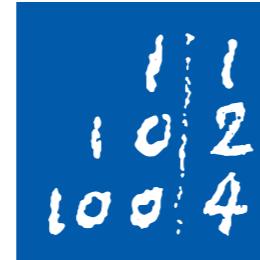


# On Type IIA Cosmology From Geometric Fluxes

Marco Zagermann  
(Leibniz University Hannover)



Based on : 0812.3551 (Caviezel, Koerber, Körs, Lüst, Wrase, M.Z.)  
0806.3458 (Caviezel, Koerber, Körs, Lüst, Tsimpis, M.Z.)  
(0812.3886 (Flauger, Paban, Robbins, Wrase))  
+ Work in progress

An important problem in string phenomenology:

Moduli stabilization

$$\Rightarrow V(\varphi^i)$$

An important problem in string phenomenology:

Moduli stabilization

$$\Rightarrow V(\varphi^i)$$

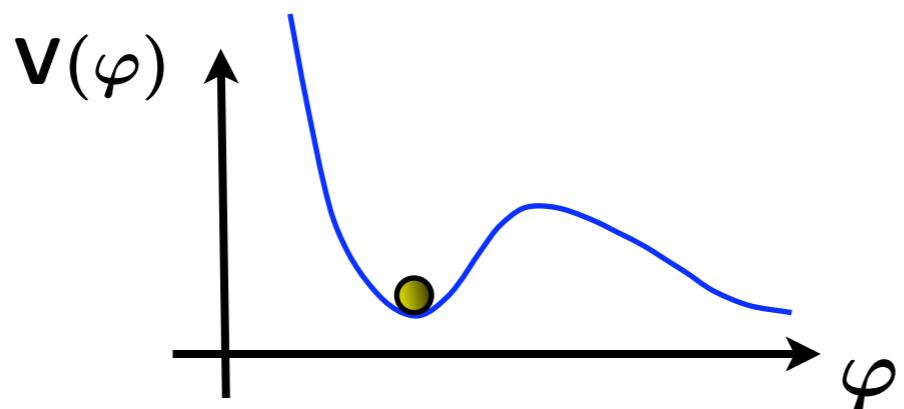
Particularly interesting:

$$V(\varphi^i) > 0$$

# Positive potential energy for:

(i) de Sitter vacua

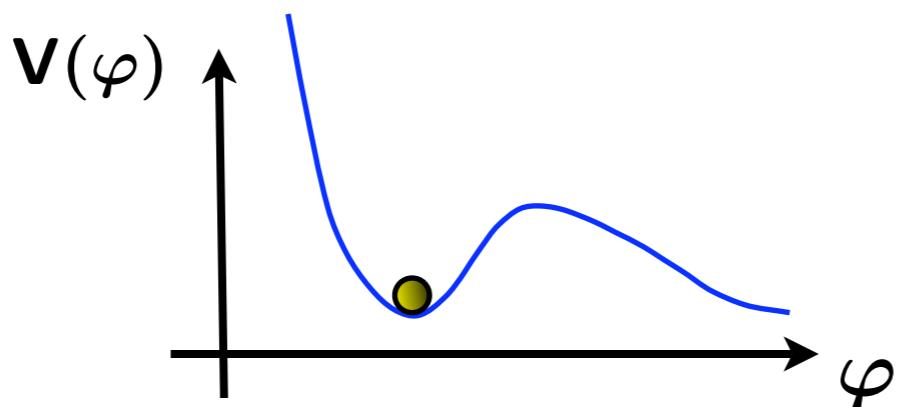
$(\Lambda > 0 \Rightarrow$  Today's accelerated expansion)



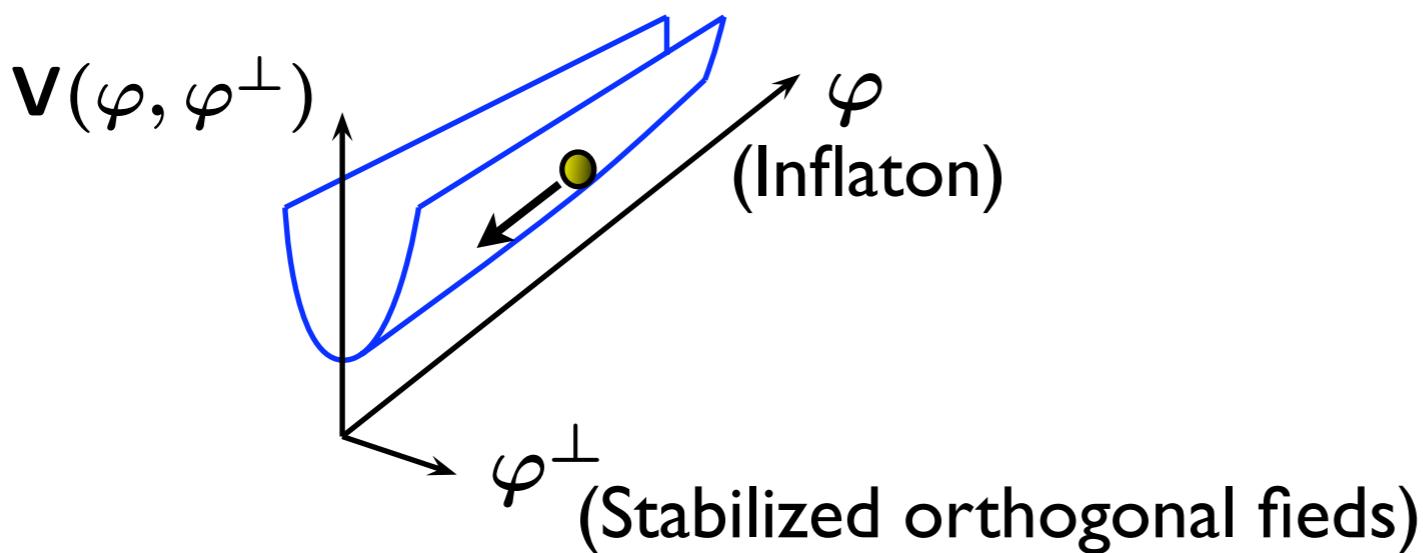
# Positive potential energy for:

(i) de Sitter vacua

$(\Lambda > 0 \Rightarrow$  Today's accelerated expansion)



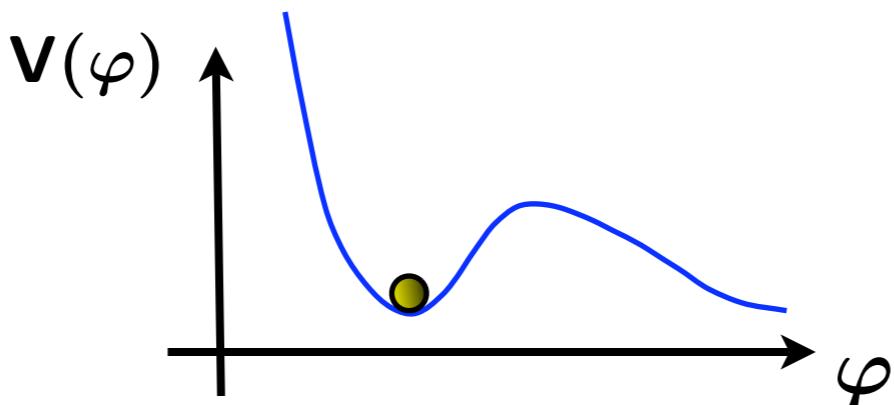
(ii) Slow-roll inflation



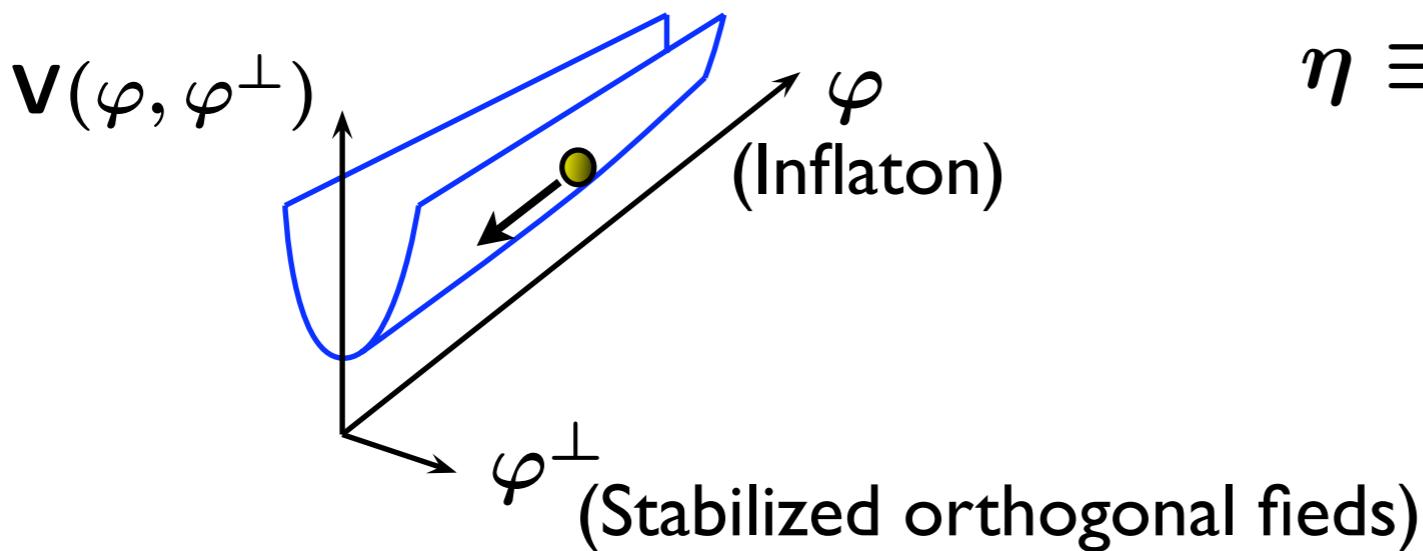
# Positive potential energy for:

(i) de Sitter vacua

$(\Lambda > 0 \Rightarrow$  Today's accelerated expansion)



(ii) Slow-roll inflation



$$\epsilon \equiv \frac{1}{2} g^{ij} \frac{(\partial_{\varphi^i} V)(\partial_{\varphi^j} V)}{V^2}$$

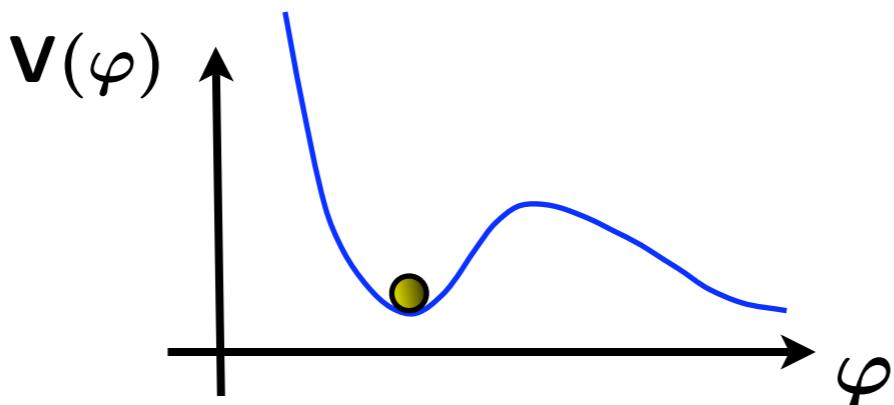
$$\eta \equiv \text{Min. eig.val.} \left( \frac{\nabla^i \partial_j V}{V} \right)$$

$$\epsilon, |\eta| \ll 1$$

# Positive potential energy for:

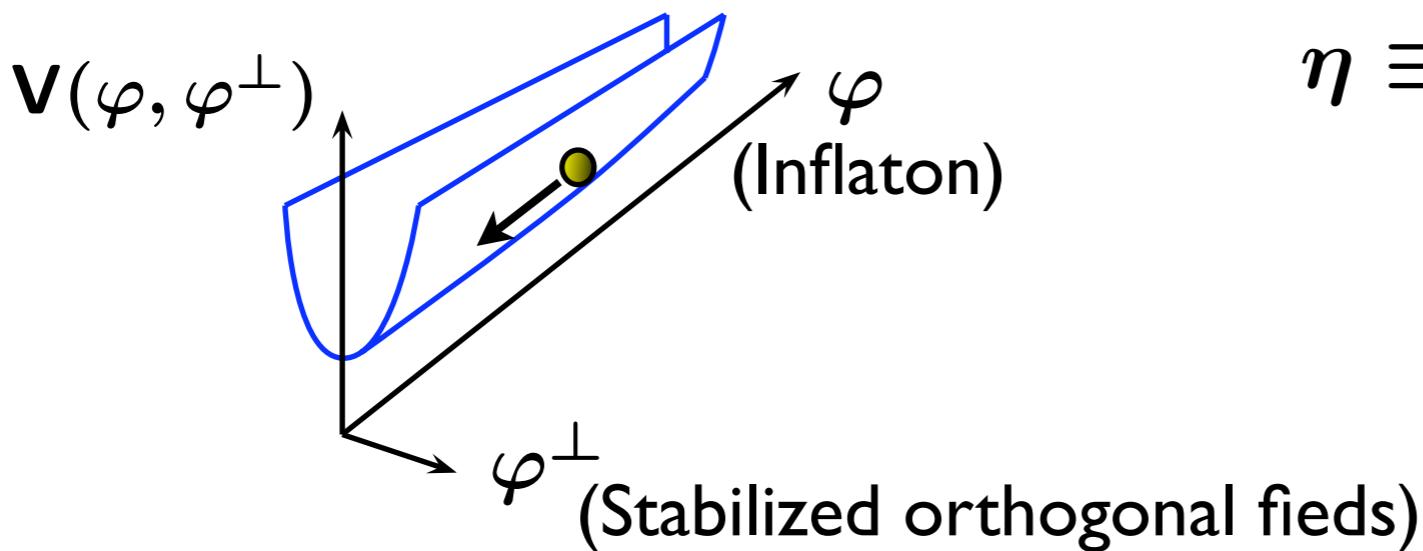
(i) de Sitter vacua

$(\Lambda > 0 \Rightarrow$  Today's accelerated expansion)



$$\epsilon = 0, \quad \eta > 0$$

(ii) Slow-roll inflation



$$\epsilon \equiv \frac{1}{2} g^{ij} \frac{(\partial_{\varphi^i} V)(\partial_{\varphi^j} V)}{V^2}$$

$$\eta \equiv \text{Min. eig.val.} \left( \frac{\nabla^i \partial_j V}{V} \right)$$

$$\epsilon, |\eta| \ll 1$$

A general problem:

Typical scalar potentials receive many contributions and corrections

## A general problem:

Typical scalar potentials receive many contributions and corrections

Often:

Subtle interplay of

classical

and

quantum effects



Easy



Hard to compute precisely

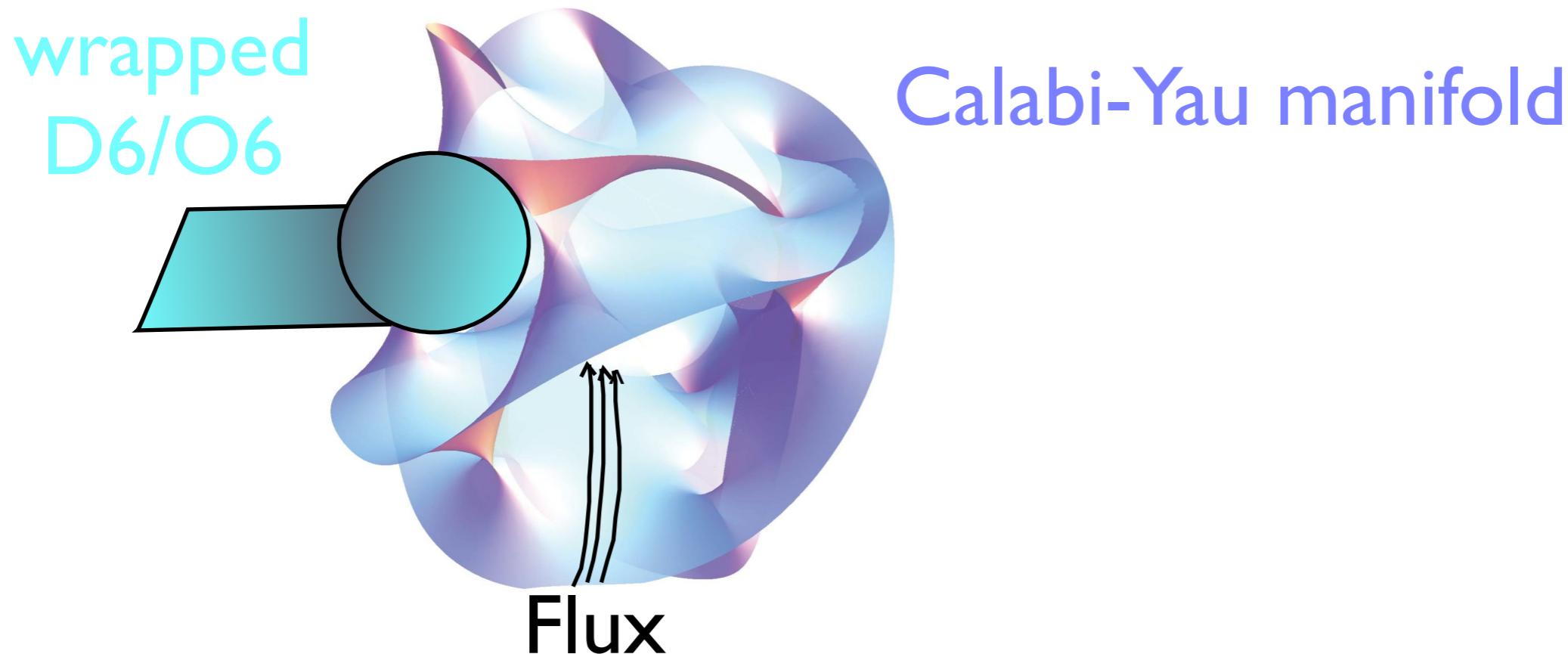
E.g. KKLT; LARGE Volume;  
KKLMMT...

Cf. also Ramos-Sánchez' talk

A nice laboratory:

## Type IIA on Calabi-Yau spaces with

- Magnetic fluxes of p-form field strengths
- D6-branes/O6-planes



## Observation:

All geometric moduli can be stabilized at tree-level

Grimm, Louis (2004); Kachru, Kashani-Poor (2004)

## Observation:

All geometric moduli can be stabilized at tree-level

Grimm, Louis (2004); Kachru, Kashani-Poor (2004)

In special cases:

- All moduli stabilized
- Parameterically controlled classical regime  
⇒ Quantum corrections small

Derendinger, Kounnas, Petropoulos, Zwirner (2004, 2005)

Villadoro, Zwirner (2005)

de Wolfe, Giryavets, Kachru, Taylor (2005)

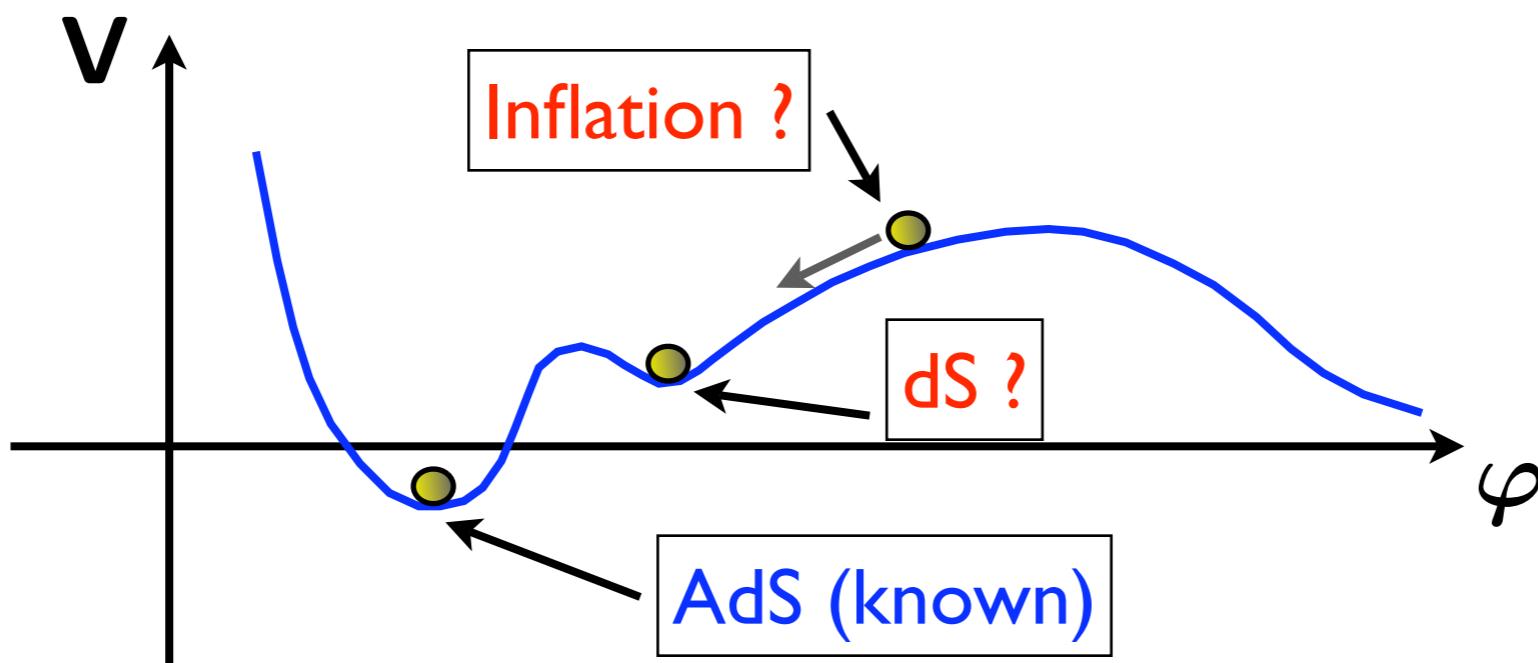
Unfortunately...

All these stabilized vacua have

$$\Lambda < 0 \ (\Rightarrow \text{AdS})$$

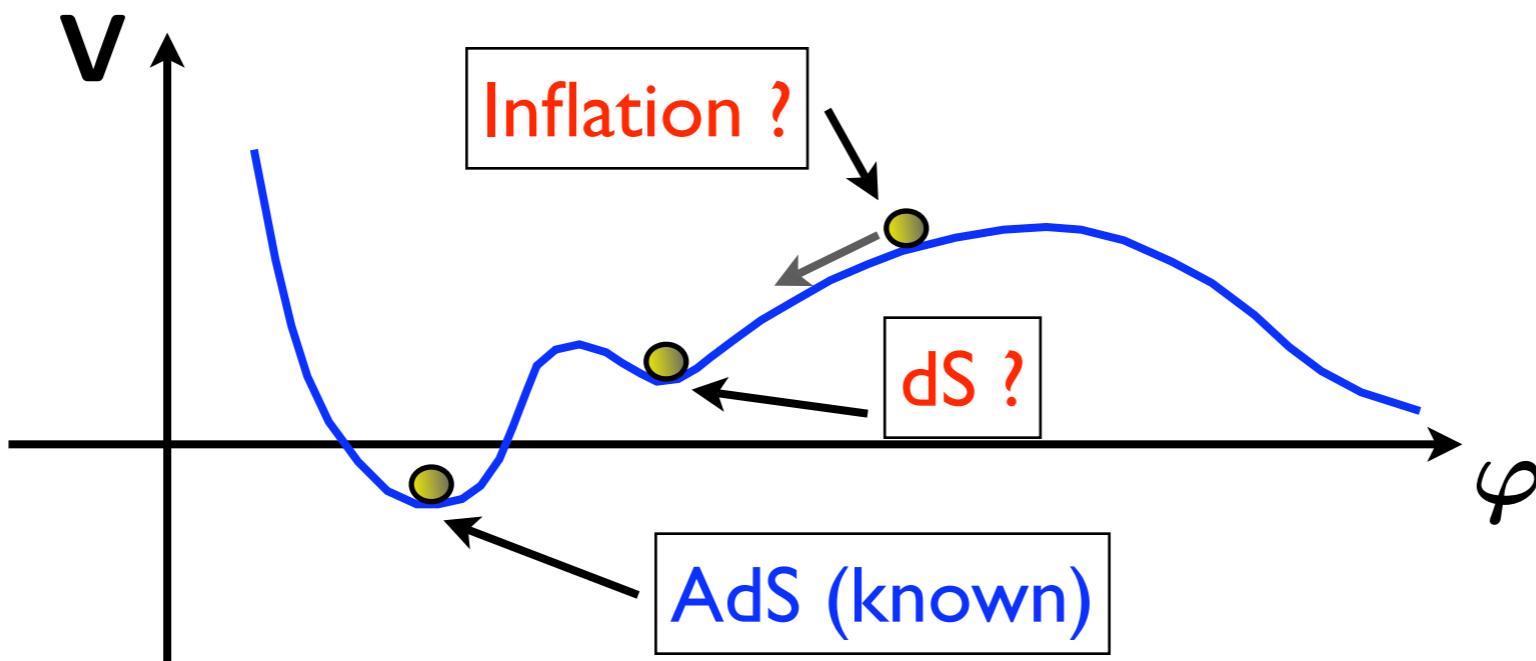
Two possibilities:

(i) Search for dS/inflation **away** from AdS vacuum

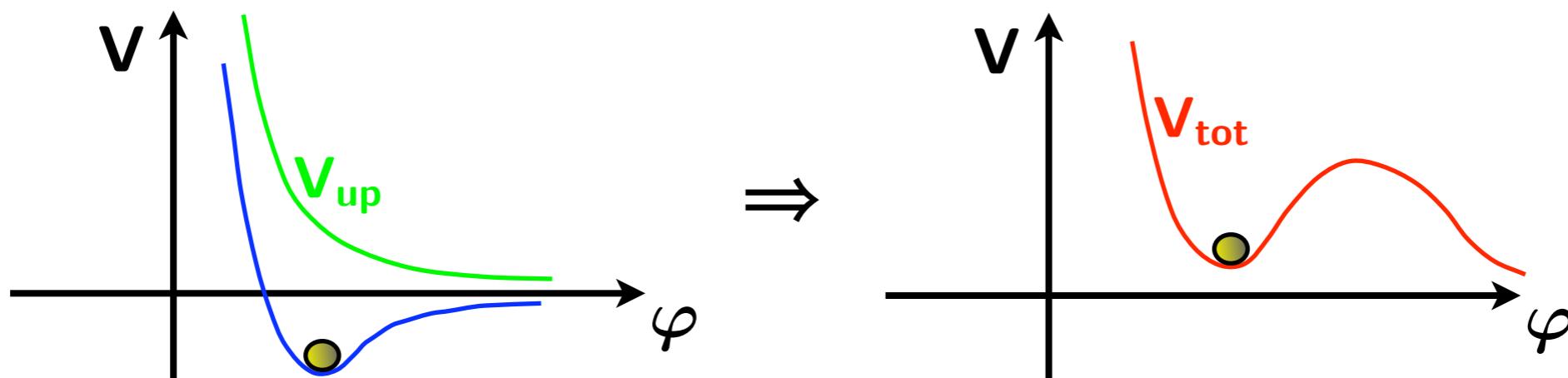


Two possibilities:

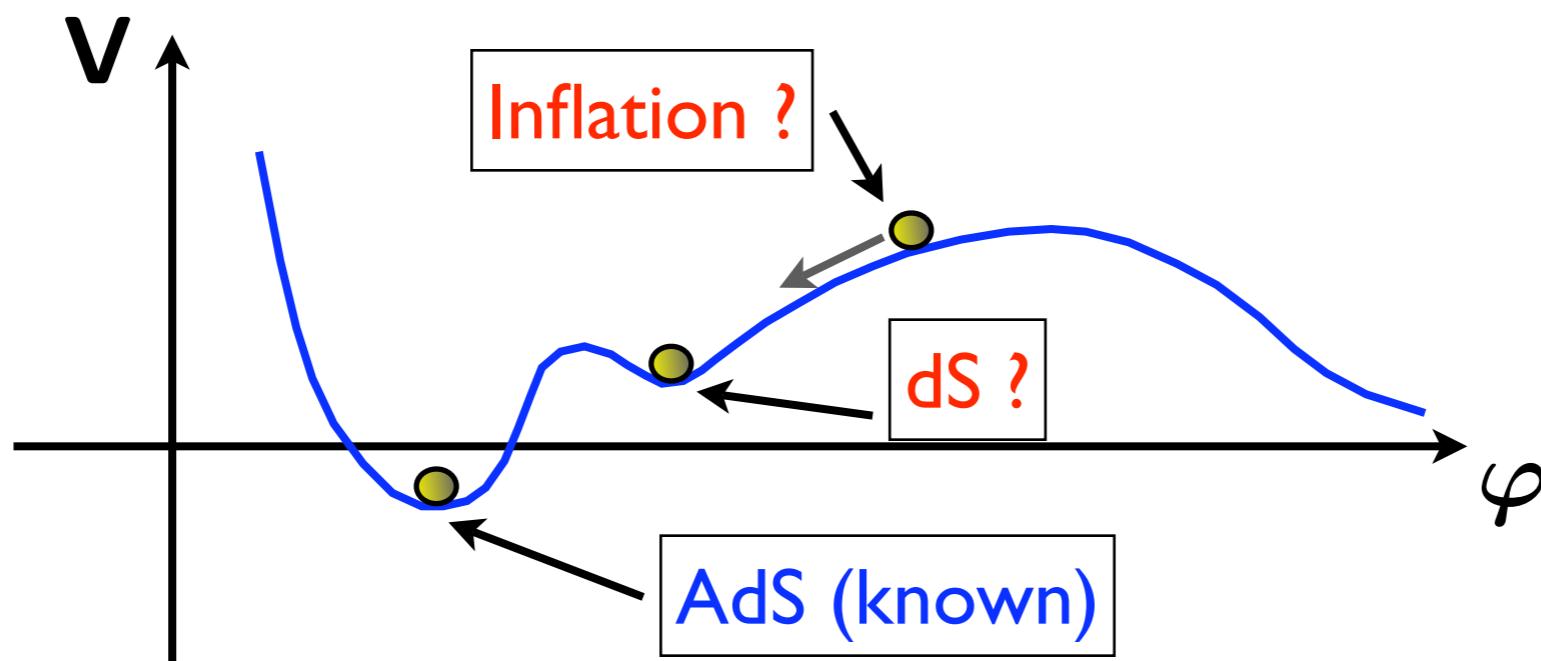
(i) Search for dS/inflation **away from AdS vacuum**



(ii) Add additional ingredients  $\Rightarrow$  “Uplift potentials”



# (i) dS or inflation away from AdS vacuum?



(i) dS or inflation away from AdS vacuum?

No-go theorem:

Classical IIA compactifications with

- $\mathcal{M}^{(6)} = \text{Calabi-Yau}$  ( $\rightarrow$  Ricci-flatness)
- O6/D6 sources
- p-form fluxes (incl. Romans' mass)

$\Rightarrow$  No de Sitter vacua and no slow-roll inflation !

Hertzberg, Kachru, Taylor, Tegmark (2007)

**Note:** Due to the O6-planes ( $\rightarrow$  negative tension),  
this goes beyond no-go theorems by

Gibbons (1985)

de Wit, Smit, Hari Dass (1987)

Maldacena-Nuñez (2000)

Cf. also Wesley, Steinhardt (2008)

Townsend, Wohlfarth (2003)

## Sketch of proof:

Consider scaling of potential w.r.t.

$$\rho \equiv (\text{Vol})^{1/3}$$

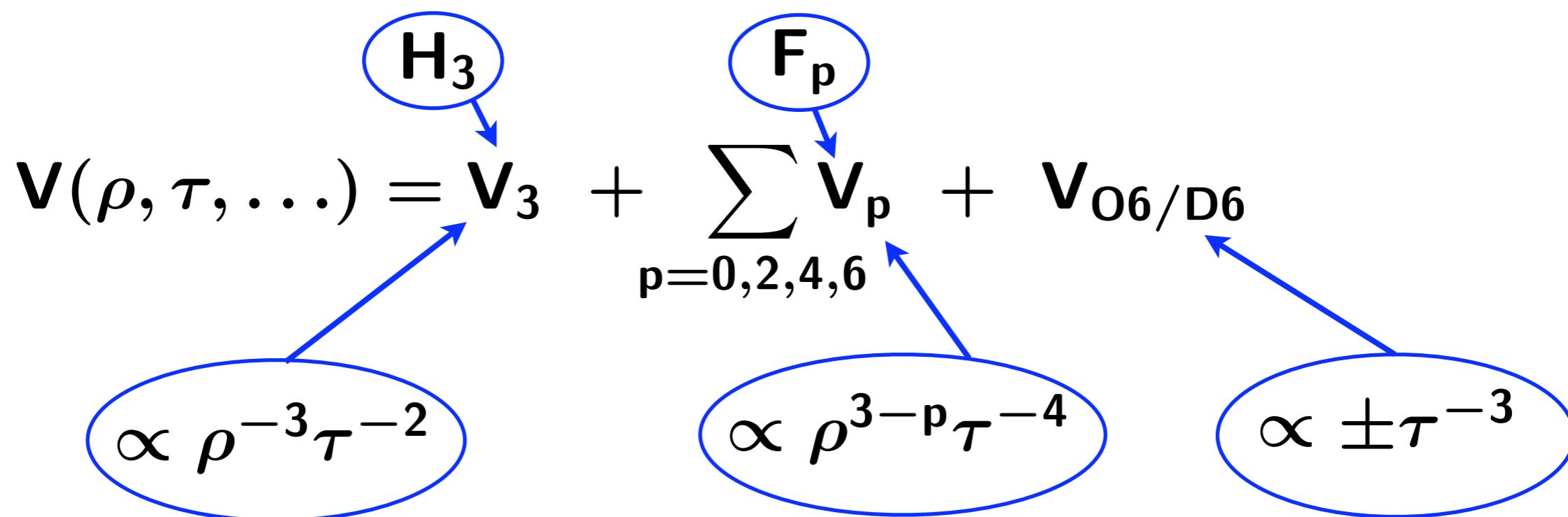
$$\tau \equiv e^{-\phi} \sqrt{\text{Vol}} \quad (\phi = 10\text{D Dilaton})$$

## Sketch of proof:

Consider scaling of potential w.r.t.

$$\rho \equiv (\text{Vol})^{1/3}$$

$$\tau \equiv e^{-\phi} \sqrt{\text{Vol}} \quad (\phi = 10\text{D Dilaton})$$



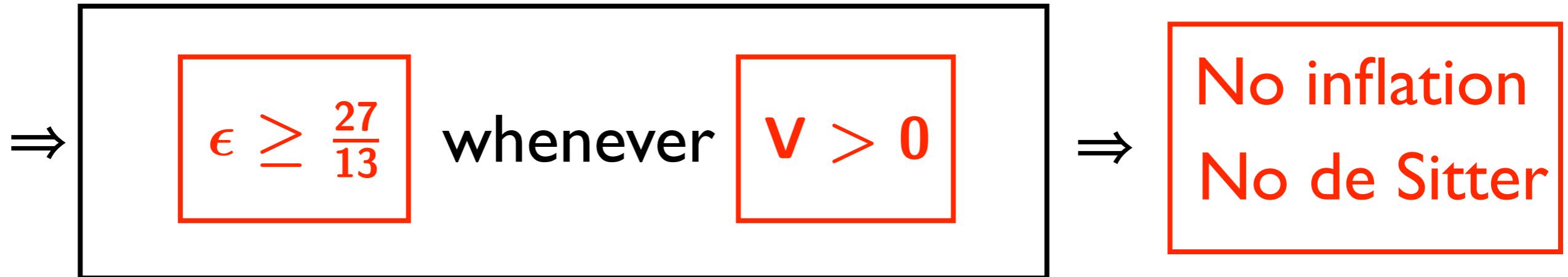
$$\Rightarrow \mathbf{DV} \equiv (-\rho\partial_\rho - 3\tau\partial_\tau)\mathbf{V} \geq 9\mathbf{V}$$

$$\mathsf{DV} \equiv (-\rho\partial_\rho - 3\tau\partial_\tau)\mathsf{V} \geq 9\mathsf{V}$$

$$\epsilon = \mathsf{V}^{-2}\left[\tfrac{(\mathsf{DV})^2}{39} + \text{(positive)}\right]$$

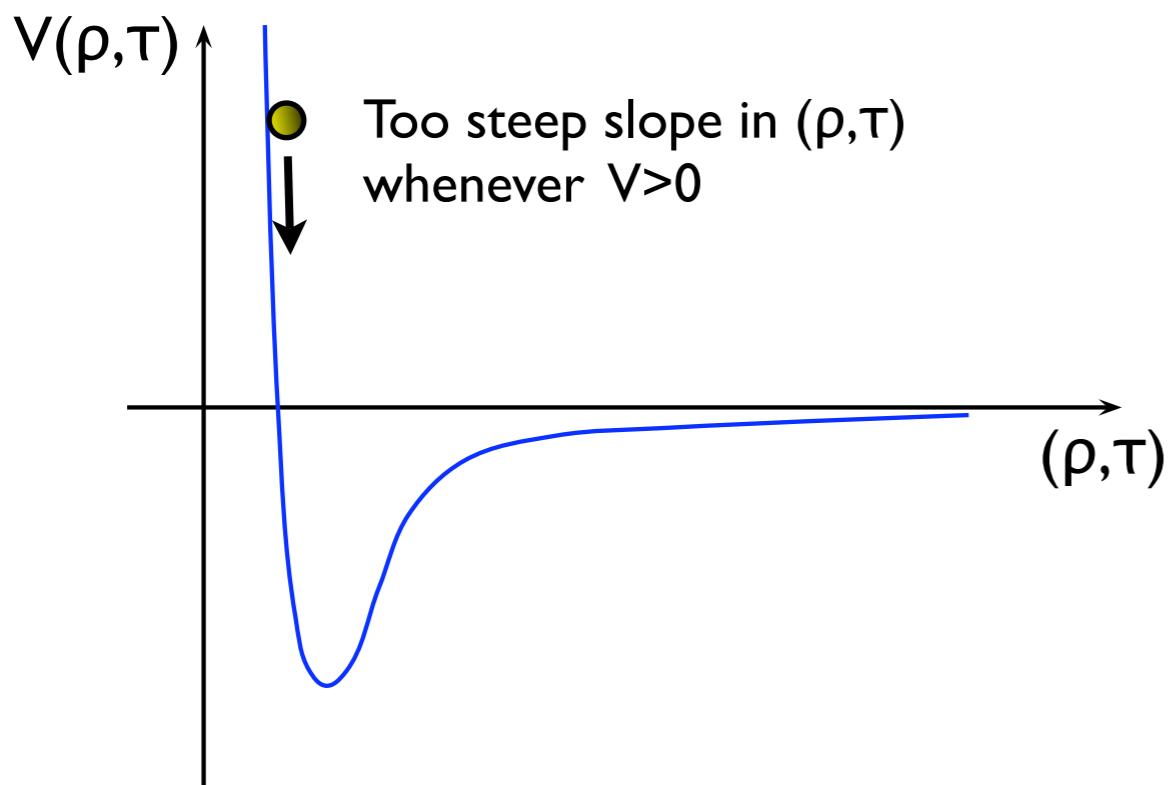
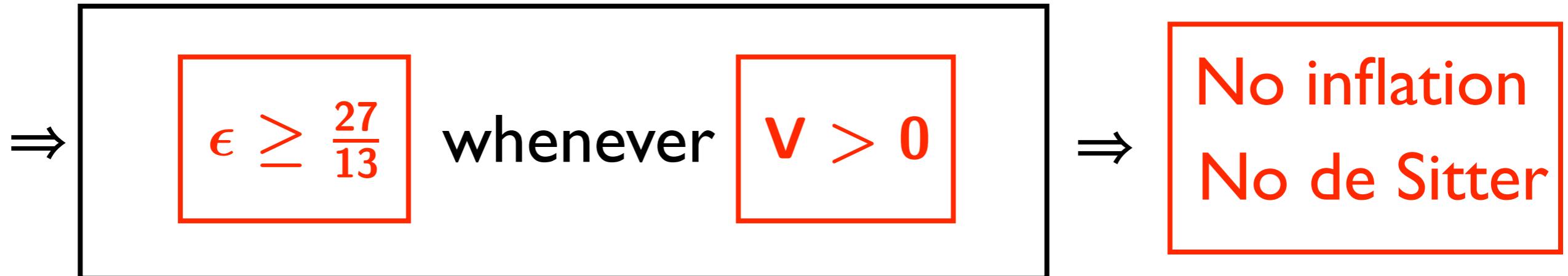
$$DV \equiv (-\rho\partial_\rho - 3\tau\partial_\tau)V \geq 9V$$

$$\epsilon = V^{-2} \left[ \frac{(DV)^2}{39} + (\text{positive}) \right]$$



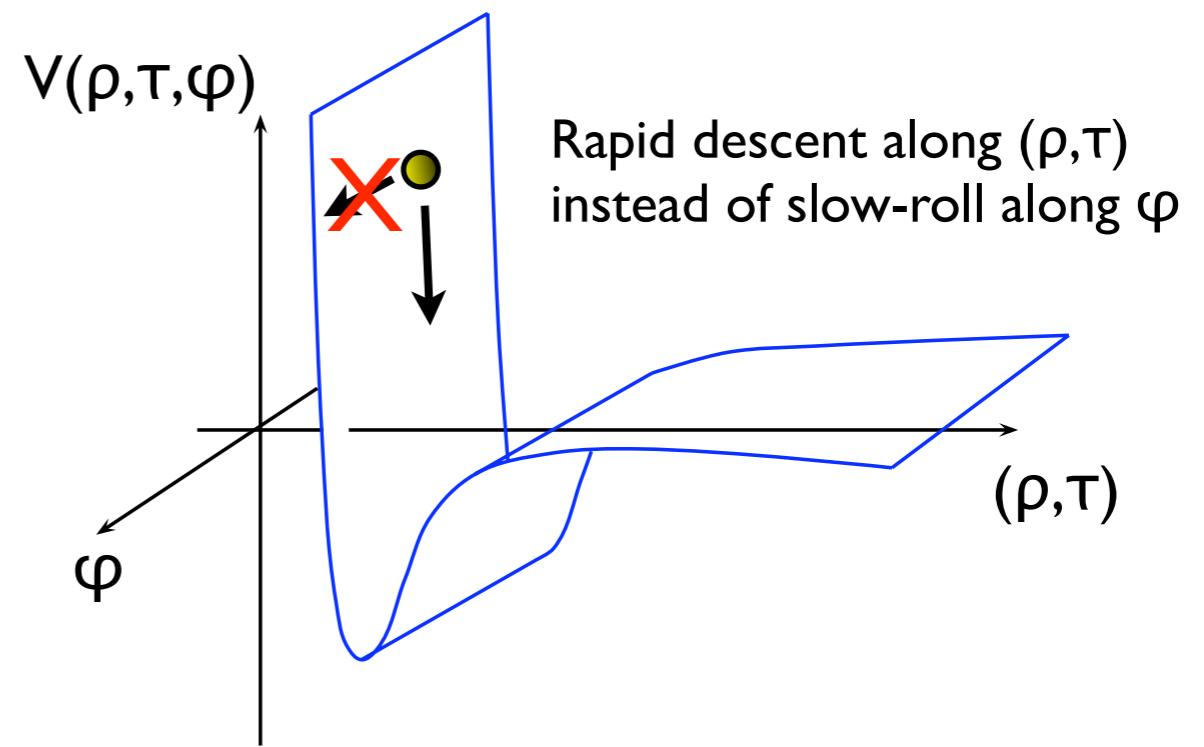
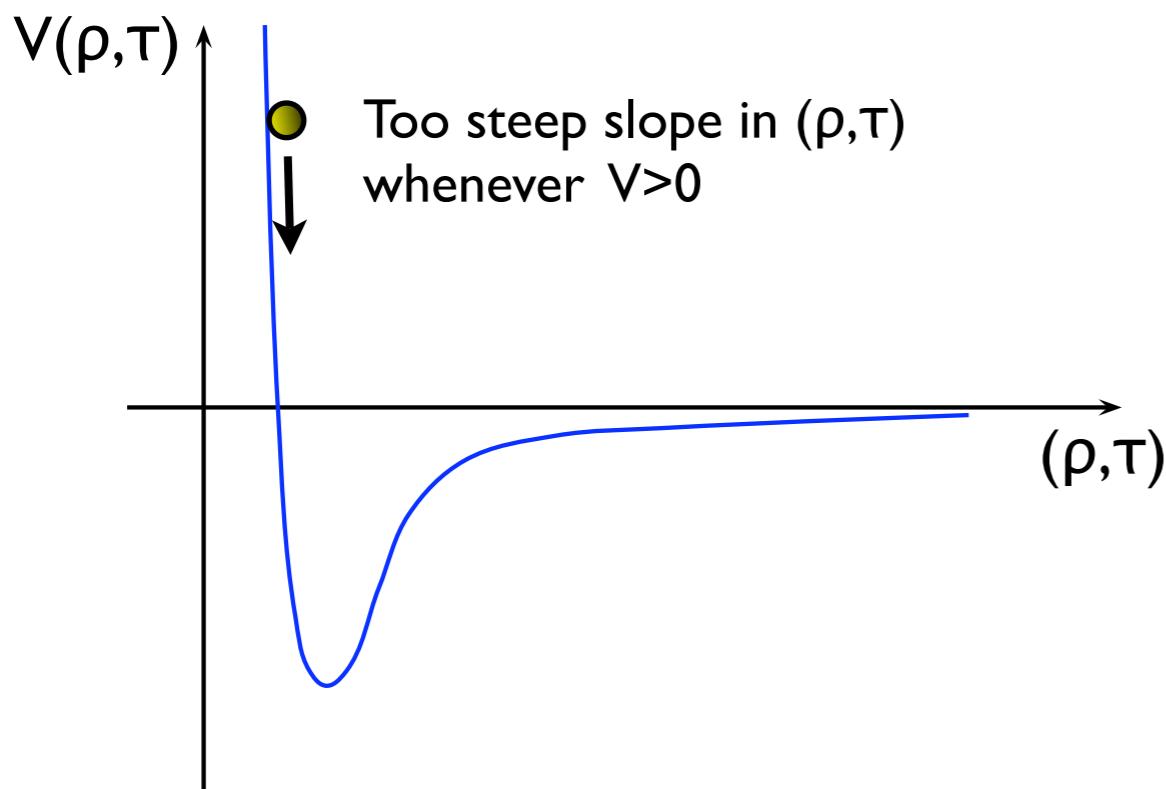
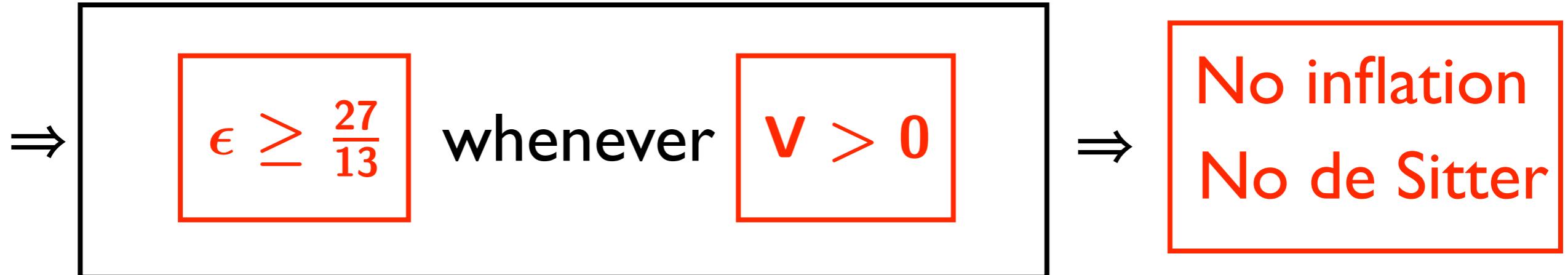
$$DV \equiv (-\rho \partial_\rho - 3\tau \partial_\tau)V \geq 9V$$

$$\epsilon = V^{-2} \left[ \frac{(DV)^2}{39} + (\text{positive}) \right]$$



$$DV \equiv (-\rho \partial_\rho - 3\tau \partial_\tau)V \geq 9V$$

$$\epsilon = V^{-2} \left[ \frac{(DV)^2}{39} + (\text{positive}) \right]$$



## Possible caveats:

- Quantum corrections      E.g. Saueressig, Theis, Vandoren (2005);  
Palti, Tasinato, Ward (2008)
- Additional classical ingredients
  - “Geometric fluxes” (“Torsion”) (= geometric twisting away from CY  $\Rightarrow R_{mn} \neq 0$ )
  - O4-planes
  - D8-branes
  - NS5-branes
  - KK5-monopoles
  - “Nongeometric fluxes”
    - 
    - 
    -

Silverstein (2007):

For classical de Sitter vacua add, e.g.:

- Geometric fluxes  
(Particular twisted torus)
- KK5-monopoles
- Fractional Chern-Simons invariants

Some issues to keep in mind:

- High SUSY breaking scale
- No large mass gap to KK modes
- Tadpole cancellation?
- Backreaction under control?

**What is the minimal controllable setup?**

**What is the minimal controllable setup?**

**Best understood extra ingredient:**

**Geometric fluxes**



**Are they sufficient ?**

## Three complementary works:

(i) Haque, Shiu, Underwood, Van Riet (2008)

$$\mathcal{M}^{(6)} = (\text{Nil}_3 \times \text{Nil}'_3)/\mathcal{O}$$

(ii) Caviezel, Koerber, Körs, Lüst, Wräse, M.Z. (2008)

$\mathcal{M}^{(6)}$  = Cosets with SU(3)-structure

(iii) Flauger, Paban, Robbins, Wräse (2008)

$\mathcal{M}^{(6)}$  = More general twisted tori

## Three complementary works:

(i) Haque, Shiu, Underwood, Van Riet (2008)

$$\mathcal{M}^{(6)} = (\text{Nil}_3 \times \text{Nil}'_3)/\mathcal{O}$$

(ii) Caviezel, Koerber, Körs, Lüst, Wräse, M.Z. (2008)

$$\mathcal{M}^{(6)} = \text{Cosets with SU}(3)\text{-structure}$$

This  
talk

(iii) Flauger, Paban, Robbins, Wräse (2008)

$$\mathcal{M}^{(6)} = \text{More general twisted tori}$$

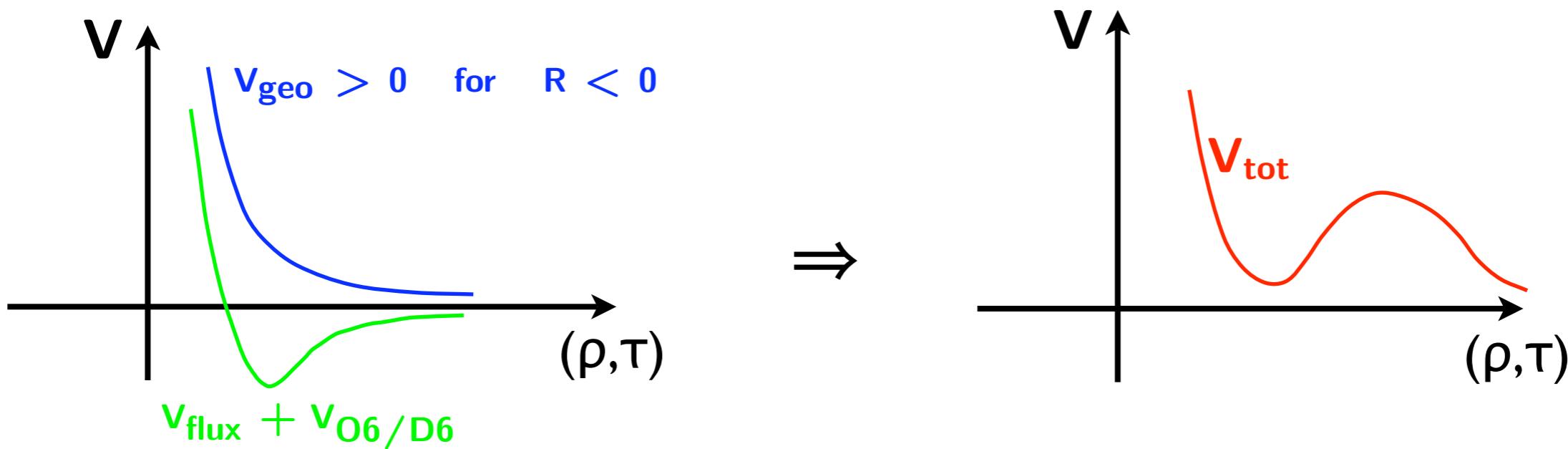
# Geometric fluxes

$$V_{\text{geo}} \propto -R \propto \rho^{-1} \tau^{-2}$$

## Geometric fluxes

$$V_{\text{geo}} \propto -R \propto \rho^{-1} \tau^{-2}$$

⇒ Effective uplift term for  $R < 0$

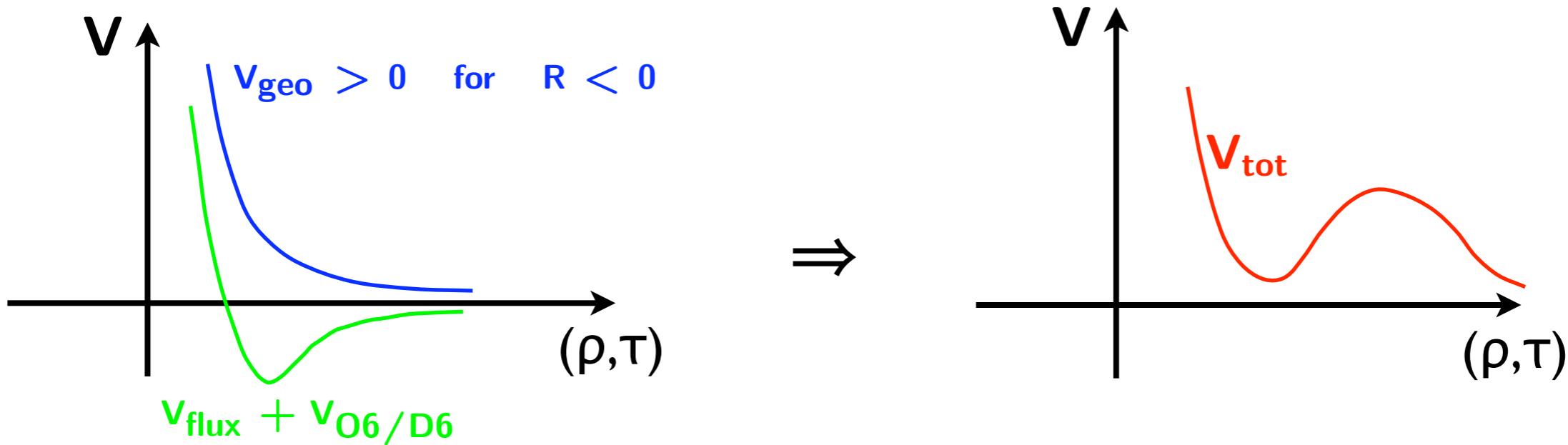


⇒ May help to remove steep fall-off in  $(\rho, \tau)$

## Geometric fluxes

$$V_{\text{geo}} \propto -R \propto \rho^{-1} \tau^{-2}$$

⇒ Effective uplift term for  $R < 0$



⇒ May help to remove steep fall-off in  $(\rho, \tau)$

But: Are all orthogonal field directions also ok?

⇒ Need a setup in which

$$\mathbf{V} = \mathbf{V}(\rho, \tau, \varphi^\perp)$$

is well understood

## A well-controlled setup:

Caviezel, Koerber, Körs, Lüst, Tsimpis, M.Z. (2008)

cf. also House,  
Palti (2005)

$\mathcal{M}^{(6)}$  = Coset space  $G/H$  with ( $G$ -invariant)  
 $SU(3)$ -structure (+ orientifolding)

## A well-controlled setup:

cf. also House,  
Caviezel, Koerber, Körs, Lüst, Tsimpis, M.Z. (2008) Palti (2005)

$\mathcal{M}^{(6)}$  = Coset space  $G/H$  with (G-invariant)  
**SU(3)-structure** (+ orientifolding)

### SU(3)-structure:

- ⇒  $\mathcal{M}^{(6)}$  admits globally well-defined spinor  $\eta$
- ⇒ 4D, N=1 supergravity action
- ⇒ For  $\nabla\eta \neq 0$ : No CY ⇒  $R_{mn} \neq 0 \Rightarrow V_{geo} \neq 0$

Cf. Tsimpis' talk

## A well-controlled setup:

cf. also House,  
Caviezel, Koerber, Körs, Lüst, Tsimpis, M.Z. (2008) Palti (2005)

$\mathcal{M}^{(6)}$  = Coset space  $G/H$  with ( $G$ -invariant)  
 $SU(3)$ -structure (+ orientifolding)

### $SU(3)$ -structure:

- ⇒  $\mathcal{M}^{(6)}$  admits globally well-defined spinor  $\eta$
- ⇒ 4D,  $N=1$  supergravity action
- ⇒ For  $\nabla\eta \neq 0$ : No CY ⇒  $R_{mn} \neq 0 \Rightarrow V_{geo} \neq 0$

Early work: Gurrieri, Louis, Micu, Waldram;  
Dall'Agata, Prezas; Lüst, Tsimpis; Behrndt, Cvetič;  
Gauntlett, Martelli, Waldram;  
Graña, Minasian, Petrini, Tomasiello;...

## A well-controlled setup:

cf. also House,  
Caviezel, Koerber, Körs, Lüst, Tsimpis, M.Z. (2008) Palti (2005)

$\mathcal{M}^{(6)}$  = Coset space  $G/H$  with ( $G$ -invariant)  
 $SU(3)$ -structure (+ orientifolding)

### $SU(3)$ -structure:

- ⇒  $\mathcal{M}^{(6)}$  admits globally well-defined spinor  $\eta$
- ⇒ 4D,  $N=1$  supergravity action
- ⇒ For  $\nabla\eta \neq 0$ : No CY ⇒  $R_{mn} \neq 0 \Rightarrow V_{geo} \neq 0$

### Problem:

$$J_{mn} \equiv i\eta_+^\dagger \gamma_{mn} \eta_+$$

$$\Omega_{mnp} \equiv \eta_-^\dagger \gamma_{mnp} \eta_+$$

$\nabla\eta \neq 0 \Rightarrow dJ, d\Omega \neq 0 \Rightarrow$  Expansion basis, moduli?

## A well-controlled setup:

cf. also House,  
Caviezel, Koerber, Körs, Lüst, Tsimpis, M.Z. (2008) Palti (2005)

$\mathcal{M}^{(6)}$  = Coset space  $G/H$  with ( $G$ -invariant)  
 $SU(3)$ -structure (+ orientifolding)

### $SU(3)$ -structure:

- ⇒  $\mathcal{M}^{(6)}$  admits globally well-defined spinor  $\eta$
- ⇒ 4D,  $N=1$  supergravity action
- ⇒ For  $\nabla\eta \neq 0$ : No CY ⇒  $R_{mn} \neq 0 \Rightarrow V_{geo} \neq 0$

### Coset space structure:

- ⇒ Natural expansion basis:  $G$ -invariant forms
- ⇒ Explicit 4D action (consistent truncation)

Cassani, Kashani-Poor (2009)

# Restriction to semisimple and Abelian group factors

⇒ **7 Models:**

Koerber, Lüst, Tsimpis (2008)

$$\frac{\mathbf{G}_2}{\mathbf{SU}(3)}, \quad \frac{\mathbf{Sp}(2)}{\mathbf{S}(\mathbf{U}(2) \times \mathbf{U}(1))}, \quad \frac{\mathbf{SU}(3)}{\mathbf{U}(1) \times \mathbf{U}(1)}, \quad \frac{\mathbf{SU}(3) \times \mathbf{U}(1)}{\mathbf{SU}(2)}, \quad \frac{\mathbf{SU}(2)^2}{\mathbf{U}(1)} \times \mathbf{U}(1)$$
$$\mathbf{SU}(2) \times \mathbf{U}(1)^3, \quad \mathbf{SU}(2) \times \mathbf{SU}(2)$$

# Restriction to semisimple and Abelian group factors

⇒ **7 Models:**

Koerber, Lüst, Tsimpis (2008)

$$\frac{G_2}{SU(3)}, \quad \frac{Sp(2)}{S(U(2) \times U(1))}, \quad \frac{SU(3)}{U(1) \times U(1)}, \quad \frac{SU(3) \times U(1)}{SU(2)}, \quad \frac{SU(2)^2}{U(1)} \times U(1)$$

$$SU(2) \times U(1)^3, \quad SU(2) \times SU(2)$$

R < 0 possible

⇒ Evade old no-go!

⇒ de Sitter or inflation?

Or are there new no-go's?

## A refined no-go theorem

Cf. Flauger, Paban,  
Robbins, Wrase (2008)

Old no-go:  $V_{\text{geo}} = 0, \quad V = V(\tau, \rho, \dots)$

$$\Rightarrow \varepsilon \geq 27/13$$

# A refined no-go theorem

Cf. Flauger, Paban,  
Robbins, Wrase (2008)

Old no-go:  $V_{\text{geo}} = 0, \quad V = V(\tau, \rho, \dots)$

$$\Rightarrow \varepsilon \geq 27/13$$

Refined no-go:  $V_{\text{geo}} \neq 0, \quad V = V(\tau, \sigma, \dots)$

Violates old no-go

Different Kähler modulus

# A refined no-go theorem

Cf. Flauger, Paban,  
Robbins, Wrase (2008)

Old no-go:

$$V_{\text{geo}} = 0, \quad V = V(\tau, \rho, \dots)$$

$$\Rightarrow \varepsilon \geq 27/13$$

Refined no-go:

$$V_{\text{geo}} \neq 0, \quad V = V(\tau, \sigma, \dots)$$

Violates old no-go

Different Kähler modulus

If:

$$(i) \quad \kappa_{ijk} = \kappa_{0ab} \quad \Rightarrow$$

$$\sigma \equiv \sqrt{\frac{\rho^3}{k^0}}$$

# A refined no-go theorem

Cf. Flauger, Paban,  
Robbins, Wrase (2008)

Old no-go:

$$V_{\text{geo}} = 0, \quad V = V(\tau, \rho, \dots)$$

$$\Rightarrow \epsilon \geq 27/13$$

Refined no-go:

$$V_{\text{geo}} \neq 0, \quad V = V(\tau, \sigma, \dots)$$

Violates old no-go

Different Kähler modulus

If:

(i)  $\kappa_{ijk} = \kappa_{0ab} \Rightarrow \boxed{\sigma \equiv \sqrt{\frac{\rho^3}{k^0}}}$

(ii)  $-\sigma \partial_\sigma [2\tau^2 \rho^3 V_{\text{geo}}] \geq 0$

$$\Rightarrow \boxed{\epsilon \geq 2} \quad \text{for} \quad \boxed{V > 0}$$

## Our cosets:

(i)  $\kappa_{ijk} = \kappa_{0ab}$  is always satisfied

⇒ Can define  $\sigma$

## Our cosets:

(i)  $\kappa_{ijk} = \kappa_{0ab}$  is always satisfied

⇒ Can define  $\sigma$

(ii)  $-\sigma \partial_\sigma [2\tau^2 \rho^3 V_{\text{geo}}] \geq 0$  is always satisfied except for

$$\mathcal{M}^{(6)} = \text{SU}(2) \times \text{SU}(2)$$

## Our cosets:

(i)  $\kappa_{ijk} = \kappa_{0ab}$  is always satisfied

$\Rightarrow$  Can define  $\sigma$

(ii)  $-\sigma \partial_\sigma [2\tau^2 \rho^3 V_{\text{geo}}] \geq 0$  is always satisfied **except for**

$$\mathcal{M}^{(6)} = \text{SU}(2) \times \text{SU}(2)$$

Only remaining candidate.  
(The others have  $\varepsilon \geq 2$ )

## SU(2)×SU(2)

Numerically:  $\epsilon \approx 0$  with  $V > 0$

But:

$$\eta \leq -2.4$$

(Large tachyonic direction)

Interestingly:

- Tachyon is combination of all moduli
- Not the tachyon of (?)

Covi, Gomez-Reino, Gross, Louis, Palma, Scrucca (2008)

Also:

- NS5, D4, D8 can **not** be added in these models
- No F-term uplifting à la “O’KKLT” possible  
Cf. Kallosh, Linde (2006), Kallosh, Soroush (2006)
- KK5-Monopole ? ⇒ Drastic modification of geometry  
See also Villadoro, Zvirner (2007)  
⇒ So far nothing really worked...

## Summary

Type IIA on CY + p-form fluxes + D6/O6:

- Tree-level moduli stabilization in AdS
- No-go against dS and inflation (HKKT)

→  $V > 0$  is too steep in  $(\rho, \tau)$

## Summary

Type IIA on CY + p-form fluxes + D6/O6:

- Tree-level moduli stabilization in AdS
- No-go against dS and inflation (HKKT)

→  $V > 0$  is too steep in  $(\rho, \tau)$

⇒ Quantum effects or/and additional classical ingredients

- Best understood: Geometric fluxes (deviation from CY)

Studied cosets with SU(3)-structure

⇒ Refined no-go in  $(\sigma, \tau)$ :  $\varepsilon \geq 2$  except for  $SU(2) \times SU(2)$

$SU(2) \times SU(2)$ :  $\varepsilon \approx 0$ , but  $\eta \leq -2.4$

Consistent with other works:

Haque, Shiu, Underwood, Van Riet (2008)

Flauger, Paban, Robbins, Wrase (2008)

⇒ Geometric fluxes may help in  $(\rho, \tau)$ -plane, but certainly do not automatically take care of all moduli ⇒ Many dangerous directions

Consistent with other works:

Haque, Shiu, Underwood, Van Riet (2008)

Flauger, Paban, Robbins, Wrane (2008)

⇒ Geometric fluxes may help in  $(\rho, \tau)$ -plane, but certainly do not automatically take care of all moduli ⇒ Many dangerous directions

More recent works:

Roest (2009), Dall'Agata, Villadoro, Zwieger (2009)

⇒ N=4 gauged supergravity

de Carlos, Guarino, Moreno (2009)

⇒ Non-geometric fluxes?

Danielsson, Haque, Shiu, Van Riet (2009)

⇒ 10D point of view

Consistent with other works:

Haque, Shiu, Underwood, Van Riet (2008)

Flauger, Paban, Robbins, Wrase (2008)

⇒ Geometric fluxes may help in  $(\rho, \tau)$ -plane, but certainly do not automatically take care of all moduli ⇒ Many dangerous directions

More recent works:

Roest (2009), Dall'Agata, Villadoro, Zwieger (2009)

⇒ N=4 gauged supergravity

de Carlos, Guarino, Moreno (2009)

⇒ Non-geometric fluxes?

Danielsson, Haque, Shiu, Van Riet (2009)

⇒ 10D point of view

So far no geometrical tree-level de Sitter...

## IIB on SU(2)-structure manifolds with

- O5/O7-planes
- $F_p$ -fluxes ( $p=1,3,5$ )

Cf. Lüst, Tsimpis (2009), Tsimpis' talk

## T-dual to IIA with D6 with possibly non-geometric fluxes

- Moduli stabilization ok
- Some new no-go theorems

Wrase, M.Z., ...

## Type IIB:

$$\begin{aligned}\rho &\equiv (\text{Vol})^{1/3} \\ \tilde{\tau} &\equiv e^{-\phi}\end{aligned}$$

$\Rightarrow$

$$\begin{aligned}V_{H_3} &\propto \tilde{\tau}^{-2} \rho^{-6} \\ V_{F_3} &\propto \tilde{\tau}^{-4} \rho^{-6} \\ V_{D3/O3} &\propto \pm \tilde{\tau}^{-3/2} \rho^{-6}\end{aligned}$$

$$\Rightarrow -\rho \partial_\rho V = 6V \Rightarrow V = 0$$

Cf. e.g. Hertzberg, Kachru, Taylor, Tegmark (2007)

$\Leftrightarrow$ No-scale structure of ISD Flux vacua in IIB

Giddings, Kachru, Polchinski (2001)

Haque, Shiu, Underwood, Van Riet (2008):

$V_{\text{geo}}$  and  $V_{F_1}$  may help

⇒ Need 1-cycles ⇒ No CY

SU(3)-structure also nontrivial  
(e.g.  $\Omega \wedge J = 0$ )

But:

∃ Explicit solutions in IIB on SU(2)-structure manifolds

E.g. Lüst, Tsimpis (2009)

Tried SU(2)-structure with O7/O5-planes

⇒ Moduli stabilization + no-go's

Wräse, M.Z., ...