

Higgs bosons: discovered and hidden, in extended supersymmetric standard models at the LHC

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Outline

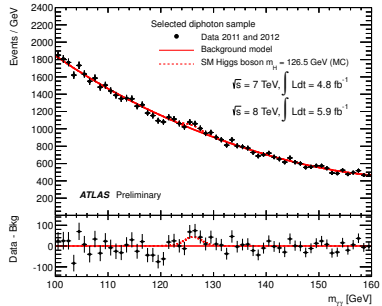
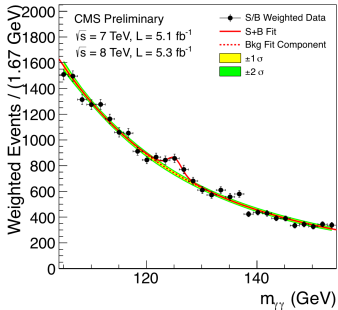
1 Plan

2 MSSM

3 TNSSM

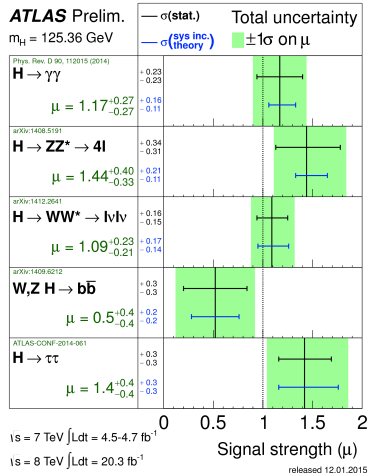
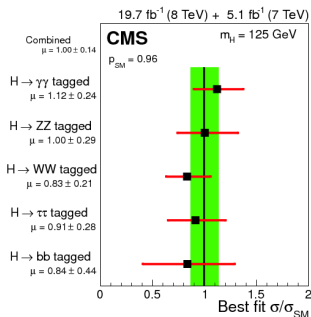
- 1 Observed Higgs boson
- 2 MSSM and its status
- 3 Triplet and singlet extended Supersymmetric scenarios
- 4 Possibility of Hidden scalar
- 5 Conclusion

Experimental constraints: LHC on 4th July, 2012



- Finally, on 4th of July 2012, we get evidence of a even-integer-spin particle similar to Higgs boson at the LHC.
- So far we have discovered Higgs boson around $\sim 125 \text{ GeV}$ in $WW^*, ZZ^*, \gamma\gamma$ modes.

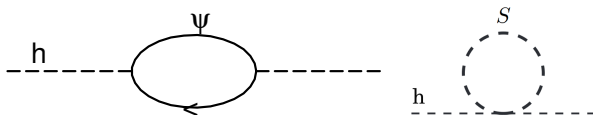
Status of Higgs at the LHC



- CMS: $H \rightarrow ZZ \Rightarrow m_H = 125.5 \pm 0.2(\text{stat})^{+0.5}_{-0.4}(\text{sys})$
- ATLAS: $H \rightarrow ZZ \Rightarrow m_H = 125.7 \pm 0.3(\text{stat}) \pm 0.3(\text{sys})$

Hierarchy problem and supersymmetry

- Higgs mass is not protected by any symmetry \Rightarrow **Hierarchy problem**.
- **Supersymmetry** protects the Higgs mass by giving possible cancellation.



- For each particle there is a super partner differing by spin $1/2$.

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Higgs sector in MSSM

- Unlike Standard Model, MSSM has five Higgs bosons,
 h, H the CP-even neutral Higgs bosons
 A the CP-odd neutral Higgs bosons
 H^\pm charged Higgs bosons
- Including R -parity, $P_R = (-1)^{3(B-L)+2s}$ conservation leads to the lightest supersymmetric particle as **Dark matter candidate**.

Lightest Higgs in MSSM

- While in the SM the Higgs mass is essentially a free parameter the lightest CP-even Higgs particle in the MSSM is bounded from above.
- At tree level the lightest CP-even Higgs (h) mass, $m_h \leq M_Z$
- For desired Higgs mass around 125 GeV, one has to look for quantum correction.

Higgs sector in MSSM

Lightest Higgs in MSSM

- Depending on the SUSY parameters that enter the radiative corrections, it is restricted to values

$$\begin{aligned} M_h^{\max} &\approx M_Z |\cos 2\beta| + \text{radiative corrections} \\ &< 110 - 135 \text{ GeV} \\ &\sim \end{aligned}$$

- SUSY particles mass and other parameters enter in the radiative correction to the Higgs sector.

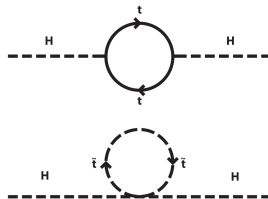
- ⇒ “Observed $M_h \approx 125 \text{ GeV}$ ” at the LHC, would place very strong constraints on the MSSM parameters through their contributions to the radiative corrections to the Higgs sector.

- With the radiative correction at one-loop the lightest Higgs mass becomes,

$$m_h^2 \underset{\sim}{<} m_Z^2 \cos^2 2\beta + \frac{3}{4\pi^2} \frac{m_t^4}{v^2} \left[\ln \frac{M_S^2}{m_t^2} + \frac{X_t^2}{M_S^2} \left(1 - \frac{X_t^2}{12M_S^2} \right) \right]$$

where $M_S = \sqrt{m_{\tilde{t}_1} m_{\tilde{t}_2}}$,

the stop mixing parameter is $X_t = A_t - \mu \cot \beta$
and for maximal mixing scenario: $X_t^{\max} = \sqrt{6} M_S$.



$$M_{\tilde{t}}^2 = \begin{pmatrix} M_L^2 & m_t X_t \\ m_t X_t & M_R^2 \end{pmatrix}.$$

Higgs sector in MSSM

Lightest Higgs in MSSM

- Lightest CP-even Higgs boson is the candidate ~ 125 GeV Higgs boson.
- In MSSM, the main loop contributions come from the third generation strong sectors,
- push the mass scale $\gtrsim 1$ TeV to up to 10 TeV depending on the models.
- In most constrained SUSY scenarios viz mSUGRA, the required SUSY mass scale is $\gtrsim 10$ TeV.
- In pMSSM, \lesssim TeV SUSY scale still a possibility but with large mixing in the stop mass matrix.
- Fine tuning is necessary for MSSM to have lightest Higgs boson mass around ~ 125 GeV.

► more

Heavy Higgs bosons

- The other CP-even and CP-odd Higgs search is still on.
- So far both CMS and ATLAS have not found any evidence of heavier Higgs bosons.

Charged Higgs bosons

- MSSM has one charged Higgs boson.
- The charged Higgs dominantly decays to $\tau\nu$
- So far no evidence has been found.

~ 125 GeV Higgs \Rightarrow very large SUSY mass scale or/and large mixings \Rightarrow fine tuning.

Extended Higgs sector!

- New scalars give additional contributions to the lightest Higgs mass
- so no large mixing and/or heavy sfermions needed
- Extended theory may solve the μ problem
- Possibility of spontaneous CP-violation.

Options ?

Triplet and/or singlet extensions.

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3 **TNSSM**

Triplet-Singlet extended Supersymmetric scenarios

- The μ_D term in the superpotential can be generated spontaneously by the introduction of a gauge singlet superfield S .
- The scale invariant superpotential with the $Y = 0$ triplet T and singlet S is given by,

$$W_S = \lambda_T H_d \cdot T H_u + \lambda_S S H_d \cdot H_u + \lambda_{TS} S \text{Tr}[T^2] + \frac{\kappa}{3} S^3.$$

- The complete Lagrangian with the soft SUSY breaking terms has an accidental Z_3 symmetry; $\phi_i \rightarrow e^{2\pi i/3} \phi_i$ [▶ more](#)
- When the neutral parts get vev,

$$\langle H_{u,d}^0 \rangle = \frac{v_{u,d}}{\sqrt{2}}, \quad \langle S \rangle = \frac{v_S}{\sqrt{2}} \quad \langle T^0 \rangle = \frac{v_T}{\sqrt{2}}$$

generate the the doublet mixing term $\mu_D = \frac{\lambda_S}{\sqrt{2}} v_S + \frac{\lambda_T}{2} v_T$

- Singlet does not couple to fermions and gauge bosons, whereas $Y = 0$ triplet only couples to W^\pm
- Experimental ρ parameter gives strong bounds, $v_T \leq 5 \text{ GeV}$.

TNSSM: Higgs bosons and mass spectrum at tree-level

- Enhanced Higgs particle spectrum

CP – even	CP – odd	charged
h_1, h_2, h_3, h_4	a_1, a_2, a_3	$h_1^\pm, h_2^\pm, h_3^\pm$

- The extra interaction terms also enhances the Higgs mass contribution at the tree-level thus reduces required SUSY the fine-tuning further.

$$m_{h_1}^2 \leq m_Z^2 (\cos^2 2\beta + \frac{\lambda_T^2}{g_L^2 + g_Y^2} \sin^2 2\beta + \frac{2\lambda_S^2}{g_L^2 + g_Y^2} \sin^2 2\beta),$$
$$\tan \beta = \frac{v_u}{v_d}$$

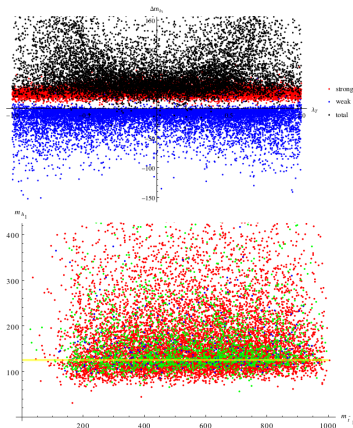
TNSSM: Higgs mass spectrum at one-loop

Effective potential at one loop

- For one-loop neutral Higgs mass spectrum we follow the effective-potential approach by [Coleman-Weinberg](#)
- The one-loop radiative corrections to the Higgs potential can be calculated using the effective potential approach

$$\Delta V = \frac{1}{64\pi^2} \text{Str} \left[\mathcal{M}^4 \left(\ln \frac{\mathcal{M}^2}{\Lambda^2} - \frac{3}{2} \right) \right].$$

where \mathcal{M} represents the field dependent mass matrices of the particles and Λ is the renormalization scale.



- The required stop mass is much lower and one should look for directed SUSY searches.

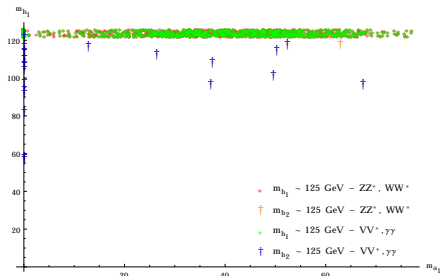
TNMSSM: Light pseudo-scalar

In the limit of a continuous symmetry

- In the limit when the A_i parameters go to zero, the discrete Z_3 symmetry of the Lagrangian is promoted to a continuous $U(1)$ symmetry

$$(\hat{H}_u, \hat{H}_d, \hat{T}, \hat{S}) \rightarrow e^{i\phi}(\hat{H}_u, \hat{H}_d, \hat{T}, \hat{S})$$

- If this symmetry is softly broken by very small parameters A_i of $\mathcal{O}(1)$ GeV, we get a very light pseudoscalar as pseudo-Nambu-Goldstone boson of the symmetry.



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Hidden Higgs phenomenology

- The possibility of very light hidden Higgs(es) is still viable in this scenario.
- Specially $h^\pm \rightarrow h_1 W^\pm$, $h_i \rightarrow a_j a_k$, $h_i \rightarrow Z a_j$ are very crucial in determining the Higgs sector.
- We are investigating $2b + 2\tau$, $2\tau + 2\mu$ and $2b + 2\mu$ final states to probe the hidden Higgs scenario taking care of the possible backgrounds.

Charged Higgs phenomenology

- Triplet type charged Higgs does not couple to fermions.
- Tree-level $h_i^\pm - W^\mp - Z$ introduces new decay channel and production channels.

Conclusions

- So far we have observed a CP-even Higgs boson around ~ 125 GeV.
- All possible Standard Model modes are yet to be discovered.
- Possibility of hidden Higgs is still a possibility.
- Observation of Charged Higgs would be a direct proof of extended Higgs sector.
- We have studied the triplet and triplet-singlet extended supersymmetric models in the context of 125 GeV Higgs.
- In the extended Higgs SUSY scenarios, lighter SUSY mass scale is still a possibility.
- $h_1^\pm \rightarrow ZW^\pm$ could be a smoking gun signature for the triplet scenario.
- A very light pseudo-scalar is still a possibility with extended Higgs sector.

Thank you

TNMSSM potential

$$\begin{aligned}
V_{soft} = & m_{H_u}^2 |H_u|^2 + m_{H_d}^2 |H_d|^2 + m_S^2 |S|^2 \\
& + m_T^2 |T|^2 + m_Q^2 |Q|^2 + m_U^2 |U|^2 + m_D^2 |D|^2 \\
& + (A_S S H_d \cdot H_u + A_\kappa S^3 + A_T H_d \cdot T \cdot H_u + A_{TS} S \text{Tr}(T^2) \\
& + A_U U H_U \cdot Q + A_D D H_D \cdot Q + h.c.),
\end{aligned}$$

while the D-terms are given by

$$V_D = \frac{1}{2} \sum_k g_k^2 (\phi_i^\dagger t_{ij}^a \phi_j)^2.$$

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