

# Heterotic Standard Models from Calabi-Yau Three-folds



Andre Lukas

University of Oxford

Workshop on Fields and Strings, Corfu, September 2011

arXiv:1106.4804 in collaboration with Lara Anderson, James Gray and Eran Palti

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# Overview

- Introduction: Heterotic Calabi-Yau compactifications
- Heterotic line bundle models
- A standard model data base
- A standard model example
- Phenomenological issues
- Conclusion and outlook

# Introduction

Data to define a heterotic N=1, d=4 Calabi-Yau compactification:

- A Calabi-Yau 3-fold  $X$
- A stable holomorphic vector bundle  $V$  on  $X$   
with structure group  $G \subset SU(n) \subset E_8$
- Anomaly condition:  $c_2(V) \leq c_2(TX)$   
(hidden sector with 5-branes and/or bundle)

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- freely acting symmetry  $\Gamma$  on  $X$ , so  $\hat{X} = X/\Gamma$  is smooth and non simply-connected
- bundle  $V$  needs to be equivariant so it descends to a bundle  $\hat{V}$  on  $\hat{X}$
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N=1, D=4  
standard-like  
model  
(hopefully)

Frequently: structure group  $G = SU(5) \subset E_8$  with commutant =  
low energy gauge group  $SU(5)$  is used.

$$\mathbf{248}_{E_8} \rightarrow [(\mathbf{1}, \mathbf{24}) \oplus (\mathbf{5}, \bar{\mathbf{10}}) \oplus (\bar{\mathbf{5}}, \mathbf{10}) \oplus (\mathbf{10}, \mathbf{5}) \oplus (\bar{\mathbf{10}}, \bar{\mathbf{5}}) \oplus (\mathbf{24}, \mathbf{1})]_{SU(5) \times SU(5)}$$



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Observation: bundle structure groups can (and often do) split at special loci in moduli space and the low-energy gauge group enhances. For example:

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Here, we study bundles with a “maximal” splitting:

$$SU(5) \rightarrow S(U(1)^5)$$

$$SU(5) \rightarrow SU(5) \times U(1)^4$$

## Key features:

- Gauge fields are Abelian and bundle is a sum of line bundles  $\rightarrow$  much easier to handle technically
- Abelian bundle still carries many of the properties of the generic non-Abelian bundle
- We have to remember that Abelian bundle usually resides in a larger, non-Abelian bundle moduli space.

## Heterotic line bundle models

CY manifold  $X$  with  $h^{1,1}(X)$  Kahler parameters  $t^i$ ,  
Kahler form  $J = t^i J_i$  and freely acting symmetry  $\Gamma$ .

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Note: A model is specified by the integer matrix  $k_a^i$ .



Recall: 4d gauge group is  $SU(5) \times S(U(1)^5) \cong SU(5) \times U(1)^4$

Label  $S(U(1)^5)$  representations by integer vectors  $\mathbf{q} = (q_1, \dots, q_5)$  with identification  $\mathbf{q} \sim \tilde{\mathbf{q}}$  iff  $\mathbf{q} - \tilde{\mathbf{q}} \in \mathbb{Z}(1, 1, 1, 1, 1)$ .

4d matter:  $10_{\mathbf{e}_a}$ ,  $\bar{10}_{-\mathbf{e}_a}$ ,  $\bar{5}_{\mathbf{e}_a+\mathbf{e}_b}$ ,  $5_{-\mathbf{e}_a-\mathbf{e}_b}$ ,  $1_{\mathbf{e}_a-\mathbf{e}_b}$ ,  $1_{-\mathbf{e}_a+\mathbf{e}_b}$  for  $a < b$ .

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families and  
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Number of each multiplet type obtained from  $H^1(X, L)$ .

bundle supersymmetry:  $D_a = k_a^i t_i - \sum_{b>a} (|C_{ab}^+|^2 - |C_{ab}^-|^2) \stackrel{!}{=} 0$

At Abelian locus  $\langle C_{ab}^\pm \rangle = 0$ , so slopes  $\mu(L_a) = k_a^i t_i \stackrel{!}{=} 0$ .

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Bundle moduli and moving away from Abelian locus:

$$0 \rightarrow L_a \rightarrow U \rightarrow L_b \rightarrow 0, \quad \text{Ext}^1(L_b, L_a) \cong H^1(X, L_a \otimes L_b^*) \ni C_{ab}^+$$

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$$\langle C_{ab}^\pm \rangle = 0 \implies U = L_a \oplus L_b$$

$$\langle C_{ab}^\pm \rangle \neq 0 \implies U \text{ is } U(2) \text{ bundle}$$

$$L_a \oplus L_b \text{ is replaced by } U$$



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$$(\text{number of massless } U(1)) = 4 - \text{rank}(k_a^i)$$

For all four  $U(1)$ s to be massive we need  $h^{1,1}(X) \geq 5$ .

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Models with massless U(1)s are still included. Such U(1)s can be broken spontaneously when moving away from the Abelian locus by switching on VEVs  $\langle C_{ab}^\pm \rangle$ .

## A standard model data base

### Arena: complete intersection CY manifolds (CICYs)

CICYs defined as common zero locus  $X = \{p_i = 0\} \subset \mathcal{A}$  of homogeneous polynomials  $p_i$  in ambient space  $\mathcal{A} = \bigotimes_{r=1}^m \mathbb{P}^{n_r}$ .

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for example: quintic  $X \sim [\mathbb{P}^4|5]$  or bi-cubic  $X \sim \left[ \begin{array}{c|c} \mathbb{P}^2 & 3 \\ \hline \mathbb{P}^2 & 3 \end{array} \right]$

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## Complete classification of about 8000 spaces

(Hubsch, Green, Lutken, Candelas 1987)

## Classification of freely-acting discrete symmetries

(Braun, 2010)

## Line bundle cohomology can be computed.

(Anderson, He, Lukas, 2008)

## The search for heterotic line bundle standard models:

- For a given CICY  $X$  with symmetry  $\Gamma$  generate many integer matrices  $(k_a^i)$  representing bundles  $V = \bigoplus_{a=1}^5 L_a$ ,  $L_a = \mathcal{O}_X(\mathbf{k}_a)$ .
- Check that  $\mu(L_a) = 0$  (bundle supersymmetric) and that  $V$  is equivariant under  $\Gamma$  ( $V$  descends to  $X/\Gamma$ ).
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- Require  $\text{ind}(L_a \otimes L_b) \leq 0$  (no chiral  $5$ ) and  $h^1(X, L_a^* \otimes L_b^*) < |\Gamma|$  (Higgs triplets can be projected out).

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- "heterotic standard model": SM gauge group (plus  $U(1)$ 's, massive or massless), three MSSM families, one or more Higgs doublets, bundle moduli (SM singlets), no exotics.

Have scanned CICYs with symmetries and  $h^{1,1}(X) \leq 5$  (60 spaces).

$(-9 + h^{1,1} \leq k_a^i \leq 9 - h^{1,1})$  (Anderson, Constantin, Gray, Lukas, Palti, in prep.)

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total number	1012
no massless U(1)	283
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Note: This is counting “upstairs” models. One upstairs model is often valid for more than one symmetry and one choice of Wilson line.

→ Total number of downstairs SMs is  $O(10000)$ .

CICY 6777: $\left(\begin{array}{c cccc} \mathbb{P}^1 & 1 & 1 & 0 & 0 \\ \mathbb{P}^1 & 0 & 0 & 0 & 2 \\ \mathbb{P}^1 & 0 & 0 & 2 & 0 \\ \mathbb{P}^1 & 2 & 0 & 0 & 0 \\ \mathbb{P}^3 & 1 & 1 & 1 & 1 \end{array}\right)_{-64}^{5,37}$	$\mathbb{Z}_2$
$(1,1,1,0,-1)(1,0,-2,-1,1)(1,-2,0,-1,1)(-1,1,1,1,-1)(-2,0,0,1,0)$	$(1,1,0,1,-1)(1,0,-1,0,0)(0,0,1,-1,0)(0,-2,-1,0,1)(-2,1,1,0,0)$
$(1,0,1,1,-1)(1,-1,0,0,0)(0,1,0,-1,0)(0,-1,-2,0,1)(-2,1,1,0,0)$	$(1,0,0,-1,0)(1,-1,-2,0,1)(0,1,1,1,-1)(0,-1,1,0,0)(-2,1,0,0,0)$
$(1,0,0,-1,0)(1,-2,-1,0,1)(0,1,1,1,-1)(0,1,-1,0,0)(-2,0,1,0,0)$	$(1,1,1,0,-1)(0,1,-1,0,0)(0,0,1,-2,0)(0,-2,-1,1,1)(-1,0,0,1,0)$
$(1,1,0,1,-1)(0,1,1,-2,0)(0,0,-1,1,0)(0,-2,-1,0,1)(-1,0,1,0,0)$	$(1,1,-1,-1,0)(0,1,1,1,-1)(0,0,-1,-2,1)(0,-2,1,1,0)(-1,0,0,1,0)$
$(1,1,1,0,-1)(0,1,0,-2,0)(0,-1,1,0,0)(0,-1,-2,1,1)(-1,0,0,1,0)$	$(1,0,1,1,-1)(0,1,1,-2,0)(0,-1,0,1,0)(0,-1,-2,0,1)(-1,1,0,0,0)$
$(1,-1,1,-1,0)(0,1,1,1,-1)(0,1,-2,1,0)(0,-1,0,-2,1)(-1,0,0,1,0)$	
CICY 6890: $\left(\begin{array}{c ccccc} \mathbb{P}^1 & 1 & 1 & 0 & 0 & 0 \\ \mathbb{P}^1 & 0 & 0 & 1 & 1 & 0 \\ \mathbb{P}^1 & 0 & 0 & 0 & 0 & 2 \\ \mathbb{P}^1 & 0 & 0 & 2 & 0 & 0 \\ \mathbb{P}^4 & 1 & 1 & 1 & 1 & 1 \end{array}\right)_{-64}^{5,37}$	$\mathbb{Z}_2$
$(1,1,1,0,-1)(1,0,-1,0,0)(1,-2,0,1,0)(-1,1,1,-1,0)(-2,0,-1,0,1)$	$(1,1,1,0,-1)(1,0,-1,0,0)(0,0,1,-2,0)(0,-1,0,1,0)(-2,0,-1,1,1)$
$(1,1,0,1,-1)(1,0,1,-2,0)(0,0,-1,1,0)(0,-1,1,0,0)(-2,0,-1,0,1)$	$(1,1,-1,-1,0)(1,0,1,1,-1)(0,0,-1,-2,1)(0,-1,0,1,0)(-2,0,1,1,0)$
$(1,1,1,0,-1)(1,-1,1,1,-1)(0,1,-2,-1,1)(0,-2,0,1,0)(-2,1,0,-1,1)$	$(1,1,0,1,-1)(1,-2,1,0,0)(0,1,-1,0,0)(0,0,1,-1,0)(-2,0,-1,0,1)$
$(1,0,1,1,-1)(1,0,-1,0,0)(0,1,0,-1,0)(0,-2,1,0,0)(-2,1,-1,0,1)$	$(1,1,1,0,-1)(1,0,0,-2,0)(0,-1,0,1,0)(-1,0,1,0,0)(-1,0,-2,1,1)$
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$(1,0,1,1,-1)(1,0,-2,1,0)(0,-1,0,1,0)(-1,1,1,-1,0)(-1,0,0,-2,1)$	$(1,0,1,-2,0)(1,-1,0,0,0)(0,1,1,1,-1)(-1,0,0,1,0)(-1,0,-2,0,1)$
$(1,0,1,1,-1)(1,-2,0,1,0)(0,1,0,-1,0)(-1,1,1,-1,0)(-1,0,-2,0,1)$	$(1,0,1,1,-1)(1,-2,0,0,0)(0,1,0,-1,0)(-1,1,-2,0,1)(-1,0,1,0,0)$
$(1,0,0,-1,0)(1,-2,1,0,0)(0,1,1,1,-1)(-1,1,0,0,0)(-1,0,-2,0,1)$	$(1,1,1,0,-1)(0,1,0,-2,0)(0,0,-1,1,0)(0,-2,-1,1,1)(-1,0,1,0,0)$
CICY 7447: $\left(\begin{array}{c cc} \mathbb{P}^1 & 1 & 1 \\ \mathbb{P}^1 & 1 & 1 \\ \mathbb{P}^1 & 1 & 1 \\ \mathbb{P}^1 & 1 & 1 \\ \mathbb{P}^1 & 1 & 1 \end{array}\right)_{-80}^{5,45}$	$\mathbb{Z}_2 \times \mathbb{Z}_2$
$(0,1,0,-2,1)(0,1,-2,1,0)(0,0,1,1,-2)(0,-1,1,0,0)(0,-1,0,0,1)$	$(1,-2,0,0,1)(0,1,-2,0,1)(0,0,1,1,-2)(0,0,1,-1,0)(-1,1,0,0,0)$
$(1,-2,0,0,1)(0,1,0,1,-2)(0,0,1,-2,1)(0,0,-1,0,1)(-1,1,0,1,-1)$	$(1,-2,-1,1,1)(0,1,1,-2,0)(0,1,-1,0,0)(0,0,1,1,-2)(-1,0,0,0,1)$

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CICY 6777: $\left( \begin{array}{c cccc} \mathbb{P}^1 & 1 & 1 & 0 & 0 \\ \mathbb{P}^1 & 0 & 0 & 0 & 2 \\ \mathbb{P}^1 & 0 & 0 & 2 & 0 \\ \mathbb{P}^1 & 2 & 0 & 0 & 0 \\ \mathbb{P}^3 & 1 & 1 & 1 & 1 \end{array} \right)_{-64}^{5,37}$	$\mathbb{Z}_2$
$(1,1,1,0,-1)(1,0,-2,-1,1)(1,-2,0,-1,1)(-1,1,1,1,-1)(-2,0,0,1,0)$	$(1,1,0,1,-1)(1,0,-1,0,0)(0,0,1,-1,0)(0,-2,-1,0,1)(-2,1,1,0,0)$
$(1,0,1,1,-1)(1,-1,0,0,0)(0,1,0,-1,0)(0,-1,-2,0,1)(-2,1,1,0,0)$	$(1,0,0,-1,0)(1,-1,-2,0,1)(0,1,1,1,-1)(0,-1,1,0,0)(-2,1,0,0,0)$
$(1,0,0,-1,0)(1,-2,-1,0,1)(0,1,1,1,-1)(0,1,-1,0,0)(-2,0,1,0,0)$	$(1,1,1,0,-1)(0,1,-1,0,0)(0,0,1,-2,0)(0,-2,-1,1,1)(-1,0,0,1,0)$
$(1,1,0,1,-1)(0,1,1,-2,0)(0,0,-1,1,0)(0,-2,-1,0,1)(-1,0,1,0,0)$	$(1,1,-1,-1,0)(0,1,1,1,-1)(0,0,-1,-2,1)(0,-2,1,1,0)(-1,0,0,1,0)$
$(1,1,1,0,-1)(0,1,0,-2,0)(0,-1,1,0,0)(0,-1,-2,1,1)(-1,0,0,1,0)$	$(1,0,1,1,-1)(0,1,1,-2,0)(0,-1,0,1,0)(0,-1,-2,0,1)(-1,1,0,0,0)$
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CICY 6890: $\left( \begin{array}{c ccccc} \mathbb{P}^1 & 1 & 1 & 0 & 0 & 0 \\ \mathbb{P}^1 & 0 & 0 & 1 & 1 & 0 \\ \mathbb{P}^1 & 0 & 0 & 0 & 0 & 2 \\ \mathbb{P}^1 & 0 & 0 & 2 & 0 & 0 \\ \mathbb{P}^4 & 1 & 1 & 1 & 1 & 1 \end{array} \right)_{-64}^{5,37}$	$\mathbb{Z}_2$
$(1,1,1,0,-1)(1,0,-1,0,0)(1,-2,0,1,0)(-1,1,1,-1,0)(-2,0,-1,0,1)$	$(1,1,1,0,-1)(1,0,-1,0,0)(0,0,1,-2,0)(0,-1,0,1,0)(-2,0,-1,1,1)$
$(1,1,0,1,-1)(1,0,1,-2,0)(0,0,-1,1,0)(0,-1,1,0,0)(-2,0,-1,0,1)$	$(1,1,-1,-1,0)(1,0,1,1,-1)(0,0,-1,-2,1)(0,-1,0,1,0)(-2,0,1,1,0)$
$(1,1,1,0,-1)(1,-1,1,1,-1)(0,1,-2,-1,1)(0,-2,0,1,0)(-2,1,0,-1,1)$	$(1,1,0,1,-1)(1,-2,1,0,0)(0,1,-1,0,0)(0,0,1,-1,0)(-2,0,-1,0,1)$
$(1,0,1,1,-1)(1,0,-1,0,0)(0,1,0,-1,0)(0,-2,1,0,0)(-2,1,-1,0,1)$	$(1,1,1,0,-1)(1,0,0,-2,0)(0,-1,0,1,0)(-1,0,1,0,0)(-1,0,-2,1,1)$
$(1,1,1,0,-1)(1,0,-2,1,0)(0,-2,-1,0,1)(-1,1,1,-1,0)(-1,0,1,0,0)$	$(1,1,1,0,-1)(1,-2,0,1,0)(0,1,0,-1,0)(-1,0,1,0,0)(-1,0,-2,0,1)$
$(1,0,1,1,-1)(1,0,-2,1,0)(0,-1,0,1,0)(-1,1,1,-1,0)(-1,0,0,-2,1)$	$(1,0,1,-2,0)(1,-1,0,0,0)(0,1,1,1,-1)(-1,0,0,1,0)(-1,0,-2,0,1)$
$(1,0,1,1,-1)(1,-2,0,1,0)(0,1,0,-1,0)(-1,1,1,-1,0)(-1,0,-2,0,1)$	$(1,0,1,1,-1)(1,-2,0,0,0)(0,1,0,-1,0)(-1,1,-2,0,1)(-1,0,1,0,0)$
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## A standard model example

$$L_1 = \mathcal{O}_X(1, 0, 0, -1, 0), \quad L_2 = \mathcal{O}_X(1, -1, -2, 0, 1), \quad L_3 = \mathcal{O}_X(0, 1, 1, 1, -1) \\ L_4 = \mathcal{O}_X(0, -1, 1, 0, 0), \quad L_5 = \mathcal{O}_X(-2, 1, 0, 0, 0)$$



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$S(U(1)^5)$  symmetry restricts operators in 4d theory.

**Note:**  $S(U(1)^5)$  non-invariant operators are perturbatively forbidden,  
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Forbidden for example.

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Forbidden for example.

## ● Dimension four proton decay away from Abelian locus

$p(C_I) \bar{5} \bar{5} 10$  can be checked explicitly, since all charges known

Again, forbidden for example.

# Phenomenological issues

$S(U(1)^5)$  symmetry restricts operators in 4d theory.

Note:  $S(U(1)^5)$  non-invariant operators are perturbatively forbidden, but allowed operators are not necessarily present.

## ● Dimension four proton decay at Abelian locus

$\bar{5}_{e_a+e_b} \bar{5}_{e_c+e_d} 10_{e_f}$  allowed if  $e_a + e_b + e_c + e_d + e_f = (1, 1, 1, 1, 1)$

Forbidden for example.

## ● Dimension four proton decay away from Abelian locus

$p(C_I) \bar{5} \bar{5} 10$  can be checked explicitly, since all charges known

Again, forbidden for example.

## ● Dimension five proton decay

$\bar{5}_{e_a+e_b} 10_{e_c} 10_{e_d} 10_{e_f}$  allowed if  $e_a + e_b + e_c + e_d + e_f = (1, 1, 1, 1, 1)$

Ok for example (also with singlet insertions).

## $\mu$ - term

$H_u H_d$  is singlet, but term absent at Abelian locus.

Away from Abelian locus:  $C_I H_u H_d$  forbidden

$C_I C_J H_u H_d$  may be allowed

For example:  $C_5 C_6 H_u H_d$  allowed, so we need  $\epsilon_5 \epsilon_6 \ll 1$ .

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For example:  $Y^u = \begin{pmatrix} \epsilon_5 & 1 & 1 \\ 1 & \epsilon_6 & \epsilon_6 \\ 1 & \epsilon_6 & \epsilon_6 \end{pmatrix}$ ,  $Y^d = \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix}$

(needs non-perturbative effects)

## Conclusion and outlook

- Heterotic line bundle models on CY manifolds are a useful and technically accessible arena for string model building.
  - We have found 1000+ (upstairs) heterotic standard models on CICYs with  $h^{1,1}(X) \leq 5$ .
  - The  $S(U(1)^5)$  symmetry restricts the 4d theory and facilitates phenomenological analysis beyond the computation of the spectrum.
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- Need a better understanding of full, non-Abelian moduli space.
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(Anderson, Gray, Lukas, Ovrut, 2010)

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*Thanks*