

Functional RG for few-body physics

Michael C Birse
The University of Manchester

Review of results from:

Schmidt and Moroz, arXiv:0910.4586

Krippa, Walet and Birse, arXiv:0911.4608

Krippa, Walet and Birse, in preparation

Outline

- Background: effective field theories in nuclear and cold-atom physics
- Functional RG
- Three-body systems (unitary limit – scaling)
- Four-body systems
- Summary

Background

Effective field theory: now mainstream tool for nuclear forces

- symmetries of underlying dynamics (QCD)
 - systematic expansions in powers of low-energy scales (momenta, pion mass ...)
 - RG methods used to derive power countings
- classify terms as perturbations around fixed points (Wilsonian approach, sharp cut-offs)

Background

Effective field theory: now mainstream tool for nuclear forces

- symmetries of underlying dynamics (QCD)
 - systematic expansions in powers of low-energy scales (momenta, pion mass ...)
 - RG methods used to derive power countings
- classify terms as perturbations around fixed points (Wilsonian approach, sharp cut-offs)

2-nucleon scattering very strong at low-energies

(inverse scattering lengths $1/a \sim 10 - 40$ MeV)

- expand around nontrivial fixed point
“unitary limit” (zero-energy bound state: scale free)
- corresponds to effective-range expansion [Bethe (1949)]
- also describes atomic systems with Feshbach resonance tuned to threshold

3-body physics in unitary limit

Momentum space: one-variable integral equation

[Skornyakov and Ter-Martirosian (1956)]

Faddeev equation in hyperspherical coordinates

$$(R^2 = |\mathbf{r}_1 - \mathbf{r}_2|^2 + |\mathbf{r}_2 - \mathbf{r}_3|^2 + |\mathbf{r}_3 - \mathbf{r}_1|^2)$$

- Schrödinger equation with $1/R^2$ potential [Efimov, 1971]

$$-\frac{1}{M} \left[\frac{d^2}{dR^2} + \frac{1}{R} \frac{d}{dR} - \frac{v^2}{R^2} \right] u(r) = p^2 u(R)$$

- hyperangular eigenvalue v^2 fixed by boundary condition
(S-waves)

$$1 = \sigma \frac{4}{\sqrt{3\pi} v} \frac{\Gamma\left(\frac{1-v}{2}\right) \Gamma\left(\frac{1+v}{2}\right)}{\Gamma\left(\frac{3}{2}\right)} \sin\left(\frac{\pi v}{6}\right)$$

- spatially symmetric: $\sigma = +1$; mixed-symmetry $\sigma = -\frac{1}{2}$
(“particle-exchange interaction” between pair and third particle)

Three bosons or more than two “species” of fermion

(spatially symmetric state, eg triton)

- attractive $1/R^2$ potential ($v^2 < 0$)
- RG flow tends to limit cycle (discrete remnant of scale invariance)

[Moroz, talk at this meeting]

→ Efimov effect: tower of bound states with constant ratio between energies

- momentum scale factor $e^{\pi/s_0} \simeq 23$ where $s_0 = 1.00624$
- leading 3-body force is marginal
- fixes starting point on cycle or energy of one bound state (Phillips line relating ^3H binding to nd scattering length)
- Efimov states now observed for various alkali atoms

Other fermionic 3-body systems less interesting

- 3-body forces **irrelevant**
(no role in very-low-energy physics)
- but noninteger anomalous dimensions
- mixed symmetry state of two or more species of fermion:
leading 3-body force scales as $Q^{4.33244}$ (Q : momentum)
- determine energy eigenvalues in harmonic traps
[Werner and Castin (2005)]

4-body systems: no exact results

- 4-body forces irrelevant, even in Efimov systems
[Platter, Hammer and Meissner (2004)]
- low-energy 4-body physics determined by 2- and 3-body forces
(Tjon line relating ^4He and ^3H ; atomic Efimov systems)

Wilsonian RG methods used to obtain these

- exact results only for problems with 2-body-like forms
2-body systems; 3-body with contact interactions
not 3-body with long-range forces; 4-body

EFTs based on contact interactions

- not well suited for standard many-body techniques
 - difficult to apply to larger systems, dense matter
- new approaches needed

Functional (“exact”) RG

RG for the Legendre-transformed effective action Γ [Wetterich (1993)]
(1-particle-irreducible generating function)

Exact evolution equation has 1-loop structure

$$\partial_k \Gamma = +\frac{i}{2} \text{Tr} \left[(\partial_k \mathbf{R}_F) \left((\Gamma^{(2)} - \mathbf{R})^{-1} \right)_{FF} \right] \\ -\frac{i}{2} \text{Tr} \left[(\partial_k \mathbf{R}_B) \left((\Gamma^{(2)} - \mathbf{R})^{-1} \right)_{BB} \right]$$

$\Gamma^{(2)}$: matrix of second derivatives of the action

$\mathbf{R}_F(q; k)$: regulator for fermion fields; $\mathbf{R}_B(q; k)$: for bosons
(suppress modes with momenta $q < k$)

$\Gamma \rightarrow$ full effective action as $k \rightarrow 0$

Example: 2 species of fermion

Fermion field: $\psi(x)$ (neutrons or spin- $\frac{1}{2}$ atoms)

Boson “dimer” field: $\phi(x)$ (strongly interacting pairs)

Truncated ansatz for action in vacuum: 2-body, local

$$\begin{aligned} & \Gamma[\psi, \psi^\dagger, \phi, \phi^\dagger; k] \\ &= \int d^4x \left[\psi(x)^\dagger \left(i\partial_0 + \frac{\nabla^2}{2M} \right) \psi(x) \right. \\ & \quad \left. + Z_\phi(k) \phi(x)^\dagger \left(i\partial_0 + \frac{\nabla^2}{4M} \right) \phi(x) - u_1(k) \phi(x)^\dagger \phi(x) \right. \\ & \quad \left. - g \left(\frac{i}{2} \phi(x)^\dagger \psi(x)^T \sigma_2 \psi(x) + \text{H.c.} \right) \right] \end{aligned}$$

g : AA→D coupling

$u_1(k)$: dimer self-energy (u_1/g^2 : only physical parameter)

$Z_\phi(k)$: dimer wave-function renormalisation

Regulators

- fermions: sharp cutoff

$$R_F(\mathbf{q}, k) = \frac{k^2 - q^2}{2M} \theta(k - q)$$

- nonrelativistic version of “optimised” cutoff [Litim (2001)]
- fastest convergence at this level of truncation
- bosons

$$R_B(\mathbf{q}, k) = c_B \frac{k^2 - q^2}{4M} \theta(k - q)$$

- optimised choice $c_B = 1$ [cf Pawłowski (2007)]
(no mismatch between fermion and boson cutoffs)

Unitary limit

- $u_1(k) \rightarrow 0$ as $k \rightarrow 0$ ($a = -Mg^2 / (4\pi u_1(0)) \rightarrow \infty$)
- $Z_\phi(k) \rightarrow 0$ as $k \rightarrow \infty$ (dimer auxiliary field at starting scale)

3-body interaction (AD scattering)

$$\Gamma[\psi, \psi^\dagger, \phi, \phi^\dagger; k] = \dots - \lambda(k) \int d^4x \psi^\dagger(x) \phi^\dagger(x) \phi(x) \psi(x)$$

Evolution driven by terms corresponding to “skeleton” diagrams



(loops with AD contact interaction and single-A exchange)

Evolution equation

$$\partial_k \lambda = \frac{28k}{125g^2 M} \lambda^2 + \frac{156}{125k} \lambda + \frac{128g^2 M}{125k^3}$$

Rescale: $\hat{\lambda} = \frac{k^2}{g^2 M} \lambda$

- dimensionless equation

$$k\partial_k \hat{\lambda} = \frac{28}{125} \hat{\lambda}^2 + \frac{406}{125} \hat{\lambda} + \frac{128}{125}$$

→ two fixed point solutions (roots of RHS)

- expand around IR stable point: $\hat{\lambda} - \hat{\lambda}_0 \propto k^{3.10355}$
- compare exact solution: $Q^{4.33244}$

→ strongly irrelevant but only agree at 30% level

Bosons

Very similar action and evolution equations

- different numerical coefficients
 - $\partial_k \lambda$ term linear in λ gets factor of -2 (cf Faddeev equation)
- rescaled equation

$$k \partial_k \hat{\lambda} = \frac{56}{125} \hat{\lambda}^2 - \frac{62}{125} \hat{\lambda} + \frac{256}{125}$$

Bosons

Very similar action and evolution equations

- different numerical coefficients
 - $\partial_k \hat{\lambda}$ term linear in $\hat{\lambda}$ gets factor of -2 (cf Faddeev equation)
- rescaled equation

$$k \partial_k \hat{\lambda} = \frac{56}{125} \hat{\lambda}^2 - \frac{62}{125} \hat{\lambda} + \frac{256}{125}$$

→ two complex roots – fixed points

- expand around either: $\hat{\lambda} - \hat{\lambda}_0 \propto k^{\pm 2i s_0}$
- imaginary exponent → limit cycle [Moroz, talk at this meeting]
- real solutions periodic under scaling k by factor e^{π/s_0}
where $s_0 = 0.92503$
- agrees with Efimov $s_0 = 1.00624$ to $\sim 5\%$

4-body interactions (2 species of fermion)

Need DD \rightarrow DD, DD \rightarrow DAA, DAA \rightarrow DAA terms

(dimer “breakup” terms allow 3-body physics to feed in)

$$\Gamma[\psi, \psi^\dagger, \phi, \phi^\dagger; k] = \dots - \int d^4x \left[\frac{1}{2} u_2(k) (\phi^\dagger \phi)^2 \right. \\ \left. + \frac{1}{4} v(k) (\phi^{\dagger 2} \phi \psi^T \psi + \text{H.c.}) \right. \\ \left. + \frac{1}{4} w(k) \phi^\dagger \phi \psi^\dagger \psi^{\dagger T} \psi^T \psi \right]$$

\rightarrow coupled evolution equations for u_2, v, w (27 distinct skeletons)

Rescaled 4-body evolution equations

- 4 fixed-point solutions
- only one IR stable
- smallest eigenvalue $\rightarrow k^{4.19149}$ (irrelevant)

Bosons

- 4 complex fixed points (since λ complex)
 - only one IR stable
 - eigenvalue with smallest real part $\rightarrow k^{0.055165+3.50440i}$
- \rightarrow only weakly irrelevant ??
- couplings flow to cycle driven by $\lambda(k)$

Summary

First full applications of functional RG to 3- and 4-body systems

- local truncation, “optimised” cutoffs
- unitary limit: scaling behaviours agree with exact 3-body qualitatively for 2 species of fermion
much more accurately for bosons (Efimov effect)
- first estimates of anomalous dimensions for 4-body forces
bosons: real part puzzlingly small

Summary

First full applications of functional RG to 3- and 4-body systems

- local truncation, “optimised” cutoffs
- unitary limit: scaling behaviours agree with exact 3-body qualitatively for 2 species of fermion
much more accurately for bosons (Efimov effect)
- first estimates of anomalous dimensions for 4-body forces
bosons: real part puzzlingly small

Future work

- away from unitary limit [Krippa, talk at this meeting]
- 4 species of fermion – nucleons
SU(4) symmetry: evolution same as either bosons or 2 species
- use these 3-, 4-body interactions as input into calculations of dense matter (nuclear and cold atomic)
[Floerchinger, talk at this meeting]