

Complex Geometry and Sigma Models

U. Lindström¹

¹Department of Physics and Astronomy
Division of Theoretical Physics
University of Uppsala

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- Formulations of Generalized Kähler Geometry
- The corresponding Sigma Models

M. Götteman, C. Hull, M. Roček, I. Ryb, R. von Unge, M. Zabzine.

$$(M, g, J_{(\pm)}, H)$$

$$J_{(\pm)}^2 = -\mathbf{1}, \quad J_{(\pm)}^t g J_{(\pm)} = g, \quad \nabla^{(\pm)} J_{(\pm)} = 0$$

$$\Gamma^{(\pm)} = \Gamma^0 \pm \frac{1}{2} g^{-1} H, \quad H = dB.$$

$$E := g + B$$

$$(M, g, J_{(\pm)})$$

$$J_{(\pm)}^2 = -\mathbf{1}, \quad J_{(\pm)}^t g J_{(\pm)} = g, \quad \omega_{(\pm)} := g J_{(\pm)}$$

$$d_{(+)}^c \omega_{(+)} + d_{(-)}^c \omega_{(-)} = 0, \quad dd_{(\pm)}^c \omega_{(\pm)} = 0,$$

$$H := d_{(+)}^c \omega_{(+)} = -d_{(-)}^c \omega_{(-)}$$

Description on $T \oplus T^*$

$$\mathcal{J}_{(1,2)}^2 = -\mathbf{1}, \quad [\mathcal{J}_{(1)}, \mathcal{J}_{(2)}] = 0, \quad \mathcal{J}_{(1,2)}^t \mathcal{I} \mathcal{J}_{(1,2)} = \mathcal{I}, \quad \mathcal{G} := -\mathcal{J}_{(1)} \mathcal{J}_{(2)}$$

$$\mathcal{J}_{(1,2)} =$$

$$\begin{pmatrix} 1 & 0 \\ B & 1 \end{pmatrix} \begin{pmatrix} J_{(+)} \pm J_{(-)} & -(\omega_{(+)}^{-1} \mp \omega_{(-)}^{-1}) \\ \omega_{(+)} \mp \omega_{(-)} & -(J_{(+)}^t \pm J_{(-)}^t) \end{pmatrix} \begin{pmatrix} 1 & 0 \\ -B & 1 \end{pmatrix}$$

$$(M, J_{(\pm)})$$

Locally, \exists “symplectic” two-forms $\mathcal{F}_{(\pm)}$ such that

$$\mathcal{F}_{(\pm)}(v, J_{(\pm)}v) > 0, \quad d(\mathcal{F}_{(+)}J_{(+)} - J_{(-)}^t\mathcal{F}_{(-)}) = 0.$$

$$\mathcal{F}_{(\pm)} = \frac{1}{2}i(B_{(\pm)}^{(2,0)} - B_{(\pm)}^{(0,2)}) \mp \omega_{(\pm)}$$

$$\mathcal{F}_{(+)} = -\frac{1}{2}E_{(+)}^t J_{(+)}, \quad \mathcal{F}_{(-)} = -\frac{1}{2}J_{(-)}^t E_{(-)}^t$$

Summary:

Geometric data: $(M, g, H, J_{(\pm)})$ or $(M, g, J_{(\pm)})$ or $(M, \mathcal{F}_{(\pm)}, J_{(\pm)})$.
 In each case, there is a complete description in terms of a Generalized Kähler potential K . Unlike the Kähler case, the expressions are non-linear in second derivatives of K . E.g.,

$$J_{(+)} = \begin{pmatrix} J & 0 \\ (K_{LR})^{-1}[J, K_{LL}] & (K_{LR})^{-1}JK_{LR} \end{pmatrix}$$

$$g = \Omega[J_{(+)}, J_{(-)}]$$

$$\mathcal{F}_{(+)} = d\lambda_{(+)} , \quad \lambda_{(+)\ell} = iK_R J (K_{LR})^{-1} K_{L\ell} , \dots$$

$$d = 2, N = (2, 2)$$

Algebra:

$$\{\mathbb{D}_\pm, \bar{\mathbb{D}}_\pm\} = i\partial_\pm$$

Constrained superfields:

$$\begin{aligned}\bar{\mathbb{D}}_\pm \phi^a &= 0, \\ \bar{\mathbb{D}}_+ \chi^{a'} &= \mathbb{D}_- \chi^{a'} = 0, \\ \bar{\mathbb{D}}_+ \mathbb{X}^\ell &= 0, \\ \bar{\mathbb{D}}_- \mathbb{X}^r &= 0.\end{aligned}$$

Notation: $c := a, \bar{a}$, $t := a', \bar{a}'$, $L := \ell, \bar{\ell}$, $R := r, \bar{r}$.

The (2, 2) formulation uses the generalized Kähler Potential.

$$S = \int \mathbb{D}_+ \bar{\mathbb{D}}_+ \mathbb{D}_- \bar{\mathbb{D}}_- K(\phi^c, \chi^t, \mathbb{X}^L, \mathbb{X}^R)$$

$$K \rightarrow K(\mathbb{X}^L, \mathbb{X}^R)$$

Reduction to (2, 1) superspace

$$\mathbb{D}_- =: D_- - iQ_- , \quad \mathbb{X}| =: X , \quad Q_- \mathbb{X}^L| =: \Psi_-^L$$

$$S = \int \mathbb{D}_+ \bar{\mathbb{D}}_+ D_- \left(K_L \Psi_-^L + K_R J D_- X^R \right)$$

$$S = i \int \mathbb{D}_+ \bar{\mathbb{D}}_+ D_- (\lambda_{(+)\alpha} D_- \varphi^\alpha + \text{c.c.})$$

which uses the “Liouville form” ($\mathcal{F}_{(+)} = d\lambda_{(+)}$)

Reduction to (1, 1) finally yields

$$S = \int D_+ D_- (D_+ X E D_- X) .$$

The reduction goes via $\mathbb{D}_+ =: D_+ - iQ_+$, $Q_+ \mathbb{X}^R| =: \Psi_+^R$ and both the auxiliary spinors Ψ_-^L and Ψ_+^R have been eliminated.

The (1, 1) formulation uses $E = g + B$ directly.