Possible string effects in 4D from a 5D anisotropic gauge theory in a mean-field background

Nikos Irges, Wuppertal U.

Based on N.I . & F. Knechtli, arXiv:0905.2757
to appear in NPB + work in progress

Corfu, Sept. 2009

Background

1. J. M. Drouffe and J. B. Zuber, Phys. Rept. 102 (1983) 1

2. J. Zinn-Justin, Quantum Field Theories and Critical Phenomena

3. I. Montvay and G. Muenster, Quantum Fields on a Lattice

4. M. Creutz, Phys. Rev. Lett. 43 (1979), 553

5. B. Svetitsky and L.G. Yaffe, Nucl. Phys. B210 (1982) 423

6. Y.K. Fu and H.B. Nielsen, Nucl. Phys. B236 (1984), 167

7. S. Ejiri, J. Kubo and M. Murata, Phys. Rev. D62 (2000), 105025

8. M. Luscher and P. Weisz, JHEP 0407:014, 2204

In 5 Dimensions...

$$A_M \xrightarrow{\mathcal{M}_5 = E_4 \times S^1} \{A_\mu, A_5\}$$

5D gauge field SU(2)

Z: 4D gauge field

h: 4D "Higgs"

S^1

$+\phi(y + 2\pi R)$

S^1

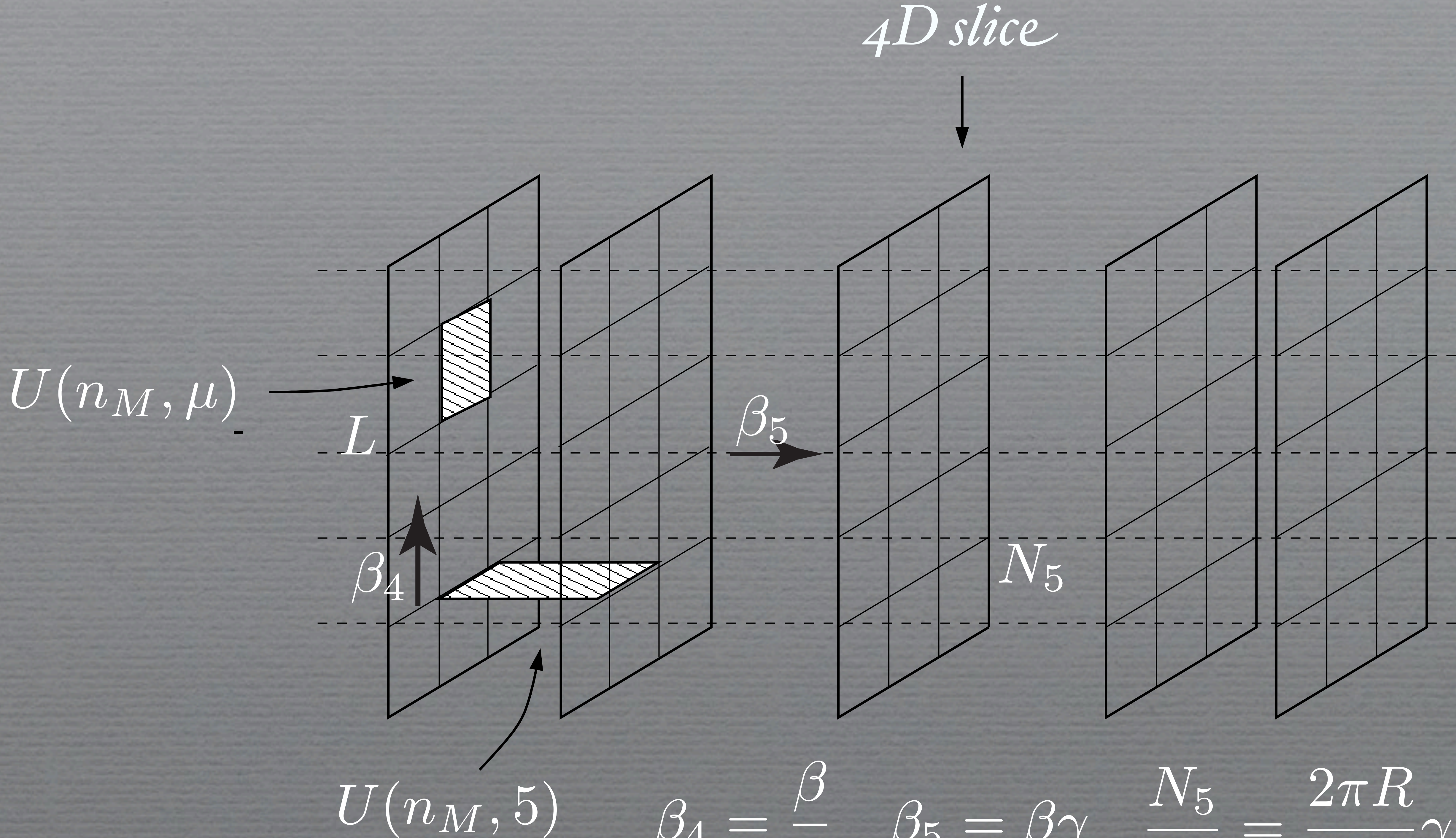
$+\psi(y + 2\pi R)$

$$\phi(y) = \sum_n \phi_n \, e^{iny/R}$$

$$\psi(y) = \sum_n \psi_x e^{iny/R}$$

4D Georgi-Glashow (G-G) model + KK

The Anisotropic Lattice



$$\beta_4 = \frac{\beta}{\gamma} \quad \beta_5 = \beta\gamma \quad \frac{N_5}{L} = \frac{2\pi R}{l}\gamma$$

$\gamma =$ anisotropy parameter

The mean-field expansion

parameters: $L,\ \beta,\ \gamma$ (we keep $N_5 = L$)

$$Z = \int DU \int DV \int DH e^{(1/N)\mathrm{Re}[\mathrm{tr}H(U-V)]} e^{-S_G[V]}$$

$$Z = \int DV \int DH e^{-Seff[V,H]}, \quad Seff = S_G[V] + u(H) + (1/N)\mathrm{Retr}HV$$

$$e^{-u(H)} = \int DU e^{(1/N)\mathrm{Retr}UH}$$

To 0'th order

The background is determined by

$$\overline{V} = -\frac{\partial u}{\partial H}\bigg|_{\overline{H}} \qquad \overline{H} = -\frac{\partial S_G[V]}{\partial V}|_{\overline{V}}$$

The free energy

$$F^{(0)} = -\frac{1}{\mathcal{N}}\ln(Z[\overline{V},\overline{H}]) = \frac{S_{\text{eff}}[\overline{V},\overline{H}]}{\mathcal{N}}$$

To 1st order

$$H = \bar{H} + h \qquad V = \bar{V} + v$$

$$Seff = Seff[\bar{V}, \bar{H}] + \frac{1}{2}\left(\left.\frac{\delta^2 Seff}{\delta H^2}\right|_{\overline{V},\overline{H}} h^2 + 2 \left.\frac{\delta^2 Seff}{\delta H \delta V}\right|_{\overline{V},\overline{H}} hv + \left.\frac{\delta^2 Seff}{\delta V^2}\right|_{\overline{V},\overline{H}} v^2 \right)$$

$$\left.\frac{\delta^2 Seff}{\delta H^2}\right|_{\overline{V},\overline{H}} h^2 = h_i K_{ij}^{(hh)} h_j = h^T K^{(hh)} h$$

$$\left.\frac{\delta^2 Seff}{\delta V^2}\right|_{\overline{V},\overline{H}} v^2 = v_i K_{ij}^{(vv)} v_j = v^T K^{(vv)} v$$

$$\left.\frac{\delta^2 Seff}{\delta V \delta H}\right|_{\overline{V},\overline{H}} v^2 = v_i K_{ij}^{(vh)} h_j = v^T K^{(vh)} h$$

The Gaussian fluctuations

$$z = \int Dv \int Dh e^{-S^{(2)}[v,h]} \qquad S^{(2)}[v,h] = \frac{1}{2}\left(h^T K^{(hh)} h + 2v^T K^{(vh)} h + v^T K^{(vv)} v\right)$$

$$z = \frac{(2\pi)^{|h|/2}(2\pi)^{|v|/2}}{\sqrt{\det[(-\mathbf{1} + K^{(hh)}K^{(vv)})]}}$$

$$Z^{(1)} = Z[\overline{V}, \overline{H}] \cdot z = e^{-S_{\mathrm{eff}}[\overline{V},\overline{H}]} \cdot z$$

The free energy to first order

$$F^{(1)} = F^{(0)} - \frac{1}{\mathcal{N}} \ln(z) = F^{(0)} + \frac{1}{2\mathcal{N}} \ln\Big[\det\left(-\mathbf{1} + K^{(hh)}K^{(vv)}\right) \Delta_{\mathrm{FP}}^{-2}\Big]$$

Observables

$$\mathcal{O}[V] = \mathcal{O}[\overline{V}] + \left.\frac{\delta \mathcal{O}}{\delta V}\right|_{\overline{V}} v + \frac{1}{2}\left.\frac{\delta^2 \mathcal{O}}{\delta V^2}\right|_{\overline{V}} v^2 + \dots$$

$$= \mathcal{O}[\overline{V}] + \frac{1}{2}\left.\frac{\delta^2 \mathcal{O}}{\delta V^2}\right|_{\overline{V}} \frac{1}{z}\int Dv \int Dh v^2 e^{-S^{(2)}[v,h]}$$

$$\langle \mathcal{O} \rangle = \frac{1}{Z}\int Dv \int Dh \left(\mathcal{O}[\overline{V}] + \frac{1}{2}\left.\frac{\delta^2 \mathcal{O}}{\delta V^2}\right|_{\overline{V}} v^2 \right) e^{-\left(Seff[\overline{V},\overline{H}]+S^{(2)}[v,h]\right)}$$

$$< v_i v_j >= \frac{1}{z}\int Dv \int Dh\, v_i v_j e^{-S^{(2)}[v,h]} = (K^{-1})_{ij}$$

$$K = -K^{(vh)} {K^{(hh)}}^{-1} K^{(vh)} + K^{(vv)}$$

$$\boxed{< \mathcal{O} >= \mathcal{O}[\overline{V}] + \frac{1}{2}\mathrm{tr}\left\{ \left.\frac{\delta^2 \mathcal{O}}{\delta V^2}\right|_{\overline{V}} K^{-1} \right\}}$$

1st order master formula

Correlators

$$C(t) = < \mathcal{O}(t_0+t)\mathcal{O}(t_0) > - < \mathcal{O}(t_0+t) >< \mathcal{O}(t_0) > = C^{(0)}(t) + C^{(1)}(t) + \cdots$$

$$C^{(0)}(t) = 0$$

$$< \mathcal{O}(t_0+t)\mathcal{O}(t_0) > = \mathcal{O}^{(0)}(t_0+t)\mathcal{O}^{(0)}(t_0) + \frac{1}{2}\mathrm{tr}\left\{ \frac{\delta^2(\mathcal{O}(t_0+t)\mathcal{O}(t_0))}{\delta^2 v} K^{-1} \right\} + \cdots$$

$$C^{(1)}(t) = \frac{1}{2}\mathrm{tr}\left\{ \frac{\delta^{(1,1)}(\mathcal{O}(t_0+t)\mathcal{O}(t_0))}{\delta^2 v} K^{-1} \right\} = \frac{1}{2}\mathrm{tr}\left\{ \frac{\tilde{\delta}^{(1,1)}(\mathcal{O}(t_0+t)\mathcal{O}(t_0))}{\delta^2 v} \tilde{K}^{-1} \right\}$$

$$C^{(1)}(t) = \sum_\lambda c_\lambda e^{-E_\lambda t} \qquad E_0 = m_H, \quad E_1 = m_H^*, \cdots$$

To 2nd order

$$S_{\rm eff} = S_{\rm eff}[\overline{V},\overline{H}] + \frac{1}{2}\left(\frac{\delta^2 S_{\rm eff}}{\delta H^2}h^2 + 2\frac{\delta^2 S_{\rm eff}}{\delta H \delta V}hv + \frac{\delta^2 S_{\rm eff}}{\delta V^2}v^2\right)$$

$$+\frac{1}{6}\left(\frac{\delta^3 S_{\rm eff}}{\delta H^3}h^3 + \frac{\delta^3 S_{\rm eff}}{\delta V^3}v^3\right) + \frac{1}{24}\left(\frac{\delta^4 S_{\rm eff}}{\delta H^4}h^4 + \frac{\delta^4 S_{\rm eff}}{\delta V^4}v^4\right) + \cdots$$

$$\mathcal{O}[V] = \mathcal{O}[\overline{V}] + \frac{\delta \mathcal{O}}{\delta V}v + \frac{1}{2}\frac{\delta^2 \mathcal{O}}{\delta V^2}v^2 + \frac{1}{6}\frac{\delta^3 \mathcal{O}}{\delta V^3}v^3 + + \frac{1}{24}\frac{\delta^4 \mathcal{O}}{\delta V^4}v^4 + \cdots$$

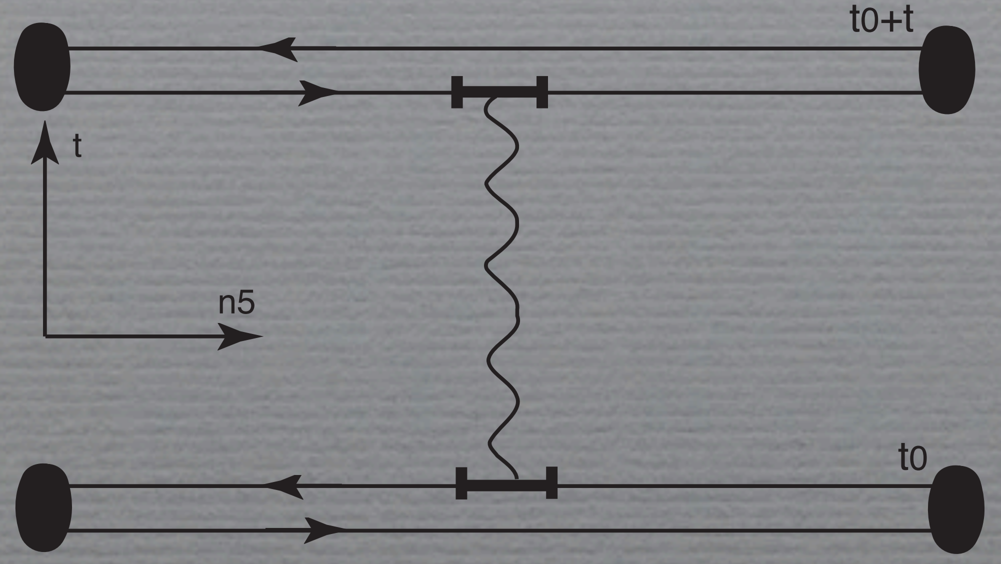
2nd order master formula

$$\boxed{\begin{aligned} &< \mathcal{O} >= \mathcal{O}[\overline{V}] + \frac{1}{2}\left(\frac{\delta^2 \mathcal{O}}{\delta V^2}\right)_{ij}\left(K^{-1}\right)_{ij} \\ &+\frac{1}{24}\sum_{i,j,l,m}\left(\frac{\delta^4 \mathcal{O}}{\delta V^4}\right)_{ijlm}\left((K^{-1})_{ij}(K^{-1})_{lm} + (K^{-1})_{il}(K^{-1})_{jm} + (K^{-1})_{im}(K^{-1})_{jl}\right)\end{aligned}}$$

$$C^{(2)}(t) = \frac{1}{24}\sum_{i,j,l,m}\left(\frac{\delta^4 \mathcal{O}^c(t)}{\delta v^4}\right)_{ijlm}\left((K^{-1})_{ij}(K^{-1})_{lm} + (K^{-1})_{il}(K^{-1})_{jm} + (K^{-1})_{im}(K^{-1})_{jl}\right)$$

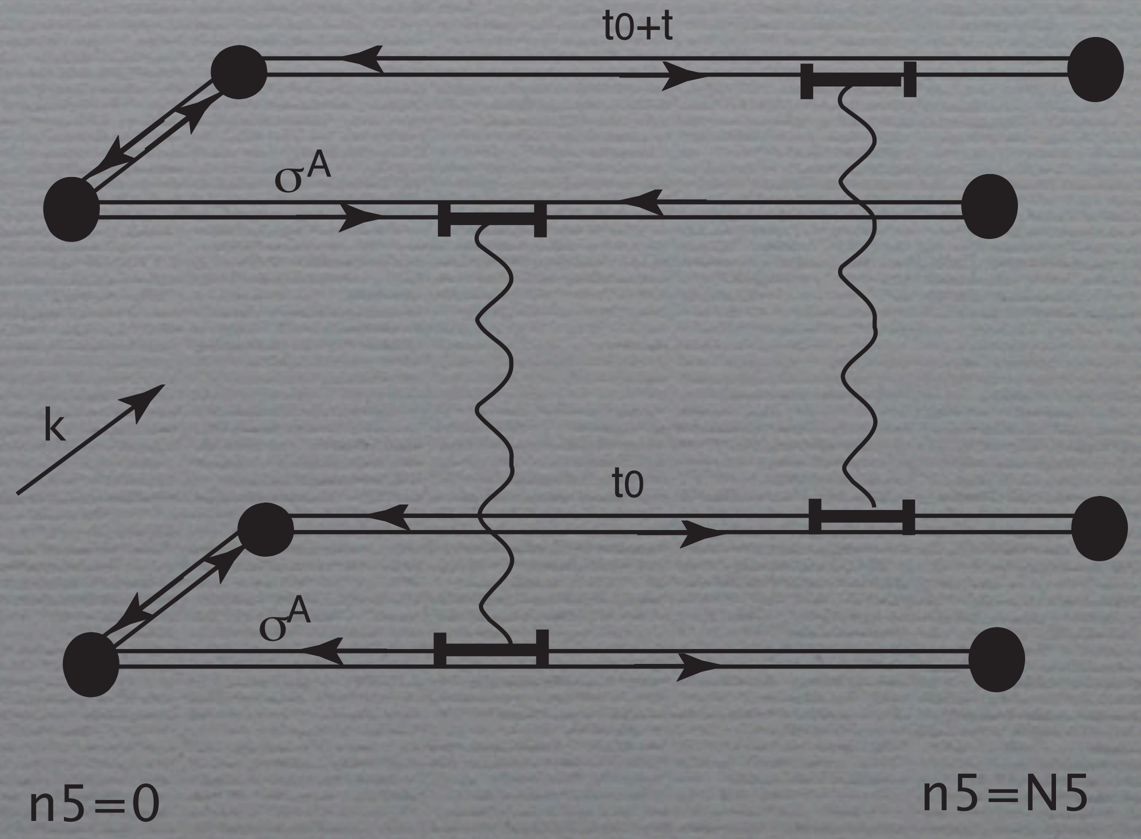
Lattice Observables (Polyakov loops)

The scalar



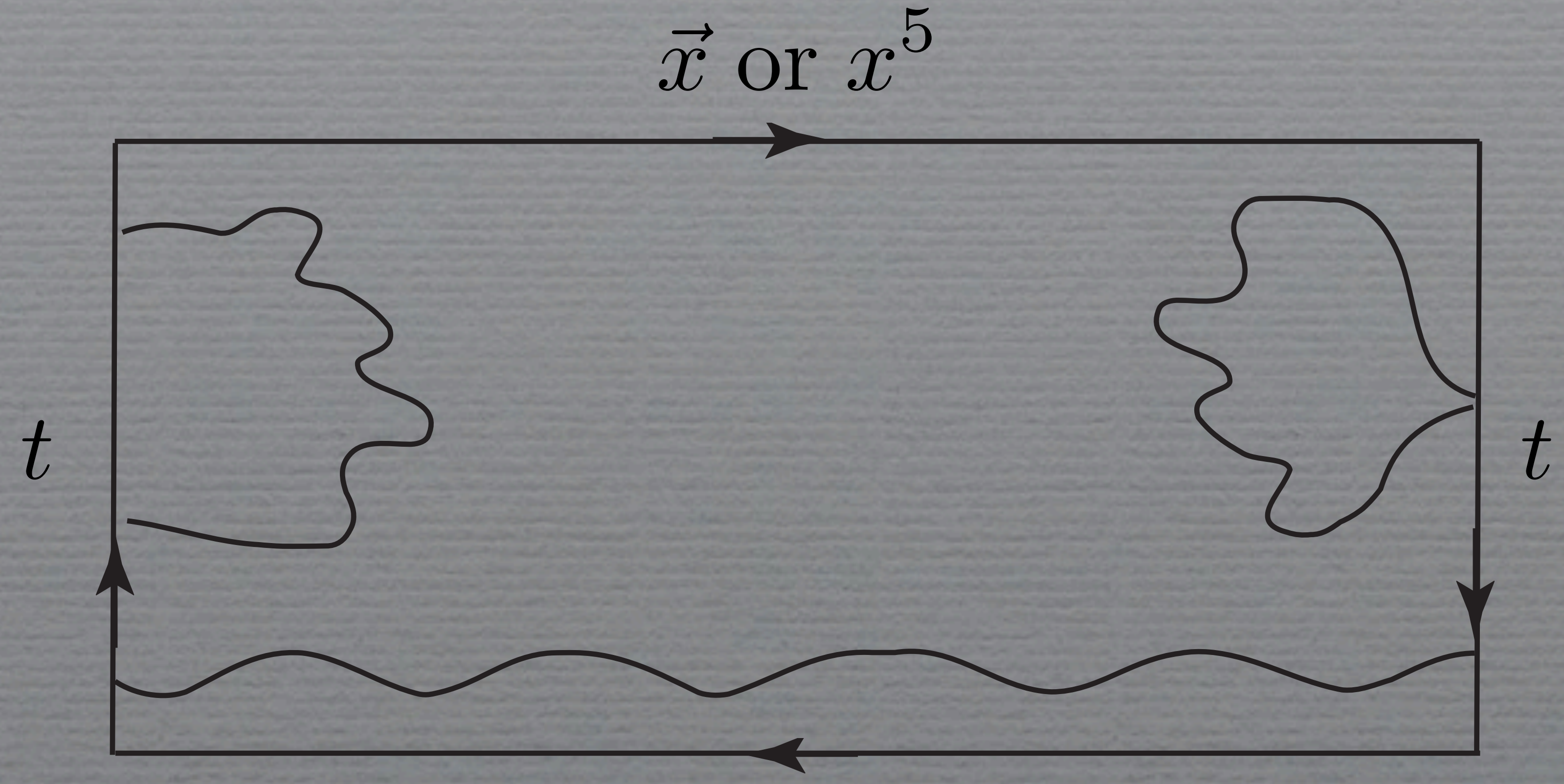
$$m = \lim_{t\to\infty} \ln \frac{C^{(1)}(t)}{C^{(1)}(t-1)}$$

The vector



$$m = \lim_{t\to\infty} \ln \frac{C^{(2)}(t)}{C^{(2)}(t-1)}$$

The Wilson loop



$$\mathrm{t} \to \infty : \qquad \mathrm{e}^{-Vt} \simeq < \mathcal{O}_W >$$

The SU(2) model

Mean-field parametrization

$$U(n,M) = u_0(n,M)\mathbf{1} + i\sum_k u_k(n,M)\sigma^k \qquad u_\alpha \in \mathbb{R}, u_\alpha u_\alpha = 1$$

$$V(n,M) = v_0(n,M)\mathbf{1} + i\sum_k v_k(n,M)\sigma^k\,, \quad v_\alpha \in \mathbb{C}$$

$$H(n,M) = h_0(n,M)\mathbf{1} - i\sum_k h_k(n,M)\sigma^k\,, \quad h_\alpha \in \mathbb{C}$$

The propagator (in momentum space)

$$\tilde{K}(p',M',\alpha';p'',M'',\alpha'') = \delta_{p'p''}\delta_{\alpha'\alpha''}C_{M'M''}(p',\alpha')$$

$$C_{M'M''}(p',\alpha') = [A\delta_{M'M''} + B_{M'M''}(1-\delta_{M'M''})]$$

$$A = -\left[\frac{1}{b_2}(1-\delta_{\alpha'0}) + \frac{1}{b_1}\delta_{\alpha'0}\right] - 2\beta\overline{v}_0^2\left[\sum_{N\neq M'}\cos(p'_N) + \frac{1}{\xi}\sin^2(p'_{M'}/2)\right]$$

$$B_{M'M''} = -4\beta\overline{v}_0{}^2\left[\delta_{\alpha'0}\cos\left(\frac{p'_{M'}}{2}\right)\cos\left(\frac{p'_{M''}}{2}\right) +(1-\delta_{\alpha'0})\sin\left(\frac{p'_{M'}}{2}\right)\sin\left(\frac{p'_{M''}}{2}\right)\right]$$

$$b_1 = -\frac{1}{\overline{h}_0 I_1(\overline{h}_0)}\left(I_2(\overline{h}_0) - \overline{h}_0\left(\frac{I_2(\overline{h}_0)^2}{I_1(\overline{h}_0)} - I_3(\overline{h}_0)\right)\right)$$

$$b_2 = -\frac{\overline{v}_0}{\overline{h}_0}$$

The static potential

$$V(r) = -2\log(\overline{v}_0) - \frac{1}{2\overline{v}_0^2}\frac{1}{L^3 N_5} \times \left\{ \sum_{p'_{M\neq 0}, p'_0=0} \left[\frac{1}{4}\sum_{N\neq 0}(2\cos(p'_N r)+2)\right] C_{00}^{-1}(p',0)\right]$$

$$+3 \sum_{p'_{M\neq 0}, p'_0=0} \left[\frac{1}{4}\sum_{N\neq 0}(2\cos(p'_N r)-2)\right] \frac{1}{C_{00}(p',1)}\right\}$$

The scalar mass

$$C_H^{(1)}(t) = \frac{1}{\mathcal{N}}(P_0^{(0)})^2 \sum_{p'_0} \cos(p'_0 t) \sum_{p'_5} |\tilde{\Delta}^{(\mathcal{N}_5)}(p'_5)|^2 \tilde{K}^{-1}\Big((p'_0,\vec{0},p'_5),5,0;(p'_0,\vec{0},p'_5),5,0\Big)$$

$$\Delta^{(m_5)}(n_5) = \sum_{r=0}^{m_5-1} \frac{\delta_{n_5 r}}{\overline{v}_0(\hat{r})}, \qquad \hat{r} = r+1/2$$

The vector mass

$$C_V^{(2)}(t) = \frac{2304}{\mathcal{N}^2}(P_0^{(0)})^4(\overline{v}_0(0))^4 \sum_{\vec{p}'} \sum_k \sin^2(p'_k)\left(\overline{\overline{K}}^{-1}(t,\vec{p}',1)\right)^2$$

$$\overline{K}^{-1}((p'_0,\vec{p}'),5,\alpha) = \sum_{p'_5,p''_5} \tilde{\Delta}^{(N_5)}(p'_5)\tilde{\Delta}^{(N_5)}(-p''_5)K^{-1}(p'',5,\alpha;p',5,\alpha)$$

$$\overline{\overline{K}}^{-1}(t,\vec{p}',\alpha) = \sum_{p'_0} e^{ip'_0 t}\overline{K}^{-1}((p'_0,\vec{p}'),5,\alpha)$$

The free energy

$$F^{(1)} = F^{(0)} + \frac{1}{2\mathcal{N}}\sum_p \ln\left[\det\left(-\mathbf{1} + \tilde{K}^{(hh)}_{\alpha'=0}\tilde{K}^{(vv)}_{\alpha'=0}\right)\det\left(-\mathbf{1} + \tilde{K}^{(hh)}_{\alpha'\neq 0}\tilde{K}^{(vv)}_{\alpha'\neq 0}\right)^3 \Delta_{\mathrm{FP}}^{-2}\right]$$

A few remarks on the calculation

1. We fix the Lorentz gauge;
We pick up a Fadeev-Popov determinant but no ghost loops at this order; Gauge dependent free energy but no worries

2. Torons appear always as 0/0 --> "0/0=0 regularization"

3. Formulas are generalized to the anisotropic lattice -->
2 independent Wilson loops since the background is different along 4D and the 5th dimension

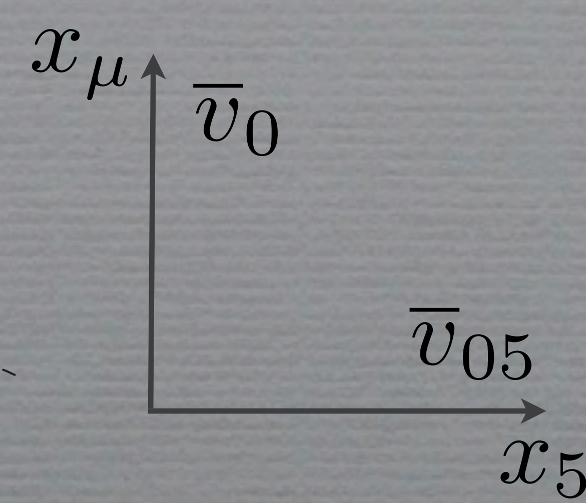
4. Analytical formulas are computed numerically

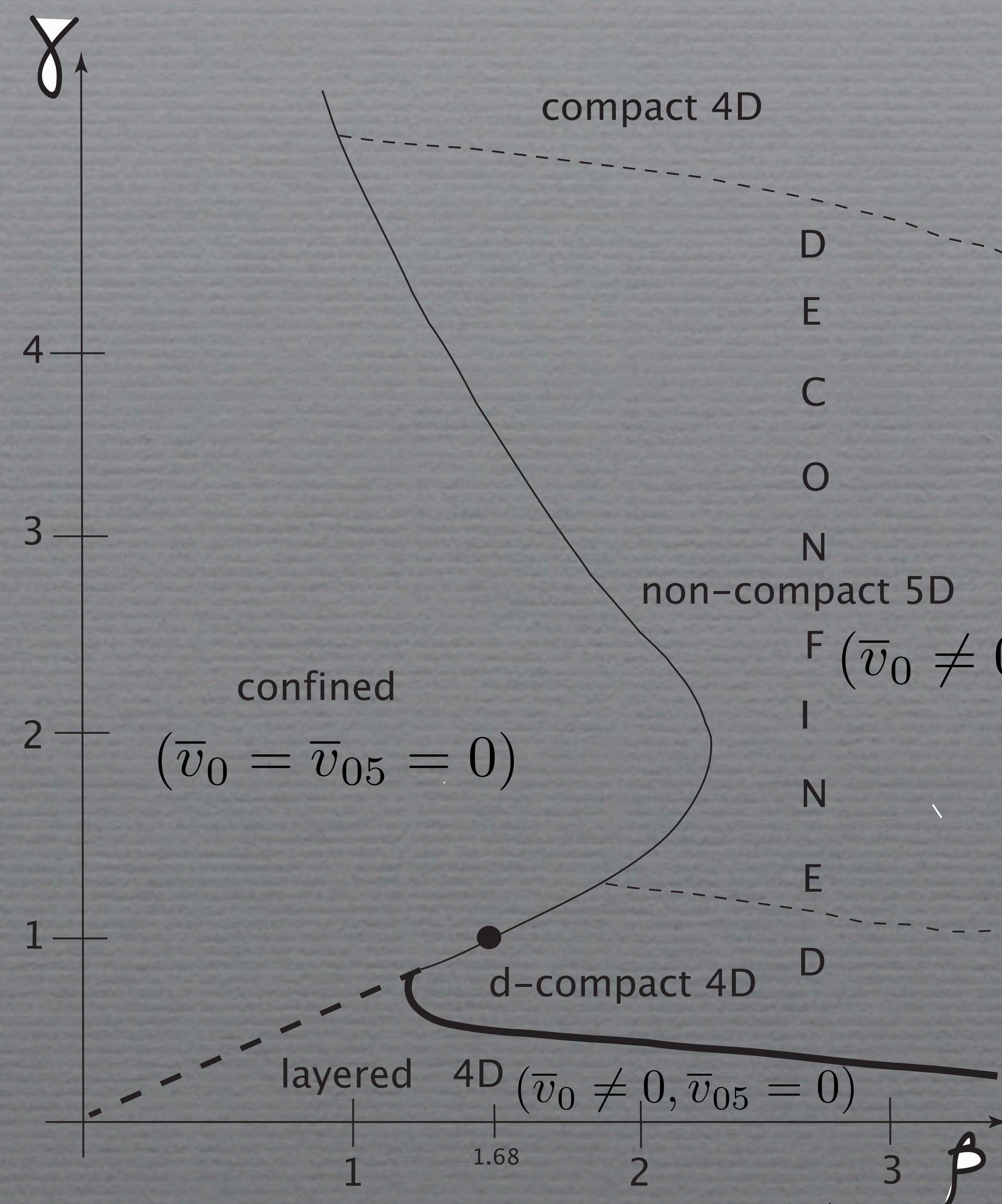
5. Our working assumption:
whenever non-trivial, physical

Keep this in mind

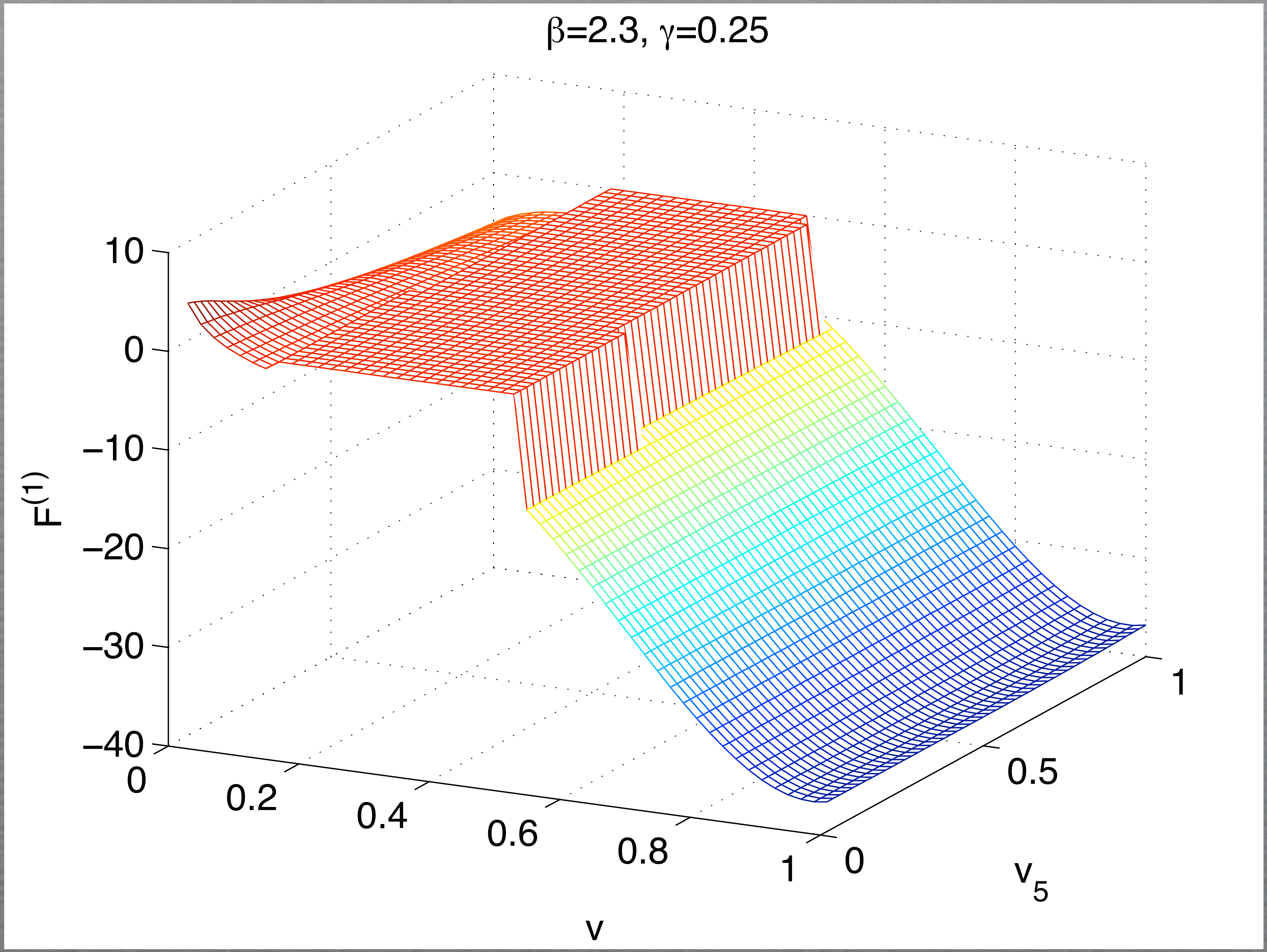
The phase diagram

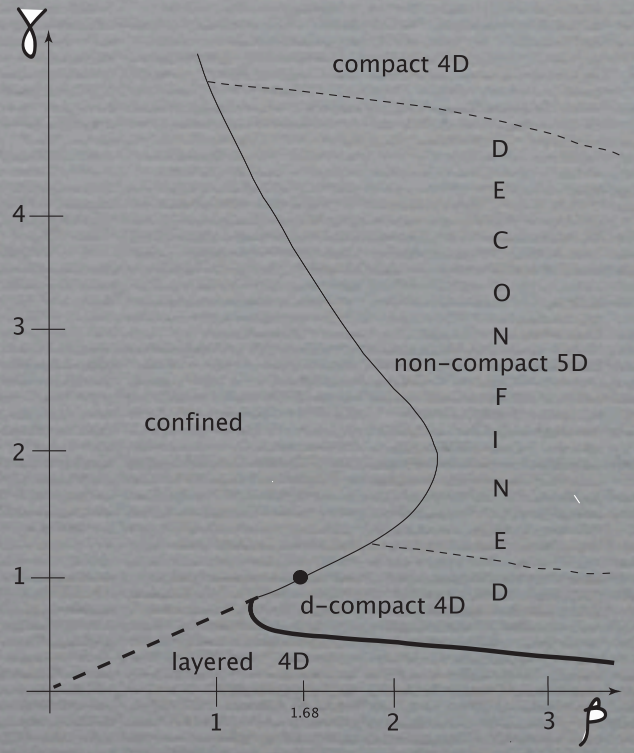
the mean-field background





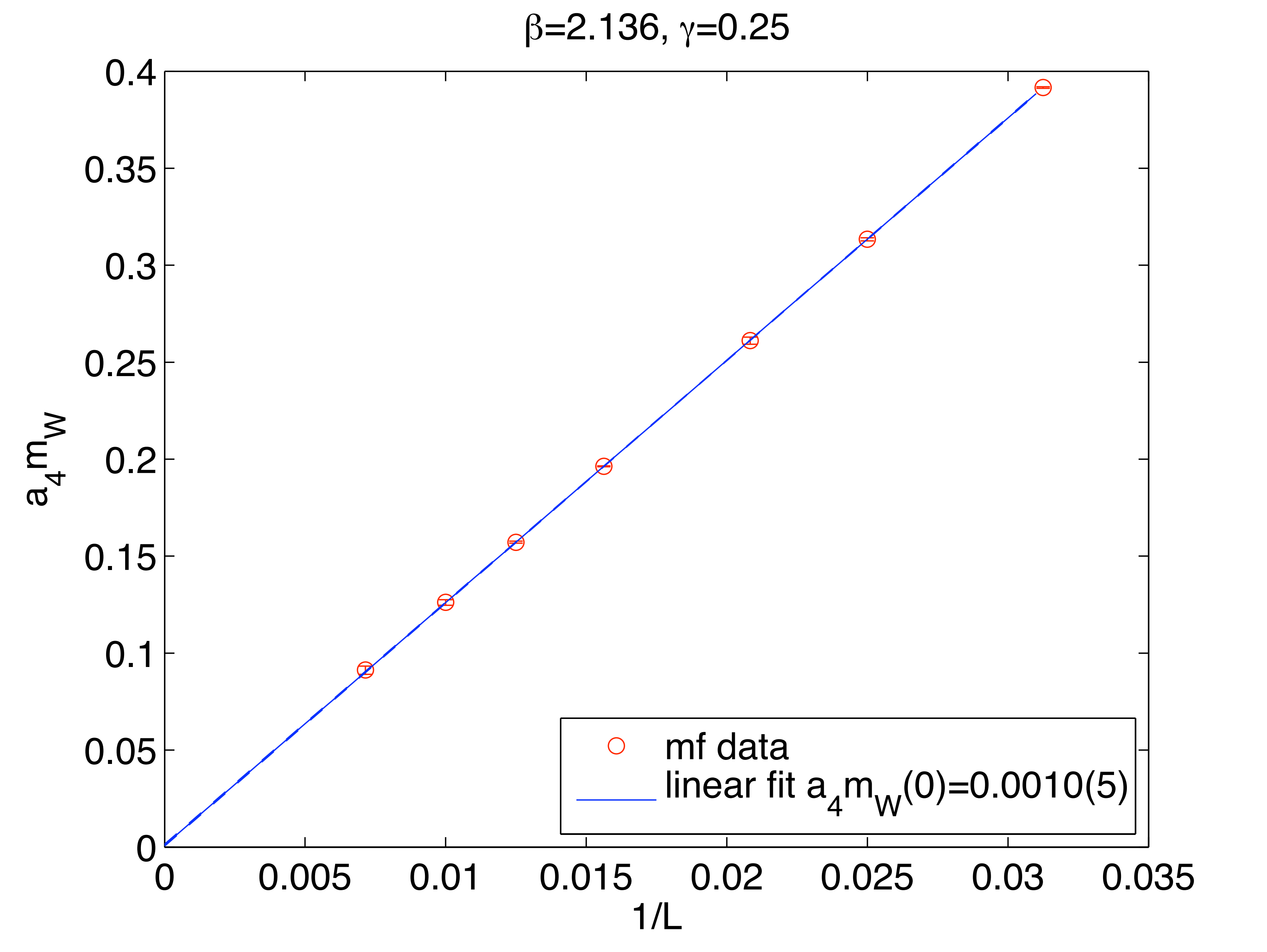
The d-compact phase is stable:

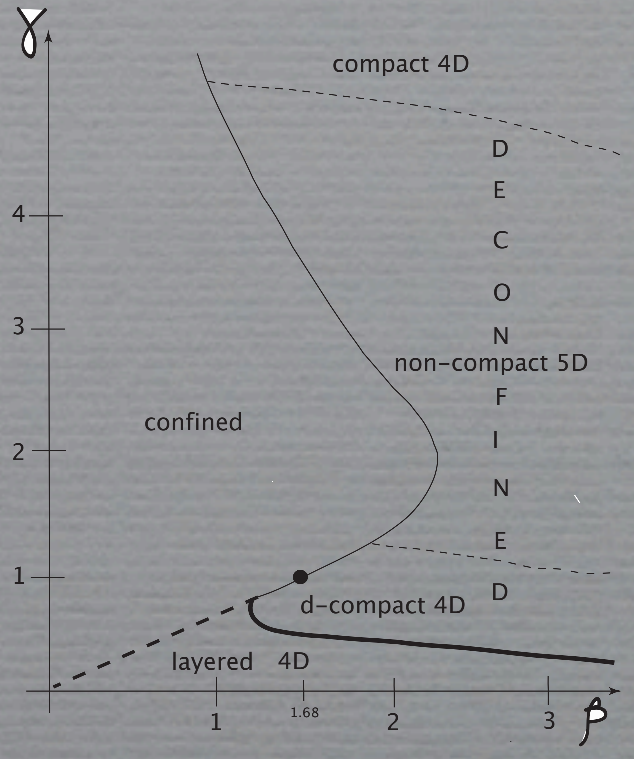




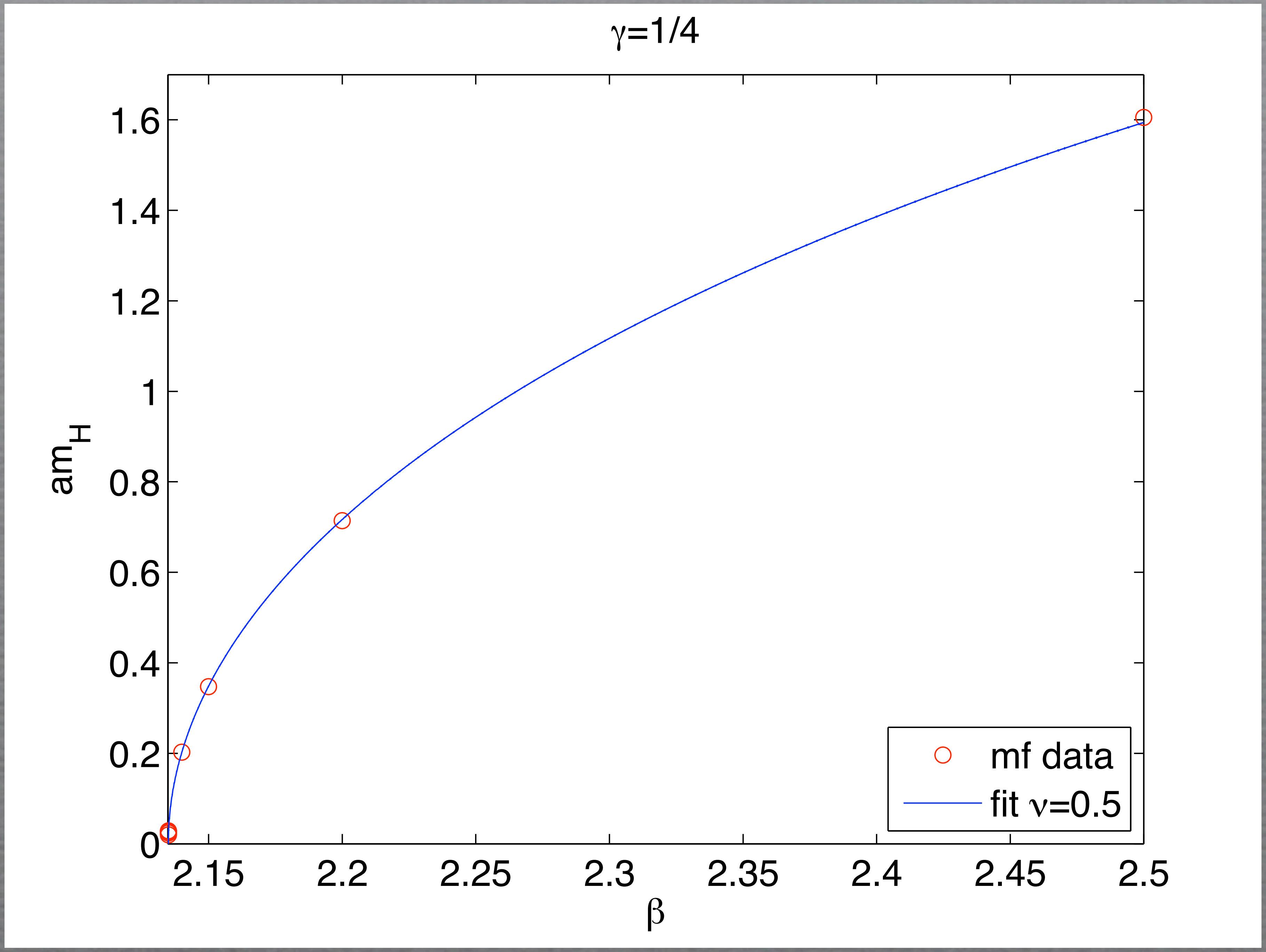
The compact phase (for small β) is unstable to this order

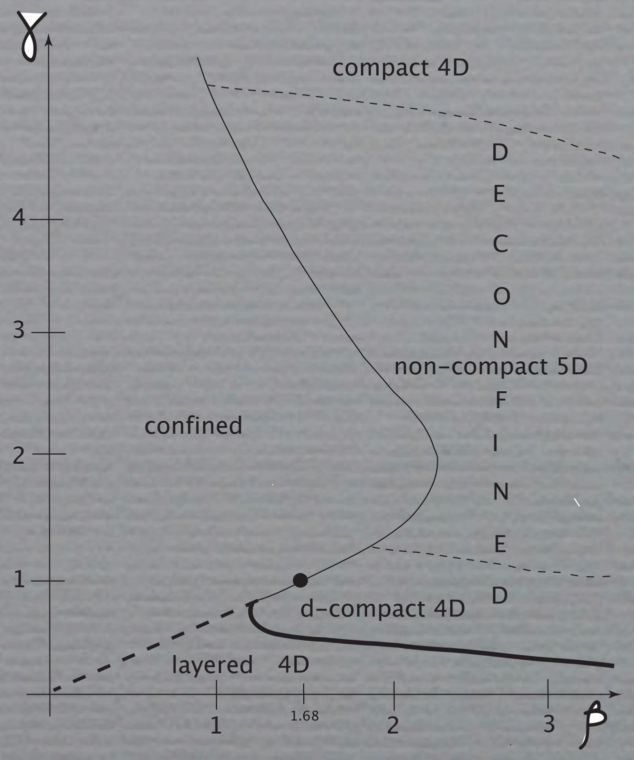
The W is massless or perhaps exponentially light:





The d-compact phase is separated from the layered phase by a 2nd order phase transition:



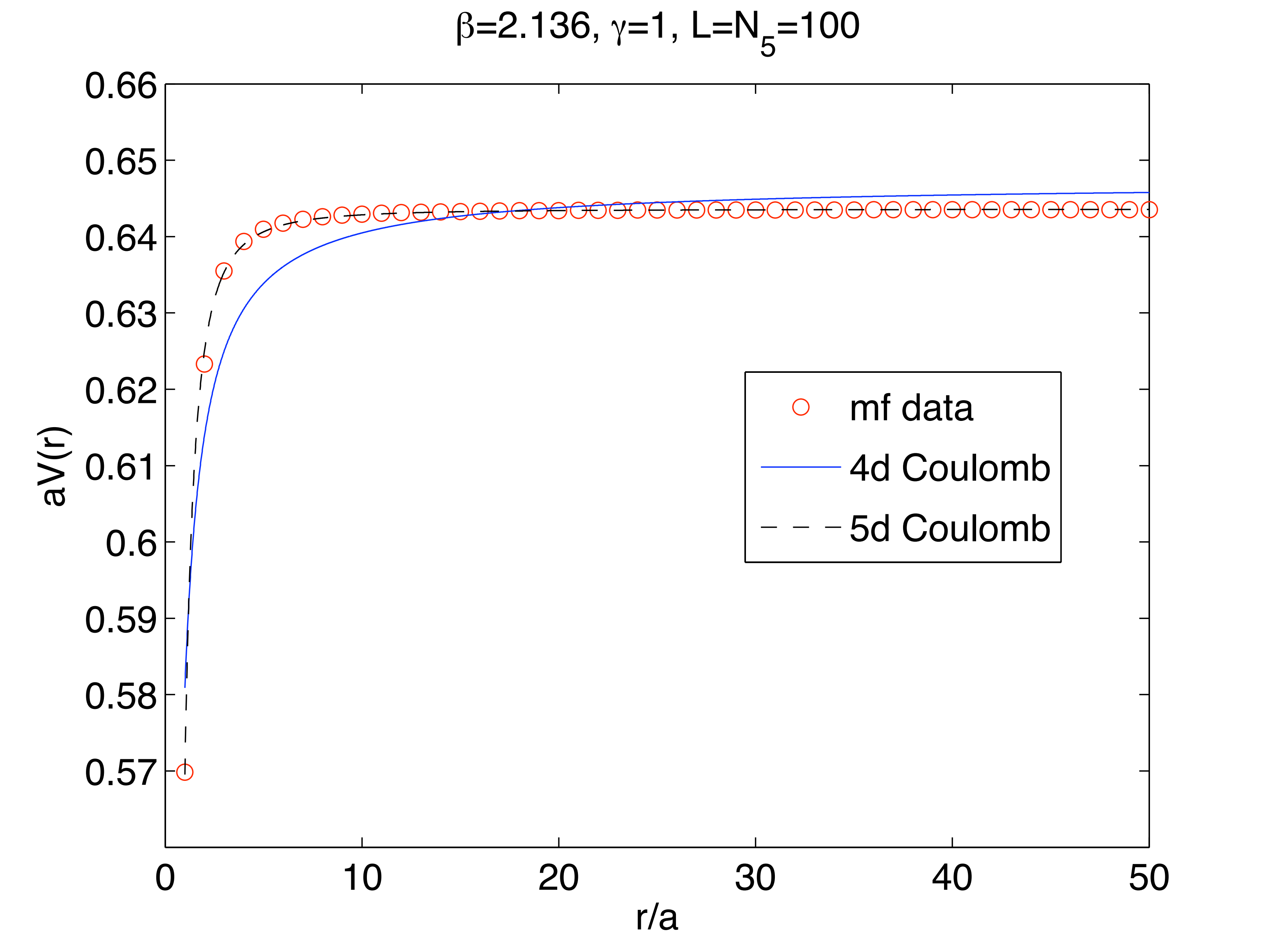


$$\nu = 1/2$$

agrees with Svetitski & Yaffe: dim reduction to 4d Ising

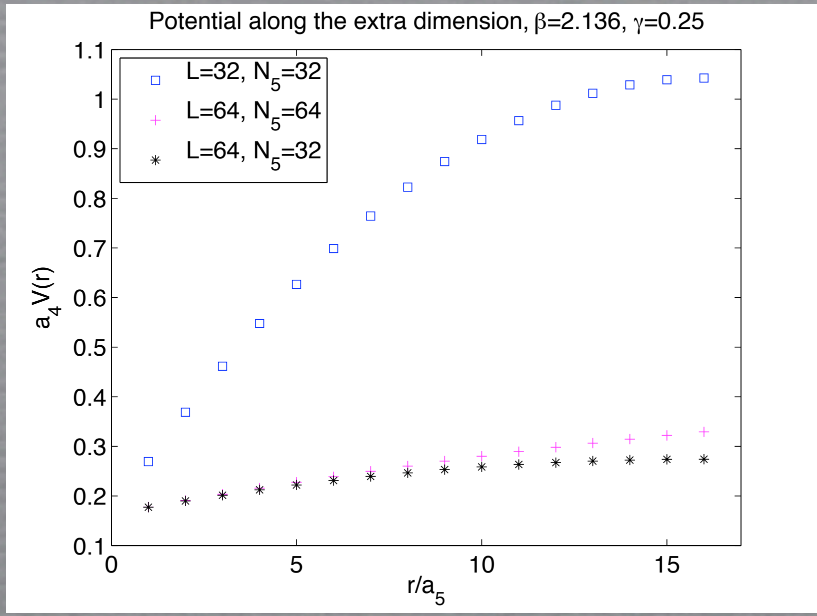
The Static Potential

The static potential on the isotropic lattice:

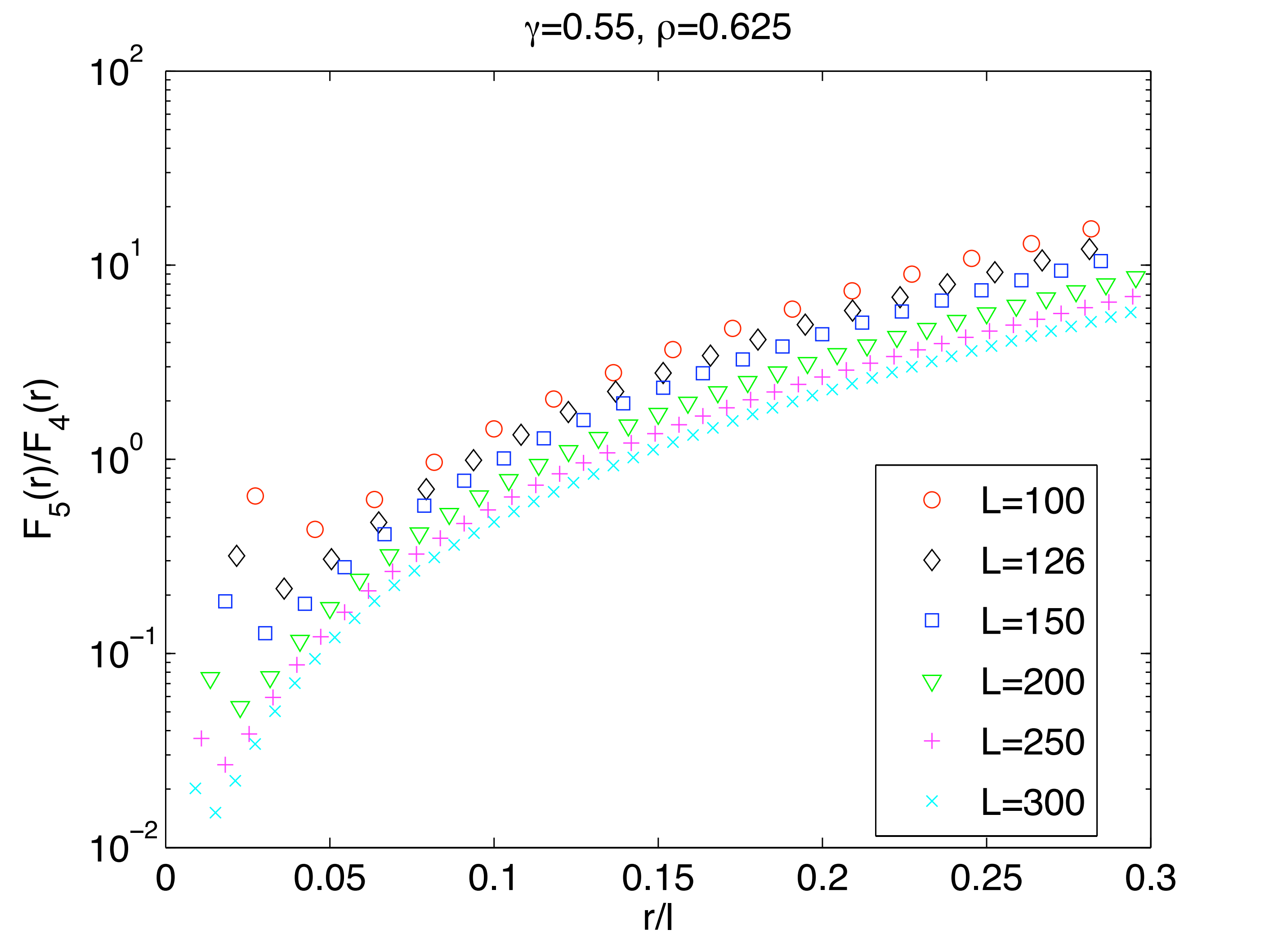


The static potential (Wilson loop) in the d-compact phase along the extra dimension:

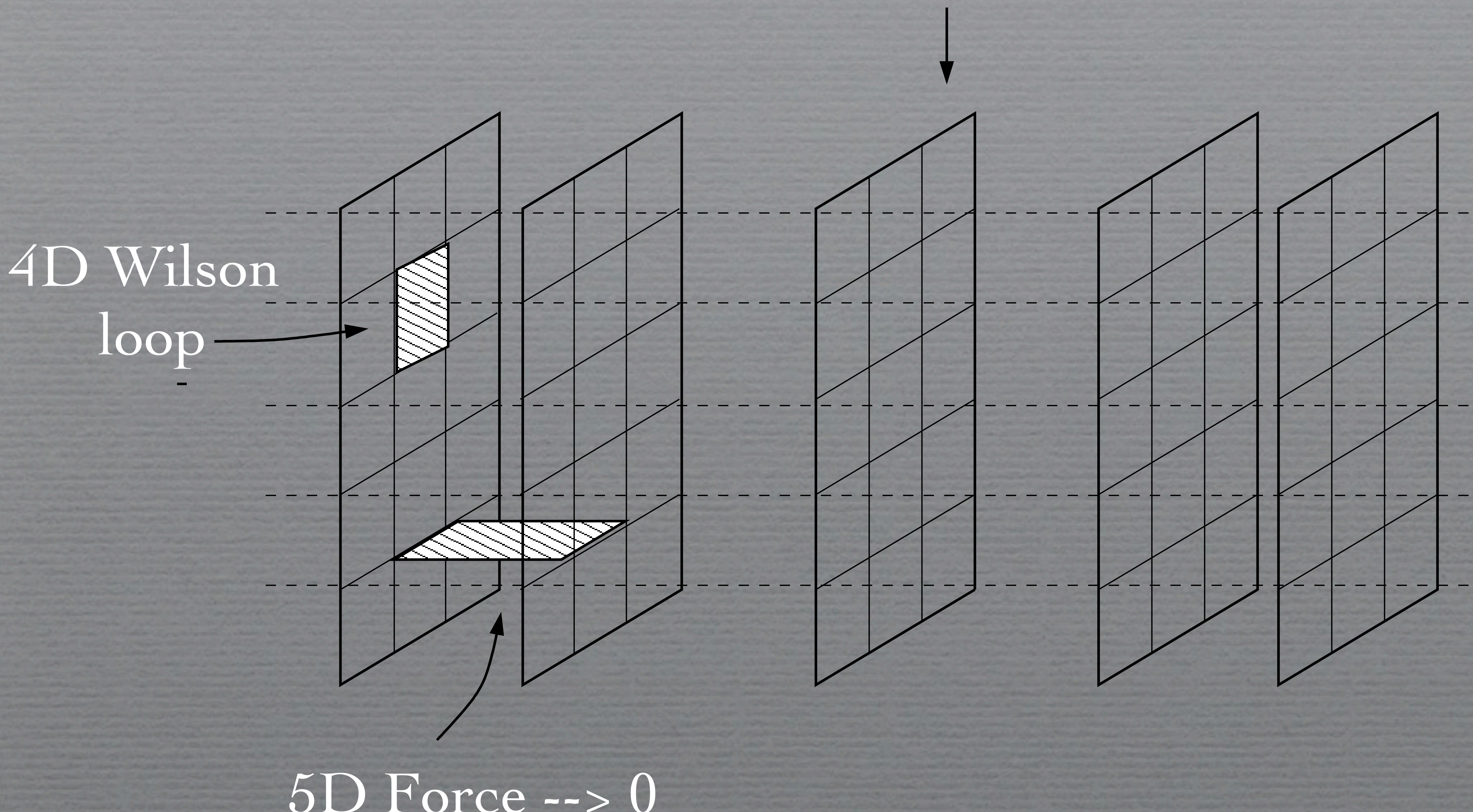
It doesn't scale with L:



The 5D force vanishes in the infinite volume-continuum limit:

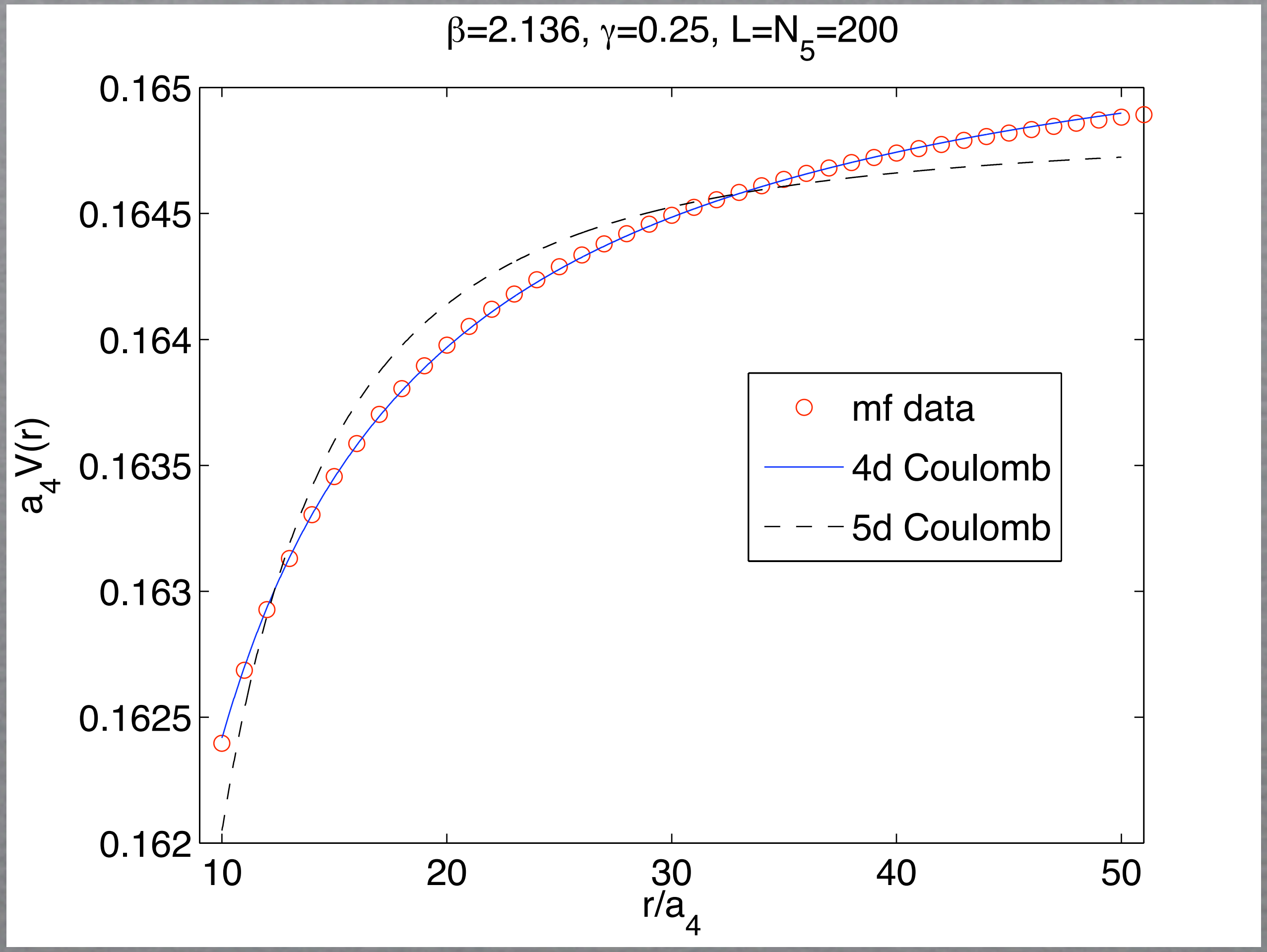


We are looking at an array of non-interacting 4D hyperplanes



The short distance static potential in the d-compact phase along the 4d hyperplanes:

β_4 large and β_5 small

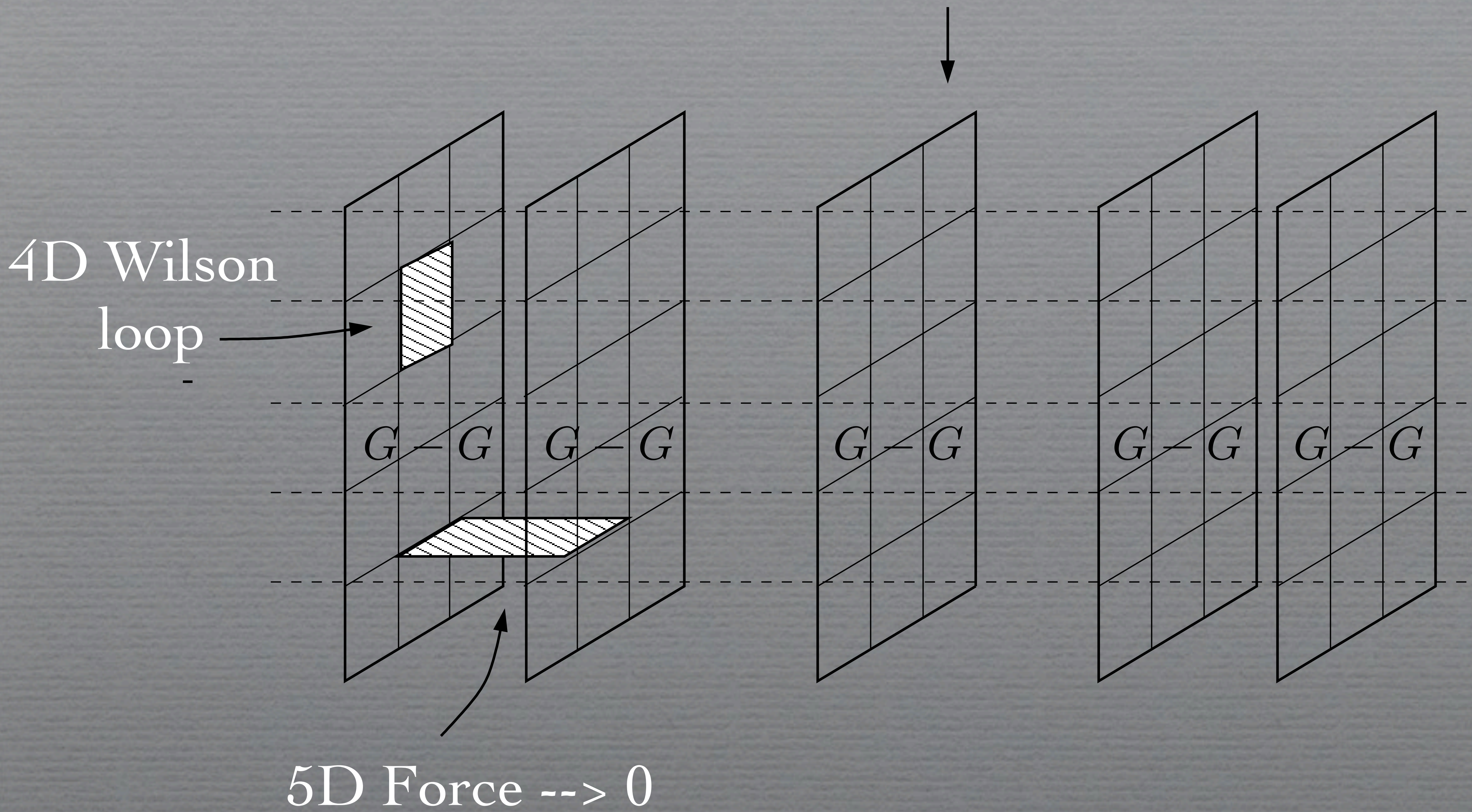


In the stability region r=10...50 the potential is clearly 4d Coulomb, large Yukawa excluded

$\longrightarrow$

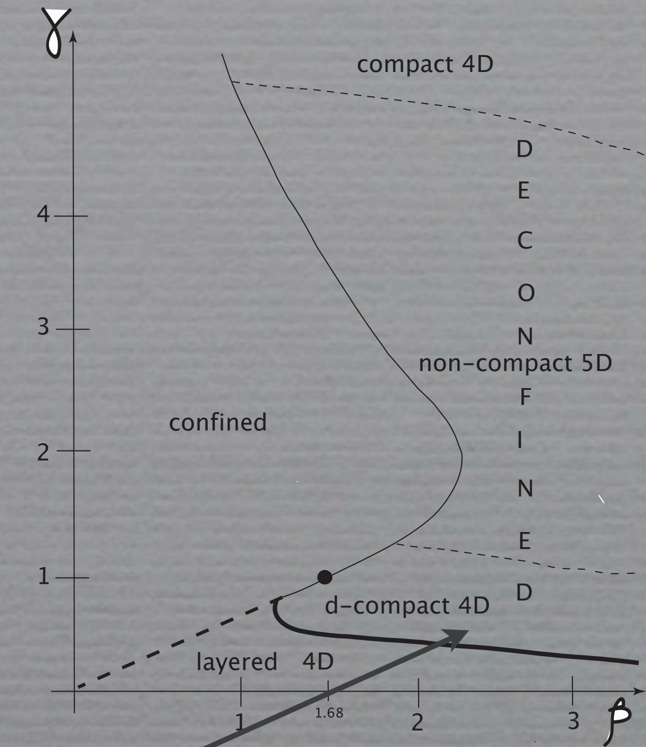
dimensional reduction in the d-compact phase

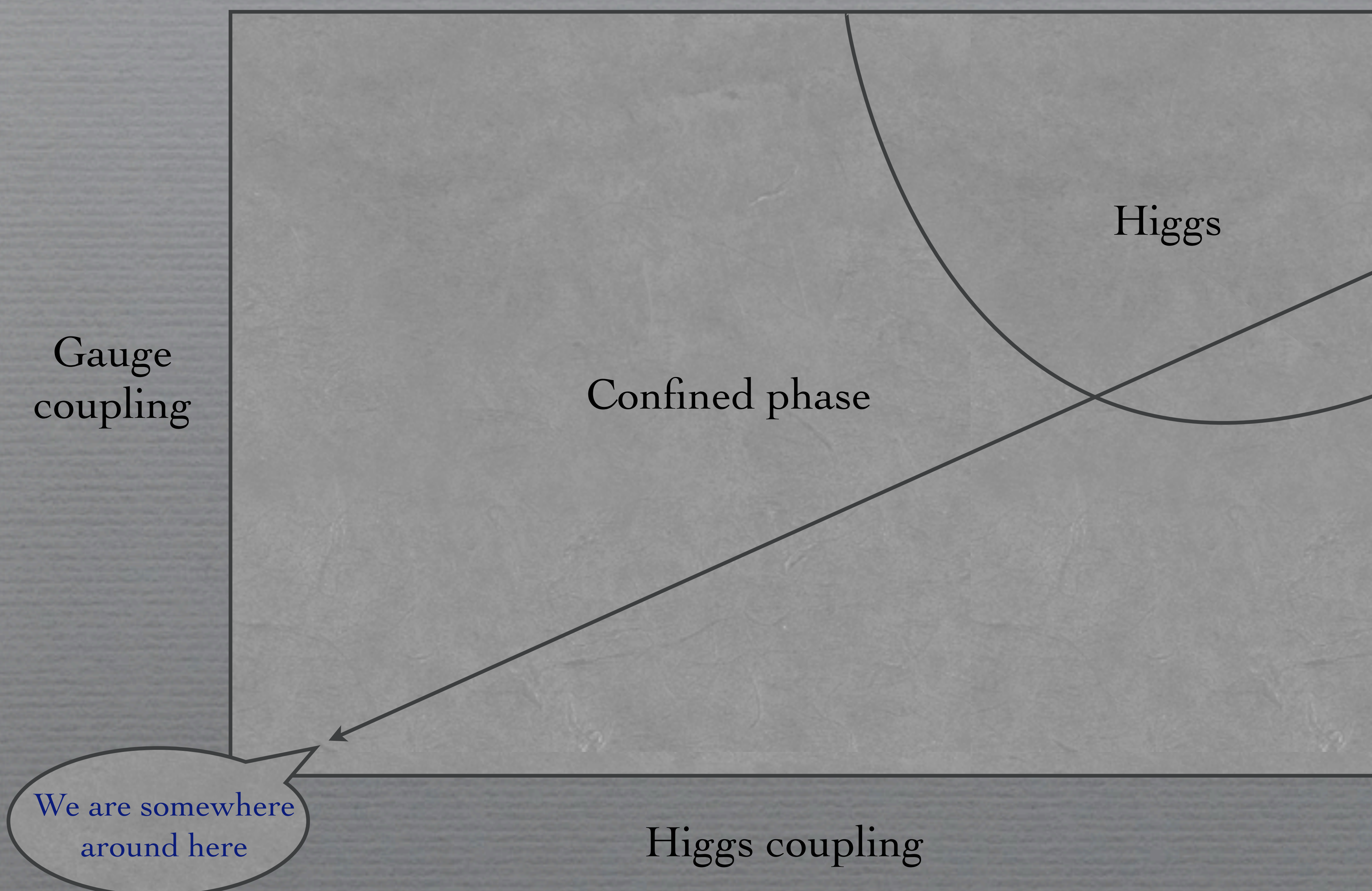
We are looking at an array of non-interacting 4D hyperplanes where gauge interactions are localized.
4D Georgi-Glashow on each of the branes



The localization is "perfect" in the infinite L limit

Remember: 4D Georgi-Glashow model + ~~KK~~

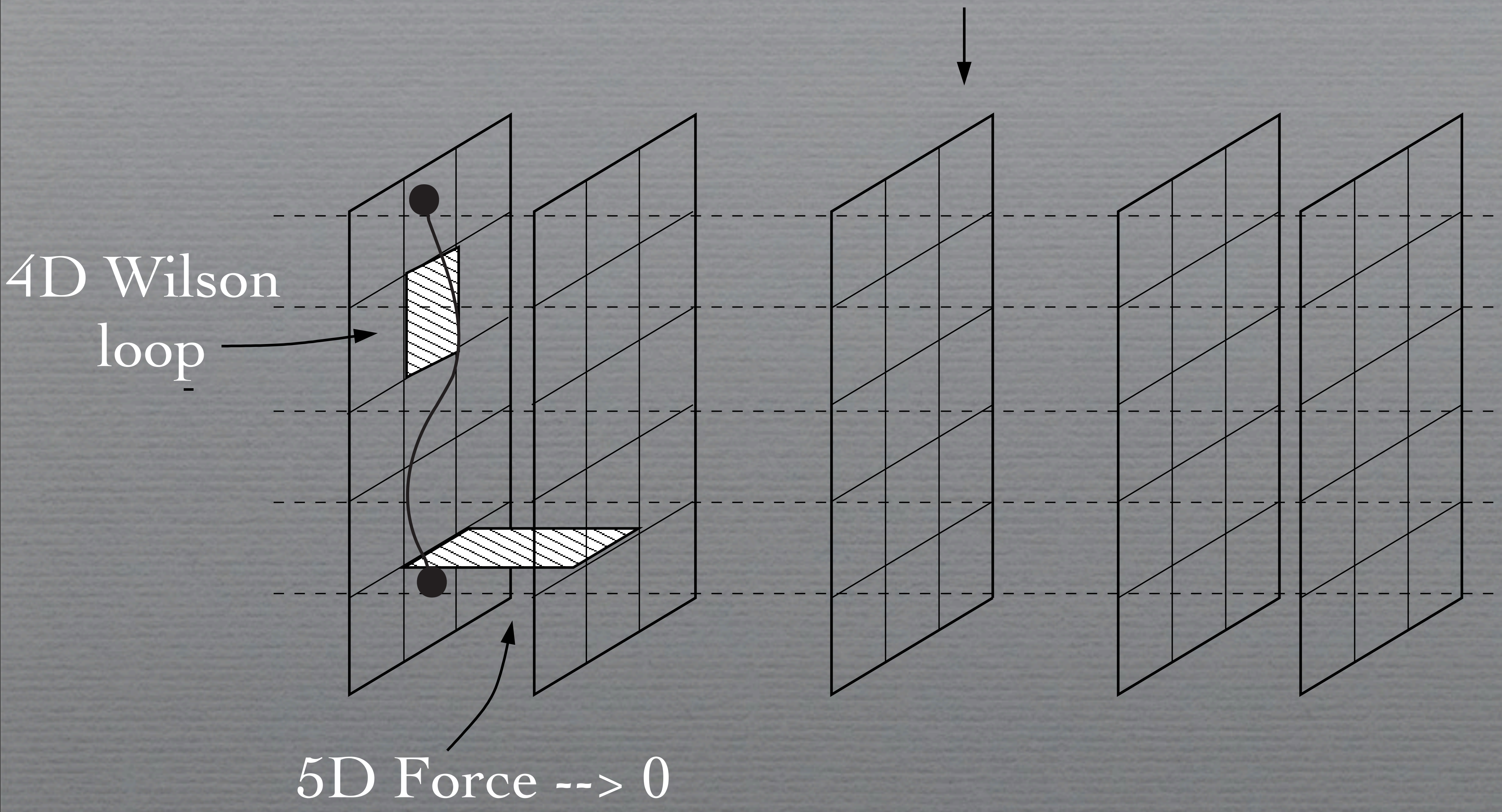




We are looking at an array of non-interacting 4D hyperplanes where gauge interactions are localized.

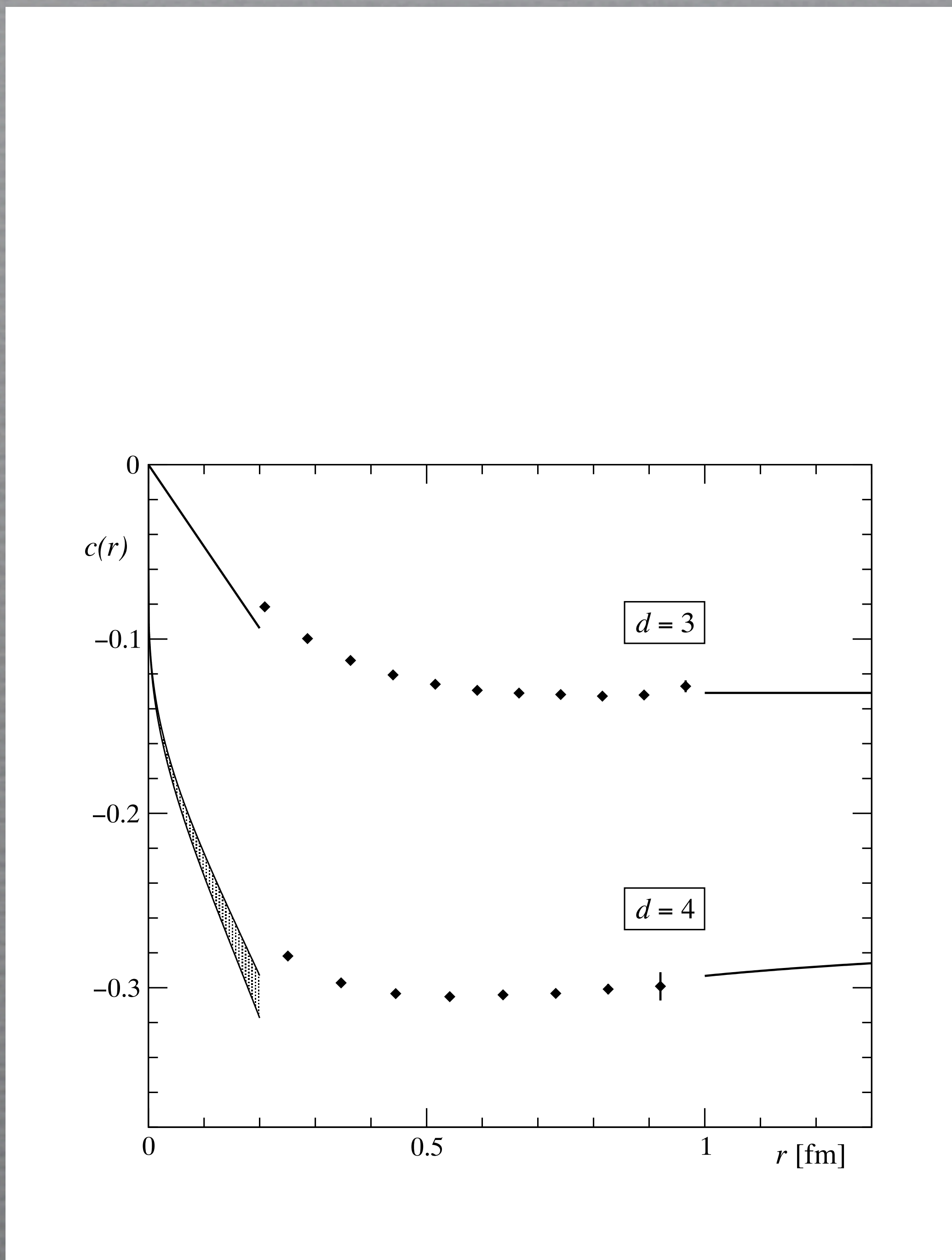
4D G-G on each of the branes.

The 4D Wilson loop at large r must describe a string

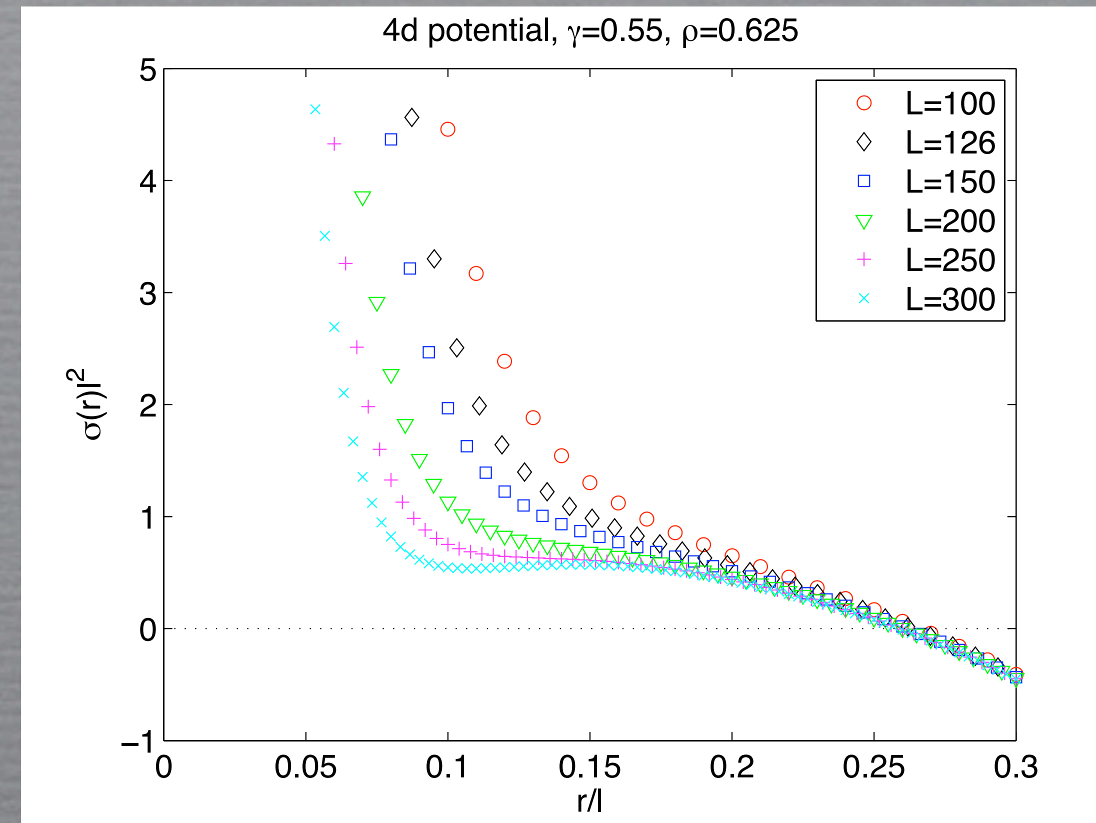


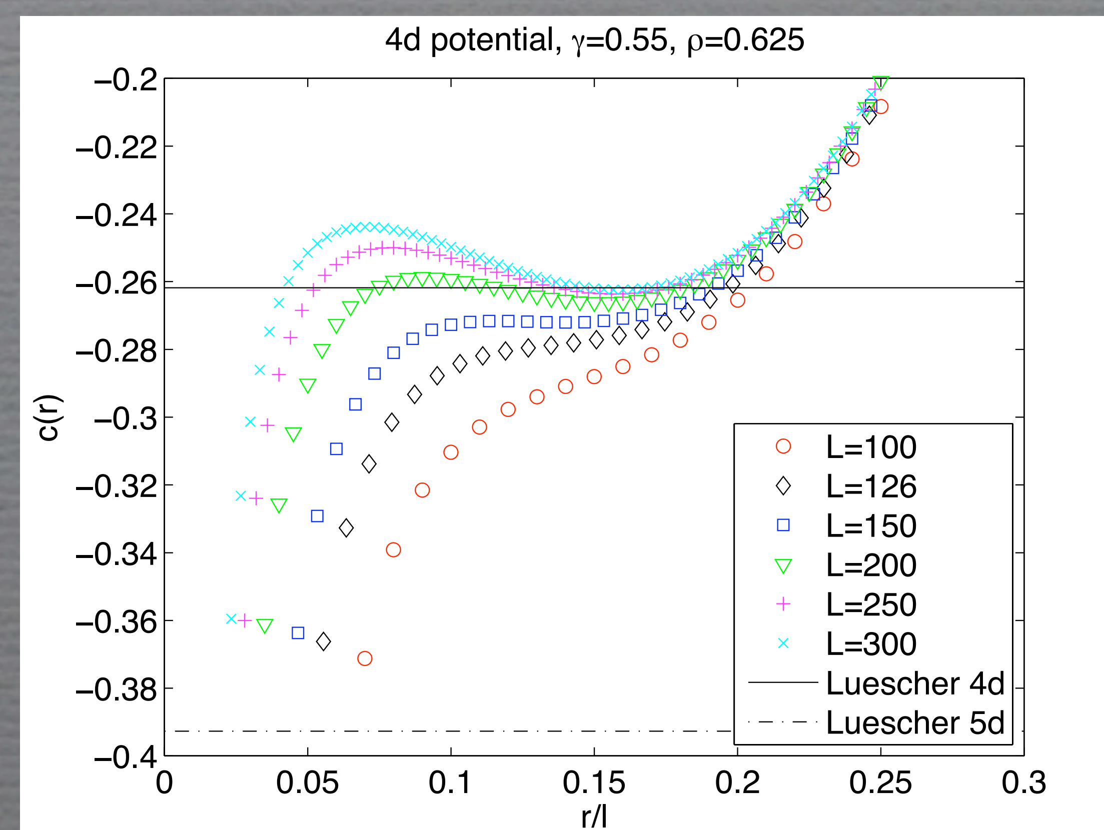
What are we trying to do (preliminary analysis)?

In the Luescher & Weisz paper:



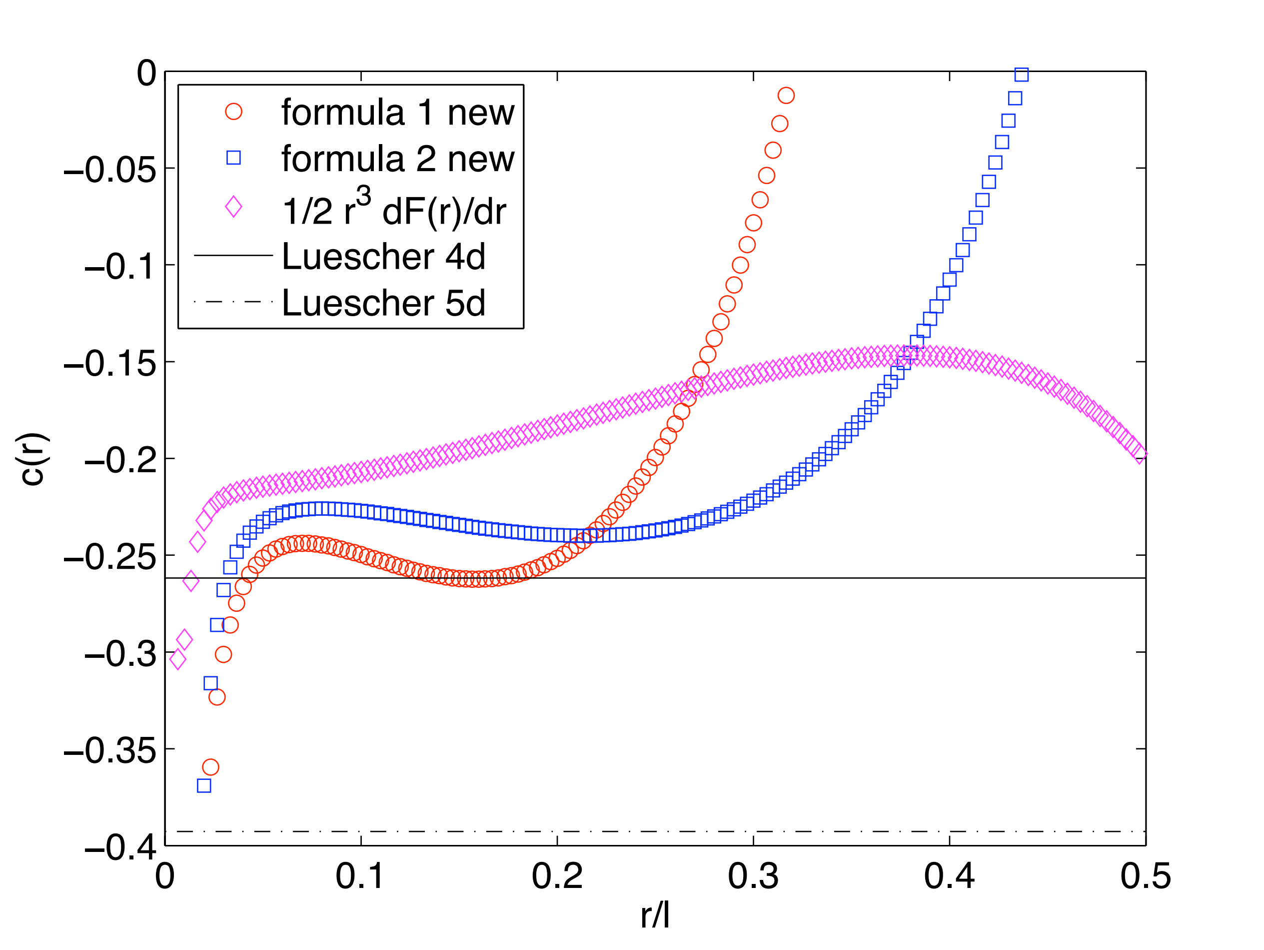
Our formalism:

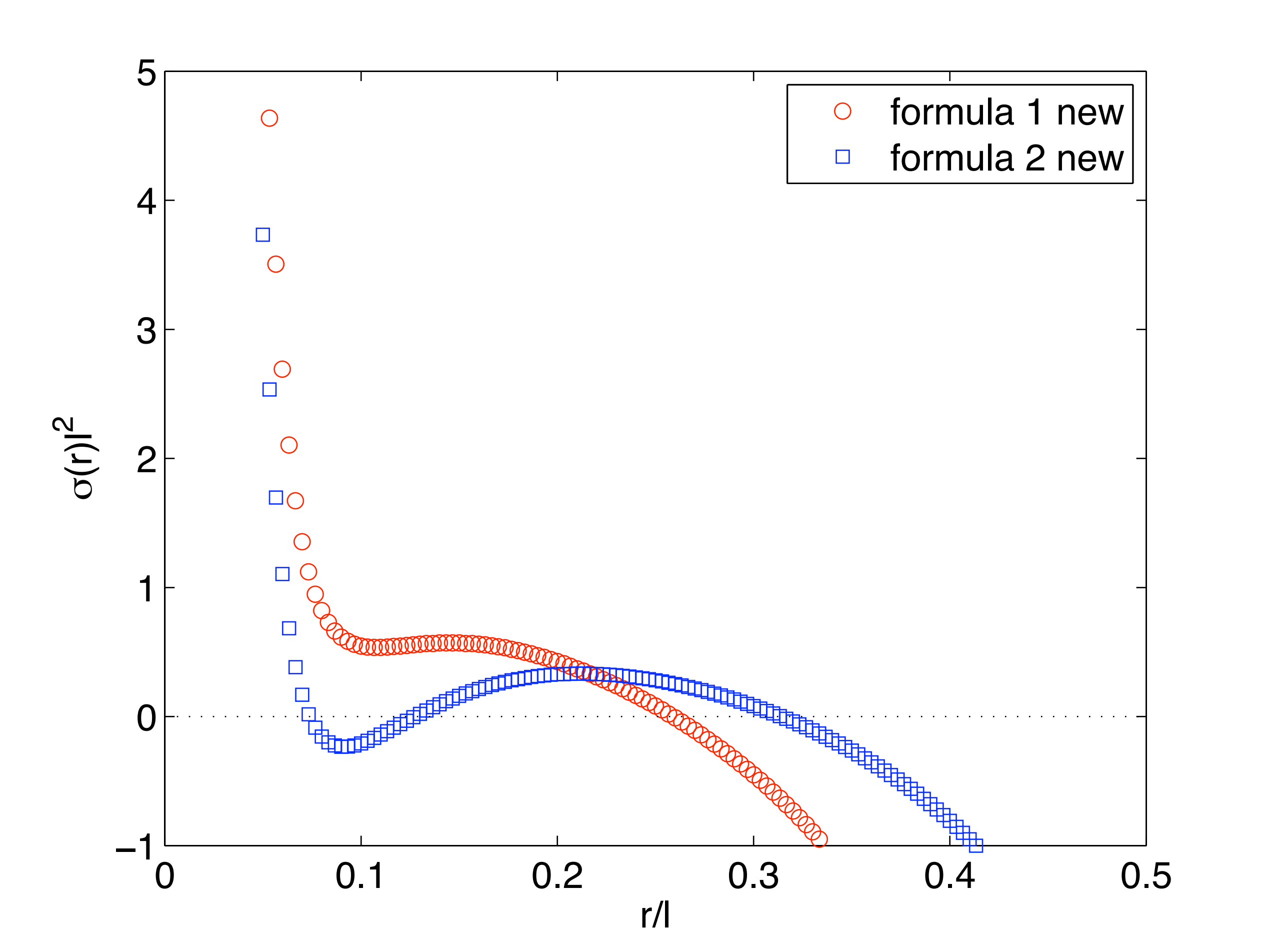




$$V(r) = \mu + \sigma r + \frac{c}{r} + \frac{d}{r^2} + l\log(r)$$

Systematics are not under control yet...
e.g. if the 1/r^2 or the log term is not included in the ansatz:





also, if we move along the line of the phase transition,
...the value of the Luescher term shifts a bit...

Conclusions

1. The non-perturbative regime of 4d gauge theories can be probed analytically by the mean-field expansion in 5d

2. The phase diagram has a 2nd order phase transition where the system reduces dimensionally to an array of non-interacting "3-branes" and where the continuum limit can be taken

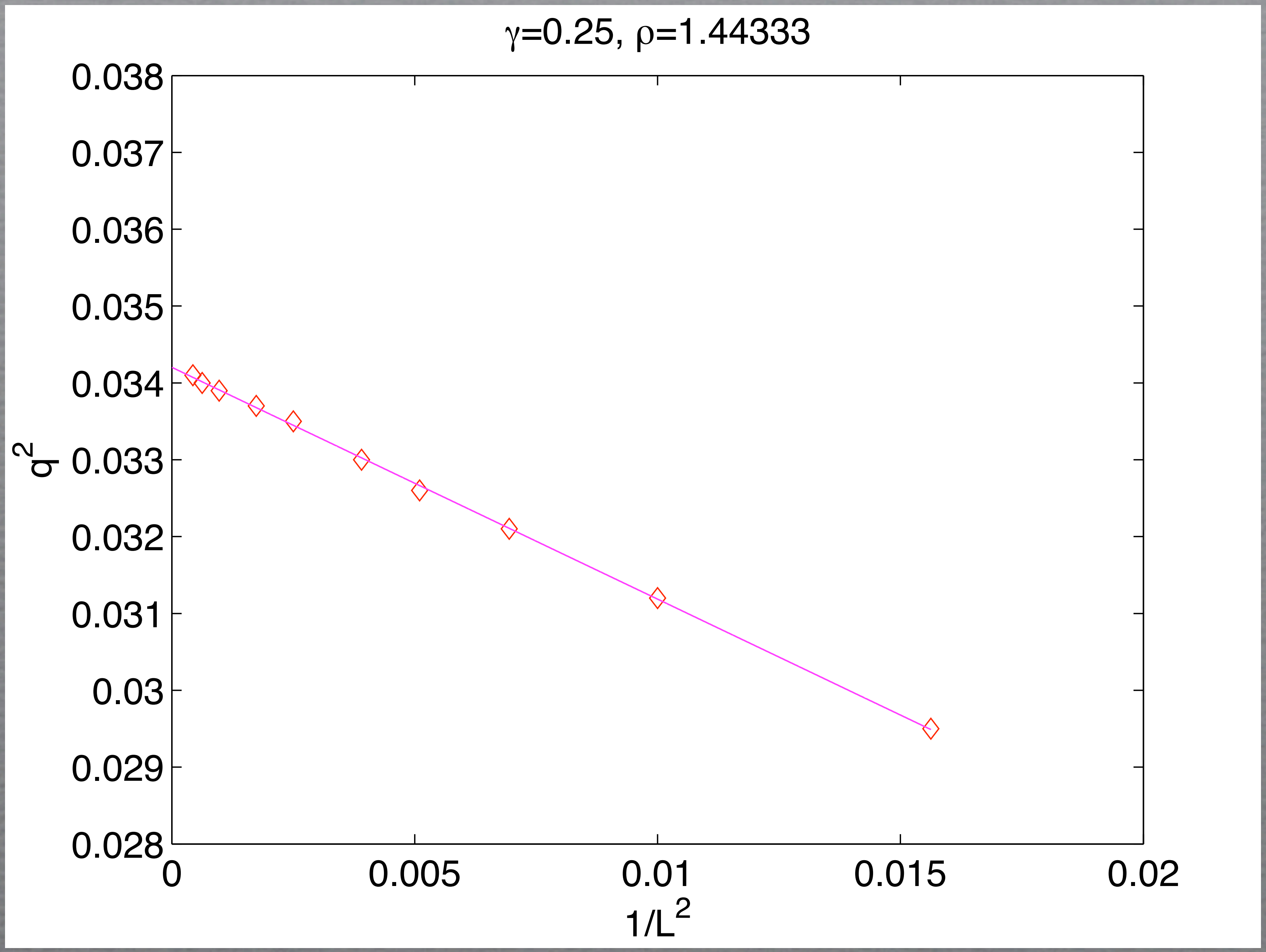
3. In a dimensionally reduced phase there are effects of confinement and of the associated string

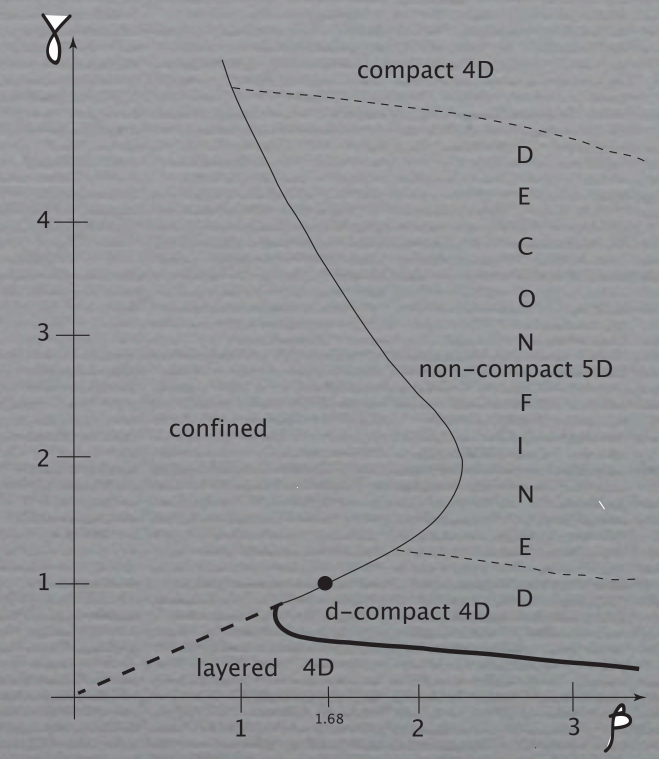
4. Controlling the systematics is in progress

I would like to thank the

Alexander von Humboldt Foundation for finacial support

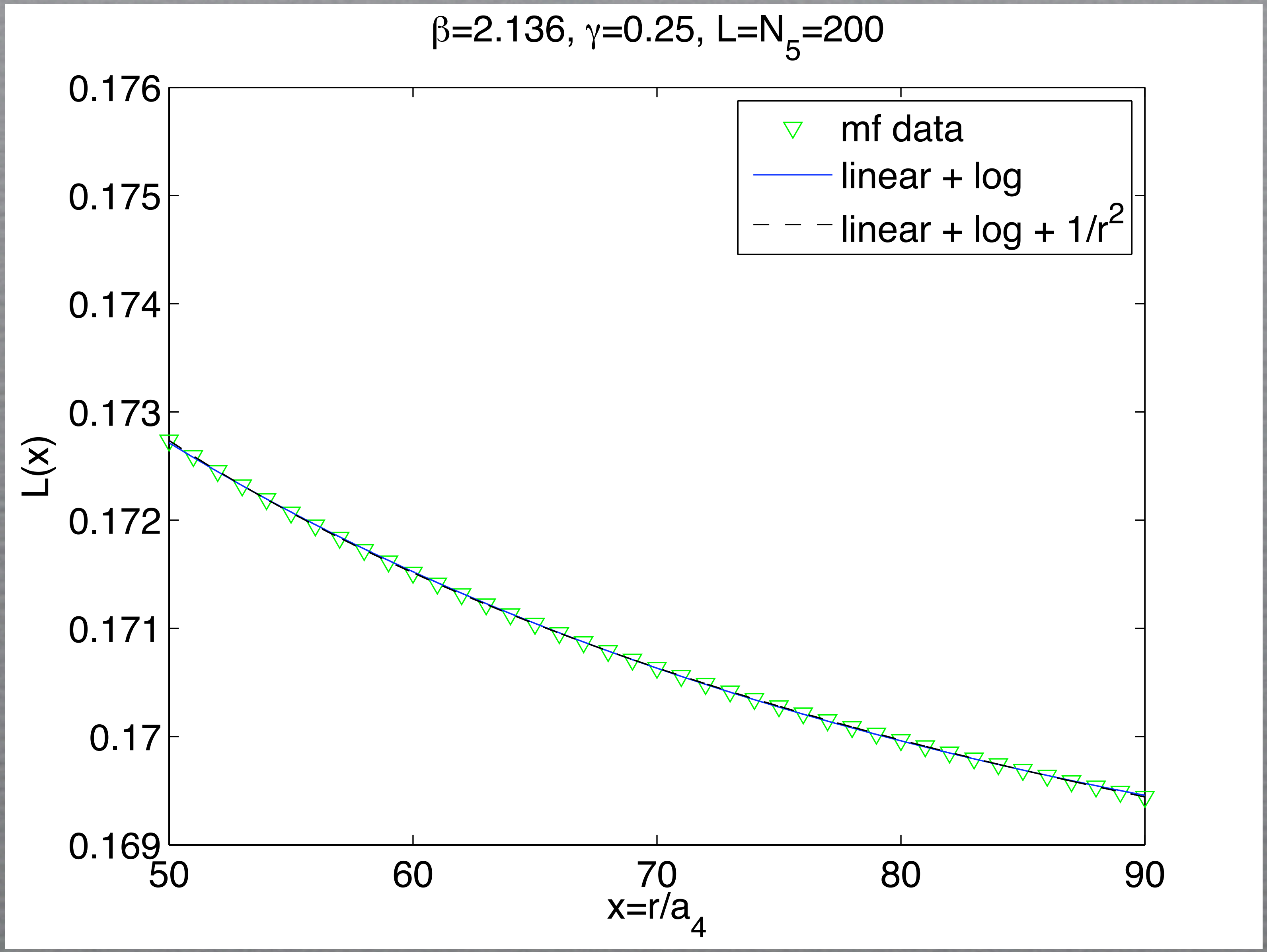
The continuum limit can be taken:





The physical charge at fixed $\rho = \dfrac{m_W}{m_H}$

The long distance static potential in the d-compact phase along the 4d hyperplanes:



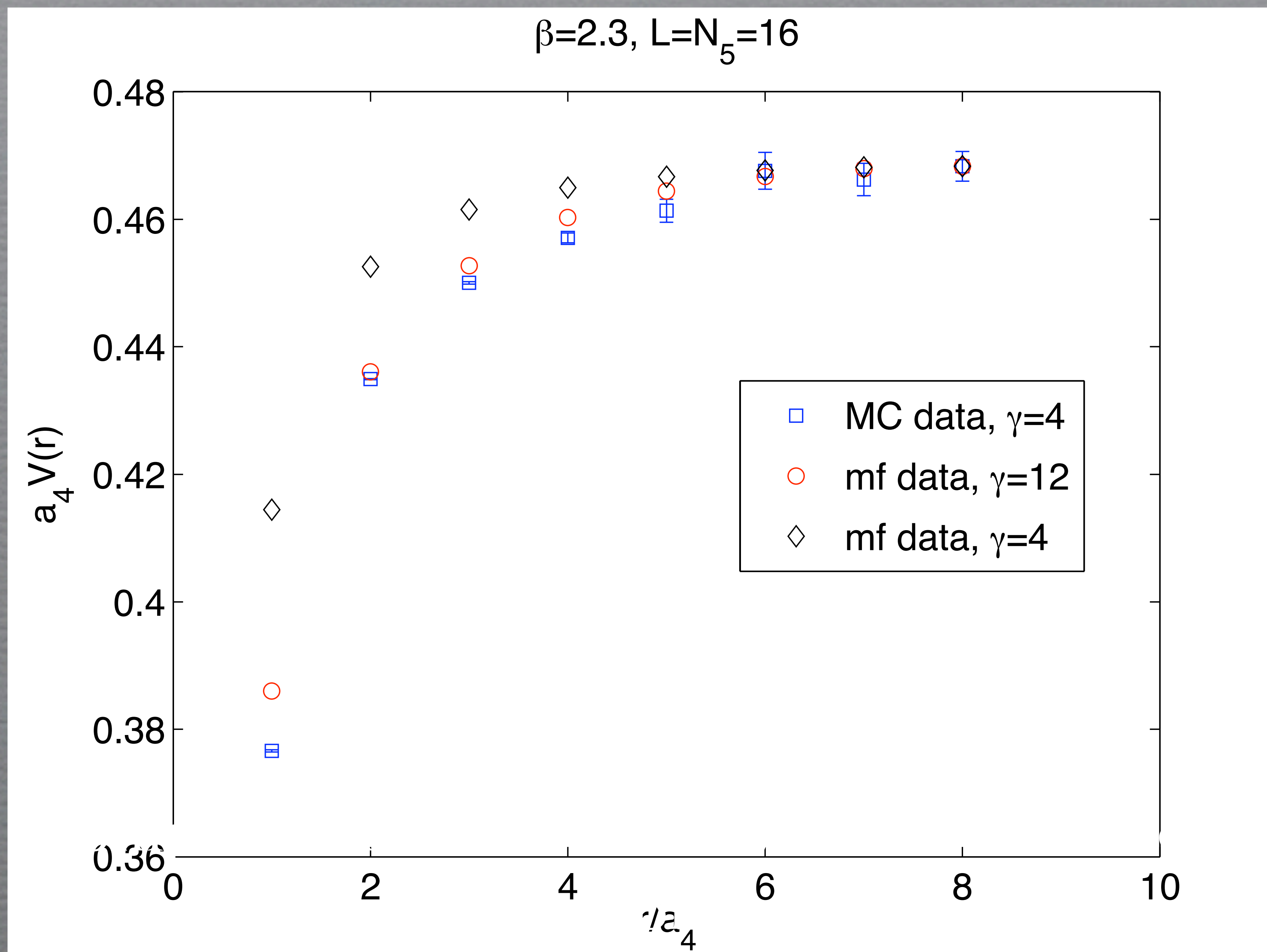
Luescher term

$$\downarrow$$

$$L(x) = a_4\left[V(r) + \pi/(8r)\right]$$

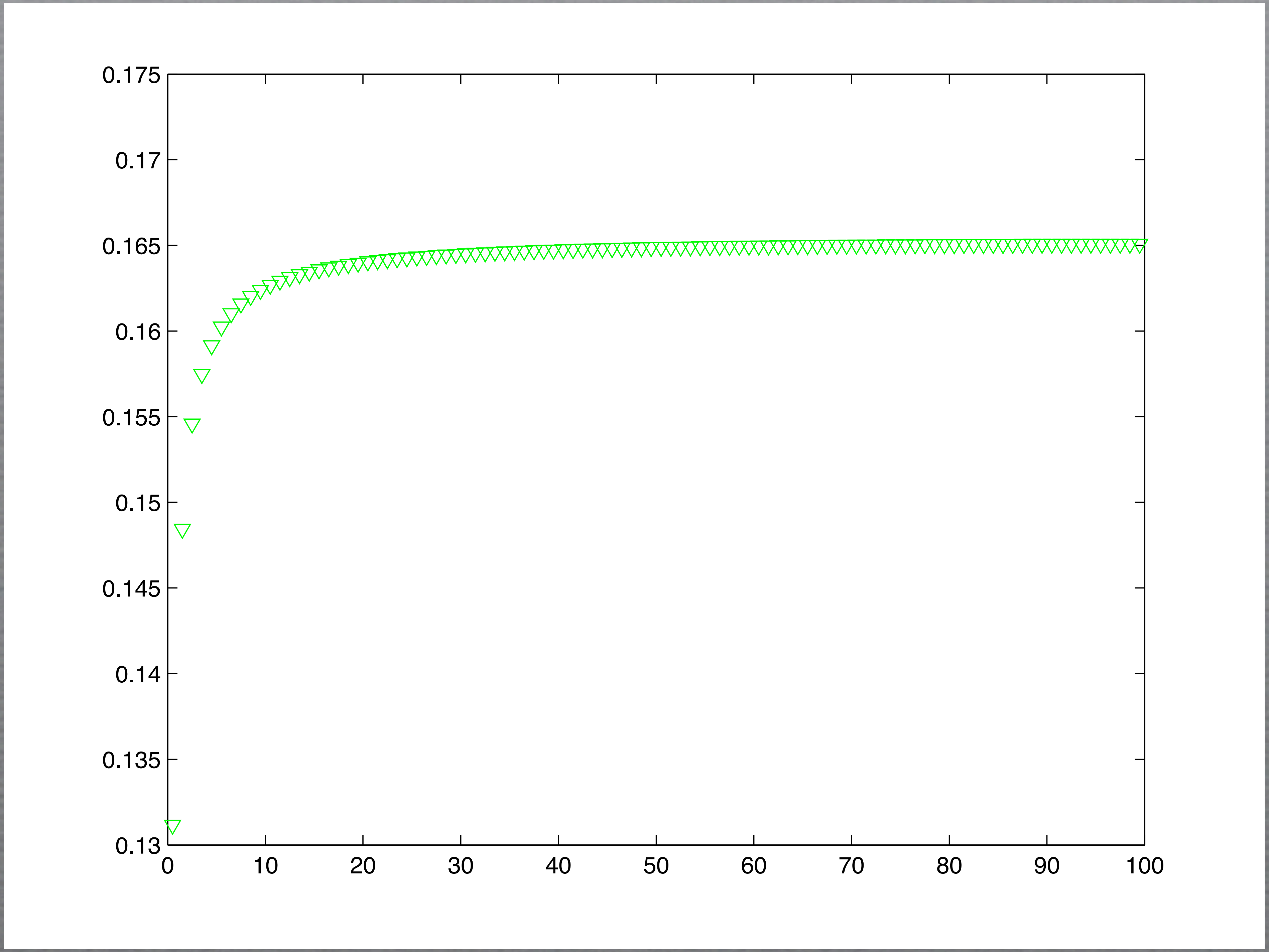
and fit to $\quad \overline{\sigma}x + \overline{\mu} + c_2 \ln(x)$

Comparison Mean-Field vs Monte Carlo: (sample)

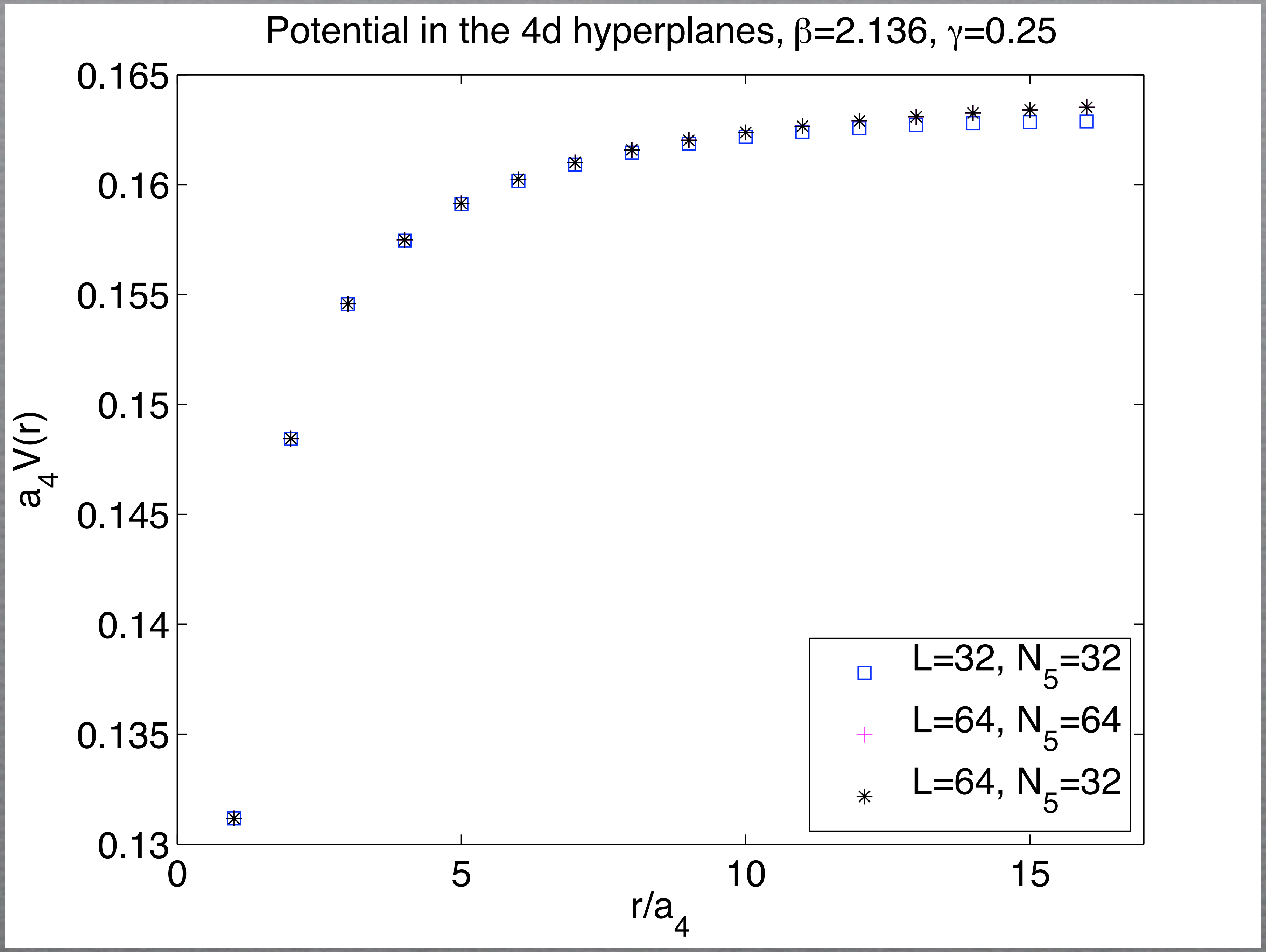


MC data generated by M. Luz

The raw data...


0.175
0.17
0.165
0.16
0.155
0.15
0.145
0.14
0.135
0.13
0
10
20
30
40
50
60
70
80
90
100

Yes, it scales with the lattice size...



No, it can not be 5D Coulomb or Yukawa: