

# BPS Saturated Objects in String Theory: K3 Elliptic Genus and Igusa Cusp Form $\chi_{10}$

Stefan Hohenegger



MAX-PLANCK-GESELLSCHAFT

Max-Planck-Institut für Physik  
Werner-Heisenberg-Institut  
München



Max-Planck-Institut für Physik  
(Werner-Heisenberg-Institut)

16th Sept. 2011

Corfu Summer Institute  
Workshop on Fields and Strings

Work in collaboration with M. Gaberdiel (ETHZ), S. Stieberger (MPI Munich) and R. Volpato (ETHZ)

synopsis of 1006.0221, 1008.3778, 1106.4315, 1108.0323

# BPS Saturated Quantities in String Theory

Studied from different viewpoints in string theory

target space

particular **effective** superstring  
**couplings** (top. amplitudes)

world-sheet

topological quantities in **CFT**  
(twisted correlators, BPS indices)

# BPS Saturated Quantities in String Theory

Studied from different viewpoints in string theory

## target space

- topological amplitudes, BPS saturated couplings

superspace **analyticity properties**

- $\mathcal{N} = 2$  : chirality
  
- $\mathcal{N} = 4$  : G-analyticity in harmonic superspace

## world-sheet

- topological objects in CFT

differential equations:

- holomorphicity eq. [Bershadsky, Cecotti, Ooguri, Vafa 1993](#)  
coupled first order relations  
[Antoniadis, SH, Narain, Sokatchev 2011](#)
- harmonicity relations  
[Berkovits, Vafa 1994](#)  
[Ooguri, Vafa 1995](#)  
[Antoniadis, SH, Narain, Sokatchev 2007](#)

# BPS Saturated Quantities in String Theory

Studied from different viewpoints in string theory

## target space

- topological amplitudes, BPS saturated couplings
- superspace analyticity properties (holomorphicity, harmonicity,...)

**mathematically:** encode top. invariants of internal manifold

## world-sheet

- topological objects in CFT
- differential equations

**mathematically:** related to BPS indices or index-like structures

# BPS Saturated Quantities in String Theory

Studied from different viewpoints in string theory

## target space

- topological amplitudes, BPS saturated couplings
- superspace analyticity properties (holomorphicity, harmonicity,...)
- topological invariants

**physically:** window into dualities, black hole entropy, pheno.

Ooguri, Strominger, Vafa 2004; Antoniadis, SH 2009  
Dabholkar, Denef, Moore, Pioline 2005

## world-sheet

- topological objects in CFT
- differential equations
- BPS indices

**physically:** algebraic properties of the spectrum of BPS states

Harvey, Moore 1995, 1996  
Gaberdiel, SH, Persson 2011; SH, Persson 2011

# BPS Saturated Quantities in String Theory

Studied from different viewpoints in string theory

## target space

- topological amplitudes, BPS saturated couplings
- superspace analyticity properties (holomorphicity, harmonicity,...)
- topological invariants
- dualities, symmetries, physical problems

## world-sheet

- topological objects in CFT
- differential equations
- BPS indices
- algebraic properties of BPS spectrum

Numerous **examples** found in different string setups!

# BPS Saturated Quantities in String Theory

Studied from different viewpoints in string theory

## target space

- topological amplitudes, BPS saturated couplings
- superspace analyticity properties (holomorphicity, harmonicity,...)
- topological invariants
- dualities, symmetries, physical problems

## world-sheet

- topological objects in CFT
- differential equations
- BPS indices
- algebraic properties of BPS spectrum

Numerous **examples** found in different string setups!

This talk focuses on one particular BPS quantity which has received a lot of attention recently: the **elliptic genus of K3**

- 1 Motivation: BPS Saturated Quantities in String Theory
- 2 World-Sheet BPS-like Structures
  - Elliptic Genus of  $K3$
  - Mathieu Moonshine
- 3  $1/4$  BPS Saturated Amplitudes
  - $1/4$  BPS Contribution to the Elliptic Genus
  - Generating Functional: Igusa Cusp Form



# Elliptic Genera of $\mathcal{N} = 2$ SCFT

**Index**-like objects in two-dimensional superconformal theories ( $\mathcal{N} \geq 2$ ) receive contributions only from (a subset of) the BPS states

# Elliptic Genera of $\mathcal{N} = 2$ SCFT

**Index**-like objects in two-dimensional superconformal theories ( $\mathcal{N} \geq 2$ ) receive contributions only from (a subset of) the BPS states

Important example given by the **elliptic genus** Schellekens, Warner 1987; Witten 1987  
Lerche, Nilsson, Schellekens, Warner 1988

$$\phi(\tau, z) := \text{Tr}_{RR} \left[ (-1)^{F+\bar{F}} y^{J_0} q^{L_0 - \frac{c}{24}} \bar{q}^{\bar{L}_0 - \frac{\bar{c}}{24}} \right], \quad \begin{aligned} q &= e^{2\pi i \tau} \\ y &= e^{2\pi i z} \end{aligned}$$

# Elliptic Genera of $\mathcal{N} = 2$ SCFT

**Index**-like objects in two-dimensional superconformal theories ( $\mathcal{N} \geq 2$ ) receive contributions only from (a subset of) the BPS states

Important example given by the **elliptic genus** Schellekens, Warner 1987; Witten 1987  
Lerche, Nilsson, Schellekens, Warner 1988

$$\phi(\tau, z) := \text{Tr}_{RR} \left[ (-1)^{F+\bar{F}} y^{J_0} q^{L_0 - \frac{c}{24}} \bar{q}^{\bar{L}_0 - \frac{\bar{c}}{24}} \right], \quad \begin{aligned} q &= e^{2\pi i \tau} \\ y &= e^{2\pi i z} \end{aligned}$$

Properties:

- holomorphic in  $\tau$  and  $z$  (right moving ground state)

# Elliptic Genera of $\mathcal{N} = 2$ SCFT

**Index**-like objects in two-dimensional superconformal theories ( $\mathcal{N} \geq 2$ ) receive contributions only from (a subset of) the BPS states

Important example given by the **elliptic genus** Schellekens, Warner 1987; Witten 1987  
Lerche, Nilsson, Schellekens, Warner 1988

$$\phi(\tau, z) := \text{Tr}_{RR} \left[ (-1)^{F+\bar{F}} y^{J_0} q^{L_0 - \frac{c}{24}} \bar{q}^{\bar{L}_0 - \frac{\bar{c}}{24}} \right], \quad \begin{aligned} q &= e^{2\pi i \tau} \\ y &= e^{2\pi i z} \end{aligned}$$

Properties:

- holomorphic in  $\tau$  and  $z$  (right moving ground state)
- weak Jacobi form ( $w = 0$ ,  $m = 1$ ) under  $SL(2, \mathbb{Z})$  Kawai, Yamada, Yang 1994

$$\phi\left(\frac{a\tau + b}{c\tau + d}, \frac{z}{c\tau + d}\right) = (c\tau + d)^w e^{2\pi i m \frac{cz^2}{c\tau + d}} \phi(\tau, z)$$

$$\phi(\tau, z + \ell\tau + \ell') = e^{-2\pi i m(\ell^2\tau + 2\ell z)} \phi(\tau, z) \quad \ell, \ell' \in \mathbb{Z}$$

# Elliptic Genera of $\mathcal{N} = 2$ SCFT

**Index**-like objects in two-dimensional superconformal theories ( $\mathcal{N} \geq 2$ ) receive contributions only from (a subset of) the BPS states

Important example given by the **elliptic genus** Schellekens, Warner 1987; Witten 1987  
Lerche, Nilsson, Schellekens, Warner 1988

$$\phi(\tau, z) := \text{Tr}_{RR} \left[ (-1)^{F+\bar{F}} y^{J_0} q^{L_0 - \frac{c}{24}} \bar{q}^{\bar{L}_0 - \frac{\bar{c}}{24}} \right], \quad \begin{aligned} q &= e^{2\pi i \tau} \\ y &= e^{2\pi i z} \end{aligned}$$

Properties:

- holomorphic in  $\tau$  and  $z$  (right moving ground state)
- weak Jacobi form ( $w = 0$ ,  $m = 1$ ) under  $SL(2, \mathbb{Z})$  Kawai, Yamada, Yang 1994

$$\phi\left(\frac{a\tau + b}{c\tau + d}, \frac{z}{c\tau + d}\right) = (c\tau + d)^w e^{2\pi i m \frac{cz^2}{c\tau + d}} \phi(\tau, z)$$

$$\phi(\tau, z + \ell\tau + \ell') = e^{-2\pi i m(\ell^2\tau + 2\ell z)} \phi(\tau, z) \quad \ell, \ell' \in \mathbb{Z}$$

- Fourier Expansion

$$\phi(\tau, z) = \sum_{n \geq 0, \ell \in \mathbb{Z}} c(n, \ell) q^n y^\ell \quad \text{with} \quad c(n, \ell) = (-1)^w c(n, -\ell)$$

# Elliptic Genera of $\sigma$ -models with $\mathcal{N} = (4, 4)$

Elliptic genus encodes basic **topological information** of SCFT-target e.g.

$$\phi(\tau, z = 0) = \chi(\text{target}) \quad (\text{Euler number})$$

# Elliptic Genera of $\sigma$ -models with $\mathcal{N} = (4, 4)$

Elliptic genus encodes basic **topological information** of SCFT-target e.g.

$$\phi(\tau, z = 0) = \chi(\text{target}) \quad (\text{Euler number})$$

In fact, **2-dimensional** CFTs with  $\mathcal{N} = (4, 4)$  SUSY and **central charge**  $c = 6$  are classified through their **elliptic genus**

$$\phi = \text{Tr}_{RR} \left( (-1)^{F+\bar{F}} y^{J_0} q^{L_0 - \frac{1}{4}} \bar{q}^{\bar{L}_0 - \frac{1}{4}} \right) = \begin{cases} 0 & \text{target } \mathbb{T}^4 \\ 8 \sum_{i=2}^4 \frac{\theta_i(\tau, z)^2}{\theta_i(\tau, 0)^2} & \text{target } K3 \end{cases}$$

# Elliptic Genera of $\sigma$ -models with $\mathcal{N} = (4, 4)$

Elliptic genus encodes basic **topological information** of SCFT-target e.g.

$$\phi(\tau, z = 0) = \chi(\text{target}) \quad (\text{Euler number})$$

In fact, **2-dimensional** CFTs with  $\mathcal{N} = (4, 4)$  SUSY and **central charge**  $c = 6$  are classified through their **elliptic genus**

$$\phi = \text{Tr}_{RR} \left( (-1)^{F+\bar{F}} y^{J_0} q^{L_0 - \frac{1}{4}} \bar{q}^{\bar{L}_0 - \frac{1}{4}} \right) = \begin{cases} 0 & \text{target } \mathbb{T}^4 \\ 8 \sum_{i=2}^4 \frac{\theta_i(\tau, z)^2}{\theta_i(\tau, 0)^2} & \text{target } K3 \end{cases}$$

Decomposition in  $\mathcal{N} = 4$  characters for  $K3$  target space [Eguchi, Taormina 1988](#)  
[Eguchi, Hikami 2009](#)

$$\phi_{K3} = 24 \text{ch}_{h=\frac{1}{4}, \ell=0} + \sum_{n=0}^{\infty} A_n \text{ch}_{h=n+\frac{1}{4}, \ell=\frac{1}{2}}, \quad \begin{aligned} \text{ch}_{h, \ell=\frac{1}{2}} &:= q^{h-\frac{3}{8}} \frac{\theta_1(\tau, z)^2}{\eta(\tau)^3} \\ \text{ch}_{h=\frac{1}{4}, \ell=0} &:= \frac{\theta_1(\tau, z)^2}{\eta(\tau)^3} \mu(\tau, z) \end{aligned}$$

where  $\mu(\tau, z)$  is an Appell-Lerch sum (**mock modular theta function**)



# Mathieu Moonshine

$$\phi_{K3} = 24 \operatorname{ch}_{h=\frac{1}{4}, \ell=0} + \sum_{n=0}^{\infty} A_n \operatorname{ch}_{h=n+\frac{1}{4}, \ell=\frac{1}{2}} \quad \begin{array}{l} A_1 = 45, \quad A_2 = 231, \\ A_3 = 770, \quad A_4 = 2277 \end{array}$$

# Mathieu Moonshine

$$\phi_{K3} = 24 \text{ch}_{h=\frac{1}{4}, \ell=0} + \sum_{n=0}^{\infty} A_n \text{ch}_{h=n+\frac{1}{4}, \ell=\frac{1}{2}} \quad \begin{array}{l} A_1 = 45, \quad A_2 = 231, \\ A_3 = 770, \quad A_4 = 2277 \end{array}$$

First few  $A_n$  are the dimensions of the smallest irreducible representations of the largest Mathieu group  $M_{24}$  (**Moonshine!**). [Eguchi, Ooguri, Tachikawa 2010](#)

Suggests  $M_{24}$ -symmetry of (sub)space of BPS states contributing to  $\phi_{K3}$

$$\mathcal{H}_{\phi_{K3}}^{\text{BPS}} = \bigoplus_n (H_n \otimes \mathcal{H}_n^{\mathcal{N}=4}) \quad \begin{array}{l} H_n \dots M_{24} \text{ repr.} \\ \mathcal{H}_n^{\mathcal{N}=4} \dots \mathcal{N} = 4 \text{ irrep.} \end{array}$$

# Mathieu Moonshine

$$\phi_{K3} = 24 \operatorname{ch}_{h=\frac{1}{4}, \ell=0} + \sum_{n=0}^{\infty} A_n \operatorname{ch}_{h=n+\frac{1}{4}, \ell=\frac{1}{2}} \quad \begin{array}{l} A_1 = 45, \quad A_2 = 231, \\ A_3 = 770, \quad A_4 = 2277 \end{array}$$

First few  $A_n$  are the dimensions of the smallest irreducible representations of the largest Mathieu group  $M_{24}$  (**Moonshine!**). [Eguchi, Ooguri, Tachikawa 2010](#)

Suggests  $M_{24}$ -symmetry of (sub)space of BPS states contributing to  $\phi_{K3}$

$$\mathcal{H}_{\phi_{K3}}^{\text{BPS}} = \bigoplus_n (H_n \otimes \mathcal{H}_n^{\mathcal{N}=4}) \quad \begin{array}{l} H_n \dots M_{24} \text{ repr.} \\ \mathcal{H}_n^{\mathcal{N}=4} \dots \mathcal{N} = 4 \text{ irrep.} \end{array}$$

Leads to **conjecture**

$$\phi_{K3} = (\dim H_{00}) \operatorname{ch}_{h=\frac{1}{4}, l=0} - (\dim H_0) \operatorname{ch}_{h=\frac{1}{4}, l=\frac{1}{2}} + \sum_{n=1}^{\infty} (\dim H_n) \operatorname{ch}_{h=n+\frac{1}{4}, l=\frac{1}{2}}$$

# Mathieu Moonshine

$$\phi_{K3} = 24 \text{ch}_{h=\frac{1}{4}, \ell=0} + \sum_{n=0}^{\infty} A_n \text{ch}_{h=n+\frac{1}{4}, \ell=\frac{1}{2}} \quad \begin{array}{l} A_1 = 45, \quad A_2 = 231, \\ A_3 = 770, \quad A_4 = 2277 \end{array}$$

First few  $A_n$  are the dimensions of the smallest irreducible representations of the largest Mathieu group  $M_{24}$  (**Moonshine!**). [Eguchi, Ooguri, Tachikawa 2010](#)

Suggests  $M_{24}$ -symmetry of (sub)space of BPS states contributing to  $\phi_{K3}$

$$\mathcal{H}_{\phi_{K3}}^{\text{BPS}} = \bigoplus_n (H_n \otimes \mathcal{H}_n^{\mathcal{N}=4}) \quad \begin{array}{l} H_n \dots M_{24} \text{ repr.} \\ \mathcal{H}_n^{\mathcal{N}=4} \dots \mathcal{N} = 4 \text{ irrep.} \end{array}$$

Leads to **conjecture**

$$\phi_{K3} = (\dim H_{00}) \text{ch}_{h=\frac{1}{4}, l=0} - (\dim H_0) \text{ch}_{h=\frac{1}{4}, l=\frac{1}{2}} + \sum_{n=1}^{\infty} (\dim H_n) \text{ch}_{h=n+\frac{1}{4}, l=\frac{1}{2}}$$

This might be a '**supersymmetric** generalization' of monstrous moonshine!

[McKay, Thompson 1978](#)

# Twining Genera

Idea of checking: **twining genera**: ( $g \in \mathbb{M}_{24}$ , 26 conjugacy classes)

$$\phi_g := \text{Tr}_{H_{00}}(g) \text{ch}_{h=\frac{1}{4}, l=0} - \text{Tr}_{H_0}(g) \text{ch}_{h=\frac{1}{4}, l=\frac{1}{2}} + \sum_{n=1}^{\infty} \text{Tr}_{H_n}(g) \text{ch}_{h=n+\frac{1}{4}, l=\frac{1}{2}}$$

# M<sub>24</sub> Character Table

1A	2A	3A	5A	4B	7A	7B	8A	6A	11A	15A	15B	14A	14B	23A	23B	12B	6B	4C	3B	2B	10A	21A	21B	4A	12A	
1	1	1	1	1	1	1	1	1	1	1	1	1	1	1	1	1	1	1	1	1	1	1	1	1	1	
23	7	5	3	3	2	2	1	1	1	0	0	0	0	0	0	-1	-1	-1	-1	-1	-1	-1	-1	-1	-1	
252	28	9	2	4	0	0	0	1	-1	-1	-1	0	0	-1	-1	0	0	0	12	2	0	0	0	4	4	
253	13	10	3	1	1	1	-1	-2	0	0	0	-1	-1	0	0	1	1	1	-11	-1	1	1	1	-3	0	
1771	-21	16	1	-5	0	0	-1	0	0	1	1	0	0	0	0	-1	-1	-1	7	11	1	0	0	3	0	
3520	64	10	0	0	-1	-1	0	-2	0	0	0	1	1	1	1	0	0	0	-8	0	0	-1	-1	0	0	
45	-3	0	0	1	$e_7^+$	$e_7^-$	-1	0	1	0	0	$-e_7^+$	$-e_7^-$	-1	-1	1	-1	1	3	5	0	$e_7^-$	$e_7^+$	-3	0	
45	-3	0	0	1	$e_7^-$	$e_7^+$	-1	0	1	0	0	$-e_7^-$	$-e_7^+$	-1	-1	1	-1	1	3	5	0	$e_7^+$	$e_7^-$	-3	0	
990	-18	0	0	2	$e_7^+$	$e_7^-$	0	0	0	0	0	$e_7^+$	$e_7^-$	1	1	1	-1	-2	3	-10	0	$e_7^-$	$e_7^+$	6	0	
990	-18	0	0	2	$e_7^-$	$e_7^+$	0	0	0	0	0	$e_7^-$	$e_7^+$	1	1	1	-1	-2	3	-10	0	$e_7^+$	$e_7^-$	6	0	
1035	-21	0	0	3	$2e_7^+$	$2e_7^-$	-1	0	1	0	0	0	0	0	0	-1	1	-1	-3	-5	0	$-e_7^-$	$-e_7^+$	3	0	
1035	-21	0	0	3	$2e_7^-$	$2e_7^+$	-1	0	1	0	0	0	0	0	0	-1	1	-1	-3	-5	0	$-e_7^+$	$-e_7^-$	3	0	
1035	27	0	0	-1	-1	-1	1	0	1	0	0	-1	-1	0	0	0	2	3	6	35	0	-1	-1	3	0	
231	7	-3	1	-1	0	0	-1	1	0	$e_{15}^+$	$e_{15}^-$	0	0	1	1	0	0	3	0	-9	1	0	0	-1	-1	
231	7	-3	1	-1	0	0	-1	1	0	$e_{15}^-$	$e_{15}^+$	0	0	1	1	0	0	3	0	-9	1	0	0	-1	-1	
770	-14	5	0	-2	0	0	0	1	0	0	0	0	0	$e_{23}^+$	$e_{23}^-$	1	1	-2	-7	10	0	0	0	2	-1	
770	-14	5	0	-2	0	0	0	1	0	0	0	0	0	$e_{23}^-$	$e_{23}^+$	1	1	-2	-7	10	0	0	0	2	-1	
483	35	6	-2	3	0	0	-1	2	-1	1	1	0	0	0	0	0	0	3	0	3	-2	0	0	3	0	
1265	49	5	0	1	-2	-2	1	1	0	0	0	0	0	0	0	0	0	-3	8	-15	0	1	1	-7	-1	
2024	8	-1	-1	0	1	1	0	-1	0	-1	-1	1	1	0	0	0	0	0	8	24	-1	1	1	8	-1	
2277	21	0	-3	1	2	2	-1	0	0	0	0	0	0	0	0	0	2	-3	6	-19	1	-1	-1	-3	0	
3312	48	0	-3	0	1	1	0	0	1	0	0	-1	-1	0	0	0	-2	0	-6	16	1	1	1	0	0	
5313	49	-15	3	-3	0	0	-1	1	0	0	0	0	0	0	0	0	0	-3	0	9	-1	0	0	1	1	
5796	-28	-9	1	4	0	0	0	-1	-1	1	1	0	0	0	0	0	0	0	0	36	1	0	0	-4	-1	
5544	-56	9	-1	0	0	0	0	1	0	-1	-1	0	0	1	1	0	0	0	0	24	-1	0	0	0	-8	1
10395	-21	0	0	-1	0	0	1	0	0	0	0	0	0	-1	-1	0	0	3	0	-45	0	0	0	3	0	

# Twining Genera

Idea of checking: **twining genera**: ( $g \in \mathbb{M}_{24}$ , 26 conjugacy classes)

$$\phi_g := \text{Tr}_{H_{00}}(g) \text{ch}_{h=\frac{1}{4}, l=0} - \text{Tr}_{H_0}(g) \text{ch}_{h=\frac{1}{4}, l=\frac{1}{2}} + \sum_{n=1}^{\infty} \text{Tr}_{H_n}(g) \text{ch}_{h=n+\frac{1}{4}, l=\frac{1}{2}}$$

CFT suggests that  $\phi_g$  is a weak Jacobi form of weight 0 and index 1 under some **congruence subgroup** of  $SL(2, \mathbb{Z})$  (possibly with multiplier system)

# Twining Genera

Idea of checking: **twining genera**: ( $g \in \mathbb{M}_{24}$ , 26 conjugacy classes)

$$\phi_g := \text{Tr}_{H_{00}}(g) \text{ch}_{h=\frac{1}{4}, l=0} - \text{Tr}_{H_0}(g) \text{ch}_{h=\frac{1}{4}, l=\frac{1}{2}} + \sum_{n=1}^{\infty} \text{Tr}_{H_n}(g) \text{ch}_{h=n+\frac{1}{4}, l=\frac{1}{2}}$$

CFT suggests that  $\phi_g$  is a weak Jacobi form of weight 0 and index 1 under some **congruence subgroup** of  $SL(2, \mathbb{Z})$  (possibly with multiplier system)

Conjectured transformation for  $g \in \mathbb{M}_{24}$  (order  $N$ ) [Gaberdiel, SH, Volpato 2010](#)

$$\phi_g \left( \frac{a\tau + b}{c\tau + d}, \frac{z}{c\tau + d} \right) = e^{\frac{2\pi i c d}{N h}} e^{\frac{2\pi i c z^2}{c\tau + d}} \phi_g(\tau, z), \quad \begin{pmatrix} a & b \\ c & d \end{pmatrix} \in \Gamma_0(N)$$



# Twining Genera

Idea of checking: **twining genera**: ( $g \in \mathbb{M}_{24}$ , 26 conjugacy classes)

$$\phi_g := \text{Tr}_{H_{00}}(g) \text{ch}_{h=\frac{1}{4}, l=0} - \text{Tr}_{H_0}(g) \text{ch}_{h=\frac{1}{4}, l=\frac{1}{2}} + \sum_{n=1}^{\infty} \text{Tr}_{H_n}(g) \text{ch}_{h=n+\frac{1}{4}, l=\frac{1}{2}}$$

CFT suggests that  $\phi_g$  is a weak Jacobi form of weight 0 and index 1 under some **congruence subgroup** of  $SL(2, \mathbb{Z})$  (possibly with multiplier system)

Conjectured transformation under  $g \in \mathbb{M}_{24}$  (order  $N$ ) [Gaberdiel, SH, Volpato 2010](#)

$$\phi_g \left( \frac{a\tau + b}{c\tau + d}, \frac{z}{c\tau + d} \right) = e^{\frac{2\pi i c d}{N h}} e^{\frac{2\pi i c z^2}{c\tau + d}} \phi_g(\tau, z), \quad \begin{pmatrix} a & b \\ c & d \end{pmatrix} \in \Gamma_0(N)$$

$$\Gamma_0(N) := \left\{ \begin{pmatrix} a & b \\ c & d \end{pmatrix} \in SL(2, \mathbb{Z}) \mid c \equiv 0 \pmod{N} \right\}, \quad h \in \mathbb{N}$$

$h$  determines multiplier system and differs for various conjugacy classes

$$h = 1 \iff g \in \mathbb{M}_{23} \subset \mathbb{M}_{24}$$

# Status of Checks of $M_{24}$ Moonshine

Available/conjectural information about twining genera

- first few Fourier coefficients [Eguchi, Ooguri, Tachikawa 2010](#)
- conjectured modular properties (particularly multiplier system)

allows for **closed expressions** for all 26 conjugacy classes

[Cheng 2010](#)  
[Gaberdiel, SH, Volpato 2010](#)  
[Eguchi, Hikami 2010](#)

$$\phi_{2A}(\tau, z) = 8 \frac{\theta_2(\tau, z)^2}{\theta_2(\tau, 0)^2}, \quad \phi_{2B}(\tau, z) = -2\theta_4(2\tau, 0)^2 \frac{\theta_1(\tau, z)^2}{\eta(\tau)^6}, \quad \dots$$

# Status of Checks of $\mathbb{M}_{24}$ Moonshine

Available/conjectural information about twining genera

- first few Fourier coefficients [Eguchi, Ooguri, Tachikawa 2010](#)
- conjectured modular properties (particularly multiplier system)

allows for **closed expressions** for all 26 conjugacy classes

[Cheng 2010](#)  
[Gaberdiel, SH, Volpato 2010](#)  
[Eguchi, Hikami 2010](#)

$$\phi_{2A}(\tau, z) = 8 \frac{\theta_2(\tau, z)^2}{\theta_2(\tau, 0)^2}, \quad \phi_{2B}(\tau, z) = -2\theta_4(2\tau, 0)^2 \frac{\theta_1(\tau, z)^2}{\eta(\tau)^6}, \quad \dots$$

Knowledge of the twining genera in turn allows to **uniquely** decompose all  $H_n$  in terms of irreducible representations of  $\mathbb{M}_{24}$ .

For  $n \leq 1000$ :  $H_n$  valid  $\mathbb{M}_{24}$  representations (decomposition into irreps with non-negative integer coefficients) [Gaberdiel, SH, Volpato 2010](#)  
[Eguchi, Hikami 2010](#)

# Status of Checks of $\mathbb{M}_{24}$ Moonshine

Available/conjectural information about twining genera

- first few Fourier coefficients [Eguchi, Ooguri, Tachikawa 2010](#)
- conjectured modular properties (particularly multiplier system)

allows for **closed expressions** for all 26 conjugacy classes

[Cheng 2010](#)  
[Gaberdiel, SH, Volpato 2010](#)  
[Eguchi, Hikami 2010](#)

$$\phi_{2A}(\tau, z) = 8 \frac{\theta_2(\tau, z)^2}{\theta_2(\tau, 0)^2}, \quad \phi_{2B}(\tau, z) = -2\theta_4(2\tau, 0)^2 \frac{\theta_1(\tau, z)^2}{\eta(\tau)^6}, \quad \dots$$

Knowledge of the twining genera in turn allows to **uniquely** decompose all  $H_n$  in terms of irreducible representations of  $\mathbb{M}_{24}$ .

For  $n \leq 1000$ :  $H_n$  valid  $\mathbb{M}_{24}$  representations (decomposition into irreps with non-negative integer coefficients) [Gaberdiel, SH, Volpato 2010](#)  
[Eguchi, Hikami 2010](#)

**No proof, but very strong evidence in favor of the conjecture!**

# Explanation of Mathieu-Moonshine?

**Challenge:** Given the massive evidence in favour of the conjecture, what is the reason for the appearance of  $M_{24}$ ?

# Explanation of Mathieu-Moonshine?

**Challenge:** Given the massive evidence in favour of the conjecture, what is the reason for the appearance of  $M_{24}$ ?

So far no **satisfactory** explanation, but progress in partial understanding

- symplectic automorph. of K3 are subgroup of  $M_{23} \subset M_{24}$  Mukai 1988  
Kondo 1998
- at given point in moduli space generically no enhancement to subgroup of  $M_{24}$  through **quantum symmetries** of K3  $\sigma$ -model, instead for distinct cases.  $\mathcal{N} = 4$  preserving sym. Gaberdiel, SH, Volpato 2011
  - ▶ subgroup of  $G' \cdot G''$ , where  $G' \subset \mathbb{Z}_2^{11}$  and  $G'' \subset M_{24}$
  - ▶ subgroup of  $5^{1+2} \cdot \mathbb{Z}_4$
  - ▶ subgroup of  $\mathbb{Z}_3^4 \cdot A_6$
  - ▶ subgroup of  $3^{1+4} \cdot \mathbb{Z}_2 \cdot G''$ , where  $G''$  is either trivial,  $\mathbb{Z}_2$ ,  $\mathbb{Z}_2^2$  or  $\mathbb{Z}_4$ .

( $p^{1+2n}$  is the extra special group of order  $p^{1+2n}$ )
- all symmetries subgroups of Conway  $Co_1$

**New insights (different approach?) needed to explain Mathieu moonshine.**

# Connection to Target Space Physics

Attempts at understanding so far have focused on the **world-sheet**.

**Question:** How much symmetry maps over to target space physics?

- Any connection of  $\mathbb{M}_{24}$  to **algebra of BPS states**? Harvey, Moore 1995, 1996  
Gaberdiel, SH, Persson 2011  
SH, Persson 2011
- Which target space structures are sensible to  $\mathbb{M}_{24}$ ?
- Can we think of  $\mathbb{M}_{24}$  as symmetry of some effective action couplings?
- Are there  $\mathcal{N} = 4$  (1/4 BPS) topological amplitudes which capture the elliptic genus?

# Connection to Target Space Physics

Attempts at understanding so far have focused on the **world-sheet**.

**Question:** How much symmetry maps over to target space physics?

- Any connection of  $\mathbb{M}_{24}$  to **algebra of BPS states**? Harvey, Moore 1995, 1996  
Gaberdiel, SH, Persson 2011  
SH, Persson 2011
- Which target space structures are sensible to  $\mathbb{M}_{24}$ ?
- Can we think of  $\mathbb{M}_{24}$  as symmetry of some effective action couplings?
- Are there  $\mathcal{N} = 4$  (1/4 BPS) topological amplitudes which capture the elliptic genus?

Will report on first steps in answering these questions in the following.

*I.e.* I will discuss a particular **BPS saturated amplitude** in **type II** string theory on  $K3 \times T^2$  which is related to the elliptic genus of  $K3$ .

( $\mathcal{N} = 2$  results for heterotic on  $K3 \times T^2$  Kawai 1995)



# Massive External States

Consider (carefully arranged) amplitude of the form [SH, Stieberger 2011](#)

$$\mathcal{F}_N = \left\langle \int d^2 z_1 V_R(z_1) \int d^2 z_2 V_R(z_2) \prod_{a=1}^N \int d^2 x_a V_M^{(-1,-1)}(p_a, x_a) \prod_{b=1}^N V_{\text{PCO}} \right\rangle$$

# Massive External States

Consider (carefully arranged) amplitude of the form [SH, Stieberger 2011](#)

$$\mathcal{F}_N = \left\langle \int d^2 z_1 V_R(z_1) \int d^2 z_2 V_R(z_2) \prod_{a=1}^N \int d^2 x_a V_M^{(-1,-1)}(p_a, x_a) \prod_{b=1}^N V_{\text{PCO}} \right\rangle$$

Key ingredients are **massive** external scalar fields (**first massive level**)

$$V_M^{(-1,-1)}(p, x) = : e^{-\varphi} \psi_3 \partial H(x) e^{-\tilde{\varphi}} \tilde{\psi}_3 \bar{\partial} X_3(\bar{x}) e^{ipX} : .$$

$X_3 \dots$  torus coordinate                       $\varphi \dots$  bosonization of superghost

$\psi_3, \tilde{\psi}_3 \dots$  free torus fermion       $\partial H(x) \dots$  bosonization of  $J_{K3}$

# Massive External States

Consider (carefully arranged) amplitude of the form [SH, Stieberger 2011](#)

$$\mathcal{F}_N = \left\langle \int d^2 z_1 V_R(z_1) \int d^2 z_2 V_R(z_2) \prod_{a=1}^N \int d^2 x_a V_M^{(-1,-1)}(p_a, x_a) \prod_{b=1}^N V_{\text{PCO}} \right\rangle$$

Key ingredients are **massive** external scalar fields (**first massive level**)

$$V_M^{(-1,-1)}(p, x) = : e^{-\varphi} \psi_3 \partial H(x) e^{-\tilde{\varphi}} \tilde{\psi}_3 \bar{\partial} X_3(\bar{x}) e^{ipX} : .$$

$X_3 \dots$  torus coordinate                       $\varphi \dots$  bosonization of superghost

$\psi_3, \tilde{\psi}_3 \dots$  free torus fermion       $\partial H(x) \dots$  bosonization of  $J_{K3}$

Explicit result

$$\begin{aligned} \mathcal{F}_N &= \int \frac{d^2 \tau}{\tau_2} \tau_2^{2N} \left\langle \prod_{a=1}^N \partial H(x_a) \right\rangle_{K3} \sum_{\Gamma^{2,2}} (P_L)^N (P_R)^{2N} q^{\frac{1}{2}|P_L|^2} \bar{q}^{\frac{1}{2}|P_R|^2} = \\ &= (\mathcal{D}_{\bar{U}})^N \int \frac{d^2 \tau}{\tau_2} \left[ \frac{\partial^N}{\partial z^N} \phi_{K3}(\tau, z) \right]_{z=0} \sum_{\Gamma^{2,2}} (\tau_2 P_R)^N q^{\frac{1}{2}|P_L|^2} \bar{q}^{\frac{1}{2}|P_R|^2} \end{aligned}$$

# Analysis of Amplitude

Consider 1-loop expression in more detail

$$\int \frac{d^2\tau}{\tau_2} \left[ \frac{\partial^N}{\partial z^N} \phi_{K3}(\tau, z) \right]_{z=0} \sum_{\Gamma^{2,2}} (\tau_2 P_R)^N q^{\frac{1}{2}|P_L|^2} \bar{q}^{\frac{1}{2}|P_R|^2}$$

- related to the  $N$ -th derivative of the **full** elliptic genus

# Analysis of Amplitude

Consider 1-loop expression in more detail

$$\int \frac{d^2\tau}{\tau_2} \left[ \frac{\partial^N}{\partial z^N} \phi_{K3}(\tau, z) \right]_{z=0} \sum_{\Gamma^{2,2}} (\tau_2 P_R)^N q^{\frac{1}{2}|P_L|^2} \bar{q}^{\frac{1}{2}|P_R|^2}$$

- related to the  $N$ -th derivative of the **full** elliptic genus
- right moving sector trivial, only contributes bosonic torus zero modes resulting in Narain momenta of  $\Gamma^{2,2}$

# Analysis of Amplitude

Consider 1-loop expression in more detail

$$\int \frac{d^2\tau}{\tau_2} \left[ \frac{\partial^N}{\partial z^N} \phi_{K3}(\tau, z) \right]_{z=0} \sum_{\Gamma^{2,2}} (\tau_2 P_R)^N q^{\frac{1}{2}|P_L|^2} \bar{q}^{\frac{1}{2}|P_R|^2}$$

- related to the  $N$ -th derivative of the **full** elliptic genus
- right moving sector trivial, only contributes bosonic torus zero modes resulting in Narain momenta of  $\Gamma^{2,2}$
- powers exactly right to ensure world-sheet **modular invariance** (required for consistency of the 1-loop amplitude)

# Analysis of Amplitude

Consider 1-loop expression in more detail

$$\int \frac{d^2\tau}{\tau_2} \left[ \frac{\partial^N}{\partial z^N} \phi_{K3}(\tau, z) \right]_{z=0} \sum_{\Gamma^{2,2}} (\tau_2 P_R)^N q^{\frac{1}{2}|P_L|^2} \bar{q}^{\frac{1}{2}|P_R|^2}$$

- related to the  $N$ -th derivative of the **full** elliptic genus
- right moving sector trivial, only contributes bosonic torus zero modes resulting in Narain momenta of  $\Gamma^{2,2}$
- powers exactly right to ensure world-sheet **modular invariance** (required for consistency of the 1-loop amplitude)

Can we make the symmetry properties of this expression more manifest?

- compute the world-sheet torus integral
- make T-duality invariance manifest by writing it in terms of modular forms of  $SL(2, \mathbb{Z})_{\mathcal{T}} \times SL(2, \mathbb{Z})_{\mathcal{U}}$
- any additional (hidden?) structures or symmetries?

# Generating Functional

Define generating functional

$$\begin{aligned}\mathbb{F}(\lambda, T, U) &:= \int \frac{d^2\tau}{\tau_2} \sum_{N=0}^{\infty} \frac{\lambda^N (\partial_z^N \phi_{K3}(\tau, z))_{z=0}}{(2T_2 U_2)^{N/2} N!} \sum_{\Gamma^{2,2}} (\tau_2 P_R)^N q^{\frac{1}{2}|P_L|^2} \bar{q}^{\frac{1}{2}|P_R|^2} \\ &= \int \frac{d^2\tau}{\tau_2} \sum_{\Gamma^{2,2}} \phi_{K3} \left( \tau, \frac{\lambda \tau_2 P_R}{\sqrt{2T_2 U_2}} \right) q^{\frac{1}{2}|P_L|^2} \bar{q}^{\frac{1}{2}|P_R|^2}\end{aligned}$$



# Generating Functional

Define generating functional

$$\begin{aligned}\mathbb{F}(\lambda, T, U) &:= \int \frac{d^2\tau}{\tau_2} \sum_{N=0}^{\infty} \frac{\lambda^N (\partial_z^N \phi_{K3}(\tau, z))_{z=0}}{(2T_2 U_2)^{N/2} N!} \sum_{\Gamma^{2,2}} (\tau_2 P_R)^N q^{\frac{1}{2}|P_L|^2} \bar{q}^{\frac{1}{2}|P_R|^2} \\ &= \int \frac{d^2\tau}{\tau_2} \sum_{\Gamma^{2,2}} \phi_{K3} \left( \tau, \frac{\lambda \tau_2 P_R}{\sqrt{2T_2 U_2}} \right) q^{\frac{1}{2}|P_L|^2} \bar{q}^{\frac{1}{2}|P_R|^2}\end{aligned}$$

Use Fourier-expansion of the elliptic genus

$$\phi_{K3}(\tau, z) = \sum_{n, \ell} c(n, \ell) q^n e^{2\pi i z \ell}, \quad \text{with} \quad c(n, \ell) = c\left(n - \frac{\ell^2}{4}\right)$$

with the first few coefficients given explicitly

$$\begin{aligned}c(0) &= 20, & c(1) &= 216, & c(2) &= 1616, & c(3) &= 8032 \\ c(-1/4) &= 2, & c(3/4) &= -128, & c(7/4) &= -1026, & c(11/4) &= -5504\end{aligned}$$

# World-sheet Torus Integral

The world-sheet torus integral for the generating functional can be directly evaluated ( $Y = T_2 U_2 - \lambda^2$ ) [Dixon, Kaplunovsky, Louis 1991](#)  
[Harvey, Moore 1995](#)

$$\mathbb{F}(\lambda, T, U) \sim \ln \left[ Y^{10} \left| e^{2\pi i(T+U+i\lambda)} \prod_{r,n',\ell>0} \left( 1 - e^{2\pi i(rT+n'U+i\ell\lambda)} \right)^{c(n'r,\ell)} \right| \right]$$

# World-sheet Torus Integral

The world-sheet torus integral for the generating functional can be directly evaluated ( $Y = T_2 U_2 - \lambda^2$ ) Dixon, Kaplunovsky, Louis 1991  
Harvey, Moore 1995

$$\mathbb{F}(\lambda, T, U) \sim \ln \left[ Y^{10} \underbrace{\left| e^{2\pi i(T+U+i\lambda)} \prod_{r,n',\ell>0} \left(1 - e^{2\pi i(rT+n'U+i\ell\lambda)}\right)^{c(n'r,\ell)} \right|}_{=\chi_{10}(T,U;\lambda)} \right]$$

Igusa cusp form of weight 10 of  $Sp(4, \mathbb{Z})$ .

(Surprising) results

- generating functional is manifestly invariant under the T-duality group
- expansion parameter  $\lambda$  enhances invariance to  $Sp(4, \mathbb{Z})$
- correct behaviour in the limit  $\lambda \rightarrow 0$

# Conclusions

In this talk I discussed the elliptic genus of  $K3$  from several (complementary) perspectives

- **world-sheet/CFT**: conjectured action of  $M_{24}$  on the space of states contributing to  $\phi_{K3}$  (**Mathieu moonshine**)
  - ▶ overwhelming 'observational' evidence
  - ▶ no satisfactory explanation
- **space-time**: identified 1/4 BPS saturated one-loop coupling which is related to  $\phi_{K3}$  in type II string theory on  $K3 \times T^2$

# Conclusions

In this talk I discussed the elliptic genus of  $K3$  from several (complementary) perspectives

- **world-sheet/CFT**: conjectured action of  $\mathbb{M}_{24}$  on the space of states contributing to  $\phi_{K3}$  (**Mathieu moonshine**)
  - ▶ overwhelming 'observational' evidence
  - ▶ no satisfactory explanation
- **space-time**: identified 1/4 BPS saturated one-loop coupling which is related to  $\phi_{K3}$  in type II string theory on  $K3 \times T^2$

Several open questions

- **world-sheet/CFT**:
  - ▶ Why  $\mathbb{M}_{24}$ ? Could there be a larger group, *i.e.*  $Co_1$ ?
  - ▶ Is there an analogue of the 'Monster Module'?
- **space-time**:
  - ▶ Can we make the  $\mathbb{M}_{24}$  action visible in target space?
  - ▶ How does the dual heterotic side look like?