

AdS/CFT and the cosmological constant problem

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CORFU

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Introduction

- In QFT vacuum energy not observable
- In Gravity vacuum energy *does* backreact
- Quantum fluctuations $\Rightarrow \rho_{1-loop} \sim M_{cut}^4$
- If we take $M_{cut} \sim M_P$ then

$$\frac{\rho_{observed}}{\rho_{1-loop}} \sim 10^{-120}$$

Cosmological constant fine-tuning

- "Renormalization"

$$\rho_{obs} = \rho_{bare} + \rho_{1-loop}$$

- Two terms must be fine-tuned, one part in 10^{120}
- Will not discuss what fixes specific ρ_{obs} or cosmic coincidence problem etc.

c.c. fine-tuning and effective field theory

- c.c. problem indicates something wrong with effective field theory

- Naturalness \Rightarrow

$$S_{eff} \sim \int \left(c_1 M_{cut}^{4-\Delta_1} \mathcal{O}_1 + c_2 M_{cut}^{4-\Delta_2} \mathcal{O}_2 + \dots \right)$$

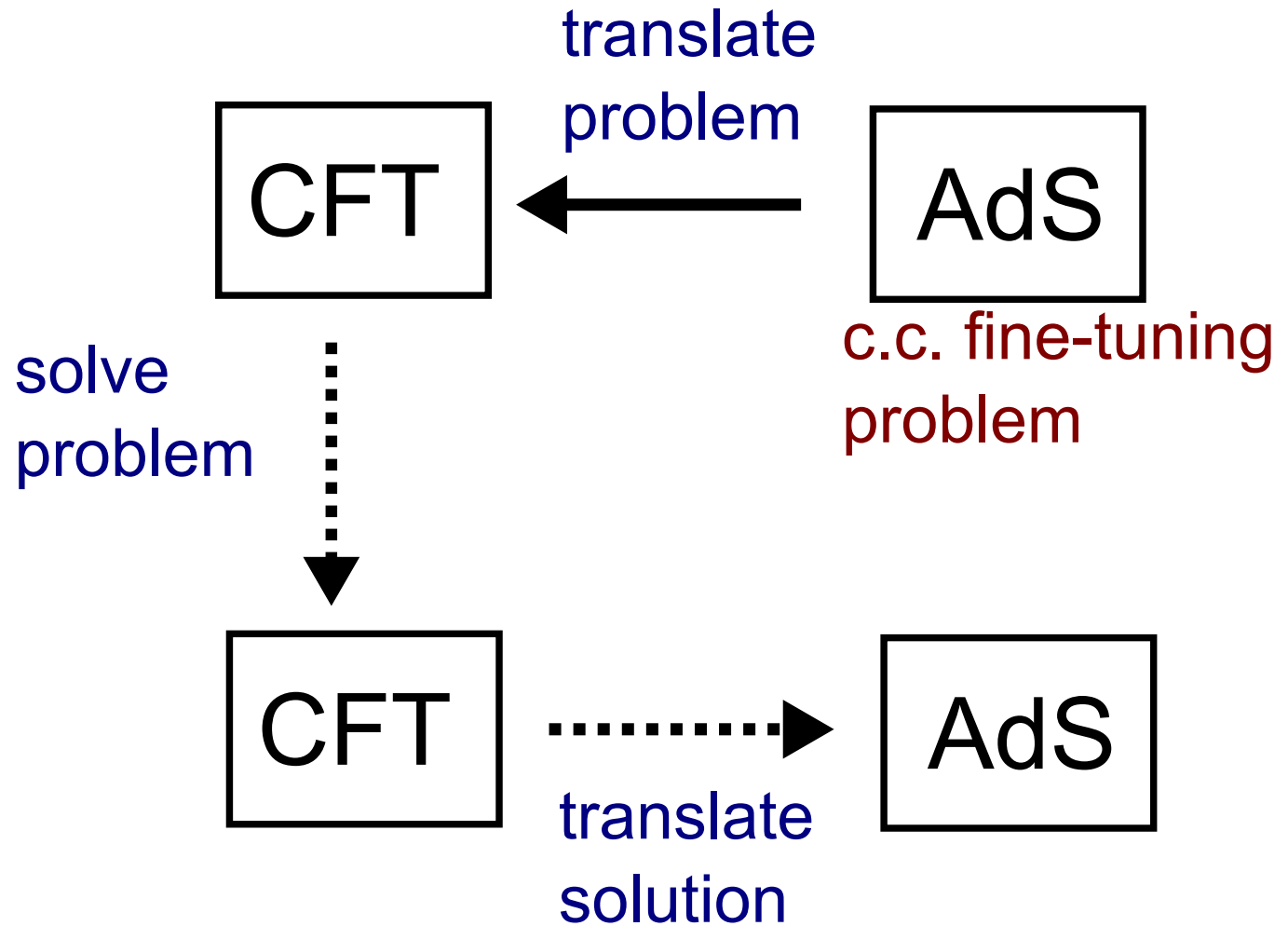
with $c_1, c_2, \dots \sim \mathcal{O}(1)$

- Vacuum energy term $\Delta = 0$ (most relevant operator) \Rightarrow should be multiplied by M_{cut}^4
- In a natural theory vacuum energy should be of order M_{cut}^4
- Is there a symmetry/dynamical mechanism to suppress value of c.c.?

Spacetime is emergent

- Why don't quantum fluctuations gravitate?
- Question independent of $\Lambda > 0$ or $\Lambda < 0$
- c.c. fine tuning also in AdS
- AdS/CFT: spacetime and gravity are EMERGENT from CFT
- Should be able to see (via the dual CFT) whether quantum fluctuations gravitate or not

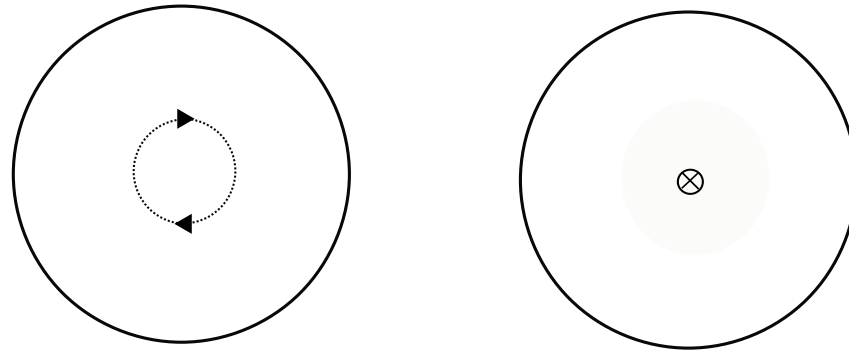
Strategy



MAIN QUESTION

How do we translate the bulk cosmological constant fine-tuning problem into a CFT question?

How do we proceed?



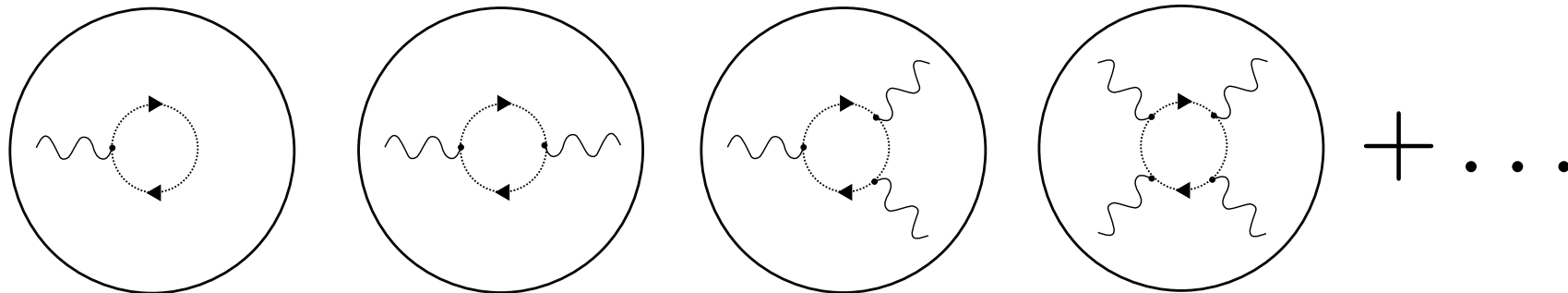
- Need to understand meaning of "1-loop diagram" and "counterterm" (\sim bare c.c.) in CFT language
- Difficult to map the bulk vacuum-to-vacuum amplitudes in CFT
- More familiar with computation of Witten diagrams (diagrams with external legs)

c.c. fine-tuning in correlation functions

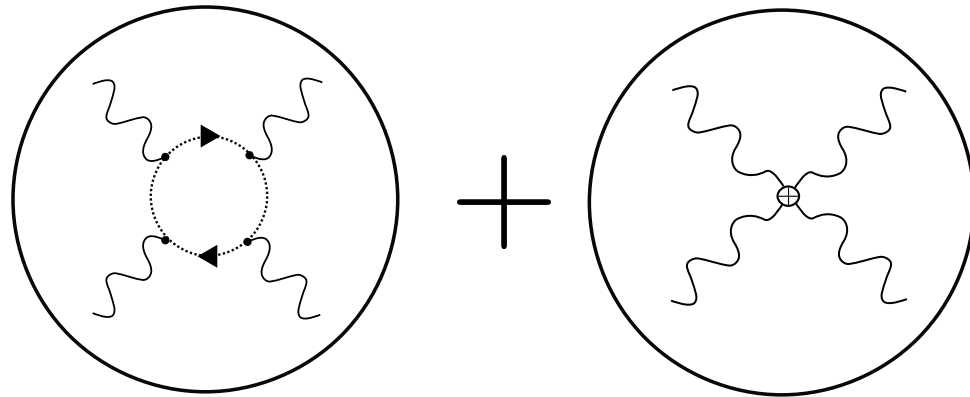
- c.c. counterterm has the form

$$\int d^4x \sqrt{g} \Lambda_{bare}$$

- Expanding \sqrt{g} generates graviton vertices with arbitrary number of external legs
- Necessary to (partly) cancel diagrams of the form



Example: graviton 4-point function



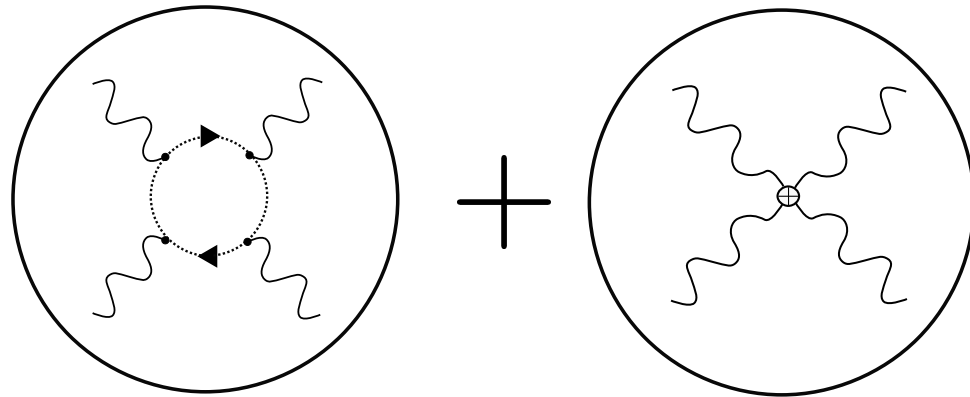
- If observed c.c. is small, then connected 4-point function should be small i.e. of the order

$$(M_P R)^{-2}$$

- If we take AdS as big as the universe, 4-point function should be of order

$$10^{-120}$$

Example: graviton 4-point function



- Estimate 1-loop diagram:

$$G_N^2 \int (d^4 p) p^4 \left(\frac{1}{p}\right)^4 \sim M_P^{-4} M_{cut}^4 \sim \mathcal{O}(1)!$$

- two diagrams above are both $\mathcal{O}(1)$ and exactly cancel up to something of order $\mathcal{O}(10^{-120})$

Summary

- c.c. fine-tuning can be understood in terms of bulk correlation functions (Witten diagrams)
- via AdS/CFT \Rightarrow c.c. fine-tuning should be visible in CFT correlation functions

Fine-tuning in CFT

- CFTs with holographic duals \Rightarrow large N expansion
- Graviton scattering in AdS \sim stress energy correlation functions in CFT
- Small (observerd) c.c. \Leftrightarrow

$$\langle \tilde{T}(x_1) \dots \tilde{T}(x_n) \rangle_{con} \sim N^{2-n}$$

Fine-tuning in the CFT

???

- IS THE LARGE N EXPANSION NATURAL OR "FINE-TUNED"??
- For example, 4-point function should be of order

$$\frac{1}{N^2}$$

is the $1/N$ suppression achieved in a **natural** way, or via cancellations between parametrically larger terms?

Splitting up of CFT correlators?



Fine-tuning in CFT

- Any CFT correlator \Rightarrow expanded in "conformal blocks"

$$\langle \tilde{T}(x_1)\tilde{T}(x_2)\tilde{T}(x_3)\tilde{T}(x_4) \rangle_{cont} = \sum_{\mathcal{A}} |C_{TT}^{\mathcal{A}}|^2 \mathbf{G}_{\mathcal{A}}(x_1, x_2, x_3, x_4)$$

Fine-tuning in CFT

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OUR CONJECTURE

Fine-tuning in CFT

- Any CFT correlator \Rightarrow expanded in "conformal blocks"

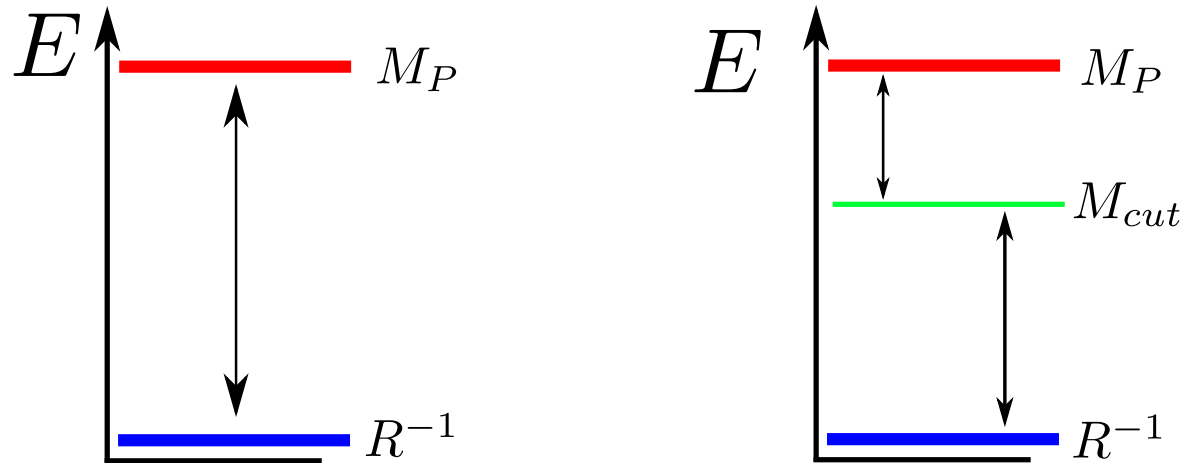
$$\langle \tilde{T}(x_1)\tilde{T}(x_2)\tilde{T}(x_3)\tilde{T}(x_4) \rangle_{cont} = \sum_A |C_{TT}^A|^2 \mathbf{G}_A(x_1, x_2, x_3, x_4)$$

OUR CONJECTURE

- Bulk c.c. fine tuning problem \Leftrightarrow "un-naturalness" of large N expansion in CFT
- Final $1/N$ suppression of correlators is achieved via delicate cancellations between (parameterically larger) conformal blocks

(Notice: $1/N$ expansion and conformal block expansion are not the same thing!)

Comments



- Hierarchy vs fine-tuning

- Fine-tuned if

$$M_{cut}^4 M_P^{-2} R^2 \gg 1$$

- Need high cutoff in bulk theory

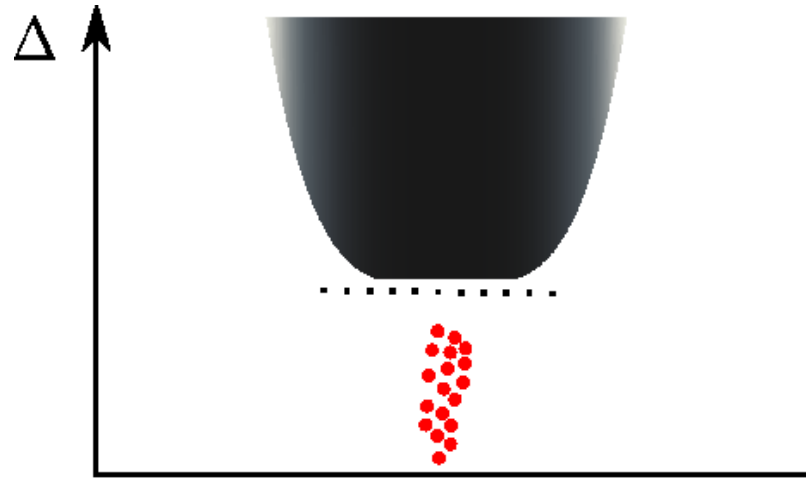
Large N gauge theories

- Large N gauge theories in 't Hooft limit \Rightarrow weakly coupled string theories (Hagedorn growth)
- String scale $M_s \sim f(\lambda)/R$ (does not scale with N), which implies

$$M_{cut} \sim M_s \ll M_P$$

- \Rightarrow Cannot pose a **sharp** c.c. fine-tuning problem
- Same for higher spin gravity (Vasiliev)

CFTs with holographic duals



1. **Large central charge c**
2. **Few low-lying operators (c -independent)**
3. **Factorization of correlators**
For semi-classical gravity
4. **Gap between spin 2 and higher**

Candidate CFTs with c.c. problem

Conditions

1. CFT must have a holographic dual (large c , few operators etc.)
2. CFT should not be supersymmetric
3. Conformal dimension Δ_{cut} where $\text{spin} > 2$ states appear must satisfy

$$\Delta_{cut} \gg c^{1/d-1}$$

- 1 and 2 but not 3: large N gauge theories, $O(N)/\text{WZW}/\text{coset}$ models, non-susy orbifolds of $\mathcal{N} = 4$
- 1 and 3 but not 2: ABJM, $k = 1, N \rightarrow \infty$
- 1 and 2 and 3?

Conformal block expansion

- Operator Product Expansion

$$\mathcal{O}_i(x)\mathcal{O}_j(y) \sim C_{ij}^k \mathcal{O}_k(y) + \dots$$

- Any n -point function can be computed by OPEs

$$\langle \mathcal{O}_i(x_1)\mathcal{O}_j(x_2)\mathcal{O}_m(x_3)\mathcal{O}_n(x_4) \rangle = \sum_k C_{12}^k C_{34}^k \mathbf{G}_k(x_1, x_2, x_3, x_4)$$

- Consistency condition: "conformal bootstrap"

$$\sum_k \begin{array}{c} 1 \\ \diagdown \\ \text{---} \\ \diagup \\ 2 \end{array} \text{---} \text{=} \sum_k \begin{array}{c} 1 \\ \diagdown \\ \text{---} \\ \diagup \\ 2 \end{array} \begin{array}{c} 4 \\ \diagdown \\ \text{---} \\ \diagup \\ 3 \end{array}$$

Witten diagrams and conformal blocks

- Interactions in AdS \Rightarrow Witten diagrams
- Witten diagram expansion \sim conformal block expansion

The diagram shows a circular Witten diagram on the left with four external legs labeled $\phi_1, \phi_2, \phi_3, \phi_4$ and an internal dashed line labeled ϕ_m . This is equated to a sum of conformal blocks on the right. The first term is a tree-level block with external legs $\mathcal{O}_1, \mathcal{O}_2, \mathcal{O}_3, \mathcal{O}_4$ and an internal line labeled \mathcal{O}_m . The second term is a sum over n of a block with external legs $\mathcal{O}_1, \mathcal{O}_2, \mathcal{O}_3, \mathcal{O}_4$ and an internal line labeled $:\mathcal{O}_1 \hat{\partial}^n \mathcal{O}_2:$. The third term is a sum over n of a block with external legs $\mathcal{O}_1, \mathcal{O}_2, \mathcal{O}_3, \mathcal{O}_4$ and an internal line labeled $:\mathcal{O}_3 \hat{\partial}^n \mathcal{O}_4:$.

- Witten diagram basis = easy solutions of conformal bootstrap

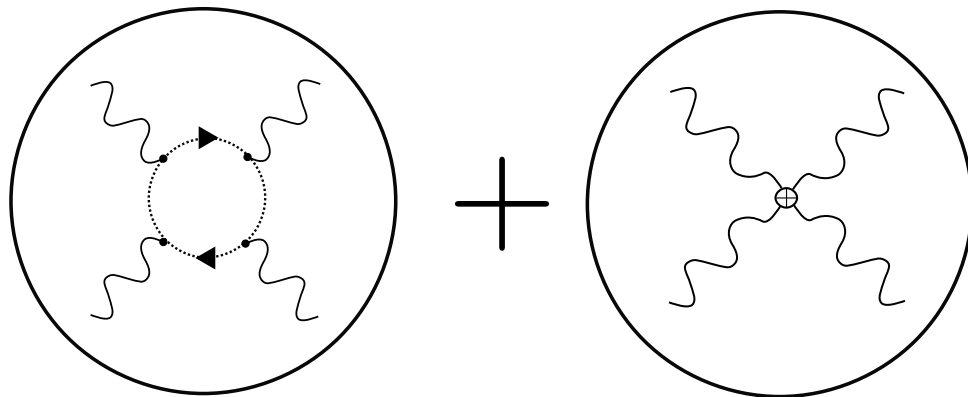
4-point function of $T_{\mu\nu}$

- It has a large N expansion

$$\langle \tilde{T}(x_1) \tilde{T}(x_2) \tilde{T}(x_3) \tilde{T}(x_4) \rangle =$$

$$= G_0(x_1, \dots, x_4) + \frac{1}{N^2} G_1(x_1, \dots, x_4) + \frac{1}{N^4} G_2(x_1, \dots, x_4) + \dots$$

- $G_2(x_1, \dots, x_4)$ gets contributions (at least) from



4-point function of $T_{\mu\nu}$

- The term $\frac{1}{N^4}G_2(x_1, \dots, x_4)$ is $\mathcal{O}(1/N^4)$. However if we split it into Witten diagrams \Rightarrow each of them is $\mathcal{O}(1)$
 - We have delicate cancellations between Witten diagrams
- \Rightarrow
- In CFT language: expand $\frac{1}{N^4}G_2(x_1, \dots, x_4)$ in conformal blocks. While $\frac{1}{N^4}G_2$ is order $\mathcal{O}(1/N^4)$, individual conformal blocks contributing to it are $\mathcal{O}(1)$ i.e. **parametrically larger**
 - **(Apparently) fine-tuned cancellations between conformal blocks**

1-loop and "counterterm" diagrams

- 1-loop Witten diagram corresponds to conformal blocks of

$$: \Psi (\partial^2)^n \partial_1 \dots \partial_l \Psi :$$

- Sum over n, l is \sim momentum integral in bulk \Rightarrow gives very large contribution
- Counterterm Witten diagram conformal block of

$$: T \partial \dots \partial T :$$

also comes with very large coefficient in order to cancel the divergence of the previous sum

- **COSMOLOGICAL CONSTANT FINE-TUNING IN ADS**

⇒

- **UN-NATURALNESS OF LARGE N EXPANSION IN CFT**
(from conformal block point of view)

Can we study this in more detail?

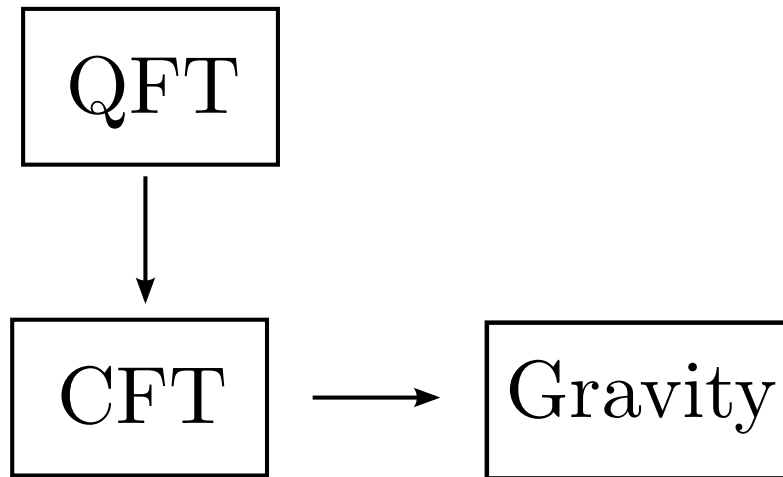
Is there a specific CFT where this fine-tuning takes place?

- No example known yet
- No such CFT exists \Leftrightarrow AdS gravity without (low energy) supersymmetry inconsistent



**IS THERE A RESOLUTION OF THE FINE-TUNING
FROM THE CFT?**

-
- Assume large N expansion looks fine tuned from conformal block point of view
 - Could there be another way to calculate correlators, where $1/N$ suppression looks natural?
 - For example in large N gauge theories, conformal block expansion: expansion in term of exchanged gauge singlets (ie. **glueballs, mesons, etc.** in intermediate channels
 - **BUT:** fundamentally correlators are computed by "double-line" diagrams in terms of **quarks and gluons**. In this description the $1/N$ expansion is manifest



QFT: description in terms of fundamental fields

CFT: description in terms of "gauge singlets" / CFT data
(Δ_i, C_{ij}^k)

While mathematically equivalent, large N expansion may look natural in first but fine-tuned in second.

-
- If true \Rightarrow Impossible to understand c.c. fine-tuning in terms of low energy fields (gravitons etc.)
 - We need to understand what they are "made out of" (in the dual gauge theory)
 - Underlying stabilizing mechanism \Rightarrow Cancellations between terms of different order in perturbation theory (unlike supersymmetry)

CONCLUSIONS

- To the extent that there is a consistent theory of AdS quantum gravity with the c.c. problem, we argued that in dual CFT the problem is manifest as an (apparent) fine-tuning of the large N expansion
- We argued that the CFT may admit an alternative description where this fine-tuning becomes naturally resolved
- WE NEED AN EXAMPLE! (does not have to be a full-fledged CFT)

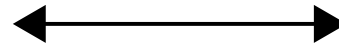


THANK YOU

Bulk

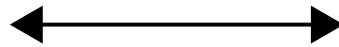
Boundary

Radial
Evolution



RG-flow

RG-flow



?