

Corfu', September '14

**The HIGGS  
and the  
EXCESSIVE success  
of the SM**

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**Roma Tre/CERN**

# LHC 7-8 TeV

A great triumph: the 126 GeV Higgs discovery



A particle apparently **just** as predicted by the SM theory

The main missing block for the experimental validation of the SM is now in place

A negative surprise: no production of new particles,  
no evidence of new physics which was expected  
on theoretical grounds

Not in ATLAS&CMS

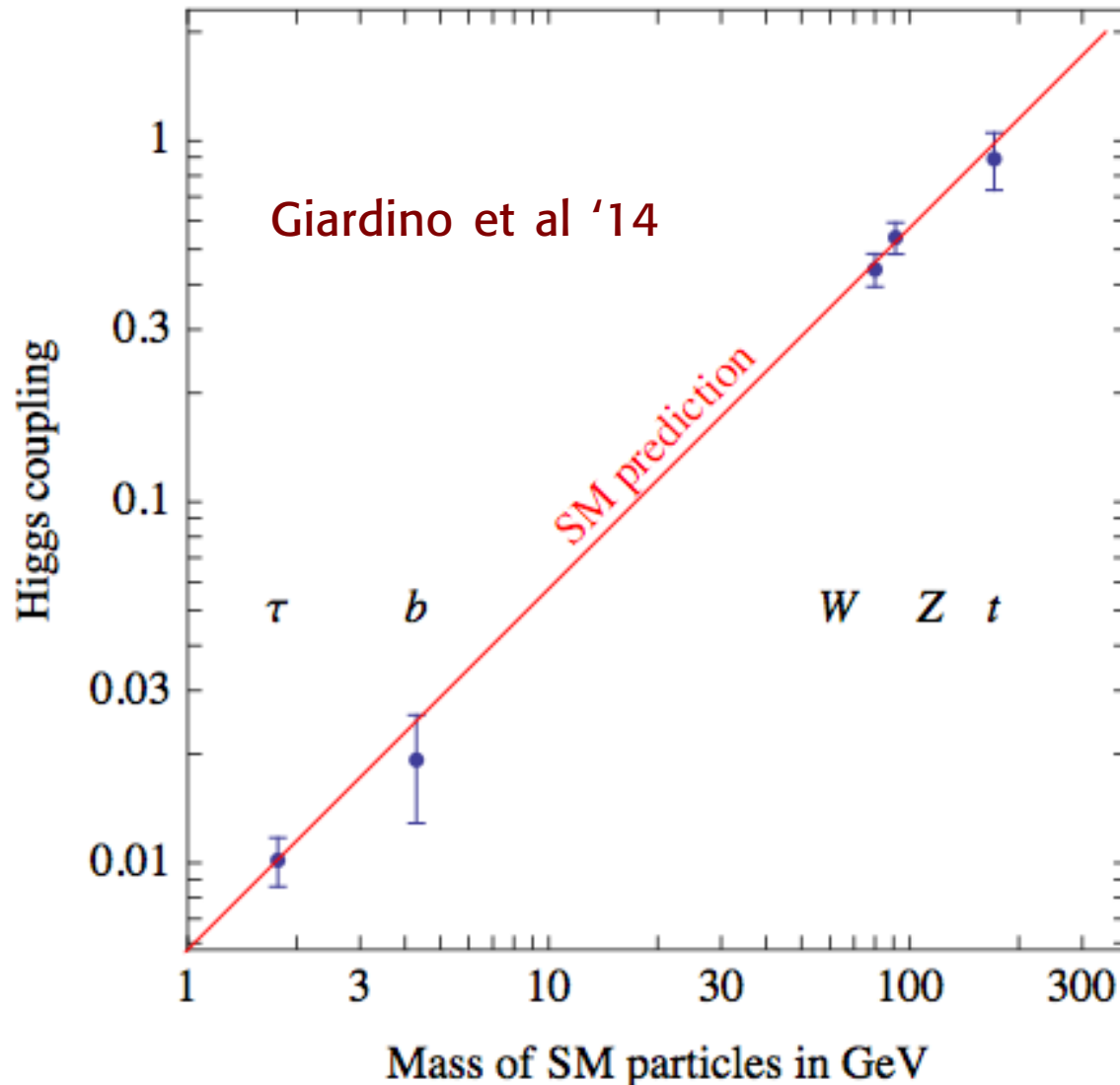
Not in Heavy Flavour decays (LHCb, ..... B-factories)

Not in  $\mu \rightarrow e \gamma$  (MEG)  $\mathcal{B} < 5.7 \times 10^{-13}$

Not in the EDM of the electron (ACME)  $|d_e| < 8.7 \times 10^{-29} e \text{ cm}$

⊕ ..... [Perhaps a deviation in  $(g-2)_\mu$ ?]

The Higgs couplings are in proportion to masses: a striking signature [plus specified,  $gg$ ,  $\gamma\gamma$ ,  $Z\gamma$  eff. couplings]



Nearly impossible to reproduce by accident

Agrees with a SM doublet: no Clebsch or mixing distortions detected

The spin-parity  $0^+$  also looks OK

It really appears as the SM Higgs particle!!!

The precise measurements of Higgs couplings are crucial to determine to what extent it is SM

It would really be astonishing if no deviation from the SM is seen!

General effective lagrangians are being studied by adding higher dim ops or introducing eff. couplings

eg Alonso et al  
Giudice et al  
Csaki et al  
Contino  
Keren-zur et al  
Falkowski et al  
Elias-Miro et al  
Pomarol, Riva.....

Almeida et al

$$\begin{aligned} \mathcal{L} = & \kappa_3 \frac{m_H^2}{2v} H^3 + \kappa_Z \frac{m_Z^2}{v} Z_\mu Z^\mu H + \kappa_W \frac{2m_W^2}{v} W_\mu^+ W^{-\mu} H \\ & + \kappa_g \frac{\alpha_s}{12\pi v} G_{\mu\nu}^a G^{a\mu\nu} H + \kappa_\gamma \frac{\alpha}{2\pi v} A_{\mu\nu} A^{\mu\nu} H + \kappa_{Z\gamma} \frac{\alpha}{\pi v} A_{\mu\nu} Z^{\mu\nu} H \\ & + \kappa_{VV} \frac{\alpha}{2\pi v} (\cos^2 \theta_W Z_{\mu\nu} Z^{\mu\nu} + 2 W_{\mu\nu}^+ W^{-\mu\nu}) H \\ & - \left( \kappa_t \sum_{f=u,c,t} \frac{m_f}{v} f \bar{f} + \kappa_b \sum_{f=d,s,b} \frac{m_f}{v} f \bar{f} + \kappa_\tau \sum_{f=e,\mu,\tau} \frac{m_f}{v} f \bar{f} \right) H. \end{aligned}$$

19.03.14

Rates:  $\mu = \sigma \text{Br} / (\sigma \text{Br})^{\text{SM}}$

E. Gross

Combined fit

Combined fit

ATLAS

$[m_H = 125.5 \text{ GeV}]$

CMS

$[m_H = 125.7 \text{ GeV}]$

$\mu_{\gamma\gamma} = 1.57^{+0.33}_{-0.28}$

$\mu_{ZZ} = 1.44^{+0.40}_{-0.35}$

$\mu_{WW} = 1.00^{+0.32}_{-0.29}$

$\mu_{\tau\tau} = 1.4^{+0.5}_{-0.4}$

$\mu_{bb} = 0.2^{+0.7}_{-0.6}$

$\mu = 1.30^{+0.18}_{-0.17}$

$\mu_{\gamma\gamma} = 0.77^{+0.27}_{-0.27}$

$\mu_{ZZ} = 0.92^{+0.28}_{-0.28}$

$\mu_{WW} = 0.68^{+0.2}_{-0.2}$

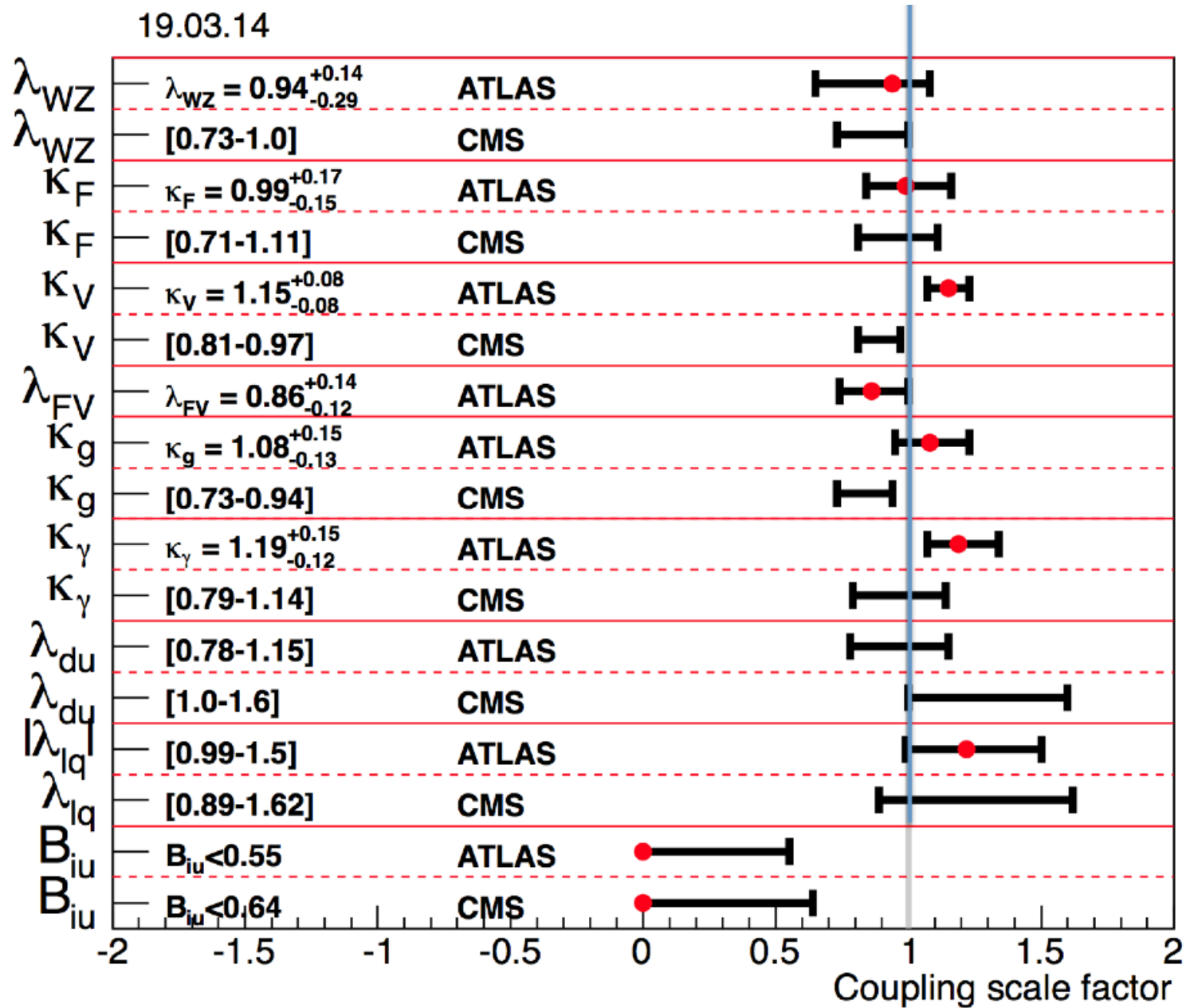
$\mu_{\tau\tau} = 1.10^{+0.41}_{-0.41}$

$\mu_{bb} = 1.15^{+0.62}_{-0.62}$

$\mu = 0.80^{+0.14}_{-0.14}$



It really looks like the SM Higgs!



For example:

MSSM: separate u and d couplings and  $|a=hVV| < 1$

Tree level formulae  
Radiative corrections  
important

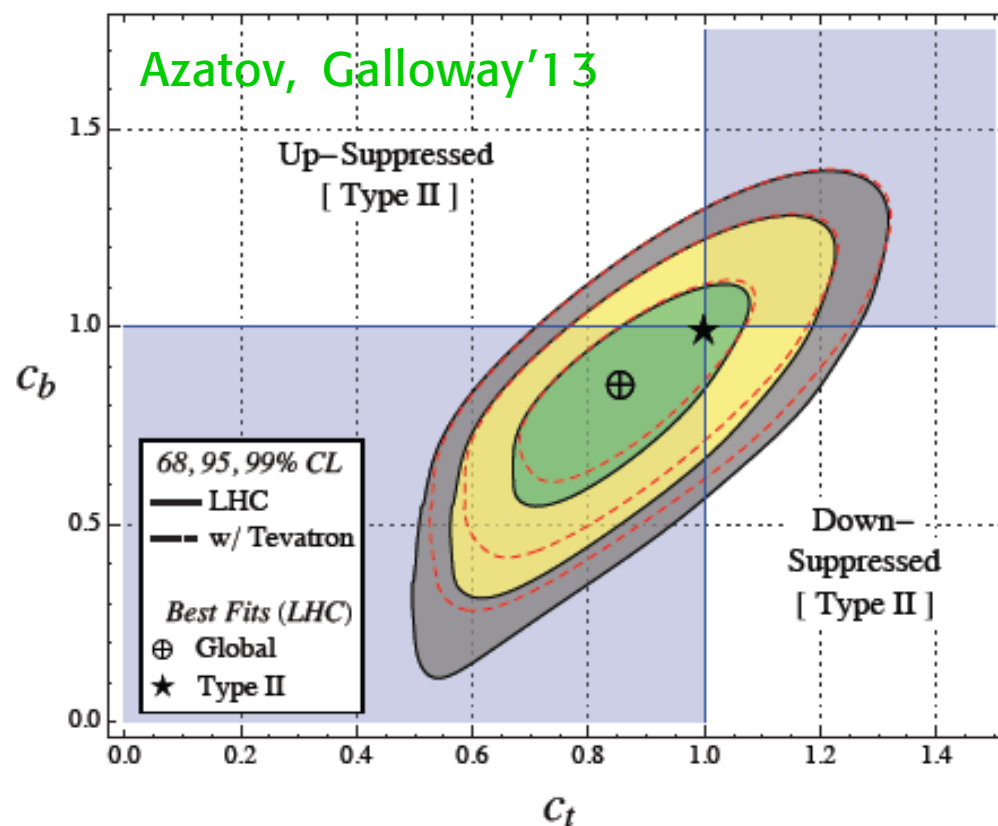
$$a = hVV = \sin(\beta - \alpha)$$

$$c_u = huu = \frac{\cos \alpha}{\sin \beta}$$

$$c_d = hdd = -\frac{\sin \alpha}{\cos \beta}$$

If  $c_u > 1$  then  $c_d < 1$   
and viceversa

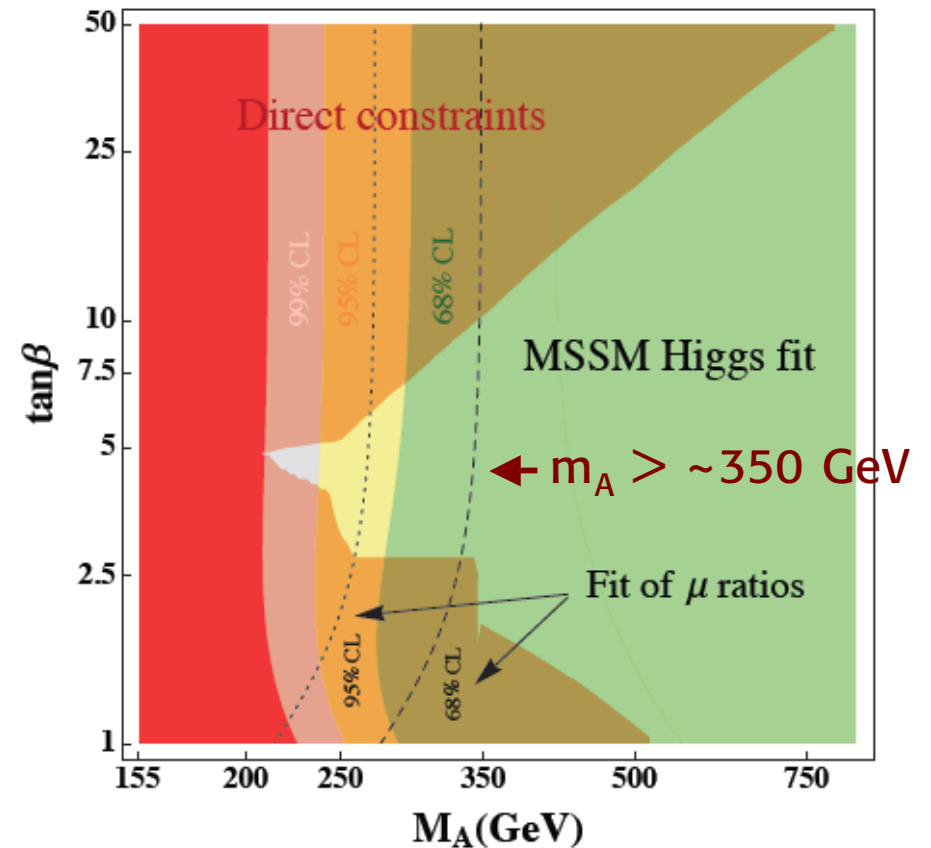
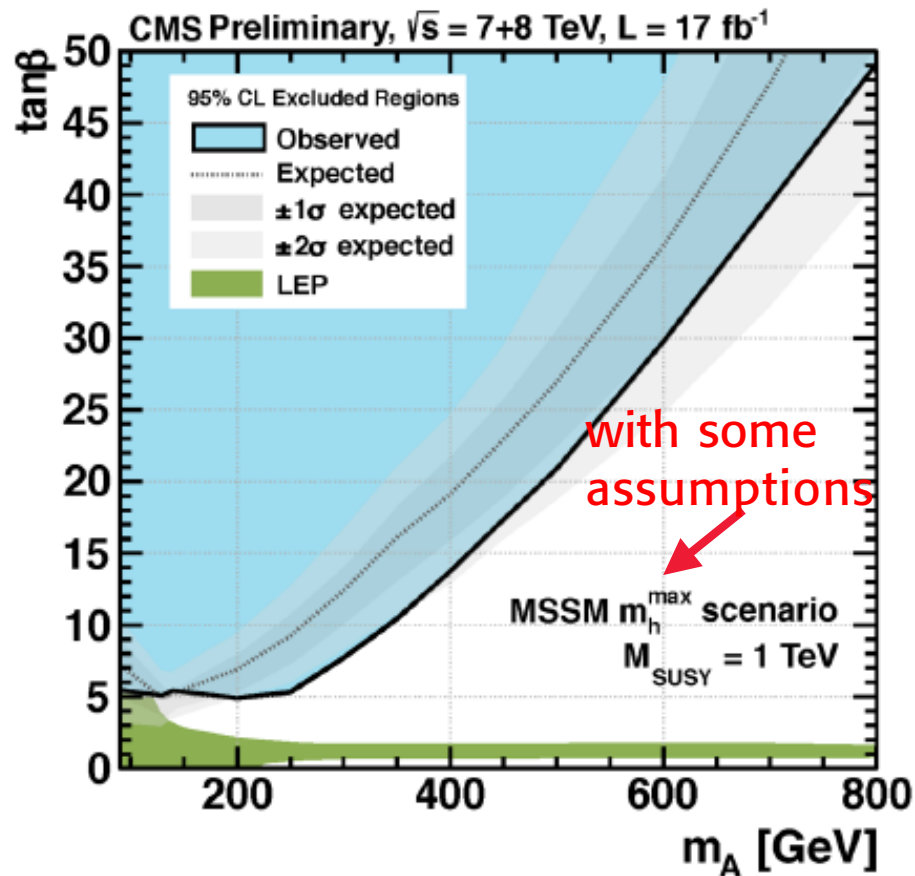
$$\tan 2\alpha = \tan 2\beta \frac{m_A^2 - m_Z^2}{m_A^2 + m_Z^2}$$



A very important open question:  
Are there more Higgs particles?  
Here we focus on MSSM

## Theorists analysis

Djouadi, Maiani et al '13



this limit has now been improved



$\tau\tau$



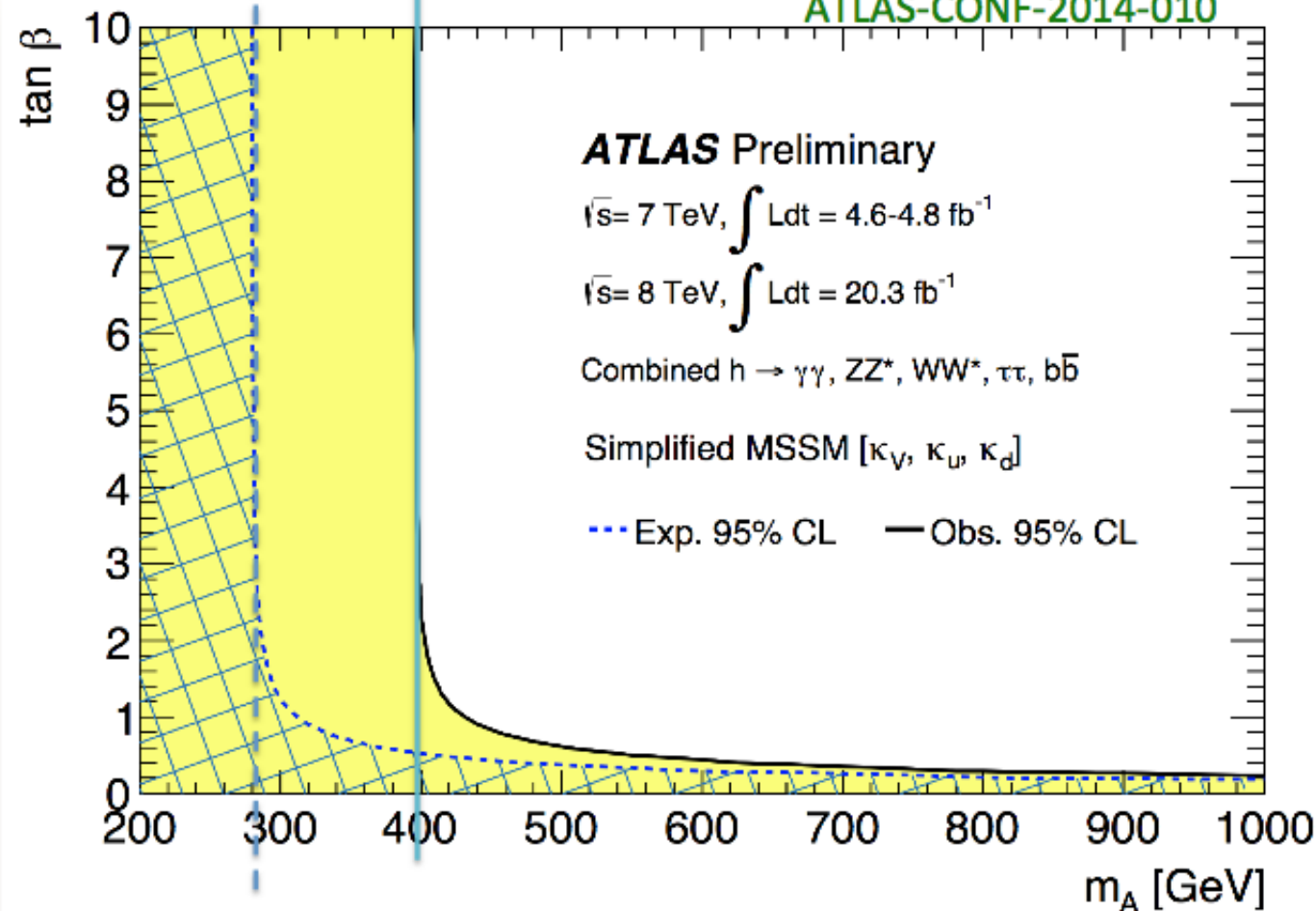


Now a better limit is obtained

E. Gross, Moriond EW '14

$$m_A > 400 \text{ GeV for } \tan\beta > 2$$

ATLAS-CONF-2014-010



Bottom line:

The issue of extra Higgs (doublets and/or singlets) is a clear priority



# Impact of the Higgs discovery

The minimal SM Higgs:  
is the simplest possible form of spont. EW symmetry breaking.

What was considered by many theorists just as a toy model,  
a temporary addendum to the gauge part of the SM,  
is now promoted to the real thing!

The only known example in physics of a fundamental,  
weakly coupled, scalar particle with VEV

→ e.g. the quartic coupling is perturbative:

$$V = -\mu^2 \phi^\dagger \phi + \frac{1}{2} \lambda (\phi^\dagger \phi)^2 \quad \phi \rightarrow v + \frac{H}{\sqrt{2}} \quad v = 174.1 \text{ GeV}$$

$$\oplus \quad m_H^2 = 2\mu^2 = 2\lambda v^2 \quad \longrightarrow \quad \frac{1}{2} \lambda \sim 0.13$$

# Higgs, unitarity and naturalness in the SM

In the SM the Higgs provides a solution to the occurrence of unitarity violations in some amplitudes ( $W_L, Z_L$  scattering)

To avoid these violations one needed either one or more Higgs particles or some new states (e.g. new vector bosons)

Something had to happen at the few TeV scale!!

While this is a theorem, once there is the Higgs, the necessity of new physics on the basis of naturalness is not a theorem but still a well motivated demand

The absence of accompanying new physics puts the issue of the relevance of **our concept of naturalness** at the forefront



# The naturalness principle

Has been and is the main motivation for new physics at the weak scale

But at present our confidence on naturalness as a guiding principle is being more and more challenged

No indirect evidence of new physics (is  $g-2$  really solid?)  
No direct evidence of new physics at the LHC

Manifestly some amount of fine tuning is imposed on us by the data. More so now after the LHC7-8 results

Does Nature really care about our concept of Naturalness? Apparently not much!

Which form of Naturalness is Natural?



The argument for naturalness is really strong...  
**except that** it has failed so far as a guiding principle

As a consequence:

We can no more be sure that within 3 or 10 or 100 TeV.....  
the solution of the hierarchy problem must be found  
--> implications for future Colliders

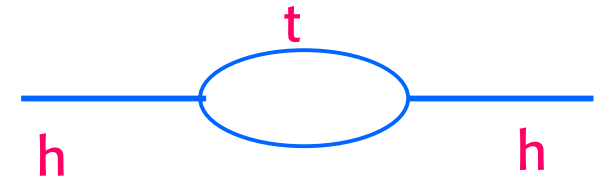
Moreover, it is true that the SM theory is renormalizable  
and completely finite and predictive

If you forget the required miraculous fine tuning  
you are not punished, you find no catastrophe!!



The naturalness argument for new physics at the EW scale is often expressed in terms of the quadratic cut-off dependence in the scalar sector

$$\delta m_h^2|_{top} = -\frac{3G_F}{2\sqrt{2}\pi} m_t^2 \Lambda^2 \sim -(0.2\Lambda)^2$$

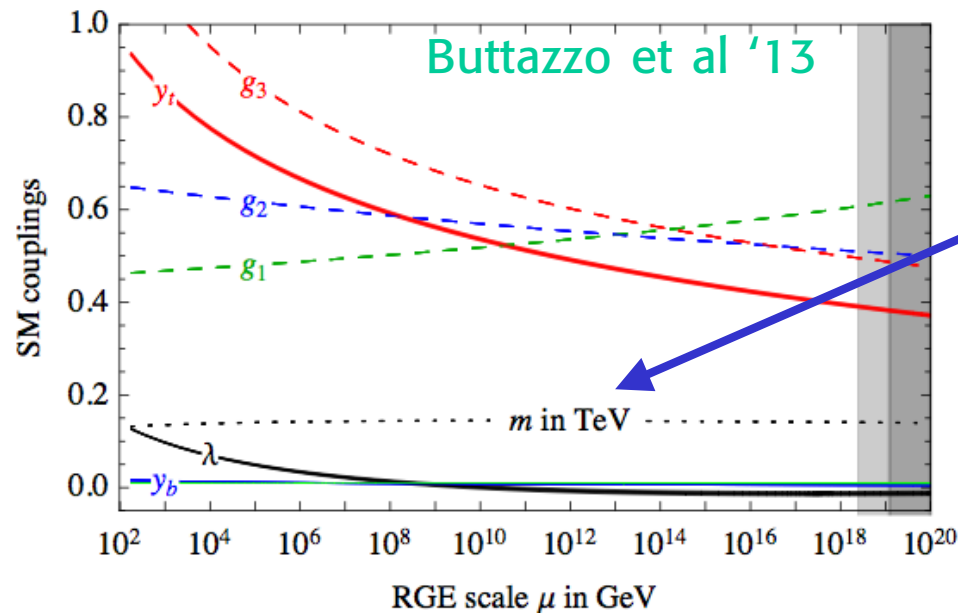


If we see the cut-off  $\Lambda$  as the scale where new physics occurs that solves the fine tuning problem, then the new physics must be nearby

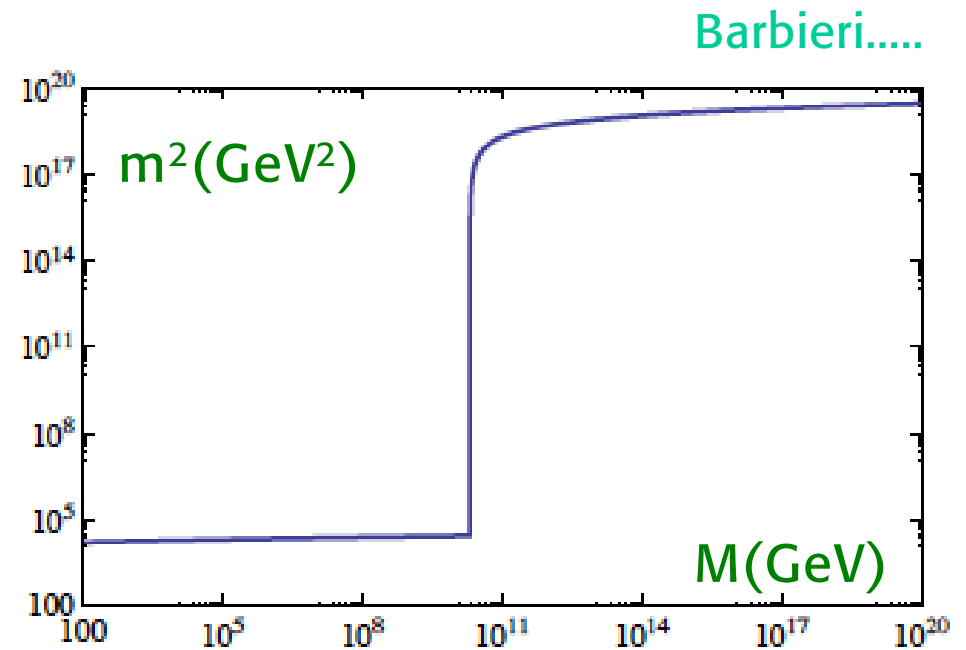
The argument can be formulated in terms of renormalized quantities with no reference to a cut-off --->  
quadratic sensitivity to thresholds at high energy



# Naturalness in a more physical language



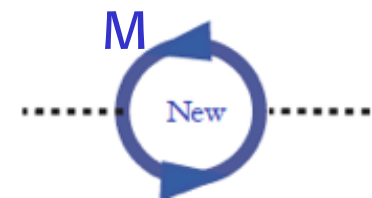
In the renormalized theory the running Higgs mass slowly evolves logarithmically



But in the presence of a threshold at  $M$  for a heavy particle coupled to the Higgs, the quadratic sensitivity produces a jump in the running mass

$$M \sim 10^{10} \text{ GeV}, \lambda_H \sim 1, \text{ jump: } m^2 \sim (\lambda_H M)^2 / (16\pi^2)$$

- Fine tuning is then needed to explain the small value of  $m$  at low energy

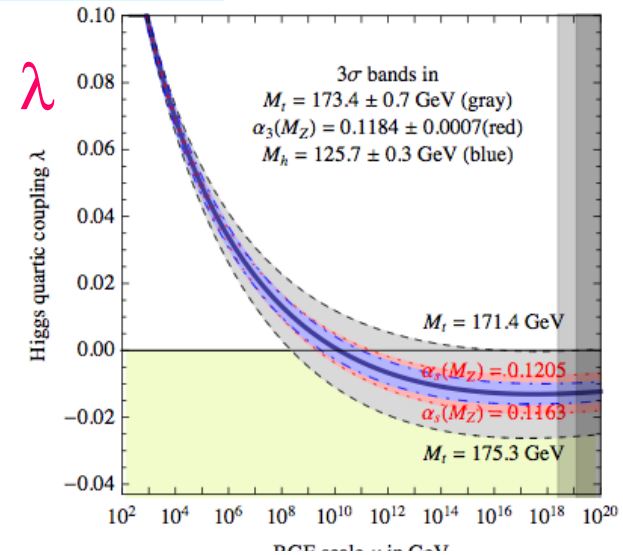
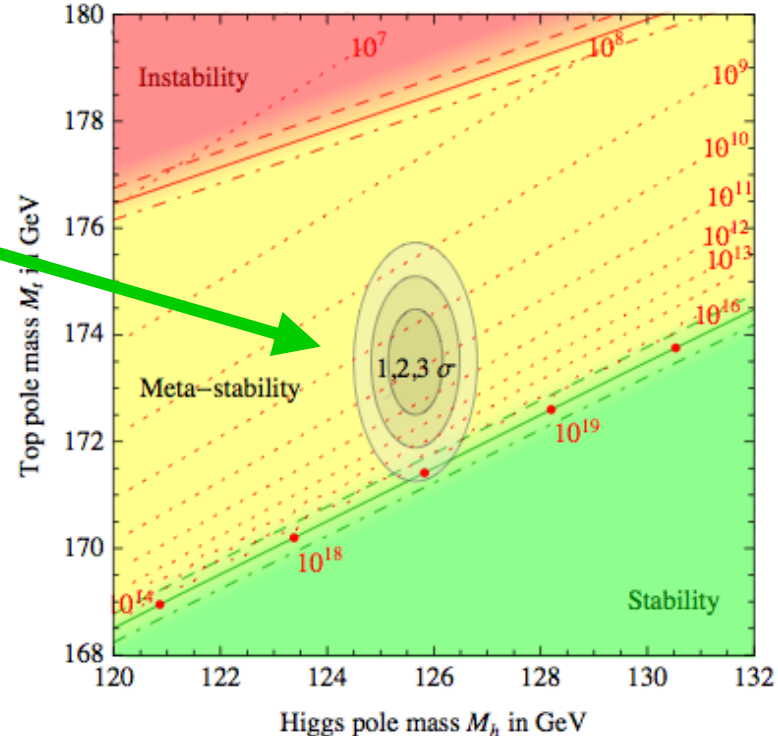
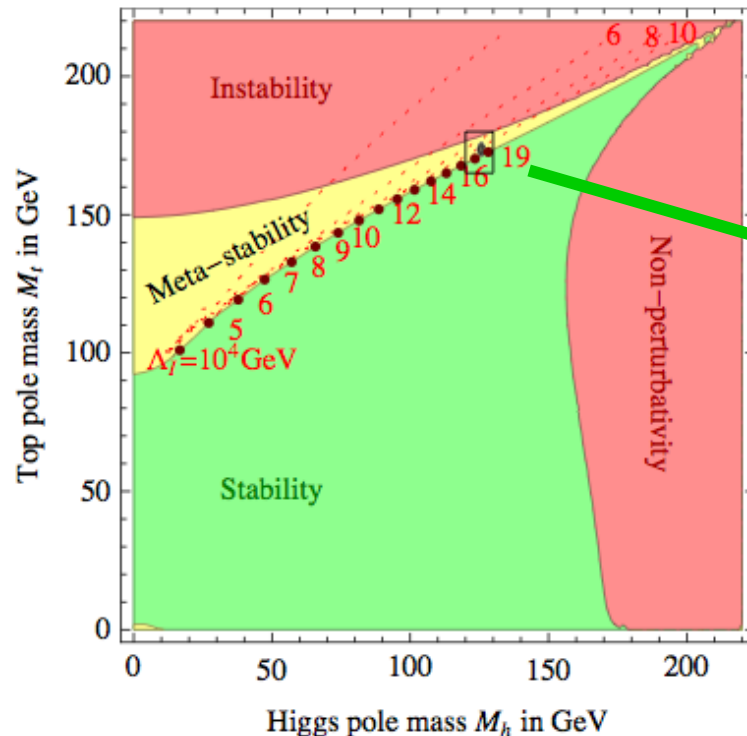


# No no-go theorem for the SM at large energies

The pure SM evolution of couplings leads to a metastable Universe

The SM evolution up to  $M_{\text{Pl}}$  leads to a narrow critical wedge: a hidden message?

Buttazzo et al '13 see also Branchina '13



$$\lambda\phi^4$$





The absence of new physics appears as a paradox to us

Still the picture repeatedly suggested by the data in the last ~20 years is simple and clear

Take the SM, extended to include Majorana neutrinos and some form of DM, as valid up to some very high energy

Thus, ignoring the FT, minimal modifications to the SM are being considered

Neutrino masses? See-Saw mechanism

Baryogenesis? Thru leptogenesis

Dark Matter? Simple WIMPs, Axions, keV sterile  $\nu$ 's.....

Coupling Unification? Some large scale threshold, e.g. non-SUSY SO(10) with an intermediate scale

GA, Meloni '13



Possibly Nature has a way, hidden to us, to realize a deeper form of naturalness at a more fundamental level

## Completing the SM with $\nu_R$

It is sufficient to introduce 3 RH gauge singlets  $\nu_R$   
[each completing a 16 of  $SO(10)$  for one generation]  
and not artificially impose that L is conserved

In the SM, in the absence of  $\nu_R$ , B and L are “accidental”  
symmetries [i.e. no renormalizable gauge invariant  
B and/or L non-conserving vertices can be built from  
the fields of the theory]

But we know that non perturbative terms (instantons)  
break B and L (not B-L) and also non renorm. operators:

Weinberg  
dim-5 operator

$$O_5 = \frac{(Hl)_i^T \lambda_{ij} (Hl)_j}{\Lambda} + h.c.$$

With Majorana  $\nu_R$  renormalizable mass terms are  
⊕ allowed by gauge symmetries and break L (and B-L)

## A very natural and appealing explanation (see-saw):

$\nu$ 's are nearly massless because they are Majorana particles and get masses through L non conserving interactions suppressed by a large scale  $M \sim M_{\text{GUT}}$

$$m_\nu \sim \frac{m^2}{M}$$

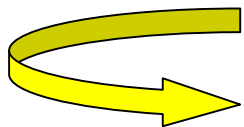
$$m: \leq m_t \sim v \sim 200 \text{ GeV}$$

$M$ : scale of L non cons.

Note:

$$m_\nu \sim (\Delta m^2_{\text{atm}})^{1/2} \sim 0.05 \text{ eV}$$

$$m \sim v \sim 200 \text{ GeV}$$



$$M \sim 10^{14} - 10^{15} \text{ GeV}$$

Observation  
of  $0\nu\beta\beta$   
would  
confirm that  $\nu$   
are Majorana

⊕ This is so impressive that, in my opinion, models with  $\nu_R$  at the EW scale or around are strongly disfavoured

# A great extra bonus of see-saw with heavy Majorana $\nu_R$ 's

Baryogenesis via Leptogenesis  
near the GUT scale



(after inflation)

Buchmuller, Yanagida,  
Plumacher, Ellis, Lola,  
Giudice et al, Fujii et al

.....

Only survives if  $\Delta(B-L)$  is not zero  
(otherwise is washed out at  $T_{ew}$  by instantons)

Decays of lightest  $\nu_R$  ( $M \sim 10^{11-12}$  GeV) satisfy Sacharov conditions

L non conserv. & CP violat.'n in  $\nu_R$  out-of-equilibrium decay:  
B-L excess survives at  $T_{ew}$  and gives the obs. B asymmetry.

Quantitative studies confirm that the range of  $m_i$  from  
 $\nu$  oscill's is compatible with BG via (thermal) LG

Buchmuller, Di Bari, Plumacher;  
Giudice et al; Pilaftsis et al;  
Hambye et al

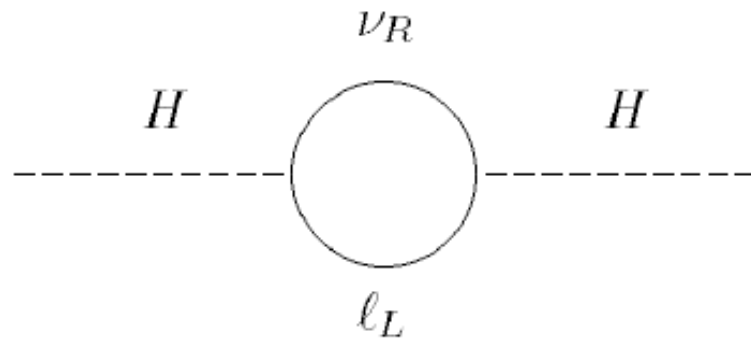


Heavy  $\nu_R$  well match with GUT's [recall the 16 of SO(10)!]

(if for naturalness SUSY is invoked, one also has the bonus that coupling unification and proton decay are OK, ...)

But so far, no SUSY or any New Physics

If only the SM + Majorana  $\nu$ 's, then heavy  $\nu_R$  are unnatural and require fine tuning:



$$\begin{aligned} & \text{for } q \gg M_R \\ \delta\mu^2 & \approx \frac{y_\nu^2}{(2\pi)^2} M_R^2 \log(q/M_R) \\ & \approx \frac{m_\nu M_R^3}{(2\pi v)^2} \log(q/M_R) \end{aligned}$$

$$\mu < 1 \text{ TeV} \longrightarrow M_R < 10^7 - 10^8 \text{ GeV}$$

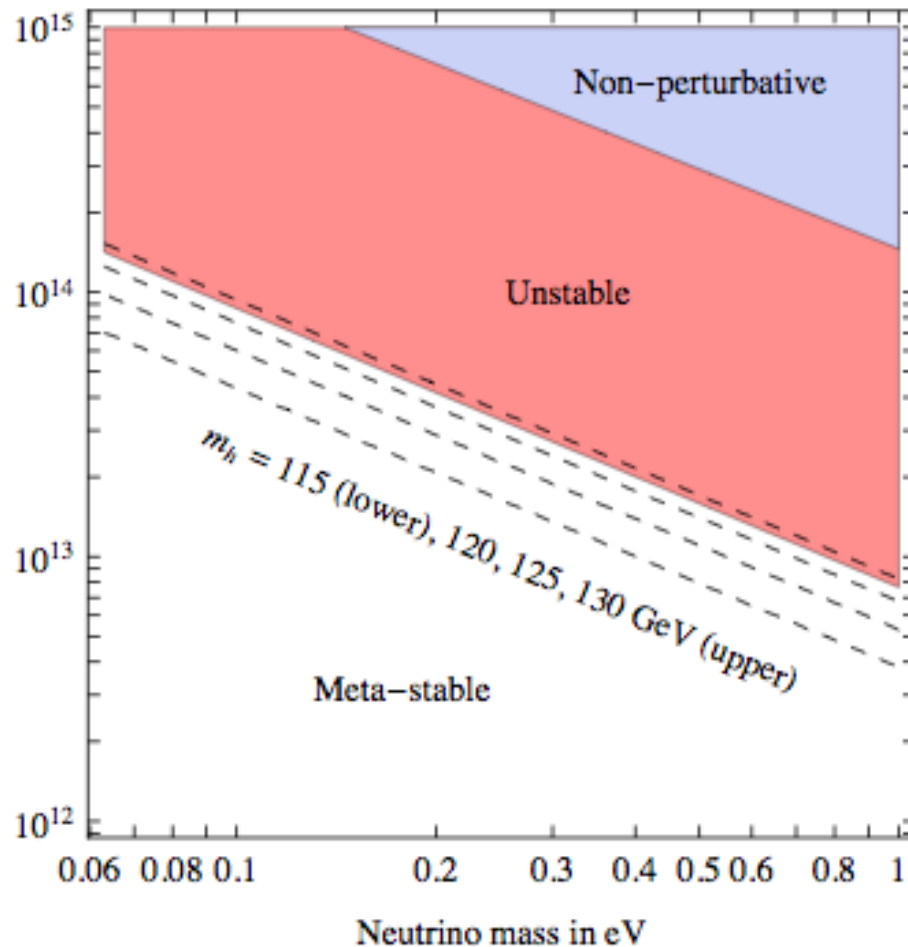


Vissani '97; Elias-Miro et al '11;  
Farina et al '13; De Gouvea et al '14

Heavy  $\nu_R$ 's further de-stabilize the vacuum

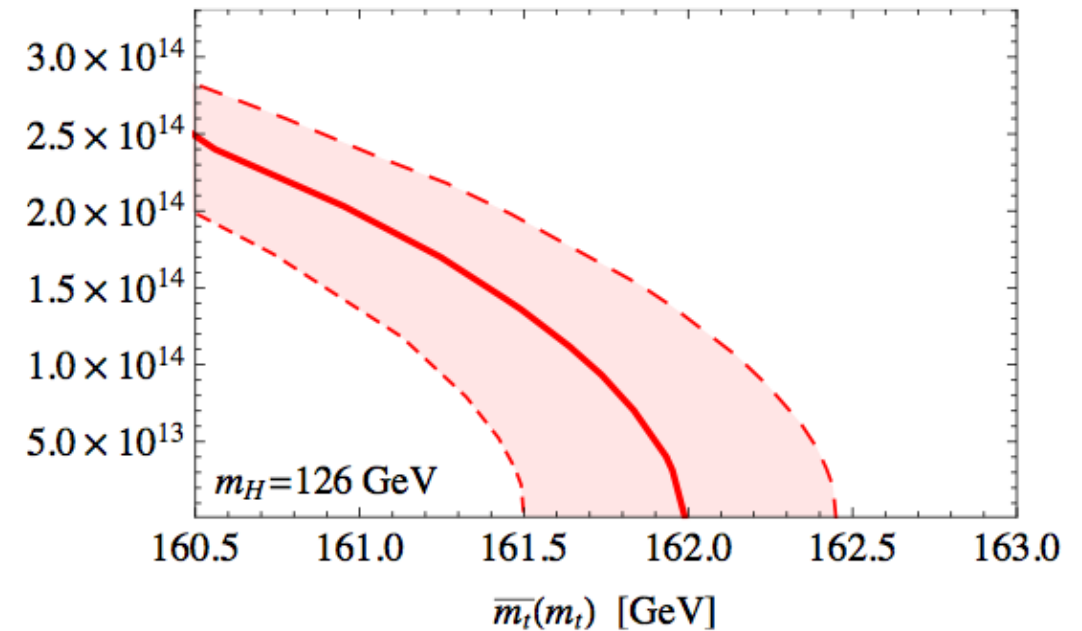
But, for  $M < 10^{14}$  GeV,  $\nu_R$ 's do not make the vacuum unstable

J. Elias-Miro' et al '11



$m_{\nu R}$  [GeV]

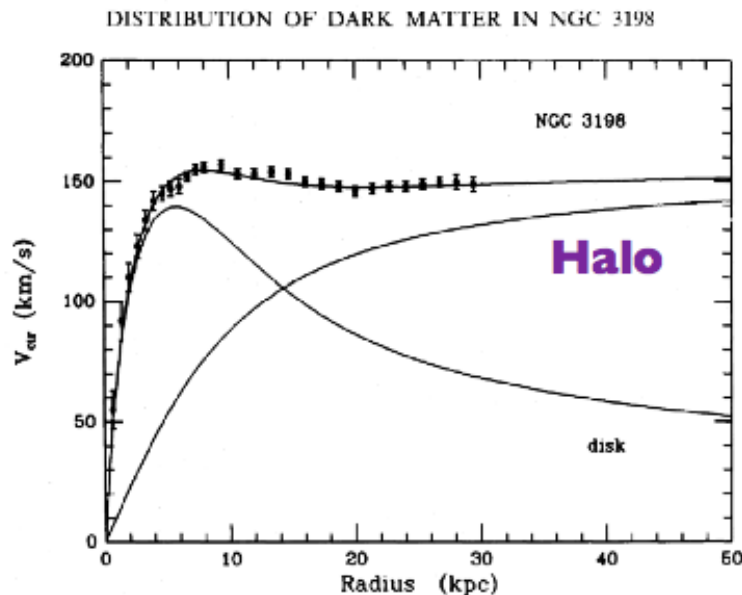
Masina'12



At present Dark Matter is **THE** crucial problem in particle physics

A by now robust evidence for Dark Matter in the Universe

Rotation of galaxies



Lensing



MACS, HST

Merging clusters of galaxies



M. Markevitch et al 2003

Cosmological evidence

anisotropies of Micro Wave Background Radiation

large scale structure

structure formation.....

e.g. Planck



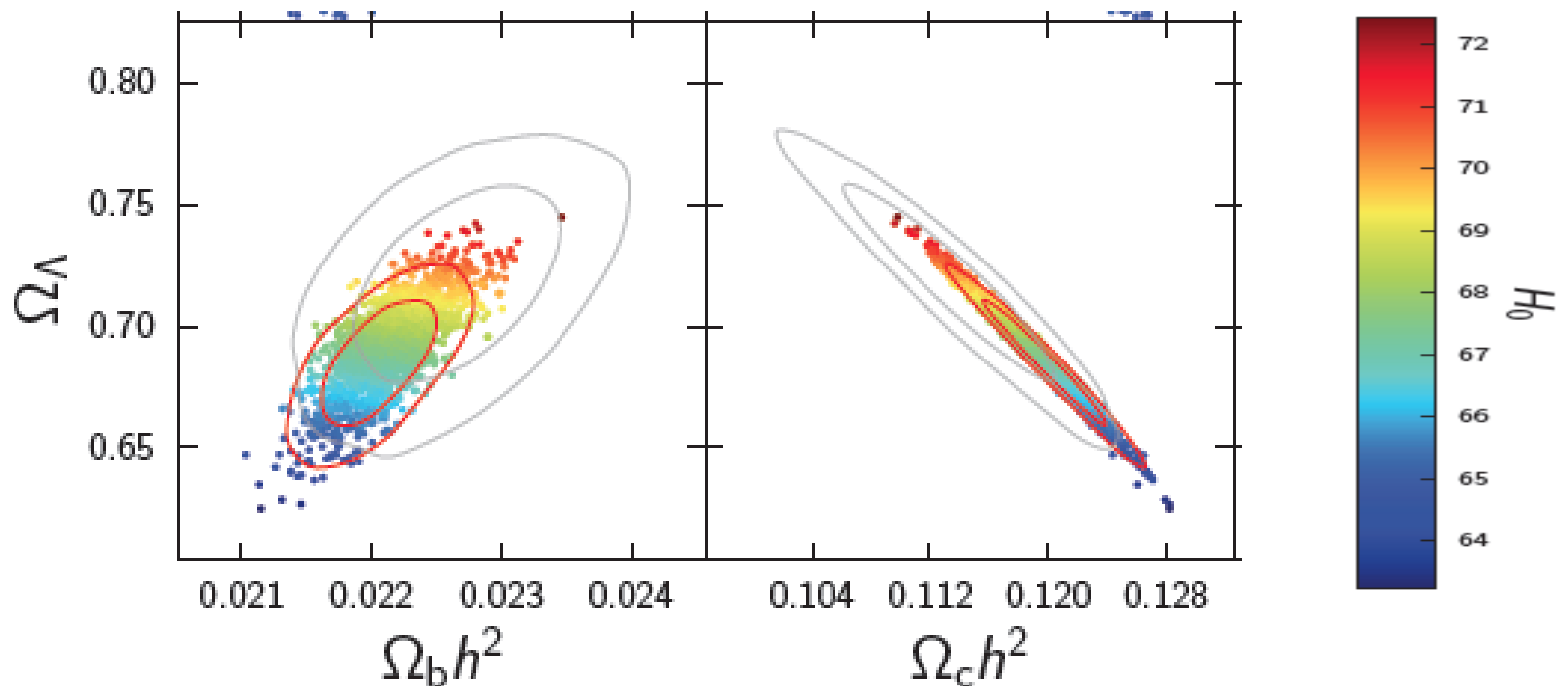
# Planck fits of DM ArXiv:1303.5076

$$H_0 = 100 h \text{ km s}^{-1} \text{ Mpc}^{-1}$$

Parameter	Planck+WP		Planck+WP+highL+BAO	
	Best fit	68% limits	Best fit	68% limits
$\Omega_b h^2$ . . . . .	0.022032	$0.02205 \pm 0.00028$	0.022161	$0.02214 \pm 0.00024$
$\Omega_c h^2$ . . . . .	0.12038	$0.1199 \pm 0.0027$	0.11889	$0.1187 \pm 0.0017$
$H_0$ . . . . .	67.04	$67.3 \pm 1.2$	67.77	$67.80 \pm 0.77$

$\Omega_c =$  cold DM density

$h \sim 0.67$





While for neutrino masses, baryogenesis... we have definite ideas on how these problems could be solved  
Dark Matter remains mysterious and is a very compelling argument for New Physics and the most pressing challenge for particle physics

A partial list of main candidates:

- WIMP's
- Axions
- keV sterile neutrinos

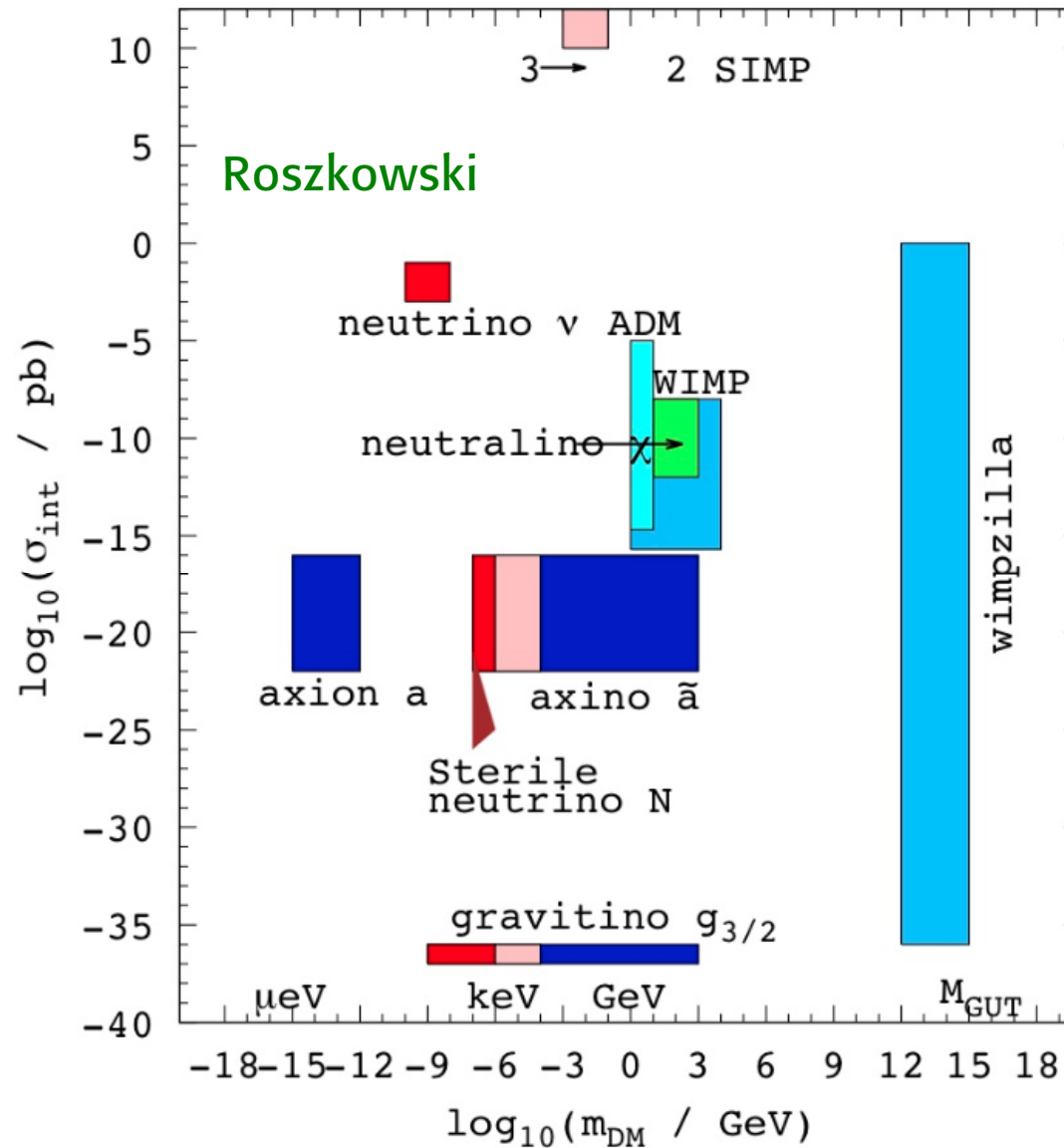
The 3 active  $\nu$ 's cannot make the whole of DM. Bounds:

- Dwarf Galaxies  $\rightarrow m > \text{few hundreds eV}$  (Tremaine-Gunn)
- Galaxies  $\rightarrow m > \text{few tens eV}$
- Hot DM also excluded by structure formation



Light sterile  $\nu$ 's ( $m \sim \text{eV}$ ) are also inadequate

In the literature the DM candidates span an enormous range of mass



**WIMP's:** Weakly Interacting Massive Particles  
with  $m \sim 10^{-1}-10^3$  GeV

WIMP's still are optimal candidates:

LHC can reach most kinds of WIMP's

For WIMP's in thermal equilibrium after inflation the density is:

$$\Omega_{\chi} h^2 \simeq \text{const.} \cdot \frac{T_0^3}{M_{\text{Pl}}^3 \langle \sigma_A v \rangle} \simeq \frac{0.1 \text{ pb} \cdot c}{\langle \sigma_A v \rangle}$$

can work for typical weak cross-sections!!!

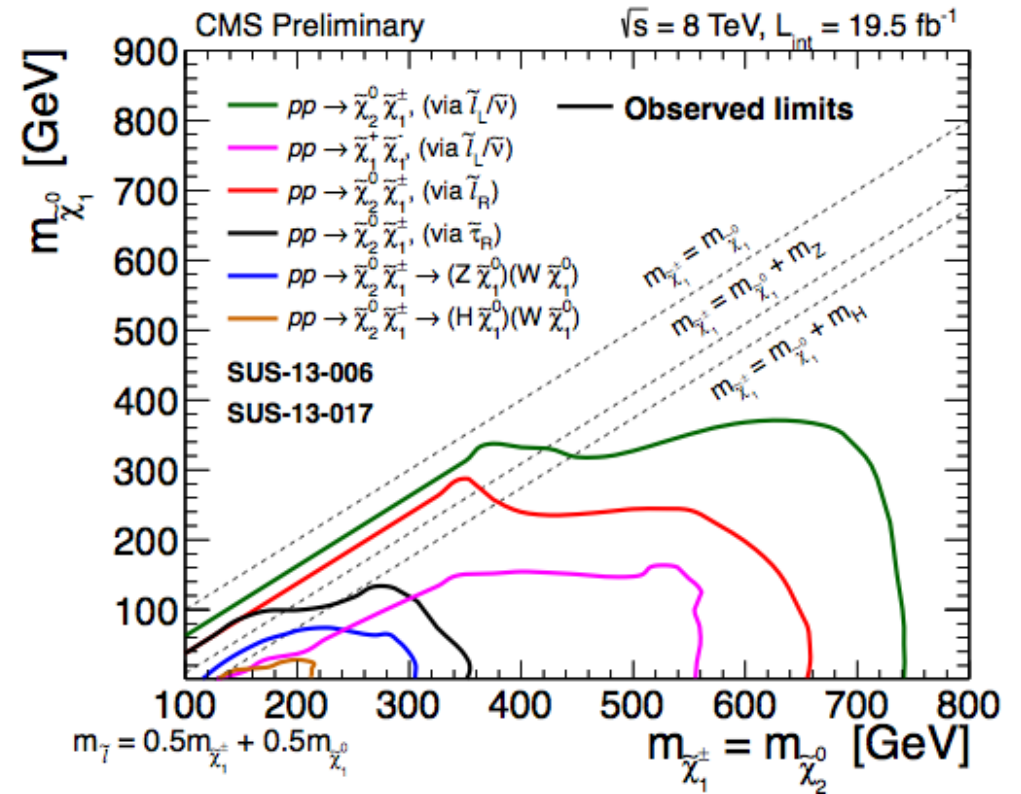
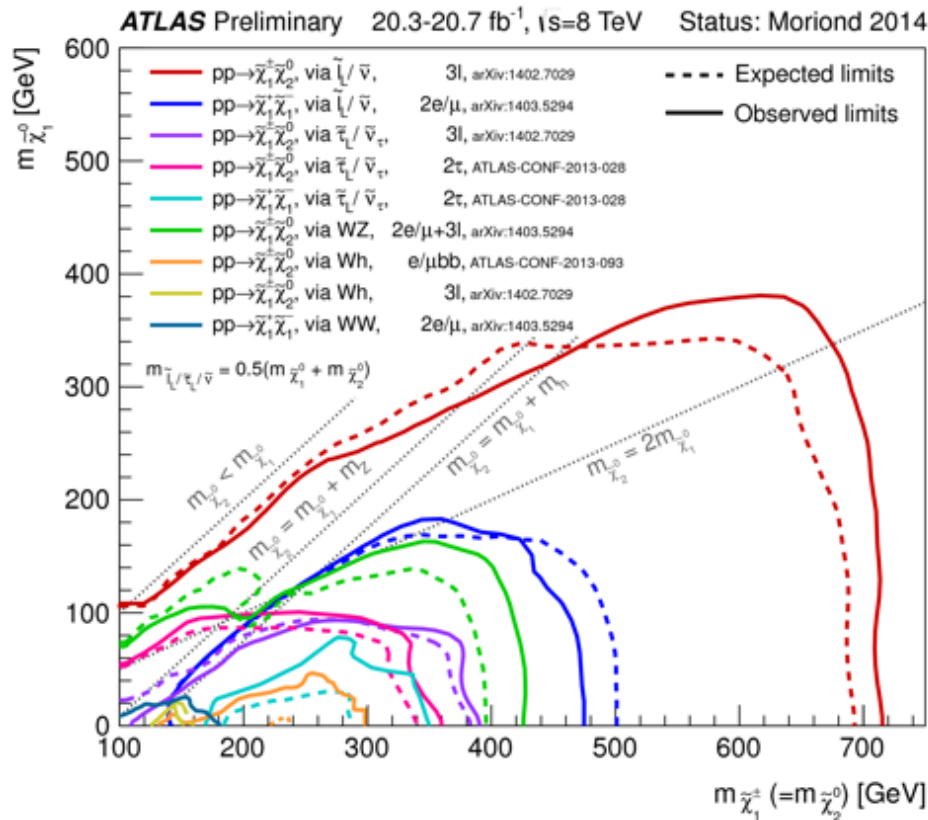
This “coincidence” is taken as a good indication in favour  
of a WIMP explanation of Dark Matter



No WIMP's have been observed at the LHC

But the limits on SUSY WIMPS (neutralinos) are not too stringent

In large regions of parameter space  $m_{\chi^0} < 350$  GeV is allowed



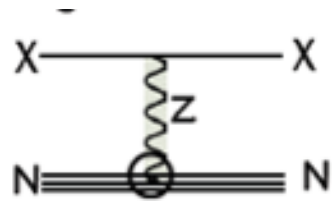
A strict bound is very low:  $m_{\chi^0} > 25$  GeV (light s-taus and higgsinos)

Calibbi et al'13

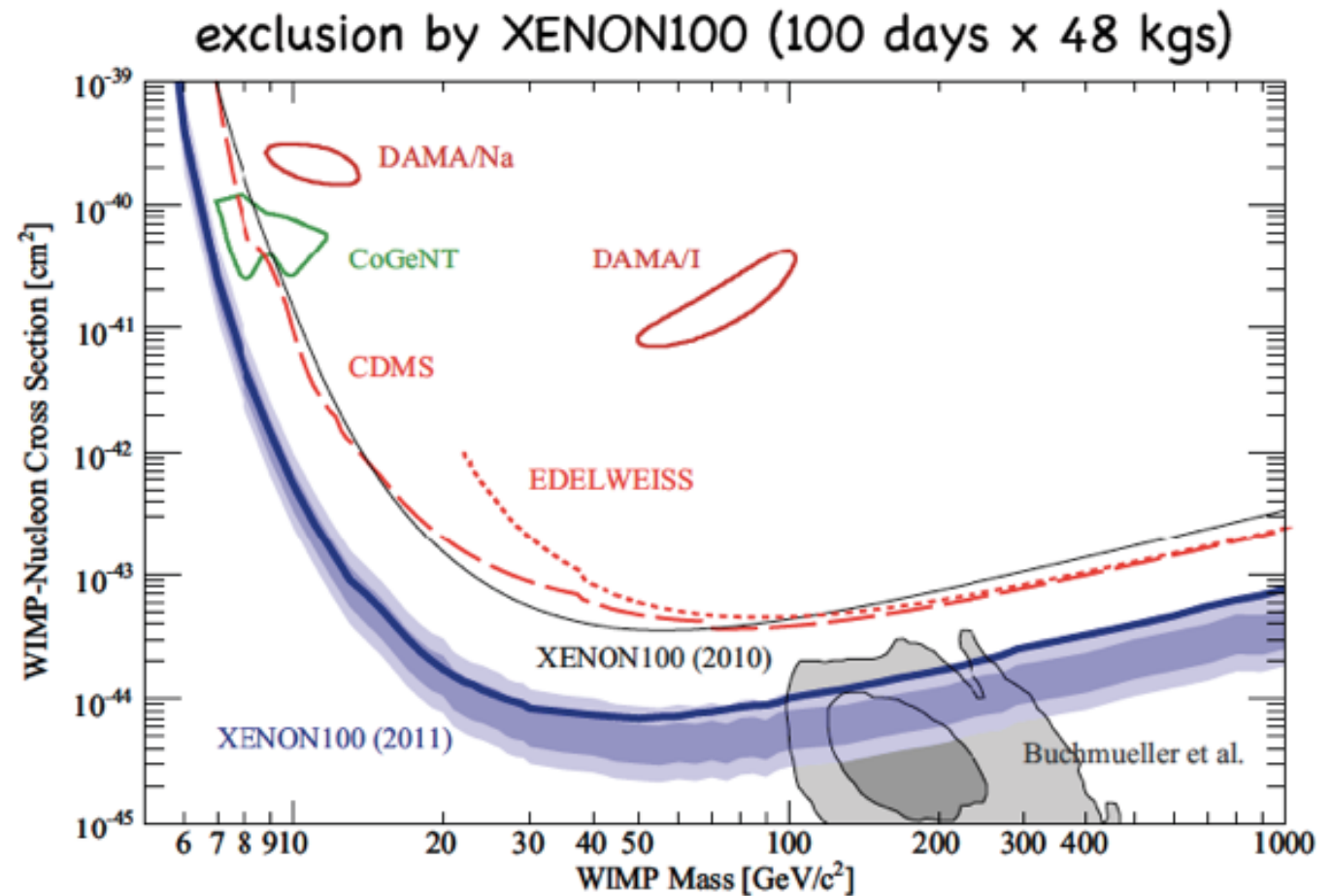
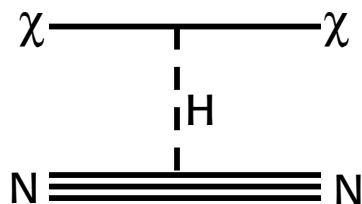
## Non accelerator searches

$$\chi N \rightarrow \chi N$$

Z exchange  
potentially  
large



125 GeV Higgs  
boson exchange  
being also  
probed now



DM coupled to Z severely limited (axial couplings less constrained)

**LUX constraints strongest**

De Simone, Giudice, Strumia '14

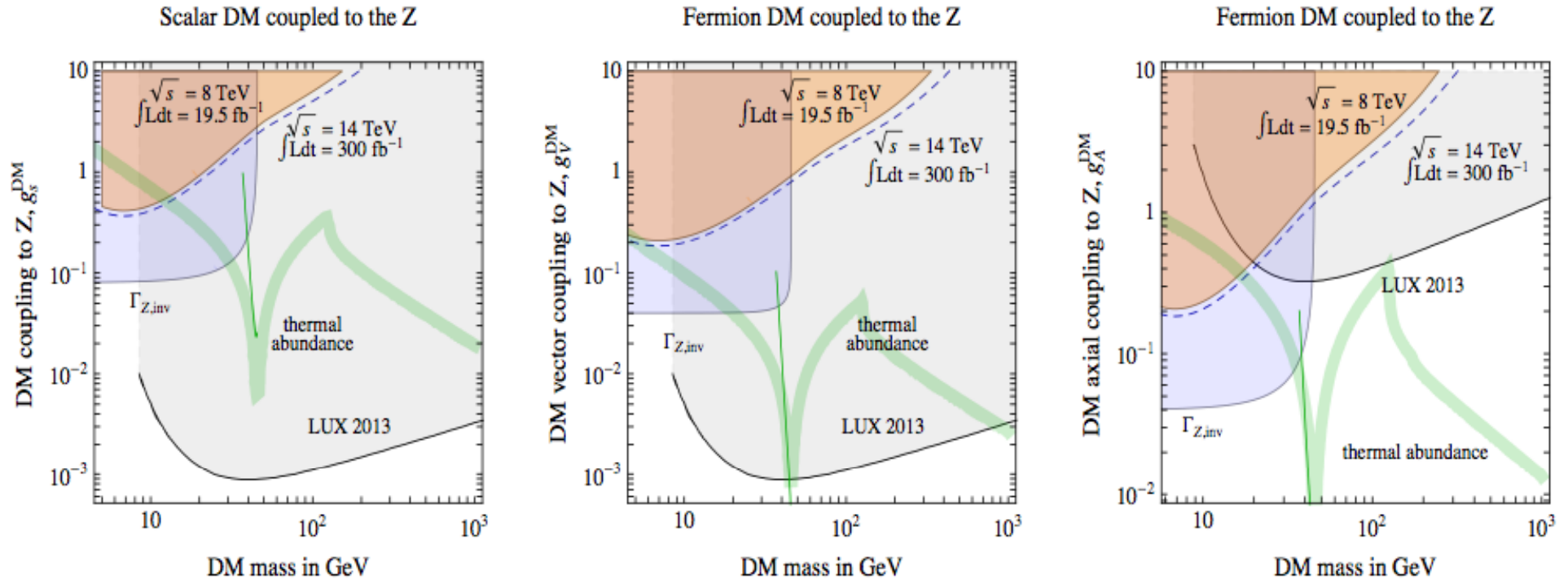


Figure 3: **DM coupled to the Z.** Regions of DM mass  $M_{\text{DM}}$  and Z couplings ( $g_s^{\text{DM}}$ ,  $g_V^{\text{DM}}$ ,  $g_A^{\text{DM}}$ ): the orange region is excluded at 90% CL by ATLAS mono-jet searches at LHC8, with forecast for LHC14 (dashed blue line); the grey region is excluded at 90% CL by LUX 2013 direct searches; the blue region is excluded by the Z-invisible width constraint  $\Gamma_{Z,\text{inv}} < 2$  MeV. The green solid curve corresponds to a thermal relic abundance via Z-coupling annihilation equal to the observed DM density (the thick curve is the off-shell estimation; the thin curve is the on-shell computation).



# DM coupled to Higgs also limited (pseudo scalar couplings less constrained)

De Simone, Giudice, Strumia '14

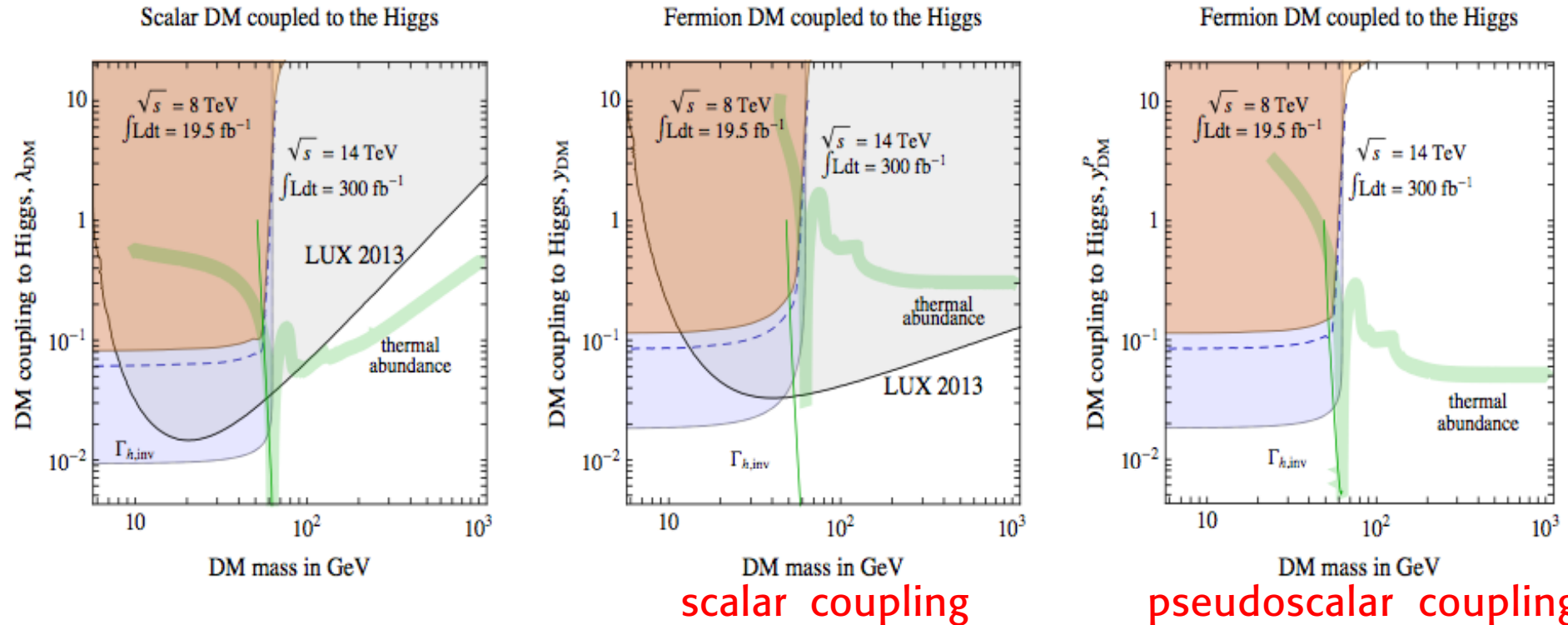
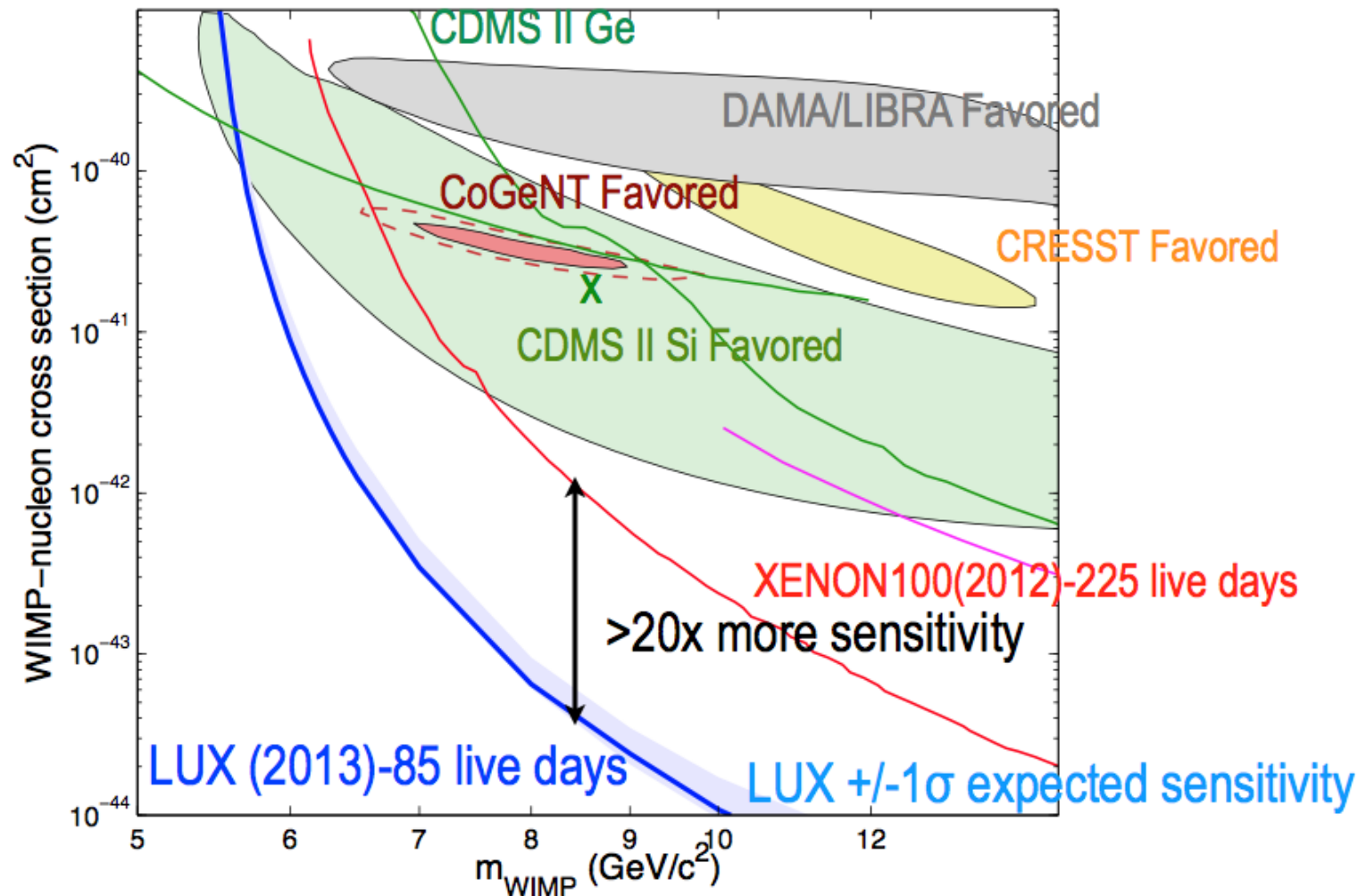


Figure 4: **DM coupled to the Higgs.** Regions of DM mass  $M_{\text{DM}}$  and Higgs couplings ( $\lambda_{\text{DM}}$ ,  $y_{\text{DM}}$ ,  $y_{\text{DM}}^P$ ): the orange region is excluded at 90% CL by ATLAS mono-jet searches at LHC8, with forecast for LHC14 (dashed blue line); the grey region is excluded at 90% CL by LUX 2013 direct searches; the blue region is excluded by the Higgs invisible width constraint  $\Gamma_{h,\text{inv}}/\Gamma_h < 20\%$ . The green solid curve corresponds to a thermal relic abundance via Higgs-coupling annihilation equal to the observed DM density (the thick curve is the off-shell estimation; the thin curve is the on-shell computation).

## Low mass $\sim 10$ GeV WIMP's?

CDMS-Si [ArXiv :1304.4279](#) 3 events in the signal region  
Now excluded by LUX [ArXiv:1310.8214](#)

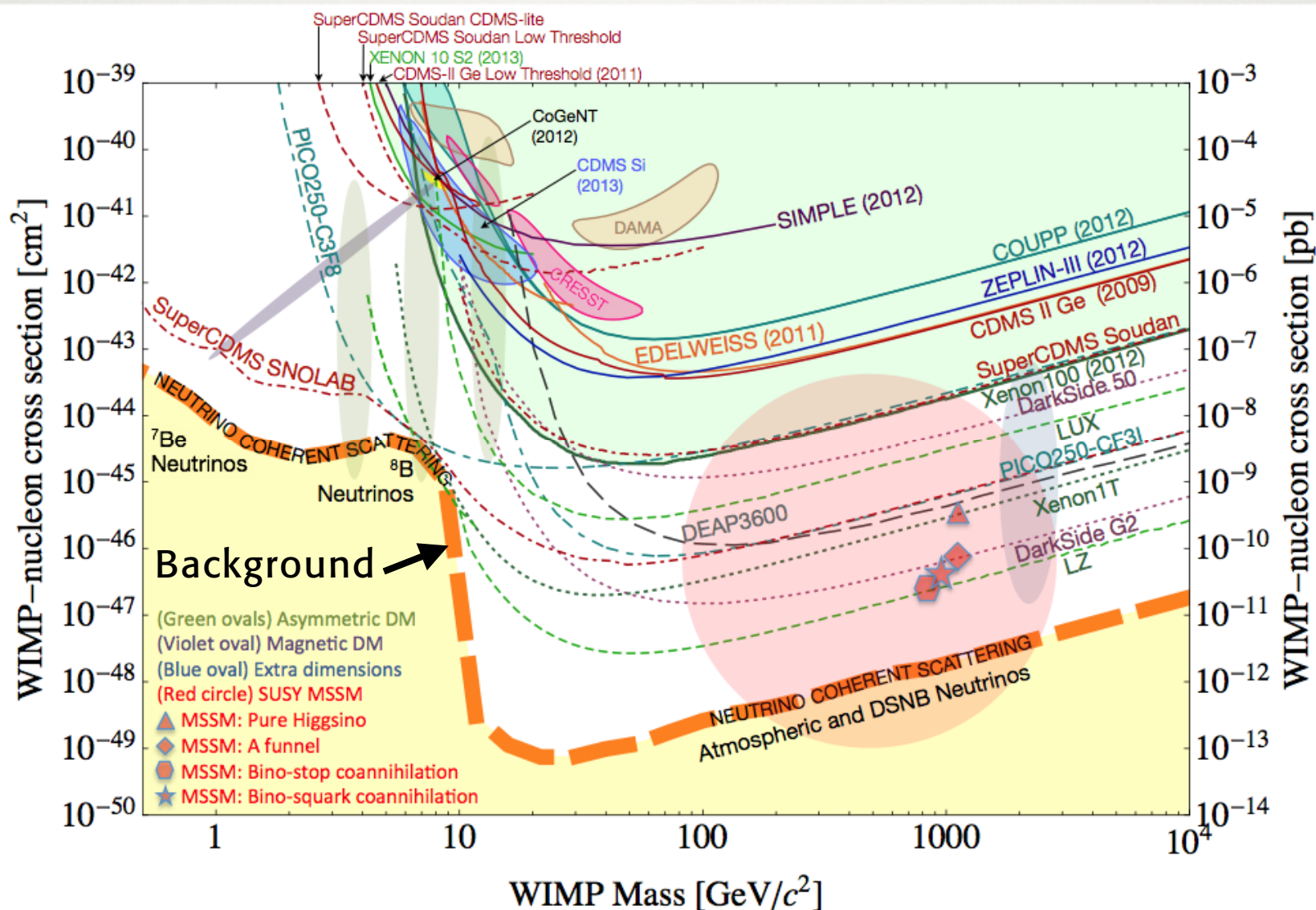


However  
there is  
still plenty of  
room  
for low mass  
WIMP's



# The WIMP non-accelerator search continues

CF1 Snowmass report, 1310.8327



# The Axion [Peccei-Quinn (PQ) solution to strong CP problem]

PQ introduce a new U(1) symmetry:  $U(1)_{PQ}$

Ex.: introduce new fermions  $\psi$  (charged colour triplets) and a scalar A

$U(1)_{PQ} :$   $\psi' = e^{i\gamma_5 \alpha} \psi$  Kim'79, Shifman, Vainshtein, Zacharov'80 (KSVZ)  
 $A' = e^{-2i\alpha} A$  No other fields are charged under  $U(1)_{PQ}$   
 $\longrightarrow M \bar{\psi} \psi$  and  $H \bar{\psi} \psi$  (H=Higgs)

The VEV  $\langle A \rangle \sim f$  spont.  
breaks  $U(1)_{PQ}$

are forbidden, while  $\lambda A \bar{\psi} \psi$  is allowed

The  $\psi$  mass is  $m \sim \lambda \langle A \rangle \sim \lambda f \longrightarrow$  new particles at scale  $f$ !

$A = |A| e^{i \frac{a}{f}}$   $a$  (the axion) is the Goldstone boson

$a' = a - 2i\alpha f$  it only has derivative couplings  
except for the  $U(1)_{PQ}$  anomaly term

$$L_{axion} = -\frac{1}{2} \partial_\mu a \partial^\mu a + L_{int}(\psi, \frac{\partial_\mu a}{f}) + [\theta + \frac{a}{f}] \frac{\alpha_s}{4\pi} Tr(F_{\alpha\beta} \tilde{F}^{\alpha\beta})$$



The only term with  $a$  and not  $\partial_\mu a$  is the potential  $V$

$$V = [\theta + \frac{a}{f}] \frac{\alpha_s}{4\pi} \text{Tr}(F_{\alpha\beta} \tilde{F}^{\alpha\beta})$$

The VEV  $\langle a \rangle$  is fixed by  $\frac{\partial V}{\partial a} = 0 \Rightarrow \frac{\alpha_s}{4\pi f} \langle \text{Tr}(F_{\alpha\beta} \tilde{F}^{\alpha\beta}) \rangle = 0$

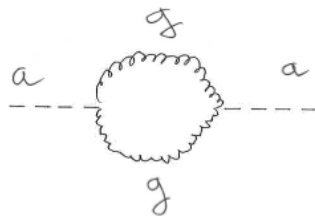
It is (not too) easy to prove that  $\langle \text{Tr}(F_{\alpha\beta} \tilde{F}^{\alpha\beta}) \rangle \propto \sin \theta_{\text{eff}} \equiv \sin(\theta + \frac{\langle a \rangle}{f})$   
 so that the coefficient of the  
 CP violating term is put to zero!

e.g. Coleman, '77; Vafa, Witten '84.....

After the shift  $a \rightarrow a'' + \langle a \rangle$  ( $a''$  is the field for perturbation theory)

we are left with the coupling  $\frac{a''}{f} \frac{\alpha_s}{4\pi} \text{Tr}(F_{\alpha\beta} \tilde{F}^{\alpha\beta})$  and no CP violation

This coupling also induces a mass for the axion (it would be massless if not for the anomalous breaking of  $U(1)_{PQ}$ )



$$m_a^2 \propto \frac{\Lambda_{QCD}^4}{f^2}$$

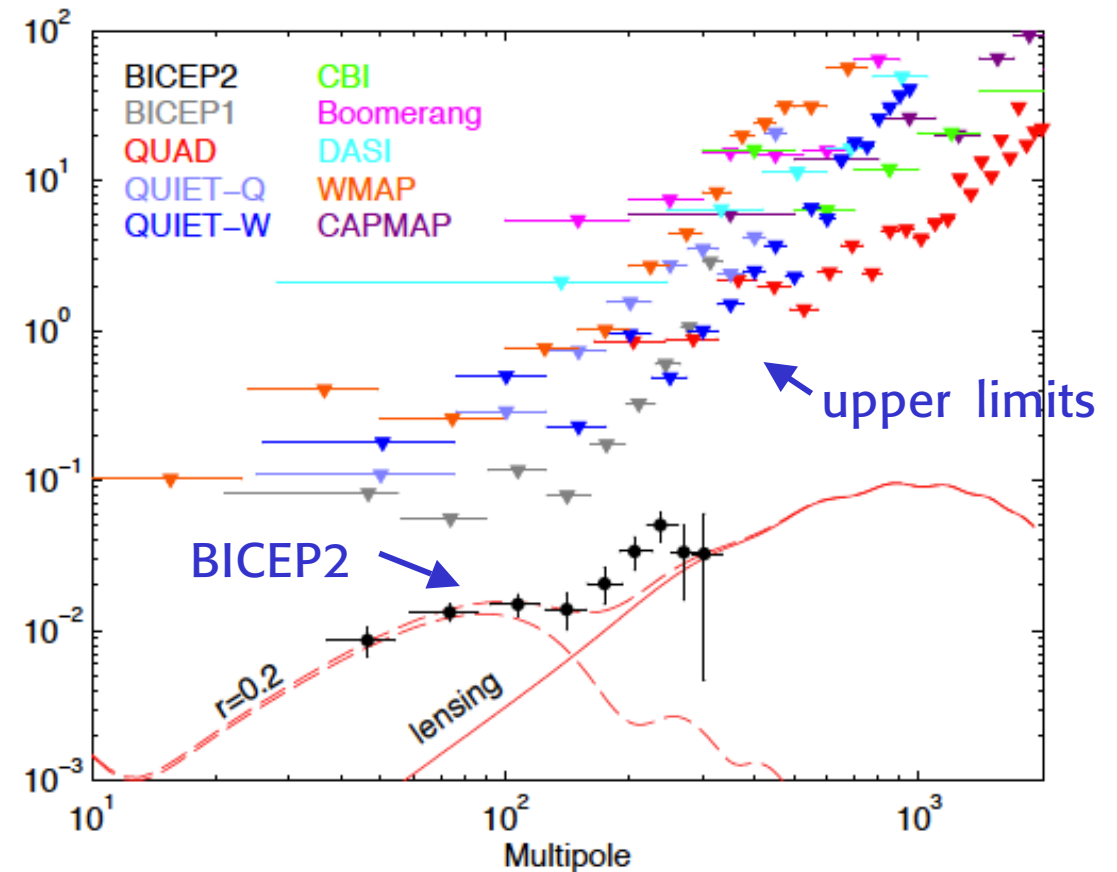
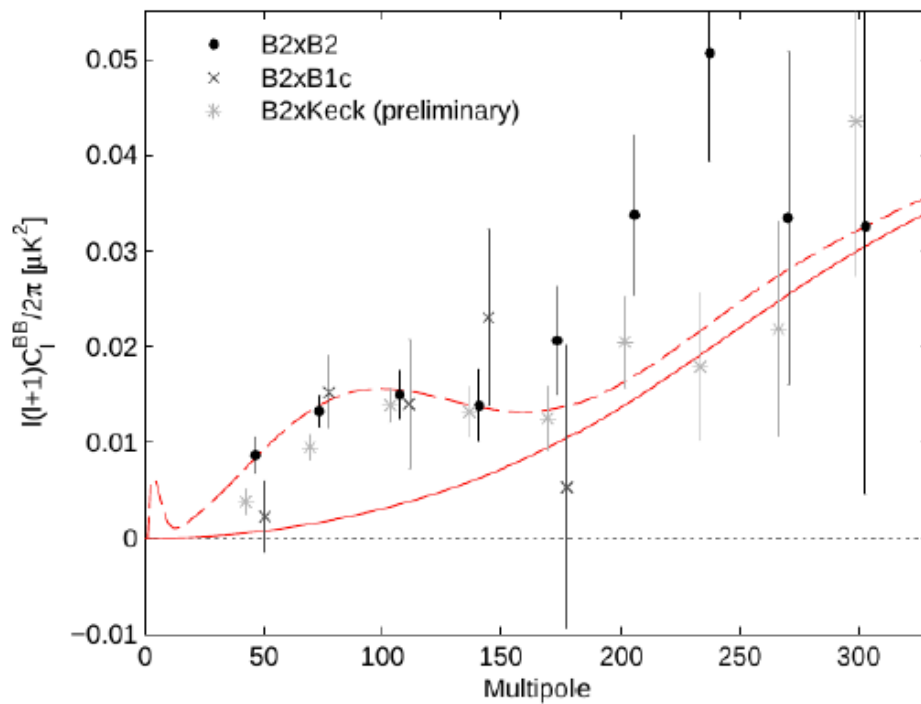
with  $f$  large,  $m_a$  is small,  
 the axion coupling is small,  
 and the  $\psi$  mass is large



The analogous coupling to photons induces the decay  $a \rightarrow \gamma\gamma$

# Sensational news from cosmology

## The BICEP2 Data

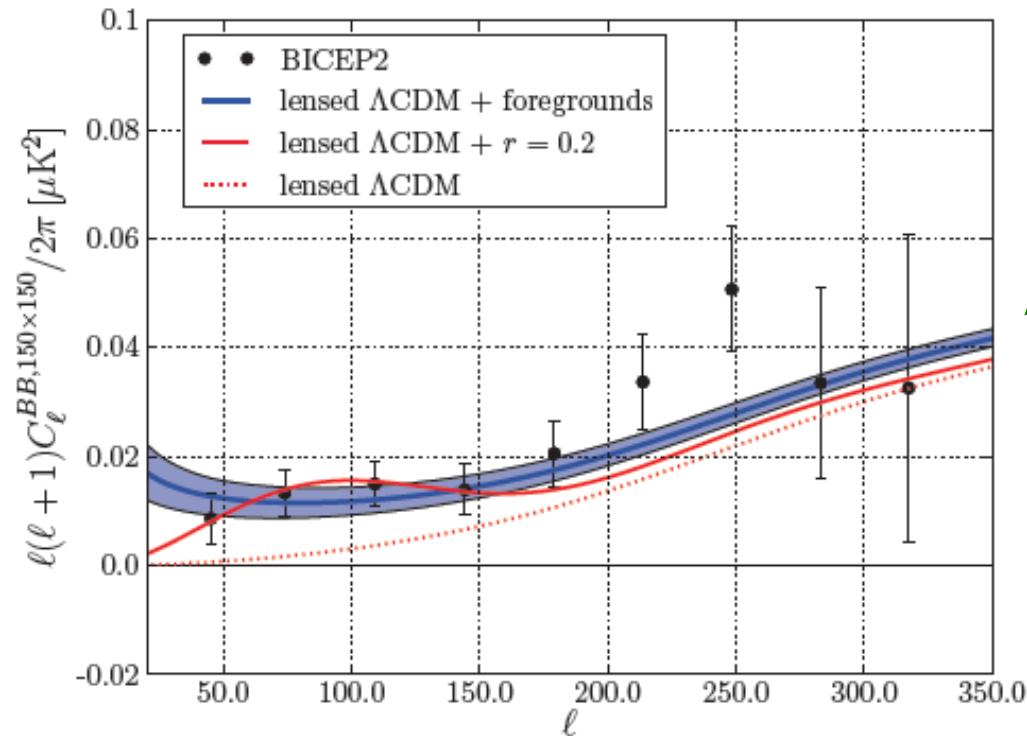


A large value of  $r = A_T/A_s \sim 0.2$  is found



Caution: it could be that it all comes from dust!

Flauger, Hill, Spergel '14



A model of  
foreground dust  
polarized emission

Planck data (released in October ? )  
and/or measurements at different frequencies  
can tell the difference

To be confirmed by other experiments  
Planck, POLARBEAR, ACTPole, Keck Array...



# Great impact on Inflation

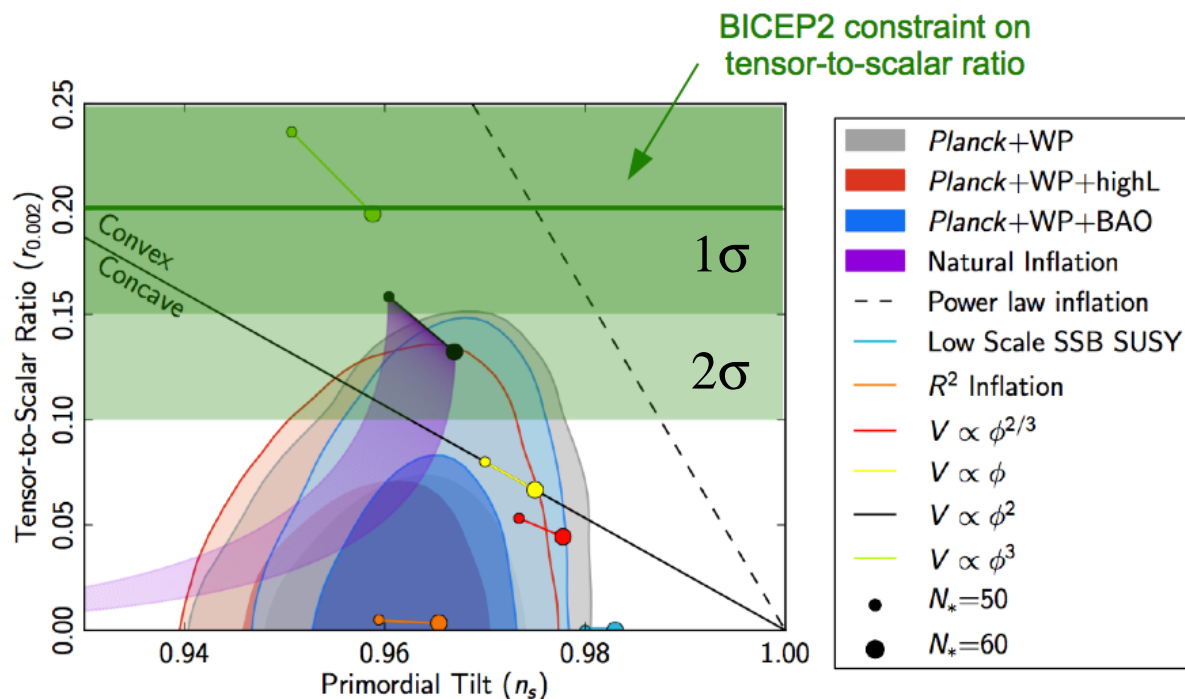
$$r = 0.20^{+0.07}_{-0.05}$$

Energy scale of inflation:

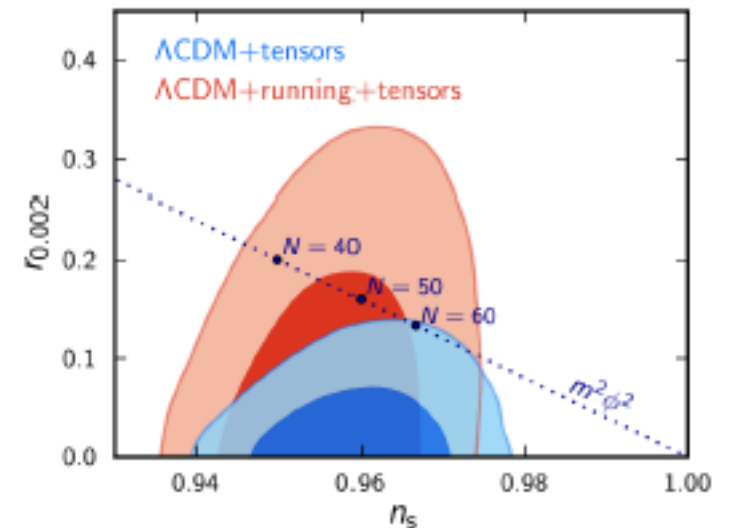
$$V_{\text{inf}}^{1/4} \approx 2.2 \cdot 10^{16} \left( \frac{r}{0.2} \right)^{1/4} \text{ GeV}$$

Tantalizing close to  $M_{\text{GUT}}$ !

Evidence of a scale below  $M_{\text{Pl}}$ ?

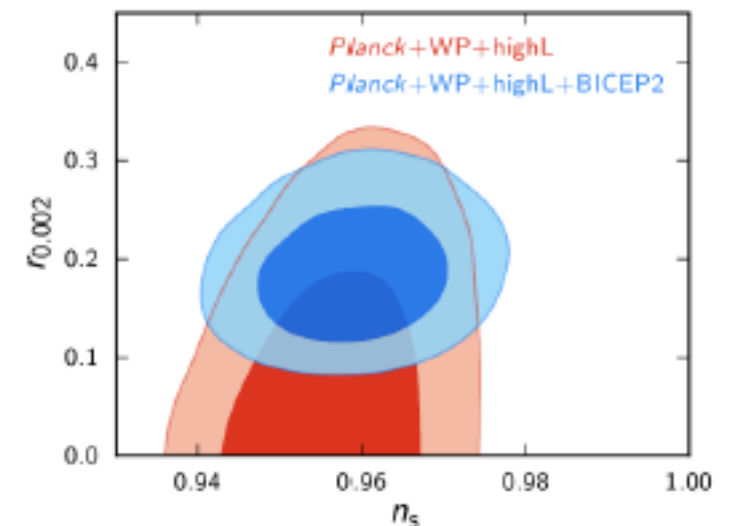


## Planck 2013

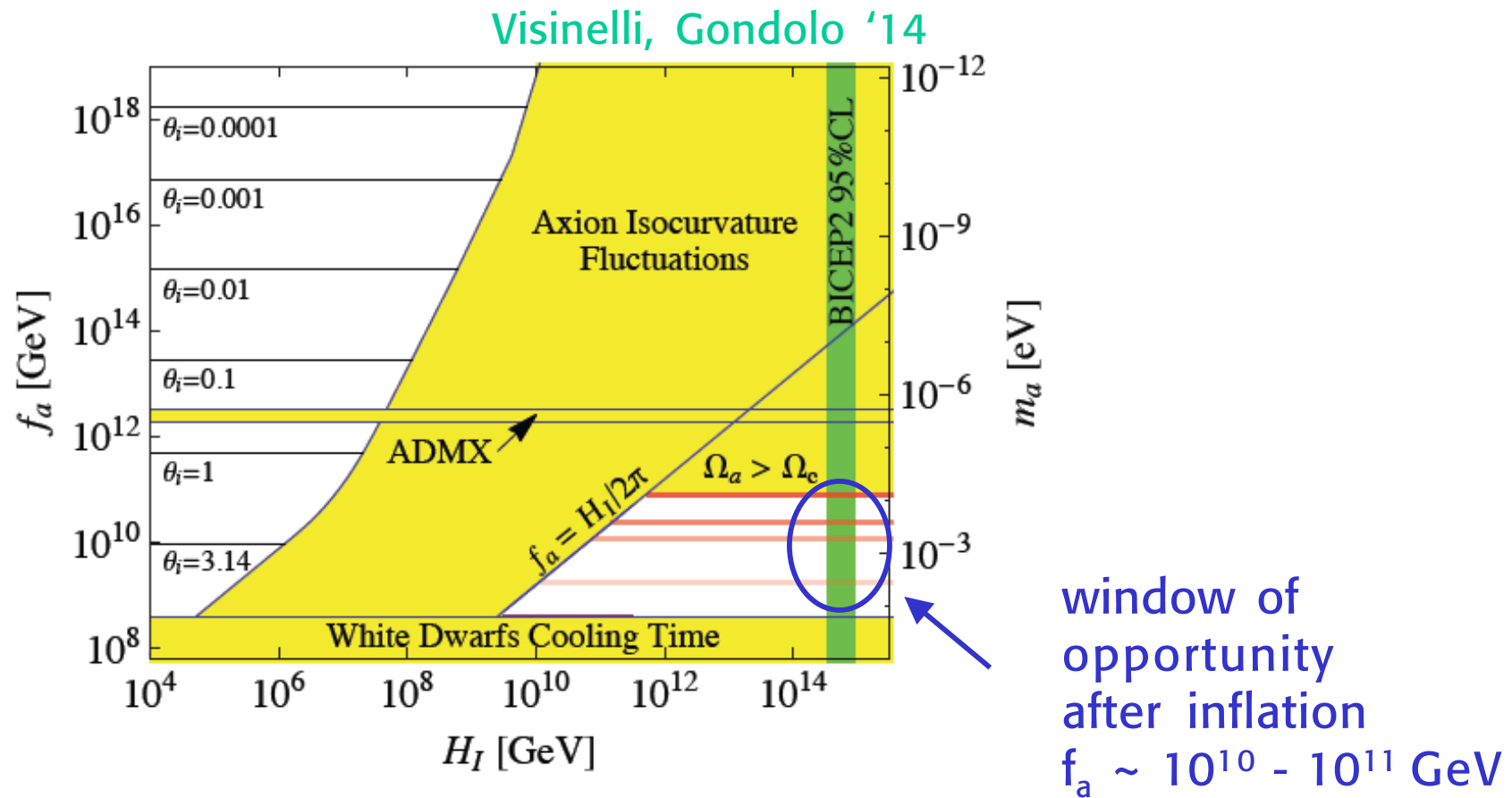


Moderate tension with Planck

## Bicep2 2014



# Implications of BICEP2 on axions

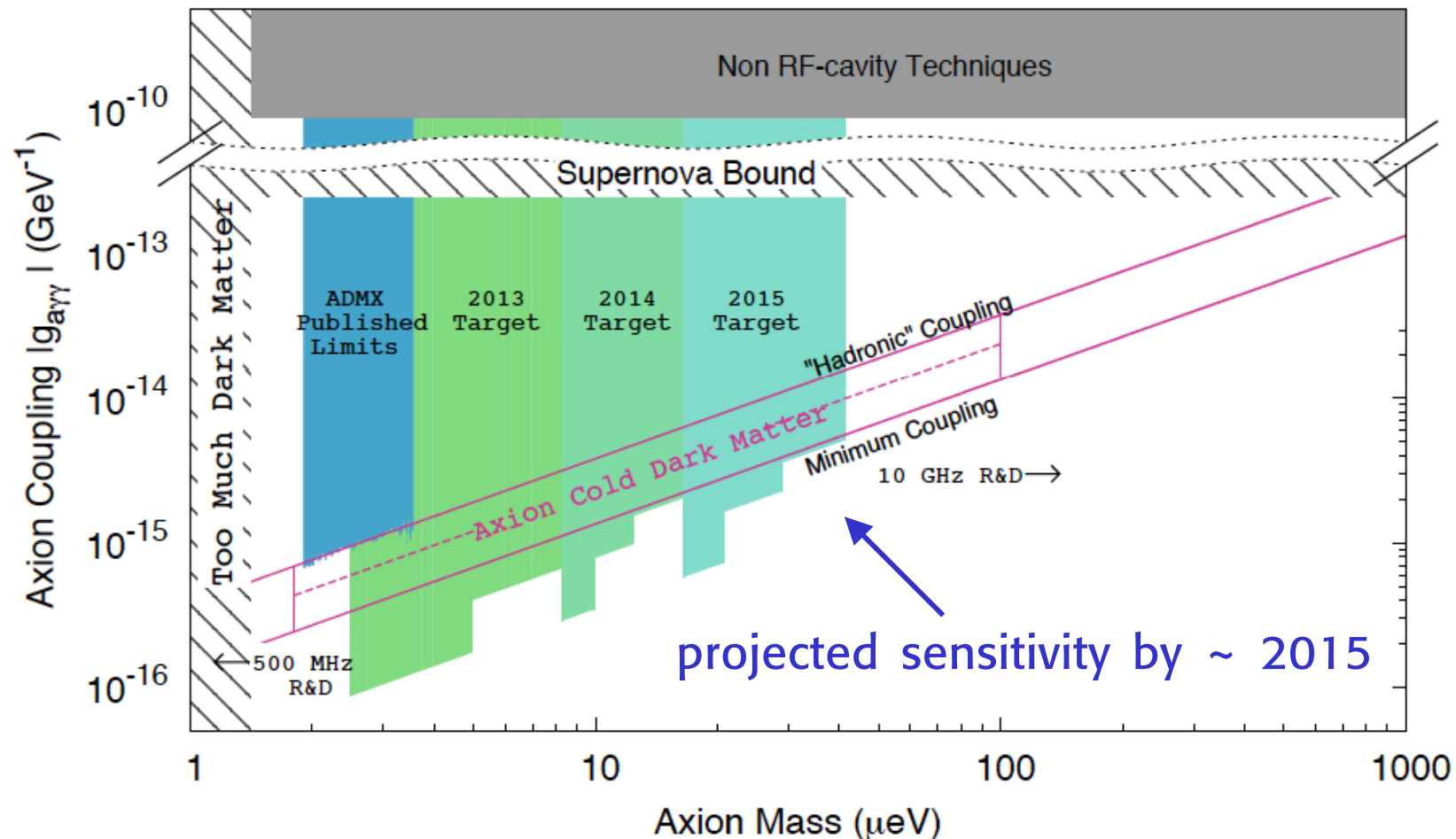




Axion searches are very important

## ADMX: the Axion Dark Matter Experiment

University of Washington at Seattle





To cope with the naturalness riddle different lines of thought have emerged

- Insist on minimizing the fine tuning: imagine suitable forms of new physics around the corner
- Opt for a total acceptance of fine tuning: the anthropic philosophy
- Accept fine tuning only up to an intermediate scale: e.g. split SUSY
- Argue that possibly there is no fine tuning: the no new threshold (up to  $M_{Pl}$ ) conjecture



# One line: insisting on minimizing the FT

- **"Stealth" Naturalness:** build models where naturalness is restored not too far from the weak scale but the related NP is arranged to be not visible so far Fine-tuning the fine-tuning-suppression mechanism?

## Two main directions

SUSY

Composite Higgs

For an orderly retreat  
simplest new ingredients are

- Compressed spectra
- Heavy first 2 generations
- NMSSM (an extra Higgs singlet)

H as PGB of extended symm.  
q and l mix with comp. ferm.  
Key role of light top partners

The last trench of natural SUSY!



# Going beyond the MSSM: an extra singlet Higgs

In a promising class of models a singlet Higgs  $S$  is added and the  $\mu$  term arises from the  $S$  VEV (the  $\mu$  problem is solved)

$$\lambda S H_u H_d$$

additional term

$$m_h^2 = M_Z^2 \cos^2 2\beta + \lambda^2 v^2 \sin^2 2\beta + \delta_t^2$$

Mixing with  $S$  can modify the Higgs mass and couplings at tree level

Hall et al '11, King et al '12, Barbieri et al '13.....

**NMSSM:**  $\lambda < \sim 0.7$  the theory remains perturbative up to  $M_{\text{GUT}}$   
(no need of large stop mixing, less fine tuning)

**$\lambda$  SUSY:**  $\lambda \sim 1 - 2$  for  $\lambda > 2$  theory non pert. at  $\sim 10$  TeV

It is not completely excluded that at 126 GeV the second heaviest is seen while the lightest escaped detection at LEP



Ellwanger '11, Belanger et al '12

- Going beyond the MSSM: **Natural SUSY**  $\longrightarrow$  Minimum for MSSM to be natural  
 $m_{\tilde{g}}, m_{\tilde{t}}, m_{\tilde{b}}, m_{\tilde{h}} < \sim 1 \text{ TeV}$

Heavy 1st, 2nd generation scalars

Flavour and CP problems improved

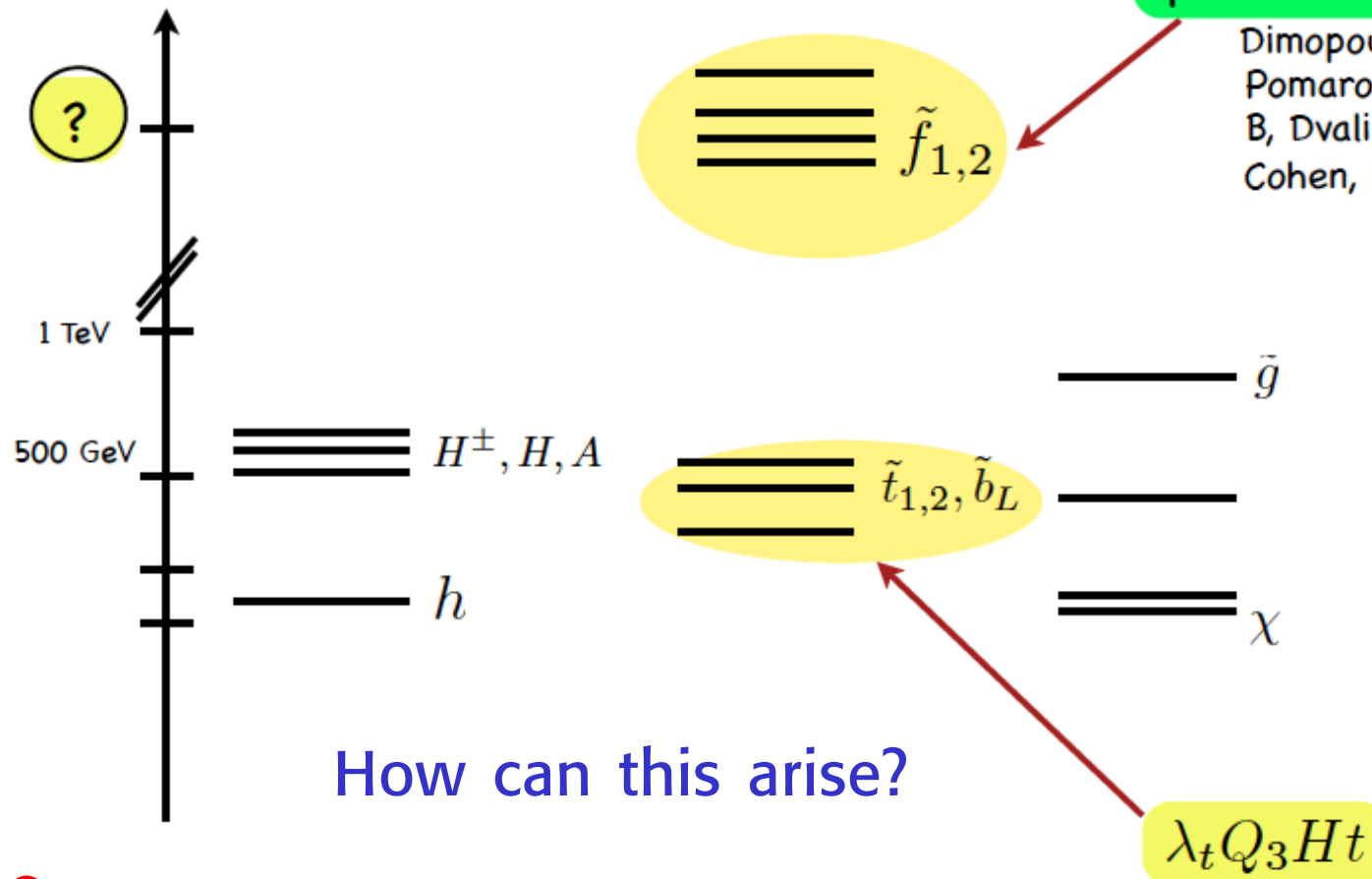
Dimopoulos, Giudice 1995  
 Pomarol, Tommasini 1995  
 B, Dvali, Hall 1995  
 Cohen, Kaplan, Nelson 1996

pioneer papers

recent papers, e.g.

Papucci et al '11  
 Brust et al '11  
 Essig et al '11  
 Katz et al '11  
 Larsen et al '12  
 Csaki et al '12  
 .....

For g-2  
 light sleptons  
 welcome



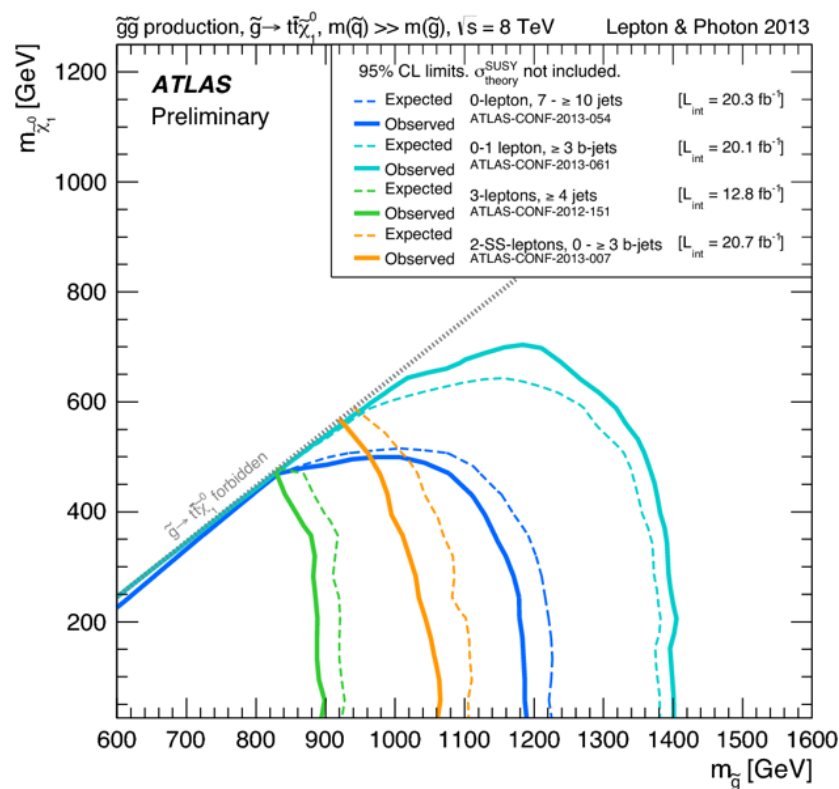
Barbieri

# Searches of light gluinos, s-top, s-bottom: already biting hard

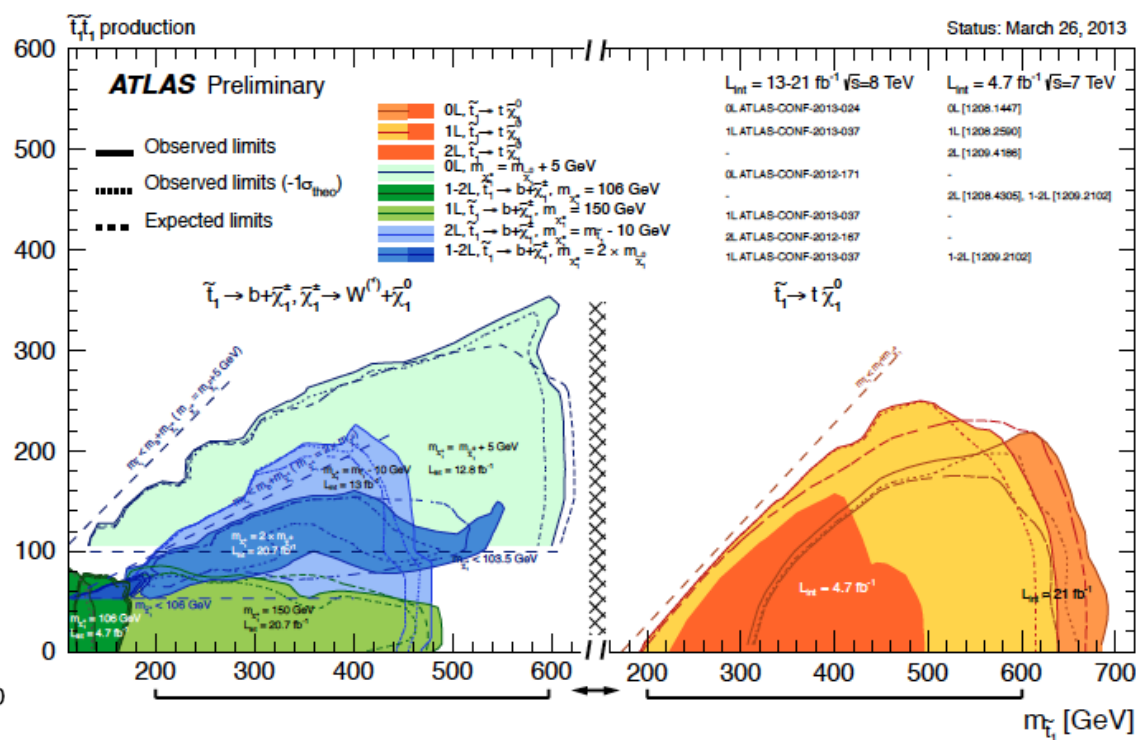
Gluino mediated s-top production:  $m_g < 1.4$  TeV excluded with some assumptions on BRs.

Direct s-top production:  $m_{stop} < 0.60$ - $0.65$  TeV excluded assuming 100% BR for either  $b\chi^+$  or  $t\chi^0$

ATLAS



gluino



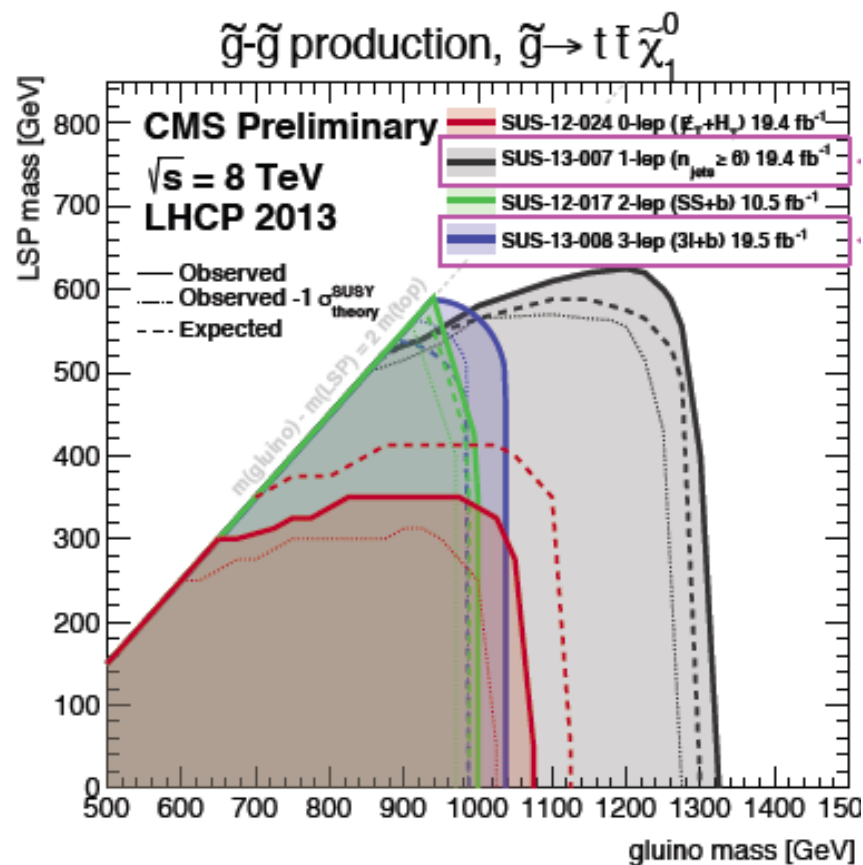
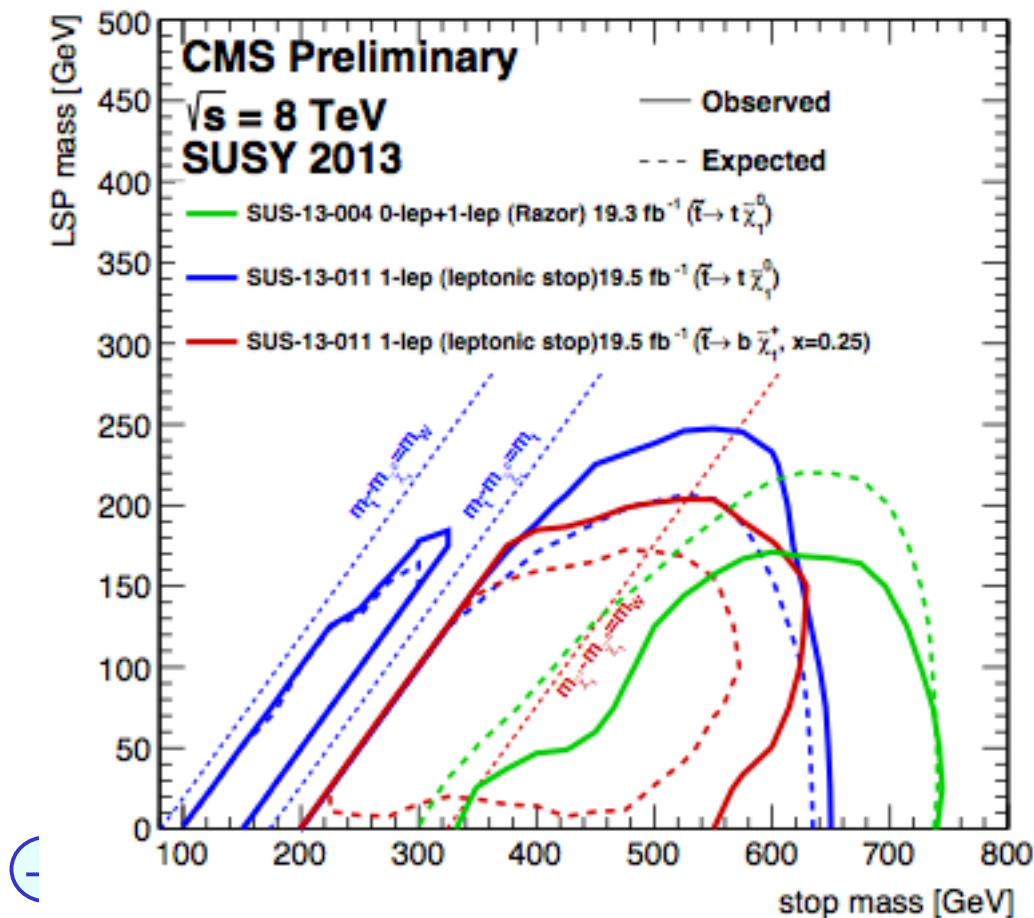
s-top

# Searches for stops, gluinos, sbottoms target natural SUSY

- Probe stops up to  $\sim 650$  GeV
- Probe gluinos up to  $\sim 1.3$  TeV
- Probe sbottoms up to  $\sim 600$  GeV

$\tilde{t}\tilde{t}$  production

CMS



# • Composite Higgs

Georgi, Kaplan '84; Kaplan '91; Agashe, Contino, Pomarol '05; Agashe et al '06; Giudice et al '07; Contino et al '07; Csaki, Falkowski, Weiler '08; Contino, Servant '08; Mrazek, Wulzer '10; Panico, Wulzer '11; De Curtis, Redi, Tesi '11; Marzocca, Serone, Shu '12; Pomarol, Riva '12; De Simone et al '12.....

The light Higgs is a bound state of a strongly interacting sector and a pseudo-Goldstone boson of an enlarged symmetry.  
eg.  $SO(5)/SO(4)$ . Can be set up in a holographic ED context.

$v \sim \text{EW scale}$        $f \sim \text{SI scale}$

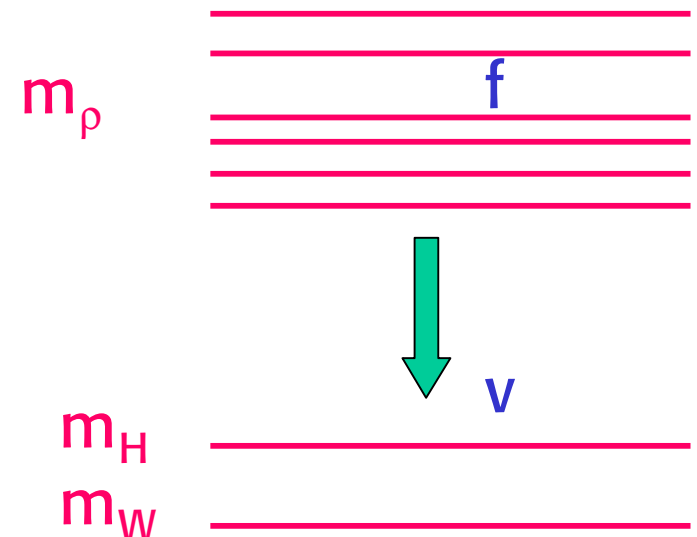
$$\sim f < m_\rho < \sim 4\pi f$$

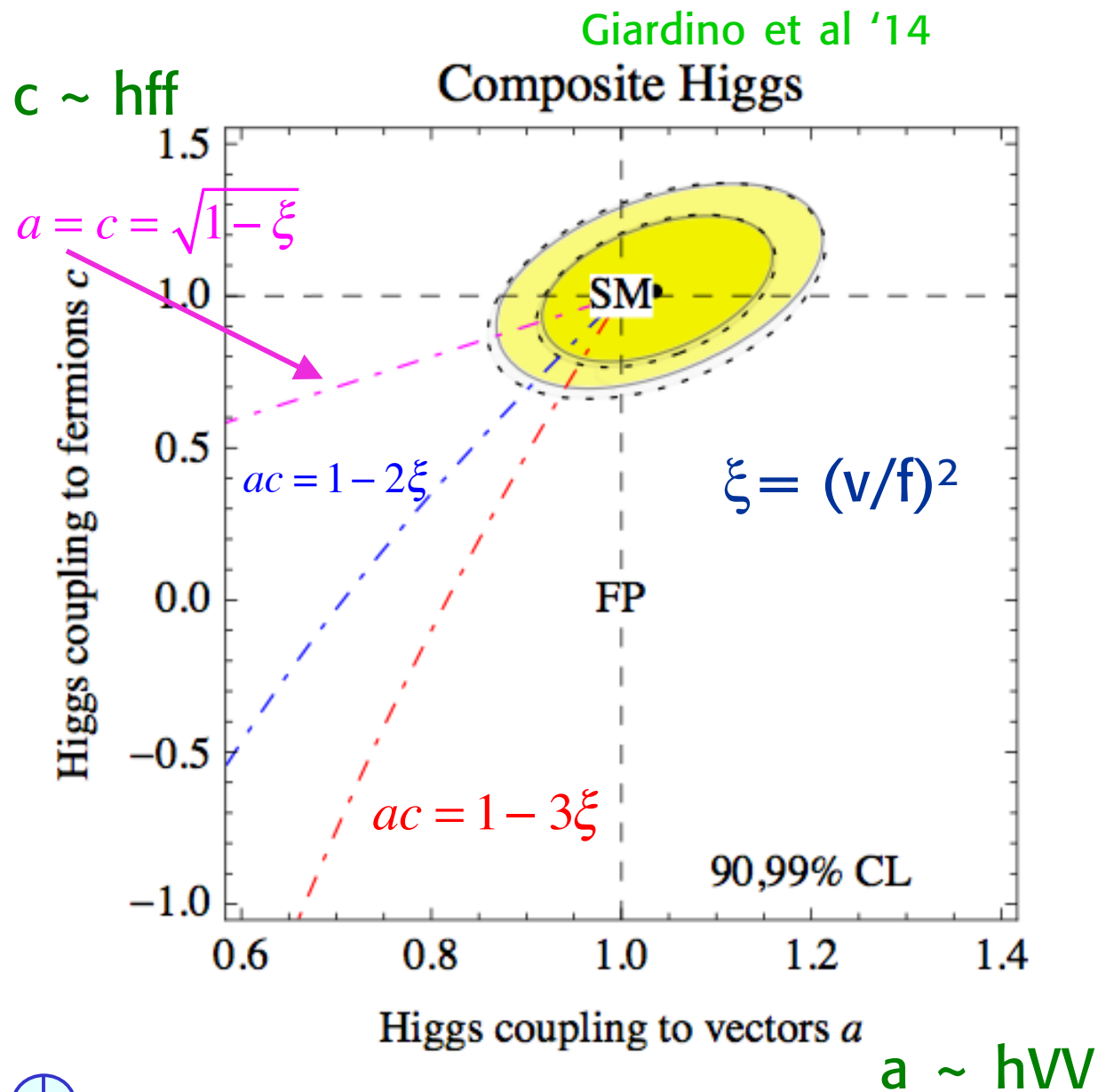
$$\xi = (v/f)^2$$

$\xi$  interpolates between SM [ $\xi \sim 0$ ]  
and some degree of compositeness

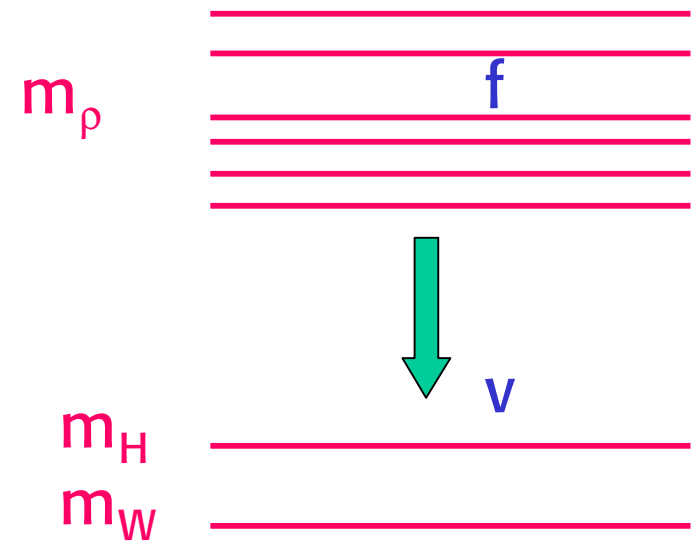
$\xi \sim 1$  similar to Technicolor

[ $\xi$  severely limited by precision EW tests  $\xi < \sim 0.2$ ]





$\xi$  severely limited by  
precision EW  
tests  $\xi < \sim 0.2$





In general composite models are more vulnerable than SUSY from EW precision tests

(for SUSY, Higgs couplings are more effective than EWPT)

No clear UV completion, no connection to GUTs

Composite models can be tested by:

- Searching for fermions of charges  $2/3$  or  $5/3$  ... that quench the bad top loop behaviour
- Measurable deviations can be expected in channels  $pp \rightarrow t\bar{t}h$ ,  $gg \rightarrow hh$  and in decays  $h \rightarrow \mu\mu$ ,  $h \rightarrow Z\gamma$

Some recent papers:

Azatov et al '13

Contino et al '13

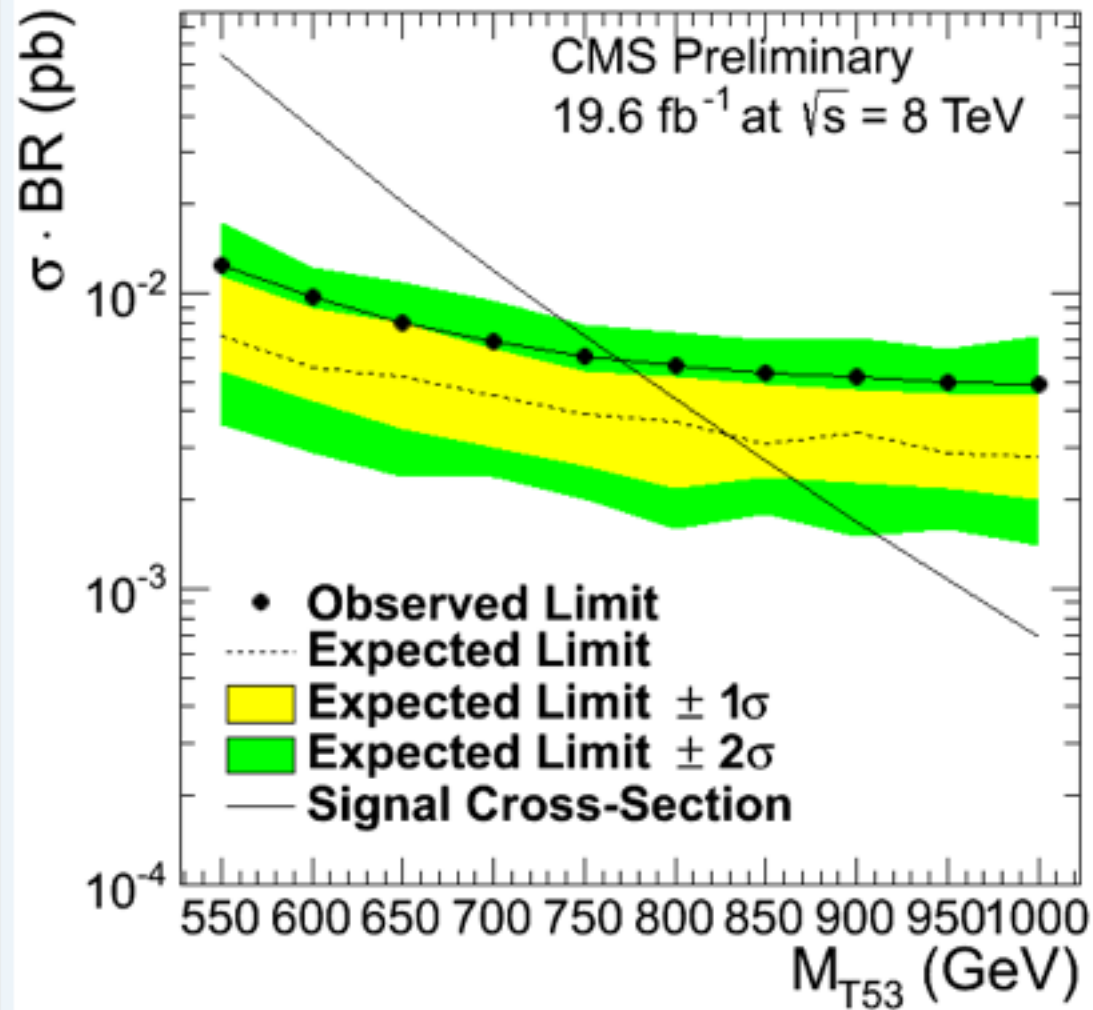
Jenkins et al '13

Grojean et al '13.....



## Searches for $t$ partners

In composite models the top loop bad behaviour is quenched by a new fermion



Expected and observed 95% C.L. limits on the  $T_{5/3}$  production cross section. The 1-sigma and 2-sigma combined statistical and systematic expected variation is shown as a yellow (light) and green (dark) band, respectively.

⊕ A  $5/3$  charged fermion cannot mix and is not pushed up

## At the other extreme: the anthropic multiverse

- The empirical value of the cosmological constant  $\Lambda_{\text{cosmo}}$  poses a tremendous, unsolved naturalness problem

While natural extensions of the SM exist, no natural explanation of the value of  $\Lambda_{\text{cosmo}}$  is known

- Yet the value of  $\Lambda_{\text{cosmo}}$  is close to the Weinberg upper bound for galaxy formation
- Possibly our Universe is just one of infinitely many continuously created from the vacuum by quantum fluctuations (multiverse)
- Different physics in different Universes according to the multitude of string theory solutions ( $\sim 10^{500}$ )

Perhaps we live in a very unlikely Universe but one that allows our existence



Given the stubborn refusal of the SM to step aside many have turned to the anthropic philosophy also for the SM

Actually applying the anthropic principle to the SM hierarchy problem is not terribly convincing

After all, we can find plenty of models that reduce the fine tuning from  $10^{14}$  to  $10^2$ . And the added ingredients do not appear to make our existence more impossible. So why make our Universe so terribly unlikely?

But there is some similarity

$\Lambda_{\text{cosmo}}$  -  $\rightarrow$  a vacuum energy density in all points of space

$v$   $\rightarrow$  a vacuum expectation value in all points of space

With larger  $\Lambda_{\text{cosmo}}$  no galaxies, with larger  $v$  no nuclear physics

⊕ The anthropic way is now being kept in mind as a possibility

# A revival of models that accept some fine tuning

## Examples:

### Split SUSY

heavy scalars, light  
gauginos and higgsinos  
(DM and Unification)

### High scale SUSY

all sparticles heavy  
 $\lambda h^4$  fixed by gauge

### Non SUSY GUT's

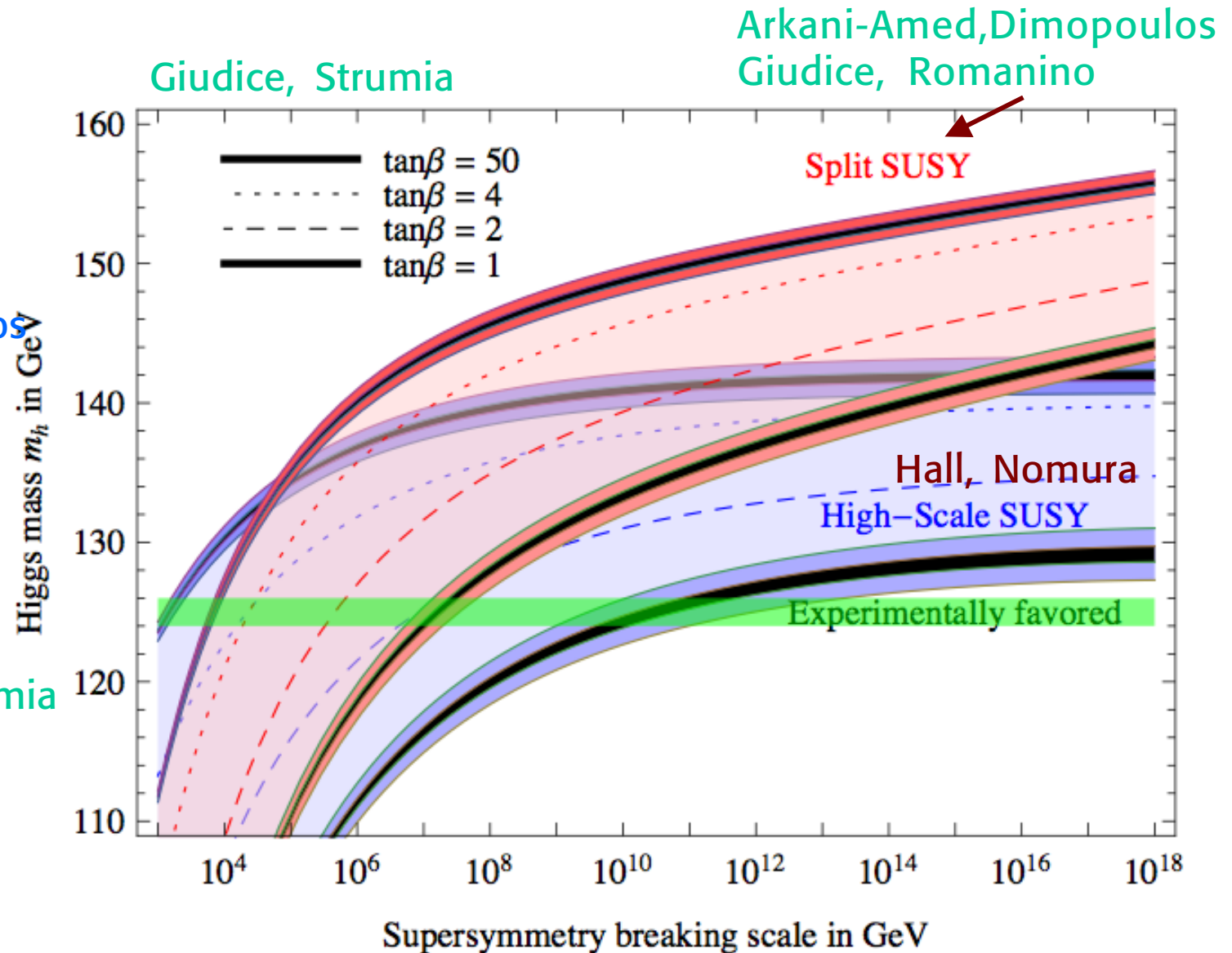
Unificaxion

Giudice, Rattazzi, Strumia

Non SUSY SO(10)

GA, Meloni

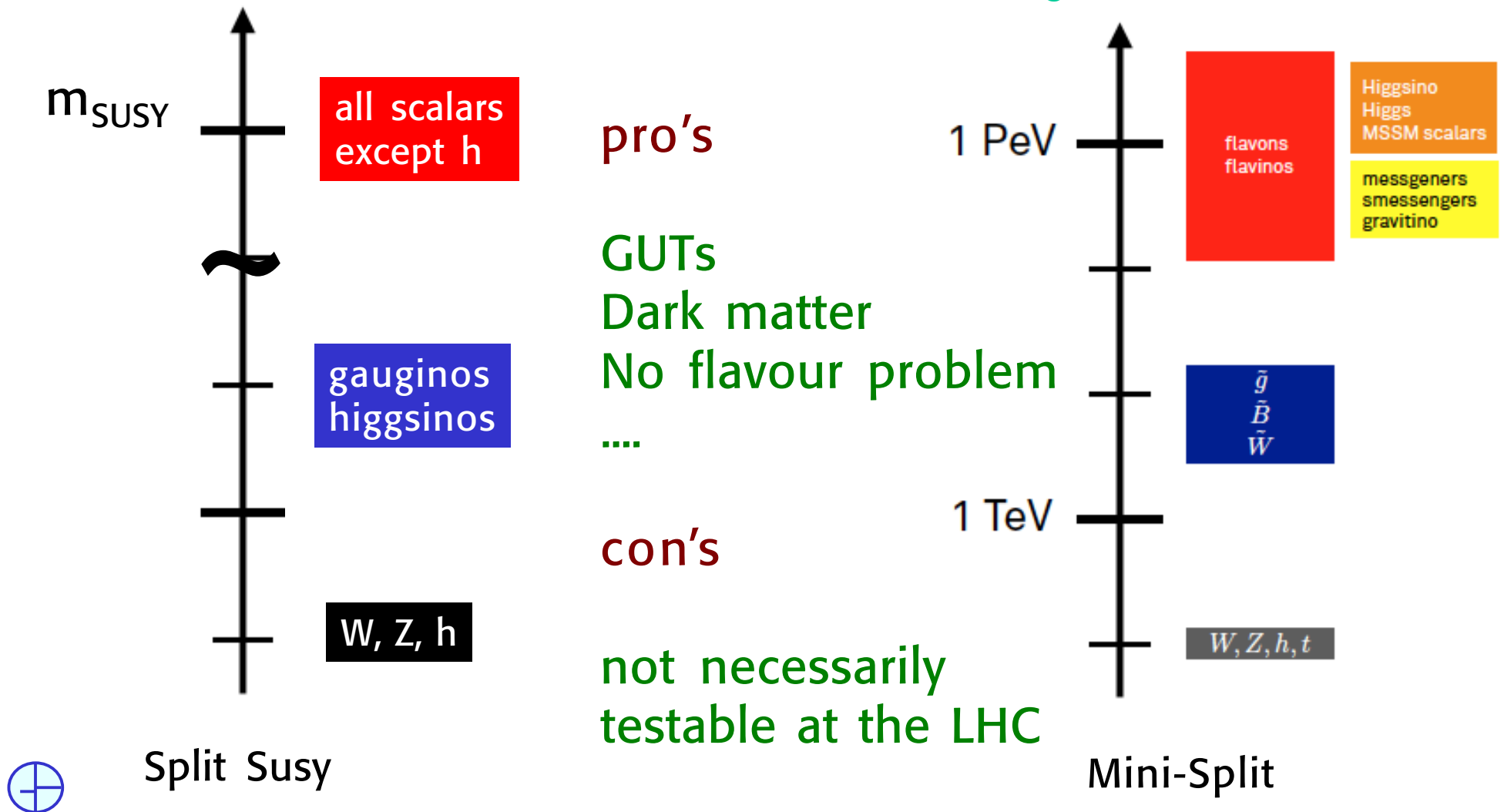
.....



# Split or Mini-Split SUSY could be a compromise: accept fine tuning but up to a point

Arkani-Hamed, Dimopoulos '04  
Giudice, Romanino '04

Baumgart, Stolarski, Zorawski '14



# Remove the FT problem: a drastic conjecture

No new thresholds between  $m_W$  and  $M_{Pl}$ ?

Shaposhnikov '07--->

And hope that gravity will somehow fix the problem of fine tuning related to the  $M_{Pl}$  threshold (with many thresholds it would be more difficult for gravity to arrange the fine tuning)

Giudice EPS'13

For this, one needs to solve all problems like Dark Matter, neutrino masses, baryogenesis.... at the EW scale

In particular no GUTs, no heavy RH neutrinos, no WIMPs..... below  $M_{Pl}$ . **A big loss!!**



# The $\nu$ MSM

Shaposhnikov et al

There are 3 RH  $\nu$ 's:  $N_1, N_2, N_3$  and the see-saw mechanism

But the  $N_i$  masses are all below the EW scale

Actually  $N_1 \sim \mathcal{O}(1-10)$  keV, and  $N_{2,3} \sim \text{GeV}$  with eV splitting

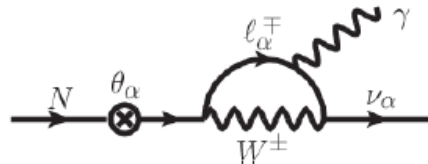
Very small Yukawa couplings are assumed to explain the small active  $\nu$  masses

$$m_\nu = \frac{y_\nu^2 v^2}{M_N}$$

The phenomenology of  $\nu$  oscillations can be reproduced

$N_1$  can explain (warm) DM

$N_{2,3}$  can explain the Baryon Asymmetry in the Universe



$N_1$  decay produces a distinct X-ray line

$$N_1 \rightarrow \nu + \gamma \quad (E_\gamma = m_{N_1}/2)$$

$$\Gamma_\gamma(m_s, \theta) = 1.38 \times 10^{-29} \text{ s}^{-1} \left( \frac{\sin^2 2\theta}{10^{-7}} \right) \left( \frac{m_s}{1 \text{ keV}} \right)^5$$

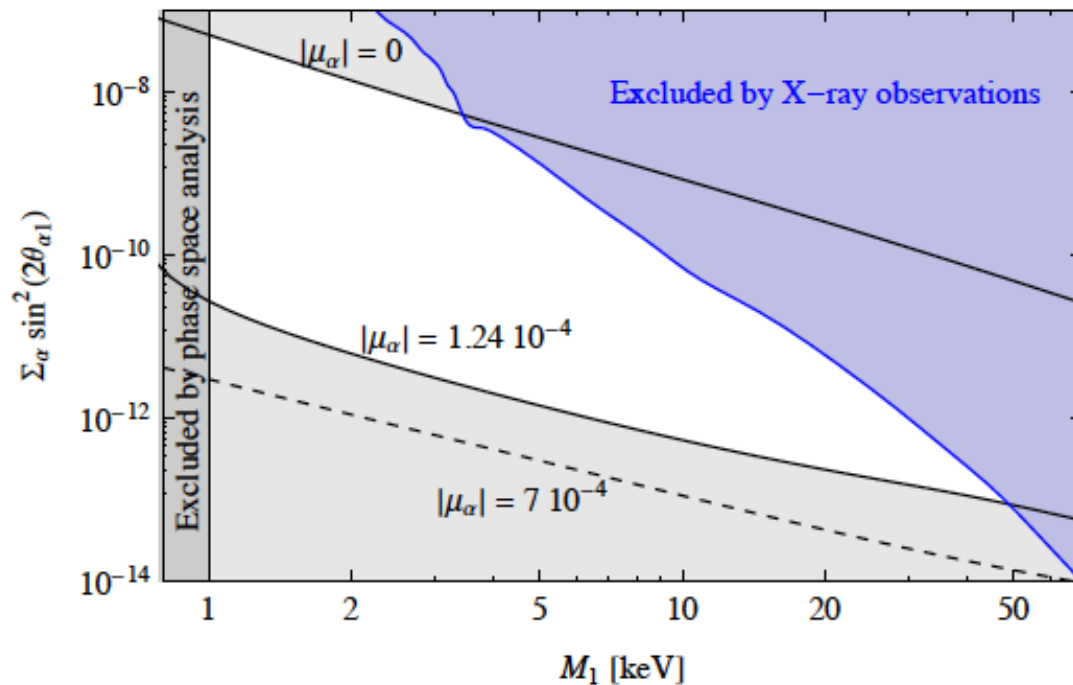
$N_{2,3}$  could be detected by dedicated accelerator experiments (eg in B decays,  $\text{Br} \sim 10^{-10}$ )

A LOI for the CERN SPS has been presented



Bonivento et al, [ArXiv:1310.1762](https://arxiv.org/abs/1310.1762)





Canetti et al '12

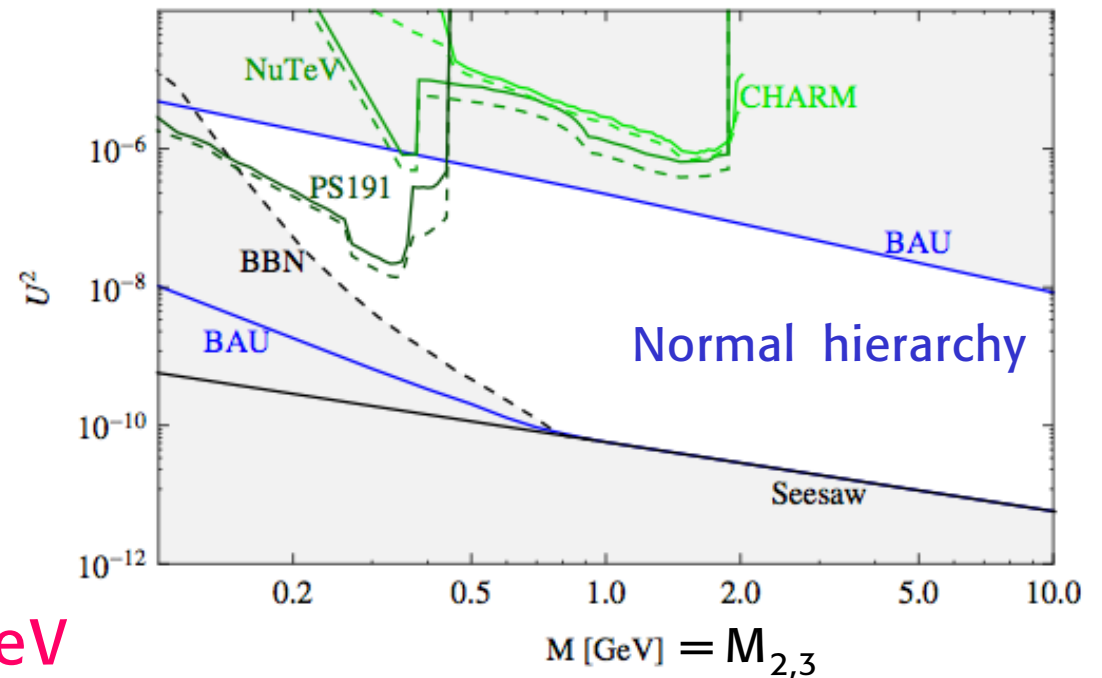
The claim is that all constraints can be satisfied

For DM one needs  $1 < M_1 < \sim 100$  keV

keV

No explanation of the mass splitting

GeV



# A $\sim 7$ keV sterile $N_1$ ?

ArXiv:1402.2301

## DETECTION OF AN UNIDENTIFIED EMISSION LINE IN THE STACKED X-RAY SPECTRUM OF GALAXY CLUSTERS

ESRA BULBUL<sup>1,2</sup>, MAXIM MARKEVITCH<sup>2</sup>, ADAM FOSTER<sup>1</sup>, RANDALL K. SMITH<sup>1</sup>, MICHAEL LOEWENSTEIN<sup>2</sup>, AND SCOTT W. RANDALL<sup>1</sup>

<sup>1</sup> Harvard-Smithsonian Center for Astrophysics, 60 Garden Street, Cambridge, MA 02138.

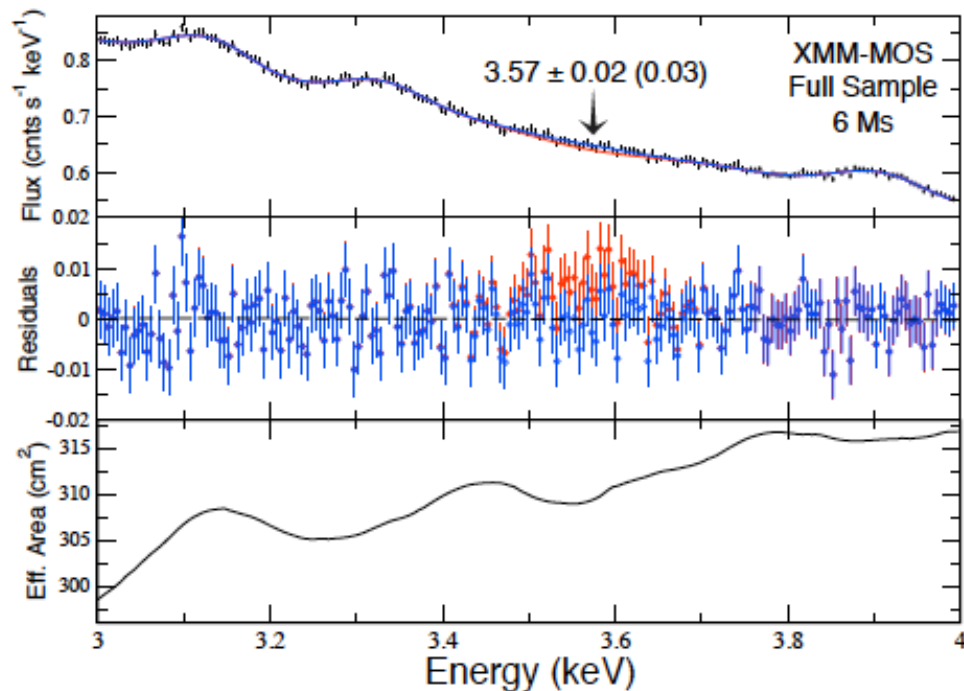
<sup>2</sup> NASA Goddard Space Flight Center, Greenbelt, MD, USA.

Submitted to *ApJ*, 2014 February 10

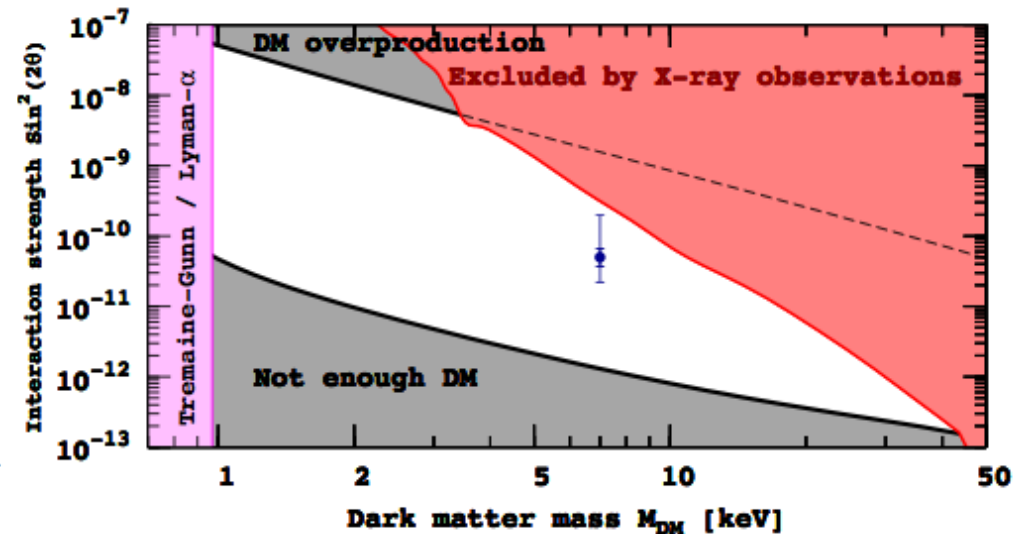
XMM-Newton X-ray observatory

### ABSTRACT

We detect a weak unidentified emission line at  $E = (3.55 - 3.57) \pm 0.03$  keV in a stacked XMM spectrum of 73 galaxy clusters spanning a redshift range  $0.01 - 0.35$ . MOS and PN observations



Independent analysis by Boyarski et al  
ArXiv:1402.4119



Confirmation from Chandra, Suzaku and eventually, Astro-H needed

The “evidence” for the 3.5 keV line has been questioned

## Questioning a 3.5 keV dark matter emission line

Signe Riemer-Sørensen<sup>1,\*</sup>

<sup>1</sup>*Institute of Theoretical Astrophysics, University of Oslo, PO 1029 Blindern, 0315 Oslo, Norway*

## Dark matter searches going bananas: the contribution of Potassium (and Chlorine) to the 3.5 keV line

Tesla Jeltema<sup>1\*</sup> and Stefano Profumo<sup>1†</sup>

<sup>1</sup>*Department of Physics and Santa Cruz Institute for Particle Physics University of California, Santa Cruz, CA 95064, USA*

but also supported: the debate is open

## Checking the dark matter origin of 3.53 keV line with the Milky Way center

A. Boyarsky<sup>1</sup>, J. Franse<sup>1,2</sup>, D. Iakubovskiy<sup>3</sup>, and O. Ruchayskiy<sup>4</sup>

<sup>1</sup>*Instituut-Lorentz for Theoretical Physics, Universiteit Leiden, Niels Bohrweg 2, Leiden, The Netherlands*

<sup>2</sup>*Leiden Observatory, Leiden University, Niels Bohrweg 2, Leiden, The Netherlands*

<sup>3</sup>*Bogolyubov Institute of Theoretical Physics, Metrologichna Str. 14-b, 03680, Kyiv, Ukraine*

<sup>4</sup>*Ecole Polytechnique Fédérale de Lausanne, FSB/ITP/LPPC, BSP, CH-1015, Lausanne, Switzerland*

## Summary

- Higgs, minimal, elementary, standard
- No new physics. Naive naturalness failed  
We expected complexity, we found simplicity
- The SM could hold up to  $M_{\text{Pl}}$   
Minimal completions of SM  
Majorana  $\nu$ 's, see-saw, leptogenesis ....
- Today the most crucial problem is Dark Matter  
WIMPS, Axions, keV  $\nu$ 's....
- Different theoretical avenues  
Insist on as minimal as possible Fine Tuning (FT)  
Stealth SUSY, nearby compositeness.....  
Accept some FT  
e.g. Split-SUSY  
Total acceptance of FT: the Anthropic metaphysics
- Denial of FT: the no-threshold philosophy  
the  $\nu$ MSM, scale invariant theories  
price: no GUTs, no heavy  $\nu_R$  ....  
But BICEP2 now makes the GUT scale to reappear!

