

Corfu Summer Institute

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TEVATRON



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OUTLINE

The Machine & the Detectors: CDF & D0
 Physics Menu
 Tevatron Part I: Heavy Flavours

- Tevatron Part II: EWK & Higgs sector
- Why and How long will Tevatron still run?

The lectures will concentrate on just a few among the many important Physics topics currently under analysis in the CDF and D0 Collaboration. This choice is of course biased but it intends to show the audience what are still the important researches ongoing at the Tevatron and competitiveness/complementariness with the forthcoming LHC results.



TEVATRON is an OLD PAL



First collisions at 1.6 TeV c.m., observed by CDF in the VTPC central innermost tracker Courtesy of G.P. Yeh, the author of this online original Display...

OCT. 15 1985!

the Tevatron is running extremely well & still improving its performances04/09/2010TEVATRON, Corfu 2010, ASN4



- proton antiproton collider
- Center of mass energy in Run II: E_{cms}=1.96 TeV
- Run II machine cycle: 396 ns
- Dominated by:

quark-antiquark interactions (LHC: gluon-gluon)

-First collisions observed Oct 13, 1985 at 1.8 TeV

-CDF experiment observed these first collisions and the first top events in 1992 (first vertex detector in hadronic machine)

- D0 started at Run I in 1992
- Both CDF and D0 discover the top in 1995

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Collider Run II Peak Luminosity

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Since over about 3 decades (!!), several generations of experimentalists have worked on building and continuously upgrading CDF with innovative and pioneering ideas making it to be still today at the forefront and with various discoveries (top, Bs mixing, single top & ...) plus important breakthroughs both in Physics and Detector techniques (2008'Panofsky award for SVX-SVT +impacts: A.Menzione +L.Ristori, Pisa)



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- Excellent coverage of the Muon and tracking η <2
- Excellent calorimetry & electron ID
- High efficiency muon trigger with Pt measurement at Level 1 (toroid)
- 2T solenoid & 1.8 T toroidal reversed weekly

Summer 2006

(OK for µ asymmetry systematics)

Trigger	CDF	D0
2 Tracks	Pt> 2GeV/c	
	Pt1+Pt2>5.5GeV/c	No
	100µm <ld12l<1mm< th=""><th></th></ld12l<1mm<>	
1-muon	No	Pt(µ)>3, 4.5GeV/c
2-muon	Pt(µ's)>1.5GeV/c	Pt(µ's)>2.0 GeV/c

Triggering: an essential issue in hadron colliders => example of a continuously upgraded & sophisticated trigger system by CDF First time triggering on tracking system at first level (XFT) First time triggering on vertex Silicon tracker => Real time Physics



Very important contribution of Tracker in CDF Trigger Architecture



Upgrading the trigger to cope with the increase in luminosity is another crucial issue, well experienced at the Tevatron (see next). It will also be of course essential at the LHC when the luminosity is going to get to 10^{32} cm⁻² s⁻¹ and higher...



Gaseous tracking chamber rebuilt from run I to run II to cope with luminosity x100 and for the first time a tracking LV1 trigger:

eXtremelyFastTracker



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Fake rejection ~8

Upgrade SVT for luminosity





Upgrade: Faster SVT components and:

32Kpatterns \rightarrow 512Kpatterns new AM.



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OTEVATRON Part I

PRESENTS

HEAVY FLAVOURS



Importance of triggers in B Physics

Di-Muon (J/ψ) Pt(μ**)** > 1.5 GeV J/ψ modes down to low Pt(J/ψ) (~0GeV)

X(3872) $\rightarrow J/\psi \pi \pi$ β_{s} in $B_{s} \rightarrow J/\psi \phi$, Ξ_{b} , Ω_{b} Observation $\Lambda_{b} \rightarrow J/\psi \Lambda$ (masses, lifetimes) $B^{0}_{s,d} \rightarrow \mu \mu$ (rare decays) $Bc \rightarrow J/\psi \pi (J/\psi | X)$

Displaced trk + 2-Track Trig. lepton (e, μ) Pt(trk) > 2 GeV IP(trk) > 120µm IP(trk) > 100 μm Pt(lepton) > 4 GeV Fully hadronic modes Semileptonic modes $B \rightarrow hh$ (CP) High statistics lifetime B_S mixing Tagging studies, mixing D^o mixing Secondary $\Sigma^{*+} \rightarrow \Lambda_{b} \pi$ Vertex В $\Lambda_{\rm b} \rightarrow \Lambda_{\rm c} \pi$ Decay Length. $P_T(B) \ge 5 GeV$ Primary $L_{xy} \ge 450 \mu m$ Vertex d = impact parameter



Hadronic B Trigger (II)

- Run I collected $O(1) B_s \rightarrow D_s \pi$ in 100 pb⁻¹ (all D_s modes)
- Run II collected ~200/100 pb⁻¹ $B_s \rightarrow D_s \pi$ ($D_s \rightarrow \Phi[\rightarrow K^+K^-]\pi$)
- Compare with only 10x integrated luminosity!



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The Tevatron is a B and charm factory

➢Bs mesons={b,s} are only produced in hadron colliders. By several aspects, they are unique probes for beyond standard model searches and CP violation phenomena.

Tevatron is fully exploiting the physics potential of Bs mesons.

≻Tevatron gives access as well to b-baryons.

 CDF has the largest world's samples of B & charm with special thanks to SVX+SVT!! (see previous slides). Important also the Particle ID with COT (dE/dX) and the TOF => π and K identification
 CDF challenges B factories (on charged final states): 40M J/ψ (about 20% from B) > 6K B⁰_s → J/ψ φ, 32 K B⁰ → J/ψ K* 50 M D⁰ → K⁻π⁺, 12 K B⁰ → K⁻π⁺ etc.



RARE DECAYS





Example of search for $B_s \rightarrow \mu \mu$ rare decay

The events are searched in a sample of dimuon triggered events a set of discriminating variables, comparing data (sidebands, red) with MC generated events (black) are studied and trained in a NN:



 $\boldsymbol{B}(\mathsf{B}^{0}_{\mathrm{s},\mathrm{d}} \rightarrow \boldsymbol{\mu}^{+} \boldsymbol{\mu}^{-}) = [\mathsf{N}_{\mathrm{s},\mathrm{d}}/\mathsf{N}_{+}] \times [\boldsymbol{\alpha}_{+}/\boldsymbol{\alpha}_{\mathrm{s}}] \times [\boldsymbol{\varepsilon}_{+}/\boldsymbol{\varepsilon}_{\mathrm{s},\mathrm{d}}] \times [1/\boldsymbol{\varepsilon}^{\mathsf{NN}}_{\mathrm{s},\mathrm{d}}] \times f_{\mathrm{u}}/f_{\mathrm{s},\mathrm{d}} \times \boldsymbol{B}(\mathsf{B}^{+} \rightarrow \mathsf{J}/\Psi \mathsf{K}^{+})$

Where N_s=number of B⁰_s \rightarrow **µ**+**µ**- at 95%CL for N observed and N_b expected backg. α stands for the acceptance, ϵ for efficiency and f for fragmentation function 04/09/2010 TEVATRON, Corfu 2010, ASN







$$B_s \rightarrow \mu^+ \mu^-$$

- 6.1 fb⁻¹ of data analyzed
- Many improvements of analysis compared to previous published result
 - Improved muon identification and trigger selection
 - Bayesian NN is used instead of Likelihood ratio method
 - Limits are calculated in several bins of BNN variable and M(μμ)

Br($B_s \to \mu^+ \mu^-$) < 5.1×10⁻⁸(95% C.L.) Expected : 4.2×10⁻⁸(95% C.L.)





b→sµ+µ-

- Suppressed decays in SM: Br ~ O(10⁻⁶)
- New Physics in penguin or box diagrams modifies the decay-kinematics





Exclusive $b \rightarrow s\mu^+\mu^-$

Current results from analysis based on 4.4 fb⁻¹:

- Trigger on dimuons (p_T^{μ} >1.5-2 GeV) forming a displaced vertex
- Look for: B⁺ \rightarrow K⁺ $\mu^+\mu^-$, B_d \rightarrow K^{*0} $\mu^+\mu^ \rightarrow$ [K⁺ π^-] $\mu^+\mu^-$, B_s \rightarrow ϕ $\mu^+\mu^ \rightarrow$ [K⁺K⁻] $\mu^+\mu^-$
- − To reject B→J/ψ (ψ') remove: 8.68 < $M_{\mu+\mu-}^2$ < 10.09 U 12.86 < $M_{\mu+\mu-}^2$ < 14.18 GeV²
- Vertex quality, PID (dE/dx and TOF) + kinematic variables combined in a Neural Network to optimize sensitivity:
 - S/ $\sqrt{(S+B)}$ for known modes
 - S/(2.5+ \sqrt{B}) for B_s mode
- Normalize rate to $B \rightarrow J/\psi h$ (h=K,K^{*}, ϕ), used as 'reference' to obtain BR



$B \rightarrow K(*) \mu^+\mu^-$ signals



		B+→K+µ+µ-	$B^0 \rightarrow K^{*0} \mu^+ \mu^-$
	CDF	[0.38±0.05±0.03]x10 ⁻⁶	[1.06±0.14±0.09]x10 ⁻⁶
	Babar	[0.41±0.16±0.02]x10 ⁻⁶	[1.35±0.40-0.37±0.10]x10 ⁻⁶
	Belle	[0.53±0.08(+0.07-0.03]x10 ⁻⁶	[1.06+0.19-0.14±0.07]x10 ⁻⁶
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$B_s \rightarrow \phi \mu^+ \mu^-$ First Observation



- Unbinned Max Likelihood fit to B_scandidate mass
 - Signal: double Gaussian from MC (width adjusted to data)
 - Background: linear shape
- 27±6 signal events (6.3 σ)
- Normalize measured rate to $BR(B_s \rightarrow J/\psi\phi)$

$$BR(B_s \to \varphi \mu^+ \mu^-) = (1.44 \pm 0.33(stat.) \pm 0.46(sys.)) \cdot 10^{-6}$$

Compare to $BR(B_s \rightarrow \varphi \mu^+ \mu^-) \sim 1.6 \cdot 10^{-6}$ C. Q. Geng, C. C. Liu, J. Phys. G29, 1103 (2003)

$\mathbf{B}_{d} \rightarrow \mathbf{K}^{*0} \mu^{+} \mu^{-}$ angular analysis

A_{FB} and K³ are extract

- Θ_{u} : helici & the oppos - Θ_{κ} : angle to the B me

*⁰ longitudinal polarization
ted from:
ity angle between
$$\mu$$
+(μ -) direction
site of B (B) direction in $\mu\mu$ rest frame
e K direction and the direction opposite
eson in the K^{*0} rest frame
 μ^+

$$\frac{1}{\Gamma} \frac{d\Gamma(B^0 \to K^{*0} \mu^+ \mu^-)}{d\cos\theta_K} = \frac{3}{2} F_L \cos^2\theta_K + \frac{3}{4} (1 - F_L)(1 - \cos^2\theta_K), \quad \text{Fit Parameters}$$

$$\frac{1}{\Gamma} \frac{d\Gamma(B^0 \to K^{*0} \mu^+ \mu^-)}{d\cos\theta_\mu} = \frac{3}{4} F_L (1 - \cos^2\theta_\mu) + \frac{3}{8} (1 - F_L)(1 + \cos^2\theta_\mu) + A_{FB} \cos\theta_\mu$$

- Correct for acceptance via detailed simulation ٠
- 5 or 6 bin in $q^2 = m_{\mu\mu}^2$ excluding ccbar resonances ٠
- Cross check with $B \rightarrow J/\psi X$ (reference) samples •

Decay Angle Analysis





A_{FB} result $B_d \rightarrow K^* \mu^+ \mu^-$



$$A_{FB}^{(th)}(1 < q^2 < 6 \,\text{GeV}^2/\text{c}^4) = -0.05_{-0.04}^{+0.03}$$
 C.Bobeth, et al. arXiv : 1006.5013

Still inconclusive, wait for more data

Neutral B_s System

- Time evolution of B_s flavor eigenstates described by Schrodinger equation:

$$i\frac{d}{dt} \begin{pmatrix} |B_s^0(t)\rangle \\ |\bar{B}_s^0(t)\rangle \end{pmatrix} = \left(\mathbf{M} - \frac{i}{2}\mathbf{\Gamma}\right) \begin{pmatrix} |B_s^0(t)\rangle \\ |\bar{B}_s^0(t)\rangle \end{pmatrix}$$

-Diagonalize mass (*M*) and decay (Γ) matrices \rightarrow mass eigenstates :

$$|B_s^H\rangle = p \,|B_s^0\rangle - q \,|\bar{B}_s^0\rangle \qquad |B_s^L\rangle = p \,|B_s^0\rangle + q \,|\bar{B}_s^0\rangle$$



- Flavor eigenstates differ from mass eigenstates and mass eigenvalues are also different:

 $\Delta m_s = m_H - m_L \approx 2/M_{12} |$ $\rightarrow B_s \text{ oscillates with frequency } \Delta m_s \\ \text{ precisely measured by} \\ \text{CDF } \Delta m_s = 17.77 + -0.12 \text{ ps}^{-1} \\ \text{DØ } \Delta m_s = 18.56 + -0.87 \text{ ps}^{-1}$



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- Mass eigenstates have different decay widths

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$$\Delta \Gamma = \Gamma_L - \Gamma_H \approx 2|\Gamma_{12}|\cos(\Phi_s) \quad \text{where} \quad \phi_s^{SM} = \arg\left(-\frac{M_{12}}{\Gamma_{12}}\right) \approx 4 \times 10^{-3}$$

$\beta_{\rm s}$ vs $\phi_{\rm s}$

- Up to now, introduced two different phases:

$$\phi_s^{SM} = \arg\left(-\frac{M_{12}}{\Gamma_{12}}\right) \approx 4 \times 10^{-3} \qquad \text{and} \qquad \beta_s^{SM} = \arg\left(-V_{ts}V_{tb}^*/V_{cs}V_{cb}^*\right) \approx 0.02$$

- New Physics affects both phases by same quantity $\phi^{
m NP}_s$

$$2\beta_s = 2\beta_s^{\rm SM} - \phi_s^{\rm NP}$$
$$\phi_s = \phi_s^{\rm SM} + \phi_s^{\rm NP}$$

- If the new physics phase ϕ_s^{NP} dominates over the SM phases $2\beta_s^{SM}$ and $\phi_s^{SM} \rightarrow$ neglect SM phases and obtain:

$$2\beta_s = -\phi_s^{\bar{\mathrm{NP}}} = -\phi_s$$

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CP Violation in $B_s \rightarrow J/\Psi \Phi$ Decays

- Analogously to the neutral B^0 system, CP violation in B_s system occurs through interference of decays with and without mixing:



- *CP* violation phase β_s in SM is predicted to be very small, $O(\lambda^2)$
- New physics particles running in the mixing diagram may enhance β_s
 - large $\beta_s \rightarrow$ clear indication of New Physics !

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Finding New Physics in Φ s

- Measurement of semileptonic CP asymmetry $\leftrightarrow A_{SL} = -|\Gamma_{12}/M_{12}|\sin \phi_s$
- \rightarrow Difficult due to smallness of the A_{SL}
- Measurement of $\Delta \Gamma \leftrightarrow \Delta \Gamma = \Gamma_L \Gamma_H = 2|\Gamma_{12}|\cos(\phi_s)$
- → Need input for $|\Gamma_{12}| \leftrightarrow$ can come from $\mathcal{B}(B_s \to D_s^{(*)}D_s^{(*)})$
- Measurement of CP violation in CP modes
- \rightarrow Golden mode pursued currently is $B_s \rightarrow J/\psi \phi$
- → φ_s not only phase entering here, but SM value small → search for sizable NP can neglect SM value

$$= 2\beta_{s} = -\phi_{s}^{\mathsf{NP}} = -\phi_{s}$$



- a_{sl}^q is the charge asymmetry of "wrong sign" semileptonic B_q^0 (q = d, s) decays:

$$a_{sl}^{q} \equiv \frac{\Gamma(\overline{B}_{q}^{0} \to \mu^{+}X) - \Gamma(B_{q}^{0} \to \mu^{-}X)}{\Gamma(\overline{B}_{q}^{0} \to \mu^{+}X) + \Gamma(B_{q}^{0} \to \mu^{-}X)}; \quad q = d, s$$

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CP violation in mixing

- Non-zero A^b_{sl} means CP violation in mixing
- Source of this type of CP violation $\overline{B}_{q} W$ complex phase ϕ_{q} of $B_{q} (q=d,s)$ mass matrix $\overline{\overline{q}} = \overline{u}, \overline{c}, \overline{t}$

$$\Delta M_q = M_H - M_L \approx 2 \left| M_q^{12} \right|$$
$$\Delta \Gamma_q = \Gamma_L - \Gamma_H \approx 2 \left| \Gamma_q^{12} \right| \cos \phi_q$$
$$\phi_q = \arg \left(-\frac{M_q^{12}}{\Gamma_q^{12}} \right)$$

$$\left\|\mathbf{M}_{q}\right\| = \begin{bmatrix} M_{q} & M_{q}^{12} \\ (M_{q}^{12})^{*} & M_{q} \end{bmatrix} - \frac{i}{2} \begin{bmatrix} \Gamma_{q} & \Gamma_{q}^{12} \\ (\Gamma_{q}^{12})^{*} & \Gamma_{q} \end{bmatrix}$$

u,c,t

q

 \overline{h}

 B_q

W

• For B_q meson, a_{sl}^q is related to the CP-violating phase ϕ_q :

$$a_{sl}^{q} = \frac{\Delta \Gamma_{q}}{\Delta M_{q}} \tan(\phi_{q})$$

SM prediction

• SM predicts very small values of ϕ_q and A^b_{sl} :

$$\phi_d^{SM} = -0.091^{+0.026}_{-0.038}$$
$$\phi_s^{SM} = 0.0042 \pm 0.0014$$
$$A_{sl}^{b,SM} = (-2.3^{+0.5}_{-0.6}) \times 10^{-4}$$

A. Lenz, U. Nierste, J. High Energy Phys. 0706, 072 (2007)

- These values are below current experimental sensitivity
- New physics contribution can significantly change these values

$$\phi_d = \phi_d^{SM} + \phi_d^{NP}$$
$$\phi_s = \phi_s^{SM} + \phi_s^{NP}$$

Non-zero A^{b}_{sl} would indicate the presence of new physics



Measurement strategy

• Measure two raw asymmetries (include µ's from all sources):

raw dimuon charge asymmetry

$$A = \frac{N(\mu^{+}\mu^{+}) - N(\mu^{-}\mu^{-})}{N(\mu^{+}\mu^{+}) + N(\mu^{-}\mu^{-})}$$

$$= (0.564 \pm 0.053)\%$$

raw inclusive muon charge asymmetry

$$a = \frac{n(\mu^+) - n(\mu^-)}{n(\mu^+) + n(\mu^-)}$$

= (0.955 ± 0.003)%

 Both asymmetries contain contributions from A^b_{sl} and detector-related background asymmetries

$$A = K A_{sl}^b + A_{bkg}$$

$$a = k A_{sl}^b + a_{bkg}$$

- contribution from A_{sl}^{b} to *a* is strongly suppressed by *k*=0.041±0.003

- Determine background contributions A_{bkg} and a_{bkg} using data with minimal input from simulation
- Exploit the correlation of background content in raw asymmetries to reduce the uncertainty on A^b_{sl}
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Test of background description

- Raw inclusive muon asymmetry *a* is dominated by the background asymmetry a_{bkg}
- a_{bkg} is measured in data
- Compare *a* and *a_{bkg}* to verify the background description
- This comparison is done as a function of muon p_T
- Good consistency between observed and expected asymmetries
 - $-\chi^2/dof = 2.4/5$ for the difference between these two distributions





Original experimental technique

- Polarities of DØ solenoid and toroid are reversed every ~2 weeks
- 4 equal sized samples with different polarities (++, --, +-, -+) Swapping Magnet Polarity
- difference in reconstruction efficiency between positive and negative particles minimized
- Reconstruction asymmetries reduced from ~1% to <0.1%
 - To be compared with raw dimuon asymmetry A= (0.564±0.053)%



Muon reconstruction asymmetry

X



Evidence for an anomalous like-sign charge asymmetry

 $A_{sl}^{b} = (-0.957 \pm 0.251 \,(\text{stat}) \pm 0.146 \,(\text{syst}))\%$

- This result differs from the SM prediction by ~3.2 σ
- A^b_{sl} produces a band in a^d_{sl} v.s. a^s_{sl} plane:

 $A_{sl}^{b} = (0.506 \pm 0.043)a_{sl}^{d} + (0.494 \pm 0.043)a_{sl}^{s}$

 Obtained result agrees well with other measurements of a^d_{sl} and a^s_{sl}



Bs mixing phase: Bs→J/ψφ best testing Lab

B_s mixing phase (or V_{ts} phase) is the last experimentally unconstrained part of CKM



BUT, additional experimental complications:

J/ $\psi \phi$: a mix of CP-even and CP-odd eigenstates, treat them separately B_s oscillates ~ 35 times faster than B⁰ sin2 β ~0.7, sin2 β _s expected about 20 times smaller

Bs mixing phase in $Bs \rightarrow J/\psi\phi$: analysis outline



EDP

Mixing phase - fit overview

Courtesy of Diego Tonelli

5. Perform unbinned maximum likelyhood fit to extract interesting quantities







New measurement of $B_s \rightarrow J/\psi \phi$ (2010)

- 5.2 fb⁻¹ of data analyzed
- ~6500 signal events
- New trigger will be soon added (increase of stat.)
- Same side flavour tagging calibrated in data
- Strong phases are free
- S wave included in the fit < 6.5% at 95% CL

 $\tau_s = 1.529 \pm 0.025 \text{ (stat)} \pm 0.012 \text{ (syst) ps}$ $\Delta \Gamma_s = 0.075 \pm 0.035 \text{ (stat)} \pm 0.01 \text{ (syst) ps}^{-1}$

Most precise measurements of $\tau(B_s)$ and $\Delta\Gamma_s$



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Evolution of the Results with increased data set

- Agreement with SM expectation increases with higher statistics

- β_{s} and $\Delta\Gamma$ allowed parameter space greatly reduced





Comparison Between Different Data Periods

- Divide 5.2 fb⁻¹ sample in three sub-samples corresponding to three public releases:
 - 0 1.4 fb⁻¹ (initial result released at the end of 2007, PRL 100, 161802 (2008)
 - 1.4 2.8 fb⁻¹ (added for 2008 ICHEP update)
 - 2.8 5.2 fb⁻¹ (added for this update)
- Previous results reproduced with updated analysis
- Clearly, improved agreement with the SM expectation comes from the second half of data $(2.8 5.2 \text{ fb}^{-1})$





B_s Lifetime and Decay Width Difference

- Assuming no *CP* violation ($\beta_s = 0$) obtain most precise measurements of lifetime τ_s and decay width difference $\Delta \Gamma_s$:

$$c au_s = 458.6 \pm 7.5 \text{ (stat.)} \pm 3.6 \text{ (syst.)} \ \mu\text{m}$$

 $au_s = 1.53 \pm 0.025 \text{ (stat.)} \pm 0.012 \text{ (syst.)} \text{ ps}$
 $\Delta\Gamma = 0.075 \pm 0.035 \text{ (stat.)} \pm 0.01 \text{ (syst.)} \ ps^{-1}$

PDG average:

$$au = (1.472^{+0.024}_{-0.026}) imes 10^{-12}$$
 s



Mixing phase – *results & prospects*



NEW: p-value is 44% (0.8 σ)

Mixing phase – *results & prospects*



Mixing phase — *results & prospects*





WHY we became interested in B Physics at Hadron Colliders

HF Physics at Hadron colliders: UA1 = precursor
 HF Physics at Hadron Colliders: CDF = pioneer:
 1st microvertex in hadronic colliders

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 Important breakthroughs and much more still to come...
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 ➢HF Physics at Hadron colliders: A goldmine to look for NP (provided the right detector & trigger!!!!), and:
 SM picture of CPV satisfactory at least at tree level in B0 and B+ decays. NP Physics at at the level of ≤ 10% effect.

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To keep NP-scale in the TeV range, Physics BSM should have a highly fine-tuned Flavour structure: Flavour problem ???



TEVATRON Part II

REVEALING THE HIGGS SECTOR

View Prom Agni Travel Office, 30 December 2900 - 10:22 -

REVEALING THE HIGGS SECTOR



This part presents some of the instrumental searches and breakings in the exploration of the Higgs sector at the Tevatron



Higgs sector demands it all from the detectors



A dedicated new Higgs trigger & N.P. searches





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The upgraded Calorimeter Trigger

courtesy of Simone Donati





The L1MET Upgrade





Effects of latest trigger upgrades on L1 and L2: 2 examples





This impacts also on the

τ- triggers

r1:lumi

Courtesy Laura Sartori



New trigger strategies for the Higgs

The calorimeter trigger upgrade together with XFT-3D, SVT upgrades significantly improves CDF reach for the Higgs (and lots of Physics topics)



Ex: New Jet clustering provides

Mode	Acceptance increase
WH → evbb	+97 %
$\begin{array}{c} \mathbf{WH} \rightarrow \\ \mu\nu\mathbf{bb} \end{array}$	+110 %
ZH → e+e-bb	+27 %
$ZH \rightarrow \mu+\mu-bb$	+60 %
$ZH \rightarrow vvbb$	+30 %
$\begin{array}{c} H \rightarrow \\ IvIv \end{array}$	+24 %



Designed and installed by: Davis, LPNHE, PISA, Rutgers, TA&M in 2002

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SM Consistency Check: M_w, M_{top}, M_H

> The W boson mass is derived from precisely measured EWK quantities



W detection at Tevatron


W-Mass Measurement at the Tevatron

•Tevatron: mainly $\overline{q}q'$ annihilation



Z decay into leptons provide an excellent calibration sample:

- $Z \rightarrow \mu \mu$ (CDF) and $Z \rightarrow ee$ (CDF, D0) events to derive recoil model

- most systematic uncertainties limited by size of Z sample (10x less Z than W) 04/09/2010 TEVATRON, Corfu 2010, ASN 73

Lepton Energy/Momentum Scale



- Calibrate calorimeter using precisely known M_z from LEP
- Detailed corrections for non instrumented regions





- CDF:
- Calibrate lepton momentum scale using Y, J/ ψ , Z data
- Calibrate calorimeter against precision tracker (E/p) and M_z



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W-Mass Results





SystematicSource	δm _w (MeV)	
Electron energy scale	34	
Electron energy resolution model	2	Imp
Electron energy nonlinearity	4	stat
W and Z electron energy loss differences	4	
Recoilmodel	6	
Electron efficiencies	5	
Backgrounds	2	1.114.5
PDF	9	
QED	7	
Bosonpt	2	
Total	37	

Improve w/ statistics

Ultimately limit precision

Tevatron Run II precision goal: $\Delta m_W < 25$ MeV/experiment





Tevatron W Mass Combination

- New Tevatron combination:
 - $m_W = 80420 \pm 31 MeV$ (0.038%)
 - Update previous CDF result to modern PDFs
 - Correct to same Γ_W
 - PDF, QED, Γ_W uncertainties correlated
- More precise than LEPII legacy: 80.367±0.033 GeV (0.04%)
- New world average (2009):

$m_W = 80399 \pm 23 \text{ MeV}$

CDF (2.4 fb⁻¹) and D0 (5fb⁻¹) are working on updates with improved analysis => expected to go \leq 25 MeV per experiment



=> W-MASS: ONE of the OUTSTANDING LEGACY OF THE TEVATRON

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The Top in itself and as a window to New Physics





TOP: A little bit of history...

1992: FIRST OBSERVATIONS!



29 Oct 1992, first observation of TOP by CDF in the Run I in a golden e μ event and the vertex detector signing b-jet! Also found in e+jets channel (see next slide)

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The Many ways to find New Physics with the Top



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Top Production at the Tevatron: dominated by q qbar



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TOP Quark Pair Production at Tevatron



- Dilepton (lepton = e or μ) (6%)
 - Small rate, small backgrounds
 - Main background: Drell-Yan
 - Highest purity
- Lepton+Jets (lepton = e or μ) (34%)
 - Good rate and manageable backgrounds
 - Main background: W+jets,
 - Good purity "Golden Channel"

- All-hadronic (46%)
 - Large rate, large background
 - Main background: QCD multijet
 - Least purity
- Hadronic Taus (tau+lepton, tau+jets) (14%)
 - Small rate and large backgrounds
 - Main background: Multijets, W+jets
 - Challenging purity



TOP MASS Measurement Methods

Two main approaches:

> The Template Method

- Choose an observable x, sensitive to m_{top}
- Predict x-distribution = f(m_{top}) using MC (templates)
- For each event evaluate a Likelihood for each m_{top} value
- Maximize the likelihood for the entire sample

Exemple of variables used:

- o m_{top} after kinematic reconstruction
- o Lepton P_t
- o Lxy: decay length of b-quarks in the event



all hadronic

> The Matrix Element Method (ME)

Uses all information from the event, integrating over the least known variables

Systematics are somewhat different for different methods

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> Systematics are somewhat different for different methods

Top mass extraction with ME method

ME Method is based on a full LO calculation of top pair production + detector effects
 Per event, the probability to be produced under assuming a given top mass via signal process is:



- Calculate main background probability in a similar way
- Build event probability by adding: normalized signal + background probabilities
- Determine top mass from likelihood fit of the event probabilities 04/09/2010
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Top Mass example: ME applied to lepton +jets

The largest uncertainty on all Top quark mass measurements is due to JET ENERGY SCALE (JES) uncertainties

The measurement of an overall jet energy scale correction **JES** on top of the standard corrections in lepton+jets is possible because of well known W mass





Best present result in a single experiment

Measurement at CDF on lepton + jets data with 5.6 fb⁻¹ analyzed data





Top Mass: Tevatron Combination, July 2010



Multivariate Analyses: BDT & BNN

Use common Object and Event Kinematics, Angular Correlations, Jet Reconstruction and Top Quark Reconstruction variables

Boosted Decision Trees (BDT)

- Recover events that fail criteria in cut-based analysis
- Boosting averages the results over many trees, improving the performance
- Uses highest ranked common 64 variables





Bayesian Neural Network (BNN)

- NN train on signal and background, producing one output discriminant
- Bayesian NN average over many networks, improving the performance
- Uses highest ranked 18-28 variables in each channel

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Multivariate analysis: Matrix Element (ME)

- Use the 4-vectors of all reconstructed leptons and jets
- Use Feynman diagrams to compute an event probability density for signal and background hypotheses
- Uses events with 2 and 3 jets only
- ME for signal (tb & tqb) and background
- Split the sample in high and low H_T (W+jets and top quark pair dominated regions) improves the performance

Without forgetting maximum Likelihood.....

⇒Training for Higgs searches; most of the developed tools were then applied to the search for Higgs(ses): see later 04/09/2010 TEVATRON, Corfu 2010, ASN





combine up to 12 different analysis channels:







combine up to 8 different analysis channels:

single top

■ F₊+jets selection :

recover badly reconstructed e, μ; include τ





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Going down to smaller cross section processes (cont'd)

- High luminosity samples allow for measurements of low cross section processes
- provided also sophisticated triggers & analyses.

DIBOSON PRODUCTION:

- CDF & D0 have observed WW,WZ, and ZZ production via leptonic decay channels
- First observations of WW/WZ production in the semi-leptonic decay channel
- These measurements are critical for validating our Higgs search techniques



Triple Gauge Coupling (TGC) interest per se and as background for Higgs searches







Constraints on the SM Higgs Boson

What we know from:

•Direct search at LEPII:

Mh > 114 GeV/c2 @95% CL

•Precision EWK meaurements (top mass, W mass, etc):

 $M_h = 89.0^{+35} - 26 \text{ GeV/c}^2$

 $M_h < 158 \text{ GeV/c}^2 @95\% \text{ CL}$



The many ways to produce a light Higgs at the Tevatron



Higgs Production at the Tevatron: the main channels



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LOW MASS HIGGS: Production and Decay



Low Mass Final States

$$WH \rightarrow \ell \nu b\bar{b}$$

$$ZH \rightarrow \ell \ell b\bar{b}$$

$$2 \text{ High } P_{T} \text{ Leptons } + b \text{ jets}$$

$$2 \text{ High } P_{T} \text{ Leptons } + b \text{ jets}$$

$$WH \rightarrow (\ell) \nu b\bar{b}$$

$$WH \rightarrow (\ell) \nu b\bar{b}$$

Efficiency for tagging b-quark jets is critical as well as rate for mis-tagging light quark jets04/09/2010TEVATRON, Corfu 2010, ASN104



The Strategy

Attack at every opportunity! And use all means!

1) Maximize acceptance

"Tight"cuts on final state objects (leptons, jets, etc.) are too costly Valuable signal efficiency is available at the cost of additional hard work! Includes triggers as well.

2) Minimize backgrounds

- Often a good model of the backgrounds are needed to do this some may come from the data itself
- 3) Maximize measurement resolution
- Improvements in jet energy resolution directly impact acceptance, backgrounds, and many other aspects of the analyses
- 4) Extract the most from the resulting events
- Use advanced algorithms like decision tree, neural networks, matrix elements, etc

Maximizing Acceptance Example: $ZH \rightarrow \ell \ell b \bar{b}$

Reconstruct Z candidate using two <u>loose</u> muons AND nonmuon trigger!

Reconstruct Z candidate using a µ and an isolated track



Typical gain from increased lepton ID ~15%

Tagging Jets from b-quarks

- Idea is to pick out jets consistent with having originated from long-lived bquark (typically involves reconstruction of secondary track vertex within jet)
- More sophisticated tools incorporating both jet shape and track variables are used to obtain additional discrimination between different flavor quark jets
- CDF uses a variety of secondary vertexing algorithms

> D0 uses NN tagger based on 7 discriminating **B-lifetime variables**

Btagging eff: ~ 50-70% Depends on Jet(ET& η) •Mistag rates ~0.3 - 6.0% Loose tagging if double tag





New: a neural network Flavor Separator
Modeling of Backgrounds



Using Advanced Algorithms

Artificial Neural Networks (ANN), Boosted Decision Trees (BDT), Matrix Element:ME

- Ex: ANN used to combine information from different kinematic variables both Energy-based and Shape-based
- Improved discrimination and less sensitive to systematic effects
- Tested using already observed physics processes: identification of top in Lepton+ jets



Final Discriminants



Applications of Advanced Algorithms

Diboson observation : $WW + WZ \rightarrow Ivjj$ Similar to $WH \rightarrow Ivbb$ Matrix Element : $Us: \sigma(WW+WZ) = 16.6+3.5-3.0 \ pb$ $SM: \sigma = 15.1 \pm 0.8 \ pb$ Single top observation $t+q \rightarrow lvb+j$ (with b-tag) Similar to WH $\rightarrow lvbb$ Neural Network : $Us : \sigma(t) = 4.70 + 1.18 - 0.93 \, pb$ $SM : \sigma = 3.46 \pm 1.8 \, pb$





Features:

- 1. Small $\sigma \diamond BR$
- 2. Several tight constraints

- ii. " E_7 " \rightarrow improve jet resol.
- 1. \sim 1 evt/6 fb⁻¹ (dbl tags)

Primary Backgrounds

$$Zbb, Zc\bar{c}, Zqq'$$

 $t\bar{t}$
 $WW + jj, WZ, ZZ$
 $Z \rightarrow \tau\tau$





140 160 180 200

Dijet Mass [GeV]

10님:

Գ

20

60

40

10

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20

40

60

80 100 120

80 100 120 140 160 180 200

Dijet Mass (Kinematic Fit) [GeV]





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 $ZH \rightarrow \ell\ell bb \ Event \ Displays$





 $M_h = 115 \; GeV/c^2$

Experiment	Luminosity	Obs/SM	Exp/SM
D0	6.2 fb ⁻¹	8.0	5.7
CDF	5.7 fb⁻¹	6.0	5.5

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$ZH \rightarrow \nu \nu b \overline{b}$



$ZH \rightarrow \nu\nu bb Background Modeling$ Events/10



divides data into 5 subsamples: Signal used to search for Higgs

- QCD Region for systematic studies
- •EWK region for modeling ewk processes
- •QCD Regions (2) to check normalization of MultiJet (MJ)



EW Control sample (two asymmetric btags) 60 Events / 25.00 D0 Preliminary (5.5 fb⁻¹) 🕂 Data 50 divides data into 4 subsamples: Top •Signal used to search for Higgs V+h.f./VV 40 •MJ-model for modeling MJ background V+I.f. 30 in signal sample 20 MJ-control to validate MJ-modeling •EWK control, enhanced in W->uv 10 04/09/2010 **TEVATRON**, Corfu **%** 50 100 150 200 250 300 **DiJet Invariant Mass (GeV)**





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NN_{QCD}



$ZH \rightarrow \nu\nu b\bar{b}$ Event Display



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 $ZH \rightarrow \nu\nu b\bar{b}$ Results



Experiment	Lum	Obs/SM	Exp/SM	$M_{\rm b} = 115 {\rm GeV}/c^2$
CDF	5.7 fb ⁻¹	2.3	4.0	
D0	6.4 fb ⁻¹	3.4	4.2	

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$WH \rightarrow \ell \nu bb \ Procedure$



CDF uses 2 methods:
•e/μ MET, 2/3 jets
•double loose or single tight tagging
•Flavor separator for single tags
•Two discriminants
1) Matrix Element
2) Bayesian Neural Network (adds in single isolated tracks as well)





 $WH \rightarrow \ell \nu bb$ Results







Experiment	Lum	Obs/SM	Exp/SM	
CDF	5.7 fb ⁻¹	3.6/4.5	3.5/3.4	$M_{h} = 11$
D0	5.3 fb ⁻¹	3.7	4.7	

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 5 GeV/c^2

Tevatron Candidate Summary, $m_H=115 GeV$



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"HIGH" MASS LIGHT HIGGS

H->WW decay in: VBF \rightarrow VWW

 $VH \rightarrow VWW \quad (SS \text{ leptons})$ $VH \rightarrow VWW \rightarrow l^{\pm}vl^{\pm}v + X$

 $gg \rightarrow H \rightarrow WW$ (OS leptons)

Both CDF and D0:

Selection: 2 opposite -signed, isolated high Pt leptons, Plus missing energy and m(ll)>15 GeV



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Why W leptonic decays?





CDF Analysis Strategy









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Adding new channels (2010)



And also extend to lower M(II)

- extend the opposite sign analysis
- M(ll) < 15 GeV



DO add the lepton+ jet(s) Channels already included In the CDF 2009 analysis

- backgrounds very different for τ samples
- Multijet and $Z \rightarrow \tau \tau$ • dominate
- Control samples allow • for cross checks of τ kinematic and ID variables
 - W+jets $(e\tau \& \mu\tau)$
 - Mulijet $(e\tau)$
 - $Z \rightarrow \tau \tau (\mu \tau)$



Explaining Limit Plot



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Theory & uncertainties

An(other) important topic & no time to discuss it

- We make use of well-motivated and state of the art gluon fusion cross-section calculations and uncertainties
 - gg → H uses NNLL + NNLO calculations
 - "Next to Next to Leading Log/Order"
 - de Florian & Grazzini (Phys.Lett.B674:291-294, 2009)
 - Soft-gluon resummation treatment
 - MSTW2008 Parton Density Function
 - Anastasiou, Boughezal, Petriello (JHEP:0904:003, 2009)
 - Proper treatment of b-quarks at NLO
 - Inclusion of two-loop electroweak effects
- For those interested in a detailed explanation of our choices and comparison with more extreme approaches :

http://tevnphwg.fnal.gov/results/SMHPubWinter2010/ gghtheoryreplies_may2010.html

Courtesy of Ben Kilmenster



Reweight PYTHIA Higgs kinematics to full NNLL calculation



SM HIGH MASS HIGGS combination



•combine all 18 high mass channels

•use bayesian approach

•weighted by sensitivity, average integrated luminosity

•5.90 fb-1 @ 165 GeV

CDF SM Higgs limit @165 GeV

 $\sigma/\sigma_{\rm SM}$ (obs) = 1.13

 σ/σ_{sm} (exp) = 1.00



•combine all 10 high mass channels		
•modified frequentist approach wit log likelihood ratio test statistic		
 weighted by sensitivity, average integrated luminosity 		
•6.10 fb ⁻¹ @ 165 GeV		
D0 SM Higgs limit @165 GeV		
σ/σ _{sm} (obs) = 1.03		
σ/σ_{SM} (exp) = 1.14		

SM HIGH MASS HIGGS combination







LLR, $\pm 1\sigma$ LLR, $\pm 2\sigma$

-- LLR, DØ Preliminary, <L> = 6.1 fb LLR LLR -4 broots and some transferrations worther we 100 110 120 130 140 150 160 170 180 190 200 m., (GeV) **SM Higgs Combination** Observed ---- Expected DØ Preliminary, <L>= 6.1 fb1 Expected ±10 Expected ±2σ Standard Model = 1,0 100 110 120 130 140 150 160 170 180 190 200 m_H (GeV)

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CDF/D0 Combinations



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Perspectives high/low mass Higgs by end Run II



Delivered luminosity now ~8.5 fb⁻¹ per experiment Tevatron will deliver ~11 fb⁻¹ per experiment by end of 2011



Finance issues

Scientific motivations

Politics Si TevatrOn Run II No? Manpower

Impact on the FNAL future

Detector aging

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Courtesy G. Punzi (PAC Augst 2010)

Wire aging in COT (main drift chamber)

- Since the recovery after adding O₂, the gain is steady.
- Example of the many things that are running well



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New amplifier board with more sensitive light receivers has been prepared as backup => no loss of ladders up to 20 fb-1 04/09/2010 TEVATRON, Corfu 2010, ASN 146 Dr. Steven Chu Secretary of Energy U.S. Department of Energy 1000 Independence Ave., SW Washington, DC 20585

Dear Dr. Chu,

Cc: Dr. Steven E. Koonin, DOE Under Secretary for Science Cc: Dr. William F. Brinkman, Director, DOE Office of Science Cc: Dr. Dennis Kovar, DOE Associate Director of Science for High Energy Physics Cc: Dr. Joseph L. Dehmer, NSF Associate Director, Physics Division Cc: Dr. Pier Oddone,

We are writing to alert you to an urgent and exceptional oppiector. Femilab concerning a concrege physics. Given the outstanding current performance of the Tevatron at Fermilab, coupled with the delays in the Large Hadron Collider at CERN, extended running of the Tevatron beyond the planned shutdown in 2011 offers an unanticipated opportunity for American science. The Tevatron has the potential to discover new phenomena – most notably the Higgs boson – more quickly than the LHC, and has unique features that complement LHC experiments. A premature shutdown would rob the U.S. program of valuable discovery opportunities.

Running the Tevatron through 2014 could produce the first direct evidence of a Higgs boson, a remarkable triumph for U.S. HEP and one of the most exciting scientific developments of this decade. The current plan to shut down the Tevatron at the end of FY11 was formulated some years ago, before the unanticipated delays at the LHC. The Tevatron now has the unique capacity to detect the expected low mass Higgs boson during the next few years. Even after LHC experiments find evidence for the Higgs boson, combining a bottom quark decay signal from the Tevatron with complementary decay signals from the LHC would help determine the properties of the Higgs particle, testing if it agrees with the simplest predictions, and furthermore help confirm the case for the Higgs particle interpretation.

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Dr. Steven Chu Secretary of Energy U.S. Department of Energy 1000 Independence Ave., SW Washington, DC 20585

Dear Dr. Chu,

We are writing to alert you to an physics. Given the outstanding c the delays in the Large Hadron (planned shutdown in 2011 offer: Tevatron has the potential to dis quickly than the LHC, and has u shutdown would rob the U.S. pri

Running the Tevatron through 2 remarkable triumph for U.S. HE decade. The current plan to shut years ago, before the unanticipat capacity to detect the expected le LHC experiments find evidence from the Tevatron with complen properties of the Higgs particle, help confirm the case for the Hig

Prof. Nima Arkani-Hamed Institute for Advanced Study Prof. Jonathan Bagger Johns Hopkins University Prof. Thomas Banks UC Santa Cruz and Rutgers University Prof. Prof. Vernon Barger University of Wisconsin Prof. Andrew Cohen Boston University Prof. Csaba Csaki Cornell University Dr. Sally Dawson Brookhaven National Laboratory Stanford University Prof. Lance J. Dixon SLAC and Stanford University Prof. Edward Farhi MIT. University of Michigan Prof. Mary K. Gaillard UC Berkelev Prof. Howard Georgi Harvard University Boston University Prof. David Gross Kavli Institute for Theoretical Physics UC Davis UC Santa Cruz Prof. Tao Han University of Wisconsin Prof. JoAnne Hewett SLAC and Stanford University

July 28, 2010

Prof. Robert L. Jaffe MIT Prof. Shamit Kachru SLAC and Stanford University Prof. Gordon L. Kane University of Michigan Prof. David B. Kaplan University of Washington Prof. Markus A. Lut. UC Davis Prof Ann Nelson University of Washington Prof. Michael Peskin SLAC and Stanford University Prof. Joseph Polchinski Kavli Institute for Theoretical Physics Prof. Pierre Ramond University of Florida Prof. Lisa Randall Harvard University Prof Martin Schmaltz Boston University Prof. Nathan Seiberg Institute for Advanced Study Prof. Eva M. Silverstein SLAC and Stanford University Prof. Raman Sundrum University of Maryland Prof. Cumrun Vafa Harvard University Prof. Carlos Wagner ANL and University of Chicago Prof. Steven Weinberg University of Texas Prof. Frank Wilczek MIT Prof. Edward Witten Institute for Advanced Study

If there is only one reason to pursue...



If there is only one reason to pursue...



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If there is only one reason to pursue...





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Tevatron vs LHC: complementarity & competition





September 1st 2010

Director's Corner



Last Friday we had a special meeting of the Physics Advisory Committee to address just one question: whether to recommend a proposed three-year extension of the Tevatron run. The Tevatron is now scheduled to close in September 2011. The <u>PAC's recommendation</u> was very strong. They recommended that we run the Tevatron in each of two proposed scenarios, one with additional funding in Fermilab's budget for Tevatron operations and one with no such extra funding. The committee's recommendation recognizes the extraordinary performance of the Tevatron and its remaining promise for the future. Extended running, combined with the predicted improvement in data analysis, would allow the Tevatron to make critical contributions to the discovery of the Higgs boson and our understanding of electroweak symmetry breaking. The combined potential of the Tevatron extension and early LHC running would constitute the optimal campaign to solve the central issue of particle physics. Many eminent members of the U.S. particle physics community have written to Secretary of Energy Steven Chu to express their support for more Tevatron running. The extension also motivates the collaborations to retain physicists and add new ones to carry out this program.

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Backup slides

AFB (TOP): backup

- Proton antiproton initial state is charge asymmetric, but strong interaction is not sensitive to the charge:
 - LO: kinematic distributions in production are charge symmetric
 - NLO 2→2, LO 2→3: expect (5—10)% asymmetry (higher order corrections are small)
 - NLO 2→3: reduce expected asymmetry significantly ⇒ strong dependence from phase space region
- SM asymmetry is small, so asymmetry is a sensitive variable to test the new physics contribution.
- A charge asymmetry can be observed as forward-backward asymmetry defined as



See e.g.:

- J. H. Kuhn and G. Rodrigo, Phys. Rev. D 59, 054017 (1999)
 - O. Antunano, J. H. Kuhn and G.Rodrigo, Phys. Rev. D 77, (2008)

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TOP A_{FB} at Tevatron: motivation

> In QCD Leading Order, top production is symmetric.

In NLO the soft Coulomb field of an incoming light quark repels the quark top to larger rapidity, while attracting the antitop quark to smaller rapidities, creating a positive asymmetry at large rapidity, as defined by the quark direction.

The MCFM NLO parton level MC (K. Ellis et al.) predicts: A_{lab} =0.038±0.006 and A_{ttbar} =0.058±0.009

The QCD terms responsible for the QCD asymmetry are proportional to the β of the top/antitop quarks in the center of mass, so the asymmetry increases with the rapidity separation of the two quarks: NLO Asymmetry Prediction from MCFM



A number of models BSM predicts much higher asymmetries ex: Z', extra dimension Etc..

Importance of A_{FB} at Tevatron: dominated by qqbar process 04/09/2010 TEVATRON, Corfu 2010, ASN

TOP Forward-Backward Asymmetry

Compare the number of top and antitop produced with momentum in a given direction:



V.B.The new results from CDF, with 5.3fb⁻¹ data distinguishes between top & anti-t Aft 04/09/2010 TEVATRON, Corfu 2010, ASN 158

