Extra dimensions and Braneworlds

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- Mass hierarchy, low energy SUSY and 126 GeV Higgs
- Live with the hierarchy
- Low scale strings and extra dimensions

H⁰ (Higgs Boson)

The observed signal is called a Higgs Boson in the following, although its detailed properties and in particular the role that the new particle plays in the context of electroweak symmetry breaking need to be further clarified. The signal was discovered in searches for a Standard Model (SM)-like Higgs. See the following section for mass limits obtained from those searches.

H⁰ MASS

VALUE (GeV)	DOCUMENT ID	TECN	COMMENT	
125.9±0.4 OUR AVERAGE				
$125.8 \pm 0.4 \pm 0.4$	¹ CHATRCHYAN 13J	CMS	pp, 7 and 8 TeV	
$126.0\pm0.4\pm0.4$	² AAD 12AI	ATLS	pp, 7 and 8 TeV	
• • • We do not use the followi	ing data for averages, fits,	limits,	etc. • • •	
$126.2 \pm 0.6 \pm 0.2$	³ CHATRCHYAN 13J	CMS	pp, 7 and 8 TeV	
$125.3 \pm 0.4 \pm 0.5$	⁴ CHATRCHYAN 12N	CMS	pp, 7 and 8 TeV	
HTTP://PDG.LBL.GOV	Page 1	Crea	ted: 7/31/2013	

particle listing

summary tables

H^{0} Mass $m = 125.9 \pm 0.4$ GeV

H⁰ signal strengths in different channels [n]

Combined Final States = 1.07 ± 0.26 (S = 1.4) *W W*^{*} Final State = 0.88 ± 0.33 (S = 1.1) *Z Z*^{*} Final State = $0.89^{+0.30}_{-0.25}$ $\gamma\gamma$ Final State = 1.65 ± 0.33 *b b* Final State = $0.5^{+0.8}_{-0.7}$ $\tau^+\tau^-$ Final State = 0.1 ± 0.7

Couplings of the new boson vs SM



exclusion : spin 2 and pseudoscalar at $\gtrsim 95\%$ CL

Agreement with Standard Model expectation at $\sim 2\,\sigma$

Big success of the Standard Theory of particle physics

Quantum Field Theory + Gauge Invariance

Very accurate description of Particle Physics at present energies

$$\mathcal{L}_{SM} = -\frac{1}{2} \text{tr} F_{\mu\nu}^{2} + \bar{\psi} \mathcal{D} \psi + \bar{\psi} Y H \psi - |DH|^{2} - V(H)$$

$$\overrightarrow{} \qquad \uparrow \qquad \uparrow$$
Three sectors: Forces Matter Higgs

Forces: uniquely determined by geometry (given the gauge group)

- Matter: fixed by the representations (discrete arbitrariness)
- Scalar Higgs sector: arbitrary (no guidance)

experiment \Rightarrow minimal: $V(H) = -\mu^2 |H|^2 + \lambda (|H|^2)^2$ Why?

 μ : the only relevant operator creating the hierarchy problem

Beyond the Standard Theory of Particle Physics: driven by the mass hierarchy problem

- Higgs mass: very sensitive to high energy physics
- quantum corrections: $\delta m_H \sim \delta M_W$ of order of UV cutoff Λ
- or in general the scale of new physics/massive particles
- stability requires adjustment of parameters at very high accuracy to keep the physical mass $(m_{H}^{tree})^2 + \delta m_{H}^2$ at the weak scale
- $\Lambda = M_{GUT}$ or $M_P \Rightarrow$ fine tuning at 28-32 decimal places !

Why gravity is so weak compared to the other interactions ?

Newton's law

$$m \bullet \longleftarrow r \longrightarrow \bullet m$$
 $F_{\text{grav}} = G_N \frac{m^2}{r^2}$ $G_N^{-1/2} = M_{\text{Planck}} = 10^{19} \text{ GeV}$
Compare with electric force: $F_{\text{el}} = \frac{e^2}{r^2} \Rightarrow$

effective dimensionless coupling $G_N m^2$ or in general $G_N E^2$ at energies E

$$E = m_{\rm proton} \Rightarrow \frac{F_{\rm grav}}{F_{\rm el}} = \frac{G_N m_{\rm proton}^2}{e^2} \simeq 10^{-40} \Rightarrow$$
 Gravity is very weak !

At what energy gravitation becomes comparable to the other interactions?

 $M_{
m Planck}\simeq 10^{19}~{
m GeV}
ightarrow {
m Planck}$ length: $10^{-33}~{
m cm}$

 10^{15} \times the LHC energy ! [39]

- Quantize gravity
- Charge quantization
- Origin of Dark matter [10]
- Origin of Dark energy
- Baryon asymmetry in the Universe
- Origin of neutrino masses

New states, beyond Higgs mechanism, lepton number violation ?

• Unification of gauge couplings ? [11]

(Hyper)charge quantization

All color singlet states have integer charges Why? $SU(3) \times SU(2) \times U(1)_Y$ $Q = T_3 + Y$

 $q = (3,2)_{1/6} \qquad q = \begin{pmatrix} u_{2/3} \\ d_{-1/3} \end{pmatrix}$ $u^{c} = (\bar{3},1)_{-2/3}$ $d^{c} = (\bar{3},1)_{1/3}$ $\ell = (1,2)_{-1/2} \qquad \ell = \begin{pmatrix} \nu_{0} \\ e_{-1} \end{pmatrix}$ $e^{c} = (1,1)_{1}$

In a non-abelian theory charges are quantized

e.g. SU(2): T_3 eigenvalues are 1/2-integers

What our Universe is made of ?

- Ordinary matter: only a tiny fraction
- Non-luminous (dark) matter: $\sim 25\%$

Natural explanation: new stable Weakly Interacting Massive Particle in the LHC energy region

• Unknown relativistic dark energy: $\sim 70\%$ [8]



Gauge coupling unification

Energy evolution of gauge couplings $\alpha_i = \frac{g_i^2}{4\pi}$:

$$\frac{d\alpha_i}{d\ln Q} = -\frac{b_i}{2\pi}\alpha_i^2 \quad \Rightarrow \quad \alpha_i^{-1}(Q) = \alpha_i^{-1}(Q_0) - \frac{b_i}{2\pi}\ln\frac{Q}{Q_0}$$

low energy data \rightarrow extrapolation at high energies:



 \Rightarrow unification at $M_{GUT} \simeq 10^{15} - 10^{16}$ GeV

Standard picture: low energy supersymmetry

every particle has a superpartner with spin differ by $1/2\,$

cancel large quantum corrections to the Higgs mass

 \Rightarrow superpartner mass splittings must be not far from the weak scale

Advantages:

- natural elementary scalars
- gauge coupling unification
- LSP: natural dark matter candidate
- radiative EWSB

Problems:

- too many parameters: soft breaking terms
- MSSM : already a % ‰ fine-tuning 'little' hierarchy problem

Natural framework: Heterotic string (or high-scale M/F) theory

ATLAS SUSY Searches* - 95% CL Lower Limits (Status: March 26, 2013)

Inclusive searches	$\begin{array}{l} \text{MSUGRA/CMSSM: 0 lap + f = 4 = } r_{\text{mis}} \\ \text{MSUGRA/CMSSM: 1 lap + f = 4 = } r_{\text{mis}} \\ \text{MSUGRA/CMSSM: 1 lap + f = 4 = } r_{\text{mis}} \\ \text{Pheno model: 0 lap + f = 4 = } r_{\text{mis}} \\ \text{Pheno model: 0 lap + f = 4 = } r_{\text{mis}} \\ \text{MSUGRA/CMSM: 1 lap + 1 = 4 = } r_{\text{mis}} \\ \text{MSSB (R NLSP): 1 = 2 + f = 4 = } r_{\text{mis}} \\ \text{GMS (M NLSP): 2 + 1 = 4 = } r_{\text{mis}} \\ \text{GM (wino NLSP): 7 + 1 = 4 = } r_{\text{mis}} \\ \text{GM (wino NLSP): 7 + 1 = 4 = } r_{\text{mis}} \\ \text{GM (mison NLSP): 7 + 1 = 4 = } r_{\text{mis}} \\ \text{GM (mison NLSP): 7 + 1 = 4 = } r_{\text{mis}} \\ \text{Gravition SP: 7 model + 4 = } r_{\text{mis}} \\ \text{Gravition SP: 7 model + 4 = } r_{\text{mis}} \\ \end{array}$	Let an " a two untractions on the test let a m" a two provides constructions let a more than a construction let a more than a construction let a more than a construction let a more than a more than a more let a more than a more than a more than let a more than a more than a more than a more let a more than a more than a more than a more let a more than a more than a more than a more than a more let a more than a more than a more than a more than a more let a more than a more than a more than a more than a more let a more than a more than a more than a more than a more let a more than a more let a more than a		$ \begin{split} \widetilde{q} &= \widetilde{g} \; \text{mass} \\ \widetilde{g} \; \text{mass} \\ \mathfrak{sg} \; \text{mass} \\ \text{mass} \; (\mathfrak{m} \widetilde{g}) < 2 \text{fiv}, \text{light} \widetilde{\chi}_{1}^{3} \\ \tilde{g} \; \text{mass} \; (\mathfrak{m} \widetilde{g}) < 2 \text{fiv}, \text{light} \widetilde{\chi}_{1}^{3} \\ (\mathfrak{m} \widetilde{g}_{1}^{2} > 20 \text{cut}, \mathfrak{m}_{1}^{2} + \frac{1}{2} \text{mig}_{1}^{2} \text{mig}_{1}^{2} \\ \text{mass} \; (\mathfrak{m} \mathfrak{u}^{-1} < 5) \\ \text{mass} \; (\mathfrak{m} \mathfrak{u}^{-1} < 5) \\ \text{mass} \; (\mathfrak{m} \mathfrak{u}^{-1} > 50 \text{GeV}) \\ (\mathfrak{m} \widetilde{\zeta}_{1}^{1} > 220 \text{GeV}) \\ \mathfrak{h} > 200 \text{GeV} \\ (\mathfrak{m} \widetilde{\zeta}_{1}^{1} > \mathfrak{r}^{2} \\ \mathfrak{m} \widetilde{\eta}_{1}^{2} \\ \text{mig} \\ \text{mig} \\ \tilde{\eta} > \mathfrak{m}^{2} \\ \text{mig} \\ \tilde{\eta} > \mathfrak{m}^{2} \\ \tilde{\eta} > \mathfrak{m}^$	ATLAS Preliminary
3rd gen. gluino mediated	$\tilde{g} \rightarrow bb\gamma^{\prime\prime}_{\gamma}$: 0 lep + 3 b-j's + $E_{\tau,mins}$ $\tilde{g} \rightarrow t\tilde{t}\gamma^{\prime\prime}_{\gamma}$: 2 SS-lep + (0-3b-)j's + $E_{\tau,mins}$ $\tilde{g} \rightarrow t\tilde{t}\gamma^{\prime\prime}_{\gamma}$: 0 lep + multi-j's + $E_{\tau,mins}$ $\tilde{g} \rightarrow t\tilde{t}\gamma^{\prime\prime}_{\gamma}$: 0 lep + 3 b-j's + $E_{\tau,mins}$	L=12.8 fb ⁻¹ , 8 TeV [ATLAS-CONF-2012-145] L=20.7 fb ⁻¹ , 8 TeV [ATLAS-CONF-2013-007] L=5.8 fb ⁻¹ , 8 TeV [ATLAS-CONF-2012-163] L=12.8 fb ⁻¹ , 8 TeV [ATLAS-CONF-2012-145]	1.24 TeV g m 900 GeV ğ mass 1.00 TeV ğ mas 1.15 TeV ğ m	mass $(m[\tilde{\chi}_1^*] < 200 \text{ GeV})$ 8 T $(any m[\tilde{\chi}_1^*])$ 8 T S $(m[\tilde{\chi}_1^*] < 300 \text{ GeV})$ 8 T ass $(m[\tilde{\chi}_1^*] < 200 \text{ GeV})$ 8 T	eV, all 2012 data eV, partial 2012 data
3rd gen. squarks direct production 11	$\begin{array}{c} bb, b \rightarrow b y_i^{(2)} (1 \text{ is } p + 2 p \text{ is } s + \varepsilon_{ $	L+12 6/1 2 16/2 (ALL-2C-004-2013-607) L-20 7 6/1 2 16/2 (B20-305, 120-301-607) L-20 7 6/1 2 16/2 (B20-305, 120-302-607) L-20 7 6/1 2 16/2 (ALL-3C-004-2013-607) L-20 7 6/1 2 16/2 (ALL-3C-004-2013-607) L-20 5 6/1 2 16/2 (ALL-3C-004-2013-607) L-20 5 6/1 2 16/2 (ALL-3C-004-2013-607) L-20 5 6/1 2 16/2 (ALL-3C-004-2013-607) L-20 7 6/1 2 16/2 (ALL-3C-004-2013-607)	Base and Datases (mg2) All operating and the second seco	< 120 GeV) 7 T)) m(t)m(z) = 150 GeV) m(t)m(z) = 10 GeV) = 0) 3 GeV) 3 GeV) 3 GeV)	ēV, all 2011 data
EW direct	$\begin{array}{c} \left(\bigcup_{i}, \bigcup_{i} - i\overline{\chi}_{i}^{0} : 2 \text{ lep } + \mathcal{E}_{\text{Trains}}^{-} \\ \tilde{\chi}_{i}^{+} \tilde{\chi}_{i}^{-} \tilde{\chi}_{i}^{-} \rightarrow \tilde{h}(\overline{V}) : 2 \text{ lep } + \mathcal{E}_{\text{Trains}} \\ \tilde{\chi}_{i}^{+} \tilde{\chi}_{i}^{-} \rightarrow \tilde{U}(\overline{V}) : 2 \tau + \mathcal{E}_{\text{Trains}} \\ \tilde{\chi}_{i}^{+} \tilde{\chi}_{i}^{-} \rightarrow \tilde{U}(\overline{V}) : \tilde{V}(\overline{V}) : 3 \text{ lep } + \mathcal{E}_{\text{Trains}} \\ \tilde{\chi}_{i}^{+} \tilde{\chi}_{i}^{-} \rightarrow \tilde{U}(\overline{V}) : \tilde{V}(\overline{V}) : 3 \text{ lep } + \mathcal{E}_{\text{Trains}} \\ \tilde{\chi}_{i}^{+} \tilde{\chi}_{i}^{-} \rightarrow \tilde{U}(\overline{V}) : \tilde{\chi}_{i}^{+} \tilde{\chi}_{i}^{-} \rightarrow \tilde{U}(\overline{V}) : 3 \text{ lep } + \mathcal{E}_{\text{Trains}} \end{array}$	L=4.7 fb ⁻¹ , 7 TeV [1208.2884] L=4.7 fb ⁻¹ , 7 TeV [1208.2884] L=20.7 fb ⁻¹ , 8 TeV [ATLAS-CONF-2013-628] L=20.7 fb ⁻¹ , 8 TeV [ATLAS-CONF-2013-628] L=20.7 fb ⁻¹ , 8 TeV [ATLAS-CONF-2013-635]	85-195 GeV Î mass (m(0 [±] ₁) = 0) 110-340 GeV \$\$\vec{\cap_{1}}{\cap_{1}}\$ mass (m(0 [±] ₁) < 10 GeV, r 180-330 GeV \$\$\vec{\cap_{1}}{\cap_{1}}\$ mass (m(0 [±] ₁) < 10 GeV, r 180-330 GeV \$\$\vec{\cap_{1}}{\cap_{1}}\$ mass (m(0 [±] ₁) < 10 GeV, r 180-330 GeV \$\$\vec{\cap_{1}}{\cap_{1}}\$ mass (m(0 [±] ₁) < 10 GeV, r 180-330 GeV \$\$\vec{\cap_{1}}{\cap_{1}}\$ mass (m(0 [±] ₁) < 10 GeV, r 180-330 GeV \$\$\vec{\cap_{1}}{\cap_{1}}\$ mass (m(0 [±] ₁) < 10 GeV, r 180-330 GeV \$\$\vec{\cap_{1}}{\cap_{1}}\$ mass (m(0 [±] ₁) < 10 GeV, r	$\begin{split} & \widehat{n}(\overline{x}) = \frac{1}{2} (m(\overline{x}_1^+) + m(\overline{x}_2^+))) \\ & \widehat{n}(\overline{x}) = \frac{1}{2} (m(\overline{x}_1^+) + m(\overline{x}_1^+))) \\ & \widehat{n}_1^+ = m(\overline{x}_2^+) + m(\overline{x}_1^+) = 0, m(\overline{x}^+) \text{ as above } \end{split}$	
Long-lived particles	$\begin{array}{l} \text{Direct} \ \widetilde{\chi}_1^* \ \widetilde{\text{pair}} \ \text{prod.} \ (\text{AMSB}) : \text{long-lived} \ \widetilde{\chi}_1^* \\ \text{Stable} \ \widetilde{g}, \ R\text{-hadrons} : \text{low} \ \beta, \ \beta\gamma \\ \text{GMSB}, \ \text{stable} \ \widetilde{\tau} : \text{low} \ \beta, \ \text{GMSB}, \ \text{stable} \ \widetilde{\tau} : \text{low} \ \beta, \ \text{GMSB}, \ \tilde{\chi}_2^* \rightarrow \gamma \widetilde{G} : \text{non-pointing photons} \\ \text{GMSB}, \ \widetilde{\chi}_1^0 \rightarrow \gamma \widetilde{G} : \text{non-pointing photons} \\ \text{ord} \ (\text{RPV}) : \ \mu + \text{heavy displaced vertex} \end{array}$	L+4.7 fb ⁻¹ , 7 TeV [1210.2832] L+4.7 fb ⁻¹ , 7 TeV [1211.1597] L+4.7 fb ⁻¹ , 7 TeV [1211.1597] L+4.7 fb ⁻¹ , 7 TeV [ATLAS-CONF-2013-016] L=4.4 fb ⁻¹ , 7 TeV [1210.7451]	220 GeV \$\tilde{\chi}_1^+\$ mass (1 < \tilde{\chi}_1^+) < 10 mass 985 GeV \$\tilde{g}\$ mass 300 GeV \$\tilde{x}\$ mass 230 GeV \$\tilde{\chi}\$ mass 230 GeV \$\tilde{\chi}\$ mass 700 GeV \$\tilde{\chi}\$ mass 700 GeV \$\tilde{\tilde{x}\$ mass	s s nm < ct < 1 m,õjdeccupled)	
Nd X X	$\begin{array}{c} LFV: pp{-}\bar{v};t,\bar{v}, \neg e{+}\mathfrak{t} \text{ resonance}\\ LFV: pp{-}\bar{v};t,\bar{v}, \neg e(\mathfrak{q}) + resonance\\ linear\;RPV\;CMSSM: 1\;lep{+}\mathcal{T} s + E_{rmax}\\ g_{\bar{v}}^{'} \rightarrow W_{\bar{v}}^{'}, \bar{v}^{'}_{\gamma} \rightarrow eev_{\mu}eu^{'} : 1\;lep{+}\mathcal{T} s + E_{rmax}\\ g_{\bar{v}}^{'} \rightarrow W_{\bar{v}}^{'}, v^{'}_{\gamma} \rightarrow rtv_{u}eev_{v} : 3\;lep{+}1\;t + E_{rmax}\\ \bar{g} \rightarrow qq : 3\;let\;resonance\;pair\\ \bar{g} \rightarrow qq : 3\;-let\;rosonance\;pair : 2\;S\;shep{+}1\;(o\;d S) = E_{rmax}\\ \bar{g} \rightarrow qq : 3\;-let\;rosonance\;pair : S\;S\;shep : q : q : S\;bhep : q : q : S\;shep : q : q : S\;shep : q : q : S\;bhep : q : q : S\;shep : q : q : S\;bhep : q : q : S\;shep : q : q : shep : she$	L=4.6 th ⁻¹ , 7 TeV [1212.1227] L=4.6 th ⁻¹ , 7 TeV [1212.1227] L=4.7 th ⁻¹ , 7 TeV [ATLAS-CONF-2013-469] L=20.7 th ⁻¹ , 8 TeV [ATLAS-CONF-2013-636] L=20.7 th ⁻¹ , 8 TeV [ATLAS-CONF-2013-667] L=20.7 th ⁻¹ , 8 TeV [ATLAS-CONF-2013-607]	1.01 TeV 1.10 TeV \v 1.10 TeV	$ \begin{array}{ c c c c c c c c c c c c c c c c c c c$	
WIMP in	Scalar gluon : 2-jet resonance päir teraction (D5, Dirac χ) : 'monojet' + Ε _{τ,niss}	L+4.6 fb ⁻¹ , 7 TeV [1210.4826] L+10.5 fb ⁻¹ , 8 TeV [ATLAS-CONF-2012-147] 10	100-227 GeV sgluon mass (incl. limit from 11 704 GeV M* scale (i 1 1	10.2693) m _y < 80 GeV, limit of < 687 GeV for D8)	10

 $^{*}Only$ a selection of the available mass limits on new states or phenomena shown. All limits quoted are observed minus 1 σ theoretical signal cross section uncertainty.

Mass scale [TeV]



I. Antoniadis (CERN)

Remarks on the value of the Higgs mass $\sim 126~\text{GeV}$

- consistent with expectation from precision tests of the SM
- favors perturbative physics quartic coupling $\lambda = m_H^2/v^2 \simeq 1/8$
- 1st elementary scalar in nature signaling perhaps more to come

Window to new physics ?

- compatible with supersymmetry
 but appears fine-tuned in its minimal version [17]
 early to draw a general conclusion before LHC13/14
 e.g. an extra singlet or split families can alleviate the fine tuning
- very important to measure its properties and couplings any deviation of its couplings to top, bottom and EW gauge bosons implies new light states involved in the EWSB altering the fine-tuning



Fine-tuning in MSSM

Upper bound on the lightest scalar mass:

$$m_h^2 \lesssim m_Z^2 \cos^2 2\beta + \frac{3}{(4\pi)^2} \frac{m_t^4}{v^2} \left[\ln \frac{m_{\tilde{t}}^2}{m_t^2} + \frac{A_t^2}{m_{\tilde{t}}^2} \left(1 - \frac{A_t^2}{12m_{\tilde{t}}^2} \right) \right] \lesssim (130 \, GeV)^2$$

 $m_h \simeq 126 \,\, {
m GeV} \, \Rightarrow \, m_{ ilde{t}} \simeq 3 \,\, {
m TeV}$ or $A_t \simeq 3 m_{ ilde{t}} \simeq 1.5 \,\, {
m TeV}$

 \Rightarrow % to a few ‰ fine-tuning

minimum of the potential:
$$m_Z^2 = 2 rac{m_1^1 - m_2^2 \tan_eta^2}{ an^2 eta - 1} \sim -2m_2^2 + \cdots$$

 $\begin{array}{ll} \mathsf{RG evolution:} & m_2^2 = & m_2^2(M_{\mathrm{GUT}}) - \frac{3\lambda_t^2}{4\pi^2}m_{\tilde{t}}^2\ln\frac{M_{\mathrm{GUT}}}{m_{\tilde{t}}} + \cdots \, {}_{\scriptscriptstyle [27]} \\ & \sim & m_2^2(M_{\mathrm{GUT}}) - \mathcal{O}(1)m_{\tilde{t}}^2 + \cdots \end{array}$

• minimize radiative corrections

 $M_{\rm GUT} \rightarrow \Lambda$: low messenger scale (gauge mediation)

$$\delta m_{\tilde{t}} = \frac{8\alpha_s}{3\pi} M_3^2 \ln \frac{\Lambda}{M_3} + \cdots$$

extend the MSSM

extra fields beyond LHC reach \rightarrow effective field theory approach

o . . .

MSSM with dim-5 and 6 operators

I.A.-Dudas-Ghilencea-Tziveloglou '08, '09, '10

parametrize new physics above MSSM by higher-dim effective operators

relevant super potential operators of dimension-5:

$$\mathcal{L}^{(5)} = \frac{1}{M} \int d^2 \theta \left(\eta_1 + \eta_2 S \right) \left(H_1 H_2 \right)^2$$

 η_1 : generated for instance by a singlet

$$W = \lambda \sigma H_1 H_2 + M \sigma^2 \quad \rightarrow \quad W_{\text{eff}} = \frac{\lambda^2}{M} (H_1 H_2)^2$$

Strumia '99 ; Brignole-Casas-Espinosa-Navarro '03 Dine-Seiberg-Thomas '07

 η_2 : corresponding soft breaking term spurion $S \equiv m_S \theta^2$

Physical consequences of MSSM₅: Scalar potential

$$\begin{split} \mathcal{V} &= m_1^2 |h_1|^2 + m_2^2 |h_2|^2 + B\mu (h_1 h_2 + \text{h.c.}) + \frac{g_2^2 + g_Y^2}{8} \left(|h_1|^2 - |h_2|^2 \right)^2 \\ &+ \left(|h_1|^2 + |h_2|^2 \right) \left(\eta_1 h_1 h_2 + \text{h.c.} \right) + \frac{1}{2} \left[\eta_2 (h_1 h_2)^2 + \text{h.c.} \right] \\ &+ \eta_1^2 |h_1 h_2|^2 \left(|h_1|^2 + |h_2|^2 \right) \end{split}$$

- $\eta_{1,2} \Rightarrow$ quartic terms along the D-flat direction $|h_1| = |h_2|$
- tree-level mass can increase significantly
- bigger parameter space for LSP being dark matter

Bernal-Blum-Nir-Losada '09

• last term $\sim \eta_1^2$: guarantees stability of the potential

but requires addition of dim-6 operators

MSSM Higss with dim-6 operators

dim-6 operators can have an independent scale from dim-5

Classification of all dim-6 contributing to the scalar potential (without SUSY) \Rightarrow

large tan β expansion: $\delta_6 m_h^2 = f v^2 + \cdots$ constant receiving contributions from several operators

$$f \sim f_0 imes \left(\mu^2/M^2, \ m_S^2/M^2, \ \mu m_S/M^2, \ v^2/M^2
ight)$$

 $m_S=1$ TeV, M=10 TeV, $f_0\sim 1-2.5$ for each operator

 $\Rightarrow m_h \simeq 103 - 119 \text{ GeV}$

 \Rightarrow MSSM with dim-5 and dim-6 operators:

possible resolution of the MSSM fine-tuning problem

Can the SM be valid at high energies?

Degrassi-Di Vita-Elias Miró-Espinosa-Giudice-Isidori-Strumia '12



Instability of the SM Higgs potential \Rightarrow metastability of the EW vacuum



 $\lambda=0$ at a scale $\geq 10^{10}~{
m GeV} \Rightarrow m_{H}=126\pm 3~{
m GeV}$

Ibanez-Valenzuela '13



If the weak scale is tuned \Rightarrow split supersymmetry is a possibility Arkani Hamed-Dimopoulos '04, Giudice-Romaninio '04

- natural splitting: gauginos, higgsinos carry R-symmetry, scalars do not
- main good properties of SUSY are maintained gauge coupling unification and dark matter candidate
- also no dangerous FCNC, CP violation, ...
- experimentally allowed Higgs mass \Rightarrow 'mini' split

 $m_S \sim {
m few}$ - thousands TeV

gauginos: a loop factor lighter than scalars ($\sim m_{3/2}$)

• natural string framework: intersecting (or magnetized) branes

IA-Dimopoulos '04

D-brane stacks are supersymmetric with massless gauginos intersections have chiral fermions with broken SUSY & massive scalars

Giudice-Strumia '11

Predicted range for the Higgs mass



An extra U(1) can also cure the instability problem Anchordoqui-IA-Goldberg-Huang-Lüst-Taylor-Vicek '12

usually associated to known global symmetries of the SM: B, L, \ldots

- B anomalous and superheavy
- B L massless at the string scale (no associated 6d anomaly) but broken at TeV by a scalar VEV with the quantum numbers of N_R
- L-violation from higher-dim operators suppressed by the string scale
- U(3) unification, Y combination \Rightarrow 2 parameters: 1 coupling + $m_{Z''}$
- perturbativity $\Rightarrow 0.5 \lesssim g_{U(1)_R} \lesssim 1$
- interesting LHC phenomenology and cosmology

Alternative answer: Low UV cutoff $\Lambda \sim \text{TeV}$

- low scale gravity \Rightarrow extra dimensions: large flat or warped
- low string scale \Rightarrow low scale gravity, ultra weak string coupling

Experimentally testable framework:

- spectacular model independent predictions
- radical change of high energy physics at the TeV scale

Moreover no little hierarchy problem:

radiative electroweak symmetry breaking with no logs [17]

 $\Lambda \sim$ a few TeV and $m_{H}^2 =$ a loop factor $imes \Lambda^2$ [46]

But unification has to be probably dropped

New Dark Matter candidates e.g. in the extra dims

Standard Model of electroweak + strong Interactions :

- Quantum Field Theory: Quantum Mechanics + Special Relativity
- Principle: gauge invariance $U(1) \times SU(2) \times SU(3)$

String theory : • Quantum Mechanics + General Relativity

point particle \rightarrow extended objects

Framework for unification of all interactions

Mass scale: String tension $M_{\rm s} \leftrightarrow \text{string size: } I_{\rm s}$



Consistent theory \Rightarrow 9 spatial dimensions !

six new dimensions of space

matter and gauge interactions may be localized

in less than 9 dimensions \Rightarrow

our universe on a membrane ? [34]

p-plane: extended in p spatial dimensions

p = 0: particle, p = 1: string,...

Extra Dimensions

how they escape observation?

finite size R

energy cost to send a signal: $E > R^{-1} \leftarrow$ compactification scale

experimental limits on their size

light signal $\Rightarrow E \gtrsim 1 \text{ TeV}$ $R \lesssim 10^{-16} \text{ cm}$

how to detect their existence?

motion in the internal space \Rightarrow mass spectrum in 3d

Kaluza and Klein 1920

Dimensions D=??



example: - one internal circular dimension

- light signal



plane waves e^{ipy} periodic under $y \rightarrow y + 2\pi R$

 \Rightarrow quantization of internal momenta: $p = \frac{n}{R}$; n = 0, 1, 2, ...

 \Rightarrow 3d: tower of Kaluza Klein particles with masses $M_n = n/R$

$$p_0^2 - \vec{p}^2 - p_5^2 = 0 \implies p^2 = p_5^2 = \frac{n^2}{R^2}$$

 $E >> R^{-1}$: emission of many massive photons \Leftrightarrow propagation in the internal space [30]

Our universe on a membrane



Two types of new dimensions:

- longitudinal: along the membrane
- transverse: "hidden" dimensions only gravitational signal $\Rightarrow R_{\perp} \lesssim 1 \text{ mm}$!

Gravity modification at submillimeter distances

Newton's law: force decreases with area



3d: force $\sim 1/r^2$ (3+*n*)d: force $\sim 1/r^{2+n}$

observable for n = 2: $1/r^4$ with r << .1 mm [38]

Adelberger et al. '06



 ${\it R}_{\perp} \lesssim$ 45 $\mu{\rm m}$ at 95% CL

• dark-energy length scale pprox 85 μ m
Experiment: Relativistic dark energy 70-75% of the observable universe

negative pressure: $p = -\rho \Rightarrow$ cosmological constant

$$R_{ab} - \frac{1}{2}Rg_{ab} + \Lambda g_{ab} = \frac{8\pi G}{c^4}T_{ab} \Rightarrow \rho_{\Lambda} = \frac{c^4\Lambda}{8\pi G} = -\rho_{\Lambda}$$

Two length scales:

- $[\Lambda] = L^{-2} \leftarrow \text{size of the observable Universe}$ $\Lambda_{obs} \simeq 0.74 \times 3H_0^2/c^2 \simeq 1.4 \times (10^{26} \text{ m})^{-2}$ Hubble parameter $\simeq 73 \text{ km s}^{-1} \text{ Mpc}^{-1}$
- $\left[\frac{\Lambda}{G} \times \frac{c^3}{\hbar}\right] = L^{-4} \leftarrow \text{dark energy length} \simeq 85 \mu \text{m}$

 \Rightarrow Gravity modification at large (cosmological) and short distances ?

Low scale gravity

Extra large \perp dimensions can explain the apparent weakness of gravity total force = observed force \times volume \perp total force $\simeq \mathcal{O}(1)$ at 1 TeV n dimensions of size R \Rightarrow volume $R^n_{\perp} = 10^{32}$ in TeV⁻¹ units n=1 : $R_\perp\simeq 10^8$ km excluded n = 2: $R_{\perp} \simeq 0.1 \text{ mm}$ (10⁻¹² GeV) possible n = 6: $R_{\perp} \simeq 10^{-13}$ mm (10⁻² GeV) • distances $> R_{\perp}$: gravity 3d however for $< R_{\perp}$: gravity (3+n)d [35]

 \bullet strong gravity at $10^{-16}~\text{cm}\leftrightarrow 10^3~\text{GeV}$

10³⁰ times stronger than thought previously!

Extra large \perp dimensions can explain the apparent weakness of gravity total force = observed force \times volume \perp [7] $\uparrow \qquad \uparrow \qquad \uparrow$ $G_N^* E^{2+n} = G_N E^2 \times V_\perp E^n$ $G_{N}^{*} = M_{*}^{-(2+n)}$: (4 + n)-dim gravitational constant total force $\simeq \mathcal{O}(1)$ at 1 TeV *n* dimensions of size R_{\perp} $\Rightarrow V_{\perp} = R_{\perp}^{n}$ $\Rightarrow 1 = E^2/M_P^2 \times (R_\perp E)^n$ for $E = M_* \simeq 1$ TeV

Connect string theory to the real world: What is the value of the string scale M_s ?

- arbitrary parameter : Planck mass $M_P \longrightarrow \text{TeV}$
- physical motivations \Rightarrow favored energy regions:

• High :
$$\left\{ \begin{array}{ll} M_P^* \simeq 10^{18} \,\, {\rm GeV} & {\rm Heterotic \ scale} \\ \\ M_{\rm GUT} \simeq 10^{16} \,\, {\rm GeV} & {\rm Unification \ scale} \end{array} \right.$$

• Intermediate : around 10^{11} GeV $(M_s^2/M_P \sim \text{TeV})$

SUSY breaking, strong CP axion, see-saw scale

• Low : TeV (hierarchy problem)

perturbative heterotic string : the most natural for SUSY and unification gravity and gauge interactions have same origin massless excitations of the closed string

But mismatch between string and GUT scales:

 $M_s = g_H M_P \simeq 50 M_{
m GUT}$ $g_H^2 \simeq lpha_{
m GUT} \simeq 1/25$ [43]

in GUTs only one prediction from 3 gauge couplings unification: $\sin^2 \theta_W$ introduce large threshold corrections or strong coupling $\rightarrow M_s \simeq M_{GUT}$ but loose predictivity

gravity + gauge kinetic terms [44]

$$\int [d^{10}x] \frac{1}{g_H^2} M_H^8 \mathcal{R}^{(10)} + \int [d^{10}x] \frac{1}{g_H^2} M_H^6 \mathcal{F}_{MN}^2 \quad \text{simplified units: } 2 = \pi = 1$$

Compactification in 4 dims on a 6-dim manifold of volume $V_6 \Rightarrow$

Type I string theory ⇒ D-brane world I.A.-Arkani-Hamed-Dimopoulos-Dyali '98

- gravity: closed strings propagating in 10 dims
- gauge interactions: open strings with their ends attached on D-branes

Dimensions of finite size: *n* transverse 6 - n parallel calculability $\Rightarrow R_{\parallel} \simeq I_{\text{string}}$; R_{\perp} arbitrary

small M_s/M_P \Rightarrow extra-large R_{\perp}

 $R_{\perp} \sim .1 - 10^{-13}$ mm for n = 2 - 6

distances $< R_{\perp}$: gravity (4+*n*)-dim \rightarrow strong at 10⁻¹⁶ cm

 $M_{\rm s} \sim 1 {
m TeV} \Rightarrow R_{\perp}^n = 10^{32} I_{\rm s}^n$ [68]

Type I/II strings: gravity and gauge interactions have different origin gravity + gauge kinetic terms $\int [d^{10}x] \frac{1}{g_s^2} M_s^8 \mathcal{R}^{(10)} + \int [d^{p+1}x] \frac{1}{g_s} M_s^6 \mathcal{F}_{MN}^2$ [42]

Compactification in 4 dims \Rightarrow

Braneworld

I.A.-Arkani-Hamed-Dimopoulos-Dvali '98

2 types of compact extra dimensions:

• parallel (d_{\parallel}): $\lesssim 10^{-16}$ cm (TeV) • transverse (\perp): $\lesssim 0.1 \text{ mm} (\text{meV})$



Origin of EW symmetry breaking?

possible answer: radiative breaking [27] I.A.-Benakli-Quiros '00 $V = \mu^2 H^{\dagger} H + \lambda (H^{\dagger} H)^2$ $\mu^2 = 0$ at tree but becomes < 0 at one loop non-susy vacuum simplest case: one scalar doublet from the same brane \Rightarrow tree-level V same as susy: $\lambda = \frac{1}{8}(g_2^2 + g'^2)$ D-terms $\mu^2 = -g^2 \varepsilon^2 M_s^2 \leftarrow \text{effective UV cutoff}$ $e^{2}(R) = \frac{R^{3}}{2\pi^{2}} \int_{0}^{\infty} dll^{3/2} \frac{\theta_{2}^{4}}{16l^{4}\eta^{12}} \left(il + \frac{1}{2}\right) \sum n^{2} e^{-2\pi n^{2}R^{2}l}$



Accelerator signatures: 4 different scales

- Gravitational radiation in the bulk \Rightarrow missing energy present LHC bounds: $M_* \gtrsim 3-5$ TeV
- Massive string vibrations \Rightarrow e.g. resonances in dijet distribution [53]

 $M_j^2 = M_0^2 + M_s^2 j$; maximal spin: j + 1

higher spin excitations of quarks and gluons with strong interactions present LHC limits: $M_s\gtrsim 5~{
m TeV}$

• Large TeV dimensions \Rightarrow KK resonances of SM gauge bosons I.A. '90

$$M_k^2 = M_0^2 + k^2/R^2$$
; $k = \pm 1, \pm 2, \dots$

experimental limits: $R^{-1} \gtrsim 0.5 - 4$ TeV (UED - localized fermions) [55]

• extra U(1)'s and anomaly induced terms

masses suppressed by a loop factor from M_s [59]

Gravitational radiation in the bulk \Rightarrow missing energy



Angular distribution \Rightarrow spin of the graviton

Collider bounds on R_{\perp} in mm				
	<i>n</i> = 2	<i>n</i> = 4	<i>n</i> = 6	
LEP 2	$4.8 imes10^{-1}$	$1.9 imes10^{-8}$	$6.8 imes10^{-11}$	
Tevatron	$5.5 imes10^{-1}$	$1.4 imes10^{-8}$	$4.1 imes 10^{-11}$	
LHC	$4.5 imes10^{-3}$	$5.6 imes10^{-10}$	$2.7 imes10^{-12}$	

present LHC bounds:

 $M_* \gtrsim 3 - 5$ TeV [51]

cooling due to graviton production e.g. $NN \rightarrow NN + \text{graviton}$ number of gravitons: $\sim (TR_{\perp})^n$ $T >> R_{\perp}^{-1}$ $\sim 10 \text{ MeV}$ \Rightarrow production rate: $P_{\text{gr}} \sim \frac{1}{M_P^2} (TR_{\perp})^n \sim \frac{T^n}{M_*^{(2+n)}}$ $P_{\text{gr}} < P_{\nu} \Rightarrow M_* |_{n=2} \gtrsim 50 \text{ TeV} \Rightarrow M_s \gtrsim 10 \text{ TeV}$



String-size black hole energy threshold : $M_{
m BH}\simeq M_s/g_s^2$

Horowitz-Polchinski '96, Meade-Randall '07

- string size black hole: $r_H \sim l_s = M_s^{-1}$
- black hole mass: $M_{\rm BH} \sim r_H^{d-3}/G_N$ $G_N \sim I_s^{d-2}g_s^2$

weakly coupled theory \Rightarrow strong gravity effects occur much above M_s , M_* $g_s \sim 0.1$ (gauge coupling) $\Rightarrow M_{\rm BH} \sim 100 M_s$

Comparison with Regge excitations : $M_j = M_s \sqrt{j} \Rightarrow$

production of $j\sim 1/g_s^4\sim 10^4$ string states before reach $M_{
m BH}$ [48]

Universal deviation from Standard Model in jet distribution

 $M_s = 2 \text{ TeV}$ Width = 15-150 GeV

Anchordoqui-Goldberg-Lüst-Nawata-Taylor-Stieberger '08 [48] [51]



Tree level superstring amplitudes involving at most 2 fermions and gluons: model independent for any compactification, # of susy's, even none no intermediate exchange of KK, windings or graviton emmission Universal sum over infinite exchange of string (Regge) excitations

Partonic Luminosity Parton luminosities in pp above TeV are dominated by gq, gg \Rightarrow model independent 10 $gq \rightarrow gq, gg \rightarrow gg, gg \rightarrow q\bar{q}$ 10 10 35

M_s(TeV)

Localized fermions (on 3-brane intersections)

 \Rightarrow single production of KK modes

I.A.-Benakli '94

• strong bounds indirect effects

• new resonances but at most n = 1

Otherwise KK momentum conservation

 \Rightarrow pair production of KK modes (universal dims)



- weak bounds
- no resonances
- lightest KK stable \Rightarrow dark matter candidate

Servant-Tait '02





UED hadron collider phenomenology

- large rates for KK-quark and KK-gluon production
- cascade decays via KK-W bosons and KK-leptons
 determine particle properties from different distributions
- missing energy from LKP: weakly interacting escaping detection
- phenomenology similar to supersymmetry

spin determination important for distinguishing SUSY and UED $\ensuremath{\left[48\right]}$

gluino	1/2	KK-gluon	1
squark	0	KK-quark	1/2
chargino	1/2	KK-W boson	1
slepton	0	KK-lepton	1/2
neutralino	1/2	KK-Z boson	1

SUSY vs UED signals at LHC

Example: jet dilepton final state

SUSY

UED



Extra U(1)'s and anomaly induced terms

masses suppressed by a loop factor

usually associated to known global symmetries of the SM

(anomalous or not) such as (combinations of)

Baryon and Lepton number, or PQ symmetry

Two kinds of massive U(1)'s: I.A.-Kiritsis-Rizos '02

- 4d anomalous U(1)'s: $M_A \simeq g_A M_s$
- 4d non-anomalous U(1)'s: (but masses related to 6d anomalies)

 $M_{NA} \simeq g_A M_s V_2 \leftarrow (6d \rightarrow 4d)$ internal space $\Rightarrow M_{NA} \ge M_A$

or massless in the absence of such anomalies [62]

D-brane embedding of the Standard Model

Generic spectrum: N coincident branes $\Rightarrow U(N)$

a-stack

```
endpoint transformation: N_a or \overline{N}_a U(1)_a charge: +1 or -1

\Rightarrow "baryon" number
```

- open strings from the same stack \Rightarrow adjoint gauge multiplets of $U(N_a)$
- stretched between two stacks \Rightarrow bifundamentals of $U(N_a) \times U(N_b)$

a-stack



non-oriented strings \Rightarrow also:

- orthogonal and symplectic groups SO(N), Sp(N)
- matter in antisymmetric + symmetric reps

Non oriented strings \Rightarrow orientifold planes

where closed strings change orientation

 \Rightarrow mirror branes identified with branes under orientifold action

• strings stretched between two mirror stacks



Standard Model on D-branes : SM⁺⁺



TeV string scale Anchordogui-IA-Goldberg-Huang-Lüst-Taylor '11

- B and L become massive due to anomalies Green-Schwarz terms
- the global symmetries remain in perturbation
 - Baryon number \Rightarrow proton stability
 - Lepton number \Rightarrow protect small neutrino masses

- Lepton number \Rightarrow process _ no Lepton number $\Rightarrow \frac{1}{M_s}LLHH \rightarrow$ Majorana mass: $\frac{\langle H \rangle^2}{M_s}LL$ $\swarrow \sim \text{GeV}$

• $B, L \Rightarrow$ extra Z's

with possible leptophobic couplings leading to CDF-type Wij events $Z' \simeq B$ lighter than 4d anomaly free $Z'' \simeq B - L$

Green-Schwarz anomaly cancellation

$$= k_I^A \sim \operatorname{Tr} Q_A Q_I^2 \to \operatorname{axion} \theta : \delta A = d\Lambda \quad \delta \theta = -m_A \Lambda$$
$$-\frac{1}{4g_I^2} F_I^2 - \frac{1}{2} (d\theta + m_A A)^2 + \frac{\theta}{m_A} k_I^A \operatorname{Tr} F_I \wedge F_I$$
cancel the anomaly

D-brane models: $U(1)_A$ gauge boson acquires a mass but global symmetry remains in perturbation theory string theory: θ = Poincaré dual of a 2-form $d\theta = *dB_2$

R-neutrinos: in the bulk

Arkani Hamed-Dimopoulos-Dvali-March Russell '98 Dienes-Dudas-Gherghetta '98 Dvali-Smirnov '98

R-neutrino: $\nu_R(x, y)$ y: bulk coordinates

$$S_{int} = g_s \int d^4 x H(x) L(x) \nu_R(x, y = 0)$$

$$\langle H \rangle = v \implies \text{mass-term:} \frac{g_s v}{R_\perp^{n/2}} \nu_L \nu_R^0 \leftarrow \text{4d zero-mode}$$

Dirac neutrino masses: $m_{\nu} \simeq \frac{g_s v}{R_{\perp}^{n/2}} \simeq v \frac{M_*}{M_p}$

 $\simeq 10^{-3} - 10^{-2}$ eV for $M_* \simeq 1 - 10$ TeV

 $m_{\nu} << 1/R_{\perp} \Rightarrow$ KK modes unaffected

microgravity experiments

- change of Newton's law at short distances detectable only in the case of two large extra dimensions
- new short range forces light scalars and gauge fields if SUSY in the bulk or broken by the compactification on the brane I.A.-Dimopoulos-Dvali '98, I.A.-Benakli-Maillard-Laugier '02 such as radion and lepton number volume suppressed mass: $(\text{TeV})^2/M_P \sim 10^{-4} \text{ eV} \rightarrow \text{mm}$ range can be experimentally tested for any number of extra dimensions
 - Light U(1) gauge bosons: no derivative couplings
 - \Rightarrow for the same mass much stronger than gravity: $\gtrsim~10^{6}$

Experimental limits on short distance forces



$$V(r) = -G \frac{m_1 m_2}{r} \left(1 + \alpha e^{-r/\lambda}\right)$$

Radion $\Rightarrow M_* \gtrsim 6$ TeV 95% CL Adelberger et al. '06 [75]

More general framework: large number of species

N particle species \Rightarrow lower quantum gravity scale : $M_*^2 = M_p^2/N$

Dvali '07, Dvali, Redi, Brustein, Veneziano, Gomez, Lüst '07-'10 derivation from: black hole evaporation or quantum information storage

 $M_* \simeq 1 \text{ TeV} \Rightarrow N \sim 10^{32} \text{ particle species }!$

- 2 ways to realize it lowering the string scale
 - Large extra dimensions SM on D-branes [43]

 $N = R_{\perp}^{n} I_{s}^{n}$: number of KK modes up to energies of order $M_{*} \simeq M_{s}$

Effective number of string modes contributing to the BH bound

 $N = \frac{1}{g_s^2}$ with $g_s \simeq 10^{-16}$ SM on NS5-branes

I.A.-Pioline '99, I.A.-Dimopoulos-Giveon '01

More general framework: large number of species

N particle species \Rightarrow lower quantum gravity scale : $M_*^2 = M_p^2/N$

Dvali '07, Dvali, Redi, Brustein, Veneziano, Gomez, Lüst '07-'10 derivation from: black hole evaporation or quantum information storage Pixel of size L containing N species storing information:



localization energy $E\gtrsim N/L \rightarrow$ Schwarzschild radius $R_s=N/(LM_p^2)$

no collapse to a black hole : $L \gtrsim R_s \Rightarrow L \gtrsim \sqrt{N}/M_p = 1/M_*$

Gauge/Gravity duality \Rightarrow toy 5d bulk model

Gravity background : near horizon geometry (holography) Maldacena '98

Analogy from D3-branes : AdS_5

NS-5 branes : $(\mathcal{M}_6 \otimes \mathbb{R}_+)$ inear dilaton background in 5d flat string-frame metric $\Phi = -\alpha |y|$ Aharony-Berkooz-Kutasov-Seiberg '98

"cut" the space of the extra dimension \Rightarrow gravity on the brane

$$S_{bulk} = \int d^4x \int_0^{r_c} dy \sqrt{-g} e^{-\Phi} \left(M_5^3 R + M_5^3 (\nabla \Phi)^2 - \Lambda \right)$$
$$S_{vis(hid)} = \int d^4x \sqrt{-g} \left(e^{-\Phi} \right) \left(L_{SM(hid)} - T_{vis(hid)} \right)$$

Tuning conditions: $T_{vis} = -T_{hid} \leftrightarrow \Lambda < 0$

Constant dilaton and AdS metric : Randal Sundrum model

spacetime = slice of AdS₅ : $ds^2 = e^{-2k|y|}\eta_{\mu\nu}dx^{\mu}dx^{\nu} + dy^2$ $k^2 \sim \Lambda/M_5^3$



• exponential hierarchy: $M_W = M_P e^{-2kr_c}$ $M_P^2 \sim M_5^3/k$ $M_5 \sim M_{GUT}$

• 4d gravity localized on the UV-brane, but KK gravitons on the IR $m_n = c_n \, k \, e^{-2kr_c} \sim \text{TeV}$ $c_n \simeq (n + 1/4)$ for large n \Rightarrow spin-2 TeV resonances in di-lepton or di-jet channels

Linear dilaton background IA-Arvanitaki-Dimopoulos-Giveon '11

dilaton $\Phi = -\alpha |y|$ and flat metric \Rightarrow

$$g_s^2 = e^{-lpha|y|}$$
 ; $ds^2 = e^{rac{2}{3}lpha|y|} \left(\eta_{\mu
u} dx^\mu dx^
u + dy^2
ight) \leftarrow$ Einstein frame

 $z \sim e^{\alpha y/3} \Rightarrow$ polynomial warp factor + log varying dilaton



• exponential hierarchy: $g_s^2 = e^{-\alpha|y|}$ $M_P^2 \sim \frac{M_5^3}{\alpha} e^{\alpha r_c}$ $\alpha \equiv k_{RS}$

4d graviton flat, KK gravitons localized near SM
LST KK graviton phenomenology

• KK spectrum :
$$m_n^2 = \left(\frac{n\pi}{r_c}\right)^2 + \frac{\alpha^2}{4}$$
; $n = 1, 2, ...$

 \Rightarrow mass gap + dense KK modes $\alpha \sim 1$ TeV $r_c^{-1} \sim 30$ GeV

• couplings :
$$\frac{1}{\Lambda_n} \sim \frac{1}{(\alpha r_c)M_5}$$

 \Rightarrow extra suppression by a factor $(\alpha r_c) \simeq 30$

• width :
$$1/(\alpha r_c)^2$$
 suppression $\sim 1 \text{ GeV}$

 \Rightarrow narrow resonant peaks in di-lepton or di-jet channels

• extrapolates between RS and flat extra dims (n = 1)

 \Rightarrow distinct experimental signals

Similar to RS using the dilaton as the Goldeberger-Wise scalar $% \left({{{\rm{S}}_{{\rm{B}}}} \right)$

add dilaton boundary potentials \Rightarrow

radion stabilization with the desired hierarchy

Radion phenomenology different from RS:

• mass spectrum: similar to the graviton KK modes

with possible lower parametrically mass gap

- ullet new radion couplings to SM fields besides to the trace of $\mathcal{T}_{\mu
 u}$
- larger coupling to the radion 0-mode relative to KK excitations
- Higgs-radion mixing \Rightarrow

branching fraction to $\gamma\gamma$ can be significantly enhanced

Conclusions

- Confirmation of the EWSB scalar at the LHC: important milestone of the LHC research program
- Precise measurement of its couplings is of primary importance
- Hint on the origin of mass hierarchy and of BSM physics
 - natural or unnatural SUSY?
 - low string scale in some realization?
 - something new and unexpected?

all options are still open

• LHC enters a new era with possible new discoveries

The LHC timeline

LS1 Machine Consolidation

LS2 Machine upgrades for high Luminosity

- Collimation
- Cryogenics
- · Injector upgrade for high intensity (lower emittance)
- · Phase I for ATLAS : Pixel upgrade, FTK, and new small wheel

LS3 Machine upgrades for high Luminosity

- Upgrade interaction region
- · Crab cavities?
- Phase II: full replacement of tracker, new trigger scheme (add L0), readout electronics.



Europe's top priority should be the exploitation of the full potential of the LHC, including the high-luminosity upgrade of the machine and detectors with a view to collecting ten times more data than in the initial design, by around 2030.

Start of LHC 2009 Run 1, 7+8 TeV, ~25 fh⁻¹ int lumi 2013/14 Prepare LHC for LS1 design E & lumi Collect ~30 fb⁻¹ per year at 13/14 TeV 2018 Phase-1 upgrade 152 ultimate lumi Twice nominal lumi at 14 TeV, ~100 fb⁻¹ per year ~2022 Phase-2 upgrade LS3 to HL-LHC ~300 fb⁻¹ per year, run up to > 3 ab^{-1} collected ~2030

IHC timeline