

# The Phase Structure of the Polyakov–Quark-Meson Model beyond Mean Field

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In Collaboration with B.-J. Schaefer and J.M. Pawłowski

arXiv: 1008.0081 [hep-ph]



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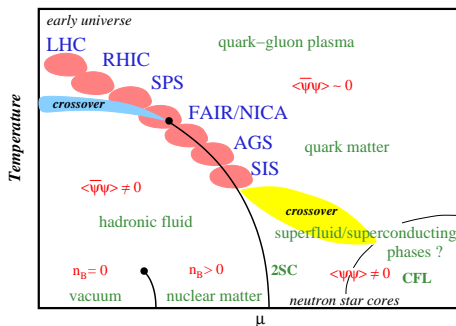
Corfu, Greece



Recipient of a DOC-fFORTE-fellowship of the Austrian Academy of Sciences  
at the Institute of Physics.



# QCD Phase Structure



(approx.) order parameters

$$\langle \bar{q}q \rangle \begin{cases} = 0 & \text{symmetric} \\ \neq 0 & \text{broken} \end{cases}$$

$$\Phi \begin{cases} = 0 & \text{confined} \\ \neq 0 & \text{deconfined} \end{cases}$$

## Chiral Symmetry

- $m_q \rightarrow 0$
- **chiral condensate**  $\langle \bar{q}q \rangle$

## $Z_{N_c}$ Center Symmetry

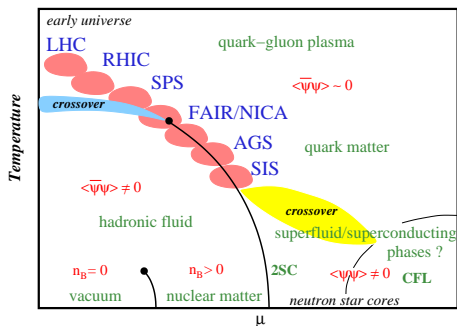
- $m_q \rightarrow \infty$
- **Polyakov loop**  $\Phi = \langle \ell(\vec{x}) \rangle_\beta$

$$\ell(\vec{x}) = \frac{1}{N_c} \text{Tr}_c \mathcal{P} \exp \left\{ i \int_0^\beta d\tau A_4(\vec{x}, \tau) \right\}$$

- **relation to confinement:**

$$\Phi \sim e^{-\beta F_q}$$

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$$\langle \bar{q}q \rangle \begin{cases} = 0 & \text{symmetric} \\ \neq 0 & \text{broken} \end{cases}$$

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## $Z_{N_c}$ Center Symmetry

- $m_q \rightarrow \infty$
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Other Phases ?

$$\Phi \begin{cases} = 0 & \text{confined} \\ \neq 0 & \text{deconfined} \end{cases}$$

# Polyakov-Quark-Meson Model

## Lagrangian

$$\begin{aligned} \mathcal{L}_{PQM} = & \bar{q} [i\cancel{D} - h(\sigma + i\gamma_5 \vec{\tau} \cdot \vec{\pi})] q + \frac{1}{2}(\partial_\mu \phi)^2 \\ & - U(\sigma, \vec{\pi}) - \mathcal{U}(\Phi, \bar{\Phi}) \end{aligned}$$

- $\phi = (\sigma, \vec{\pi}) \dots O(4)$ -representation of the meson field ( $N_f = 2$ )
- $\cancel{D}(\Phi) = \gamma_\mu \partial_\mu - i g \gamma_0 A_0(\Phi)$
- $g \dots$  gauge coupling
- $h \dots$  Yukawa coupling

## Meson Potential

$$U(\sigma, \vec{\pi}) = \frac{\lambda}{4}(\sigma^2 + \vec{\pi}^2 - v^2)^2 - c\sigma$$

# Polyakov Loop Potential

[C. Ratti, M.A. Thaler, W. Weise, Phys.Rev. **D73**, 014019 (2006)]

## Polynomial Ansatz

$$\frac{\mathcal{U}(\Phi, \bar{\Phi})}{T^4} = -\frac{b_2(T)}{2} \Phi \bar{\Phi} - \frac{b_3}{6} (\Phi^3 + \bar{\Phi}^3) + \frac{b_4}{4} (\Phi \bar{\Phi})^2$$

coefficients fitted to lattice data (pure glue):

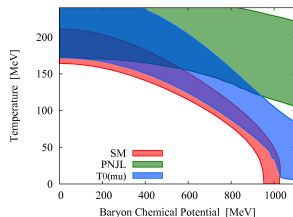
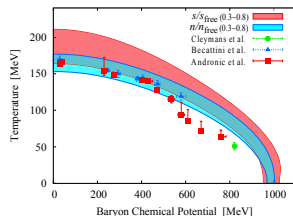
$$b_2(T) = a_0 + a_1 \left( \frac{T_0}{T} \right) + a_2 \left( \frac{T_0}{T} \right)^2 + a_3 \left( \frac{T_0}{T} \right)^3$$

| $a_0$ | $a_1$ | $a_2$ | $a_3$ | $b_3$ | $b_4$ |
|-------|-------|-------|-------|-------|-------|
| 6.75  | -1.95 | 2.625 | -7.44 | 0.75  | 7.5   |

$$T_0 = \begin{cases} 270 \text{ MeV (pure glue) ?} \\ 208 \text{ MeV ?} \\ \text{something else ?} \end{cases}$$

# $T_0(\mu)$ - One Motivation: Experiment

- experimental information on the QCD phase diagram:  
**chemical freezeout points**
- *not* raw data, but interpretation using  
**Statistical Model**
- increase of entropy (red band)  
and density suggests position of phase  
transition
- PNJL computation with  $T_0 = 200$  MeV  
inconsistent (green band)
- polynomial ansatz for  $T_0(\mu)$   
 $\Rightarrow$  greater overlap (blue band)



[in these plots:  $N_c = N_f = 3$ ]

[K. Fukushima, arXiv:1006.2596]

# $T_0(\mu)$ - Another Motivation: Theory

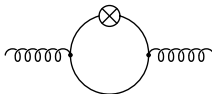
[B.-J. Schaefer, J.M. Pawłowski, J. Wambach, Phys.Rev. **D76**, 074023 (2007)]

**FRG flow** for QCD:

→ Talks by Jan Pawłowski and Lisa Haas

$$\partial_t \Gamma_k[\phi] = \frac{1}{2} \text{gluons} - \text{ghosts} - \text{quarks} + \frac{1}{2} \text{mesons}$$

dynamical quarks modify the gluon contribution:

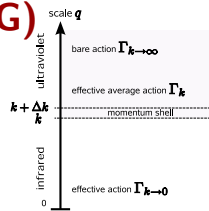


**Polyakov Loop potential:**  
from pure YM contribution

$$\partial_t \Gamma_k[\phi] = \frac{1}{2} \text{gluons} - \text{ghosts}$$

$$T_0 \rightarrow T_0(N_f, \mu) = T_\tau e^{-1/(\alpha_0 b(N_f, \mu))}$$

# Functional Renormalization Group (FRG)



## Flow Equation

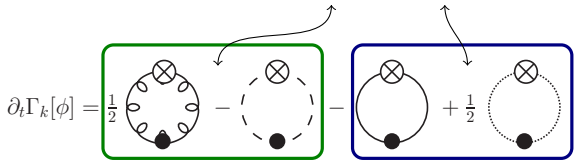
[C. Wetterich, 1993]

$$\partial_t \Gamma_k[\varphi] = \frac{1}{2} \text{Tr} \left\{ \partial_t R_{k,B} \left( \Gamma_k^{(2,0)}[\varphi] + R_{k,B} \right)^{-1} \right\}$$

## Polyakov–Quark–Meson Truncation

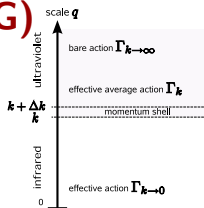
$$\Gamma_k = \int d^4x \left\{ \bar{q} (\not{D} + \mu\gamma_0 + ih(\sigma + i\gamma_5 \vec{\tau}\vec{\pi})) q + \frac{1}{2} (\partial_\mu \phi)^2 + \Omega_k[\sigma, \vec{\pi}, \Phi, \bar{\Phi}] \right\}$$

at initial scale  $\Lambda$ :  $\Omega_\Lambda[\sigma, \vec{\pi}, \Phi, \bar{\Phi}] = \mathcal{U}(\Phi, \bar{\Phi}) + U(\sigma, \vec{\pi}) + \Omega_\Lambda^\infty[\sigma, \vec{\pi}, \Phi, \bar{\Phi}]$





# Functional Renormalization Group (FRG)



## Flow Equation

[C. Wetterich, 1993]

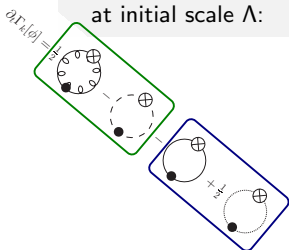
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$$\Omega_\Lambda[\sigma, \vec{\pi}, \Phi, \bar{\Phi}] = \mathcal{U}(\Phi, \bar{\Phi}) + U(\sigma, \vec{\pi}) + \Omega_\Lambda^\infty[\sigma, \vec{\pi}, \Phi, \bar{\Phi}]$$



$$\partial_k \Omega_\Lambda^k = \frac{-N_c N_f k^4}{3\pi^2 E_q} \left[ 1 - N_q(\Phi, \bar{\Phi}) - N_{\bar{q}}(\Phi, \bar{\Phi}) \right]$$

[J. Braun, K. Schwenzer, H.J. Pirner, Phys.Rev. **D70**, 085016 (2004)]  
 [V. Skokov, B. Stokic, B. Friman, Phys.Rev. **C82**, 015206 (2010)]

# PQM Flow Equation

[V. Skokov, B. Stokic, B. Friman, Phys.Rev. C82, 015206 (2010)]

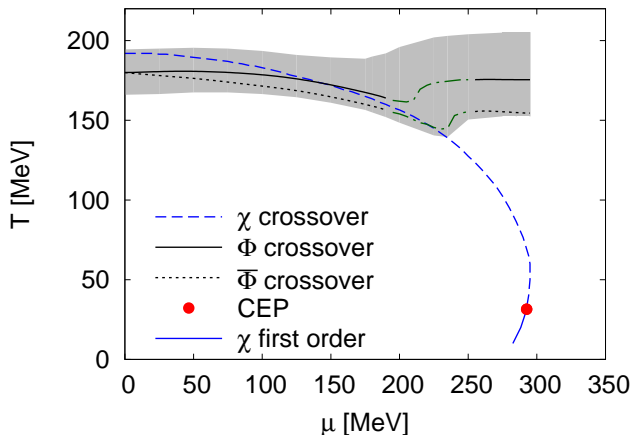
$$\partial_k \Omega_k(T, \mu) = \frac{k^4}{12\pi^2} \left[ \frac{3}{E_\pi} \coth\left(\frac{E_\pi}{2T}\right) + \frac{1}{E_\sigma} \coth\left(\frac{E_\sigma}{2T}\right) - \frac{2\nu_q}{E_q} \{1 - N_q(T, \mu; \Phi, \bar{\Phi}) - N_{\bar{q}}(T, \mu; \Phi, \bar{\Phi})\} \right]$$

$$N_q(T, \mu; \Phi, \bar{\Phi}) = \frac{1 + 2\bar{\Phi}e^{(E_q - \mu)/T} + \Phi e^{2(E_q - \mu)/T}}{1 + 3\bar{\Phi}e^{(E_q - \mu)/T} + 3\Phi e^{2(E_q - \mu)/T} + e^{3(E_q - \mu)/T}}$$

$$N_{\bar{q}}(T, \mu; \Phi, \bar{\Phi}) \equiv N_q(T, -\mu; \bar{\Phi}, \Phi)$$

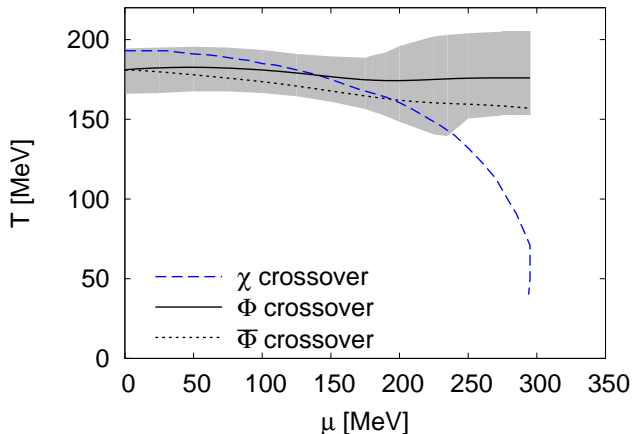
$$E_\pi = \sqrt{k^2 + 2\Omega'_k}, \quad E_\sigma = \sqrt{k^2 + 2\Omega'_k + 4\sigma^2\Omega''_k}, \quad \nu_q = 2N_c N_f$$

# Phase Structure $T_0 = 208 \text{ MeV}$ const.



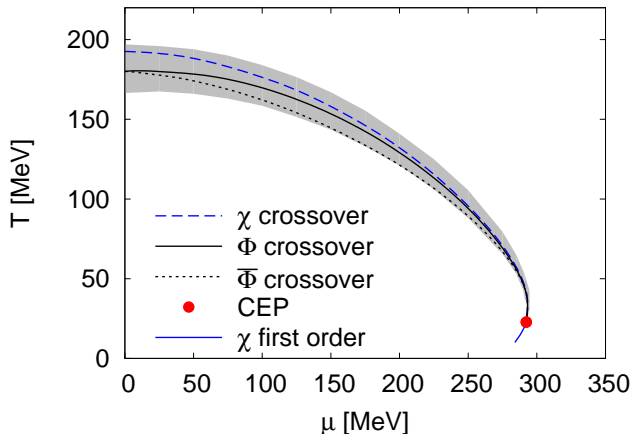
[TKH, J.M. Pawłowski, B.-J. Schaefer, arXiv:1008.0081]

# Phase Structure $T_0 = 208 \text{ MeV}$ const.



[TKH, J.M. Pawłowski, B.-J. Schaefer and M. Wagner, Work in Progress]

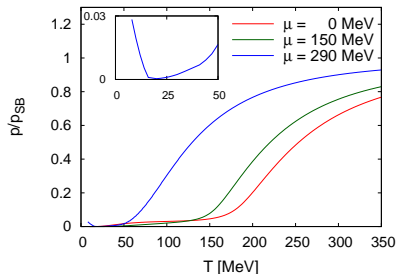
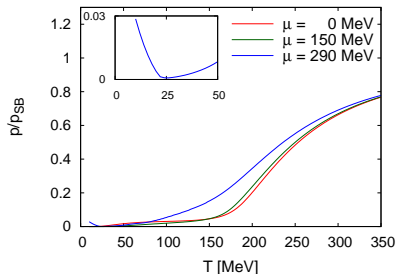
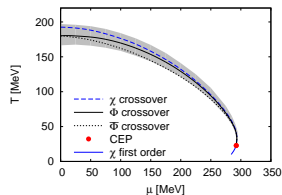
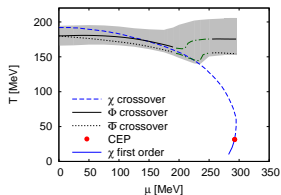
# Phase Structure $T_0(\mu)$ , $T_0(0) = 208 \text{ MeV}$



[TKH, J.M. Pawłowski, B.-J. Schaefer, arXiv:1008.0081]

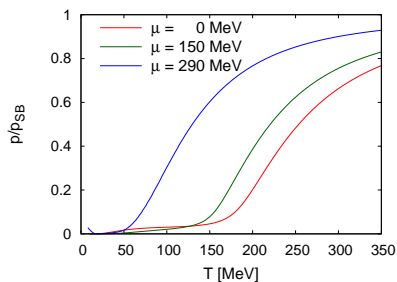
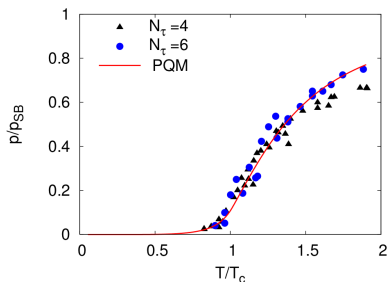
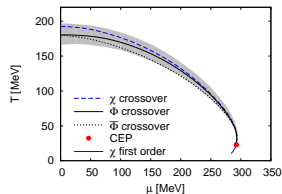
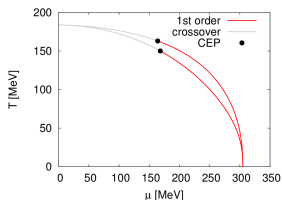
# Normalized Pressure

[TKH, J.M. Pawłowski, B.-J. Schaefer, arXiv:1008.0081]



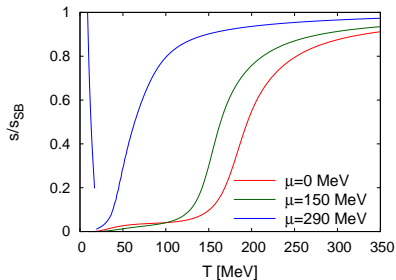
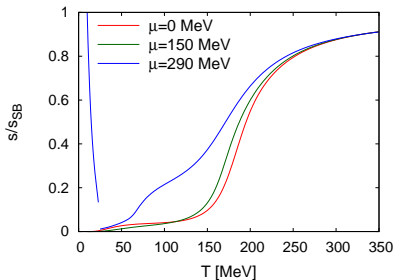
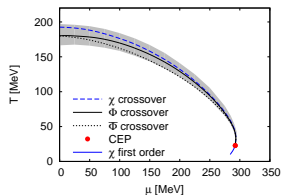
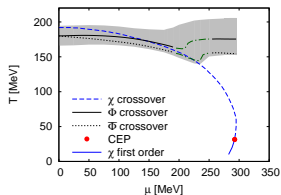
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 [B.-J. Schaefer, J.M. Pawłowski, J. Wambach, Phys.Rev. **D76**, 074023 (2007)]



# Normalized Entropy Density

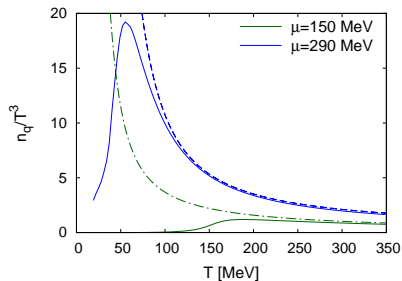
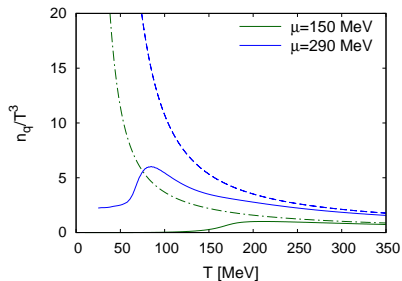
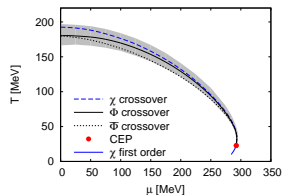
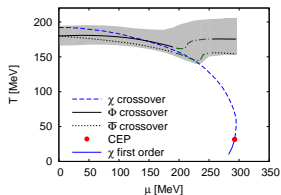
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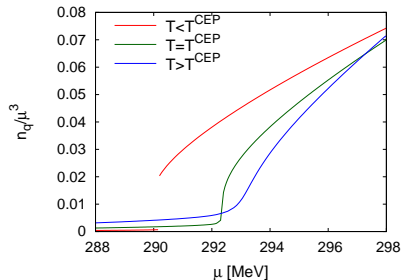
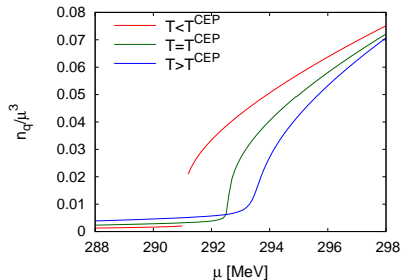
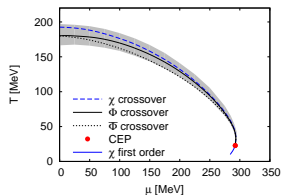
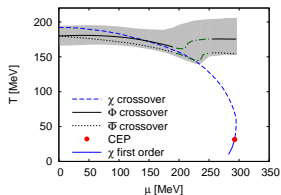
# Quark Number Density

[TKH, J.M. Pawłowski, B.-J. Schaefer, arXiv:1008.0081]



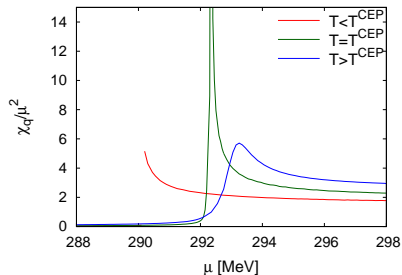
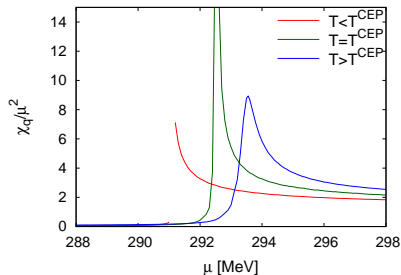
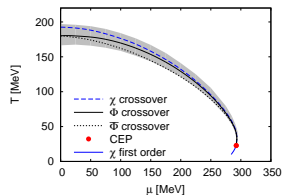
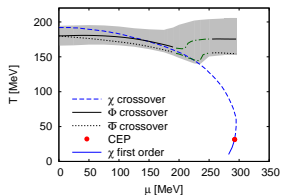
# Vicinity of the CEP

[TKH, J.M. Pawłowski, B.-J. Schaefer, arXiv:1008.0081]



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## Summary

- PQM model beyond Mean Field
  - quark-meson fluctuations included within FRG approach
  - Important feature: back-reaction to gluonic sector

$$T_0 \rightarrow T_0(N_f, \mu)$$

- Modifications of the Phase Structure
  - fluctuations push CEP downwards
  - $T_0(\mu)$ : chiral and deconfinement transitions coincide
  - $\rightsquigarrow$  no quarkyonic phase
- Thermodynamics
  - agree well with lattice studies at  $\mu = 0$

