

# An introduction to extra dimensions and string phenomenology

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*EISA*  
*European Institute for Sciences and Their Applications*



- 1 Physical context and motivations
- 2 High string scale - heterotic string
- 3 High string scale - type II and D-branes
- 4 Low string scale and large extra dimensions
- 5 Experimental predictions

- *An Introduction to perturbative and nonperturbative string theory*  
Ignatios Antoniadis, Guillaume Ovarlez  
e-Print: hep-th/9906108
- *Topics on String Phenomenology*  
I. Antoniadis  
e-Print: arXiv:0710.4267 [hep-th]
- *String theory in a nutshell*  
E. Kiritsis  
Princeton University Press, 2007
- *String theory and particle physics: An introduction to string phenomenology*  
Luis E. Ibanez, Angel M. Uranga  
Published in Cambridge, UK: Univ. Pr. (2012) 673 p

# Fundamental interactions

force	range	intensity of 2 protons	intensity at $10^{-16}$ cm
Gravitation	$\infty$	$10^{-38}$	$10^{-30}$
Electromagnetic	$\infty$	$10^{-2}$	$10^{-2}$
Weak (radioactivity $\beta$ )	$10^{-15}$ cm	$10^{-5}$	$10^{-2}$
Strong (nuclear forces)	$10^{-12}$ cm	1	$10^{-1}$

At what distance, gravitation becomes comparable to the other interactions?

Planck length:  $10^{-33}$  cm  $\rightarrow M_{\text{Planck}} \simeq 10^{15}$   $\times$  the LHC energy!

# Newton's law

$$m \bullet \leftarrow r \rightarrow \bullet m \quad F_{\text{grav}} = G_N \frac{m^2}{r^2} \quad G_N^{-1/2} = M_{\text{Planck}} = 10^{19} \text{ GeV}$$

Compare with electric force:  $F_{\text{el}} = \frac{e^2}{r^2} \Rightarrow$

effective dimensionless coupling  $G_N m^2$  or in general  $G_N E^2$  at energies  $E$

$$E = m_{\text{proton}} \Rightarrow \frac{F_{\text{grav}}}{F_{\text{el}}} = \frac{G_N m_{\text{proton}}^2}{e^2} \simeq 10^{-40} \quad [17]$$

$\Rightarrow$  Gravity is very weak !

# Standard Model of **electroweak** + **strong** forces

- Quantum Field Theory    Quantum Mechanics + Special Relativity
- Principle: gauge invariance     $U(1) \times SU(2) \times SU(3)$

Very accurate description of physics at present energies    17 parameters

- 1 mediators of **gauge** interactions (**vectors**): **photon**,  $W^\pm$ ,  $Z$  + **8 gluons**
- 2 matter (**fermions**): (**leptons** + quarks)  $\times 3$   
**electron, positron, neutrino**    (up, down) **3 colors**
- 3 Higgs sector: new **scalar(s)** particle(s)

# Electroweak symmetry : spontaneously broken

$SU(2) \times U(1) \rightarrow U(1)_{\text{photon}} \Rightarrow W^{\pm}, Z^0$  massive, photon massless

↙  
observed at LEP

a new particle is needed : Higgs boson (scalar)

- break the EW symmetry at  $\sim 100$  GeV
- generate mass for all elementary particles

through their interaction with the Higgs field

Englert-Brout-Higgs mechanism

Englert-Brout; Higgs; Guralnik-Hagen-Kibble '64

**Its discovery was one of the main goals of LHC**

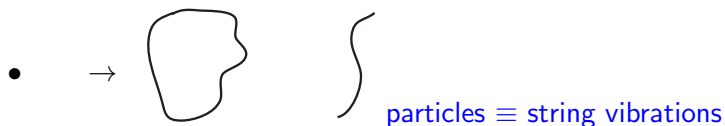
# Beyond the Standard Model : Why?

- to include gravity in a consistent quantum theory  
longstanding dream of unification of all fundamental forces of Nature
- origin of electroweak (EW) symmetry breaking  
what is behind the Brout-Englert-Higgs mechanism?
- hierarchy of masses and force intensities     EW/gravity  $\sim 10^{32}$   
stability at the quantum level  $\Rightarrow$   
fine-tuning of parameters in 32 decimal places!
- neutrino masses and oscillations
- origin of Dark Matter in the Universe



# String theory: Quantum Mechanics + General Relativity

point particle  $\rightarrow$  extended objects



- quantum gravity
- framework of unification of all interactions
- “ultimate” theory:
  - ultraviolet finite
  - no free parameters

mass scale (tension):  $M_{\text{string}} \leftrightarrow$  size:  $l_{\text{string}}$

rigid string : known particles (massless)

vibrations : infinity of massive particles

# Strings and extra dimensions

Consistent theory  $\Rightarrow$  9 spatial dimensions !

**six new dimensions of space**

matter and gauge interactions may be localized  
in less than 9 dimensions  $\Rightarrow$

**our universe on a membrane ?** [14]

$p$ -plane: extended in  $p$  spatial dimensions

$p = 0$ : particle,  $p = 1$ : string, ...

# Extra Dimensions

## how they escape observation?

finite size  $R$

Kaluza and Klein 1920

energy cost to send a signal:

$$E > R^{-1} \leftarrow \text{compactification scale}$$

## experimental limits on their size

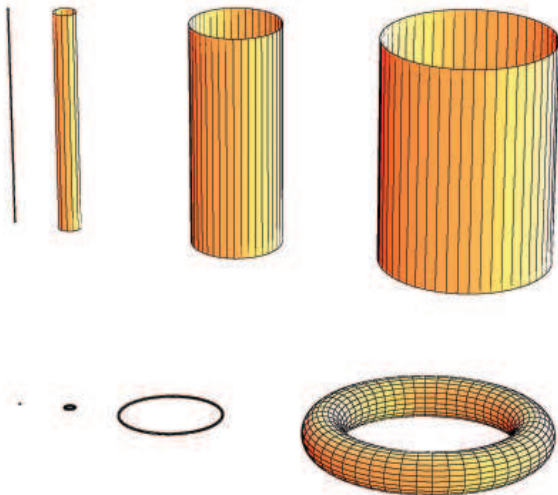
$$\text{light signal} \Rightarrow E \gtrsim 1 \text{ TeV}$$

$$R \lesssim 10^{-16} \text{ cm}$$

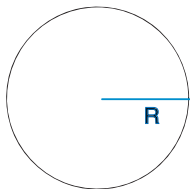
## how to detect their existence?

motion in the internal space  $\Rightarrow$  mass spectrum in 3d

# Dimensions $D=??$



example: - one internal circular dimension  
- light signal



plane waves  $e^{ipy}$  periodic under  $y \rightarrow y + 2\pi R$

$\Rightarrow$  quantization of internal momenta:  $p = \frac{n}{R}$ ;  $n = 0, 1, 2, \dots$

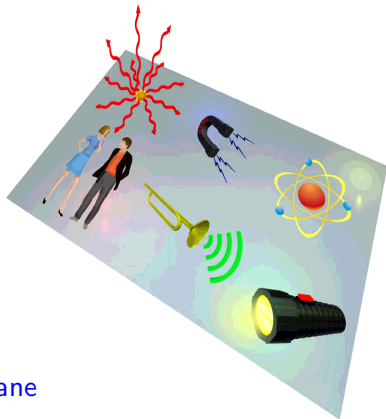
$\Rightarrow$  3d: tower of Kaluza Klein particles with masses  $M_n = n/R$

$$p_0^2 - \vec{p}^2 - p_5^2 = 0 \Rightarrow p^2 = p_5^2 = \frac{n^2}{R^2}$$

$E \gg R^{-1}$  : emission of many massive photons

$\Leftrightarrow$  propagation in the internal space [10]

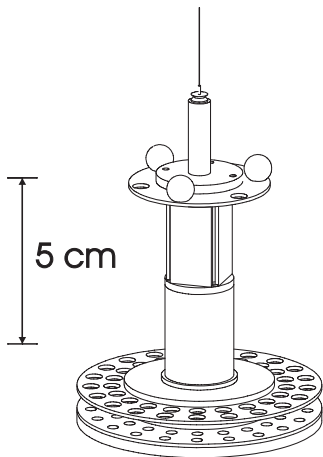
# Our universe on a membrane



Two types of new dimensions:

- longitudinal: **along the membrane**
- transverse: **“hidden” dimensions**

only gravitational signal  $\Rightarrow R_{\perp} \lesssim 1 \text{ mm} !$



$R_{\perp} \lesssim 45 \mu\text{m}$  at 95% CL

- dark-energy length scale  $\approx 85 \mu\text{m}$

# Low scale gravity

Extra large  $\perp$  dimensions can explain the apparent weakness of gravity

total force = observed force  $\times$  volume  $\perp$

total force  $\simeq \mathcal{O}(1)$  at 1 TeV

$n$  dimensions of size  $R_{\perp}$

$n = 1 : R_{\perp} \simeq 10^8$  km

excluded

$n = 2 : R_{\perp} \simeq 0.1$  mm  $(10^{-12}$  GeV)

possible

$n = 6 : R_{\perp} \simeq 10^{-13}$  mm  $(10^{-2}$  GeV)

- distances  $> R_{\perp}$  : gravity 3d

however for  $< R_{\perp}$  : gravity  $(3+n)$ d

- strong gravity at  $10^{-16}$  cm  $\leftrightarrow$   $10^3$  GeV

$10^{30}$  times stronger than thought previously! [19]



# Low scale gravity

Extra large  $\perp$  dimensions can explain the apparent weakness of gravity

total force = observed force  $\times$  volume  $\perp$  [5]

$$\begin{array}{ccccc} \uparrow & & \uparrow & & \uparrow \\ G_N^* E^{2+n} & = & G_N E^2 & \times & V_{\perp} E^n \end{array}$$

$G_N^* = M_*^{-(2+n)}$  :  $(4+n)$ -dim gravitational constant

total force  $\simeq \mathcal{O}(1)$  at 1 TeV

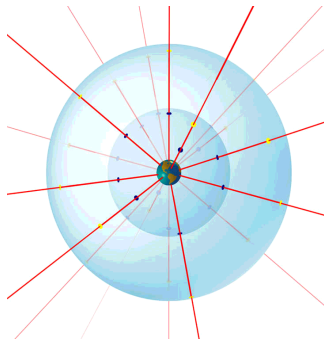
$n$  dimensions of size  $R_{\perp}$

$$\Rightarrow V_{\perp} = R_{\perp}^n$$

$$\Rightarrow M_P^2 = M_*^{2+n} R_{\perp}^n \text{ for } M_* \simeq 1 \text{ TeV} \Rightarrow (R_{\perp} M_*)^n \sim 10^{32}$$

# Gravity modification at submillimeter distances

**Newton's law:** force decreases with area



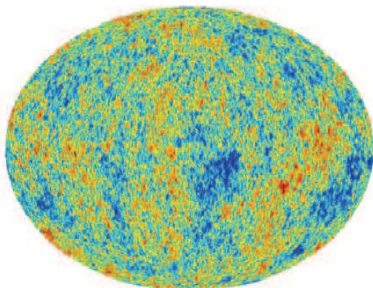
3d: force  $\sim 1/r^2$

$(3+n)$ d: force  $\sim 1/r^{2+n}$

observable for  $n = 2$ :  $1/r^4$  with  $r \ll .1$  mm [16]

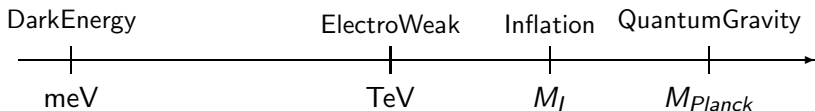
# Connect string theory to the real world

- Is string theory a tool for strong coupling dynamics  
or a theory of fundamental forces?
- If theory of Nature can string theory describe  
both particle physics and cosmology?



# Problem of scales

- describe high energy (SUSY?) extension of the Standard Model  
unification of all fundamental interactions
  - incorporate Dark Energy  
simplest case: infinitesimal (tuneable) +ve cosmological constant
  - describe possible accelerated expanding phase of our universe  
models of inflation (approximate de Sitter)
- ⇒ 3 very different scales besides  $M_{Planck}$  :



# At what energies strings may be observed?

Very different answers depending mainly on the value of the string scale  $M_s$

Before 1994:  $M_s \simeq M_{\text{Planck}} \sim 10^{18}$  GeV     $l_s \simeq 10^{-32}$  cm    After 1994:

- arbitrary parameter : Planck mass  $M_P \rightarrow$  TeV

- physical motivations  $\Rightarrow$  favored energy regions:

• High :  $\begin{cases} M_P^* \simeq 10^{18} \text{ GeV} & \text{Heterotic scale} \\ M_{\text{GUT}} \simeq 10^{16} \text{ GeV} & \text{Unification scale} \end{cases}$

• Intermediate : around  $10^{11}$  GeV ( $M_s^2/M_P \sim \text{TeV}$ )

SUSY breaking, strong CP axion, see-saw scale

• Low : TeV (hierarchy problem)

# High string scale

perturbative heterotic string : the most natural for SUSY and unification

gravity and gauge interactions have same origin

massless excitations of the closed string

But mismatch between string and GUT scales:

$$M_s = g_H M_P \simeq 50 M_{\text{GUT}} \quad g_H^2 \simeq \alpha_{\text{GUT}} \simeq 1/25 \quad [34]$$

in GUTs only one prediction from 3 gauge couplings unification:  $\sin^2 \theta_W$

introduce large threshold corrections or strong coupling  $\rightarrow M_s \simeq M_{\text{GUT}}$

but loose predictivity [25]

# Heterotic string

gravity + gauge kinetic terms [35]

$$\int [d^{10}x] \frac{1}{g_H^2} M_H^8 \mathcal{R}^{(10)} + \int [d^{10}x] \frac{1}{g_H^2} M_H^6 \mathcal{F}_{MN}^2 \quad \text{simplified units: } 2 = \pi = 1$$

Compactification in 4 dims on a 6-dim manifold of volume  $V_6 \Rightarrow$

$$\int [d^4x] \frac{V_6}{g_H^2} M_H^8 \mathcal{R}^{(4)} + \int [d^4x] \frac{V_6}{g_H^2} M_H^6 \mathcal{F}_{\mu\nu}^2$$

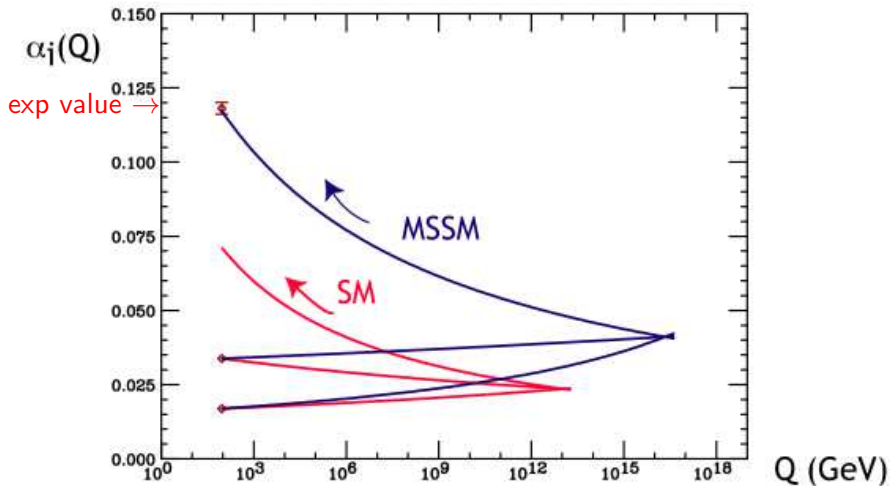
$$\begin{array}{ccc} \parallel & \parallel & \Rightarrow \\ M_P^2 & 1/g^2 & \end{array}$$

$$M_P^2 = \frac{1}{g^2} M_H^2 \quad \frac{1}{g^2} = \frac{1}{g_H^2} V_6 M_H^6 \quad \Rightarrow \quad M_H = g M_P \quad g_H = g \sqrt{V_6} M_H^3$$

$$g_H \lesssim 1 \Rightarrow V_6 \sim \text{string size}$$

# GUT prediction of QCD coupling

input  $\alpha_{em}, \sin^2 \theta_W \Rightarrow$  output  $\alpha_3$  [35]





# Open strings and D-branes

string propagation in space-time  $\Rightarrow$  2-dim world-sheet  $(\tau, \sigma)$   $X^\mu(\tau, \sigma)$

$\tau$ : time,  $\sigma \in [0, \pi]$ : spatial extension of the string

closed strings  $\Rightarrow \sigma$ : periodic  $X^\mu(\tau, 0) = X^\mu(\tau, \pi)$

open string  $\Rightarrow$  endpoints:  $\sigma = 0, \pi$  world-sheet boundaries  
they also carry gauge charges

D-branes = hypersurfaces where open strings can end

$D_p$ -brane: parallel dimensions:  $X^1, \dots, X^p$  (also time  $X^0$ )

$\partial_\sigma X^\mu = 0$  at  $\sigma = 0$  normal derivative vanishes

Newmann boundary conditions  $\Rightarrow$  free propagation along the boundary

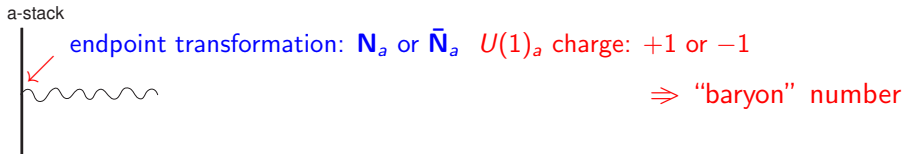
transverse dimensions:  $X^{p+1}, \dots, X^9$

$X^\mu = X_0^\mu$  at  $\sigma = 0$  ( $\partial_\tau X^\mu = 0$  at  $\sigma = 0$ )

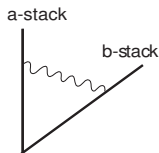
Dirichlet conditions: endpoint fixed at the boundary

# D-brane spectrum

Generic spectrum:  $N$  coincident branes  $\Rightarrow U(N)$



- open strings from the same stack  $\Rightarrow$  adjoint gauge multiplets of  $U(N_a)$
- stretched between two stacks  $\Rightarrow$  bifundamentals of  $U(N_a) \times U(N_b)$



non-oriented strings  $\Rightarrow$  also:

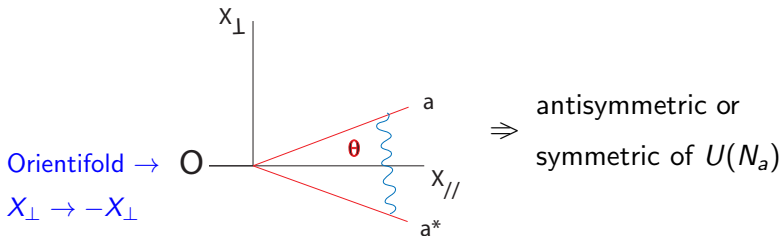
- orthogonal and symplectic groups  $SO(N)$ ,  $Sp(N)$
- matter in antisymmetric + symmetric reps

# Non oriented strings $\Rightarrow$ orientifold planes

where closed strings change orientation

$\Rightarrow$  mirror branes identified with branes under orientifold action

- strings stretched between two mirror stacks



# Minimal Standard Model embedding

General analysis using 3 brane stacks [52]

$$\Rightarrow U(3) \times U(2) \times U(1)$$

antiquarks  $u^c, d^c$   $(\bar{3}, 1)$  :

antisymmetric of  $U(3)$  or bifundamental  $U(3) \leftrightarrow U(1)$

$\Rightarrow$  3 models: antisymmetric is  $u^c, d^c$  or none

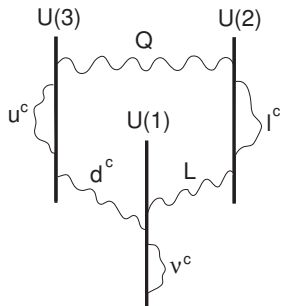
$N_i$  stack of D-branes:  $U(N_i) = SU(N_i) \times U(1)_i$

gauge couplings:  $\alpha_{N_i} = \frac{g_{N_i}^2}{4\pi}$  and  $\alpha_i$

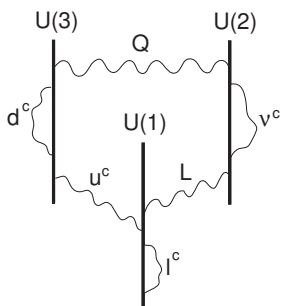
normalization:  $\text{Tr } T^a T^b = \frac{1}{2} \delta^{ab} \Rightarrow \alpha_i = \frac{\alpha_{N_i}}{2N_i}$

$$Y = c_1 Q_1 + c_2 Q_2 + c_3 Q_3 \Rightarrow \frac{1}{g_Y^2} = \frac{2c_1^2}{g_1^2} + \frac{4c_2^2}{g_2^2} + \frac{6c_3^2}{g_3^2}$$

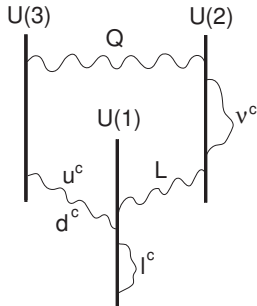
$$\sin^2 \theta_W = \frac{g_Y^2}{g_2^2 + g_Y^2} = \frac{1}{g_2^2/g_Y^2 + 1} = \frac{1}{1 + 4c_2^2 + 2c_1^2 g_2^2/g_1^2 + 6c_3^2 g_2^2/g_3^2}$$



**Model A**



**Model B**

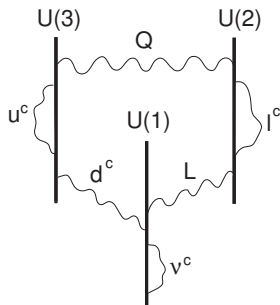


**Model C**

$Q$	$(\mathbf{3}, \mathbf{2}; 1, 1, 0)_{1/6}$
$u^c$	$(\bar{\mathbf{3}}, \mathbf{1}; 2, 0, 0)_{-2/3}$
$d^c$	$(\bar{\mathbf{3}}, \mathbf{1}; -1, 0, \varepsilon_d)_{1/3}$
$L$	$(\mathbf{1}, \mathbf{2}; 0, -1, \varepsilon_L)_{-1/2}$
$l^c$	$(\mathbf{1}, \mathbf{1}; 0, 2, 0)_1$
$\nu^c$	$(\mathbf{1}, \mathbf{1}; 0, 0, 2\varepsilon_\nu)_0$

$Q$	$(\mathbf{3}, \mathbf{2}; 1, \varepsilon_Q, 0)_{1/6}$
$u^c$	$(\bar{\mathbf{3}}, \mathbf{1}; -1, 0, 1)_{-2/3}$
$d^c$	$(\bar{\mathbf{3}}, \mathbf{1}; 2, 0, 0)_{1/3}$
$L$	$(\mathbf{1}, \mathbf{2}; 0, \varepsilon_L, 1)_{-1/2}$
$l^c$	$(\mathbf{1}, \mathbf{1}; 0, 0, -2)_1$
$\nu^c$	$(\mathbf{1}, \mathbf{1}; 0, 2\varepsilon_\nu, 0)_0$

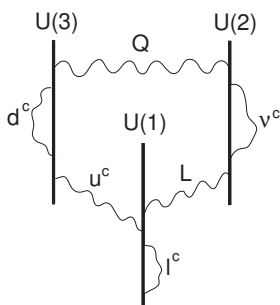
$Q$	$(\mathbf{3}, \mathbf{2}; 1, \varepsilon_Q, 0)_{1/6}$
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$l^c$	$(\mathbf{1}, \mathbf{1}; 0, 0, -2)_1$
$\nu^c$	$(\mathbf{1}, \mathbf{1}; 0, 2\varepsilon_\nu, 0)_0$



**Model A**

$$Y_A = -\frac{1}{3}Q_3 + \frac{1}{2}Q_2$$

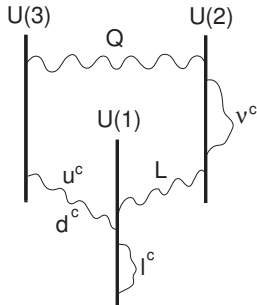
$$\sin^2 \theta_W = \frac{1}{2 + 2\alpha_2/3\alpha_3} \Big|_{\alpha_2=\alpha_3} = \frac{3}{8}$$



**Model B**

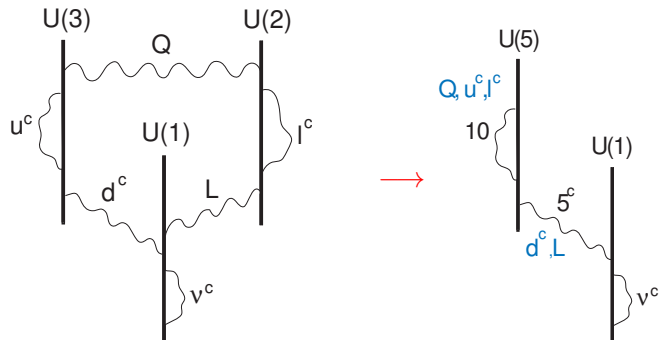
$$Y_{B,C} = \frac{1}{6}Q_3 - \frac{1}{2}Q_1$$

$$\frac{1}{1 + \alpha_2/2\alpha_1 + \alpha_2/6\alpha_3} \Big|_{\alpha_2=\alpha_3} = \frac{6}{7 + 3\alpha_2/\alpha_1}$$



**Model C**

# $SU(5)$ GUT





# Intersecting branes: 'perfect' for SM embedding

product of unitary gauge groups (brane stacks) and bi-fundamental reps  
but no unification: no prediction for  $M_s$ , independent gauge couplings

however GUTs: problematic:

- no perturbative  $SO(10)$  spinors
- no top-quark Yukawa coupling in  $SU(5)$ :  $10 10 5_H$   
 $SU(5)$  is part of  $U(5) \Rightarrow U(1)$  charges :  $10$  charge 2 ;  $5_H$  charge  $\pm 1$   
 $\Rightarrow$  cannot balance charges with  $SU(5)$  singlets  
can be generated by D-brane instantons but ...

→ Non-perturbative M/F-theory models:

combine good properties of heterotic and intersecting branes

but lack exact description for systematic studies

# Type I string theory $\Rightarrow$ D-brane world

I.A.-Arkani-Hamed-Dimopoulos-Dvali '98

- gravity: closed strings propagating in 10 dims
- gauge interactions: open strings with their ends attached on D-branes

Dimensions of finite size:  $n$  transverse  $6 - n$  parallel [36]

calculability  $\Rightarrow R_{\parallel} \simeq l_{\text{string}}$  ;  $R_{\perp}$  arbitrary

$$M_p^2 \simeq \frac{1}{g_s^2} M_s^{2+n} R_{\perp}^n \quad g_s = \alpha : \text{weak string coupling [22]}$$

Planck mass in  $4 + n$  dims:  $M_*^{2+n}$

$$M_s \sim 1 \text{ TeV} \Rightarrow R_{\perp}^n = 10^{32} l_s^n \quad \text{small } M_s/M_p \Rightarrow \text{extra-large } R_{\perp}$$

$$R_{\perp} \sim .1 - 10^{-13} \text{ mm for } n = 2 - 6$$

distances  $< R_{\perp}$  : gravity  $(4+n)$ -dim  $\rightarrow$  strong at  $10^{-16}$  cm

Type I/II strings: gravity and gauge interactions have different origin

gravity + gauge kinetic terms

$$\int [d^{10}x] \frac{1}{g_s^2} M_s^8 \mathcal{R}^{(10)} + \int [d^{p+1}x] \frac{1}{g_s} M_s^{p-3} \mathcal{F}_{MN}^2 \quad [23]$$

Compactification in 4 dims  $\Rightarrow$

$$\int [d^4x] \frac{V_6}{g_s^2} M_s^8 \mathcal{R}^{(4)} + \int [d^4x] \frac{V_{\parallel}}{g_s} M_s^{p-3} \mathcal{F}_{\mu\nu}^2 \quad V_6 = V_{\parallel} V_{\perp}$$

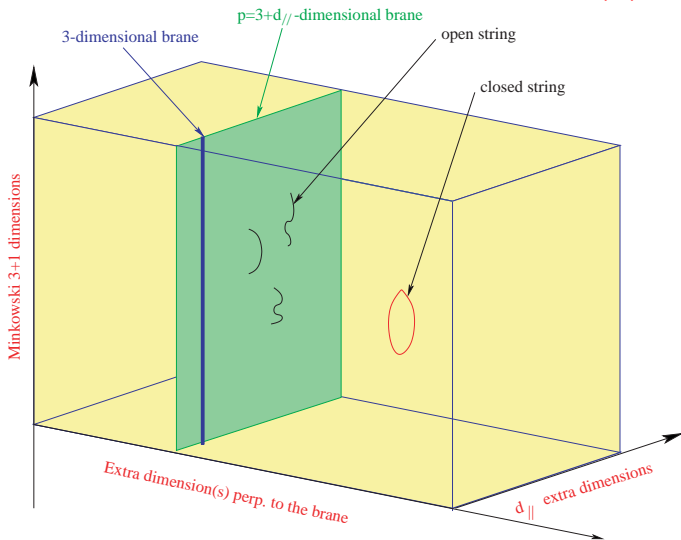
$$\begin{array}{ccc} \parallel & \parallel & \Rightarrow \\ M_P^2 & 1/g^2 & \\ \Rightarrow & g_s = g^2 V_{\parallel} M_s^{p-3} \lesssim 1 \Rightarrow V_{\parallel} \sim \text{string size} & \end{array}$$

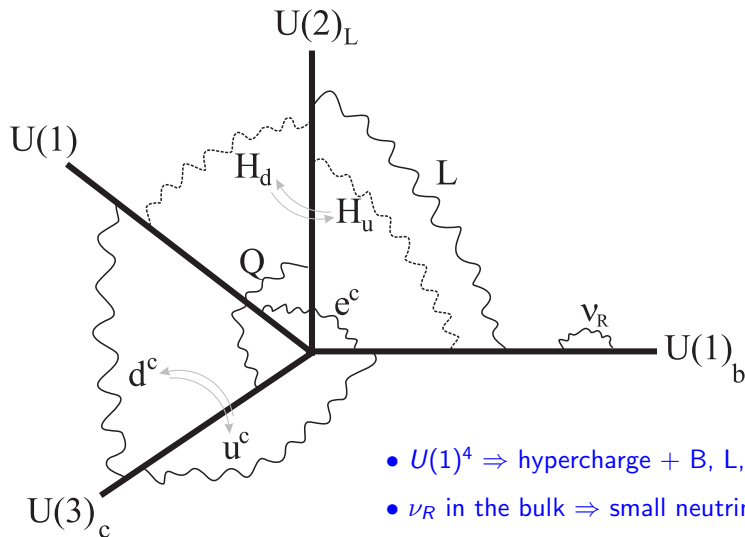
$$\Rightarrow M_P^2 = \frac{V_{\perp}}{g_s^2} M_s^{2+n} \quad g_s \simeq g^2$$

# Braneworld

2 types of compact extra dimensions:

- parallel ( $d_{\parallel}$ ):  $\lesssim 10^{-16}$  cm (TeV) [34]
- transverse ( $\perp$ ):  $\lesssim 0.1$  mm (meV)





- $U(1)^4 \Rightarrow$  hypercharge + B, L, PQ global
- $\nu_R$  in the bulk  $\Rightarrow$  small neutrino masses

# R-neutrinos: in the bulk

Arkani Hamed-Dimopoulos-Dvali-March Russell '98

Dienes-Dudas-Gherghetta '98 Dvali-Smirnov '98

R-neutrino:  $\nu_R(x, y)$   $y$ : bulk coordinates

$$S_{int} = g_s \int d^4x H(x) L(x) \nu_R(x, y=0)$$

$\langle H \rangle = v \Rightarrow$  mass-term:  $\frac{g_s v}{R_\perp^{n/2}} \nu_L \nu_R^0 \leftarrow$  4d zero-mode

Dirac neutrino masses:  $m_\nu \simeq \frac{g_s v}{R_\perp^{n/2}} \simeq v \frac{M_*}{M_p}$

$\simeq 10^{-3} - 10^{-2}$  eV for  $M_* \simeq 1 - 10$  TeV

$m_\nu \ll 1/R_\perp \Rightarrow$  KK modes unaffected

# Experimental predictions

- No little hierarchy problem:

radiative electroweak symmetry breaking with no logs

$$\Lambda \sim \text{a few TeV} \quad \text{and} \quad m_H^2 = \text{a loop factor} \times \Lambda^2$$

- particle accelerators

- Large TeV dimensions    seen by gauge interactions

- Extra large hidden dimensions    transverse  $\Rightarrow$  strong gravity

- other accelerator signatures

- microgravity experiments

- gravity modifications at short distances

- new submillimeter forces

# Accelerator signatures: 4 different scales

- Gravitational radiation in the bulk  $\Rightarrow$  missing energy

present LHC bounds:  $M_* \gtrsim 4 - 9$  TeV

- Massive string vibrations  $\Rightarrow$  e.g. resonances in dijet distribution [42]

$$M_j^2 = M_0^2 + M_s^2 j \quad ; \quad \text{maximal spin : } j + 1$$

higher spin excitations of quarks and gluons with strong interactions

present LHC limits:  $M_s \gtrsim 7$  TeV

- Large TeV dimensions  $\Rightarrow$  KK resonances of SM gauge bosons I.A. '90

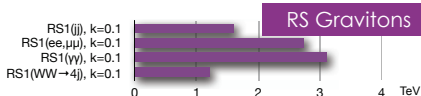
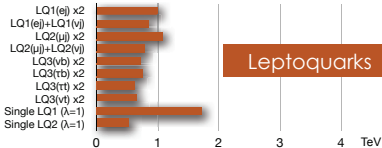
$$M_k^2 = M_0^2 + k^2/R^2 \quad ; \quad k = \pm 1, \pm 2, \dots$$

experimental limits:  $R^{-1} \gtrsim 0.5 - 4$  TeV (UED - localized fermions) [45]

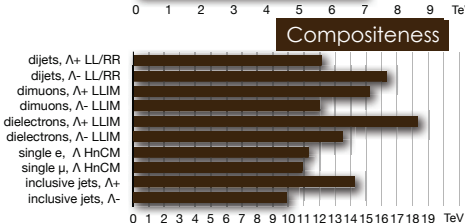
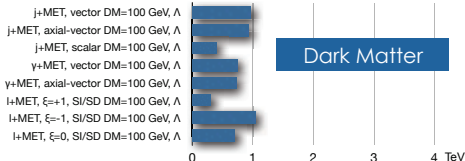
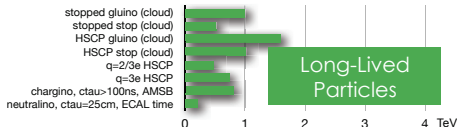
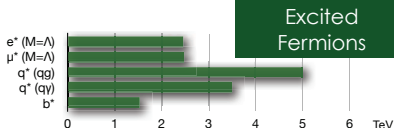
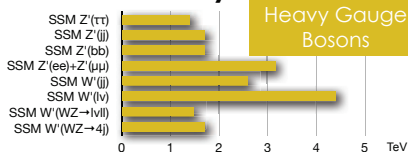
- extra  $U(1)$ 's and anomaly induced terms

masses suppressed by a loop factor from  $M_s$  [50]





# CMS Preliminary

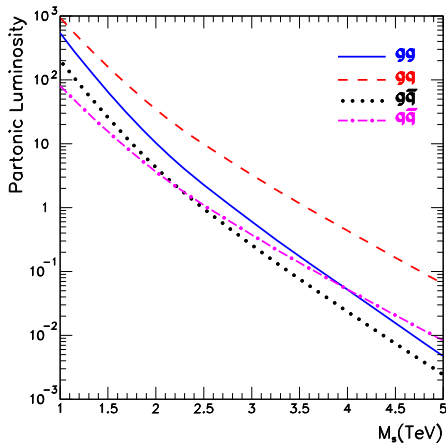


Tree level superstring amplitudes involving at most 2 fermions and gluons:  
 model independent for any compactification, # of susy's, even none  
 no intermediate exchange of KK, windings or graviton emission  
 Universal sum over infinite exchange of string (Regge) excitations

Parton luminosities in pp above TeV  
 are dominated by  $gq$ ,  $gg$

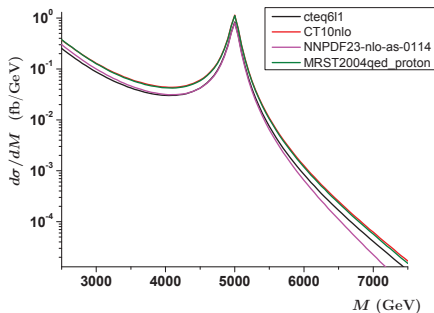
⇒ model independent

$gq \rightarrow gq$ ,  $gg \rightarrow gg$ ,  $gg \rightarrow q\bar{q}$



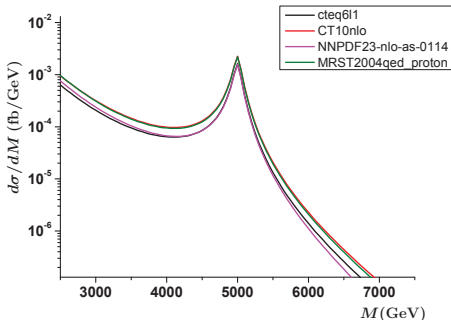
# String Resonances production at Hadron Colliders

I.A.-Anchordoqui-Dai-Feng-Goldberg-Huang-Lüst-Stojkovic-Taylor '14



$M_s = 5$  TeV:

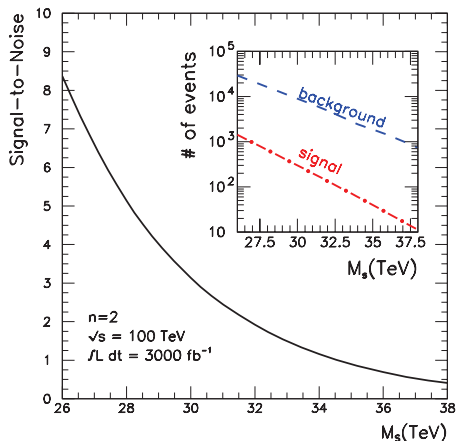
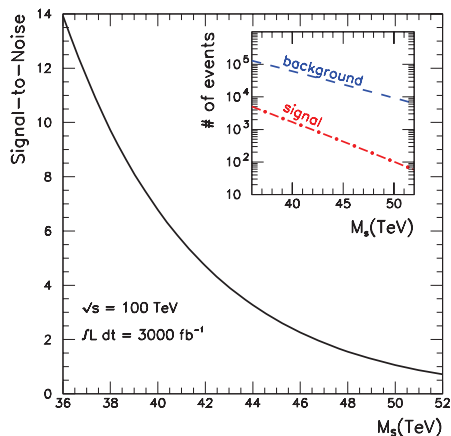
dijet at LHC14



$\gamma$ +jet

# String Resonances production at Hadron Colliders

I.A.-Anchordoqui-Dai-Feng-Goldberg-Huang-Lüst-Stojkovic-Taylor '14

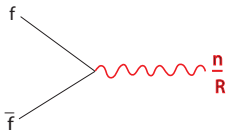


[40]

## Localized fermions (on 3-brane intersections)

⇒ single production of KK modes

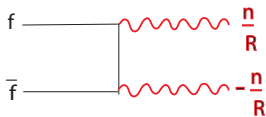
I.A.-Benakli '94



- strong bounds indirect effects:  $R^{-1} \gtrsim 3 \text{ TeV}$
- new resonances but at most  $n = 1$

## Otherwise KK momentum conservation [47]

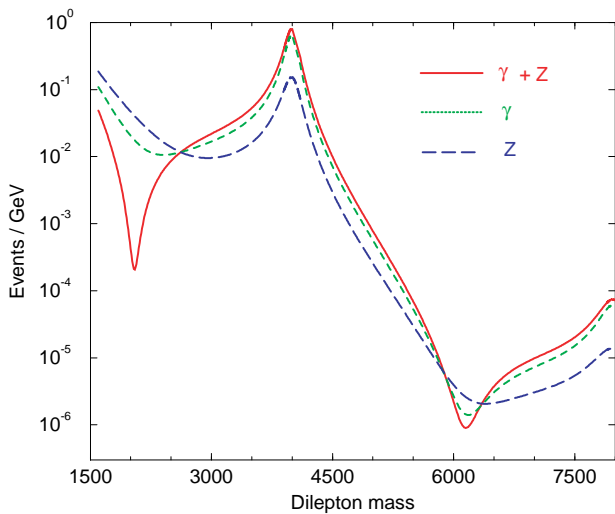
⇒ pair production of KK modes (universal dims)



- weak bounds  $R^{-1} \gtrsim 500 \text{ GeV}$
- no resonances
- lightest KK stable ⇒ dark matter candidate

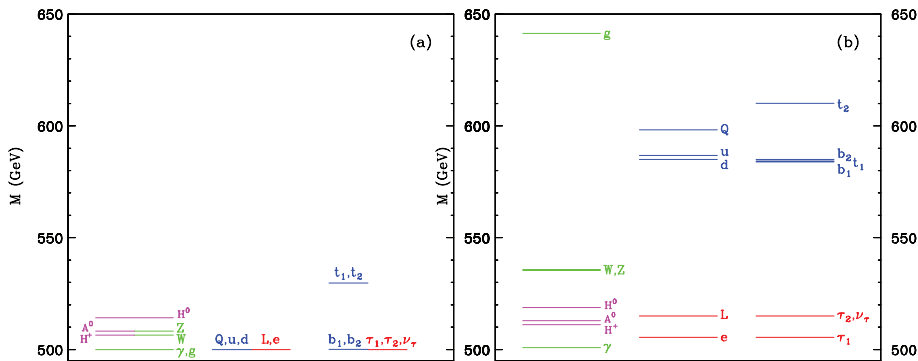
Servant-Tait '02

$$R^{-1} = 4 \text{ TeV}$$



# Universal extra dimensions (UED) : Mass spectrum

Radiative corrections  $\Rightarrow$  mass shifts that lift degeneracy at lowest KK level  
 divergent sum over KK modes in the loop  $\Rightarrow$  cutoff scale  $\Lambda \simeq 10/R$



# UED hadron collider phenomenology

- large rates for KK-quark and KK-gluon production
- cascade decays via KK- $W$  bosons and KK-leptons  
determine particle properties from different distributions
- missing energy from LKP: weakly interacting escaping detection
- phenomenology similar to supersymmetry

spin determination important for distinguishing SUSY and UED [40]

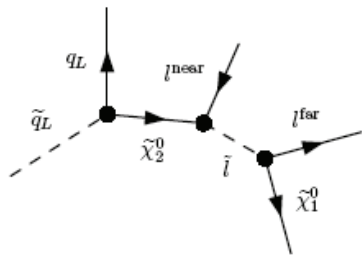
gluino	1/2	KK-gluon	1
squark	0	KK-quark	1/2
chargino	1/2	KK- $W$ boson	1
slepton	0	KK-lepton	1/2
neutralino	1/2	KK- $Z$ boson	1



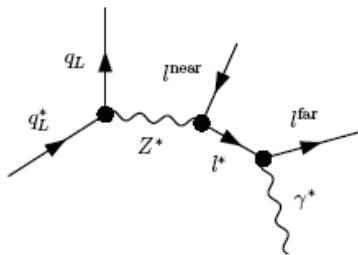
# SUSY vs UED signals at LHC

Example: jet dilepton final state

SUSY



UED



# Extra $U(1)$ 's and anomaly induced terms

masses suppressed by a loop factor

usually associated to known global symmetries of the SM

(anomalous or not) such as (combinations of)

Baryon and Lepton number, or PQ symmetry

Two kinds of massive  $U(1)$ 's:

I.A.-Kiritsis-Rizos '02

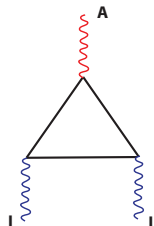
- 4d anomalous  $U(1)$ 's:  $M_A \simeq g_A M_s$

- 4d non-anomalous  $U(1)$ 's: (but masses related to 6d anomalies)

$$M_{NA} \simeq g_A M_s V_2 \leftarrow (6d \rightarrow 4d) \text{ internal space} \Rightarrow M_{NA} \geq M_A$$

or massless in the absence of such anomalies

# Green-Schwarz anomaly cancellation



$$= k_I^A \sim \text{Tr} Q_A Q_I^2 \rightarrow \text{axion } \theta : \delta A = d\Lambda \quad \delta\theta = -m_A \Lambda$$

$$-\frac{1}{4g_I^2} F_I^2 - \frac{1}{2} (d\theta + m_A A)^2 + \frac{\theta}{m_A} k_I^A \text{Tr} F_I \wedge F_I$$

cancel the anomaly

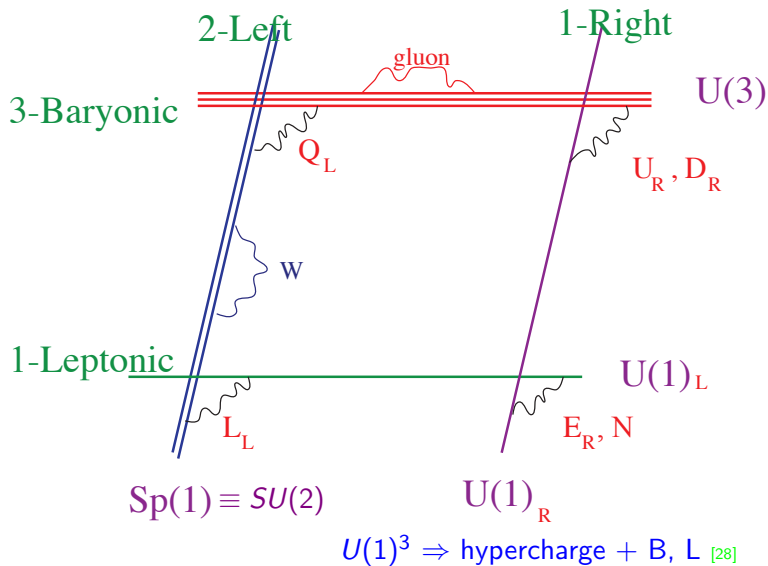
D-brane models:  $U(1)_A$  gauge boson acquires a mass

but global symmetry remains in perturbation theory

GS anomaly cancellation  $\Rightarrow$  extra scalars and axion-like particles (ALP)

- coupled to gauge kinetic terms
- lighter than the string scale (masses loop-factor suppressed)

# Standard Model on D-branes : SM<sup>++</sup>




- $B$  and  $L$  become massive due to anomalies

Green-Schwarz terms

- the global symmetries remain in perturbation

- Baryon number  $\Rightarrow$  proton stability

- Lepton number  $\Rightarrow$  protect small neutrino masses

no Lepton number  $\Rightarrow \frac{1}{M_s} LLHH \rightarrow$  Majorana mass:  $\frac{\langle H \rangle^2}{M_s} LL$   


- $B, L \Rightarrow$  extra  $Z$ 's

# Exotic $U(1)$ anomaly induced couplings

I.A.-Boyarsky-Ruchayskiy '06, '07

Non trivial anomaly cancellation  $\rightarrow$  new **dimensionless** couplings

mixed  $U(1)$  anomalies

$\Rightarrow Z'$  may couple to SM gauge bosons with no mass suppression

$$A \wedge X \wedge F_X \quad A \equiv Z', X \equiv W, Y$$

2 axionic phases:  $A \rightarrow \theta_A, X \rightarrow \theta_X \equiv$  SM Higgs  $\Rightarrow$

$$D\theta_A \wedge D\theta_X \wedge F_X \rightarrow \mathcal{L}_{\text{eff}} = c_1 D\theta_A \frac{H^\dagger DH}{|H|^2} F_Y + c_2 D\theta_A \frac{HF_W DH^\dagger}{|H|^2}$$

D'Hoker-Farhi type terms

$$c_2 \rightarrow AW^+W^- \quad c_1 \rightarrow AZY \quad (AZ\gamma, AZZ) \quad \text{vertices}$$

$$\Gamma(Z' \rightarrow ZZ) = \frac{c_1^2 \sin^2 \theta_W M_{Z'}^3}{192\pi M_Z^2} \left(1 - \frac{4M_Z^2}{M_{Z'}^2}\right)^{5/2}$$

$$\Gamma(Z' \rightarrow W^+W^-) = \frac{c_2^2 M_{Z'}^3}{48\pi M_W^2} \left(1 - \frac{4M_W^2}{M_{Z'}^2}\right)^{5/2}$$

$$\Gamma(Z' \rightarrow Z\gamma) = \frac{c_1^2 \cos^2 \theta_W M_{Z'}^3}{96\pi M_Z^2} \left(1 - \frac{M_Z^2}{M_{Z'}^2}\right)^3 \left(1 + \frac{M_Z^2}{M_{Z'}^2}\right)$$

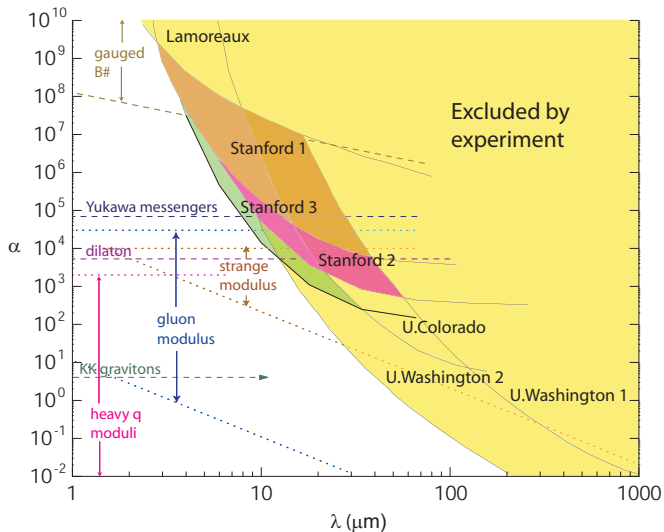
# microgravity experiments

- change of Newton's law at short distances
  - detectable only in the case of two large extra dimensions
- new short range forces
  - light scalars and gauge fields if SUSY in the bulk
    - or broken by the compactification on the brane
  - I.A.-Dimopoulos-Dvali '98, I.A.-Benakli-Maillard-Laugier '02
  - such as radion and lepton number
  - volume suppressed mass:  $(\text{TeV})^2/M_P \sim 10^{-4} \text{ eV} \rightarrow \text{mm range}$
  - can be experimentally tested for any number of extra dimensions
  - Light  $U(1)$  gauge bosons: no derivative couplings
    - $\Rightarrow$  for the same mass much stronger than gravity:  $\gtrsim 10^6$



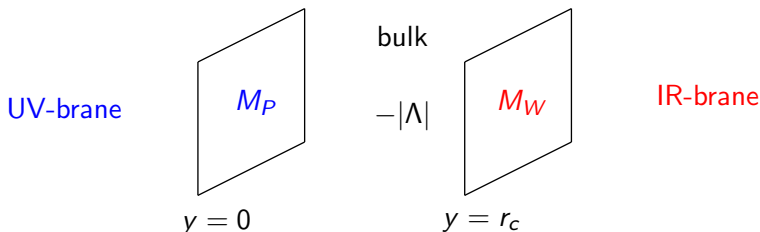
# Experimental limits on short distance forces

$$V(r) = -G \frac{m_1 m_2}{r} (1 + \alpha e^{-r/\lambda})$$



# Randal Sundrum models

spacetime = slice of  $\text{AdS}_5$  :  $ds^2 = e^{-2k|y|} \eta_{\mu\nu} dx^\mu dx^\nu + dy^2$   $k^2 \sim \Lambda/M_5^3$



- fine-tuned tensions:  $T = -T' = 24M_5^3 k$  [64]
- exponential hierarchy:  $M_W = M_P e^{-2kr_c}$   $M_P^2 \sim M_5^3/k$   
 $M_5 \sim M_{\text{GUT}}$
- 4d gravity localized on the UV-brane, but KK gravitons on the IR

- main prediction: spin-2 resonances at the TeV scale

$$m_n = c_n k e^{-2kr_c} \sim \text{TeV} \quad c_n \simeq (n + 1/4) \text{ for large } n$$

⇒ spin-2 TeV resonances in di-lepton or di-jet channels

- weakly coupled for  $m_n < M_5 e^{-2kr_c} \Rightarrow k < M_5$
- viable models: SM gauge bosons in the bulk, Higgs on the IR-brane
- AdS/CFT duals to strongly coupled 4d field theories

composite Higgs models, technicolor-type  $g_{YM} = M_5/k > 1$

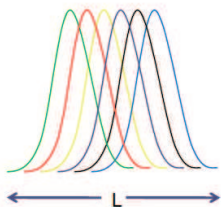
# More general framework: large number of species

$N$  particle species  $\Rightarrow$  lower quantum gravity scale :  $M_*^2 = M_p^2/N$

Dvali '07, Dvali, Redi, Brustein, Veneziano, Gomez, Lüst '07-'10

derivation from: black hole evaporation or quantum information storage

Pixel of size  $L$  containing  $N$  species storing information:



localization energy  $E \gtrsim N/L \rightarrow$

Schwarzschild radius  $R_s = N/(LM_p^2)$

no collapse to a black hole :  $L \gtrsim R_s \Rightarrow L \gtrsim \sqrt{N}/M_p = 1/M_*$

$M_* \simeq 1 \text{ TeV} \Rightarrow N \sim 10^{32}$  particle species !

# More general framework: large number of species

$N$  particle species  $\Rightarrow$  lower quantum gravity scale :  $M_*^2 = M_p^2/N$

Dvali '07, Dvali, Redi, Brustein, Veneziano, Gomez, Lüst '07-'10

derivation from: black hole evaporation or quantum information storage

$$M_* \simeq 1 \text{ TeV} \Rightarrow N \sim 10^{32} \text{ particle species !}$$

- Large extra dimensions      SM on D-branes

$N = R_{\perp}^n I_s^n$  : number of KK modes up to energies of order  $M_* \simeq M_s$

- $M_s \sim \text{TeV} \Rightarrow$  low inflation scale

allowed by the data since cosmological observables are dimensionless in units of the effective gravity scale

I.A.-Patil '14

# Cosmological observables

Power spectrum of temperature anisotropies

(adiabatic curvature perturbations  $\mathcal{R}$ )

$$\mathcal{P}_{\mathcal{R}} = \frac{H^2}{8\pi^2 M_*^2 \epsilon} \simeq \mathcal{A} \times 10^{-10} \quad ; \quad \mathcal{A} \approx 22$$

$\swarrow$   
 $-\dot{H}/H^2$

Power spectrum of primordial tensor anisotropies  $\mathcal{P}_t = 2 \frac{H^2}{\pi^2 M_*^2}$

$\Rightarrow$  tensor to scalar ratio  $r = \mathcal{P}_t / \mathcal{P}_{\mathcal{R}} = 16\epsilon$

measurement of  $\mathcal{A}$  and  $r \Rightarrow$  fix the scale of inflation

$$H \text{ in terms of } M_* \quad : \quad \frac{H}{M_*} = \left( \frac{\pi^2 \mathcal{A} r}{2 \times 10^{10}} \right)^{1/2} \equiv \Upsilon \approx 1.05 \sqrt{r} \times 10^{-4}$$

- $M_*$  may be different than  $M_{Planck}$  at the time of inflation

Explicit realisation:

Flat extra dimensions: obstruction due to the de Sitter bound:

$$M_{\text{spin } 2}^2 \geq 2H^2$$

Higuchi '87

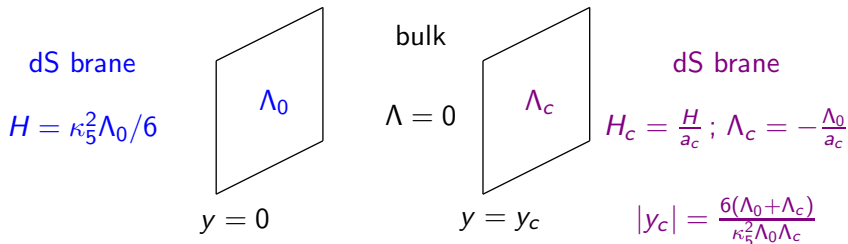
⇒ no KK-excitations with mass less than Hubble scale

Kleban-Mirbabayi-Porrati '15

5D brane-world realisation: empty bulk with two boundary dS branes [58]

$$ds^2 = \frac{(1 - H|y|)^2}{H^2\tau^2} (-d\tau^2 + dx_1^2 + dx_2^2 + dx_3^2) + dy^2$$

$|y| < 1/H$  : avoid Riddler horizon       $a(y) = 1 - H|y| > 0$





# spectrum and couplings

4d Planck mass:  $M_{Pl}^2 \sim 2/(H\kappa_5^2)$   $y_c$  large

Spectrum:

0-mode (4d graviton): wave function  $\phi_0 \sim (H/2)^{1/2} e^{-2z}$   $z \equiv -\ln a(y) > 0$

KK-modes:  $m_n^2 = H_y^2 \left( \frac{9}{4} + \pi^2 \frac{n^2}{z_c^2} \right)$   $H_y \equiv H/a(y)$ : Hubble constant at  $y$

wave functions  $\phi_n \sim \frac{H}{2m_n} \left( \frac{H}{z_c} \right)^{1/2} e^{-z/2} \left[ 3 \sin \left( \frac{n\pi z}{z_c} \right) - \frac{2\pi n}{z_c} \cos \left( \frac{n\pi z}{z_c} \right) \right]$

$\Rightarrow$  KK-modes couple much stronger than 0-mode at  $y_c$ :

$$|\phi_n|/|\phi_0| = \sqrt{2} \frac{\pi n}{z_c} \frac{H_c}{m_n} e^{3z_c/2}$$

similar result for bulk scalars

# Power spectrum

0-mode as before:

$$\mathcal{P}_0 = \frac{2}{\pi^2} \frac{H_c^2}{M_{Pl}^2}$$

KK-modes:

$$\mathcal{P}_n = \mathcal{P}_0 \frac{2\pi^2 n^2}{z_c^3} e^{3z_c} \left(\frac{H_c}{m_n}\right)^3 \left(\frac{k}{a_{dS} H_c}\right)^3 \simeq \mathcal{P}_0 \frac{2\pi^2 n^2}{z_c^3} \left(\frac{H_c}{m_n}\right)^3 e^{3(z_c - N)}$$

$N$ : number of e-foldings

Riotto '02

$$\mathcal{P}_0 \lesssim \mathcal{P}_n \Rightarrow e^N \lesssim e^{z_c}/z_c$$

satisfied for TeV scale inflation ( $N \gtrsim 35$ ,  $H_c \sim M_5 \sim \text{TeV}$ )

# Power spectrum

$$\mathcal{P}_{\text{KK}} = \sum_n^{m_n < M_5} \mathcal{P}_n = \mathcal{P}_0 e^{3(z_c - N)} \frac{2}{\pi} \ln \frac{M_5 z_c}{\pi H_c}$$
$$\gtrsim \mathcal{P}_0 \Rightarrow N \lesssim z_c$$

Allowed range of parameters:

$$M_{\text{Pl}}^2 \simeq \frac{M_5^3}{H_c} e^{z_c} \Rightarrow e^{z_c} \simeq \frac{M_{\text{Pl}}^2 H_c}{M_5^3} \lesssim \frac{M_{\text{Pl}}^2}{M_5^2} \quad H_c < M_5$$

$$e^{N_{\text{min}}} = 10^{13} \times \frac{H_c}{1 \text{ GeV}} \lesssim e^{z_c}$$

$$\Rightarrow 1 \text{ TeV} \lesssim H_c < M_5 \lesssim 10^8 \text{ GeV}$$

# Conclusions

String theory has many appealing properties:

- it provides a consistent quantization of gravity
- it gives a framework of unification of all interactions
- it inspired most of BSM new ideas
- it also inspired new results in mathematics
- it is a tool for strong coupling dynamics
- it has spectacular predictions if its scale is accessible to accelerators

It remains to be seen if it is a Theory of Nature