

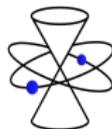
Renormalization in Tensorial Group Field Theories

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Open Problems in Theoretical Physics: the Issue of Quantum Space-Time

Corfu Summer Institute '15
Corfu, Greece



COST Action MP 1405
Quantum Structure of Spacetime

Outline

1 Introduction

- Quantizing Gravity: From Matrices to Tensors

2 Tensor Models & Graphs

- Invariant Tensor Models
- Degree of a tensor graph

3 Tensorial Group Field Theory and Renormalization

- Defining Tensor-like QFTs
- Overview of perturbatively renormalizable models
- The Functional Renormalization Program
- FRG equation for TGFT
- Application to a rank 3 model

4 Conclusion: Open Questions/Next Challenges

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Quantizing Gravity

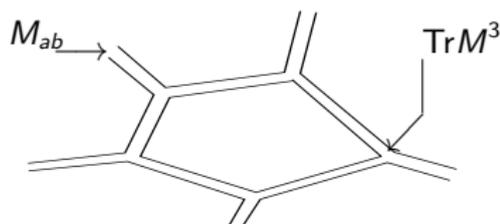
- “Ordinary” way \leadsto Several issues. The Einstein theory of gravity is perturbatively divergent and nonrenormalizable [DeWitt PR '67, Goroff & Sagnotti, NPB '86].
- Alternative scenarios (coupling gravity to other fields, Asymptotic Safety, ...) and more “daring” ones (extra-dimensions, susy, background field independent methods, ...).
- Mid 80's: *Matrix Models* [Di Francesco et al., PR '95] prove to be a solvable framework and concrete realization of an “emergent gravity” scenario.
- Matrix Models and Random 2D geometry: To “replace” the sum over topologies and geometries of 2D surfaces in a path integral, by a sum over random triangulations of surfaces.

Quantizing Gravity: From Matrices ...

- Probability measures for matrices M of large size N and describe 2D gravity. Archetype:

$$Z_{\text{matrix}} = \int dM e^{-\frac{1}{2} \text{Tr} M^2 + \frac{\lambda}{\sqrt{N}} \text{Tr} M^3}$$

• A triangulated surface \equiv A Feynman ribbon graph:



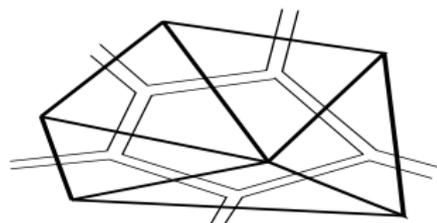
- 't Hooft's Large N limit: Planar graphs' sector \equiv surfaces of genus 0.
- Stat. Mech.: \exists phase transition ($N \rightarrow \infty; \lambda \rightarrow \lambda_c$) \rightsquigarrow a continuum limit (infinitely refined Riemann surfaces) as a 2D theory of gravity (Liouville gravity (gravity + CF matter)).

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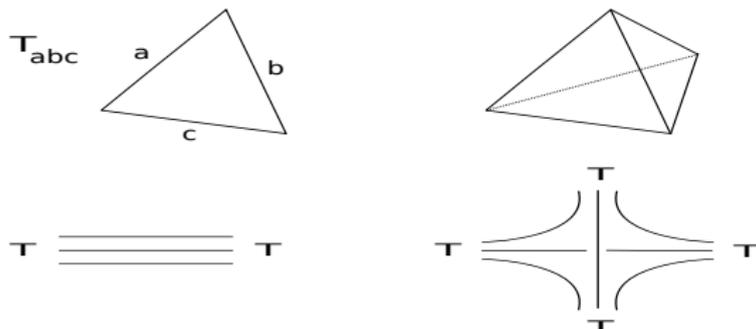
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... to Tensors ...

- Tensor Models (TM) of rank D : Tool for randomizing geometry in dimension D
- Basic building blocks $(D - 1)$ -simplexes & Interaction forms a D -simplex;
- For e.g. in 3D:

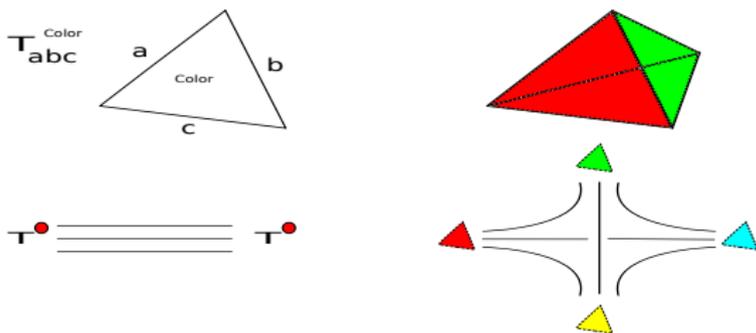


- Some results [Ambjorn et al. '91, Sasakura '91, Gross '92, Boulatov-Ooguri '92]
- Ambjorn et al., '91', phase transition (numerics).
- $1/N$ expansion missing \Rightarrow several exact results of MM difficult to extend to TM.
- TM need improvement(s): $C + DT$'s, Boulatov-Ooguri models, Group Field Theory [Oriti, '05-].

... and Colored Tensor Models

- '10 Gurau's $1/N$ expansion for colored TM [Gurau, AHP, '11]

3D:



- triangulate better objects (pseudo-manifolds) [Gurau, CMP '11]
- $\exists 1/N$ and Leading graphs triangulate only spheres in any D [Gurau, AHP '11; Bonzom, Gurau, Rivasseau '15]
- have computable phase transition, and critical exponent [Bonzom, Gurau, Riello, Rivasseau, NPB, '11];
[Benedetti, Bonzom, Carrozza, Dartois, Delpouve, Gurau, Lionni, Rivasseau, Ryan];
- Define, renormalizable field theories called TGFTs [BG & Rivasseau, '11; Carrozza, Oriti, Rivasseau '12; Samary & Vignes-Tourneret '12; BG, '13—; Avohou, Benedetti, Lahoche, Krajewski, Martini, Toriumi].

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Tensor Models [Gurau, '10, '11]

- Study of probability measures of random tensor spaces + Geometric/Topological/Physical inputs.

- A covariant complex tensor T_{p_1, \dots, p_d} with transformation rule

$$T_{p_1, \dots, p_d}^U = \sum_{q_k} U_{p_1 q_1}^{(1)} \dots U_{p_d q_d}^{(d)} T_{q_1, \dots, q_d}, \quad U^{(a)} \in U(N_a) \quad (1)$$

- G/T/Physics input: T is viewed as a $(d-1)$ -simplex.
- **Tensor Invariance** for defining the interactions

$$S_b^{\text{int}}(T, \bar{T}) = S_b^{\text{int}}(T^U, \bar{T}^U) = \text{Tr}_b(\bar{T} \cdot T \dots \bar{T} \cdot T) \quad (2)$$

b a colored graph encoding the contraction pattern; S_b^{int} "is" a gluing of simplexes and represents a d -simplex. Ex:

- Counting TM invariants is counting the number of branched covers of S^2 ! (JBG & S. Ramgoolam [AIHP D, '14])

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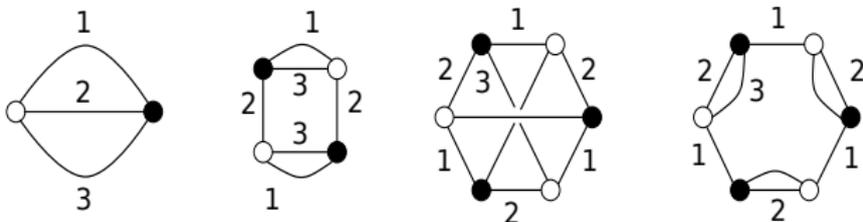
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Action/Partition function

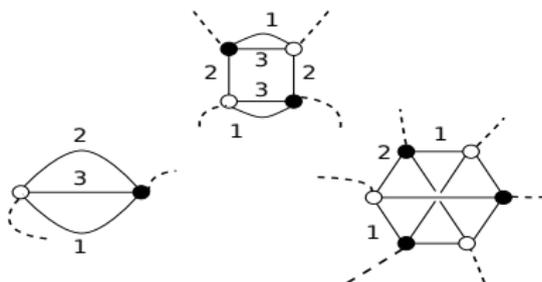
- Invariant action:

$$S[T, \bar{T}, \{\lambda_{\mathbf{b}}\}_{\mathbf{b}}] = \sum_{\mathbf{b}} \lambda_{\mathbf{b}} S_{\mathbf{b}}^{\text{int}}(T, \bar{T}), \quad (3)$$

- Partition function

$$Z[\{\lambda_{\mathbf{b}}\}_{\mathbf{b}}] = \int \prod_{p_i} [dT_{p_1, \dots, p_d} d\bar{T}_{p_1, \dots, p_d}] e^{-S[T, \bar{T}, \{\lambda_{\mathbf{b}}\}_{\mathbf{b}}]} \quad (4)$$

- Expanding the partition function: colored Feynman graphs



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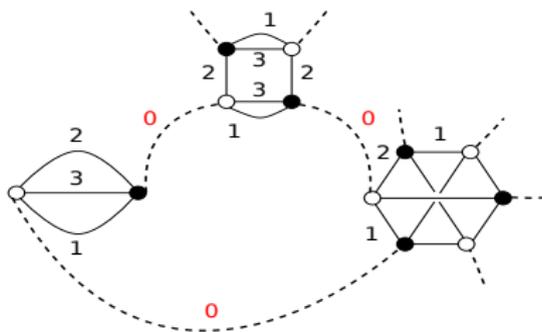
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Foundation questions:

- Tensor: Only a discretization tool or a “real” quanta of some “thing” ?

Big bang is identified with **geometrogenesis**, i.e. emergence of classical space-time.

Pre-space (analog of space-time before condensation) is treated as a **physical transplanckian early phase of the universe** (not just as a mathematical artefact) [Konopka '06, Oriti, '08].

- Divergences \Rightarrow Need of **Renormalization**. [Rivasseau: The Tensor Track '11, '12, '13]
Renormalization group is a guiding/selecting thread in theory space.

- Does it work for matrices ? The Grosse-Wulkenhaar model [Grosse-Wulkenhaar, '02-'03]: Renormalizable and Asymptotically Safe model (without any extra-symmetry required) and induces important NEW developments in Field Theory.

Defining Tensor-like QFTs

BG & Rivasseau, Commun. Math. Phys. '12; BG, Commun. Math. Phys. '14

- **Simple TM**: A complex tensor T_{p_1, \dots, p_d} with “Tensor Invariance” for defining the interactions and kinetic term

$$S^{\text{kin}}(T, \bar{T}) = \sum_{p_s} \bar{T}_{p_1, \dots, p_d} \left(\sum_s (p_s)^2 + \mu \right) T_{p_1, \dots, p_d} \quad (5)$$

$p_s^2 \equiv \Delta$. More generally: $(p_s)^{2a}$, $a > 0$.

- **Gauge invariant models (gi)** [Gurau, Bonzom, Carrozza, BG, Girelli, Livine, Lahoche, Oriti, Ousmane Samary, Rivasseau, Vignes-Tourneret] **Multi-orientable models**: [Tanasa, Dartois, Rivasseau]
- Summing over arbitrary high momenta may imply divergent amplitudes.

Results: Several Renormalizable Models and UV-Asymptotically Free [BG, 12'– '13]

Multi-scale Renormalizable TGFTs: $T : U(1)^d, SU(2)^d \rightarrow \mathbb{C}$

TGFT (type)	G_D	$\Phi^{k_{\max}}$	d	Renormalizability	UV behavior
	$U(1)$	Φ^6	4	Just-	AS (WiP with T. Koslowski)
	$U(1)$	Φ^4	5	Just-	AF
	$U(1)^2$	Φ^4	4	Just-	AF
	$U(1)$	Φ^{2k}	3	Super-	-
gi-	$U(1)$	Φ^4	6	Just-	AF
gi-	$U(1)$	Φ^6	5	Just-	-
gi-	$SU(2)^3$	Φ^6	3	Just-	AS ??
gi-	$U(1)$	Φ^{2k}	4	Super-	-
gi-	$U(1)$	Φ^4	5	Super-	-

Table: Updated list of tensorial renormalizable models and their features (AF \equiv asymptotically free).

Beyond perturbation: the Functional Method

- **A goal for TGFT:** Achieve a universal scenario for an emergent spacetime. One of the most robust mechanism for this is to see this emergence through one or several phase transitions, **running towards the IR**.
- But probing the IR, seeking and including nonperturbative effects & Deliver phase diagrams: FRG prominent tool.
- **Today's goal:** (1) Show that the FRG methods can be defined on TGFT models. (2) Applied on a rank 3 model, results suggest a phase transition **but difficult to ascertain**.

Today: **YES** to the 2 first questions !



- Already some results: [Krajewski & Toriumi, '14], [Benedetti, BG, Oriti JHEP '15], [Martini, Master Thesis '15], [BG, Martini, Oriti, '15], [Benedetti, Lahoche, '15], [BG, Koslowski, 1?].

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Wetterich-Morris formalism for TGFT

[Benedetti, BG, Oriti, JHEP '14]

- Rank d model: $\phi_{\mathbf{P}}$, \mathbf{P} a generic multi-index and an action

$$S = \sum_{\mathbf{P}, \mathbf{P}'} \frac{Z}{2} \phi_{\mathbf{P}} K(\mathbf{P}; \mathbf{P}') \phi_{\mathbf{P}'} + \sum_{\mathbf{b}} \lambda_{\mathbf{b}} \text{Tr}_{\mathbf{b}}(\mathcal{V} \phi^{\mathbf{n}_{\mathbf{b}}}). \quad (6)$$

- Cut-offed partition function $Z_{\Lambda}[J] = \int d\mu_{\Lambda}(\phi) e^{-S[\phi] + \text{Tr}_2(J \cdot \phi)} = e^{W_{\Lambda}[J]}$,
- One-parameter family ΔS_N of cut-off functions suppressing low energy modes up to N

$$\Delta S_N = \sum_{\mathbf{P}, \mathbf{P}'} \phi_{\mathbf{P}} \cdot R_N(\mathbf{P}; \mathbf{P}') \cdot \phi_{\mathbf{P}'},$$
$$\mathcal{Z}_N[J] = e^{W_N[J]} = \int_N d\phi e^{-S(\phi) - \Delta S_N(\phi) + \text{Tr}_2(J \cdot \phi)}. \quad (7)$$

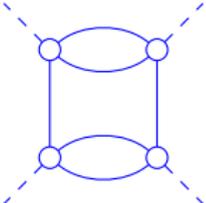
- Effective action & Legendre transform $\Gamma_N = \sup_J \left(\text{Tr}_2(J \cdot \varphi) - W_N(J) \right) - \Delta S_N(\varphi)$, $\varphi = \langle \phi \rangle$, leads to the Wetterich Equation

$$\partial_t \Gamma_N(\varphi) = \frac{1}{2} \overline{\text{Tr}} \left(\partial_t R_N \cdot [\Gamma_N^{(2)} + R_N]^{-1} \right) \quad (8)$$

A rank 3 tensor model

- $\phi : U(1)^3 \rightarrow \mathbb{R}$, with Fourier components $\phi_{\mathbf{P}}$, $\mathbf{P} = (p_i)$, $p_i \in \mathbb{Z}$. UV & IR refer to $U(1)$.
- Action (a theory of spheres)

$$S(\phi) = \frac{Z}{2} \text{Tr}_2(\phi \cdot K \cdot \phi) + \frac{m}{2} \text{Tr}_2(\phi^2) + S^{\text{int}},$$

$$K(p_i; p_i) = \frac{1}{d} \sum_{i=1}^d |p_i| \delta_{p_i, p'_i}, \quad S^{\text{int}} = \frac{\lambda_N}{4} \left[\begin{array}{c} \text{Diagram} \\ + \text{Sym} \end{array} \right] \quad (9)$$


- Truncation of the effective action:

$$\Gamma_N(\varphi) = \frac{Z_N}{2} \text{Tr}_2(\varphi \cdot K \cdot \varphi) + \frac{m_N}{2} \text{Tr}_2(\varphi^2) + S^{\text{int}},$$

$$R_N(p_i; p'_i) = \delta_{p_i, p'_i} Z_N \left(N - \sum_{i=1}^d |p_i| \right) \Theta \left(N - \sum_{i=1}^3 |p_i| \right) \quad (10)$$

β -functions of a rank 3 tensor model

$$t = \ln N, \quad \eta := \partial_t \ln Z_N, \quad m_N = Z_N \mu_N, \quad \lambda_N = Z_N^2 \bar{\lambda}_N$$

$$\eta = \frac{54\lambda_N N(2N+1)}{3(N+\mu_N)^2 - 2\lambda_N(5+18N+27N^2)} \quad (11)$$

$$\partial_t \mu_N = \frac{-\lambda_N N}{(N+\mu_N)^2} \left[\eta(18N^2 + 9N + 4) + (54N^2 + 36N + 9) \right] - \eta \mu_N \quad (12)$$

$$\partial_t \lambda_N = \frac{2\lambda_N^2 N}{(N+\mu_N)^3} \left[\frac{1}{3}\eta(18N^2 + 45N + 25) + 2(9N^2 + 18N + 9) \right] \quad (13)$$

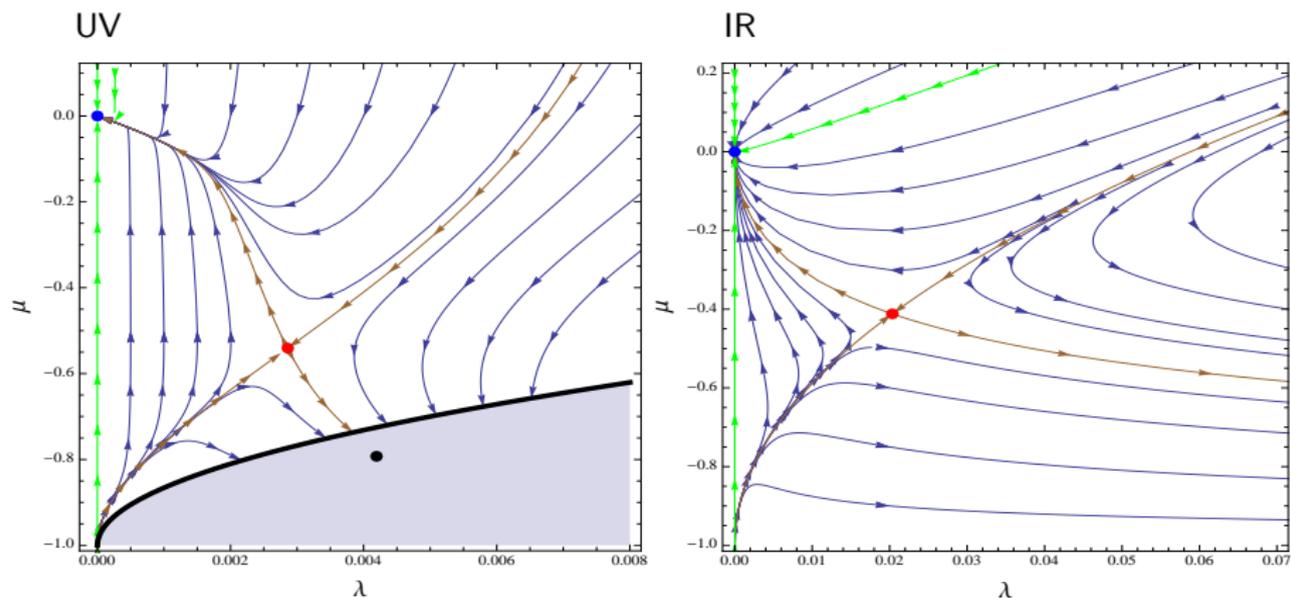
- Explicit dependency in N , a **non-autonomous system** ! A implicit scale (the radius of the group manifold) and the combinatorics of the interaction.

- What replaces the usual notion of **Canonical Dimension**?

- In the UV \implies needs a certain scaling: $\mu_N = N\tilde{\mu}_N$, $\lambda_N = \tilde{\lambda}_N$

- In the IR ($N=0$) \implies needs another scaling: $\eta \rightarrow 0$, $\mu_N = N\tilde{\mu}_N$, $\lambda_N = N^2\tilde{\lambda}_N$

Plots

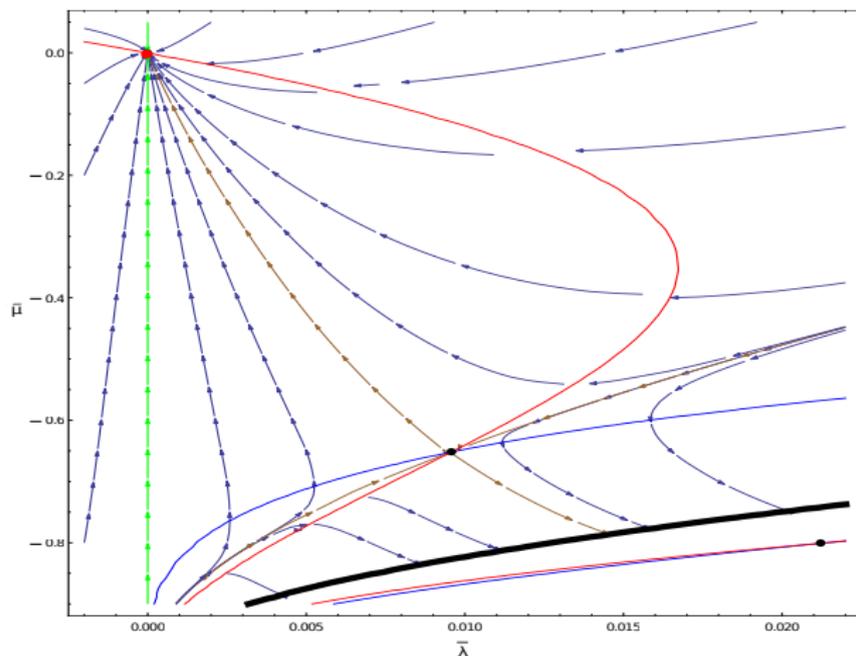


- The Gaussian FPs have a relevant and marginally relevant directions. Confirm the [UV asymptotic freedom](#) of the rank 3 model [BG and Samary, 2012].
- The N Gaussian FPs have one relevant and one irrelevant directions (similar to the Wilson-Fisher fixed point of scalar field theory in dimension 3, but the models are radically \neq).

Flash: TGFT on noncompact $\phi : \mathbb{R}^3 \rightarrow \mathbb{C}$

[BG, Martini, Oriti, 1508.01855] (1st RG analysis of TGFT in the noncompact case ☺)

- Issue of IR divergences can be treated using discretization of the background manifold and then a thermodynamic limit;



- Rank 3 Gauge invariant TGFT has a NGFP in the IR.

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Open questions/Next Challenges

- Tensor models generalize matrix models and Colored TMs is a rich and computable framework.
- Field Theory: There exist several models which are perturbatively renormalizable, some AF and some AS.
 - ~ Open questions • Understand better the theory space; • There exist certainly different types of TGFT than the one present here and Renormalization analysis works as well: How can we understand the theory space?
 - ~ Challenge: • Hunting nonperturbative and universal effects (FRG methods, recovering phase diagrams); prove that there is phase transition; investigate precisely the possible phases; order of parameters, symmetry breaking mechanism;
- Connection with Field Theory on non-associative geometrical spaces [wip with Ousmane Samary]

$$\begin{array}{ccc} [x^\mu, x^\nu] = i\Theta^{\mu\nu} & \longrightarrow & [x^\mu, x^\nu, x^\rho] = \Theta^{\mu,\nu,\rho} \\ \text{Noncommutative geometry} & \longrightarrow & \text{Nonassociative geometry} \\ \downarrow & & \downarrow \\ \text{Matrix/GW models} & \longrightarrow & \text{Tensorial field models} \end{array} \quad (14)$$