Higgs pair production in the Composite Higgs model

Ramona Gröber in coll. with Margarete Mühlleitner, see JHEP 06 (2011) 020 [arXiv: 1012.1562] | 05.09.2011

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Motivation

- Hierarchy problem
  $\rightarrow$ in Composite Higgs models: Higgs mass generated dynamically through loops of SM fermions and gauge bosons
  $\rightarrow$ Higgs mass is naturally light

- Higgs pair production:
  Measurement of triple Higgs coupling $\rightarrow$ reconstruction of Higgs potential

The Composite Higgs model

- Higgs boson is a bound state from a strongly-interacting sector [Kaplan, Georgi, Galison, Dugan; Dimopoulos, Preskill; Banks]
- Higgs boson emerges as a pseudo-Nambu-Goldstone boson from a global symmetry

$$G \xrightarrow{f} H \supset H_0 = SU(2)_L \times U(1)_Y$$

- G is explicitely broken by couplings of SM fields to strong sector $\rightarrow$ loops of SM fermions and gauge bosons generate Higgs potential
- Higgs boson is naturally separated by a mass gap from the other usual resonances
- In an effective low-energy description only the Higgs couplings to the SM fields are modified and depend on $\xi = \frac{v^2}{f^2}$ with $\xi \in [0, 1]$ [Giudice,Grojean,Pomarol,Rattazzi]

Minimal Composite Higgs models [Agashe, Contino, Pomarol; Contino, Da Rold, Pomarol]

- Realized as 5D theory on AdS spacetime $\rightarrow$ holographic interpretation in terms of 4D strongly coupled field theories
- Global symmetry $G = SO(5) \times U(1)_X$ broken down to
  - $SO(4) \times U(1)_X$ on IR-brane
  - $SU(2)_L \times U(1)_Y$ on UV-brane
- Higgs-fermion and Higgs self-couplings depend on the representation of the fermions:
  - spinorial representation: MCHM_4
  - fundamental representation: MCHM_5

Feynman rules

	MCHM_4	MCHM_5
HVV	$\sqrt{1-\xi}$	$\sqrt{1-\xi}$
$HHVV$	$1-2\xi$	$1-2\xi$
HHH	$\sqrt{1-\xi}$	$\frac{1-2\xi}{\sqrt{1-\xi}}$
$Hf\bar{f}$	$\sqrt{1-\xi}$	$\frac{1-2\xi}{\sqrt{1-\xi}}$

New coupling: $g_{HHf\bar{f}} = \frac{m_f}{v^2}\xi$ in MCHM_4

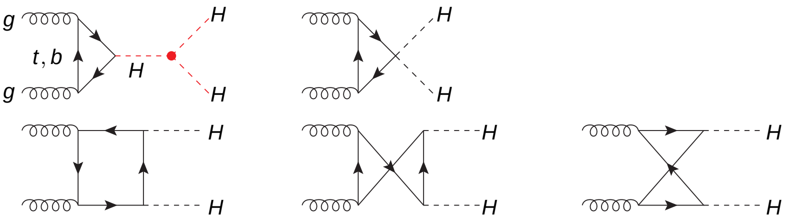
and $g_{HHf\bar{f}} = \frac{m_f}{v^2}4\xi$ in MCHM_5

In SM limit $\xi \to 0$ the new coupling vanishes.

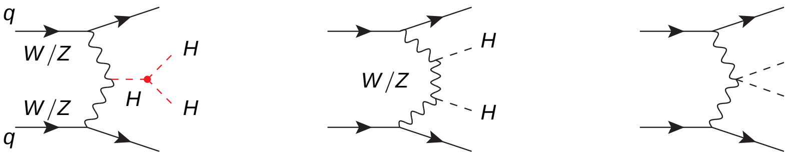
Higgs pair production at the LHC

- Higgs self-couplings are necessary to investigate Higgs potential.
- In multi-Higgs-production Higgs self-couplings can be measured.
- Most important processes in SM at the LHC:
  - Gluon fusion
  - Vector boson fusion
  - Double Higgs-strahlung

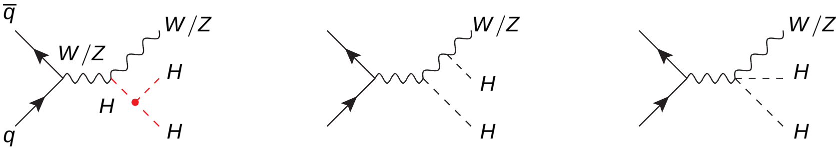
Gluon fusion:

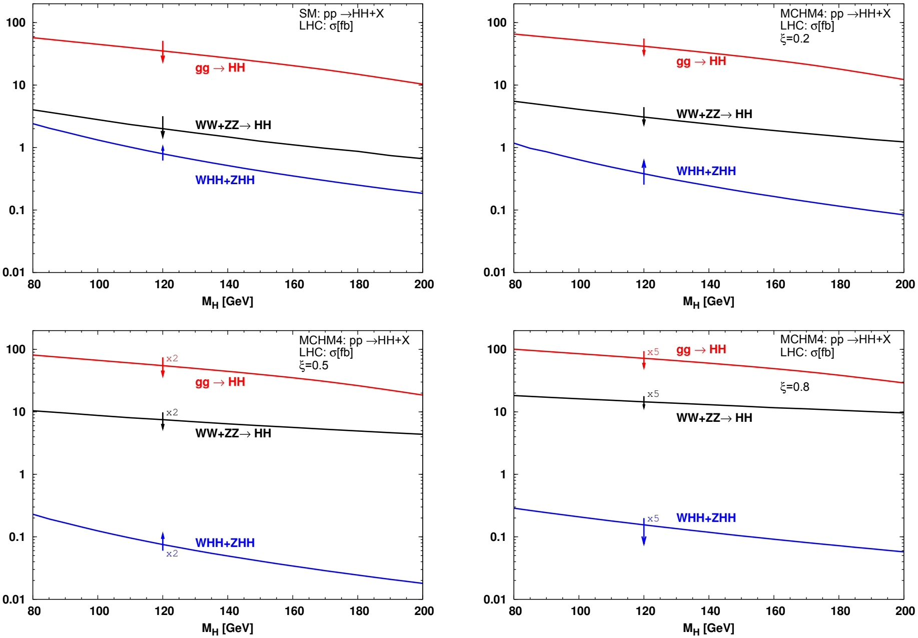


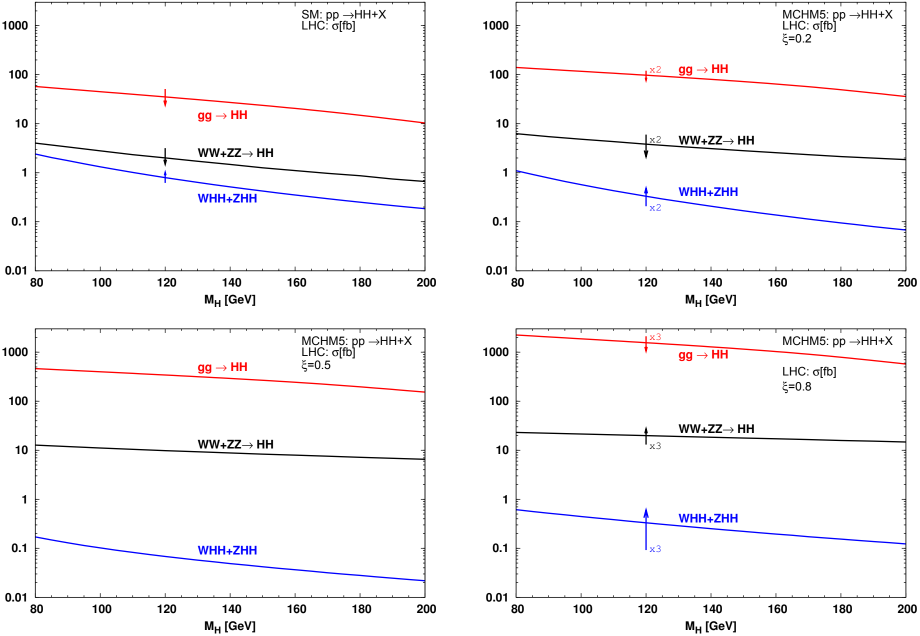
Vector boson fusion:



Double Higgs-strahlung:

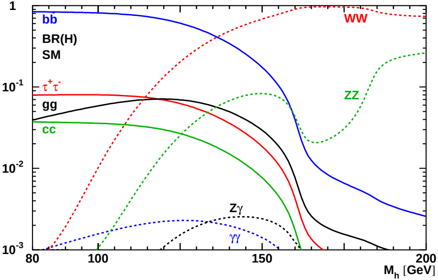



SM: pp →HH+X
LHC: σ[fb]
gg → HH
WW+ZZ→ HH
WHH+ZHH
MCHM4: pp →HH+X
LHC: σ[fb]
ξ=0.2
gg → HH
WW+ZZ→ HH
WHH+ZHH
MCHM4: pp →HH+X
LHC: σ[fb]
ξ=0.5
x2
gg → HH
x2
WW+ZZ→ HH
WHH+ZHH
x2
MCHM4: pp →HH+X
LHC: σ[fb]
x5
gg → HH
ξ=0.8
x5
WW+ZZ→ HH
x5
WHH+ZHH
100
10
1
0.1
0.01
80
100
120
140
160
180
200
M_H [GeV]


SM: pp →HH+X
LHC: σ[fb]
gg → HH
WW+ZZ→ HH
WHH+ZHH
MCHM5: pp →HH+X
LHC: σ[fb]
ξ=0.2
x2
MCHM5: pp →HH+X
LHC: σ[fb]
ξ=0.5
MCHM5: pp →HH+X
LHC: σ[fb]
ξ=0.8
x3
1000
100
10
1
0.1
0.01
80
100
120
140
160
180
200
M_H [GeV]

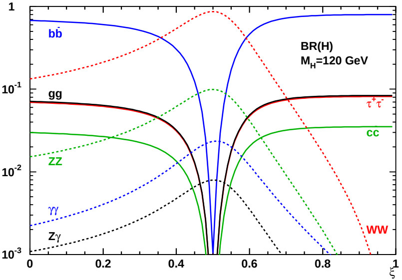
Branching ratios [Espinosa, Grojean, Mühlleitner]

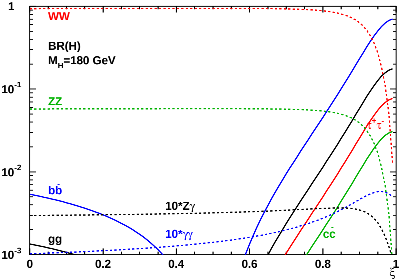
MCHM_4:



	MCHM_4	MCHM_5
HVV	$\sqrt{1-\xi}$	$\sqrt{1-\xi}$
$Hf\bar{f}$	$\sqrt{1-\xi}$	$\frac{1-2\xi}{\sqrt{1-\xi}}$

MCHM_5:





Sensitivities on $\lambda_{HHH} \neq 0$ in gluon fusion

In the SM the decay channel $b\bar{b}\gamma\gamma$ for low Higgs masses and $W^+W^-W^+W^-$ for Higgs masses above the gauge boson threshold have best prospects

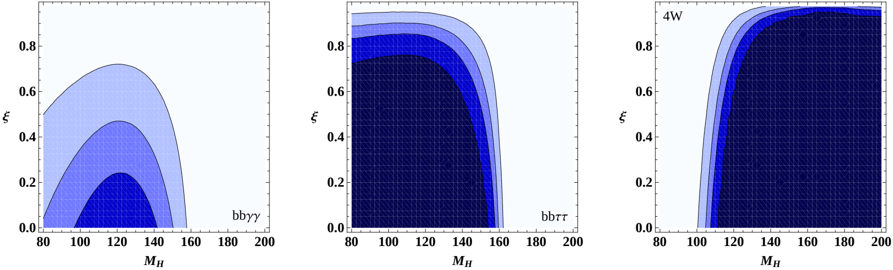
[Baur,Plehn,Rainwater]

Assumption: EWSB is realized in the framework of MCHM (ξ can be measured with accuracy of $\sim 20\%$ [Bock, Lafaye, Plehn, Rauch, Zerwas, Zerwas])

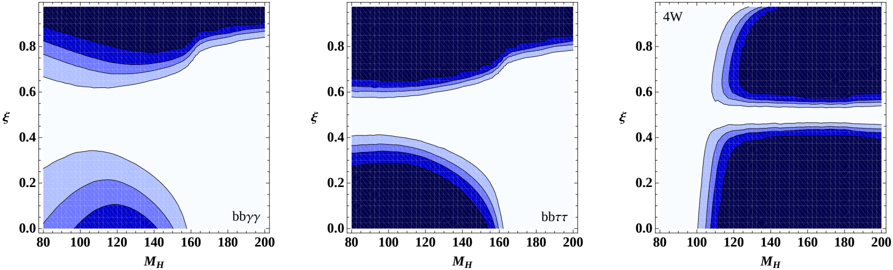
Criterion:

$$S + \beta\sqrt{S} < S_{\lambda_{HHH}=0} \quad \text{or} \quad S - \beta\sqrt{S} > S_{\lambda_{HHH}=0} \qquad (\beta = 1, 2, 3, 5)$$

MCHM$_4$ for $\int \mathcal{L} = 300\,\text{fb}^{-1}$:



MCHM$_5$ for $\int \mathcal{L} = 300\,\mathrm{fb}^{-1}$:



Above gauge boson threshold: Prospects of measuring non-vanishing λ_{HHH} are good.

Below gauge boson threshold: $b\bar{b}\tau^+\tau^-$ might be possible.

Feasibility

- Sensitivity areas present ideal case $\rightarrow$ for a full analysis: Detector properties and background processes must be investigated
- But e.g. for the $4W$ final state with $W^+W^-W^+W^- \rightarrow (jjl^\pm\nu)(jjl^\pm\nu)$:
  In SM: For 150 GeV $\leq M_H \leq$ 200 GeV a vanishing triple Higgs coupling can be excluded at 95% C.L. for 300 fb^{-1}
- Dominating background ($\mathcal{O}(\alpha_s \alpha^3)$)
  $\rightarrow$ no Higgs in intermediate state
  $\rightarrow$ leading background processes are the same as in SM
- But: Signal cross section is larger as in the SM!

Conclusion

- Composite Higgs models are an interesting alternative to the SM
- Cross sections for Higgs pair production via gluon fusion and vector boson fusion are enhanced compared to the SM for all ξ
- Cross section for Higgs pair production via double Higgs-strahlung is reduced compared to the SM
- Branching ratios in MCHM_5 are changed compared to the SM and depend on ξ.
- Prospects of measuring non-vanishing λ_{HHH} for large Higgs masses in $W^+W^-W^+W^-$ decay channel good, for low Higgs masses $b\bar{b}\tau^+\tau^-$ might be possible
- Measuring λ_{HHH} with higher precision might be possible for higher luminosities.

Constraints from electroweak data

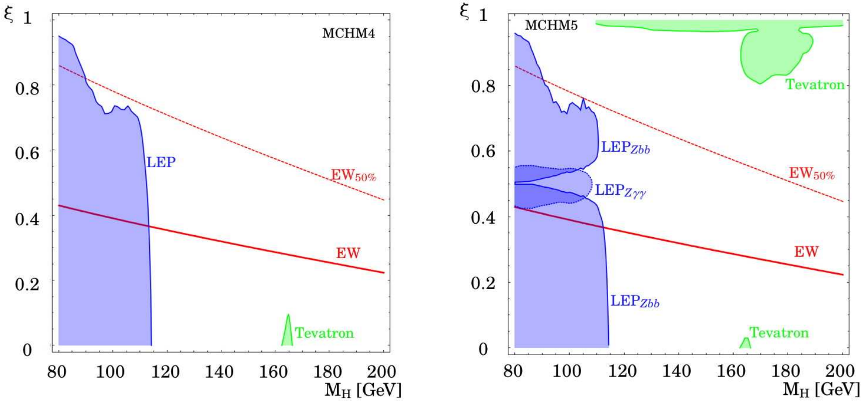
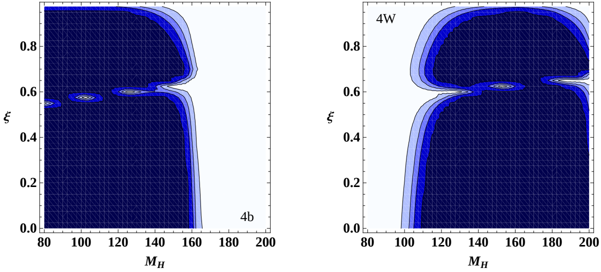


Figure: See [Espinosa,Grojean, Mühlleitner]

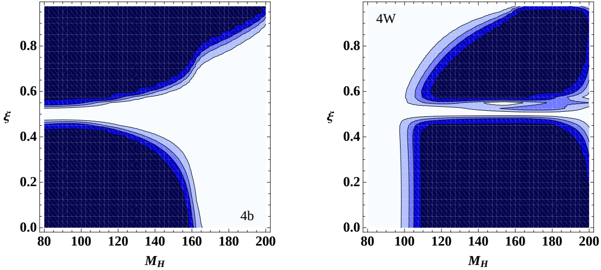
For constraints in model with partially composite top, see [Barbieri, Bellazzini, Rychkov, Varagnolo; Anastasiou, Furlan, Santiago; Gillioz; Lodone]

Double Higgs-strahlung at the ILC ($\sqrt{s} = 0.5$ TeV, $\int \mathcal{L} = 500\,\mathbf{fb}^{-1}$)

MCHM$_4$:

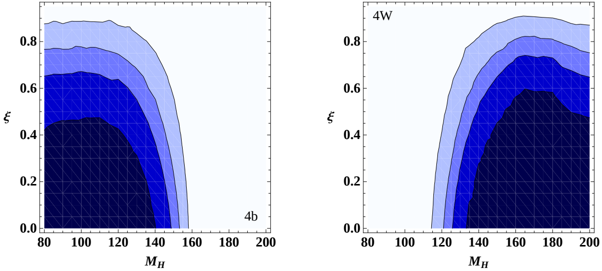


MCHM$_5$:

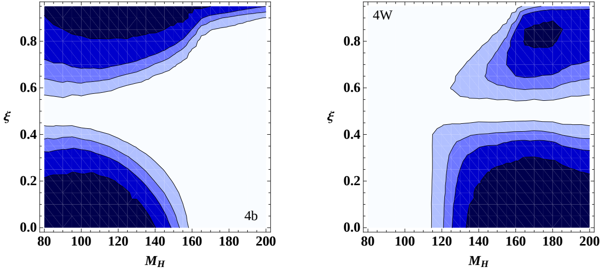


W boson fusion at the ILC ($\sqrt{s} = 0.5$ TeV, $\int \mathcal{L} = 500\,\mathbf{fb}^{-1}$)

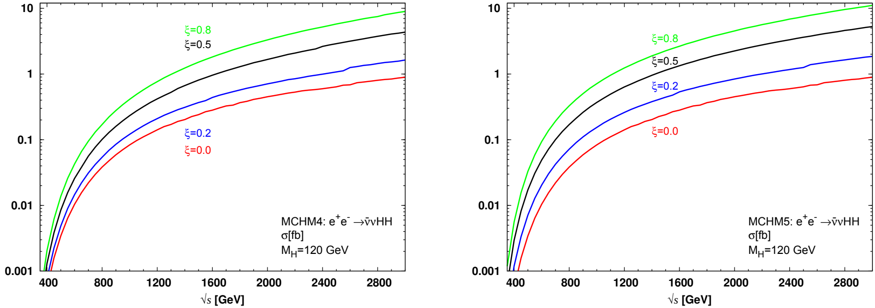
MCHM$_4$:



MCHM$_5$:



Energy behaviour of W fusion at the ILC



$$\mathcal{M} \overset{s\to\infty}{\longrightarrow} -\sqrt{2}G_F s\,\xi$$

SILH [Giudice, Grojean, Pomarol, Rattazzi]

New kinetic terms $\rightarrow$ field redefinition

$$\Phi = \frac{v}{\sqrt{2}} \begin{pmatrix} 0 \\ \frac{1}{v}\left(H - \frac{c_H \xi}{2}\left(H + \frac{H^2}{v} + \frac{H^3}{3v^2}\right)\right) + 1 \end{pmatrix}$$

$\rightarrow$ redefinition of Higgs and fermion masses and vacuum expectation value

$$\frac{v^2}{2} = \frac{\mu^2}{2\lambda}\left(1 - \frac{3}{4} c_6\, \xi\right) = \frac{(246\,\text{GeV})^2}{2} \;.$$

$$m_f = \frac{1}{\sqrt{2}} y_f v \left(1 + \frac{\xi}{2} c_y\right) \;,$$

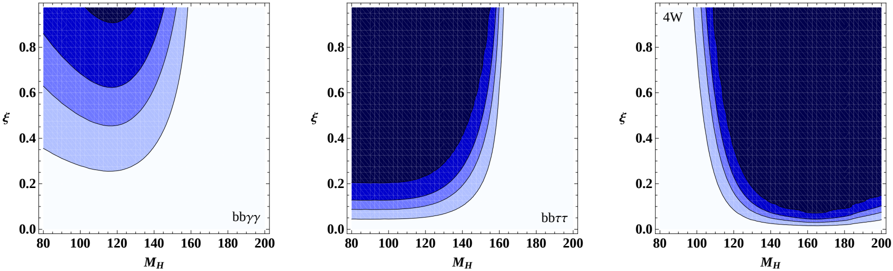
$$\frac{M_H^2}{2} = \frac{\left(\sqrt{2\lambda} v\right)^2}{2}\left(1 - c_H\, \xi + \frac{3}{4} c_6\, \xi\right) \;.$$

$$g_{HVV} = g_{HVV}^{SM}\left(1 - \frac{c_H}{2}\xi\right) \qquad \text{with } g_{HWW}^{SM} = g\,M_W \text{ and } g_{HZZ}^{SM} = \frac{g\,M_Z}{2}$$

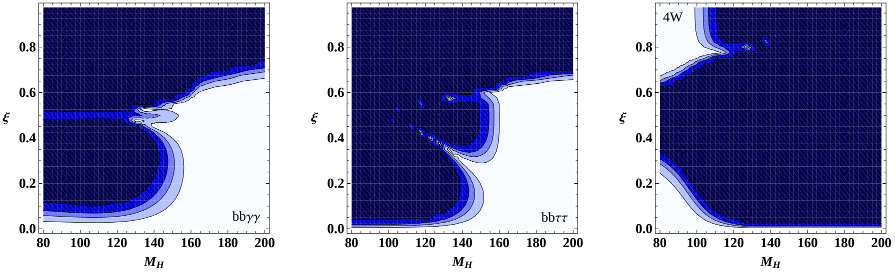
$$g_{Hff} = g_{Hff}^{SM}\left(1 - \xi\frac{c_H}{2} + \xi c_y\right) \qquad \text{with } g_{Hff}^{SM} = \frac{g\,m_f}{2\,M_W}$$

$$g_{HHH} = g_{HHH}^{SM}\left(1 + \xi c_6 - \xi\frac{3\,c_H}{2}\right) \qquad \text{with } g_{HHH}^{SM} = \frac{3\,g\,M_H^2}{2\,M_W}$$

Can the MCHM$_4$ be distinguished from the SM for $\int \mathcal{L} = 300\,\mathbf{fb}^{-1}$?

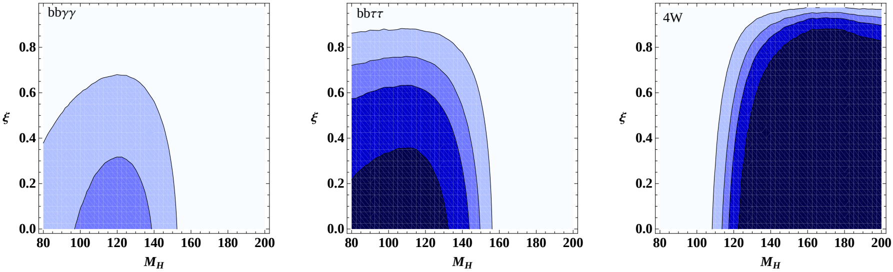


Can the MCHM$_5$ be distinguished from the SM for $\int \mathcal{L} = 300\,\textbf{fb}^{-1}$?



Can λ_{HHH} be measured with 30% precision in $MCHM_4$?

$\int \mathcal{L} = 3000\,\text{fb}^{-1}$ for $b\bar{b}\gamma\gamma$, $\int \mathcal{L} = 600\,\text{fb}^{-1}$ for $b\bar{b}\tau^+\tau^-$ and $4W$



Can λ_{HHH} be measured with 30% precision in MCHM$_5$?

$\int \mathcal{L} = 3000\,\text{fb}^{-1}$ for $b\bar{b}\gamma\gamma$, $\int \mathcal{L} = 600\,\text{fb}^{-1}$ for $b\bar{b}\tau^+\tau^-$ and $4W$