

# What is the world made of?

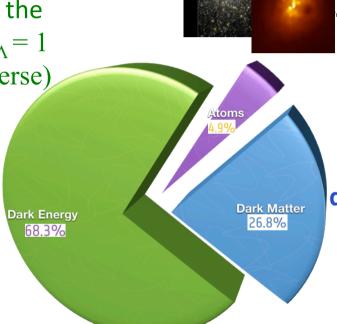
Mainly geometrical evidence:

 $\Lambda \sim O(H_0^2), H_0 \sim 10^{-42} \, \text{GeV}$ 

... dark energy is inferred from the

'cosmic sum rule':  $\Omega_{\rm m} + \Omega_{\rm k} + \Omega_{\Lambda} = 1$ 

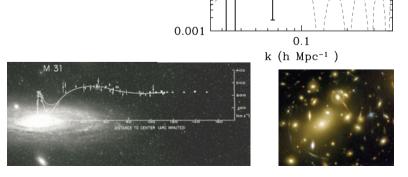
(assuming a homogeneous universe)



Baryons (but *no* anti-baryons)

Both geometrical *and* dynamical evidence for **dark matter** (*if* GR is valid)

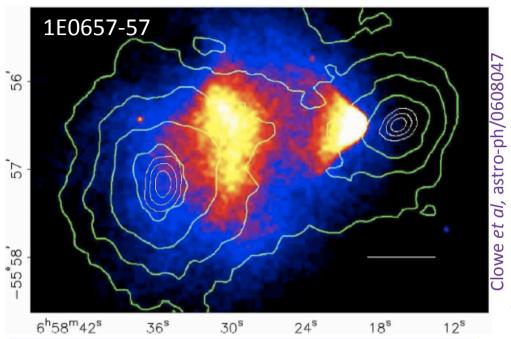
Both the baryon asymmetry and dark matter require new physics beyond the Standard  $SU(3)_{c}xSU(2)_{L}xU(1)_{Y}$  Model ... dark energy is even more mysterious (but still lacks compelling dynamical evidence)



 $k^{3} P(k)/2\pi^{2}$ 

0.01

#### What can astrophysics tell us about dark matter interactions?



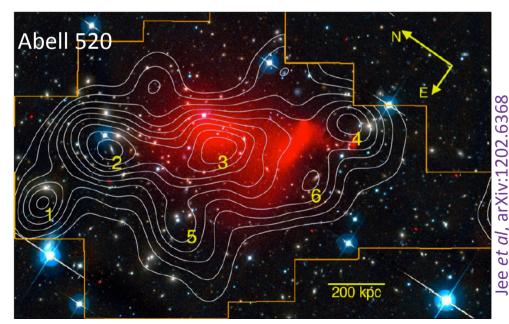
The 'Bullet Cluster' is often cited as evidence that dark matter is *collisionless* ... in actual fact it sets a rather *weak* limit on self-interactions:  $\sigma \lesssim 2 \text{x} 10^{-24} \text{ cm}^2/\text{GeV}$ 

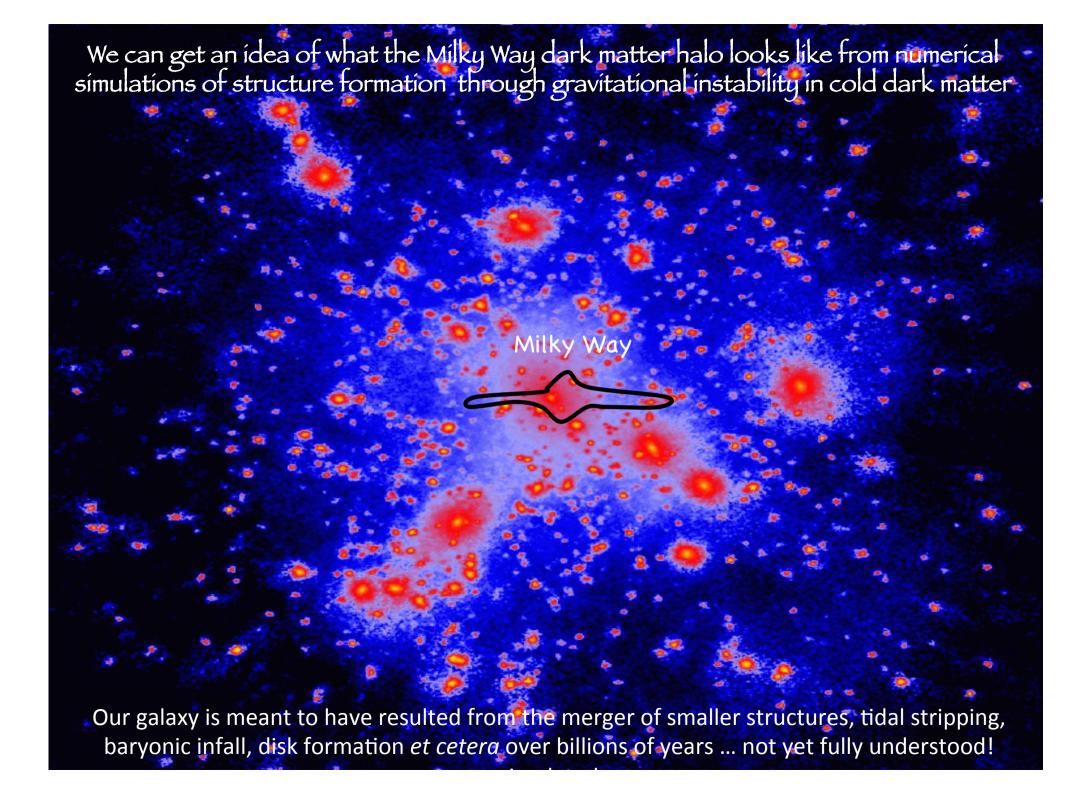
Moreover it poses a *challenge* for  $\Lambda$ CDM cosmology: why is the relative velocity so high (>3000 km/s on a scale of 5 Mpc)?

Nine other colliding clusters have been found ... the odds are *tiny* in a Gaussian density field!

But in Abell 520, the DM concentration is partly *coincident* with the X-ray emitting gas implying that DM is *self-interacting* with:  $\sigma \approx 2 \times 10^{-24} \text{ cm}^2/\text{GeV}$  (Jee *et al*, 1401.3356)

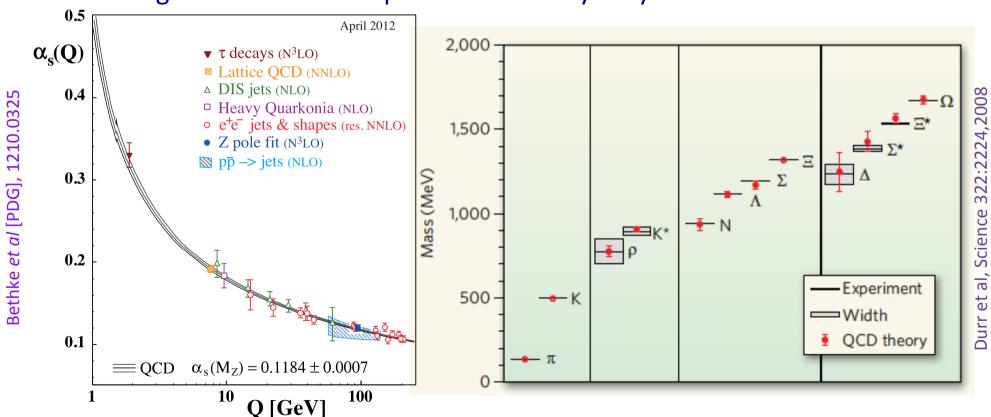
This result is contested ... in any case the separation between DM and galaxies will be *time-dependent* and sensitive to whether the self-interactions are contact or long-range (Frandsen *et al*, 1308.3419) ... so data from gravitational lensing can in principle discriminate between DM particle candidates





Mass scale	Particle	Symmetry/ Quantum #	Stability	Production	Abundance
$\Lambda_{ ext{QCD}}$	Nucleons	Baryon number	$\tau > 10^{33}$ yr	'freeze-out' from thermal equilibrium	$\Omega_{ m B}$ $\sim$ 10 <sup>-10</sup> $cf$ . observed $\Omega_{ m B}$ $\sim$ 0.05

We have a good theoretical explanation for why baryons are massive and stable



We understand the dynamics of QCD ... and can calculate the mass spectrum

#### Nevertheless we get the cosmology of baryons badly wrong!

$$\dot{n} + 3Hn = -\langle \sigma v \rangle (n^2 - n_{\rm T}^2)$$

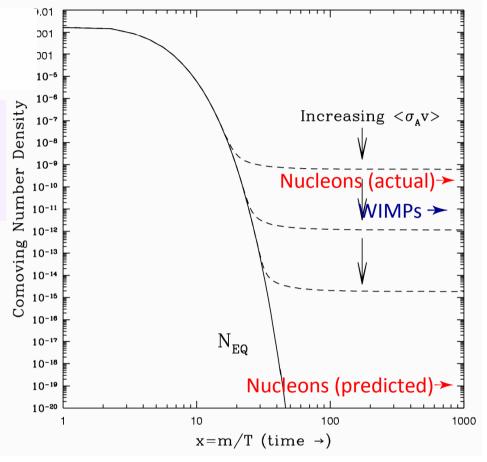
Chemical equilibrium is maintained as long as the annihilation rate exceeds the Hubble expansion rate

'Freeze-out' occurs when annihilation rate:

$$\Gamma = n\sigma v \sim m_N^{3/2} T^{3/2} e^{-m_N/T} \frac{1}{m_\pi^2}$$

becomes comparable to the expansion rate

$$H \sim rac{\sqrt{g}T^2}{M_{
m P}}$$
 where  $g$   $\sim$  # relativistic species



i.e. 'freeze-out' occurs at 
$$T\sim m_N/45$$
, with:  $\frac{n_N}{n_\gamma}=\frac{n_{\bar{N}}}{n_\gamma}\sim 10^{-19}$ 

However the observed ratio is  $6x10^9$  times bigger for baryons, and there seem to be **no antibaryons**, so we must invoke an initial baryon asymmetry:  $\frac{n_B - n_{\bar{B}}}{10^{-9}} \sim 10^{-9}$ 

Why do we not call this the 'baryon disaster'? (cf. 'WIMP miracle'!)  $^{n_B+n_{ar{B}}}$ 

## To make the baryon asymmetry requires a lot of new physics:

- ➢ B-number violation
  - > CP violation
- Departure for thermal equilibrium

The SM does allow *B*-number violation (through non-perturbative - sphaleron-mediated – processes) ... but *CP*-violation is *too weak* and  $SU(2)_L \times U(1)_Y$  breaking is a 'cross-over' (*not* out-of-equilibrium)

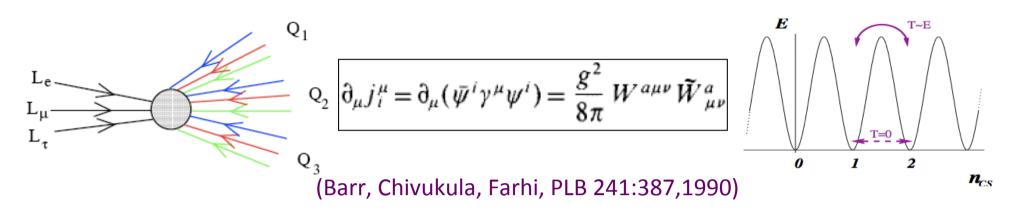
Hence the generation of the observed matter-antimatter asymmetry requires new BSM physics (could be related to neutrino mass if this arises from violation of lepton number → leptogenesis)

'See-saw': 
$$\mathcal{L} = \mathcal{L}_{SM} + \lambda_{\alpha J}^* \overline{\ell}_{\alpha} \cdot HN_J - \frac{1}{2} \overline{N_J} M_J N_J^c \qquad \lambda M^{-1} \lambda^{\mathrm{T}} \langle H^0 \rangle^2 = [m_{\nu}]$$

$$\nu_{L\alpha} \xrightarrow{m_D^{\alpha A}} \xrightarrow{M_A} \xrightarrow{m_D^{\beta A}} \nu_{L\beta}$$

$$\Delta m_{atm}^2 = m_3^2 - m_2^2 \simeq 2.6 \times 10^{-3} \mathrm{eV}^2 \qquad \Delta m_{\odot}^2 = m_2^2 - m_1^2 \simeq 7.9 \times 10^{-5} \mathrm{eV}^2$$

# Asymmetric baryonic matter

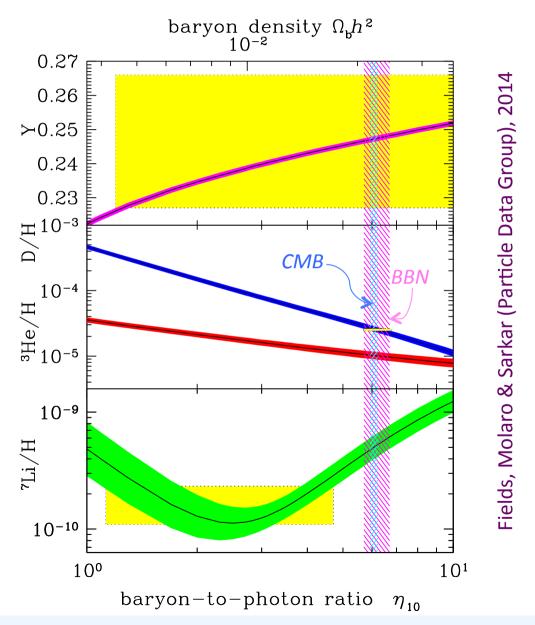


Any primordial lepton asymmetry (e.g. from *out*-of-equilibrium decays of the right-handed *N* ) would be redistributed by *B+L* violating processes in the SM (which *conserve B-L*) amongst *all* fermions – in particular **baryons** - which couple to the electroweak anomaly

Although **leptogenesis** may never be directly testable, evidence for a *Majorana* neutrino mass from observation of neutrino*less*  $\beta\beta$ -decay would provide powerful support for the idea

... in any case we accept that the only kind of matter which we are certain *exists*, originated *non-thermally* in the early universe

Although vastly overabundant compared to the natural expectation, baryons cannot close the universe (BBN & CMB concordance)



... the dark matter *must* therefore be mainly *non*-baryonic

The Standard  $SU(3)_c \times SU(2)_L \times U(1)_Y$  Model (viewed as an effective field theory up to some high energy cut-off scale M) accurately describes all microphysics

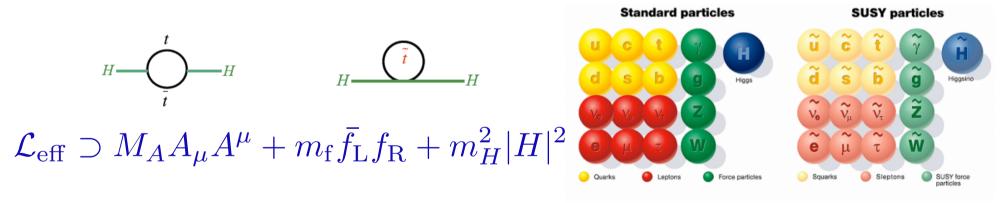
New physics beyond the SM  $\Rightarrow$  non-renormalisable operators suppressed by  $M^n$  which 'decouple' as  $M \to M_P$  (... so neutrino mass is small, proton decay is slow etc)

But as M is raised, the effects of the **super-renormalisable operators** are *exacerbated* One solution for Higgs mass divergence  $\rightarrow$  'softly broken' supersymmetry at  $M \sim 1$  TeV

This provides new possibilities for baryogenesis as well as a good candidate for dark matter – the lightest supersymmetric particle (typically the neutralino  $\chi$ ), if it is cosmologically stable because of a conserved quantum number (R-parity)

This has been the target of *most* dark matter searches, whether using nuclear recoil detectors or looking for cosmic annihilation products, or missing  $E_{\rm T}$  signals at colliders

Mass scale	Particle	Symmetry/ Quantum #	Stability	Production	Abundance
$\Lambda_{ ext{QCD}}$	Nucleons	Baryon number	$ au \geq 10^{33}   ext{yr}$	'freeze-out' from thermal equilibrium Asymmetric baryogenesis	$\Omega_{ m B}$ $\sim$ 10 <sup>-10</sup> $\it cf.$ observed $\Omega_{ m B}$ $\sim$ 0.05
$oldsymbol{\Lambda_{Fermi}} \sim \ G_{\mathrm{F}}^{-1/2}$	Neutralino?	<i>R</i> -parity?	Violated? (matter parity <i>adequate</i> to ensure <i>p</i> stability)	'freeze-out' from thermal equilibrium	$\Omega_{\rm LSP}\!\sim 0.3$



For (softly broken) **supersymmetry** we have the 'WIMP miracle':

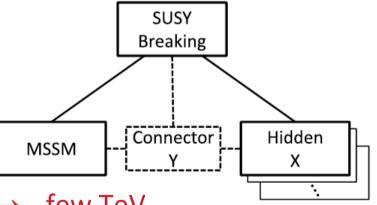
$$\Omega_{\chi} h^2 \simeq \frac{3 \times 10^{-27} \text{cm}^{-3} \text{s}^{-1}}{\langle \sigma_{\text{ann}} v \rangle_{T=T_f}} \simeq 0.1 \text{ , since } \langle \sigma_{\text{ann}} v \rangle \sim \frac{g_{\chi}^4}{16\pi^2 m_{\chi}^2} \approx 3 \times 10^{-26} \text{cm}^3 \text{s}^{-1}$$

But why should a thermal relic have an abundance comparable to non thermal relic baryons?

Mass scale	Particle	Symmetry/ Quantum #	Stability	Production	Abundance
$\Lambda_{ ext{QCD}}$	Nucleons	Baryon number	$\tau > 10^{33} \text{ yr}$	'freeze-out' from thermal equilibrium Asymmetric baryogenesis	$\Omega_{ m B}{\sim}10^{\text{-}10}$ cf. observed $\Omega_{ m B}{\sim}0.05$
$\Lambda_{ m Fermi}$ $\sim$ $G_{ m F}^{-1/2}$	Neutralino?	<i>R</i> -parity?	Violated? (matter parity <i>adequate</i> for <i>p</i> stability)	'freeze-out' from thermal equilibrium	$\Omega_{\rm LSP} \sim 0.3$

(GMSB) Hidden sector matter also provides the 'WIMPless miracle' (Feng & Kumar, 0803.4196)

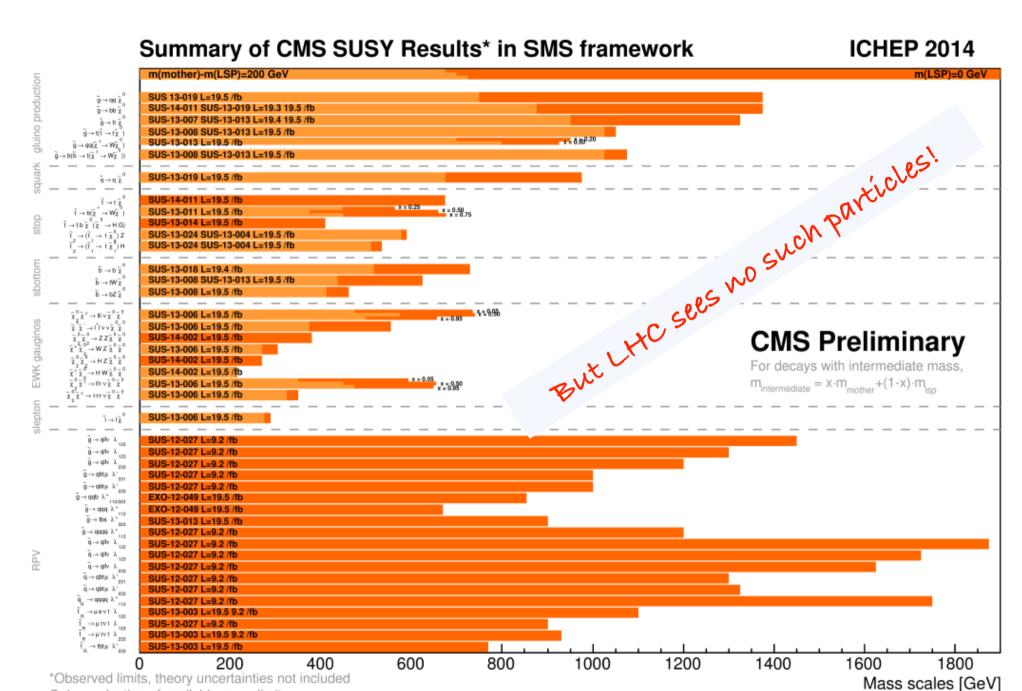
... because: 
$$g_h^2/m_h \sim g_\chi^2/m_\chi \sim F/16\pi^2 M$$



Such dark matter can have *any* mass:  $\sim 0.1$  GeV  $\rightarrow \sim$  few TeV

$$\Omega_{\chi} h^2 \simeq \frac{3 \times 10^{-27} \text{cm}^{-3} \text{s}^{-1}}{\langle \sigma_{\text{ann}} v \rangle_{T=T_f}} \simeq 0.1 \text{ , since } \langle \sigma_{\text{ann}} v \rangle \sim \frac{g_{\chi}^4}{16\pi^2 m_{\chi}^2} \approx 3 \times 10^{-26} \text{cm}^3 \text{s}^{-1}$$

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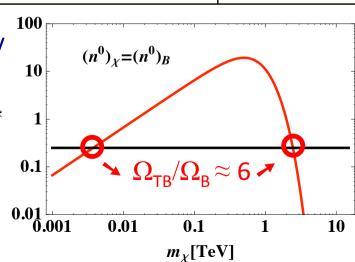


\*Observed limits, theory uncertainties not included Only a selection of available mass limits Probe \*up to\* the quoted mass limit

Mass scale	Particle	Symmetry/ Quantum #	Stability	Production	Abundance
$\Lambda_{ m QCD}$	$\Lambda_{ m QCD}$ Nucleons		τ > 10 <sup>33</sup> yr (dim-6 OK)	'Freeze-out' from thermal equilibrium Asymmetric baryogenesis (how?)	$\Omega_{ m B}{\sim}10^{\text{-}10}\text{cf.}$ observed $\Omega_{ m B}{\sim}0.05$
$\Lambda_{ m QCD}$ , ~ $6\Lambda_{ m QCD}$	Dark baryon?	$U(1)_{\mathrm{DB}}$	plausible	Asymmetric (like the observed baryons)	$\Omega_{\mathrm{DB}} \sim 0.3$
$\Lambda_{ m Fermi} \sim \ { m G_F}^{-1/2}$	Neutralino?	<i>R</i> -parity	violated?	'Freeze-out' from thermal equilibrium	$\Omega_{\rm LSP} \sim 0.3$
F	Technibaryon?	(walking) Technicolour	$ au \sim 10^{18}   ext{yr}$ $e^+$ excess?	Asymmetric (like the observed baryons)	$\Omega_{\rm TB} \sim 0.3$

A new particle can naturally *share* in the *B/L* asymmetry if it couples to the *W* ... linking dark to baryonic matter!

For example a O(TeV) mass **technibaryon** can be the dark matter (Nussinov 1985) ... another possibility is a  $\sim\!\!6$  GeV mass 'dark baryon' in a hidden sector (Gelmini, Hall & Lin 1986, Kaplan 1992):  $\Omega_\chi = (m_\chi \mathcal{N}_\chi/m_\mathrm{B} \mathcal{N}_\mathrm{B})\Omega_B$ 

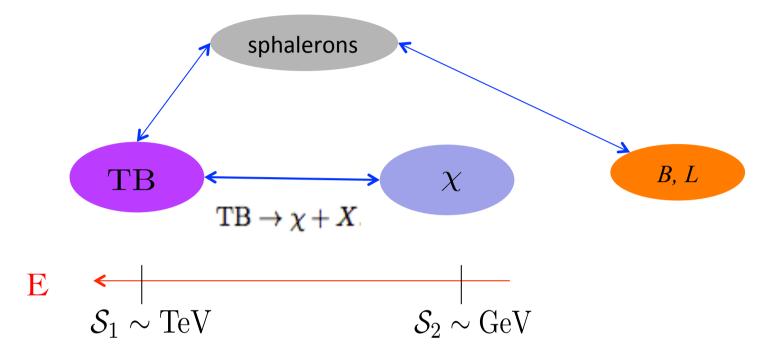


Status: ICHEP 2014

 $\int \mathcal{L} dt = (1.0 - 20.3) \text{ fb}^{-1} \qquad \sqrt{s} = 7, 8 \text{ TeV}$ 

		Model	ℓ, γ	Jets	E <sub>T</sub> miss	∫£ dt[fb	<sup>-1</sup> ] Mass limit	, at = (1.0 - 20.3) 1b	Reference
	Extra dimensions	ADD $G_{KK}+g/q$ ADD non-resonant $\ell\ell$ ADD QBH $\to \ell q$ ADD QBH ADD BH high $N_{trk}$ ADD BH high $\sum p_T$ RS1 $G_{KK} \to \ell\ell$ RS1 $G_{KK} \to \ell \ell$ Bulk RS $G_{KK} \to ZZ \to \ell\ell qq$ Bulk RS $G_{KK} \to HH \to b\bar{b}b\bar{b}$ Bulk RS $g_{KK} \to t\bar{t}$ $S^1/Z_2$ ED UED	$\begin{array}{c} -\\ 2e, \mu\\ 1e, \mu\\ -\\ 2\mu(SS)\\ \geq 1e, \mu\\ 2e, \mu\\ 2e, \mu\\ -\\ 1e, \mu\\ 2e, \mu\\ 2e, \mu\\ \end{array}$	1-2 j - 1 j 2 j - ≥ 2 j - 2 j / 1 J 4 b ≥ 1 b, ≥ 1 J	Yes	4.7 20.3 20.3 20.3 20.3 20.3 20.3 4.7 20.3 19.5 14.3 5.0 4.8	$\begin{array}{cccccccccccccccccccccccccccccccccccc$	n = 2 n = 3 HLZ n = 6 n = 6 $n = 6$ , $M_D = 1.5$ TeV, non-rot BH $n = 6$ , $M_D = 1.5$ TeV, non-rot BH n = 6, $n = 1.5$ TeV, non-rot BH	1210.4491 ATLAS-CONF-2014-030 1311.2006 to be submitted to PRD 1308.4075 1405.4254 1405.4123 1208.2880 ATLAS-CONF-2014-039 ATLAS-CONF-2014-005 ATLAS-CONF-2013-052 1209.2535 ATLAS-CONF-2012-072
	Gauge bosons	$\begin{array}{l} \operatorname{SSM} Z' \to \ell\ell \\ \operatorname{SSM} Z' \to \tau\tau \\ \operatorname{SSM} W' \to \ell\nu \\ \operatorname{EGM} W' \to WZ \to \ell\nu  \ell'  \ell' \\ \operatorname{EGM} W' \to WZ \to qq\ell\ell \\ \operatorname{LRSM} W'_R \to t\overline{b} \\ \operatorname{LRSM} W'_R \to t\overline{b} \end{array}$	2 e, μ 2 τ 1 e, μ 3 e, μ 2 e, μ 1 e, μ 0 e, μ	- - - 2 j / 1 J 2 b, 0-1 j ≥ 1 b, 1 J		20.3 19.5 20.3 20.3 20.3 14.3 20.3	GKK mass  gKK mass  gKK mass  2.0 TeV  M <sub>KK</sub> ≈ R <sup>-1</sup> Compact scale R <sup>-1</sup> 2.9 TeV  2.9 TeV  2.9 TeV  2.9 TeV  2.1 TeV  W' mass  1.52 TeV  W' mass  1.59 TeV  M.  M.  M.  LQ mass  LQ		1405.4123 ATLAS-CONF-2013-066 ATLAS-CONF-2014-017 1406.4456 ATLAS-CONF-2014-039 ATLAS-CONF-2013-050 to be submitted to EPJC
	Ö	Cl qqqq Cl qqℓℓ Cl uutt	– 2 e, μ 2 e, μ (SS)	2 j - ) ≥ 1 b, ≥ 1	– j Yes	4.8 20.3 14.3	7.6 TeV A A 3.3 TeV	$\eta = +1$ 21.6 TeV $\eta_{LL} = -1$ $ C  = 1$	1210.1718 ATLAS-CONF-2014-030 ATLAS-CONF-2013-051
	DM	EFT D5 operator (Dirac) EFT D9 operator (Dirac)	0 e, μ 0 e, μ	1-2 j 1 J, ≤ 1 j	Yes Yes	10.5 20.3	M. 2.4 TeV	at 90% CL for $m(\chi) < 80$ GeV at 90% CL for $m(\chi) < 100$ GeV	ATLAS-CONF-2012-147 1309.4017
	70	Scalar LQ 1 <sup>st</sup> gen Scalar LQ 2 <sup>nd</sup> gen Scalar LQ 3 <sup>rd</sup> gen	2 e 2 μ 1 e, μ, 1 τ	≥ 2 j ≥ 2 j 1 b, 1 j	- - -	1.0 1.0 4.7	LQ mass LQ mass LQ mass J34 GeV	$\beta = 1$ $\beta = 1$ $\beta = 1$	1112.4828 1203.3172 1303.0526
Поэм	quarks	Vector-like quark $\overline{TT} \to Ht + X$ Vector-like quark $\overline{TT} \to Wb + X$ Vector-like quark $\overline{TT} \to Zt + X$ Vector-like quark $BB \to Zb + X$ Vector-like quark $BB \to Wt + X$	2/≥3 e, μ 2/≥3 e, μ	≥2/≥1 b	j Yes - -	14.3 14.3 20.2 20.3 14.3	790 GeV 670 GeV 735 GeV 755 GeV 720 GeV	T in (T,B) doublet isospin singlet T in (T,B) doublet B in (B,Y) doublet B in (T,B) doublet	ATLAS-CONF-2013-018 ATLAS-CONF-2013-060 ATLAS-CONF-2014-036 ATLAS-CONF-2014-036 ATLAS-CONF-2013-051
P. Carrier	fermions	Excited quark $q^* \rightarrow q\gamma$ Excited quark $q^* \rightarrow qg$ Excited quark $b^* \rightarrow Wt$ Excited lepton $\ell^* \rightarrow \ell\gamma$	1 γ - 1 or 2 e, μ 2 e, μ, 1 γ	1 j 2 j 1 b, 2 j or 1	– – Ij Yes –	20.3 20.3 4.7 13.0	q* mass         3.5 TeV           q* mass         4.09 TeV           b* mass         870 GeV           ℓ* mass         2.2 TeV	only $u^*$ and $d^*$ , $\Lambda = m(q^*)$ only $u^*$ and $d^*$ , $\Lambda = m(q^*)$ left-handed coupling $\Lambda = 2.2 \text{ TeV}$	1309.3230 to be submitted to PRD 1301.1583 1308.1364
	Other	LSTC $a_T \to W \gamma$ LRSM Majorana $\nu$ Type III Seesaw Higgs triplet $H^{\pm\pm} \to \ell\ell$ Multi-charged particles Magnetic monopoles	1 e, μ, 1 γ 2 e, μ 2 e, μ 2 e, μ (SS)	2 j - ) - - -	Yes - - - - -	20.3 2.1 5.8 4.7 4.4 2.0	aT mass         960 GeV           Nº mass         1.5 TeV           N± mass         245 GeV           H± mass         409 GeV           multi-charged particle mass         490 GeV           monopole mass         882 GeV	$m(W_R)=2$ TeV, no mixing $ V_e $ =0.055, $ V_\mu $ =0.063, $ V_r $ =0 DY production, BR( $H^{++} \rightarrow \ell\ell$ )=1 DY production, $ q =4e$ DY production, $ g =1g_D$	to be submitted to PLB 1203.5420 ATLAS-CONF-2013-019 1210.5070 1301.5272 1207.6411
			<b>√</b> s =	7 TeV	<b>√</b> S =	8 TeV	10 <sup>-1</sup> 1 1	0 Mass scale [TeV]	

## Why may we not have seen these particles yet?



- $S_1$  States (constituents) carry weak charges and are connected to sphalerons
- States are SM singlets (in a hidden sector/hidden valley) but directly connected to the  $S_1$  sector (with scale separation TeV  $\rightarrow$  GeV because of different  $\beta$ -function)
- ${
  m TB} o \chi + {
  m X}$  is in equilibrium until  $T \lesssim T_{
  m sph}$ , then  $\chi$  decouples and becomes DM The  $S_1$  states do couple to the SM (so should show up at LHC14!

NB: Such asymmetric dark matter would naturally have strong self-interactions

#### Axion dark matter

$$\mathcal{L}_{ ext{eff}} = M^4 + M^2 \Phi^2$$
 super-renormalisable  $+ (D\Phi)^2 + ar{\Psi} \not D\Psi + F^2 + ar{\Psi}\Psi\Phi + \Phi^2 + heta_{ ext{QCD}}F ilde{F}$  renormalisable  $+ ar{\Psi}\Psi\Phi\Phi + ar{\Psi}\Psi\Psi\Psi + \dots$  non-renormalisable

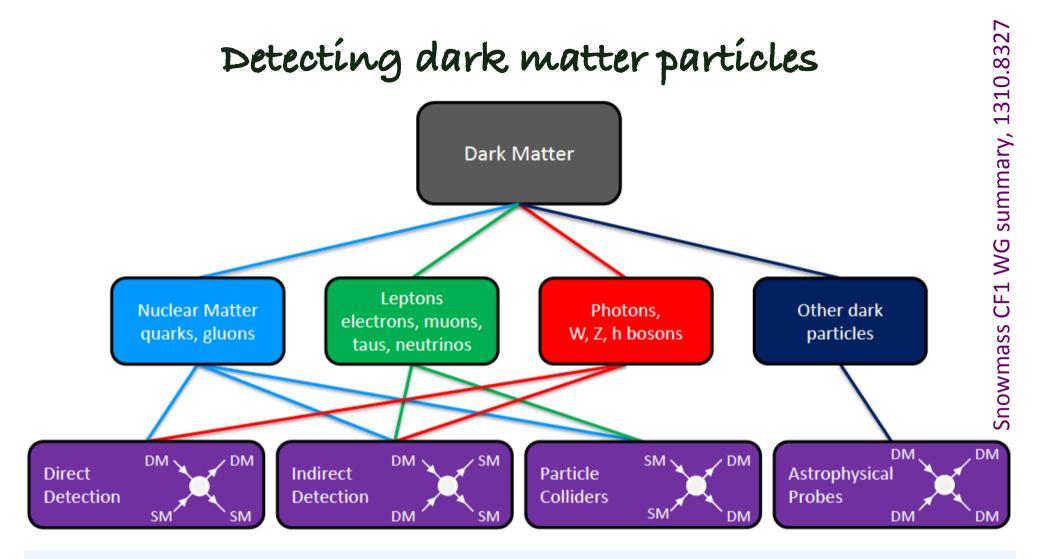
The SM admits a term which would lead to CP violation in strong interactions, hence an (unobserved) electric dipole moment for neutrons  $\rightarrow$  requires  $\theta_{QCD} < 10^{-9}$ 

To achieve this without fine-tuning,  $\theta_{\rm QCD}$  must be made a dynamical parameter, through the introduction of a new  $U(1)_{\rm Peccei-Quinn}$  symmetry which must be broken ... the resulting (pseudo) Nambu-Goldstone boson is the **axion** which acquires a small mass through its mixing with the pion (the pNGB of QCD):  $m_a = m_\pi \, (f_\pi/f_{\rm PO})$ 

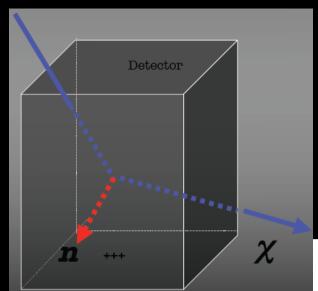
The coherent oscillations of relic axions contain energy density that behaves like CDM with  $\Omega_{\rm a}h^2\sim 10^{11}~{
m GeV}/f_{
m PO}$  ... however the natural P-Q scale is probably  $f_{
m PO}\sim 10^{18}~{
m GeV}$ 

Hence axion dark matter would typically need to be significantly diluted i.e. its relic abundance is *not* predictable (or seek anthropic explanation for why  $\theta_{QCD}$  is small?)

Mass scale	Lightest stable particle	Symmetry/ Quantum #	Stability ensured?	Production	Abundance
$\Lambda_{ extsf{QCD}}$	Nucleons	Baryon number	$ au > 10^{33}  ext{ yr}$	'Freeze-out' from equilibrium	$\Omega_{ m B}{\sim}10^{-10}~cf.$ observed
$\Lambda_{ ext{QCD'}} \ \sim 5 \Lambda_{ ext{QCD}}$	Dark baryon?	$U(1)_{DB}$	plausible	Asymmetric baryogenesis Asymmetric (like) observed baryons)	$\Omega_{ m B}$ $\sim 0.05$ $\Omega_{ m DB}$ $\sim 0.3$
$oldsymbol{\Lambda_{Fermi}} \sim G_{\scriptscriptstyle{\mathrm{F}}}^{ ext{-}1/2}$	Neutralino?	R-parity (walking)	violated?	'freeze-out' from equilibrium	$\Omega_{\mathrm{LSP}}{\sim}0.3$
$\sim O_{\mathrm{F}}$	Technibaryon?	Techni- colour	1018 yr	Asymmetric (like observed baryons)	$\Omega_{ m TB}\!\sim$ 0.3
$\Lambda_{ m hidden\ sector} \sim (\Lambda_{ m F} M_{ m P})^{1/2}$	Crypton? hidden valley?	Discrete (vely model- dependent)	τ ≥ 10.8 γγ	Varying gravitational field during inflation	$\Omega_{ m X}{\sim}0.3?$
$\Lambda_{ extsf{see-saw}} \sim \Lambda_{ extsf{Fermi}}^2/\Lambda_{ extsf{B-L}}$	Neutrinos	Lenton number	Stable <sub>.</sub>	Thermal (abundance ~ CMB photons)	$\Omega_{_{\mathrm{V}}} > 0.003$
$M_{ m string}/M_{ m Planck}$	Kajuza-Klein states?	? Peccei-	?	?	?
	Axions	Quinn	stable	Field oscillations	$\Omega_{ m a}$ » 1!



The technologies and working environments vary *extremely* widely – from nuclear recoil detectors in shielded underground laboratories, to ground-based + satellite telescopes for detecting annihilation radiation, to searches using collider detectors for missing  $E_{\text{T signals}}$  + 'monojets', to astronomical gravitational lensing studies of colliding galaxy clusters ...

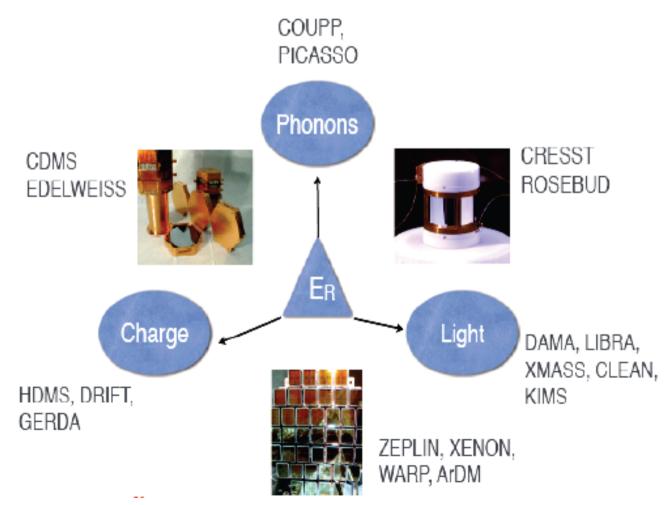


The recoil is detected via the ionization (charge), scintillation (light), and sound (phonons) → heat

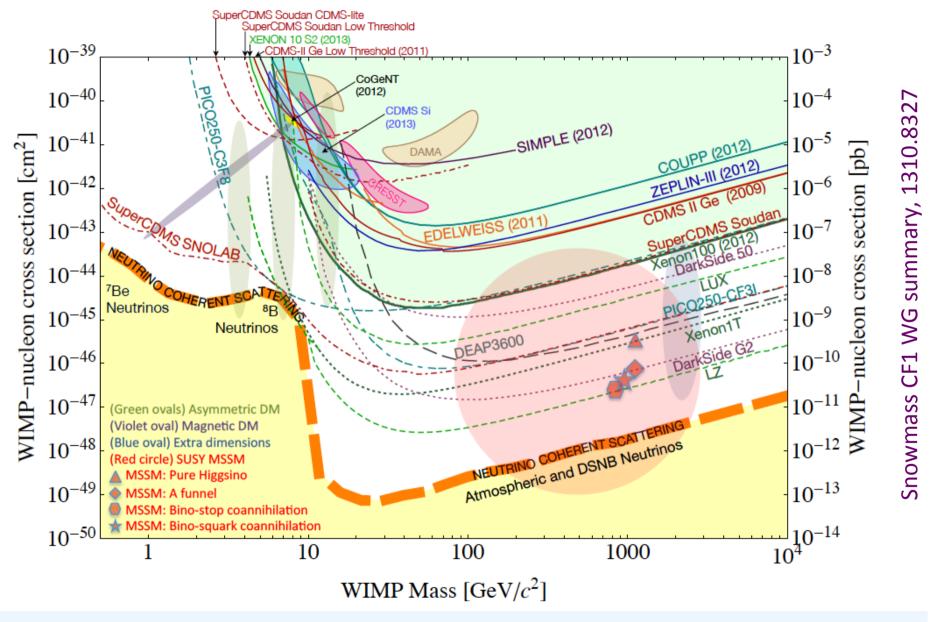
Experiments usually measure more than one channel to discriminate against the much bigger electron recoil background

(Very different techniques required to detect axions)

A passing dark matter particle orbiting in the Galaxy (at ~300 km/s) can scatter off a nucleus in an underground detector ... the expected rate is *very* low (<< 1 event/kg/yr)



#### For ~25 years there has been a world-wide race on to detect dark matter



Several claims for putative signals have apparently been ruled out by more sensitive experiments ... but are we making a fair comparison?

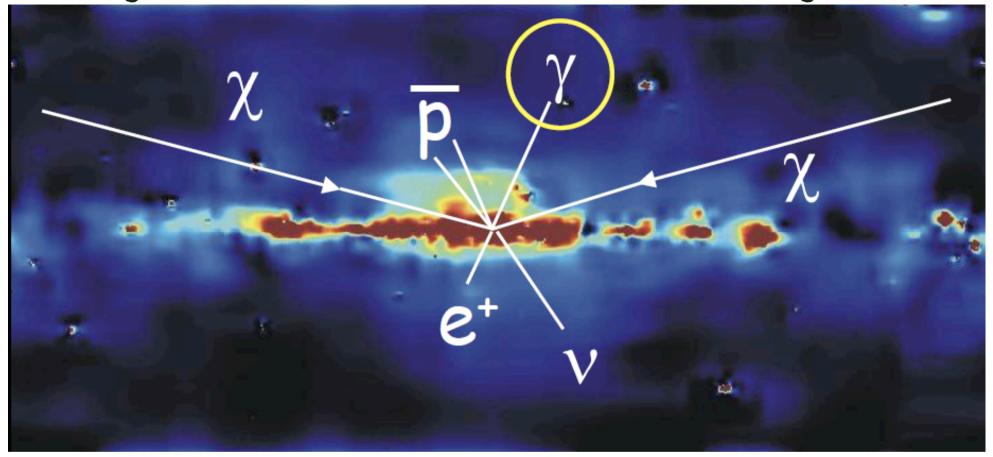
#### There are many ambiguities in interpreting the measured recoil rate:

$$\frac{\mathrm{d}R}{\mathrm{d}E_R}(E_R,t) = M_{\mathrm{tar}} \underbrace{\frac{\rho_\chi}{2\,m_\chi\mu^2}}_{\mathbf{m}_\chi\mu^2} \underbrace{\frac{(f_pZ + f_n(A-Z))^2}{f_n^2} \sigma_n F^2(E_R) \int\limits_{v_{\mathrm{Nin}}}^{\infty} \mathrm{d}^3v \; \frac{f_{\mathrm{local}}(\vec{v},t)}{v}}_{\mathrm{Nuclear physics}}$$

- ★ Dark matter may interact *differently* with neutrons and protons (Giulani, hep-ph/0504157) if the mediator is a (new) vector boson, so e.g. the events seen by CDMS-Si can be quite consistent with the upper limits set by XENON100/LUX.
- ★ Moreover different experiments are sensitive to *different* regions of the (uncertain) dark matter velocity distribution, hence apparently inconsistent results (e.g. CoGeNT and CRESST) can be reconciled by departing from the *assumed* isotropic Maxwellian form (Fox *et al*, 1011.1915, Frandsen *et al*, 1111.0292, Del Nobile *et al*, 1306.5273).
- ★ Then there are experimental uncertainties (efficiencies, energy resolution, instrumental backgrounds) ... as well as uncertainties in translating measured energies into recoil energies (channelling, quenching), uncertain nuclear form factors ...

No *single* experiment can either confirm or rule out dark matter (... also not a good strategy to look just under the supersymmetric lamp post!)

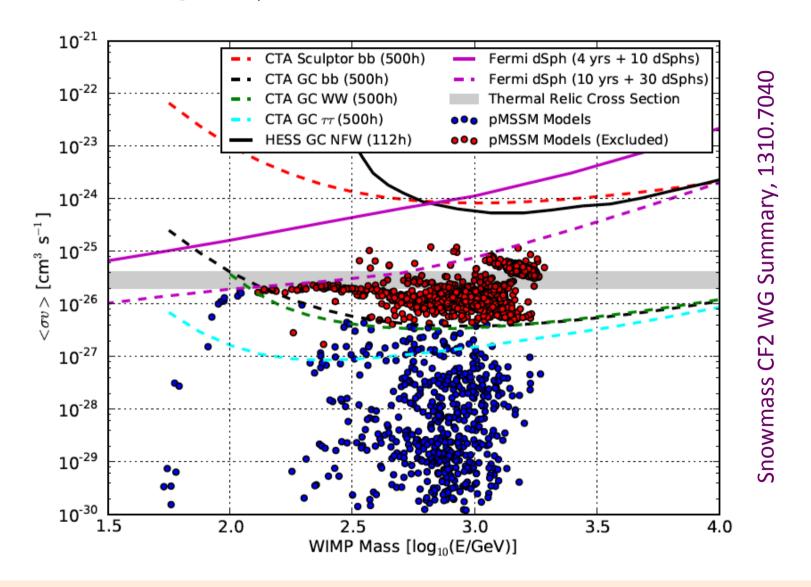
Many techniques for indirect detection ... and many claims!



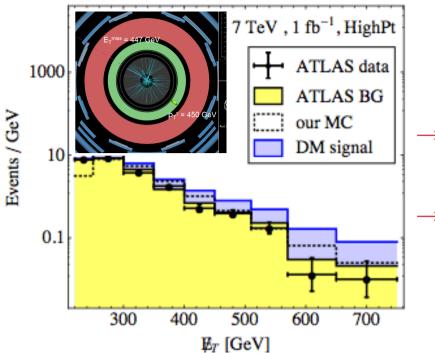
The *PAMELA/AMS-02* 'excess' ( $e^+$ ), *WMAP/Planck* 'haze' (radio), *Fermi* 'bubbles' + GC 'excess' + 130 GeV line (all  $\gamma$ -ray) ... have all been ascribed to dark matter annihilations

These probe dark matter *elsewhere* in the Galaxy so complement direct detection experiments ... but we are just beginning to understand the astrophysical backgrounds

# Prospects are good however for probing down to the expected annihilation signal for a thermal relic with Fermi and CTA



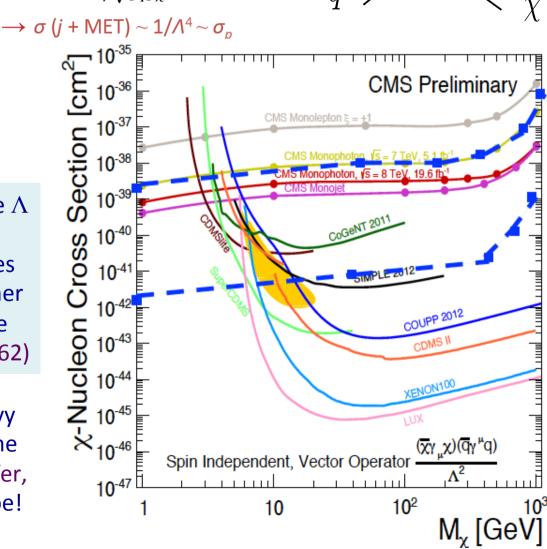
Moreover low energy extensions of IceCube (*DeepCore*, *PINGU*) will improve the sensitivity for detecting neutrinos from dark matter trapped in the Sun



However these bounds require the scale  $\Lambda$  of the effective operator to exceed ~0.7 TeV, while perturbative unitarity requires  $g_q$ ,  $g_\chi$  <  $\sqrt{4\pi}$  i.e.  $m_R$  < 2 TeV ... so for higher energy collisions *cannot* rely on effective operator description (Fox *et al*, 1203.1662)

NB: For scalar-mediated processes, heavy quark loops can significantly enhance the monojet cross-section (Haisch, Kahlhoefer, Unwin, 1208.4605) – very sensitive probe!

'Monojet' events at colliders directly measure the coupling of dark matter to SM particles, e.g.



## Summary

Experimental situation reminiscent of searches in the '80s for temperature fluctuations in the CMB ... there were clear theoretical predictions but only upper limits on detection (causing near crisis for theory)

Finally breakthrough in 1992 that transformed cosmology!

Theoretical expectations for dark matter are not as clear but there are several *complementary* experimental approaches and there has been impressive recent progress

There are bound to be false alarms but it is a reasonable expectation that the nature of dark matter will soon be determined experimentally