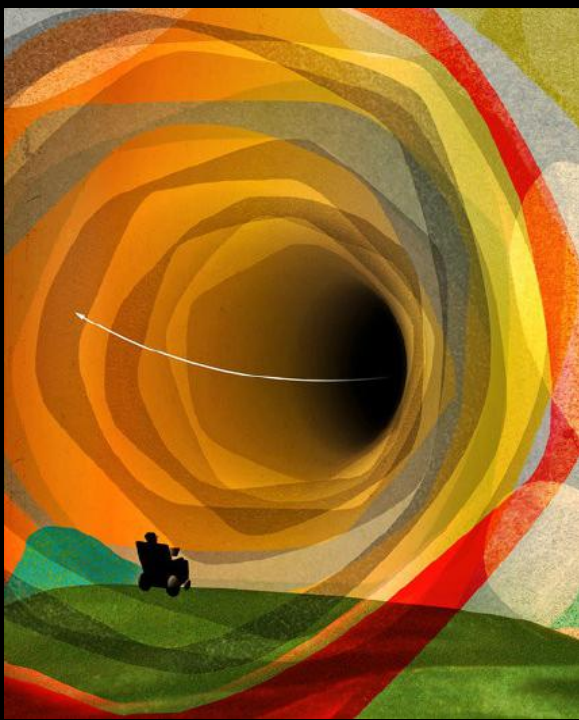
● Higgs boson

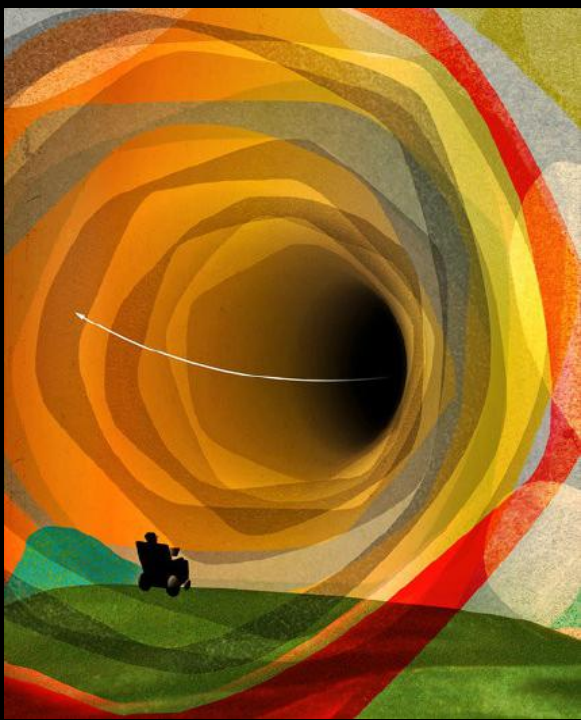


Higgs not the last brick ! ...

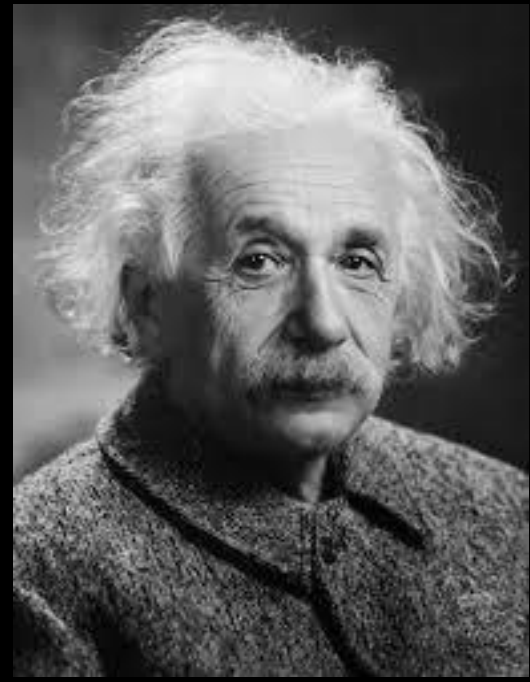


does not include gravity

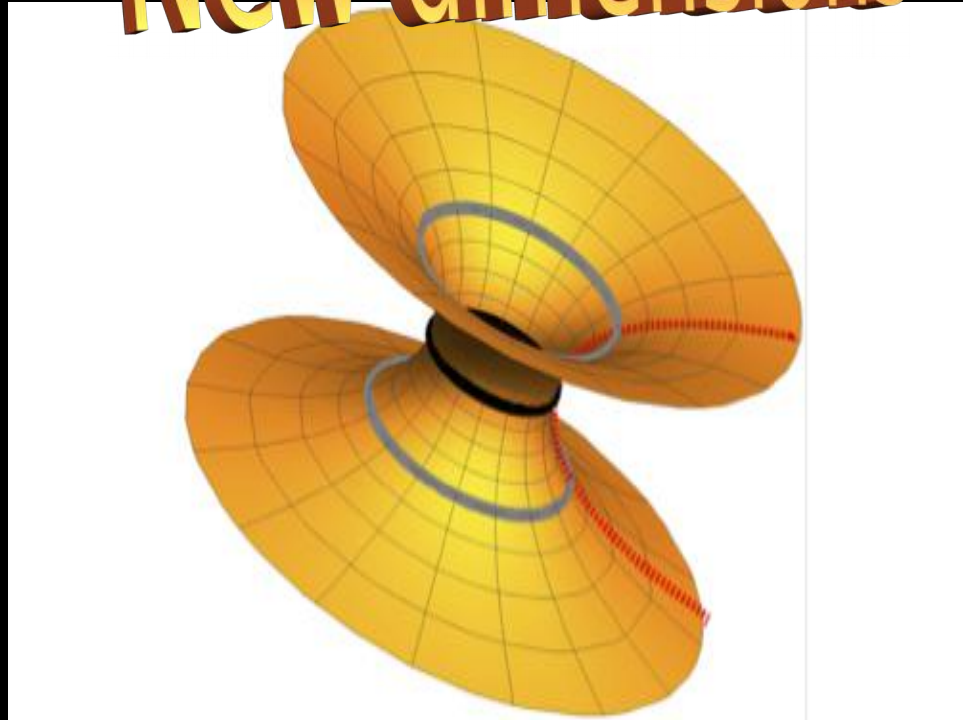




Curved space Gravity



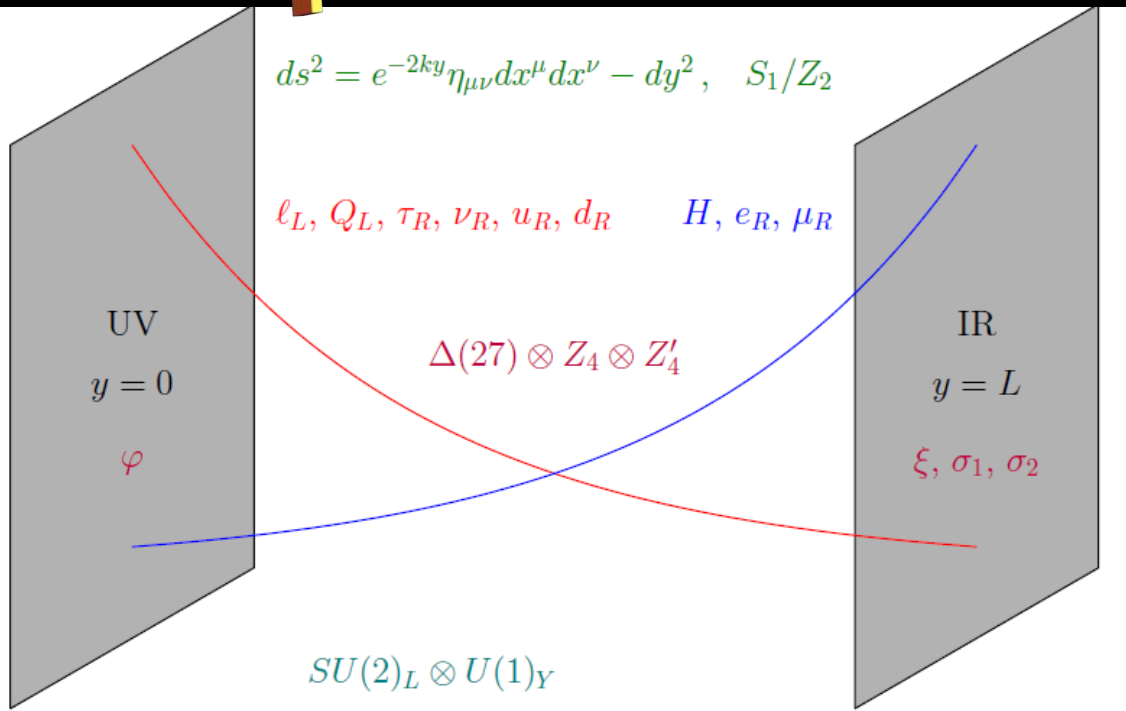
New dimensions



Warped flavor

Chen et al arXiv:1509.06683

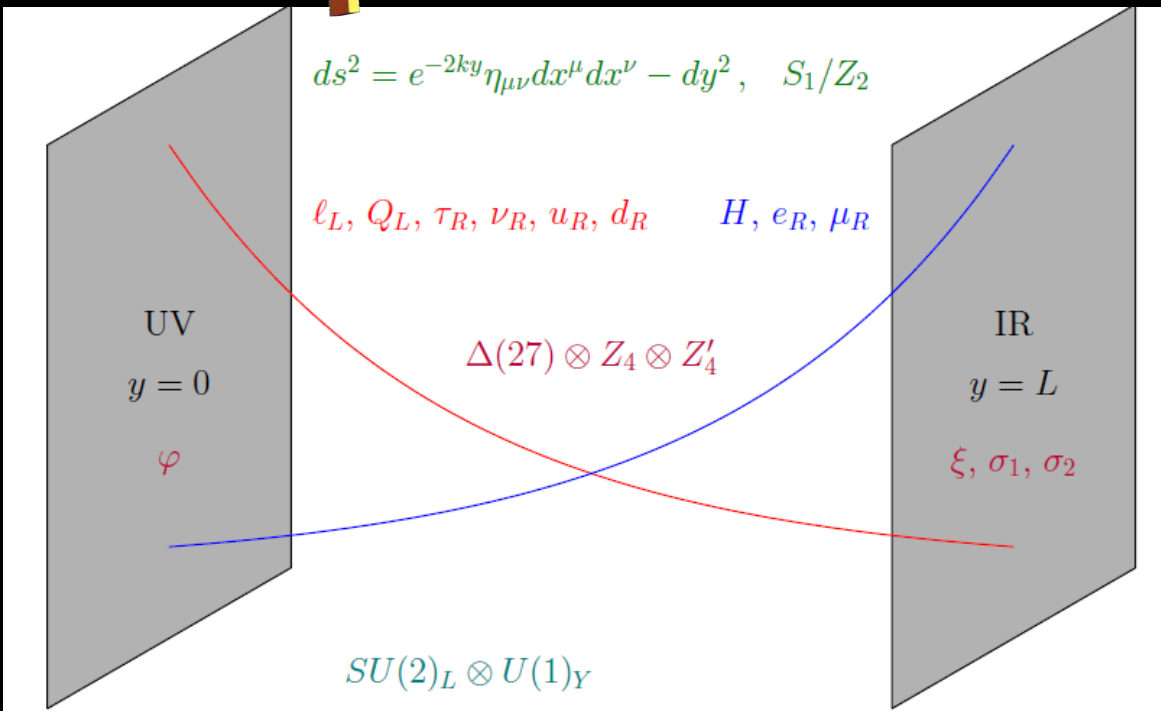
JHEP01(2016)007



Warped flavor

Chen et al arXiv:1509.06683

JHEP01(2016)007



Mass hierarchies in principle explained by judicious choices of the bulk parameters

$$\sin^2 \theta_{12} \cos^2 \theta_{13} = 1/3$$

4 neutrino mixing angles & CP phase predicted in terms of 2

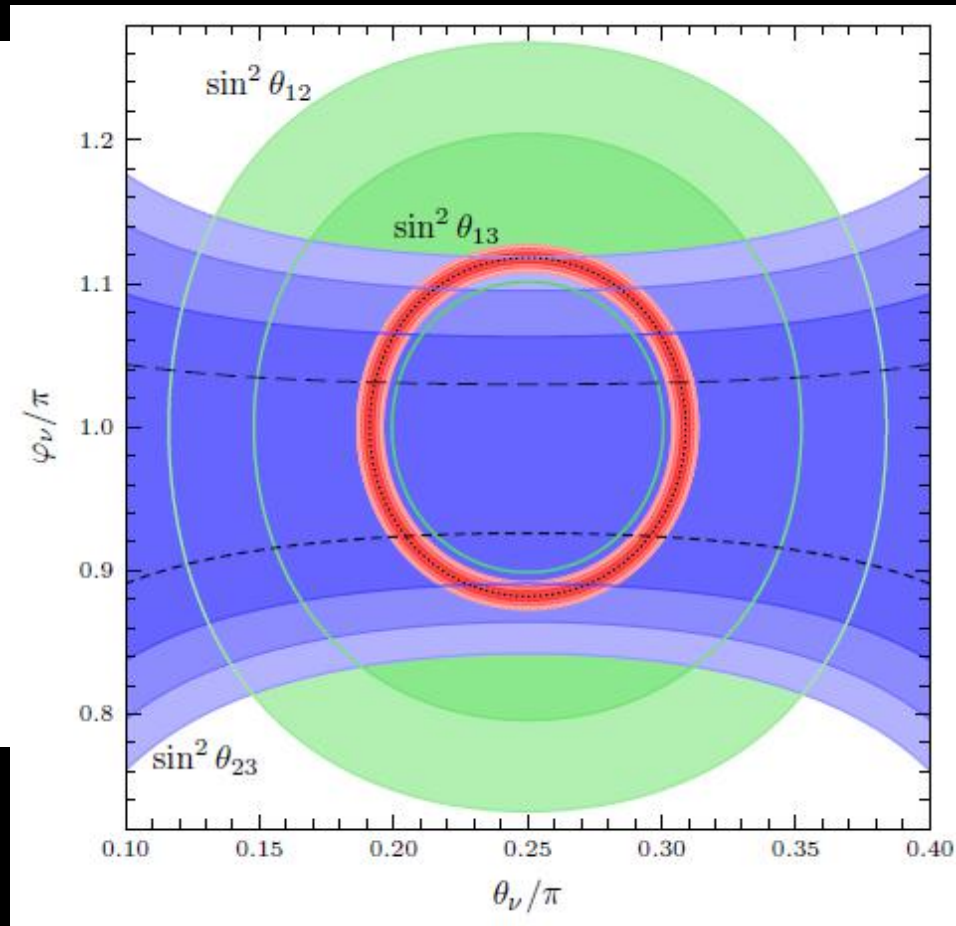
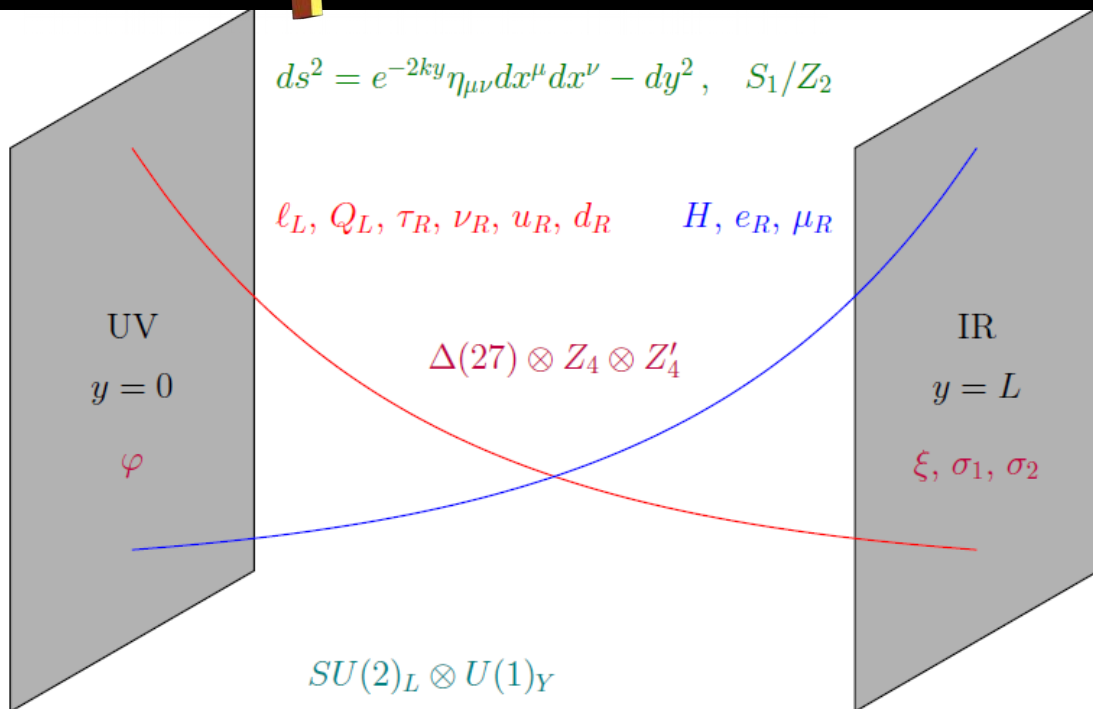
$$\sin^2 \theta_{12} = (2 - \sin 2\theta_\nu \cos \varphi_\nu)^{-1}, \quad \sin^2 \theta_{23} = \frac{1 - \sin 2\theta_\nu \sin(\pi/6 - \varphi_\nu)}{2 - \sin 2\theta_\nu \cos \varphi_\nu}$$

$$\sin^2 \theta_{13} = \frac{1}{3} (1 + \sin 2\theta_\nu \cos \varphi_\nu), \quad J_{\text{CP}} = -\frac{1}{6\sqrt{3}} \cos 2\theta_\nu.$$

Warped flavor

Chen et al arXiv:1509.06683

JHEP01(2016)007



Mass hierarchies in principle explained by judicious choices of the bulk parameters

$$\sin^2 \theta_{12} \cos^2 \theta_{13} = 1/3$$

4 neutrino mixing angles & CP phase predicted in terms of 2

$$\sin^2 \theta_{12} = (2 - \sin 2\theta_\nu \cos \varphi_\nu)^{-1},$$

$$\sin^2 \theta_{23} = \frac{1 - \sin 2\theta_\nu \sin(\pi/6 - \varphi_\nu)}{2 - \sin 2\theta_\nu \cos \varphi_\nu}$$

$$\sin^2 \theta_{13} = \frac{1}{3} (1 + \sin 2\theta_\nu \cos \varphi_\nu),$$

$$J_{\text{CP}} = -\frac{1}{6\sqrt{3}} \cos 2\theta_\nu.$$

331 from strings

10.1016/j.physletb.2016.06.015

No conventional GUT embedding :

<http://arxiv.org/abs/arXiv:1608.05334>

string completion Quiver setup

L and B conserved : no proton decay, no RPV ...

neutron-antineutron oscillations from exotic instantons

Dirac seesaw

Addazi et al arXiv:1604.02117

331 from strings

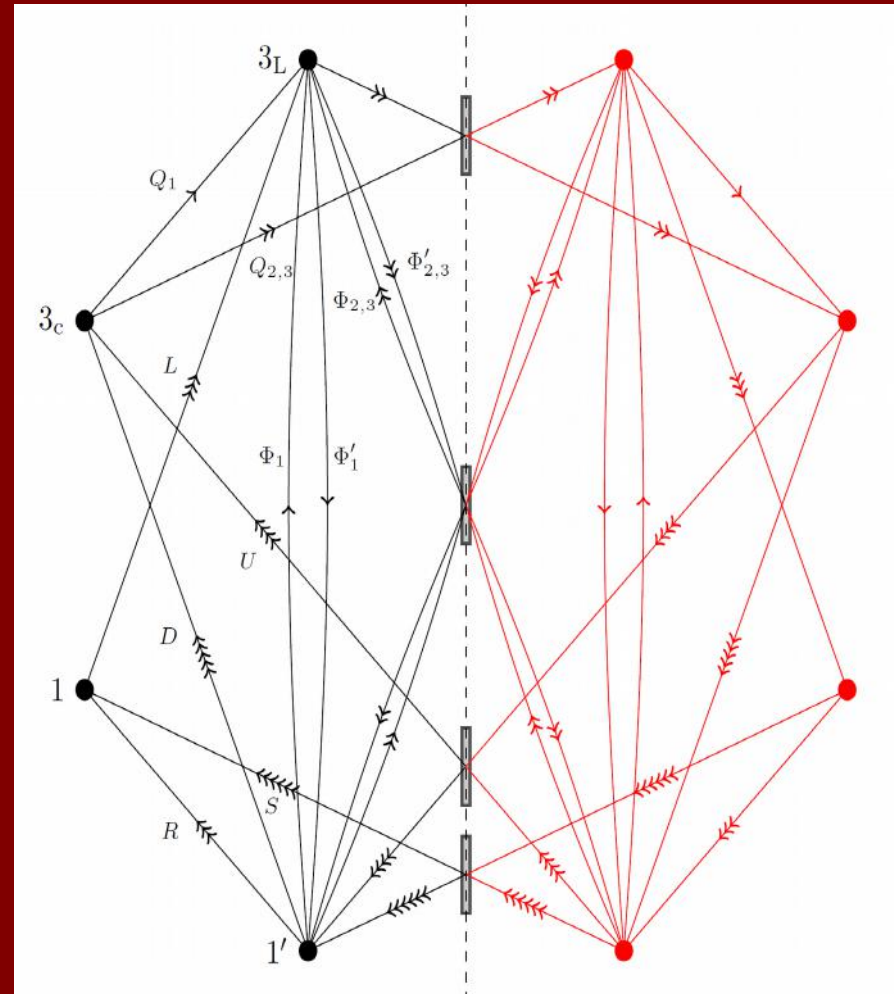
10.1016/j.physletb.2016.06.015

No conventional GUT embedding :

<http://arxiv.org/abs/arXiv:1608.05334>

string completion Quiver setup

L and B conserved : no proton decay, no RPV ...



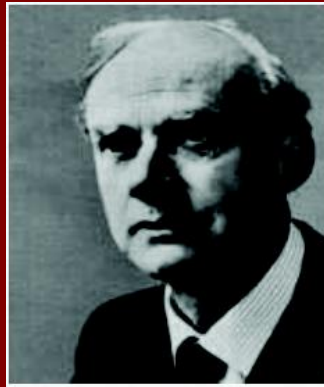
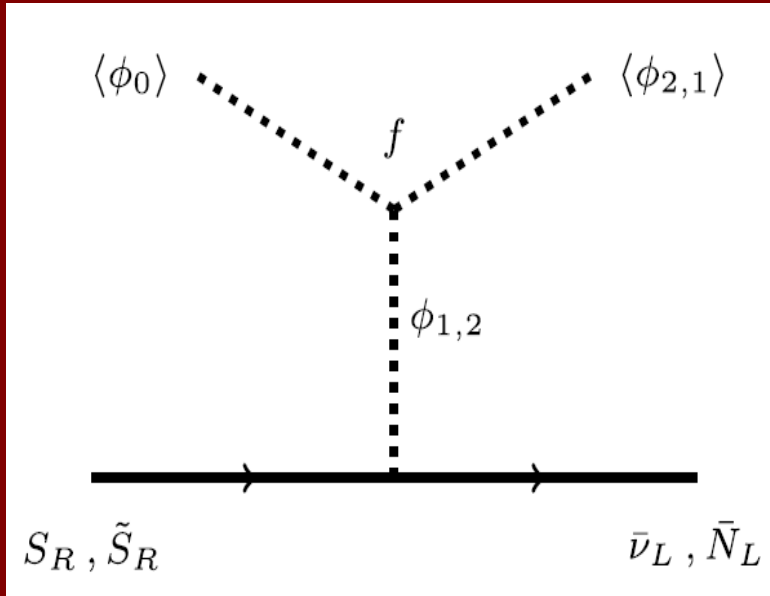
neutron-antineutron oscillations from exotic instantons

Dirac seesaw

Addazi et al arXiv:1604.02117

331 from strings

10.1016/j.physletb.2016.06.015



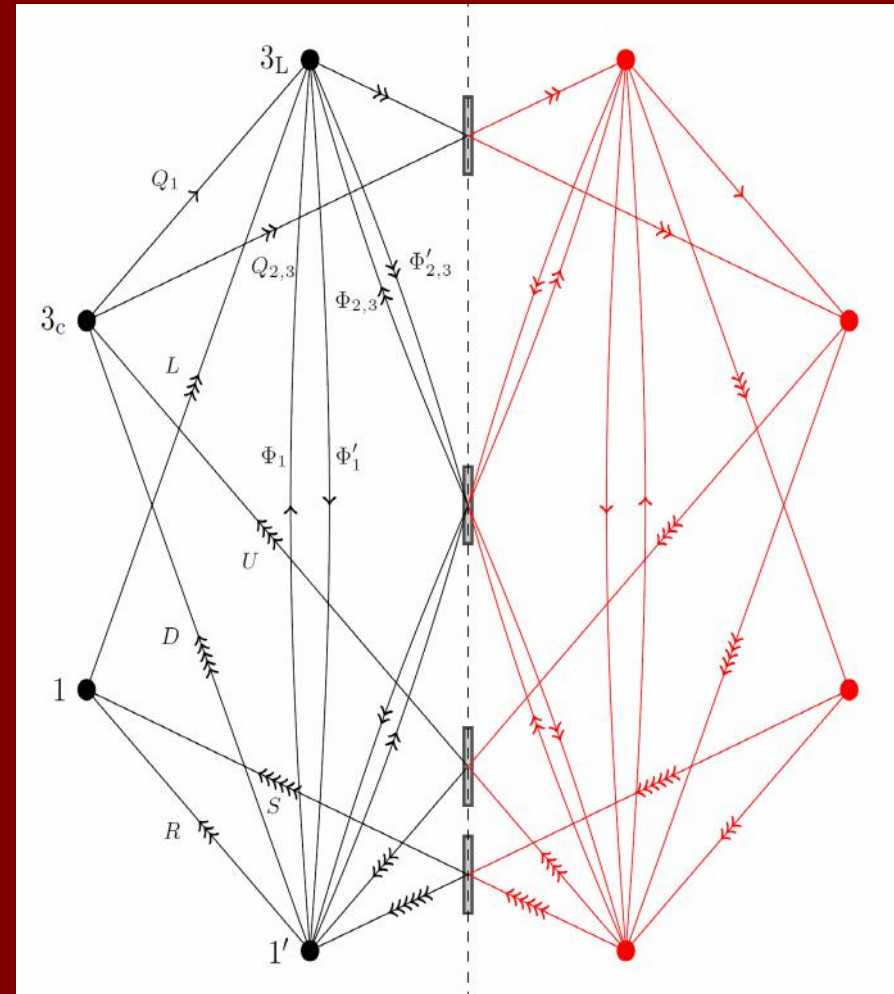
Physics Letters B 755 (2016) 363–366

No conventional GUT embedding :

<http://arxiv.org/abs/arXiv:1608.05334>

string completion Quiver setup

L and B conserved : no proton decay, no RPV ...

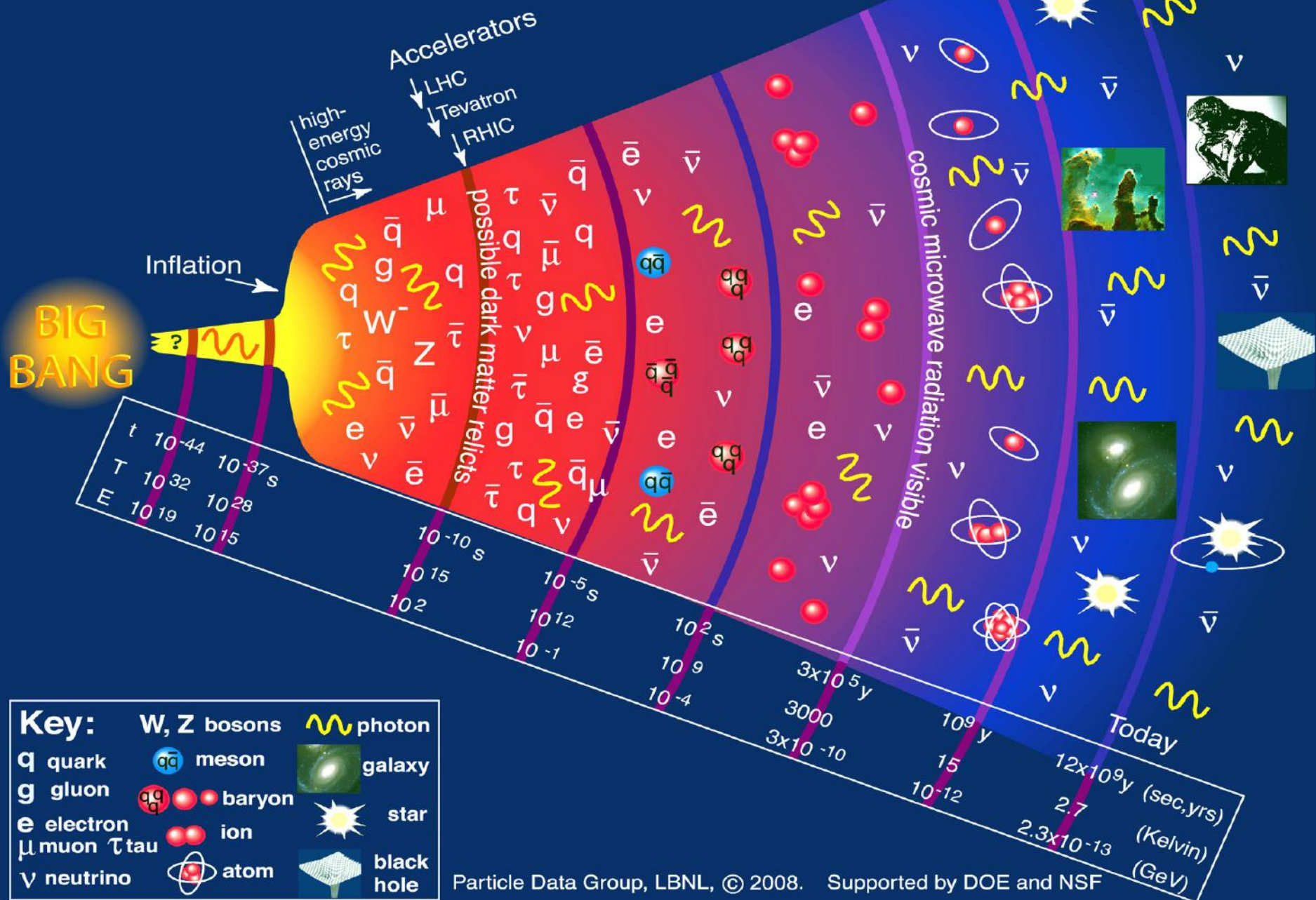


neutron-antineutron oscillations from exotic instantons

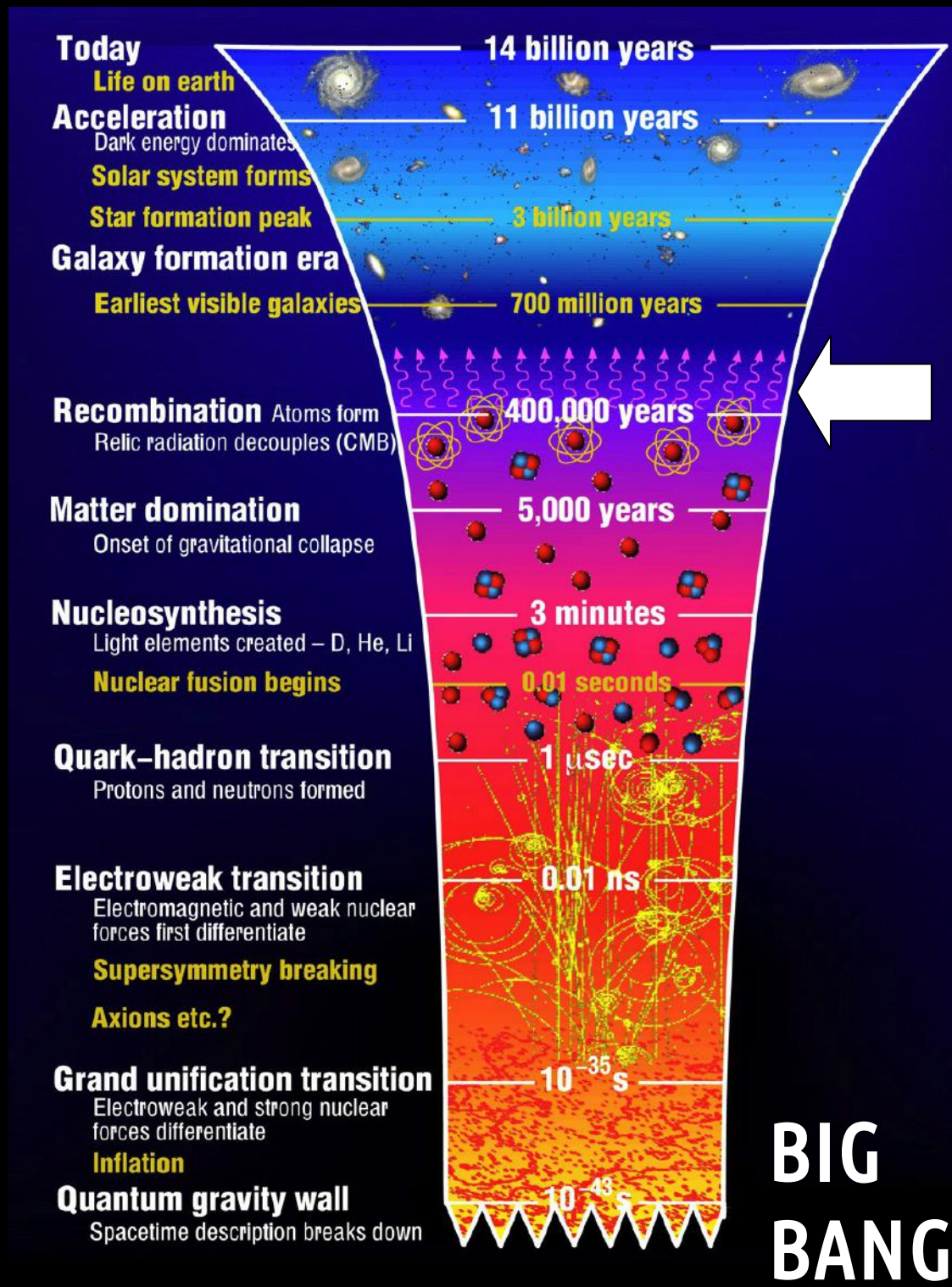


How about cosmology?

History of the Universe

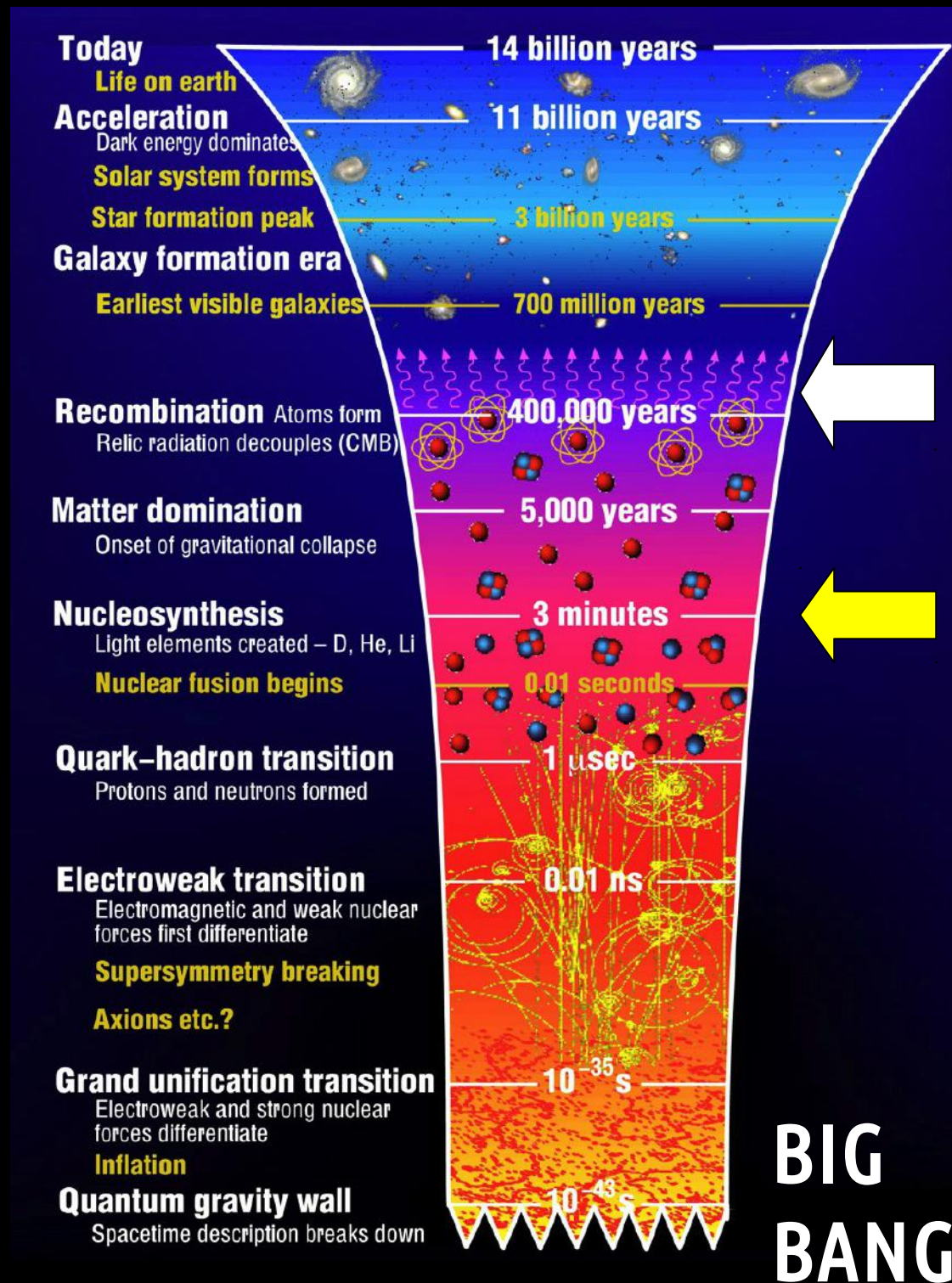


Neutrinos affect the CMB
and large scale structure
in the Universe ...



Neutrinos affect the CMB and large scale structure in the Universe ...

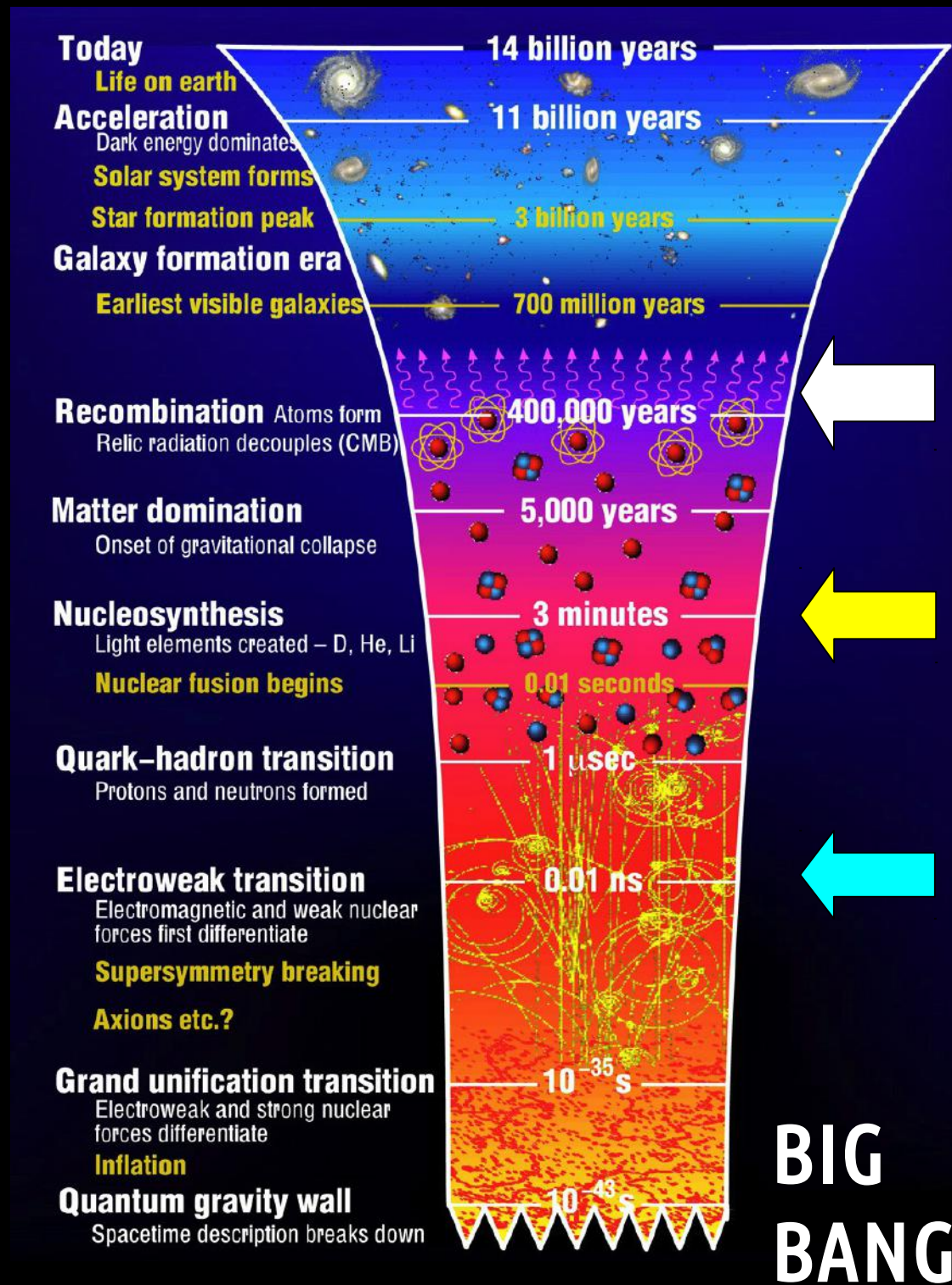
are key in the synthesis of light elements



Neutrinos affect the CMB and large scale structure in the Universe ...

are key in the synthesis of light elements

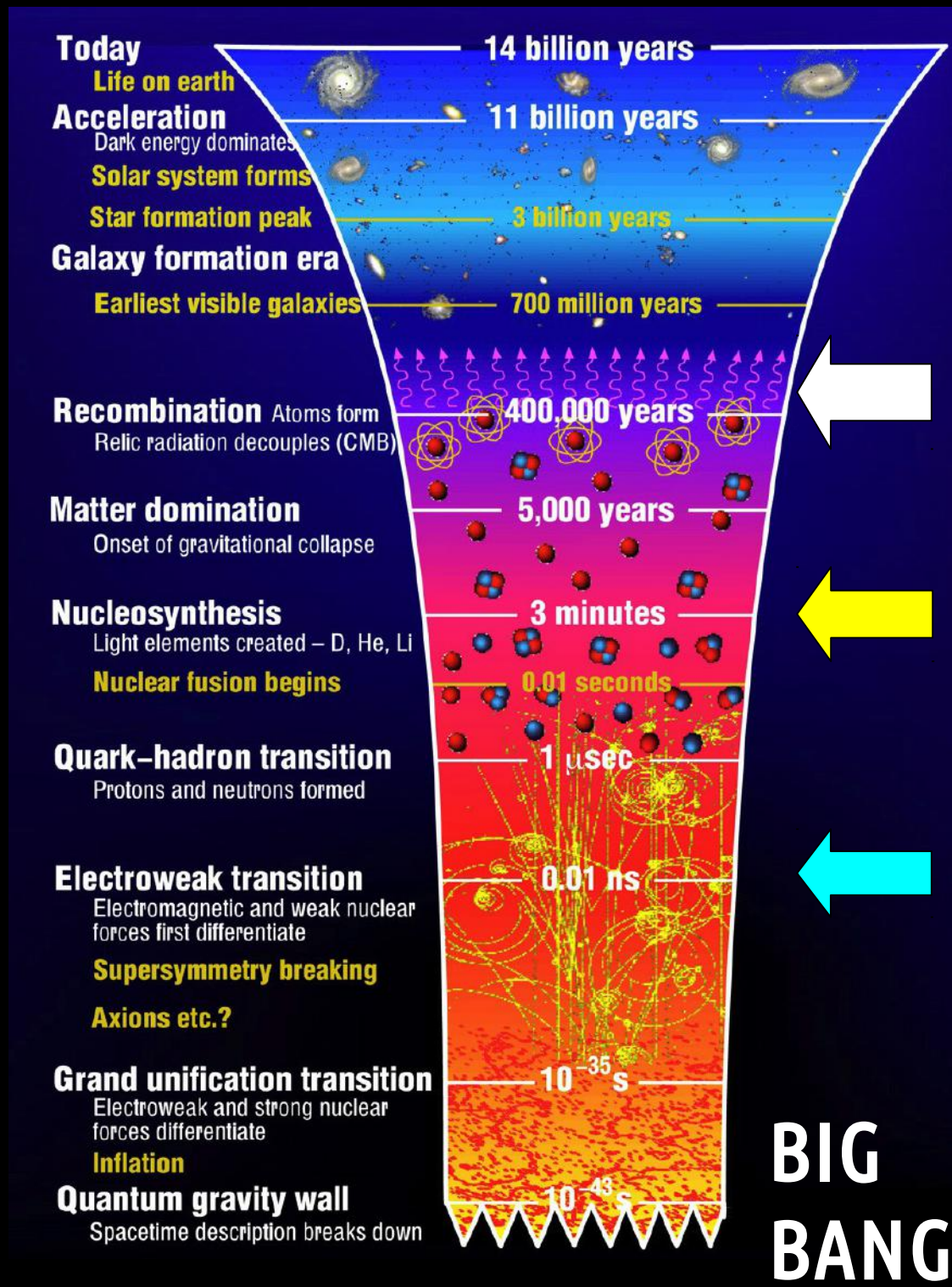
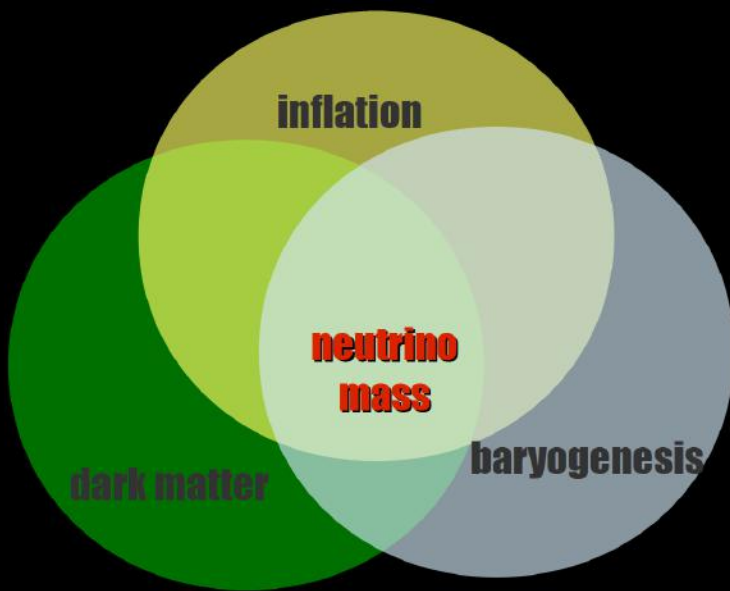
can “probe” the Universe earlier than photons ...



Neutrinos affect the CMB and large scale structure in the Universe ...

are key in the synthesis of light elements

can “probe” the Universe earlier than photons ...





The road to new physics



neutrino masses

dark matter

baryon asymmetry

Inflation

Dark energy



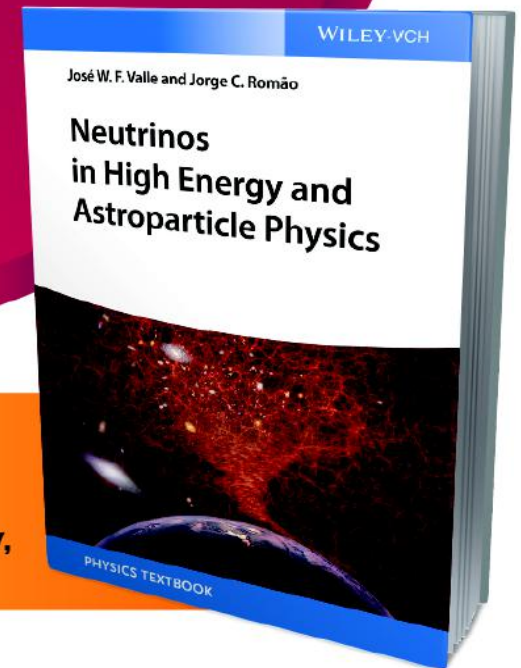
Thank you

Neutrinos in High Energy and Astroparticle Physics

*Jose Wagner Furtado Valle,
Jorge Romao*

ISBN: 978-3-527-41197-9
448 pages
February 2015

A self-contained modern advanced textbook on the role of neutrinos in astrophysics and cosmology, and high energy physics



- Written by two renowned and well-established authors in the field.
- Bridges the gap between neutrino theory and supersymmetric model building, so far missing in the current literature.
- Includes a thorough discussion of varieties of seesaw mechanism, with or without supersymmetry.
- Each chapter includes chapter summaries and further reading lists.
- Full problem sets throughout and appendices with useful tables and equations.

\$ 124.00 | £ 73.50 | € 99.90

Please note that all prices are correct at time of going to press but are subject to change without notice

WILEY-VCH

WILEY

Seesaw inflation & majoron dark matter

$$\sigma = \frac{1}{\sqrt{2}} (\langle \sigma \rangle + \rho + iJ)$$

NEUTRINO MASSES

DARK MATTER

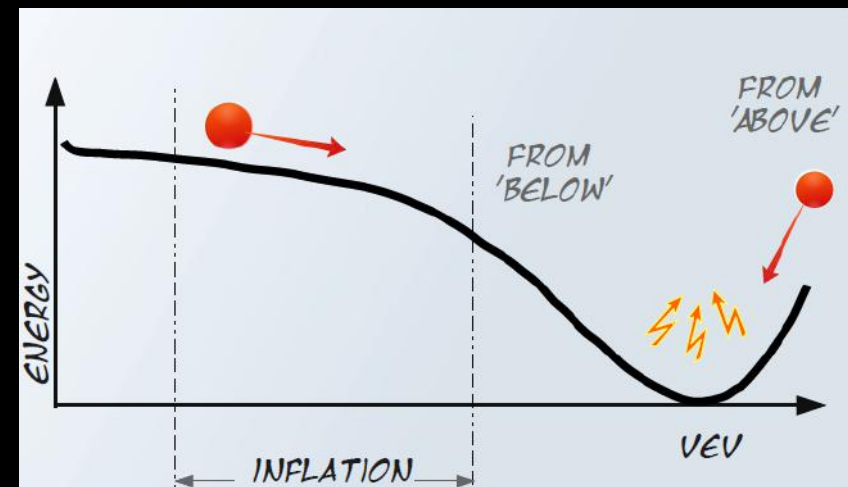
INFLATON

Boucenna et al arXiv:1405.2332

PRD90 (2014) 05502

type-I seesaw **Leptogenesis**

Aristizabal et al arXiv:1405.4706



Seesaw inflation & majoron dark matter

$$\sigma = \frac{1}{\sqrt{2}} (\langle \sigma \rangle + \rho + iJ)$$

NEUTRINO MASSES

DARK MATTER

INFLATON

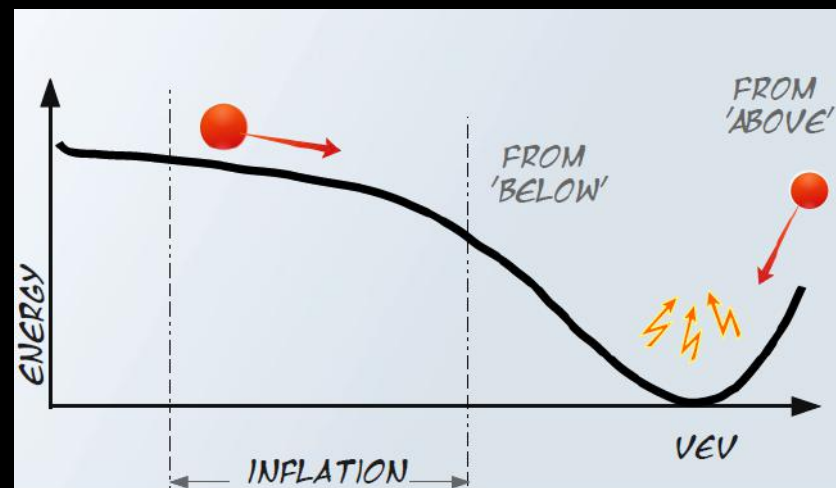
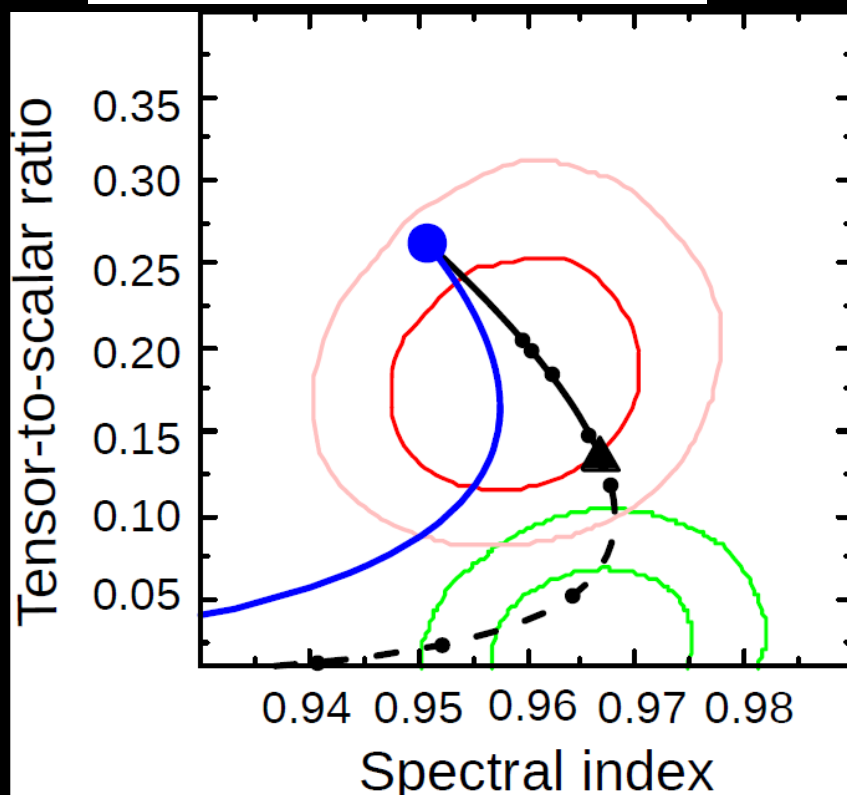
Boucenna et al arXiv:1405.2332

PRD90 (2014) 05502

type-I seesaw **Leptogenesis**

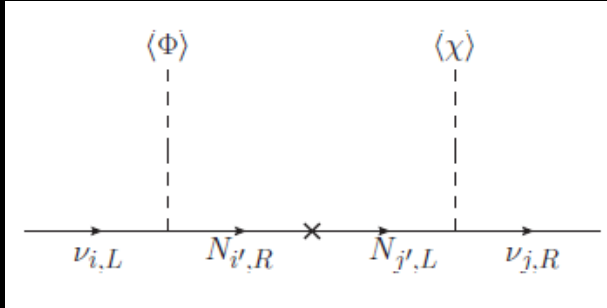
Aristizabal et al arXiv:1405.4706

Quartic versus Higgs Inflation



<http://arxiv.org/pdf/1502.00612v1>

Dark Matter Stability from Dirac nature of neutrinos



Chiulia et al

arXiv:1606.04543

arXiv:1606.06904

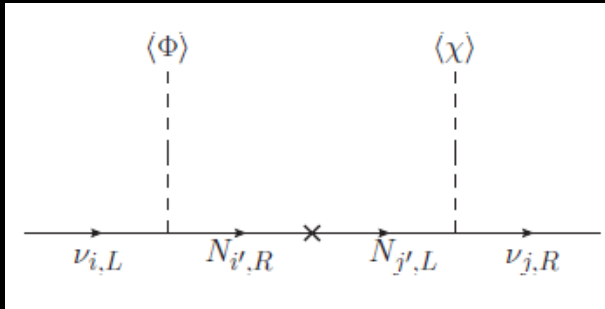
Lepton Quarticity vs Lepton number

Dark Matter Stability from Dirac nature of neutrinos

Chiulia et al

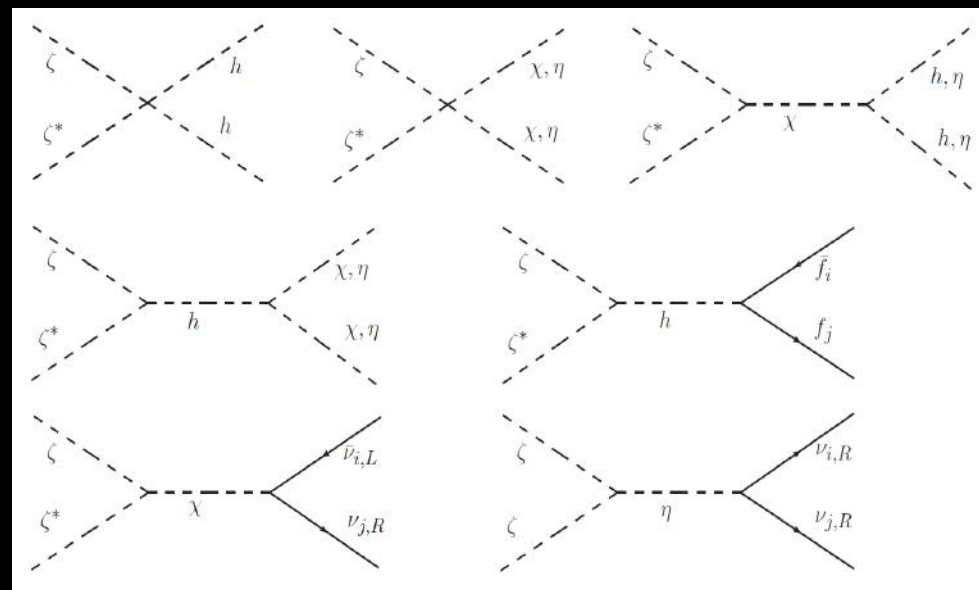
arXiv:1606.04543

arXiv:1606.06904



Lepton Quarticity vs Lepton number

Non SUSY WIMP

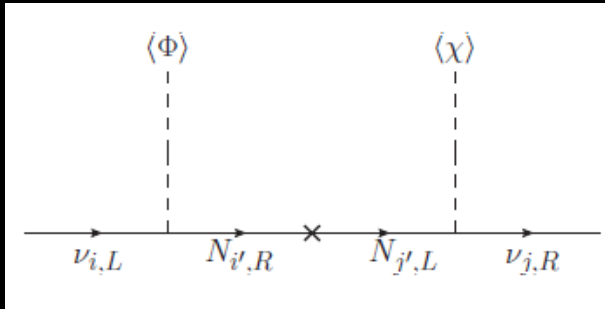


Dark Matter Stability from Dirac nature of neutrinos

Chiulia et al

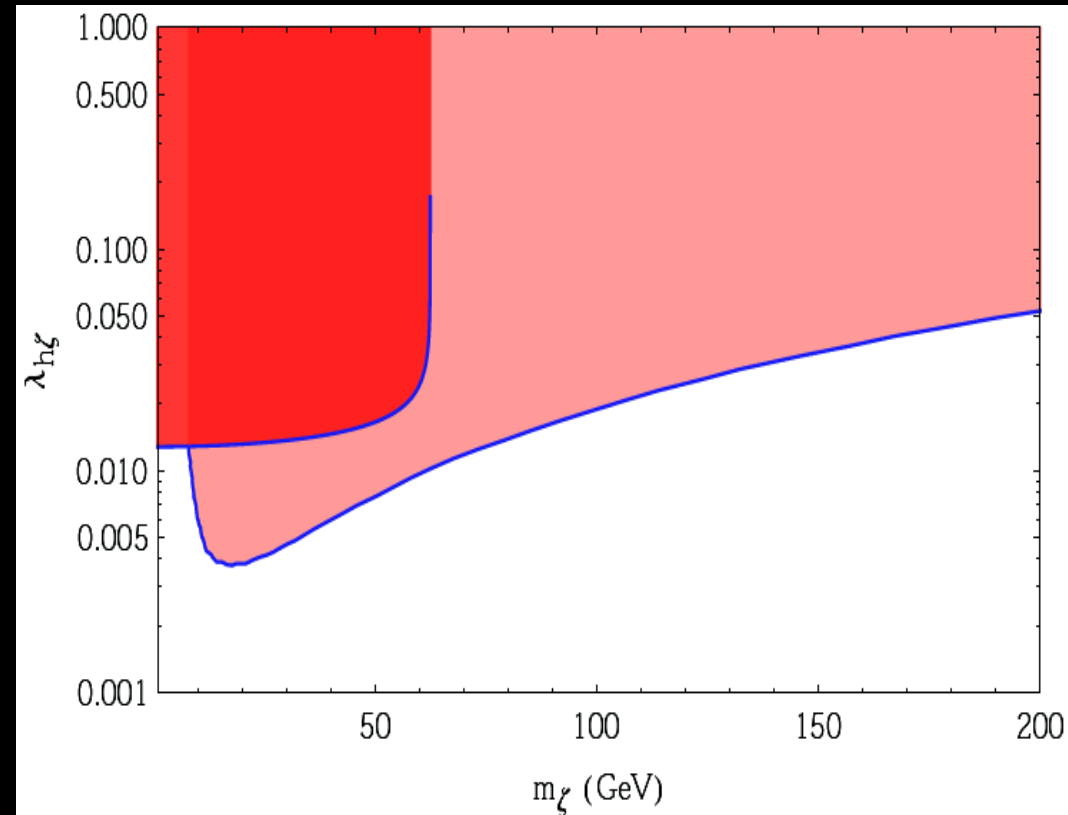
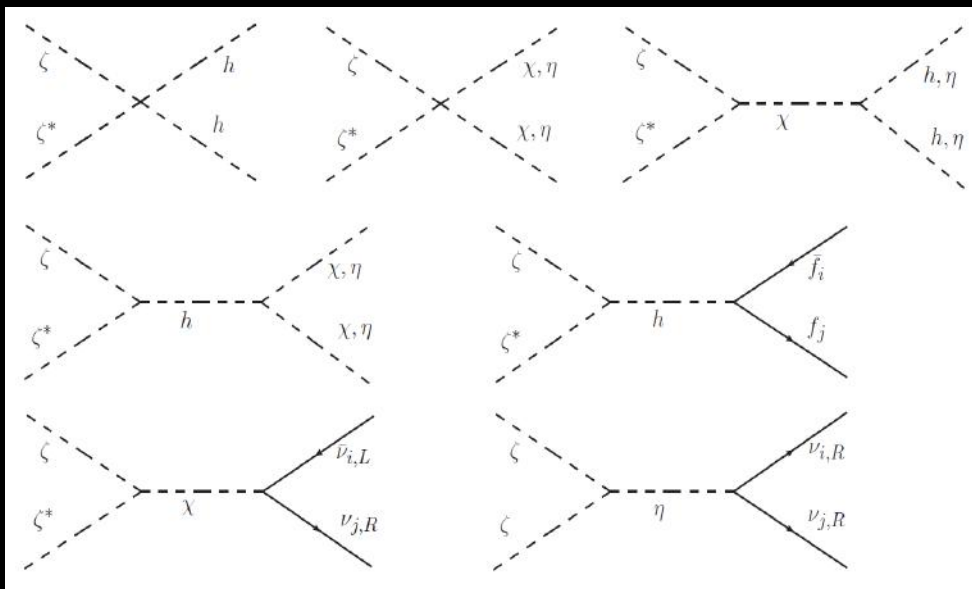
arXiv:1606.04543

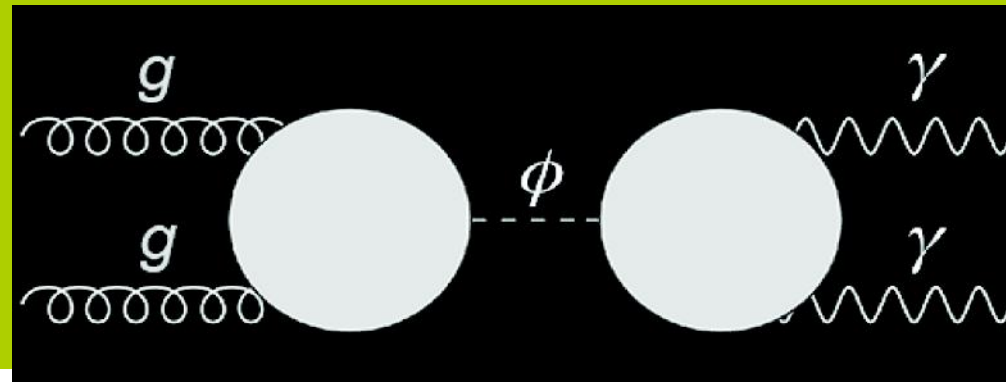
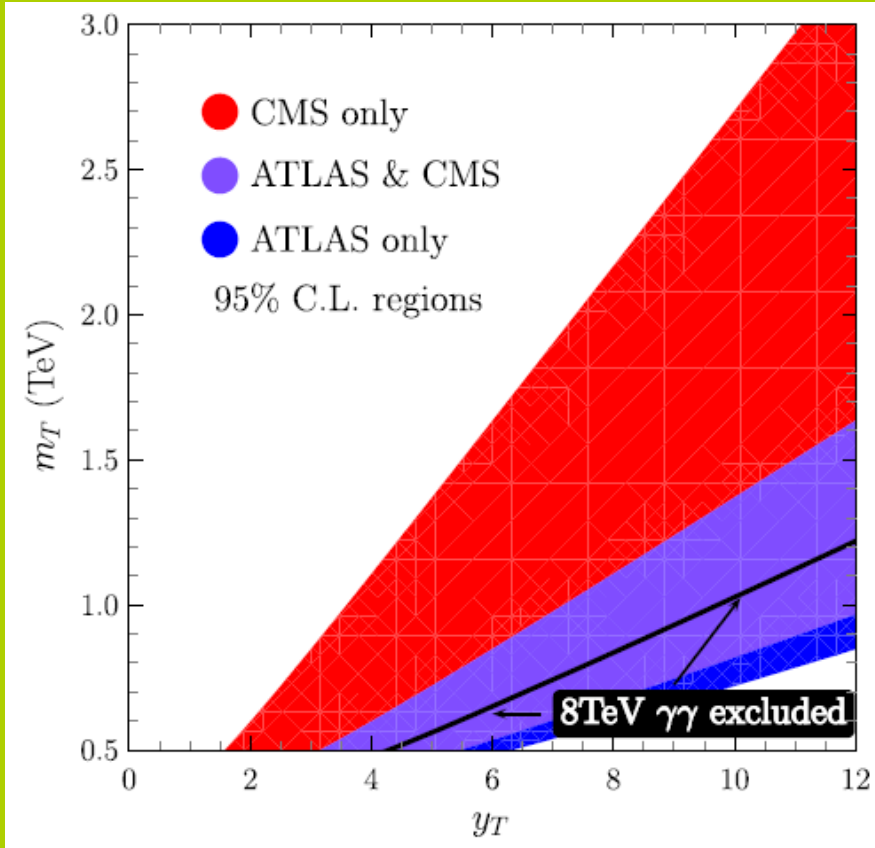
arXiv:1606.06904



Lepton Quarticity vs Lepton number

Non SUSY WIMP





Di-photon anomaly As a flavon

