

The generalized Ginsparg-Wilson relation for lattice supersymmetry and its consequences

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Outline

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- 2 The Ginsparg-Wilson-relation for SUSY
- 3 Constraints for lattice actions

Motivation

Why is lattice supersymmetry interesting?

- SUSY is a way to extend Poincaré symmetry.
- Part of the AdS/CFT correspondence.
- May be an ingredient of a GUT.
- Lattice yields results for analytically non-accessible aspects.
- Way to go beyond perturbation theory.

⇒ Fully supersymmetric lattice theory needed.

A fundamental problem

- Two SUSY generators anti-commute to derivative operator.
 \implies Successive SUSY transformations yield translation.
 \implies SUSY action invariant under (infinitesimal) translations.

Invariance in the continuum. Consider an action

$$S = \int dx \alpha(x)\beta(x)\gamma(x).$$

Small translation of the fields α , β , γ by δx yields

$$\delta S = \delta x \int dx [(\partial_x \alpha)\beta\gamma + \alpha(\partial_x \beta)\gamma + \alpha\beta(\partial_x \gamma)] = \delta x \int dx \partial_x (\alpha\beta\gamma) = 0.$$

In general, no Leibniz rule for lattice difference operator ∇ :

$$\nabla(\alpha_m \beta_m \gamma_m) \neq (\nabla \alpha)_m \beta_m \gamma_m + \alpha_m (\nabla \beta)_m \gamma_m + \alpha_m \beta_m (\nabla \gamma)_m;$$

Typical lattice action: no invariance under small translation.

\implies Not possible to obtain total derivatives like in continuum.

\implies How to construct invariant actions?

No-Go-Theorem: Leibniz rule only recovered by both non-local action and lattice derivative operator ([Bergner; 2010](#)).

\implies numerically expensive, ansatz not applicable to any theory;

State of the art: SUSY theories with one conserved supercharge.

A way out

Similar problem in lattice QCD with chiral symmetry:
Chiral lattice action has 'doublers' (two zero eigenvalues of difference operator, doubling of fermion species).

Introduction of modified symmetry relation ([Ginsparg, Wilson; 1982](#)):

$$\{D, \gamma_5\} = 2aD\gamma_5D, \text{ lattice spacing } a.$$

Relation for fermionic action obtained by block-spin-transformation.

Solution: overlap operator ([Neuberger; 1998](#)),

$$D_{\text{overlap}} = 1 + \gamma_5 \text{sgn}(\gamma_5 D).$$

GW-relation can be generalized for any linear symmetry.
 Consider lattice action $S[\phi]$ with fields ϕ_i , and infinitesimal symmetry operation

$$\phi_i \rightarrow \phi_i + \epsilon M_{ij} \phi_j,$$

with generator M .

Generalized GW-relation (Bergner, Bruckmann, Pawłowski; 2008):

$$\delta S \sim M_{ij} \phi_j \frac{\partial S}{\partial \phi_i} = (M \alpha^{-1})_{ij} \left[\frac{\partial S}{\partial \phi_j} \frac{\partial S}{\partial \phi_i} - \frac{\partial^2 S}{\partial \phi_j \partial \phi_i} \right],$$

with SUSY-breaking blocking matrix α .

Derived from a Wilsonian effective action ansatz:

$$e^{-S[\phi]} = S \det^{1/2} \alpha \int d\varphi e^{-\frac{1}{2}(\phi - \phi_f[\varphi]) \alpha (\phi - \phi_f[\varphi]) - S_d[\varphi]},$$

Application to lattice SUSY

Non-linear equation; mixes different orders of field products.

- For interacting SUSY action, not solved yet.
- Non-polynomial solutions likely.

Consider lattice supersymmetry operators M, \bar{M} .

$$\{M, \bar{M}\} \sim \nabla, \text{ with lattice derivative } \nabla.$$

One can show in general (Bergner, Bruckmann, Echigo, Igarashi, Pawłowski, Schierenberg: work in progress):

Lattice action fulfills relation for M and \bar{M} ;

\implies Relation is also fulfilled for anti-commutator ∇ .

$$\delta S \sim \nabla_{ij} \phi_j \frac{\partial S}{\partial \phi_i} = (\nabla \alpha^{-1})_{ij} \left[\frac{\partial S}{\partial \phi_j} \frac{\partial S}{\partial \phi_i} - \frac{\partial^2 S}{\partial \phi_j \partial \phi_i} \right];$$

$$\delta S \sim \nabla_{ij} \phi_j \frac{\partial S}{\partial \phi_i} = (\nabla \alpha^{-1})_{ij} \left[\frac{\partial S}{\partial \phi_j} \frac{\partial S}{\partial \phi_i} - \frac{\partial^2 S}{\partial \phi_j \partial \phi_i} \right];$$

Properties of α (and α^{-1}):

$$\alpha_{ij} = (-1)^{|i||j|} \alpha_{ji}; \quad |i| = \begin{cases} 0 & \text{for } i \text{ bosonic,} \\ 1 & \text{for } i \text{ fermionic.} \end{cases}$$

Bosonic and fermionic refers to product $\phi_i \alpha_{ij} \phi_j$.

Symmetrization:

$$(\nabla \alpha^{-1})_{ij} \rightarrow \frac{1}{2} (\nabla \alpha^{-1})_{ij} + \frac{1}{2} (-1)^{|i||j|} (\nabla \alpha^{-1})_{ji};$$

Natural assumptions:

- translational invariant operator ∇ and blocking α (commuting with each other);
- antisymmetric ∇ ;

Then:

$$\begin{aligned} (\nabla\alpha^{-1})_{ij} + (-1)^{|i||j|} (\nabla\alpha^{-1})_{ji} &= \nabla_{ik}\alpha_{kj}^{-1} + (-1)^{|i||j|}\alpha_{ki}^{-1}\nabla_{jk} \\ &= \nabla_{ik}\alpha_{kj}^{-1} - (-1)^{|i||j|+|i||k|}\alpha_{ik}^{-1}\nabla_{kj} = \nabla_{ik}\alpha_{kj}^{-1} - \nabla_{ik}\alpha_{kj}^{-1} = 0. \end{aligned}$$

Generalized GW-relation becomes

$$\delta\mathcal{S} \sim \nabla_{ij}\phi_j \frac{\partial\mathcal{S}}{\partial\phi_i} = 0.$$

\implies Naive relation for translational invariance, same problems like before.

Constraints for lattice actions

Consequences for supersymmetric lattice actions?

- With typical assumptions, generalized GW-relation yields no improvement.
- Relax some of our assumptions.

Possibilities:

- Choose α and/or ∇ translational non-invariant.
- Give ∇ a symmetric part.

1. Drop the translational invariance of α .
 \implies Resulting action not invariant under shifts of a lattice spacing.
2. Drop translational invariance of ∇ .
3. Give ∇ a symmetric part.
 \implies Violates a condition arising from blocking procedure.
 \implies Action very likely to be non-invariant like above.

Translational non-invariant action seems inevitable.

Summary

No-Go-Theorem: Local lattice action with SUSY in conventional way is impossible.

Way out: Generalized Ginsparg-Wilson-relation for SUSY.

However, relation also has to hold for anti-commutator of two supercharges

$$\{M, \bar{M}\} \sim \nabla.$$

Need to relax some natural assumptions to solve it.

\implies Translational non-invariant actions have to be taken into account.

Thanks for the attention.